Boltzmann equation in 2D SSL with cross magnetic and electric fields

Rev. 12

1 Analytical treatment

We start from following Boltzmann equation

$$\frac{\partial f}{\partial t} + \frac{e}{\hbar}\left(\mathbf{E} + \mathbf{v} \times \mathbf{B}\right)\frac{\partial f}{\partial \mathbf{k}} + \mathbf{v}(\mathbf{k})\frac{\partial f}{\partial r} = \left(\frac{\partial f}{\partial t}\right)_{st} \tag{1}$$

We consider case when $f(...)$ is spatially homogeneous, i.e. it depends only on $\mathbf{k}$, then

$$\frac{\partial f}{\partial t} + \frac{e}{\hbar}\left(\mathbf{E} + \mathbf{v} \times \mathbf{B}\right)\frac{\partial f}{\partial \mathbf{k}} = \frac{f_0 - f}{\tau} \tag{2}$$

Electric field is along x-axis and magnetic field is along z-axis.

$$\mathbf{E} = (E, 0, 0) \tag{3}$$
$$\mathbf{B} = (0, 0, B) \tag{4}$$
$$\mathbf{v} \times \mathbf{B} = (v_y B, -v_x B, 0)) \tag{5}$$

And will make following substitution

$$t \to \tau t \tag{6}$$

which gives following form to the Boltzmann equation

$$\frac{\partial f}{\partial t} + \frac{ed\tau}{\hbar}(E + v_y B)\frac{\partial f}{\partial(k_x d)} - \frac{ed\tau}{\hbar}v_x B\frac{\partial f}{\partial(k_y d)} = f_0 - f \tag{7}$$

where

$$\mathbf{v} = \frac{1}{\hbar}\frac{\partial \varepsilon}{\partial \mathbf{k}} \tag{8}$$

and ε is an energy of an electron in the zone, which in tight-binding apporiximation takes form in the first zone

$$\varepsilon = \frac{\hbar^2 k_y^2}{2m} - \frac{\Delta_1}{2}\cos(k_x d) \tag{9}$$

therefore

$$v_x = \frac{\Delta_1 d}{2\hbar}\sin(dp_x/h) \tag{10}$$
$$v_y = \frac{\hbar k_y}{m} \tag{11}$$

And Boltzmann equations now takes following form

$$\frac{\partial f}{\partial t} + \left(\frac{ed\tau}{\hbar}E + \frac{ed\tau}{m}k_y B\right)\frac{\partial f}{\partial(k_x d)} - \frac{\Delta_1 e d^2 \tau}{2\hbar^2}B\sin(k_x d)\frac{\partial f}{\partial(k_y d)} = f_0 - f \tag{12}$$

Now we can define

$$E_* = \frac{\hbar}{ed\tau} \tag{13}$$
$$m_x = \frac{2\hbar^2}{\Delta_1 d^2} \tag{14}$$
$$\tag{15}$$

$$\boxed{\alpha = m/m_x} \tag{16}$$

Then, with cyclotron frequency defined as

$$\omega_c = \frac{eB}{\sqrt{mm_x}} \tag{17}$$

we get following

$$\frac{\partial f}{\partial t} + \left(\frac{E}{E_*} + \frac{\omega_c \tau}{\sqrt{\alpha}} k_y d\right) \frac{\partial f}{\partial (k_x d)} - \sqrt{\alpha}\omega_c \tau \sin(k_x d) \frac{\partial f}{\partial (k_y d)} = f_0 - f \tag{18}$$

And if we know define

$$\phi_x = k_x d \tag{19}$$

$$\phi_y = \frac{k_y d}{\sqrt{\alpha}} \tag{20}$$

$$\tilde{E} = E/E_* \tag{21}$$

$$\tilde{B} = \omega_c \tau \tag{22}$$

Boltzmann equations takes form

$$\boxed{\frac{\partial f}{\partial t} + \left(\tilde{E} + \tilde{B}\phi_y\right) \frac{\partial f}{\partial \phi_x} - \tilde{B}\sin(\phi_x) \frac{\partial f}{\partial \phi_y} = f_0 - f} \tag{23}$$

Alternatively, if we define Bloch frequency as

$$\omega_B = \frac{edE}{\hbar} \tag{24}$$

it can be written as

$$\frac{\partial f}{\partial t} + \left(\omega_B \tau + \omega_c \tau \phi_y\right) \frac{\partial f}{\partial \phi_x} - \omega_c \tau \sin(\phi_x) \frac{\partial f}{\partial \phi_y} = f_0 - f \tag{25}$$

Ratio of Bloch and cyclotron oscillation frequencies will be important to us down the road. For f_0 we use Boltzmann distribution

$$f_0 \propto e^{-\frac{\varepsilon}{k_b T}} \tag{26}$$

which in our variables ϕ_x and ϕ_y takes form

$$f_0 = C \exp\left\{\mu \cos(\phi_x) - \frac{\mu}{2}\phi_y^2\right\} \tag{27}$$

$$\boxed{\mu = \frac{\Delta_1}{2k_b T}} \tag{28}$$

Now, constant C must be such that dimensionless norm of f_0 is 1. Therefore

$$\frac{1}{C} = \frac{d^2}{\hbar^2} \int_{-\infty}^{+\infty} \mathrm{d}p_y \int_{p_x \in \mathrm{BZ}} \mathrm{d}p_x \exp\left\{\mu \cos(\phi_x) - \frac{\mu}{2}\phi_y^2\right\} \tag{29}$$

$$= \sqrt{\alpha} \int_{-\infty}^{+\infty} e^{-\mu \phi_y^2/2} \mathrm{d}\phi_y \int_{-\pi}^{\pi} e^{\mu \cos(\phi_x)} \mathrm{d}\phi_x \tag{30}$$

$$= 2\pi I_0(\mu) \sqrt{\frac{2\pi\alpha}{\mu}} \tag{31}$$

Thus full form of equilibrium distribution is

$$\boxed{f_0 = \frac{1}{2\pi I_0(\mu)} \sqrt{\frac{\mu}{2\pi\alpha}} \exp\left\{\mu \cos(\phi_x) - \frac{\mu}{2}\phi_y^2\right\}} \tag{32}$$

This way, in total, we have four free parameters defining our system, μ and α that characterize lattice parameters such as lattice period and band width and temperature, and $\tilde{E}$, $\tilde{B}$ that specify external fields. One of the parameters that we will be calculating is drift velocity v_{dr}, which we will define like this

$$v_{dr} =\frac{2d}{\Delta_1\hbar}\iint\frac{\partial\varepsilon}{\partial p_x}f(p_x,p_y)\mathrm{d}p_x\mathrm{d}p_y \tag{33}$$

$$=\sqrt{\alpha}\int_{-\pi}^{\pi}\mathrm{d}\phi_x\int_{-\infty}^{+\infty}\mathrm{d}\phi_y\sin(\phi_x)f(\phi_x,\phi_y) \tag{34}$$

When magentic field is 0 and E is constant, i.e. E_{dc}, analytic solution Boltzmann equation is well known and as well as expression for drift velocity

$$v_{dr} =\left\{\sqrt{\alpha}\int_{-\pi}^{\pi}\mathrm{d}\phi_x\int_{-\infty}^{+\infty}\mathrm{d}\phi_y\sin(\phi_x)f_0(\phi_x,\phi_y)\right\}\frac{E_{dc}/E_*}{1+(E_{dc}/E_*)^2} \tag{35}$$

$$=\frac{I_1(\mu)}{I_0(\mu)}\frac{E_{dc}/E_*}{1+(E_{dc}/E_*)^2} \tag{36}$$

From which follows that peak value of v_{dr} is at $E_{dc}/E_* = 1$ and is

$$v_p = \frac{I_1(\mu)}{2I_0(\mu)} \tag{37}$$

Later, down the road, we will be plotting not v_{dr}, but v_{dr}/v_p, which for dc electric field only take very simple form, known as "Esaki-Tsu" equation

$$\frac{v_{dr}}{v_p} = 2\frac{E_{dc}/E_*}{1+(E_{dc}/E_*)^2} \tag{38}$$

In general we will be applying a/c emf in the form

$$\tilde{E} = \tilde{E}_{dc} + \tilde{E}_\omega\cos(\omega t) \tag{39}$$

and in case when magnetic field is not applied we also have analitic expression for v_{dr}, which is known as "Tien-Gordon" equation. Analytic expression, "Taker formulae", is also known for absorption, which we will define like this

$$A = \left\langle\frac{v_{dr}}{v_p}\cos(\omega t)\right\rangle_t \tag{40}$$

In case when magnetic field is appied, however, analytic expression for absorption is not known. And that is the quantity we are most interested in.

Now due to periodicty of $f(\phi_x,\phi_y)$ along ϕ_x with period 2π and additionally $f_0(-\phi_x,\phi_y) = f_0(\phi_x,\phi_y)$, it makes sense for us to expand f and f_0 into Fourier series.

$$f_0 = \sum_{n=0}^{\infty}a_n^{(0)}\cos(n\phi_x) \tag{41}$$

$$f = \sum_{n=0}^{\infty}a_n\cos(n\phi_x) + b_n\sin(n\phi_x) \tag{42}$$

where coeffients $a_n^{(0)}$, a_n and b_n in general will depend on ϕ_y and last two also on time t. Note, that $a_{n<0}\equiv 0$ and $b_{n<1}\equiv 0$. And with the from of f_0, as selected above, $a_n^{(0)}$ becomes

$$a_n^{(0)} =\frac{\sigma(n)}{\pi}\int_{-\pi}^{\pi}f_0(\phi_x,\phi_y)\cos(n\phi_x)\mathrm{d}\phi_x \tag{43}$$

$$=\frac{\sigma(n)I_n(\mu)}{\pi I_0(\mu)}\sqrt{\frac{\mu}{2\pi\alpha}}\exp\left\{-\frac{\mu}{2}\phi_y^2\right\} \tag{44}$$

where

$$\sigma(n) = \begin{cases} 1/2 & : n = 0 \\ 1 & : n \neq 1 \end{cases} \tag{45}$$

Norm of $f(\phi_x, \phi_y)$ must always be one, i.e.

$$\sqrt{\alpha} \int_{-\pi}^{\pi} \mathrm{d}\phi_x \int_{-\infty}^{+\infty} \mathrm{d}\phi_y f(\phi_x, \phi_y) = 1 \tag{46}$$

which in fourier representation takes form

$$\boxed{2\pi\sqrt{\alpha} \int_{-\infty}^{+\infty} a_0(\phi_y)\mathrm{d}\phi_y = 1} \tag{47}$$

