

## 1 Analytical treatment

We start from following Boltzmann equation

$$\frac{\partial f}{\partial t} + \frac{e}{\hbar} (\mathbf{E} + \mathbf{v} \times \mathbf{B}) \frac{\partial f}{\partial \mathbf{k}} + \mathbf{v}(\mathbf{k}) \frac{\partial f}{\partial r} = \left( \frac{\partial f}{\partial t} \right)_{st} \quad (1)$$

We consider case when  $f(\dots)$  is spatially homogeneous, i.e. it depends only on  $\mathbf{k}$ , then

$$\frac{\partial f}{\partial t} + \frac{e}{\hbar} (\mathbf{E} + \mathbf{v} \times \mathbf{B}) \frac{\partial f}{\partial \mathbf{k}} = \frac{f_0 - f}{\tau} \quad (2)$$

Electric field is along x-axis and magnetic field is along z-axis.

$$\mathbf{E} = (E, 0, 0) \quad (3)$$

$$\mathbf{B} = (0, 0, B) \quad (4)$$

$$\mathbf{v} \times \mathbf{B} = (v_y B, -v_x B, 0) \quad (5)$$

And will make following substitution

$$t \rightarrow \tau t \quad (6)$$

which gives following form to the Boltzmann equation

$$\frac{\partial f}{\partial t} + \frac{ed\tau}{\hbar} (E + v_y B) \frac{\partial f}{\partial (k_x d)} - \frac{ed\tau}{\hbar} v_x B \frac{\partial f}{\partial (k_y d)} = f_0 - f \quad (7)$$

where

$$\mathbf{v} = \frac{1}{\hbar} \frac{\partial \varepsilon}{\partial \mathbf{k}} \quad (8)$$

and  $\varepsilon$  is an energy of an electron in the zone, which in tight-binding approximation takes form in the first zone

$$\varepsilon = \frac{\hbar^2 k_y^2}{2m} - \frac{\Delta_1}{2} \cos(k_x d) \quad (9)$$

therefore

$$v_x = \frac{\Delta_1 d}{2\hbar} \sin(dp_x/h) \quad (10)$$

$$v_y = \frac{\hbar k_y}{m} \quad (11)$$

And Boltzmann equations now takes following form

$$\frac{\partial f}{\partial t} + \left( \frac{ed\tau}{\hbar} E + \frac{ed\tau}{m} k_y B \right) \frac{\partial f}{\partial (k_x d)} - \frac{\Delta_1 ed^2 \tau}{2\hbar^2} B \sin(k_x d) \frac{\partial f}{\partial (k_y d)} = f_0 - f \quad (12)$$

Now we can define

$$E_* = \frac{\hbar}{ed\tau} \quad (13)$$

$$m_x = \frac{2\hbar^2}{\Delta_1 d^2} \quad (14)$$

$$(15)$$

$$\boxed{\alpha = m/m_x} \quad (16)$$

Then, with cyclotron frequency defined as

$$\omega_c = \frac{eB}{\sqrt{mm_x}} \quad (17)$$

we get following

$$\frac{\partial f}{\partial t} + \left( \frac{E}{E_*} + \frac{\omega_c \tau}{\sqrt{\alpha}} k_y d \right) \frac{\partial f}{\partial(k_x d)} - \sqrt{\alpha} \omega_c \tau \sin(k_x d) \frac{\partial f}{\partial(k_y d)} = f_0 - f \quad (18)$$

And if we know define

$$\phi_x = k_x d \quad (19)$$

$$\phi_y = \frac{k_y d}{\sqrt{\alpha}} \quad (20)$$

$$\tilde{E} = E/E_* \quad (21)$$

$$\tilde{B} = \omega_c \tau \quad (22)$$

Boltzmann equations takes form

$$\boxed{\frac{\partial f}{\partial t} + (\tilde{E} + \tilde{B} \phi_y) \frac{\partial f}{\partial \phi_x} - \tilde{B} \sin(\phi_x) \frac{\partial f}{\partial \phi_y} = f_0 - f} \quad (23)$$

Alternatively, if we define Bloch frequency as

$$\omega_B = \frac{edE}{\hbar} \quad (24)$$

it can be written as

$$\frac{\partial f}{\partial t} + (\omega_B \tau + \omega_c \tau \phi_y) \frac{\partial f}{\partial \phi_x} - \omega_c \tau \sin(\phi_x) \frac{\partial f}{\partial \phi_y} = f_0 - f \quad (25)$$

Ratio of Bloch and cyclotron oscillation frequencies will be important to us down the road. For  $f_0$  we use Boltzmann distribution

$$f_0 \propto e^{-\frac{\varepsilon}{k_b T}} \quad (26)$$

which in our variables  $\phi_x$  and  $\phi_y$  takes form

$$f_0 = C \exp \left\{ \mu \cos(\phi_x) - \frac{\mu}{2} \phi_y^2 \right\} \quad (27)$$

$$\boxed{\mu = \frac{\Delta_1}{2k_b T}} \quad (28)$$

Now, constant  $C$  must be such that dimensionless norm of  $f_0$  is 1. Therefore

$$\frac{1}{C} = \frac{d^2}{\hbar^2} \int_{-\infty}^{+\infty} dp_y \int_{p_x \in \text{BZ}} dp_x \exp \left\{ \mu \cos(\phi_x) - \frac{\mu}{2} \phi_y^2 \right\} \quad (29)$$

$$= \sqrt{\alpha} \int_{-\infty}^{+\infty} e^{-\mu \phi_y^2 / 2} d\phi_y \int_{-\pi}^{\pi} e^{\mu \cos(\phi_x)} d\phi_x \quad (30)$$

$$= 2\pi I_0(\mu) \sqrt{\frac{2\pi\alpha}{\mu}} \quad (31)$$

Thus full form of equilibrium distribution is

$$\boxed{f_0 = \frac{1}{2\pi I_0(\mu)} \sqrt{\frac{\mu}{2\pi\alpha}} \exp \left\{ \mu \cos(\phi_x) - \frac{\mu}{2} \phi_y^2 \right\}} \quad (32)$$

This way, in total, we have four free parameters defining our system,  $\mu$  and  $\alpha$  that characterize lattice parameters such as lattice period and band width and temperature, and  $\tilde{E}$ ,  $\tilde{B}$  that specify external fields. One of the parameters that we will be calculating is drift velocity  $v_{dr}$ , which we will define like this

$$v_{dr} = \frac{2d}{\Delta_1 \hbar} \iint \frac{\partial \varepsilon}{\partial p_x} f(p_x, p_y) dp_x dp_y \quad (33)$$

$$= \sqrt{\alpha} \int_{-\pi}^{\pi} d\phi_x \int_{-\infty}^{+\infty} d\phi_y \sin(\phi_x) f(\phi_x, \phi_y) \quad (34)$$

When magnetic field is 0 and  $E$  is constant, i.e.  $E_{dc}$ , analytic solution Boltzmann equation is well known and as well as expression for drift velocity

$$v_{dr} = \left\{ \sqrt{\alpha} \int_{-\pi}^{\pi} d\phi_x \int_{-\infty}^{+\infty} d\phi_y \sin(\phi_x) f_0(\phi_x, \phi_y) \right\} \frac{E_{dc}/E_*}{1 + (E_{dc}/E_*)^2} \quad (35)$$

$$= \frac{I_1(\mu)}{I_0(\mu)} \frac{E_{dc}/E_*}{1 + (E_{dc}/E_*)^2} \quad (36)$$

From which follows that peak value of  $v_{dr}$  is at  $E_{dc}/E_* = 1$  and is

$$v_p = \frac{I_1(\mu)}{2I_0(\mu)} \quad (37)$$

Later, down the road, we will be plotting not  $v_{dr}$ , but  $v_{dr}/v_p$ , which for dc electric field only take very simple form, known as "Esaki-Tsu" equation

$$\frac{v_{dr}}{v_p} = 2 \frac{E_{dc}/E_*}{1 + (E_{dc}/E_*)^2} \quad (38)$$

In general we will be applying a/c emf in the form

$$\tilde{E} = \tilde{E}_{dc} + \tilde{E}_{\omega} \cos(\omega t) \quad (39)$$

and in case when magnetic field is not applied we also have analytic expression for  $v_{dr}$ , which is known as "Tien-Gordon" equation. Analytic expression, "Taker formulae", is also known for absorption, which we will define like this

$$A = \left\langle \frac{v_{dr}}{v_p} \cos(\omega t) \right\rangle_t \quad (40)$$

In case when magnetic field is applied, however, analytic expression for absorption is not known. And that is the quantity we are most interested in.

Now due to periodicity of  $f(\phi_x, \phi_y)$  along  $\phi_x$  with period  $2\pi$  and additionally  $f_0(-\phi_x, \phi_y) = f_0(\phi_x, \phi_y)$ , it makes sense for us to expand  $f$  and  $f_0$  into Fourier series.

$$f_0 = \sum_{n=0}^{\infty} a_n^{(0)} \cos(n\phi_x) \quad (41)$$

$$f = \sum_{n=0}^{\infty} a_n \cos(n\phi_x) + b_n \sin(n\phi_x) \quad (42)$$

where coefficients  $a_n^{(0)}$ ,  $a_n$  and  $b_n$  in general will depend on  $\phi_y$  and last two also on time  $t$ . Note, that  $a_{n<0} \equiv 0$  and  $b_{n<1} \equiv 0$ . And with the form of  $f_0$ , as selected above,  $a_n^{(0)}$  becomes

$$a_n^{(0)} = \frac{\sigma(n)}{\pi} \int_{-\pi}^{\pi} f_0(\phi_x, \phi_y) \cos(n\phi_x) d\phi_x \quad (43)$$

$$= \frac{\sigma(n) I_n(\mu)}{\pi I_0(\mu)} \sqrt{\frac{\mu}{2\pi\alpha}} \exp \left\{ -\frac{\mu}{2} \phi_y^2 \right\} \quad (44)$$

where

$$\sigma(n) = \begin{cases} 1/2 & : n = 0 \\ 1 & : n \neq 0 \end{cases} \quad (45)$$

Norm of  $f(\phi_x, \phi_y)$  must always be one, i.e.