Once we move on to the numerical calculations this equation can be used to check correctness. And from equation (34) it is clear that only b_1 will survice. And calculation of v_{dr} is done through following equation

$$\begin{aligned} v_{dr} =& \sqrt{\alpha} \int_{-\infty}^{+\infty} \mathrm{d}\phi_y \int_{-\pi}^{+\pi} \mathrm{d}\phi_x \sin(\phi_x) b_1(\phi_y) \sin(\phi_y) && (48) \\ =& \pi\sqrt{\alpha} \int_{-\infty}^{+\infty} b_1(\phi_y)\mathrm{d}\phi_y && (49) \end{aligned}$$

and in view of definition of peak value of v_{dr} in eq. (37)

$$\boxed{\frac{v_{dr}}{v_p} = \frac{2I_0(\mu)\pi\sqrt{\alpha}}{I_1(\mu)} \int_{-\infty}^{+\infty} b_1(\phi_y)\mathrm{d}\phi_y} \tag{50}$$

In addition to drift velocity along x-axis we can look at drift velocity alogn y-axis, which we will define like this

$$\begin{aligned} v_y =& \frac{2}{\Delta_1} \iint \frac{\partial \varepsilon}{\partial p_y} f(\phi_x, \phi_y)\mathrm{d}p_x\mathrm{d}p_y && (51) \\ =& \int_{-\pi}^{\pi} \mathrm{d}\phi_x \int_{-\infty}^{+\infty} \phi_y f(\phi_x, \phi_y)\mathrm{d}\phi_y && (52) \\ =& 2\pi \int_{-\infty}^{+\infty} a_0(\phi_y)\phi_y\mathrm{d}\phi_y && (53) \end{aligned}$$

However just as with drift velocity along x-axis we will be working with ratio of v_y and peak velocity v_p.

$$\boxed{\frac{v_y}{v_p} = \frac{4\pi I_0(\mu)}{I_1(\mu)} \int_{-\infty}^{+\infty} a_0(\phi_y)\phi_y\mathrm{d}\phi_y} \tag{54}$$

This way we accounted for meaning of a_0 and b_1. Let us now take a look at a_1. Effective mass of an electron in the x direction is given by

$$m_{x,\mathbf{k}}^{-1} = \frac{1}{\hbar^2} \frac{\partial^2 \varepsilon}{\partial k_x^2} \tag{55}$$

We will be working however with ratio of electron mass to $m_{x,\mathbf{k}}$. And in our approximation of energy (9) that takes form

$$\begin{aligned} \frac{m}{m_{x,\mathbf{k}}} =& \frac{\Delta_1 d^2 m}{2\hbar^2} \cos(k_x d) && (56) \\ =& \alpha \cos(\phi_x) && (57) \end{aligned}$$

To gather back this value from $f(\phi_x, \phi_y)$ we have to integrate over $\{p_x, p_y\}$ and to maintain dimensionlessness of ratio of these masses, this integration will take form

$$\frac{m}{m_{x,\mathbf{k}}} = \alpha^{3/2}\frac{d^2}{\hbar^2}\iint f(\phi_x, \phi_y)\cos(k_x d)\mathrm{d}p_x\mathrm{d}p_y \tag{58}$$

$$= \alpha^{3/2}\int_{-\pi}^{\pi}\mathrm{d}\phi_x\int_{-\infty}^{+\infty}\mathrm{d}\phi_y f(\phi_x, \phi_y)\cos(\phi_x) \tag{59}$$

$$= \alpha^{3/2}\int_{-\pi}^{\pi}\mathrm{d}\phi_x\int_{-\infty}^{+\infty}\mathrm{d}\phi_y a_1(\phi_y)\cos(\phi_x)\cos(\phi_x) \tag{60}$$

Finally giving us following

$$\boxed{\frac{m}{m_{x,\mathbf{k}}} = \pi\alpha^{3/2}\int_{-\infty}^{+\infty} a_1(\phi_y)\mathrm{d}\phi_y} \tag{61}$$

Now using fourier representation of $f(\phi_x, \phi_y)$ and $f_0(\phi_y)$ (equations 41, 42) we can rewrite Boltzmann equation (23) like this

$$\sum_{(n)}\left\{\frac{\partial a_n}{\partial t}\cos(n\phi_x) + \frac{\partial b_n}{\partial t}\sin(n\phi_x) = a_n^{(0)}\cos(n\phi_x) - a_n\cos(n\phi_x) - b_n\sin(n\phi_x) + \right.$$
$$\left. n(\tilde{E}+\tilde{B}\phi_y)(a_n\sin(n\phi_x) - b_n\cos(n\phi_x)) + \tilde{B}\frac{\partial a_n}{\partial \phi_y}\sin(\phi_x)\cos(n\phi_x) + \tilde{B}\frac{\partial b_n}{\partial \phi_y}\sin(\phi_x)\sin(n\phi_x)\right\} \tag{62}$$

In absence of magnetic field there is no mixing of different harmonic, however when $\tilde{B}$ is not 0 then harmonics will be come mixed due to presence of $\sin(\phi_x)\cos(n\phi_x)$ and $\sin(\phi_x)\sin(n\phi_x)$, since

$$\sin(\phi_x)\cos(n\phi_x) = \{\sin((n+1)\phi_x) - \sin((n-1)\phi_x)\}/2 \tag{63}$$

$$\sin(\phi_x)\sin(n\phi_x) = \{\cos((n-1)\phi_x) - \cos((n+1)\phi_x)\}/2 \tag{64}$$

And using this equations, after some manipultion of symbols, combining elements with the same harmonics, and noting special treatment of b_1 we get

$$\frac{\partial a_n}{\partial t} = a_n^{(0)} - a_n - n(\tilde{E}+\tilde{B}\phi_y)b_n + \tilde{B}\left(\frac{\partial b_{n+1}}{\partial \phi_y} - \frac{\partial b_{n-1}}{\partial \phi_y}\right) \tag{65}$$

$$\frac{\partial b_n}{\partial t} = -b_n - n(\tilde{E}+\tilde{B}\phi_y)a_n + \tilde{B}\left(\chi(n)\frac{\partial a_{n-1}}{\partial \phi_y} - \frac{\partial a_{n+1}}{\partial \phi_y}\right) \tag{66}$$

where

$$\chi(n) = \begin{cases} 2 & : n = 1 \\ 1 & : n \neq 1 \end{cases} \tag{67}$$

2 Correspondence with classical pendilum

In the limit where dissipation is absent instead of Boltzmann equation we can use semiclassical equations of motion working individual electrons. Such situation corresponds to $\tau = \infty$ and initial shape of $f(\phi_x, \phi_x)$ being delta function.

$$\hbar\frac{\mathrm{d}\mathbf{k}}{\mathrm{d}t} = e\mathbf{E} + e\mathbf{v}\times\mathbf{B} \tag{68}$$

$$\mathbf{v}(\mathbf{k}) = \frac{1}{\hbar}\frac{\partial\varepsilon}{\partial\mathbf{k}} \tag{69}$$

And in view of specific values of $\mathbf{E}$, $\mathbf{B}$ and ε that we are using in our problem, these equations transform into following

$$\frac{\mathrm{d}\phi_x}{\mathrm{d}t} = \tilde{E} + \tilde{B}\phi_y \tag{70}$$

$$\frac{\mathrm{d}\phi_y}{\mathrm{d}t} = -\tilde{B}\sin(\phi_x) \tag{71}$$

And taking second derivative of (70), assuming that $\tilde{E}$ is defined by (39) we arrive to second order ODE, which correspond to classical externally driven pendilum.

$$\frac{\mathrm{d}^2\phi_x}{\mathrm{d}t^2} + \tilde{B}^2\sin(\phi_x) = -\omega\tilde{E}_\omega\sin(\omega t) \tag{72}$$

In general this equation can exibit very complicated behaviour and even chaos, but can be solved in the simplest case when $\tilde{E}_\omega = 0$, in which case peroid of oscillation is defined by magnetic field $\tilde{B}$ and electric field $\tilde{E}$, which sets initial velocity of ϕ_x, i.e. $\frac{\mathrm{d}\phi_x}{\mathrm{d}t}$ at $t = 0$, assuming that ϕ_x and ϕ_y at $t = 0$ are also 0. In case of $E_\omega = 0$ we exact correspondence of our system to a classical pendilum, for which we do have conservation of energy.

$$\frac{1}{2}\left(\frac{\mathrm{d}\phi_x}{\mathrm{d}t}\right)^2 + 2\tilde{B}^2\sin^2(\phi_x/2) = E_{total} \tag{73}$$

If we define $E_p = 2\tilde{B}^2$, which corresponds to "peak" then when $E_{total} = E_p$ then such pendilum can reach inverted position. When ϕ_x and ϕ_y are equal zero $E_{total} = \tilde{E}_{dc}^2/2$ and thus for E_p there is corresponding $E_{sx} = 2\tilde{B}$, which defines separatrix dividing motion of the pendilum between vibrational ($\tilde{E}_{dc} < E_{sx}$) and rotational ($\tilde{E}_{dc} > E_{sx}$). When our system gives peak of negative absorption corresponding to the first and second cases we speak of "Cycloron" and "Bloch" resonances respectively. Expresion for frequency Ω of these natural oscilations is well known and in our variables it takes form

$$\Omega = \frac{\pi E_{sx}}{4\mathrm{K}(\tilde{E}_{dc}/E_{sx})} \qquad \text{for } \tilde{E}_{dc} < E_{sx} \tag{74}$$

$$\Omega = \frac{\pi E_{dc}}{2\mathrm{K}(E_{sx}/\tilde{E}_{dc})} \qquad \text{for } \tilde{E}_{dc} > E_{sx} \tag{75}$$

where $\mathrm{K}(x)$ is a complete elliptic integral of the first kind.

3 Numerical solution

Straightforward application of method of finite differences to (65) and (66) leads to either unstable or i.e. computationally intensive equations. To combat this problem we am using several methods at once. First, we are going to discretize a_n and b_n along time and ϕ_y axes.

$$a^{t \leftarrow \text{time step}}_{n,m \leftarrow \phi_y \text{lattice step}} \tag{76}$$

and n is "harmonic number". So, here we are going to do some tricky things. We are going to write two forms of equations (65, 66). One using forward differences and one using partial backward differences, i.e. on the right side of equal sign we are going to write partial derivatives at time t while everything else at time $t+1$ and will follow standard procedure of CrankNicolson scheme by adding these two, forward and backward differences equations. First, forward differencing scheme

$$\begin{aligned} a^{t+1}_{n,m} - a^t_{n,m} =& a^{(0)}_{n,m}\Delta t - a^t_{n,m}\Delta t - 2b^t_{n,m}\mu^t_{n,m} + \\ &+ \frac{\alpha B\Delta t}{2\Delta\phi}(b^t_{n+1,m+1} - b^t_{n+1,m-1} - b^t_{n-1,m+1} + b^t_{n-1,m-1}) \end{aligned} \tag{77}$$

$$\begin{aligned} b^{t+1}_{n,m} - b^t_{n,m} =& -b^t_{n,m}\Delta t + 2a^t_{n,m}\mu^t_{n,m} + \\ &+ \frac{\alpha B\Delta t}{2\Delta\phi}(\chi(n)[a^t_{n-1,m+1} - a^t_{n-1,m-1}] - a^t_{n+1,m+1} + a^t_{n+1,m-1}) \end{aligned} \tag{78}$$

And then partial backward differencing scheme

$$\begin{aligned}
a^{t+1}_{n,m} - a^t_{n,m} =& a^{(0)}_{n,m}\Delta t - a^{t+1}_{n,m}\Delta t - 2b^{t+1}_{n,m}\mu^{t+1}_{n,m} + \\
&+ \frac{\alpha B \Delta t}{2\Delta\phi}(b^t_{n+1,m+1} - b^t_{n+1,m-1} - b^t_{n-1,m+1} + b^t_{n-1,m-1}) \qquad (79) \\
b^{t+1}_{n,m} - b^t_{n,m} =& - b^{t+1}_{n,m}\Delta t + 2a^{t+1}_{n,m}\mu^{t+1}_{n,m} + \\
&+ \frac{\alpha B \Delta t}{2\Delta\phi}(\chi(n)[a^t_{n-1,m+1} - a^t_{n-1,m-1}] - a^t_{n+1,m+1} + a^t_{n+1,m-1}) \qquad (80)
\end{aligned}$$

where

$$\begin{aligned}
\beta^t_m =& E^t + B^t\phi_y(m) \qquad (81) \\
\mu^t_{n,m} =& n\beta^t_m \Delta t \qquad (82)
\end{aligned}$$

And application of CrankNicolson scheme leads to

$$\begin{aligned}
a^{t+1}_{n,m} =& \frac{g^t_{n,m}\nu - h^t_{n,m}\mu^{t+1}_{n,m}}{\nu^2 + \left(\mu^{t+1}_{n,m}\right)^2} \qquad (83) \\
b^{t+1}_{n,m} =& \frac{g^t_{n,m}\mu^{t+1}_{n,m} - h^t_{n,m}\nu}{\nu^2 + \left(\mu^{t+1}_{n,m}\right)^2} \qquad (84)
\end{aligned}$$

where

$$\begin{aligned}
\nu =& 1 + \Delta t/2 \qquad (85) \\
\xi =& 1 - \Delta t/2 \qquad (86)
\end{aligned}$$

$$\begin{aligned}
g^t_{n,m} =& a^t_{n,m}\xi - b^t_{n,m}\mu^t_{n,m} + A^t_{n,m} + a^{(0)}_{n,m}\Delta t \qquad (87) \\
h^t_{n,m} =& b^t_{n,m}\xi + a^t_{n,m}\mu^t_{n,m} + B^t_{n,m} \qquad (88) \\
A^t_{n,m} =& \frac{\alpha B \Delta t}{2\Delta\phi}(\chi(n)[a^t_{n-1,m+1} - a^t_{n-1,m-1}] - a^t_{n+1,m+1} + a^t_{n+1,m-1}) \qquad (89) \\
B^t_{n,m} =& \frac{\alpha B \Delta t}{2\Delta\phi}(b^t_{n+1,m+1} - b^t_{n+1,m-1} - b^t_{n-1,m+1} + b^t_{n-1,m-1}) \qquad (90)
\end{aligned}$$

These equation (83, 84)can be formally written in the form

$$\begin{aligned}
\mathbf{r}^{t+1}_{n,m} =& \mathbf{T}(\mathbf{r}^t_{n,m}; A^t_{n,m}, B^t_{n,m}) \qquad (91) \\
\mathbf{r}^t_{n,m} =& (a^t_{n,m}, b^t_{n,m}) \qquad (92)
\end{aligned}$$

Where $\mathbf{T}$ is an operator that allows us to step from time step t to $t+1$, separated by time interval Δt. Using this operation as is leads to, only, conditionally stable numerical system, because $A^t_{n,m}$ and $B^t_{n,m}$ are taken at time t, which means that we have to take time step dt to be very, very small. To make time step much larger without solving implicit equations, to calculate $a^{t+1}_{n,m}$ and $b^{t+1}_{n,m}$ we will utilize variation of leap-frog method, by using two staggered girds.

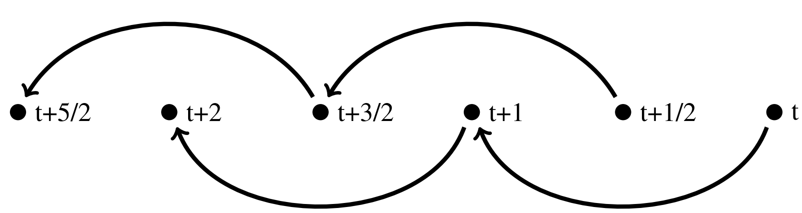


Figure 1

And writing (91) in the form

$$\mathbf{r}_{n,m}^{t+1} = \mathbf{T}(\mathbf{r}_{n,m}^{t}; A_{n,m}^{t+1/2}, B_{n,m}^{t+1/2}) \tag{93}$$

$$\mathbf{r}_{n,m}^{t+3/2} = \mathbf{T}(\mathbf{r}_{n,m}^{t+1/2}; A_{n,m}^{t+1}, B_{n,m}^{t+1}) \tag{94}$$

Use of equations (93) and (94) is indicated by lower and upper arrows in fig 1. To start this system we calculate $\mathbf{r}_{n,m}^{1/2}$ by using eq. (91) and then shift to leap-frog method.

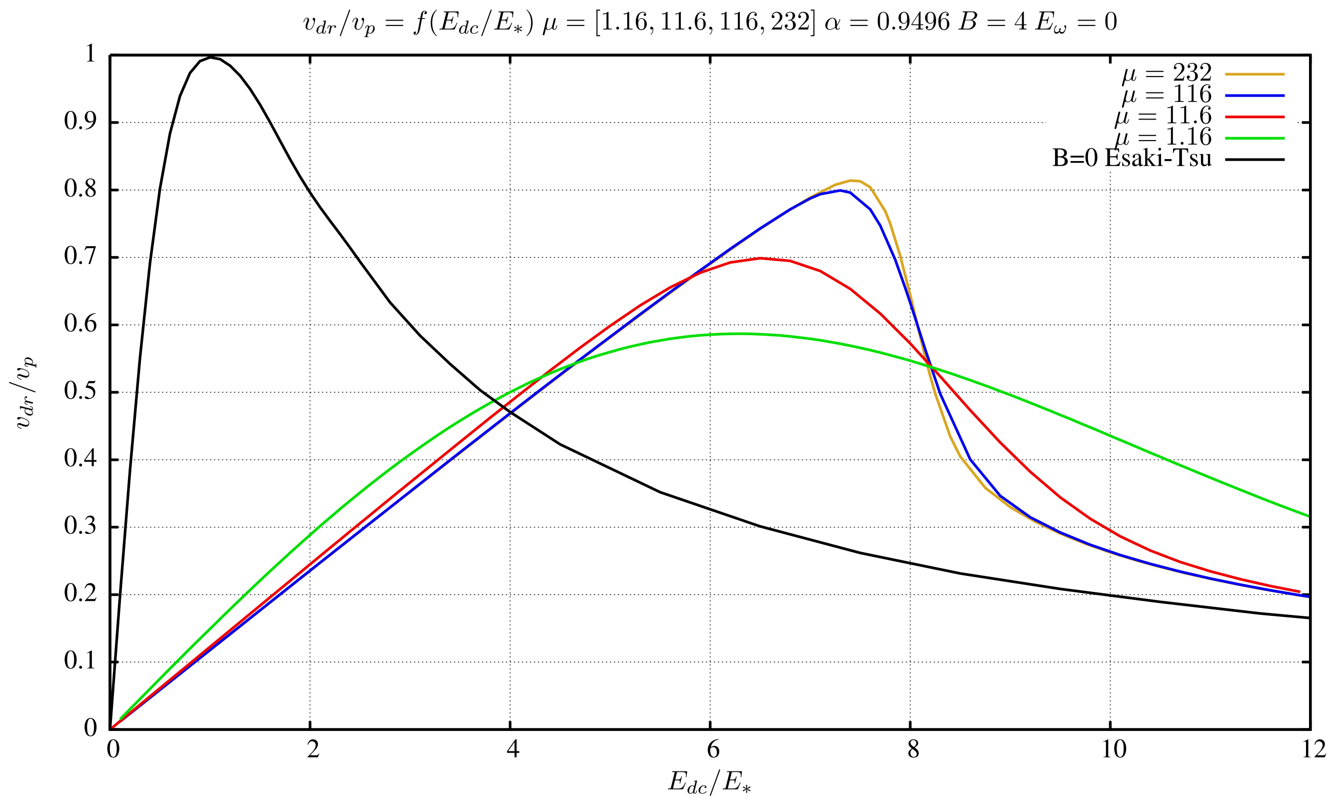
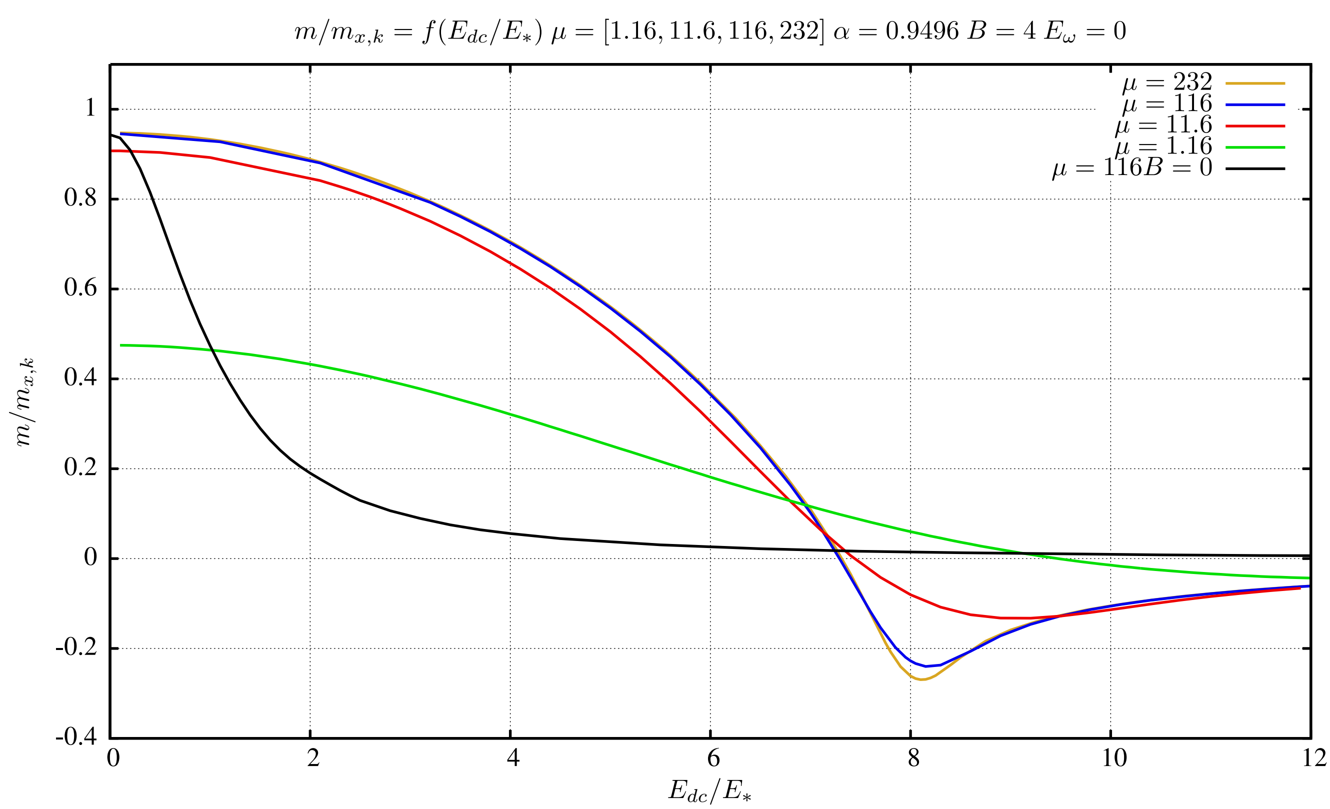