$$\sqrt{\alpha} \int_{-\pi}^{\pi} d\phi_x \int_{-\infty}^{+\infty} d\phi_y f(\phi_x, \phi_y) = 1 \quad (46)$$

which in fourier representation takes form

$$\boxed{2\pi\sqrt{\alpha} \int_{-\infty}^{+\infty} a_0(\phi_y) d\phi_y = 1} \quad (47)$$

Once we move on to the numerical calculations this equation can be used to check correctness. And from equation (34) it is clear that only  $b_1$  will survive. And calculation of  $v_{dr}$  is done through following equation

$$v_{dr} = \sqrt{\alpha} \int_{-\infty}^{+\infty} d\phi_y \int_{-\pi}^{+\pi} d\phi_x \sin(\phi_x) b_1(\phi_y) \sin(\phi_y) \quad (48)$$

$$= \pi\sqrt{\alpha} \int_{-\infty}^{+\infty} b_1(\phi_y) d\phi_y \quad (49)$$

and in view of definition of peak value of  $v_{dr}$  in eq. (37)

$$\boxed{\frac{v_{dr}}{v_p} = \frac{2I_0(\mu)\pi\sqrt{\alpha}}{I_1(\mu)} \int_{-\infty}^{+\infty} b_1(\phi_y) d\phi_y} \quad (50)$$

In addition to drift velocity along  $x$ -axis we can look at drift velocity along  $y$ -axis, which we will define like this

$$v_y = \frac{2}{\Delta_1} \iint \frac{\partial \varepsilon}{\partial p_y} f(\phi_x, \phi_y) dp_x dp_y \quad (51)$$

$$= \int_{-\pi}^{\pi} d\phi_x \int_{-\infty}^{+\infty} \phi_y f(\phi_x, \phi_y) d\phi_y \quad (52)$$

$$= 2\pi \int_{-\infty}^{+\infty} a_0(\phi_y) \phi_y d\phi_y \quad (53)$$

However just as with drift velocity along  $x$ -axis we will be working with ratio of  $v_y$  and peak velocity  $v_p$ .

$$\boxed{\frac{v_y}{v_p} = \frac{4\pi I_0(\mu)}{I_1(\mu)} \int_{-\infty}^{+\infty} a_0(\phi_y) \phi_y d\phi_y} \quad (54)$$

This way we accounted for meaning of  $a_0$  and  $b_1$ . Let us now take a look at  $a_1$ . Effective mass of an electron in the  $x$  direction is given by

$$m_{x,\mathbf{k}}^{-1} = \frac{1}{\hbar^2} \frac{\partial^2 \varepsilon}{\partial k_x^2} \quad (55)$$

We will be working however with ratio of electron mass to  $m_{x,\mathbf{k}}$ . And in our approximation of energy (9) that takes form

$$\frac{m}{m_{x,\mathbf{k}}} = \frac{\Delta_1 d^2 m}{2\hbar^2} \cos(k_x d) \quad (56)$$

$$= \alpha \cos(\phi_x) \quad (57)$$

To gather back this value from  $f(\phi_x, \phi_y)$  we have to integrate over  $\{p_x, p_y\}$  and to maintain dimensionlessness of ratio of these masses, this integration will take form

$$\frac{m}{m_{x,\mathbf{k}}} = \alpha^{3/2} \frac{d^2}{\hbar^2} \iint f(\phi_x, \phi_y) \cos(k_x d) d p_x d p_y \quad (58)$$

$$= \alpha^{3/2} \int_{-\pi}^{\pi} d\phi_x \int_{-\infty}^{+\infty} d\phi_y f(\phi_x, \phi_y) \cos(\phi_x) \quad (59)$$

$$= \alpha^{3/2} \int_{-\pi}^{\pi} d\phi_x \int_{-\infty}^{+\infty} d\phi_y a_1(\phi_y) \cos(\phi_x) \cos(\phi_x) \quad (60)$$

Finally giving us following

$$\boxed{\frac{m}{m_{x,\mathbf{k}}} = \pi \alpha^{3/2} \int_{-\infty}^{+\infty} a_1(\phi_y) d\phi_y} \quad (61)$$

Now using fourier representation of  $f(\phi_x, \phi_y)$  and  $f_0(\phi_y)$  (equations 41, 42) we can rewrite Boltzmann equation (23) like this

$$\sum_{(n)} \left\{ \frac{\partial a_n}{\partial t} \cos(n\phi_x) + \frac{\partial b_n}{\partial t} \sin(n\phi_x) = a_n^{(0)} \cos(n\phi_x) - a_n \cos(n\phi_x) - b_n \sin(n\phi_x) + \right. \\ \left. n(\tilde{E} + \tilde{B}\phi_y)(a_n \sin(n\phi_x) - b_n \cos(n\phi_x)) + \tilde{B} \frac{\partial a_n}{\partial \phi_y} \sin(\phi_x) \cos(n\phi_x) + \tilde{B} \frac{\partial b_n}{\partial \phi_y} \sin(\phi_x) \sin(n\phi_x) \right\} \quad (62)$$

In absence of magnetic field there is no mixing of different harmonic, however when  $\tilde{B}$  is not 0 then harmonics will be come mixed due to presence of  $\sin(\phi_x) \cos(n\phi_x)$  and  $\sin(\phi_x) \sin(n\phi_x)$ , since

$$\sin(\phi_x) \cos(n\phi_x) = \{\sin((n+1)\phi_x) - \sin((n-1)\phi_x)\} / 2 \quad (63)$$

$$\sin(\phi_x) \sin(n\phi_x) = \{\cos((n-1)\phi_x) - \cos((n+1)\phi_x)\} / 2 \quad (64)$$

And using this equations, after some manipulation of symbols, combining elements with the same harmonics, and noting special treatment of  $b_1$  we get

$$\frac{\partial a_n}{\partial t} = a_n^{(0)} - a_n - n(\tilde{E} + \tilde{B}\phi_y)b_n + \tilde{B} \left( \frac{\partial b_{n+1}}{\partial \phi_y} - \frac{\partial b_{n-1}}{\partial \phi_y} \right) \quad (65)$$

$$\frac{\partial b_n}{\partial t} = -b_n - n(\tilde{E} + \tilde{B}\phi_y)a_n + \tilde{B} \left( \chi(n) \frac{\partial a_{n-1}}{\partial \phi_y} - \frac{\partial a_{n+1}}{\partial \phi_y} \right) \quad (66)$$

where

$$\chi(n) = \begin{cases} 2 & : n = 1 \\ 1 & : n \neq 1 \end{cases} \quad (67)$$

## 2 Correspondence with classical pendulum

In the limit where dissipation is absent instead of Boltzmann equation we can use semiclassical equations of motion working individual electrons. Such situation corresponds to  $\tau = \infty$  and initial shape of  $f(\phi_x, \phi_y)$  being delta function.

$$\hbar \frac{d\mathbf{k}}{dt} = e\mathbf{E} + e\mathbf{v} \times \mathbf{B} \quad (68)$$

$$\mathbf{v}(\mathbf{k}) = \frac{1}{\hbar} \frac{\partial \varepsilon}{\partial \mathbf{k}} \quad (69)$$

And in view of specific values of  $\mathbf{E}$ ,  $\mathbf{B}$  and  $\varepsilon$  that we are using in our problem, these equations transform into following

$$\frac{d\phi_x}{dt} = \tilde{E} + \tilde{B}\phi_y \quad (70)$$

$$\frac{d\phi_y}{dt} = -\tilde{B}\sin(\phi_x) \quad (71)$$

And taking second derivative of (70), assuming that  $\tilde{E}$  is defined by (39) we arrive to second order ODE, which correspond to classical externally driven pendulum.

$$\frac{d^2\phi_x}{dt^2} + \tilde{B}^2 \sin(\phi_x) = -\omega \tilde{E}_\omega \sin(\omega t) \quad (72)$$

In general this equation can exhibit very complicated behaviour and even chaos, but can be solved in the simplest case when  $\tilde{E}_\omega = 0$ , in which case period of oscillation is defined by magnetic field  $\tilde{B}$  and electric field  $\tilde{E}$ , which sets initial velocity of  $\phi_x$ , i.e.  $\frac{d\phi_x}{dt}$  at  $t = 0$ , assuming that  $\phi_x$  and  $\phi_y$  at  $t = 0$  are also 0. In case of  $E_\omega = 0$  we exact correspondence of our system to a classical pendulum, for which we do have conservation of energy.

$$\frac{1}{2} \left( \frac{d\phi_x}{dt} \right)^2 + 2\tilde{B}^2 \sin^2(\phi_x/2) = E_{total} \quad (73)$$

If we define  $E_p = 2\tilde{B}^2$ , which corresponds to "peak" then when  $E_{total} = E_p$  then such pendulum can reach inverted position. When  $\phi_x$  and  $\phi_y$  are equal zero  $E_{total} = \tilde{E}_{dc}^2/2$  and thus for  $E_p$  there is corresponding  $E_{sx} = 2\tilde{B}$ , which defines separatrix dividing motion of the pendulum between vibrational ( $\tilde{E}_{dc} < E_{sx}$ ) and rotational ( $\tilde{E}_{dc} > E_{sx}$ ). When our system gives peak of negative absorption corresponding to the first and second cases we speak of "Cyclotron" and "Bloch" resonances respectively. Expression for frequency  $\Omega$  of these natural oscillations is well known and in our variables it takes form

$$\Omega = \frac{\pi E_{sx}}{4K(\tilde{E}_{dc}/E_{sx})} \quad \text{for } \tilde{E}_{dc} < E_{sx} \quad (74)$$

$$\Omega = \frac{\pi E_{dc}}{2K(E_{sx}/\tilde{E}_{dc})} \quad \text{for } \tilde{E}_{dc} > E_{sx} \quad (75)$$

where  $K(x)$  is a complete elliptic integral of the first kind.