Figure 2



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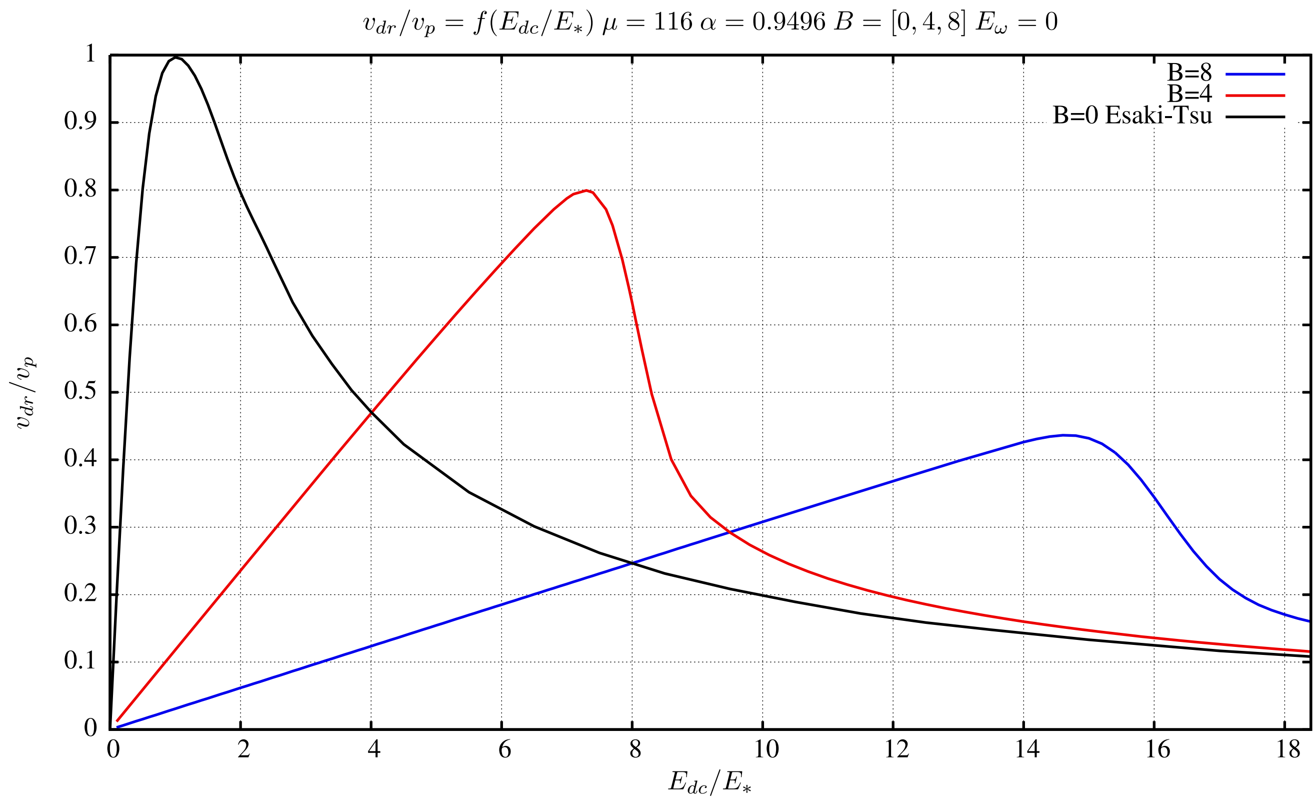

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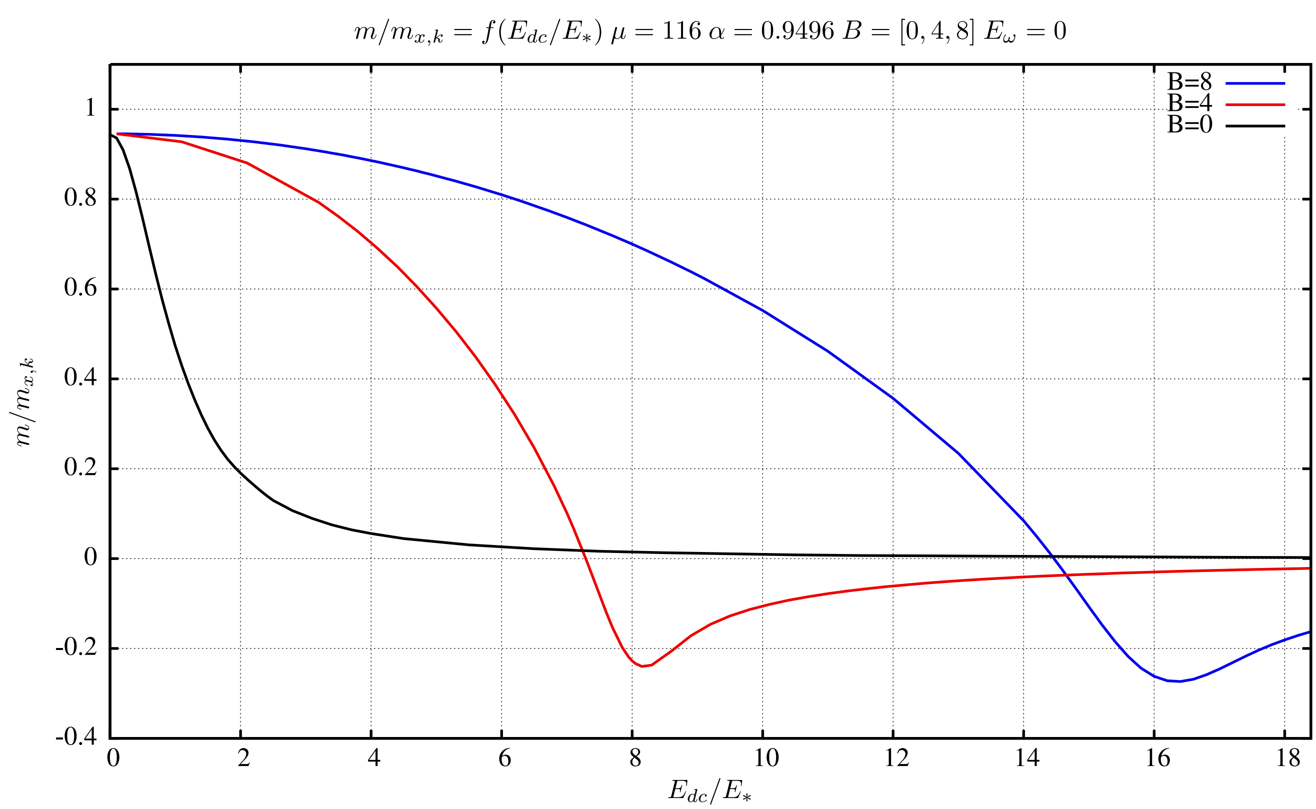

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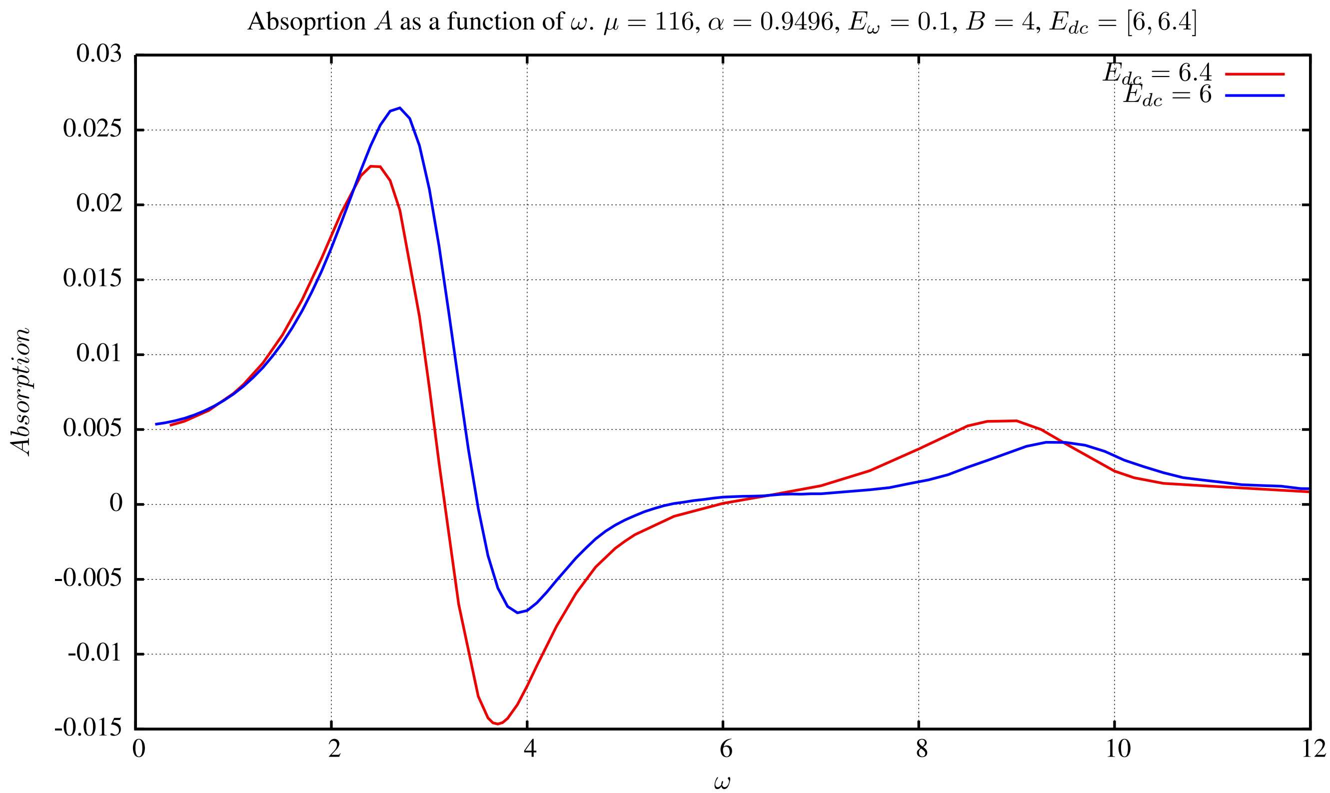

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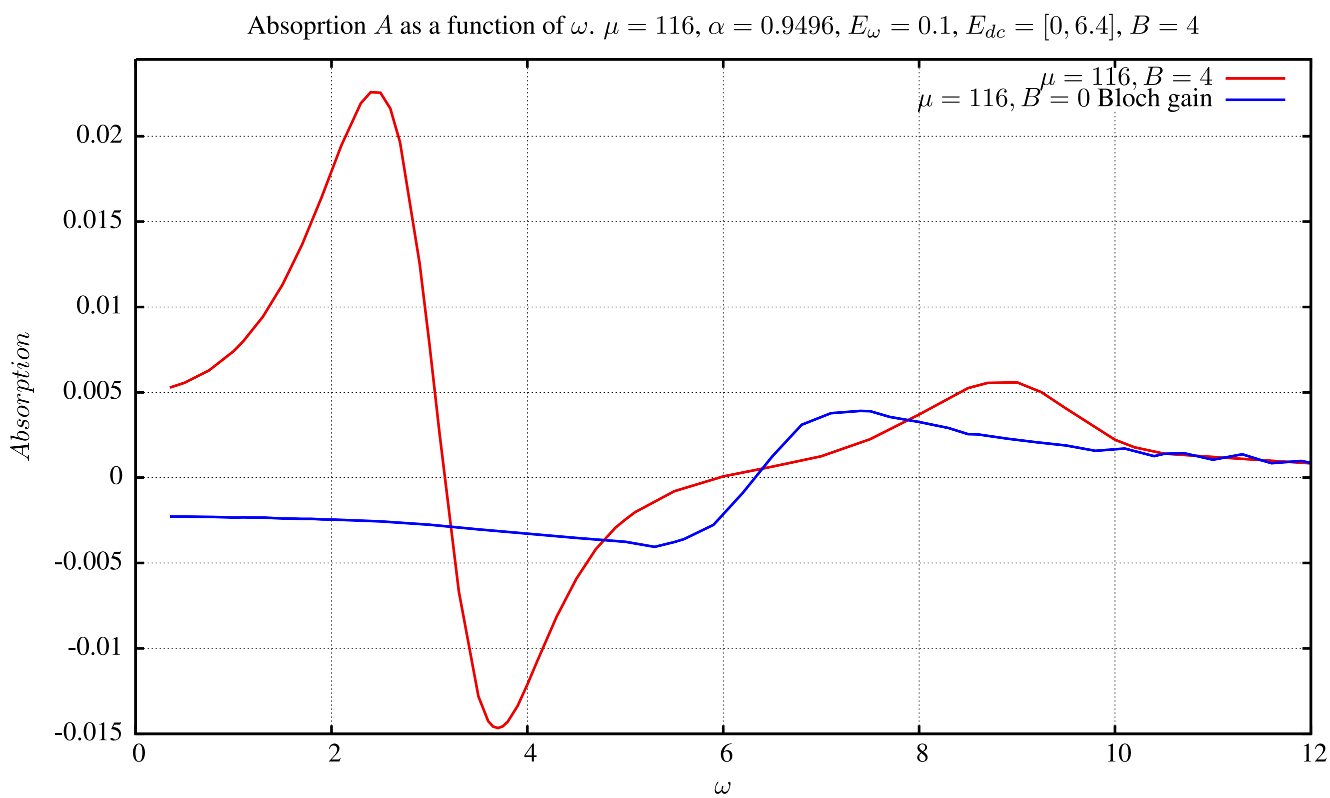