### 3 Numerical solution

Straightforward application of method of finite differences to (65) and (66) leads to either unstable or i.e. computationally intensive equations. To combat this problem we are using several methods at once. First, we are going to discretize  $a_n$  and  $b_n$  along time and  $\phi_y$  axes.

$$\begin{aligned} t &\leftarrow \text{time step} \\ a_{n,m} &\leftarrow \phi_y \text{ lattice step} \end{aligned} \quad (76)$$

and  $n$  is "harmonic number". So, here we are going to do some tricky things. We are going to write two forms of equations (65, 66). One using forward differences and one using partial backward differences, i.e. on the right side of equal sign we are going to write partial derivatives at time  $t$  while everything else at time  $t+1$  and will follow standard procedure of CrankNicolson scheme by adding these two, forward and backward differences equations. First, forward differencing scheme

$$\begin{aligned} a_{n,m}^{t+1} - a_{n,m}^t &= a_{n,m}^{(0)} \Delta t - a_{n,m}^t \Delta t - 2b_{n,m}^t \mu_{n,m}^t + \\ &+ \frac{\alpha B \Delta t}{2\Delta\phi} (b_{n+1,m+1}^t - b_{n+1,m-1}^t - b_{n-1,m+1}^t + b_{n-1,m-1}^t) \end{aligned} \quad (77)$$

$$\begin{aligned} b_{n,m}^{t+1} - b_{n,m}^t &= -b_{n,m}^t \Delta t + 2a_{n,m}^t \mu_{n,m}^t + \\ &+ \frac{\alpha B \Delta t}{2\Delta\phi} (\chi(n)[a_{n-1,m+1}^t - a_{n-1,m-1}^t] - a_{n+1,m+1}^t + a_{n+1,m-1}^t) \end{aligned} \quad (78)$$

And then partial backward differencing scheme

$$a_{n,m}^{t+1} - a_{n,m}^t = a_{n,m}^{(0)} \Delta t - a_{n,m}^{t+1} \Delta t - 2b_{n,m}^{t+1} \mu_{n,m}^{t+1} + \frac{\alpha B \Delta t}{2\Delta\phi} (b_{n+1,m+1}^t - b_{n+1,m-1}^t - b_{n-1,m+1}^t + b_{n-1,m-1}^t) \quad (79)$$

$$b_{n,m}^{t+1} - b_{n,m}^t = -b_{n,m}^{t+1} \Delta t + 2a_{n,m}^{t+1} \mu_{n,m}^{t+1} + \frac{\alpha B \Delta t}{2\Delta\phi} (\chi(n)[a_{n-1,m+1}^t - a_{n-1,m-1}^t] - a_{n+1,m+1}^t + a_{n+1,m-1}^t) \quad (80)$$

where

$$\beta_m^t = E^t + B^t \phi_y(m) \quad (81)$$

$$\mu_{n,m}^t = n \beta_m^t \Delta t \quad (82)$$

And application of CrankNicolson scheme leads to

$$a_{n,m}^{t+1} = \frac{g_{n,m}^t \nu - h_{n,m}^t \mu_{n,m}^{t+1}}{\nu^2 + (\mu_{n,m}^{t+1})^2} \quad (83)$$

$$b_{n,m}^{t+1} = \frac{g_{n,m}^t \mu_{n,m}^{t+1} - h_{n,m}^t \nu}{\nu^2 + (\mu_{n,m}^{t+1})^2} \quad (84)$$

where

$$\nu = 1 + \Delta t/2 \quad (85)$$

$$\xi = 1 - \Delta t/2 \quad (86)$$

$$g_{n,m}^t = a_{n,m}^t \xi - b_{n,m}^t \mu_{n,m}^t + A_{n,m}^t + a_{n,m}^{(0)} \Delta t \quad (87)$$

$$h_{n,m}^t = b_{n,m}^t \xi + a_{n,m}^t \mu_{n,m}^t + B_{n,m}^t \quad (88)$$

$$A_{n,m}^t = \frac{\alpha B \Delta t}{2\Delta\phi} (\chi(n)[a_{n-1,m+1}^t - a_{n-1,m-1}^t] - a_{n+1,m+1}^t + a_{n+1,m-1}^t) \quad (89)$$

$$B_{n,m}^t = \frac{\alpha B \Delta t}{2\Delta\phi} (b_{n+1,m+1}^t - b_{n+1,m-1}^t - b_{n-1,m+1}^t + b_{n-1,m-1}^t) \quad (90)$$

These equation (83, 84) can be formally written in the form

$$\mathbf{r}_{n,m}^{t+1} = \mathbf{T}(\mathbf{r}_{n,m}^t; A_{n,m}^t, B_{n,m}^t) \quad (91)$$

$$\mathbf{r}_{n,m}^t = (a_{n,m}^t, b_{n,m}^t) \quad (92)$$

Where  $\mathbf{T}$  is an operator that allows us to step from time step  $t$  to  $t + 1$ , separated by time interval  $\Delta t$ . Using this operation as is leads to, only, conditionally stable numerical system, because  $A_{n,m}^t$  and  $B_{n,m}^t$  are taken at time  $t$ , which means that we have to take time step  $dt$  to be very, very small. To make time step much larger without solving implicit equations, to calculate  $a_{n,m}^{t+1}$  and  $b_{n,m}^{t+1}$  we will utilize variation of leap-frog method, by using two staggered grids.

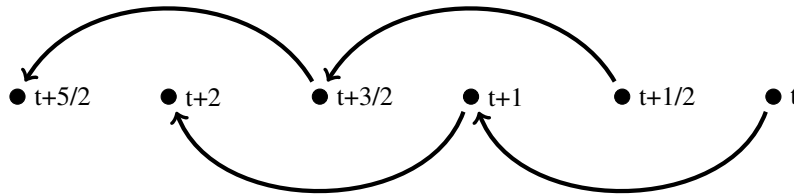


Figure 1

And writing (91) in the form

$$\mathbf{r}_{n,m}^{t+1} = \mathbf{T}(\mathbf{r}_{n,m}^t; A_{n,m}^{t+1/2}, B_{n,m}^{t+1/2}) \quad (93)$$

$$\mathbf{r}_{n,m}^{t+3/2} = \mathbf{T}(\mathbf{r}_{n,m}^{t+1/2}; A_{n,m}^{t+1}, B_{n,m}^{t+1}) \quad (94)$$

Use of equations (93) and (94) is indicated by lower and upper arrows in fig 1. To start this system we calculate  $\mathbf{r}_{n,m}^{1/2}$  by using eq. (91) and then shift to leap-frog method.