Figure 7

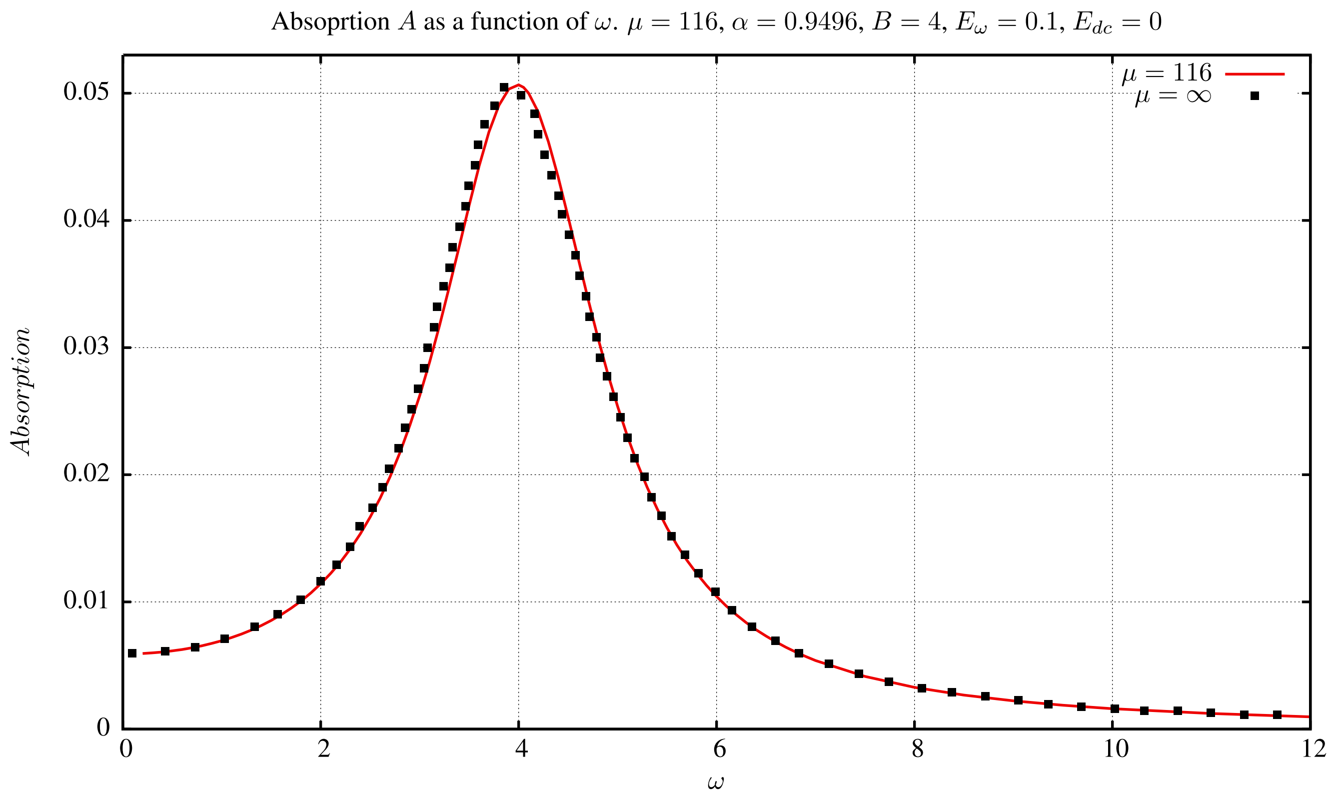


Figure 8

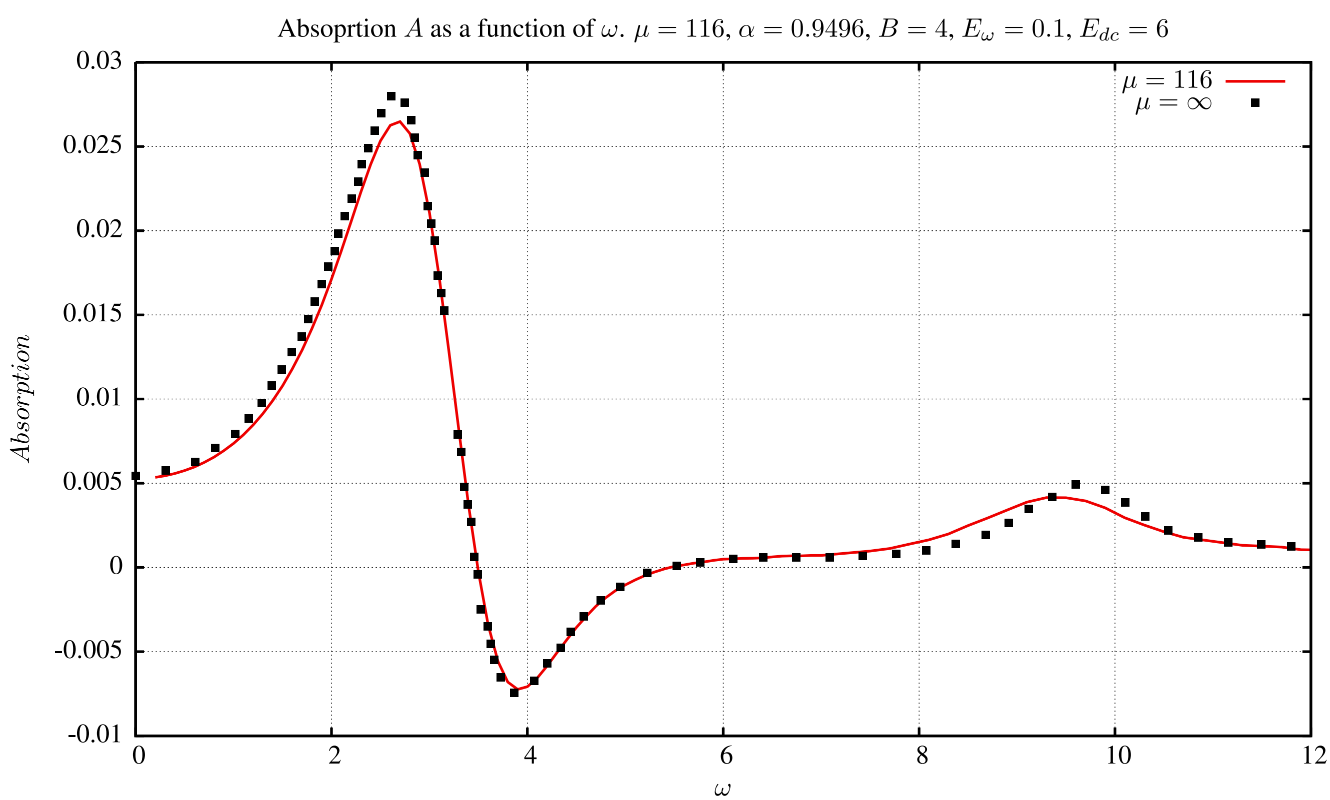
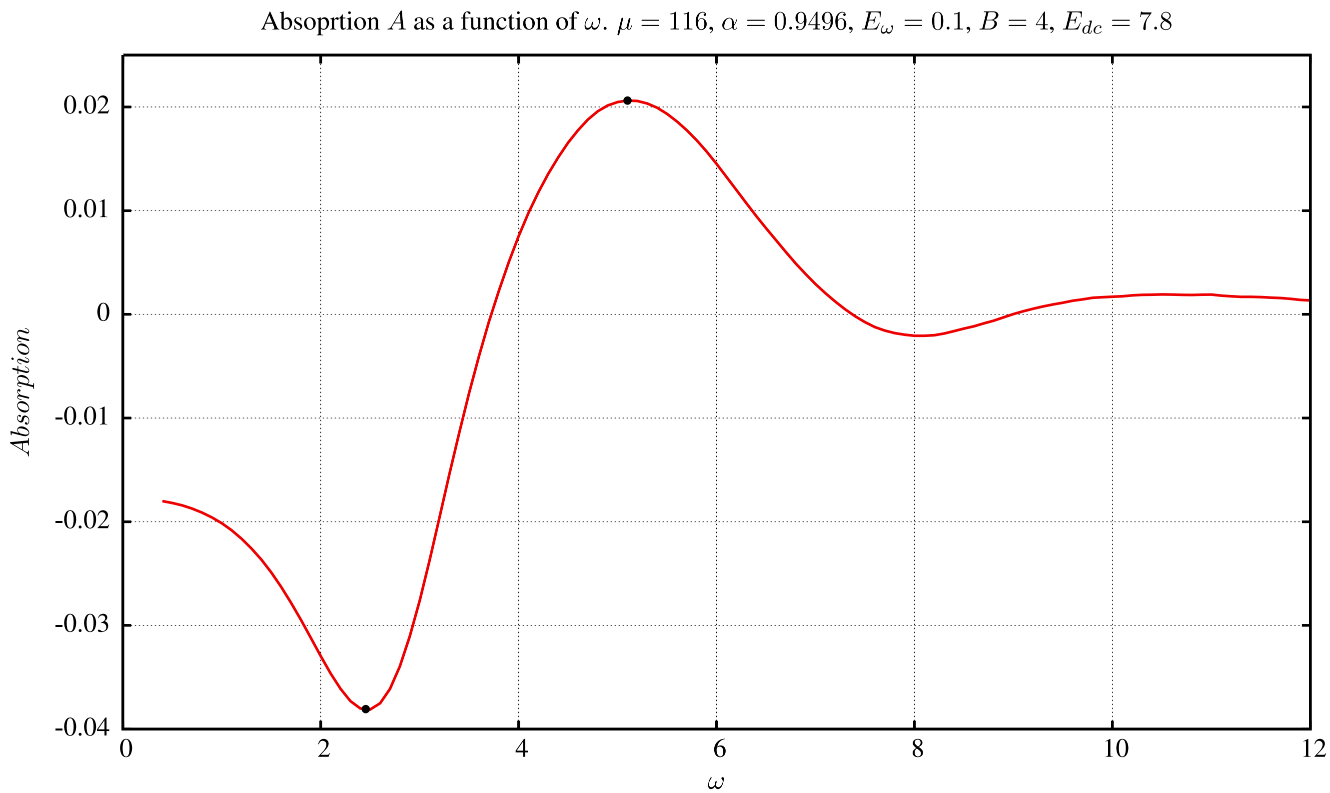


Figure 9



Absoprtion A as a function of ω. $\mu = 116$, $\alpha = 0.9496$, $E_\omega = 0.1$, $B = 4$, $E_{dc} = 7.8$
0.02
0.01
0
-0.01
-0.02
-0.03
-0.04
0
2
4
6
8
10
12
$Absorption$
ω