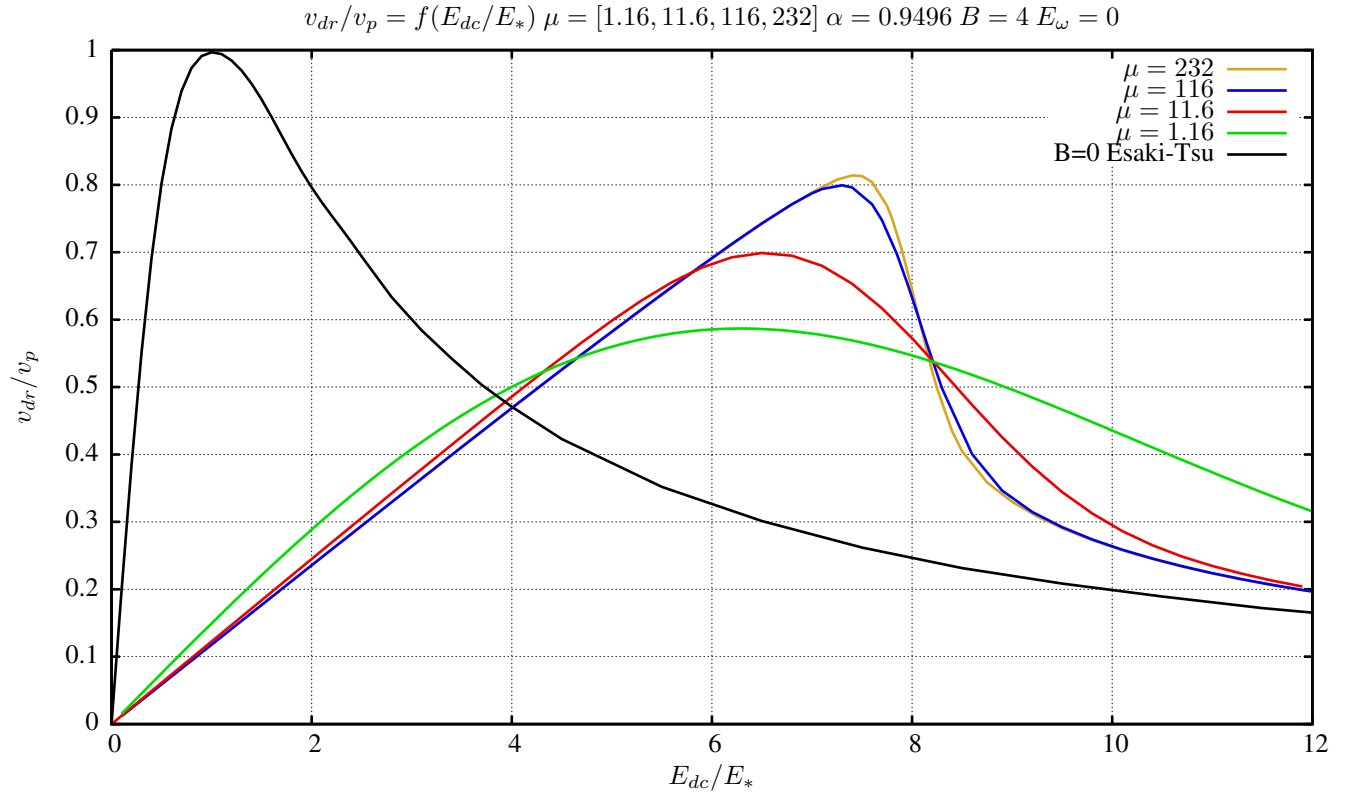


Figure 2

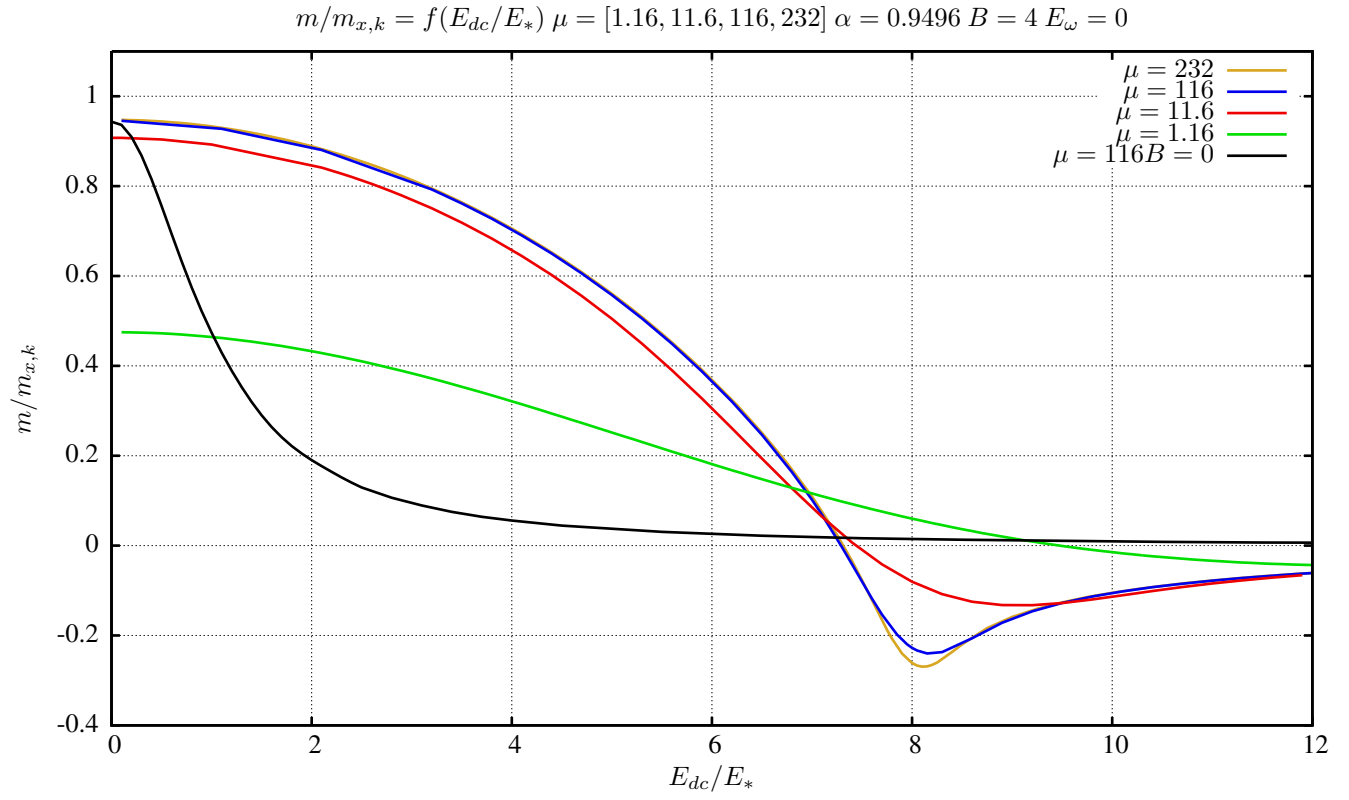


Figure 3

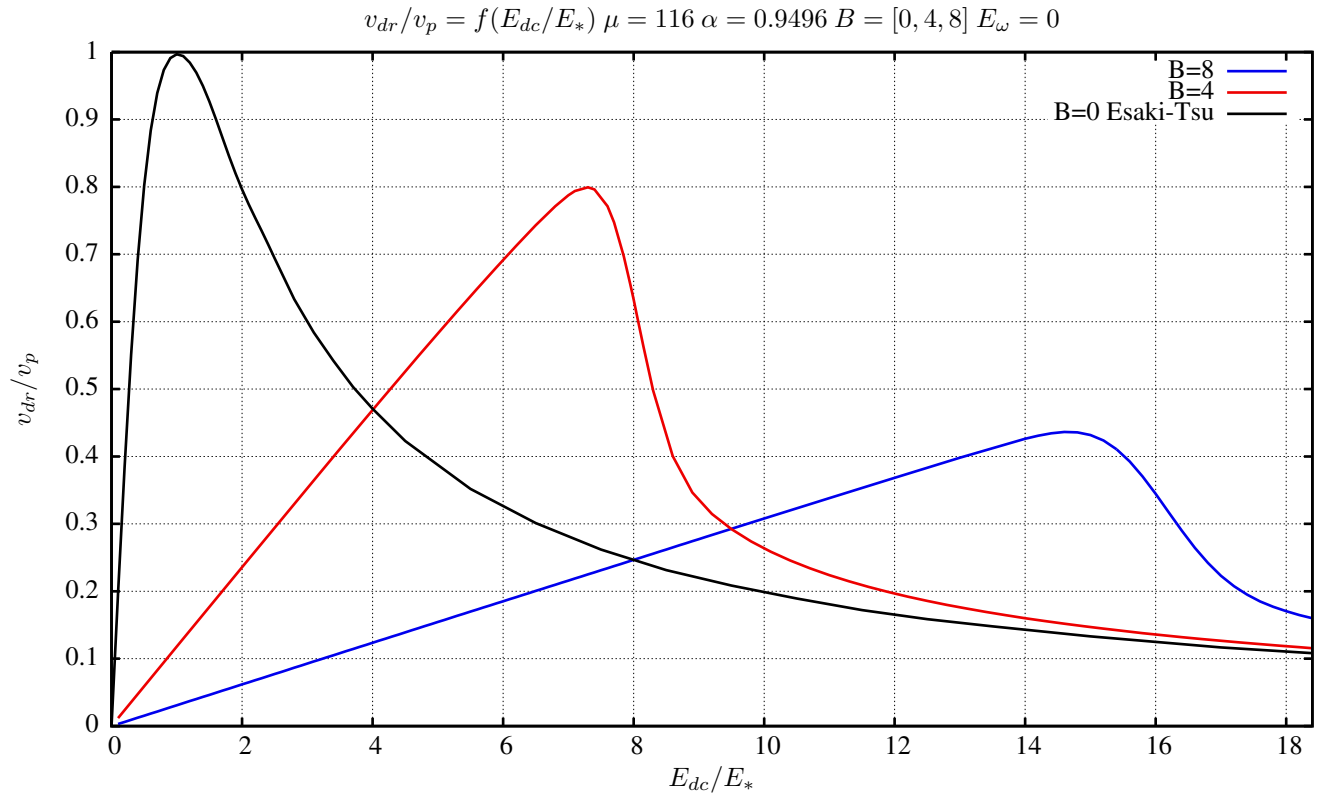


Figure 4

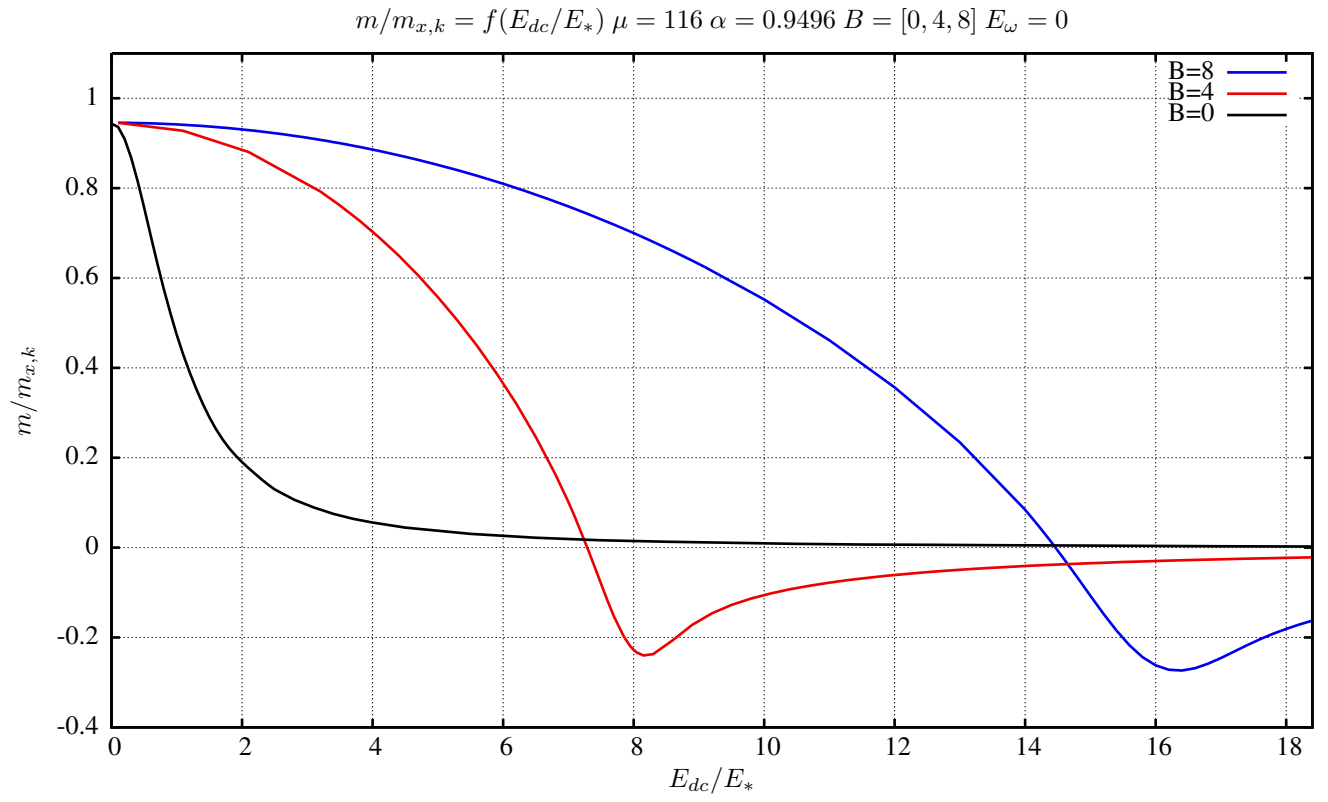


Figure 5

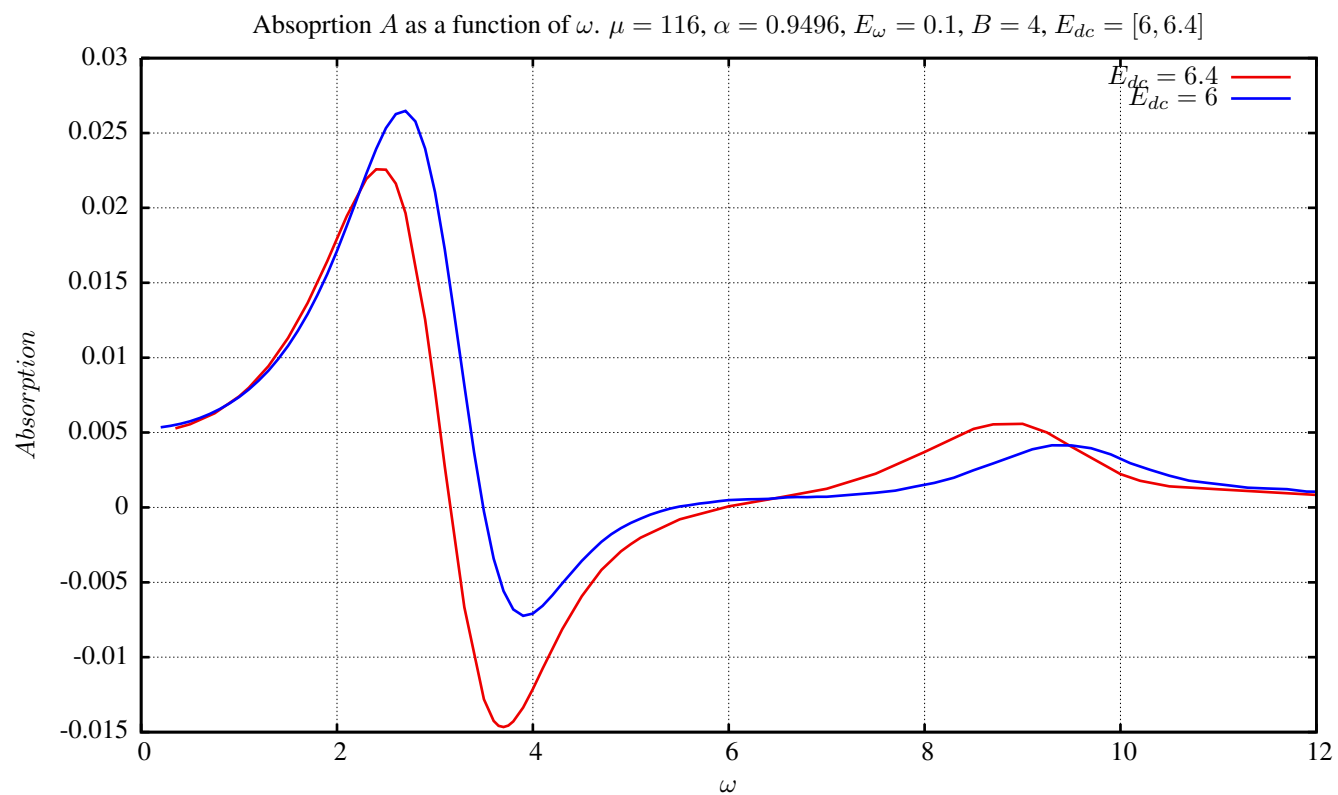


Figure 6

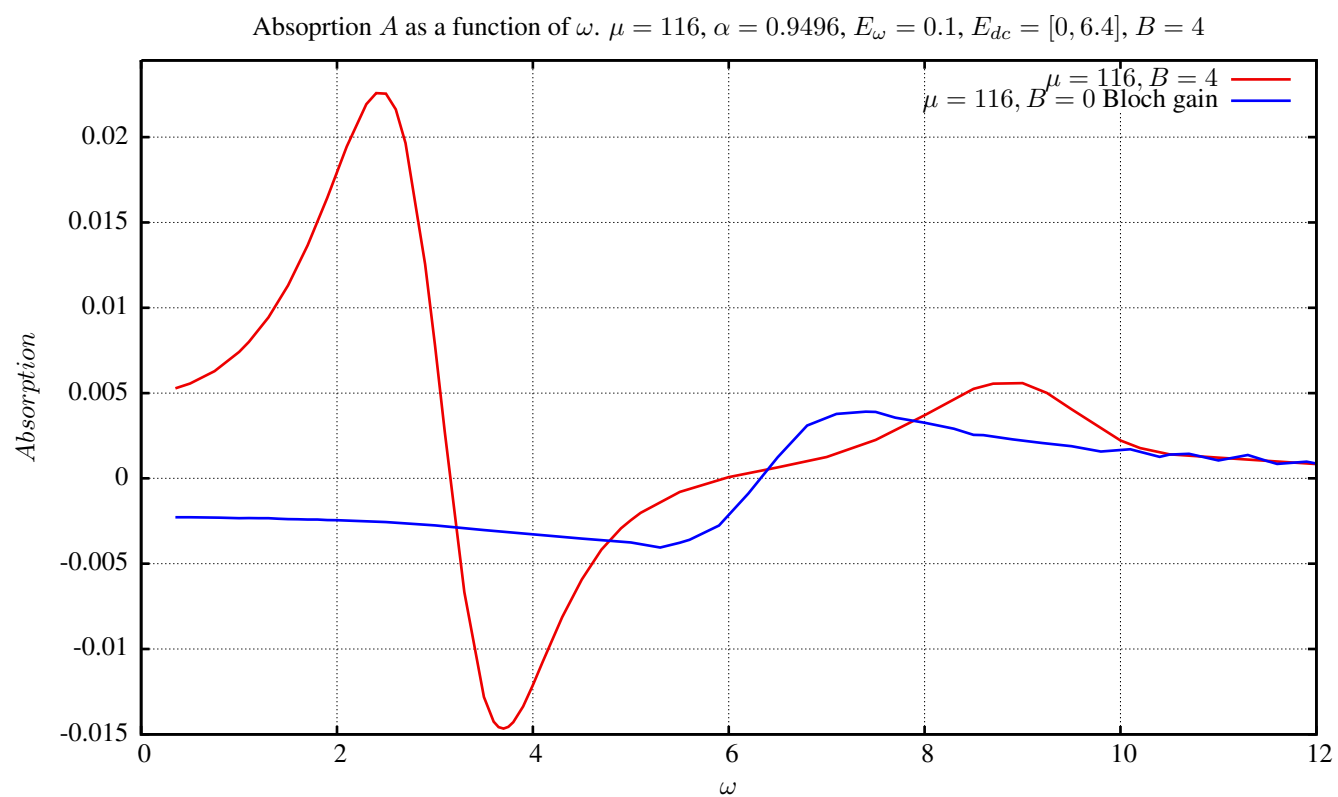


Figure 7

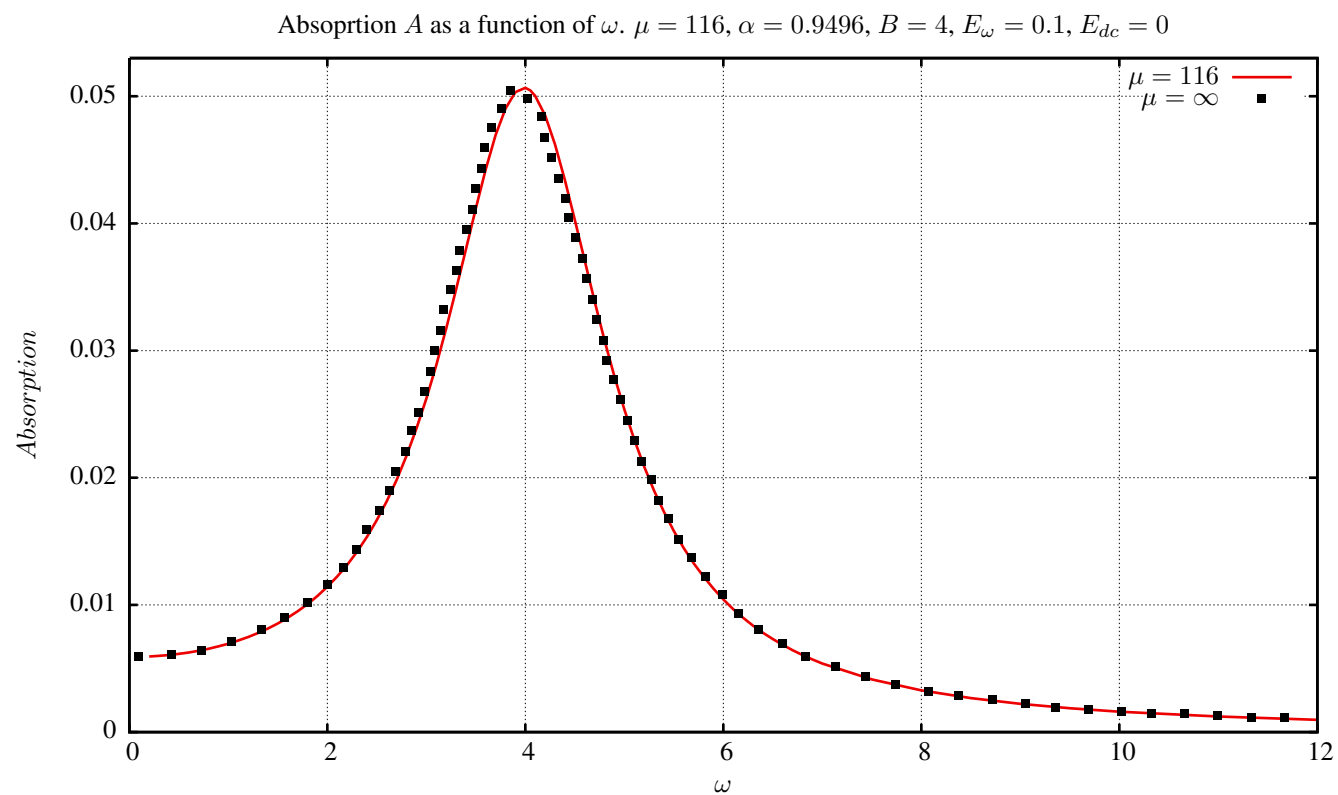


Figure 8

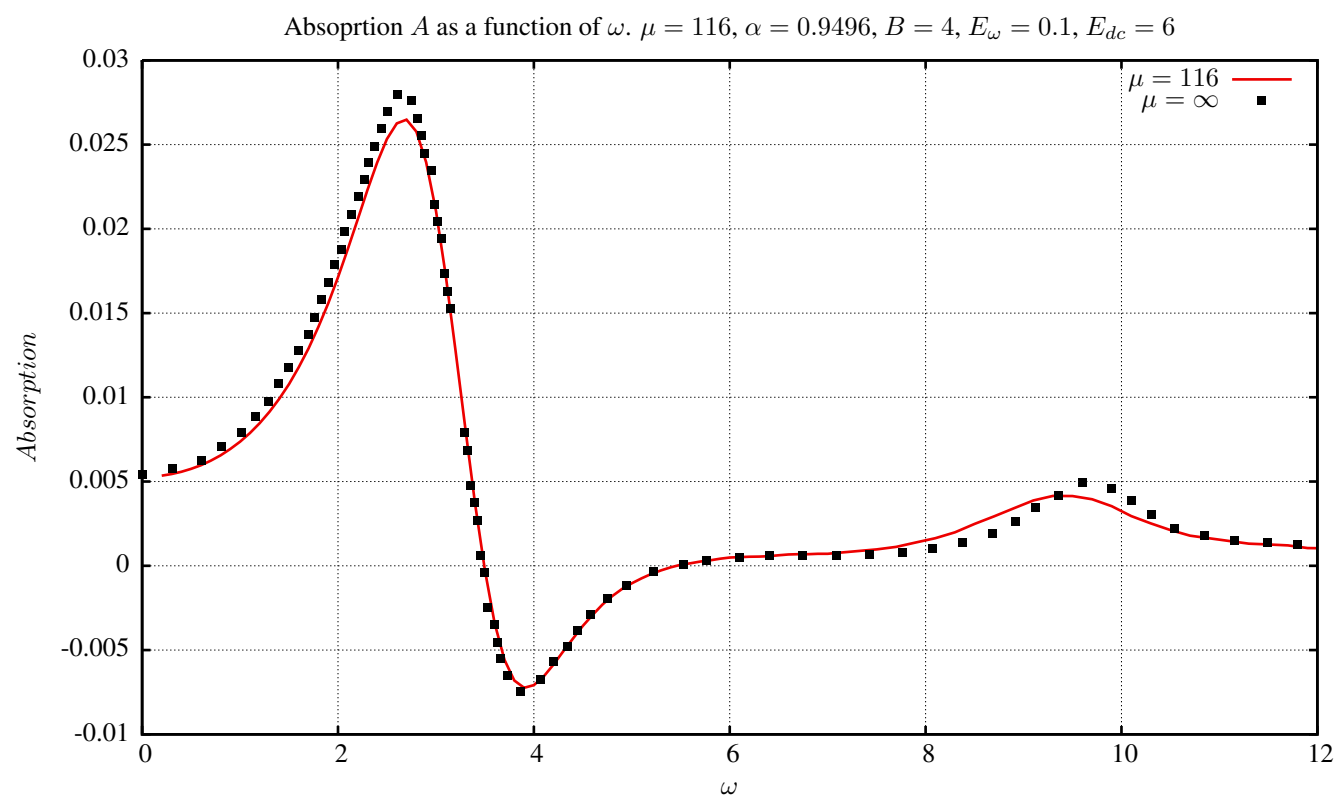


Figure 9

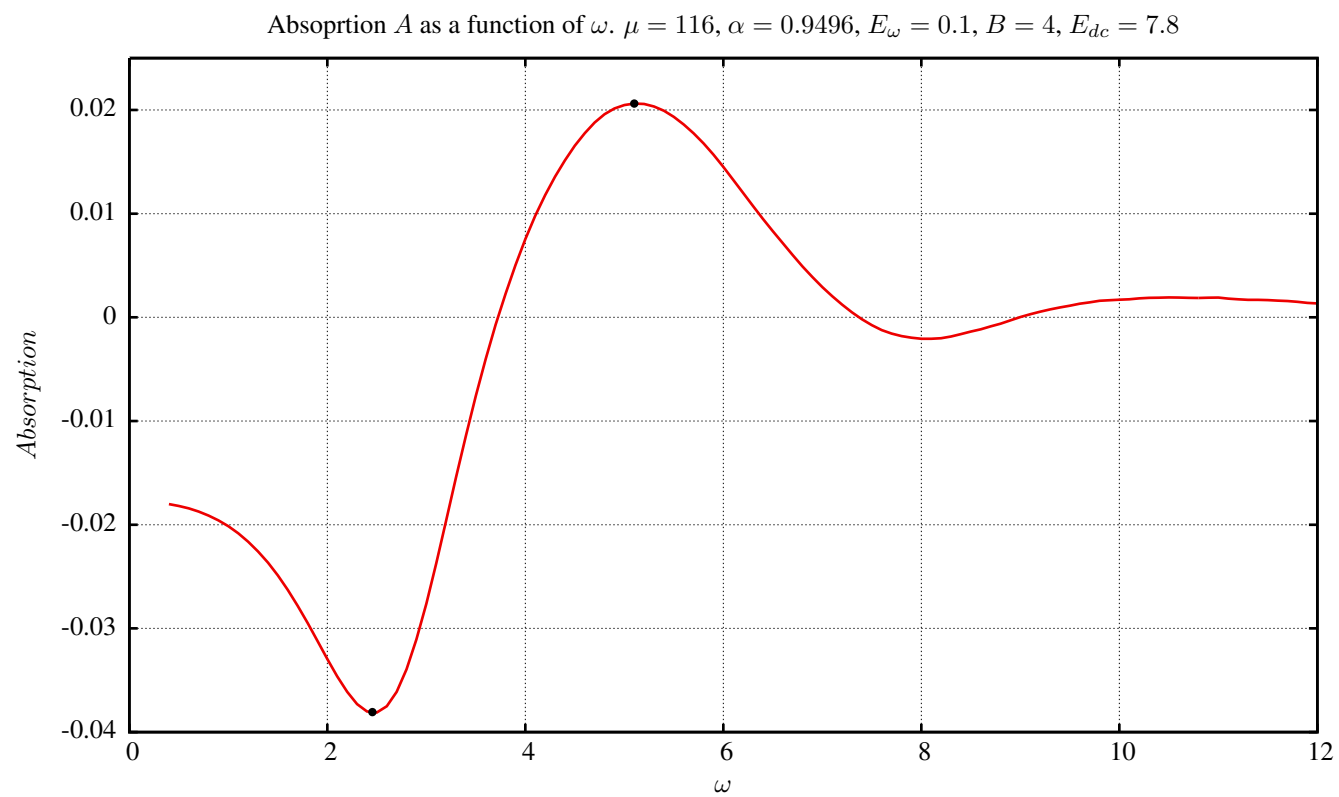


Figure 10