Figure 10

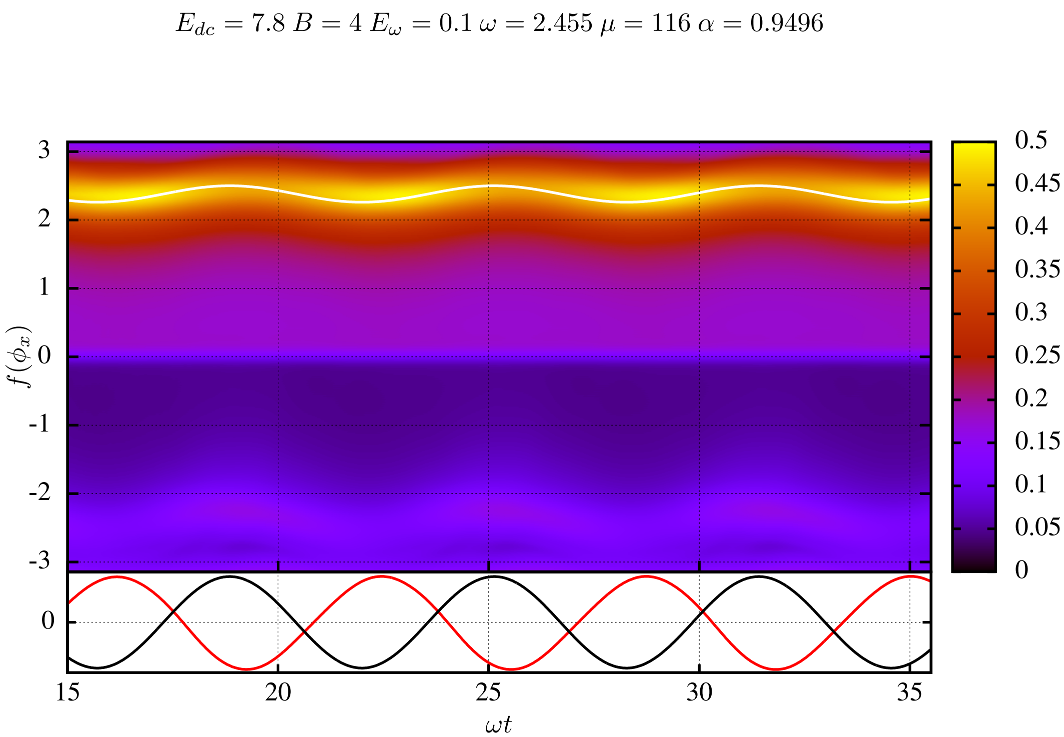


Figure 11

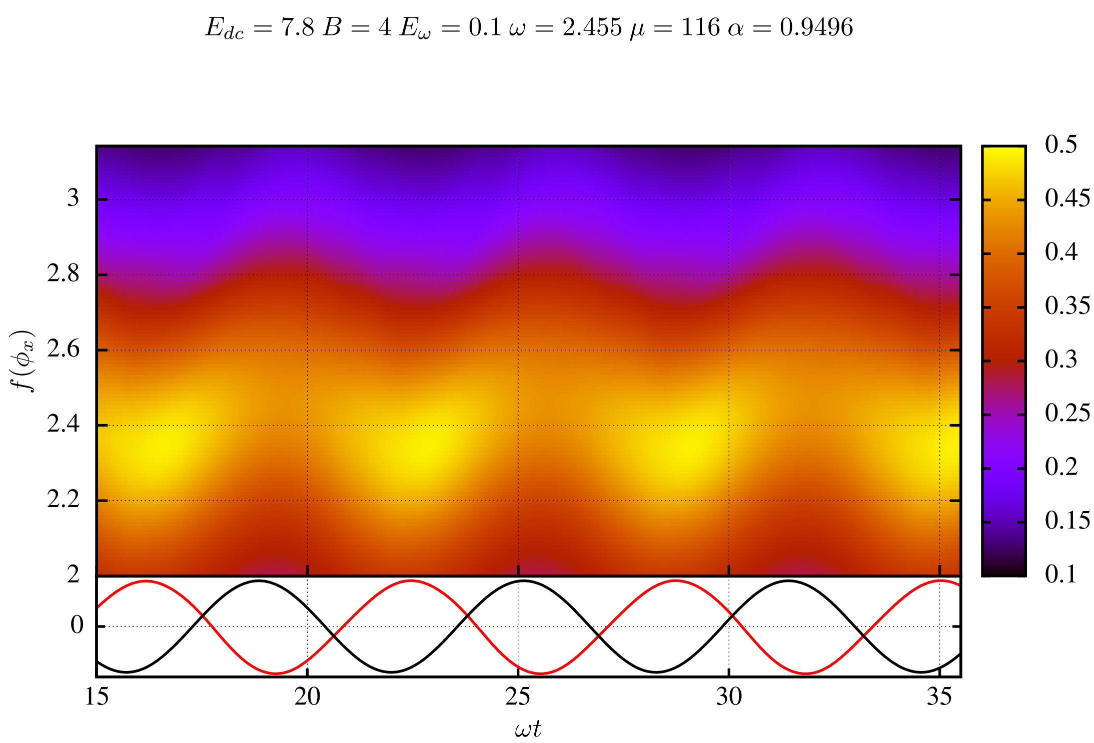


Figure 12

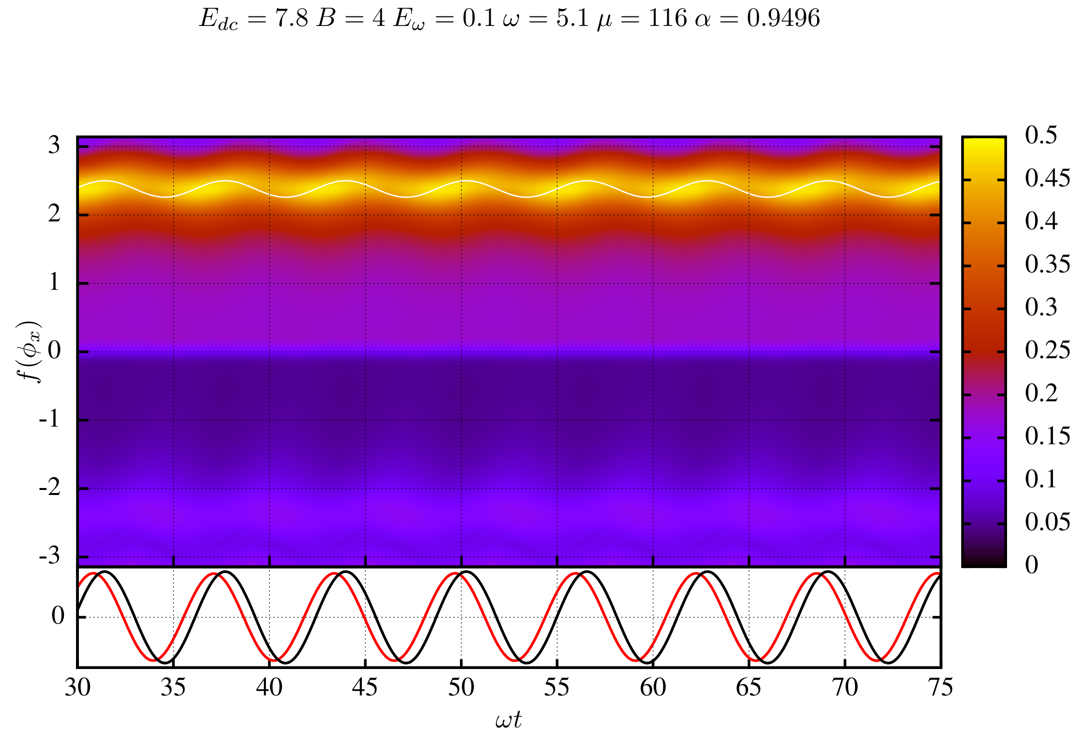


Figure 13

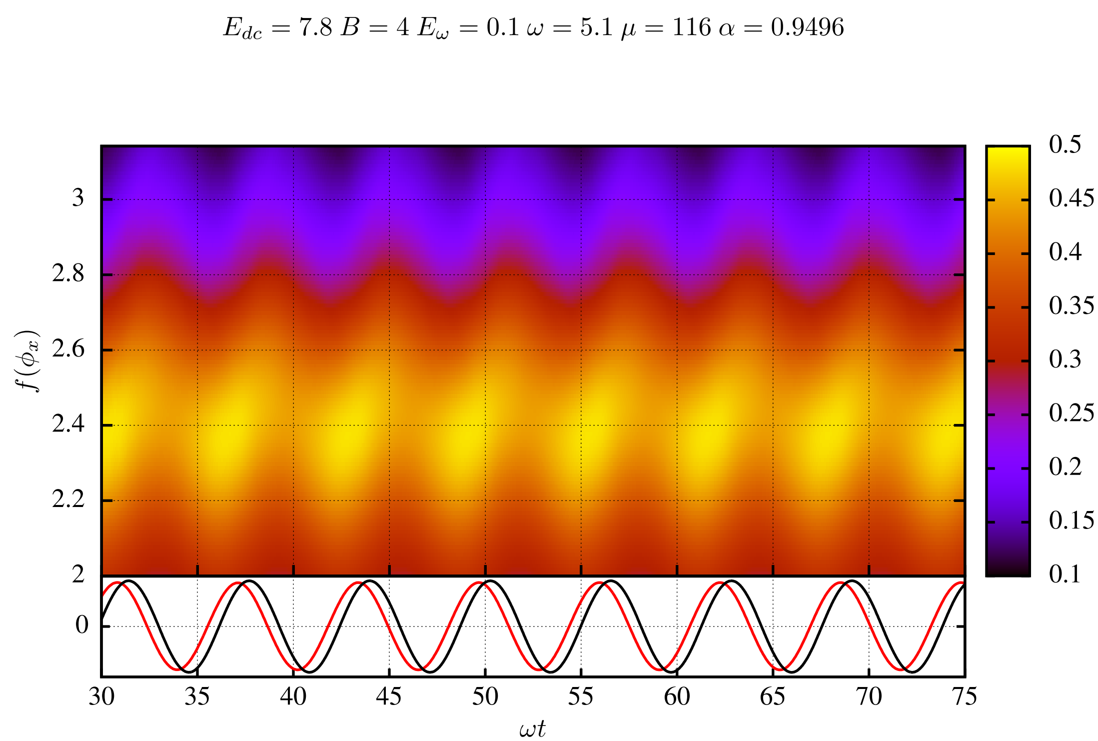


Figure 14

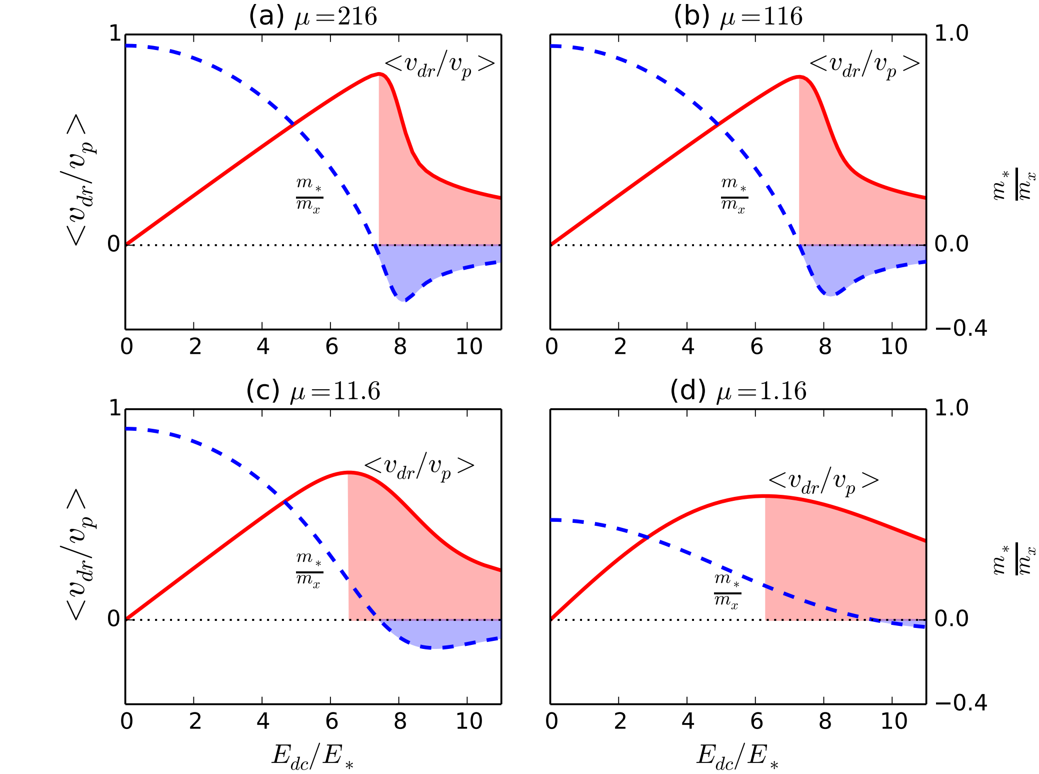


Figure 15: $E_\omega = 0.1$, $\omega = 2.455$, $B = 4$, $\alpha = 0.9496$, $\mu = [1.16, 11.6, 116, 216]$

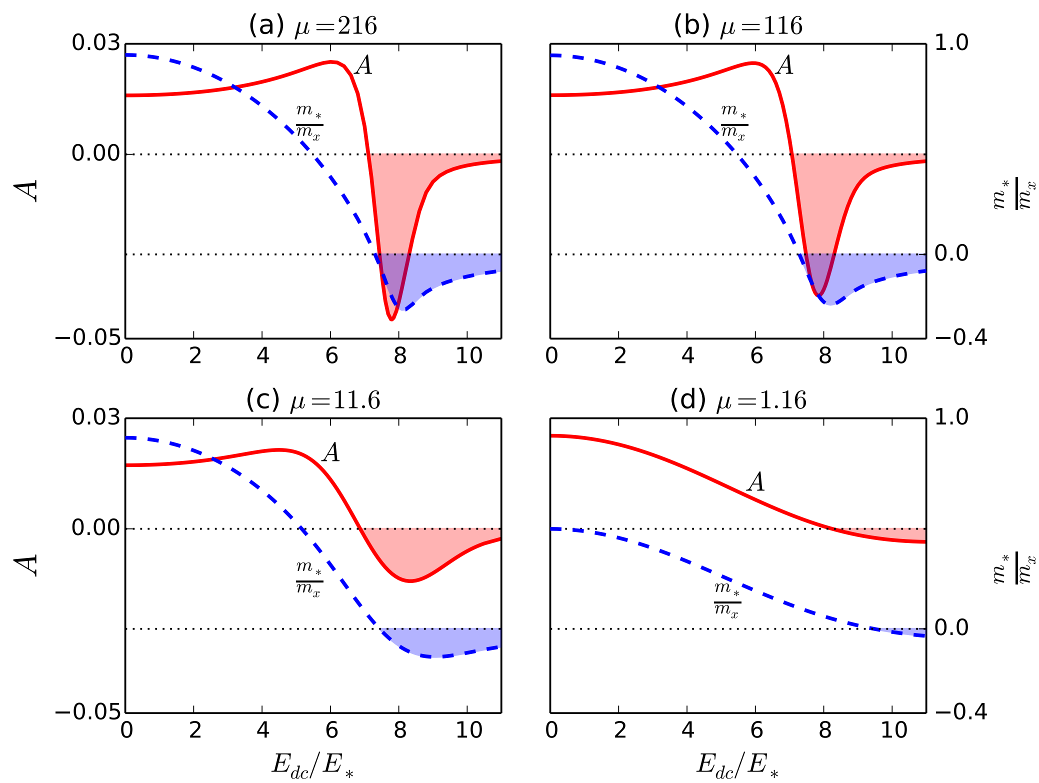


Figure 16: $E_\omega = 0.1$, $\omega = 2.455$, $B = 4$, $\alpha = 0.9496$, $\mu = [1.16, 11.6, 116, 216]$