$$E_{dc} = 7.8 \quad B = 4 \quad E_{\omega} = 0.1 \quad \omega = 2.455 \quad \mu = 116 \quad \alpha = 0.9496$$

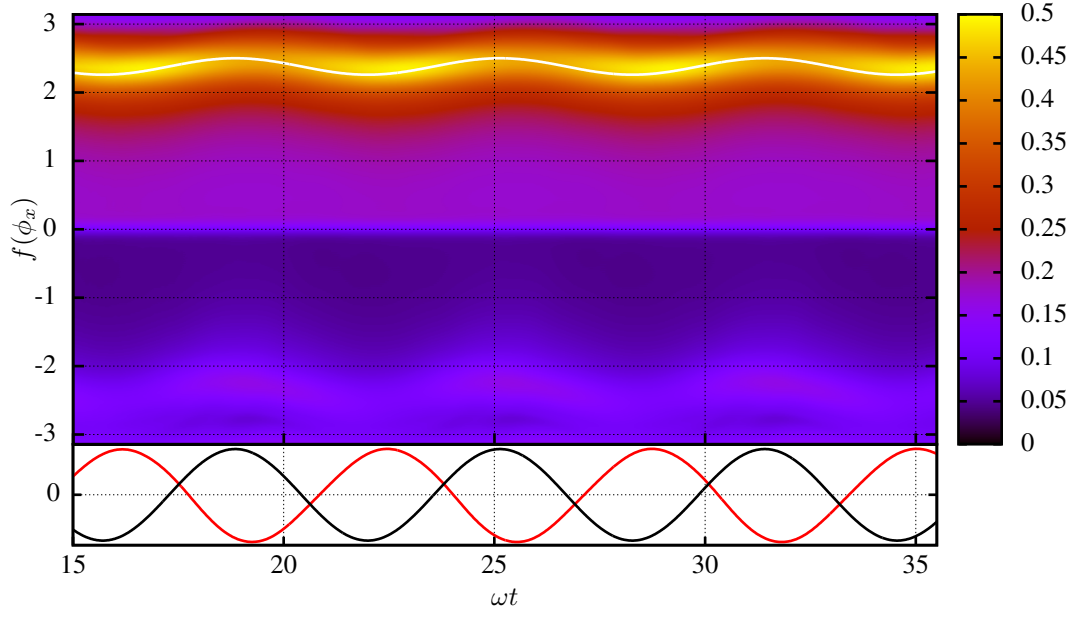


Figure 11

$$E_{dc} = 7.8 \quad B = 4 \quad E_{\omega} = 0.1 \quad \omega = 2.455 \quad \mu = 116 \quad \alpha = 0.9496$$