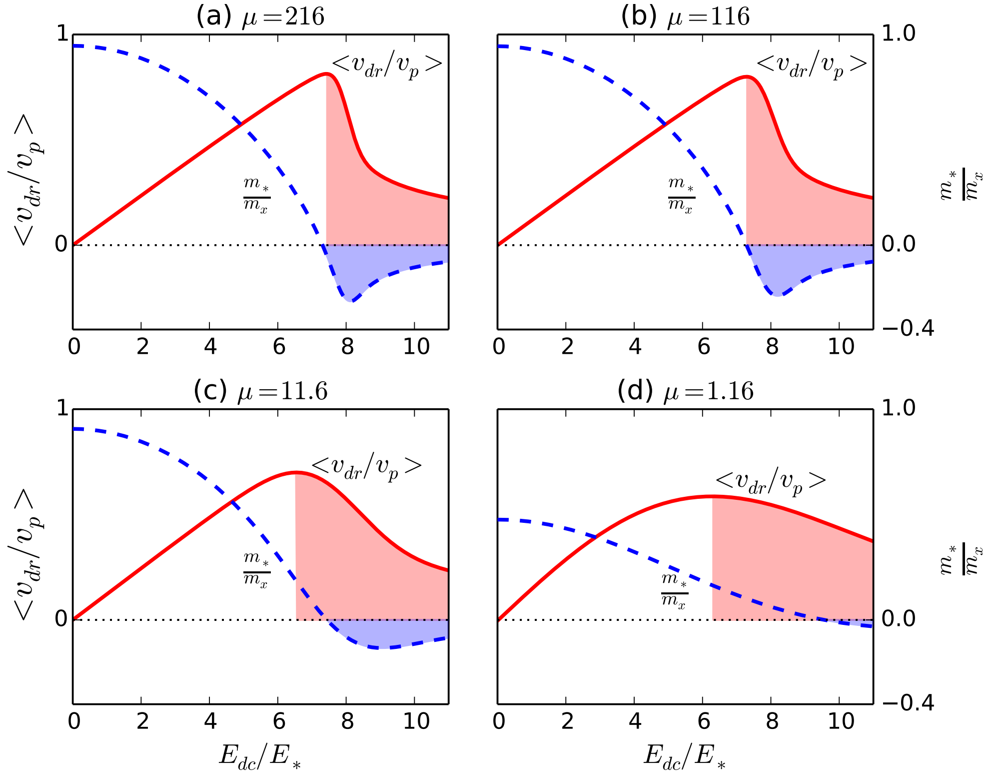


Figure 17: $E_\omega = 0.1$, $\omega = 4$, $B = 4$, $\alpha = 0.9496$, $\mu = [1.16, 11.6, 116, 216]$

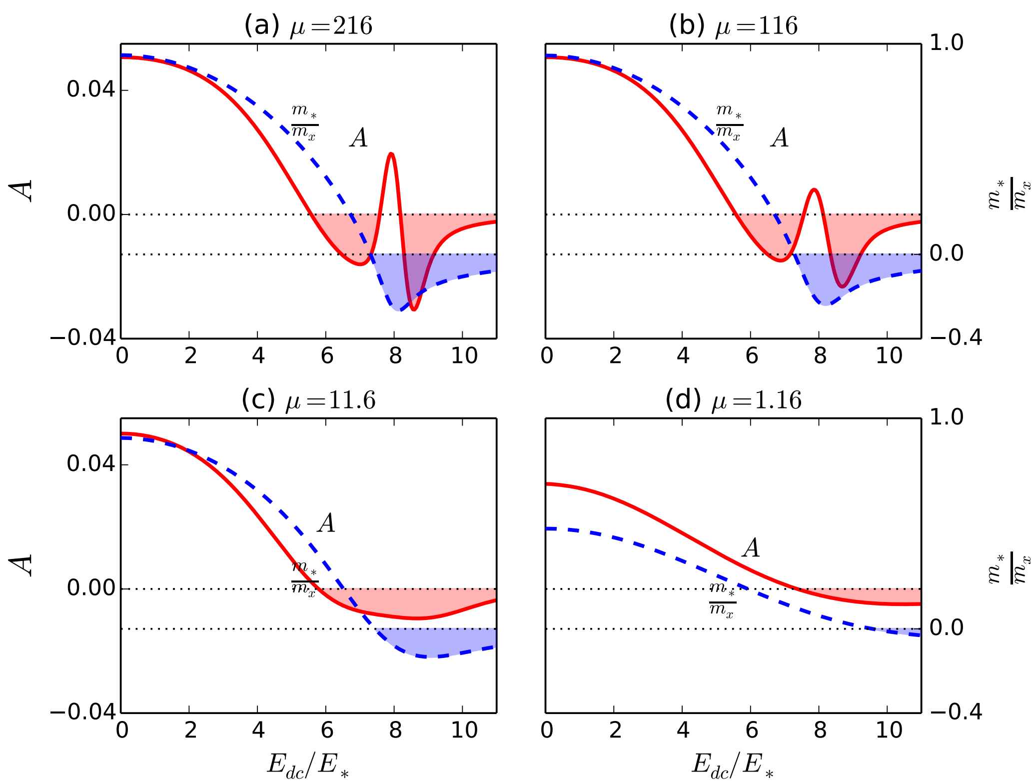


Figure 18: $E_\omega = 0.1$, $\omega = 4$, $B = 4$, $\alpha = 0.9496$, $\mu = [1.16, 11.6, 116, 216]$

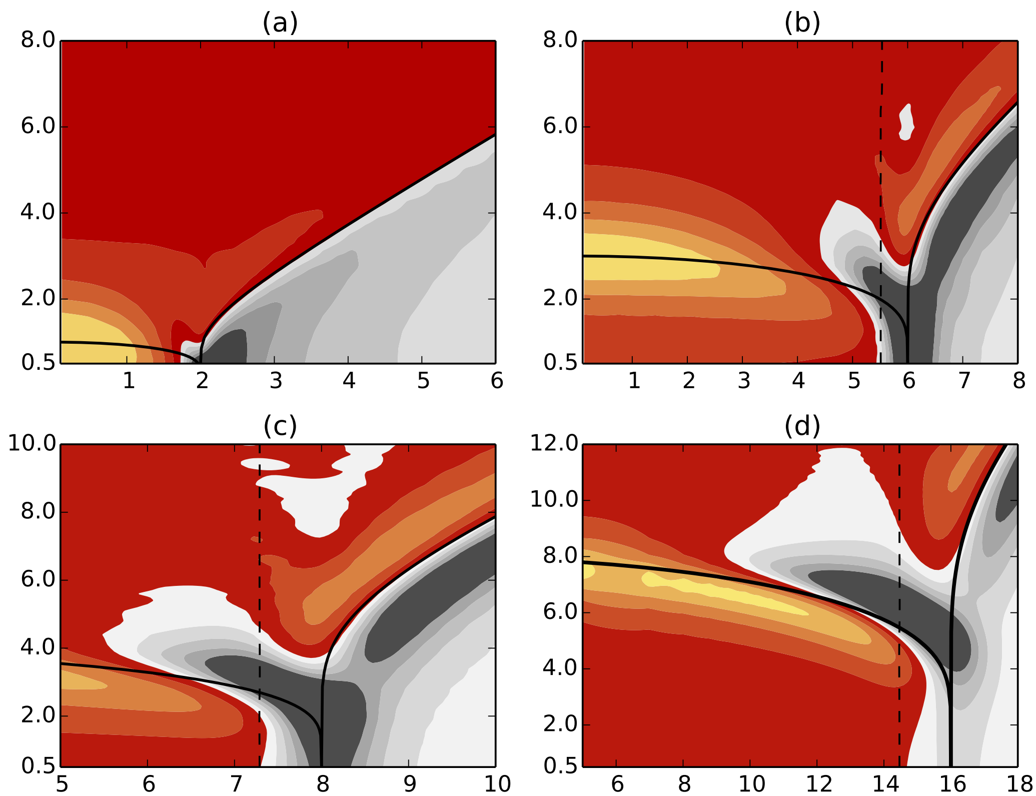


Figure 19: Maps of absoprtion for $E_\omega = 0.1$, $\alpha = 0.9496$, $\mu = 116$, $\tilde{B} = [1, 3, 4, 8]$ (plots (a), (b), (c) and (d)) as a function of $\tilde{E}$ (on the x-axis) and ω (on the y-axis)

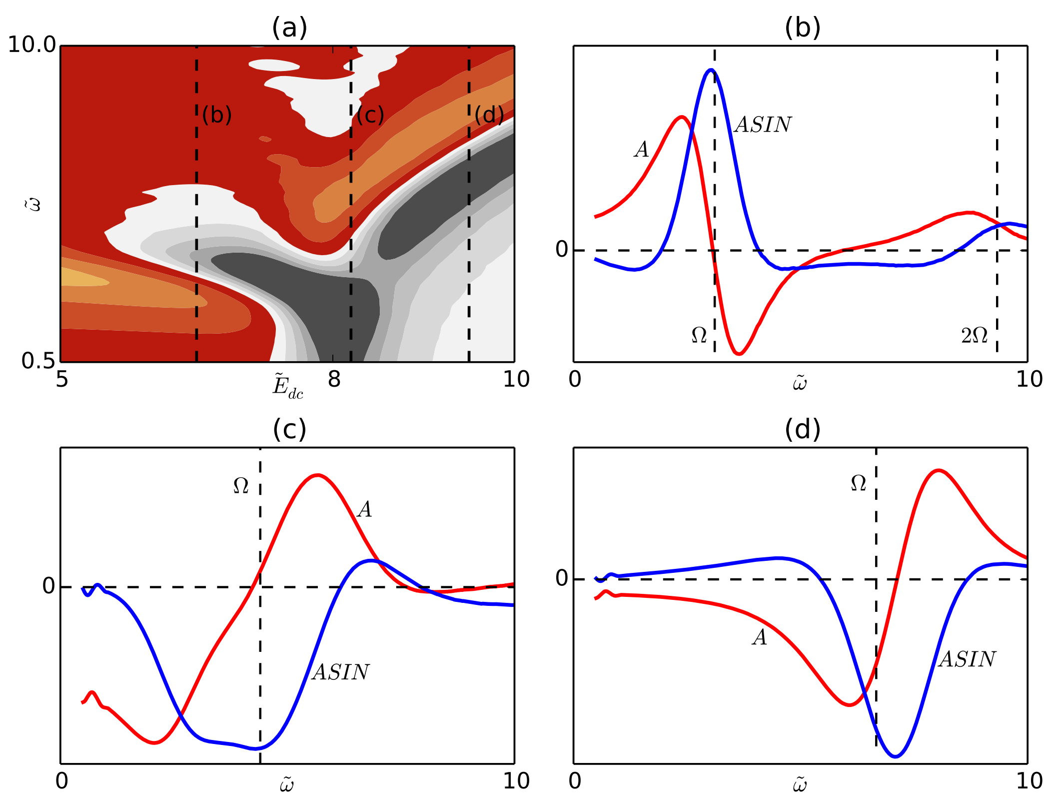


Figure 20: For map of absorption (a) for $\tilde{B} = 4$ we are looking at just the crossections of the map showing absorption A and $ASIN$ components as a functions of ω at three different values of E_{dc}. Plot (b) corresponds to $E_{dc} = 6.5$, (c) $E_{dc} = 8.2$ and (d) $E_{dc} = 9.5$