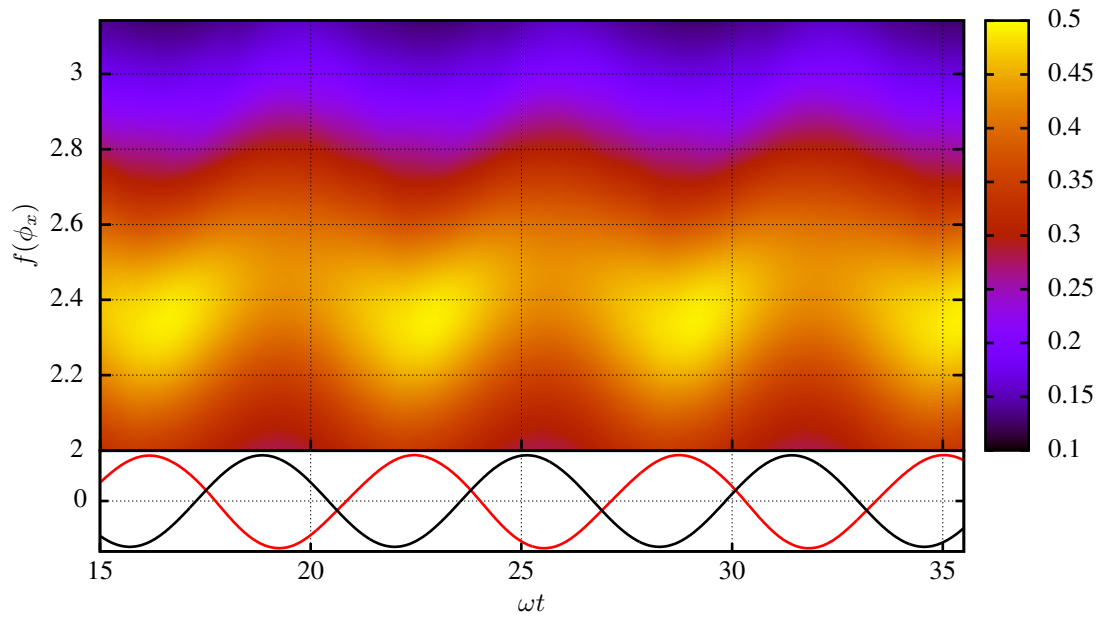


Figure 12

$$E_{dc} = 7.8 \quad B = 4 \quad E_{\omega} = 0.1 \quad \omega = 5.1 \quad \mu = 116 \quad \alpha = 0.9496$$

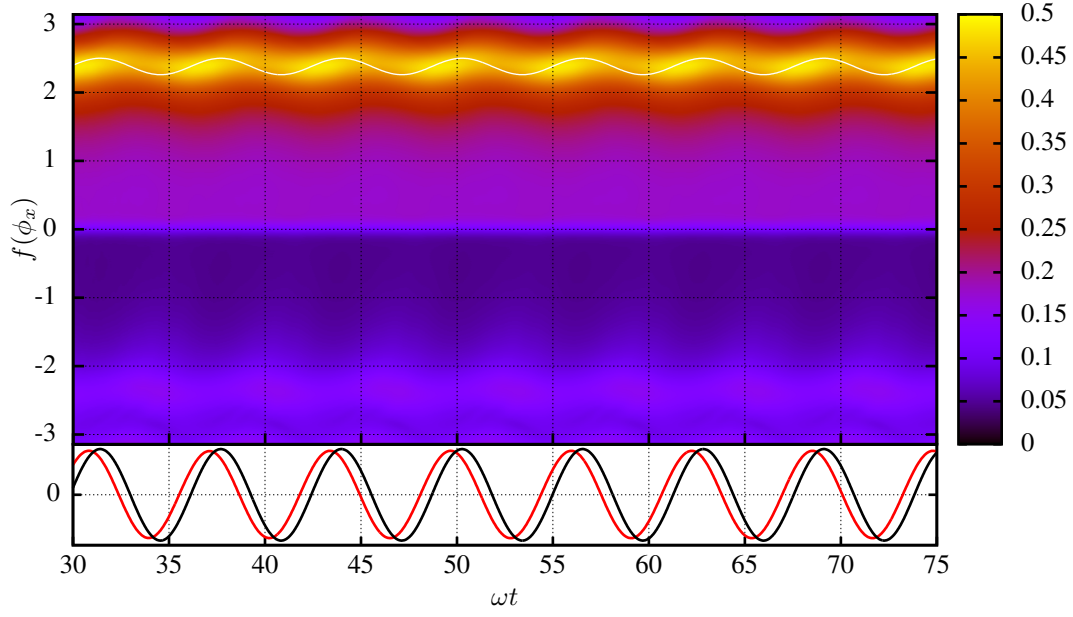


Figure 13

$$E_{dc} = 7.8 \quad B = 4 \quad E_{\omega} = 0.1 \quad \omega = 5.1 \quad \mu = 116 \quad \alpha = 0.9496$$

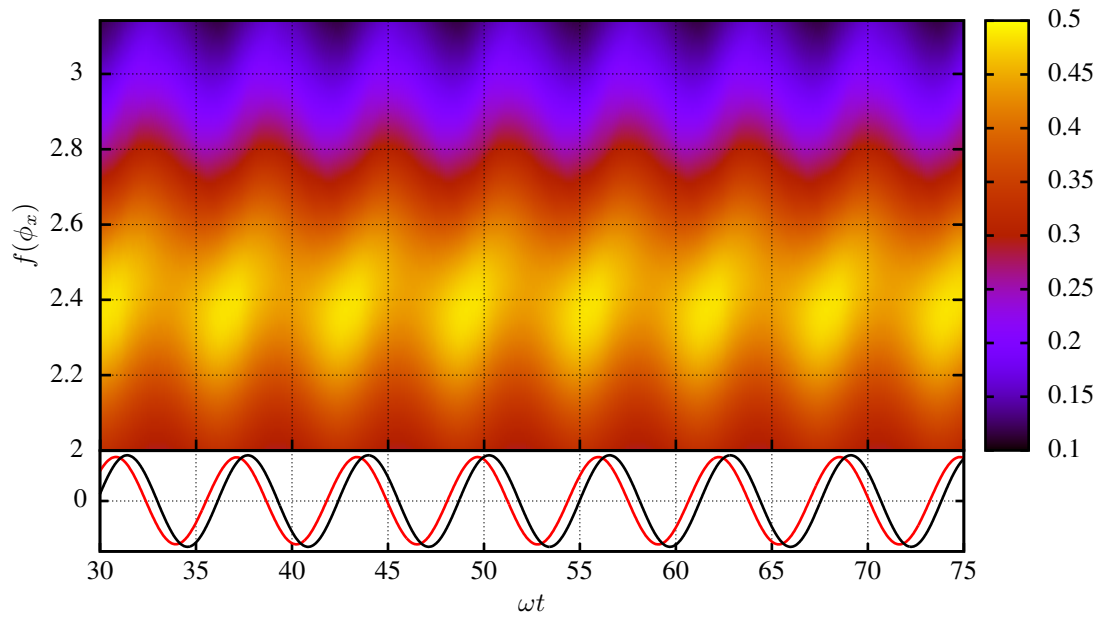


Figure 14

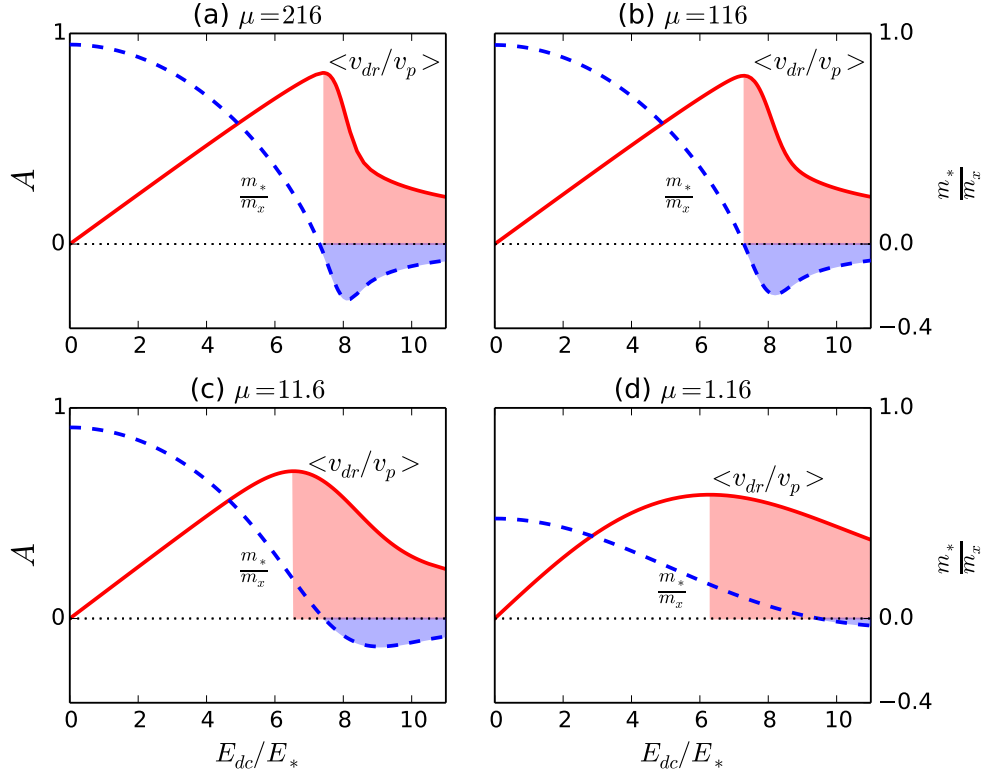


Figure 15:  $E_\omega = 0.1$ ,  $\omega = 2.455$ ,  $B = 4$ ,  $\alpha = 0.9496$ ,  $\mu = [1.16, 11.6, 116, 216]$

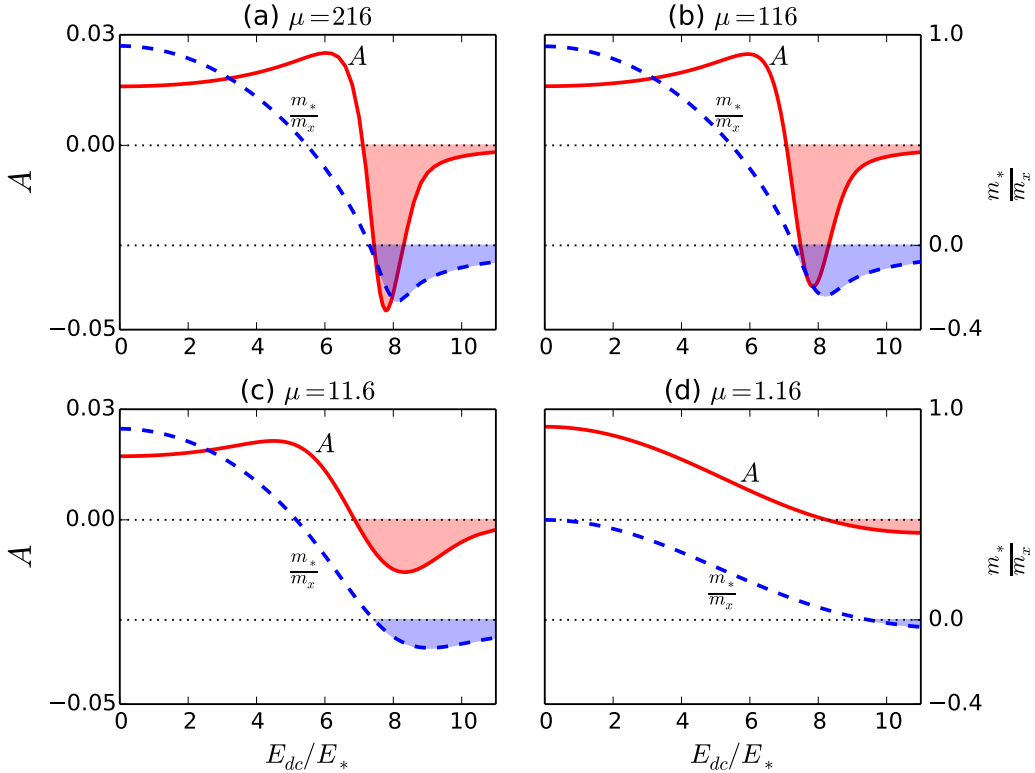


Figure 16:  $E_\omega = 0.1$ ,  $\omega = 2.455$ ,  $B = 4$ ,  $\alpha = 0.9496$ ,  $\mu = [1.16, 11.6, 116, 216]$



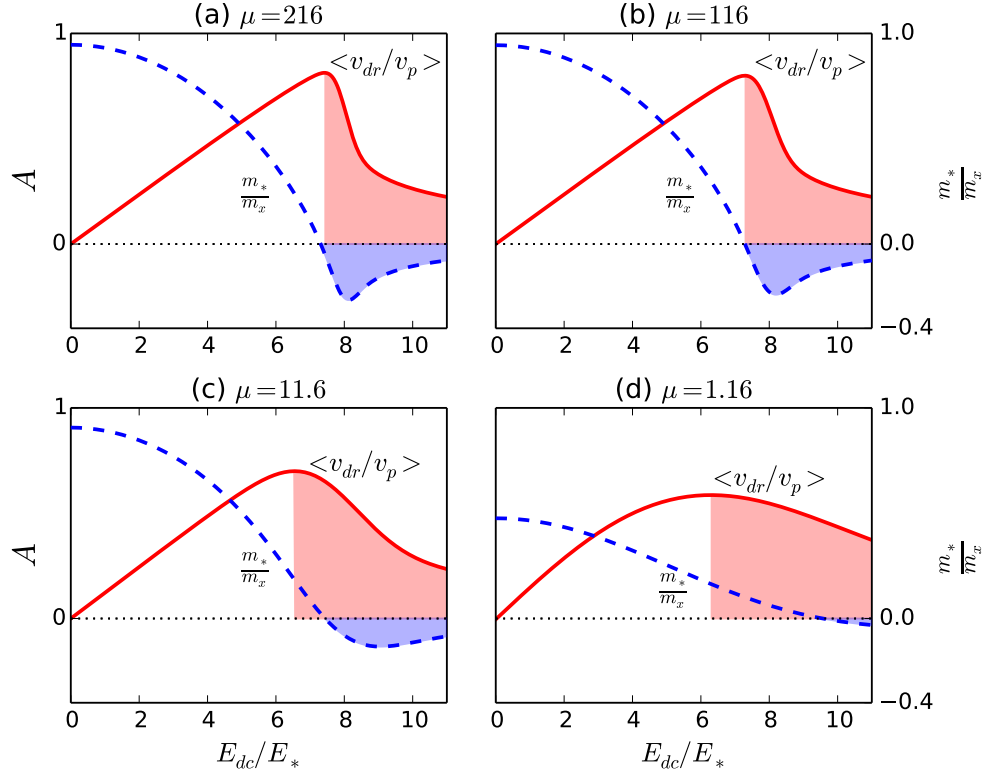


Figure 17:  $E_\omega = 0.1$ ,  $\omega = 4$ ,  $B = 4$ ,  $\alpha = 0.9496$ ,  $\mu = [1.16, 11.6, 116, 216]$

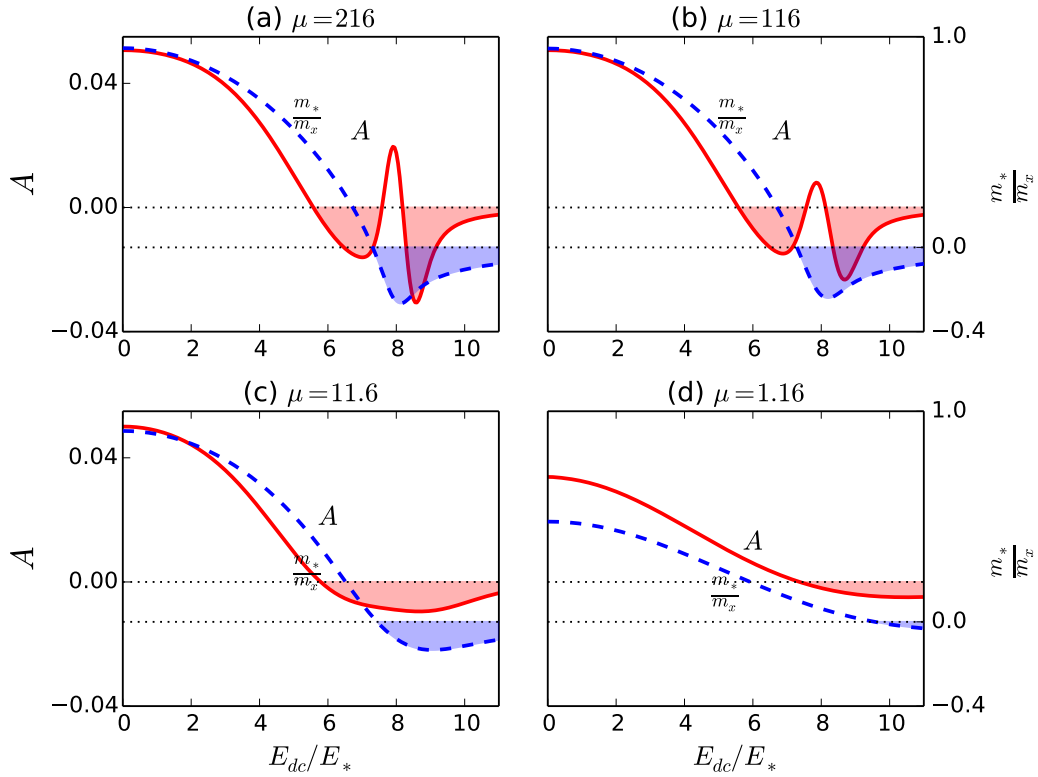


Figure 18:  $E_\omega = 0.1$ ,  $\omega = 4$ ,  $B = 4$ ,  $\alpha = 0.9496$ ,  $\mu = [1.16, 11.6, 116, 216]$

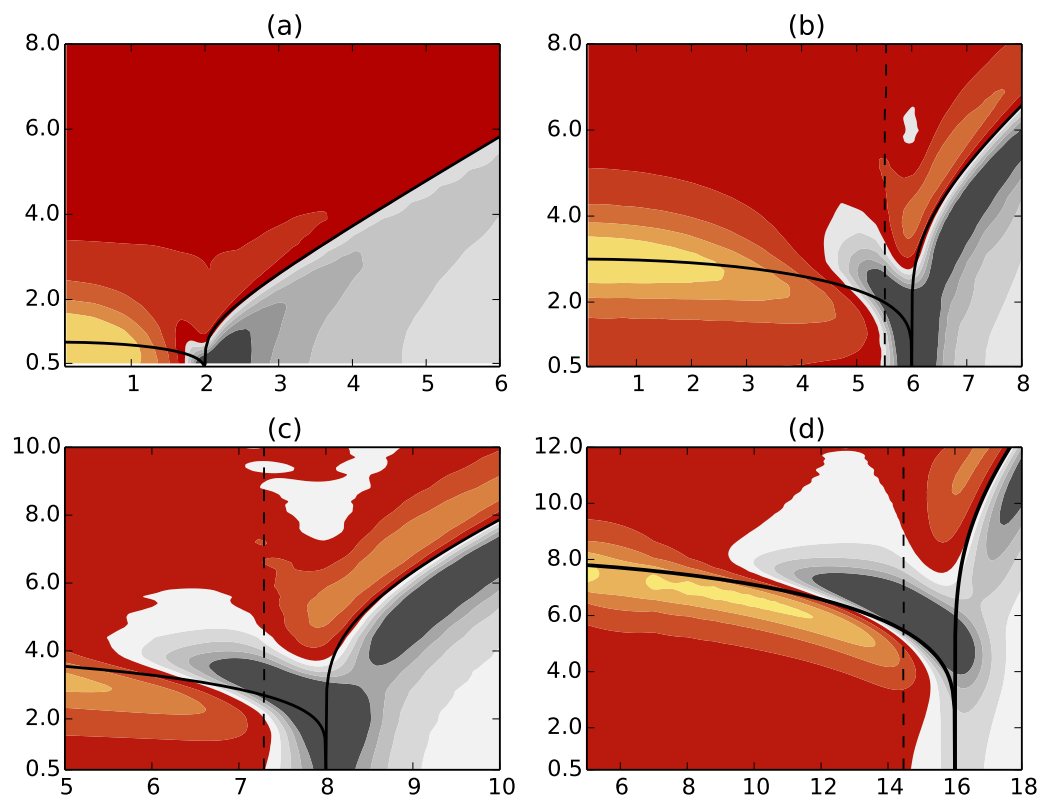


Figure 19: Maps of absorption for  $E_\omega = 0.1$ ,  $\alpha = 0.9496$ ,  $\mu = 116$ ,  $\tilde{B} = [1, 3, 4, 8]$  (plots (a), (b), (c) and (c)) as a function of  $\tilde{E}$  (on the x-axis) and  $\omega$  (on the y-axis)