

Unit 5

Electromagnetism

GRADIENT

It is defined for a scalar field/function $f(x, y, z)$ and is given by

$$\vec{\nabla} f(x, y, z) = \frac{\partial f}{\partial x} \hat{i} + \frac{\partial f}{\partial y} \hat{j} + \frac{\partial f}{\partial z} \hat{k}$$

Physical Significance: The Gradient of a scalar field $\vec{\nabla} f(x, y, z)$ represents both the magnitude and direction of the maximum increase of the scalar field in space.

E.g. Grad of T - Direction of heat flow

Grad of V - Direction of electric field

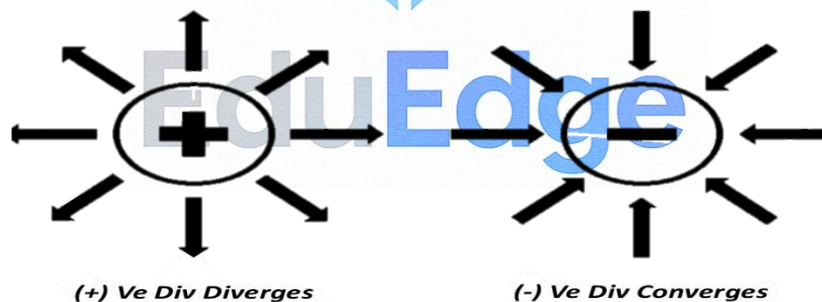
DIVERGENCE

It is defined for a vector field \vec{A} & given by

$$\vec{\nabla} \cdot \vec{A} = \frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z}$$

Physical Significance: Divergence of \vec{A} is the net outward flux per unit volume over a closed surface.

The divergence of a vector field \vec{A} at a given point is a measure of how much the vector field \vec{A} spread out i.e. Diverges.



If \vec{V} velocity of fluid flow

$\vec{\nabla} \cdot \vec{V}$ = Net rate of flow of fluid out of a small region.

CURL

Curl of a vector field is defined as

$$\vec{\nabla} \times \vec{A} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ A_x & A_y & A_z \end{vmatrix} = \hat{i} \left(\frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z} \right) - \hat{j} \left(\frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x} \right) + \hat{k} \left(\frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right)$$

Physical Significance: The Curl of a vector field \vec{A} is the measure of the rotation or circulation of the field at a given point.

E.g. Fluid flow of water

Properties VECTOR ALGEBRA

1. Gauss Divergence Theorem $\oint_s \vec{A} \cdot d\vec{S} = \int_V (\vec{\nabla} \cdot \vec{A}) dV$
2. Stokes Theorem $\oint_l \vec{A} \cdot d\vec{l} = \int_s (\vec{\nabla} \times \vec{A}) \cdot d\vec{S}$
3. Vector Identity $\vec{\nabla} \times (\vec{\nabla} \times \vec{A}) = \vec{\nabla}(\vec{\nabla} \cdot \vec{A}) - \vec{\nabla}^2 \vec{A}$

DIVERGENCE AND CURL OF ELECTRIC FIELD

Divergence

According to this law, the total electric flux passing through a closed surface is equal to the $\frac{1}{\epsilon_0}$ times of the total charge in closed in the surface.

$$\oint_s \vec{E} \cdot d\vec{S} = \frac{Q}{\epsilon_0}$$

$$\oint_s \vec{E} \cdot d\vec{S} = \frac{1}{\epsilon_0} \int_V \rho dV \quad \text{..... (1)} \quad [\text{Integral Form}]$$

Using gauss divergence theorem,

$$\oint_s \vec{E} \cdot d\vec{S} = \int_V (\vec{\nabla} \cdot \vec{E}) dV \quad \text{..... (2)}$$

By eq 1 and 2,

$$\int_V (\vec{\nabla} \cdot \vec{E}) dV = \frac{1}{\epsilon_0} \int_V \rho dV$$

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0} \quad [\text{Differential Form}]$$

Curl

We know that, point charge in static Electric Field

$$\oint_l \vec{E} \cdot d\vec{l} = 0 \quad \text{..... (1)} \quad [\text{Integral Form}]$$

Using stokes theorem,

$$\oint_l \vec{E} \cdot d\vec{l} = \int_s (\vec{\nabla} \times \vec{E}) \cdot d\vec{S} \quad \text{..... (2)}$$

By eq 1 and 2,

$$\vec{\nabla} \times \vec{E} = 0 \quad [\text{Differential Form}]$$

DIVERGENCE AND CURL OF MAGNETIC FIELD

Divergence

According to this law, the total magnetic field through closed surface is zero.

$$\oint_s \vec{B} \cdot d\vec{S} = 0 \quad \text{..... (1)} \quad [\text{Integral Form}]$$

Using gauss divergence theorem,

$$\oint_s \vec{B} \cdot d\vec{S} = \int_V (\vec{\nabla} \cdot \vec{B}) dV \quad \text{..... (2)}$$

By eq 1 and 2,

$$\int_V (\vec{\nabla} \cdot \vec{B}) dV = 0$$

$$\vec{\nabla} \cdot \vec{B} = 0 \quad [\text{Differential Form}]$$

Curl

According to this law, the line integral of the magnetic field induction over a closed path is equal to μ times the total current passing through the loop.

$$\oint_l \vec{B} \cdot d\vec{l} = \mu_0 I_{(total)} \quad \dots\dots\dots (1)$$

If J is current density, then

$$I_{(total)} = \oint_s \vec{J} \cdot d\vec{S} \quad \dots\dots\dots (2)$$

From eq 1 and 2,

$$\boxed{\oint_l \vec{B} \cdot d\vec{l} = \mu_0 \oint_s \vec{J} \cdot d\vec{S}} \quad \dots\dots\dots (3) \quad [\text{Integral Form}]$$

Using stokes theorem,

$$\oint_l \vec{B} \cdot d\vec{l} = \oint_s (\vec{\nabla} \times \vec{B}) \cdot d\vec{S} \quad \dots\dots\dots (4)$$

By eq 3 and 4,

$$\oint_s (\vec{\nabla} \times \vec{B}) \cdot d\vec{S} = \mu_0 \oint_s \vec{J} \cdot d\vec{S}$$
$$\boxed{\vec{\nabla} \times \vec{B} = \mu_0 \vec{J}} \quad [\text{Differential Form}]$$

FARADAY LAW

First Law

According to this law, whenever a magnetic flux linked with a closed loop changes and emf is induced in the loop.

Second Law

According to this law, the magnitude of induced emf in a closed loop is directly proportional to the rate of change in magnetic flux linked with the loop.

If the loop has N turns, then

$$\epsilon = - \frac{d\phi}{dt} \quad \text{where } [\because \phi = \vec{B} \cdot \vec{A}]$$
$$\boxed{\epsilon = -N \frac{d\phi}{dt}}$$

EQUATION OF CONTINUITY

An equation that expresses the equality of incoming and outgoing charges V is called continuity equation.

Physical Significance: It implies that in absence of source of charge the rate of increase of charge in a conductor is equals to a rate of flow of charge into the conductor. continuity equation represents the conservation of charge.

We know that,

$$I = - \frac{\partial Q}{\partial t} = - \frac{\partial}{\partial t} \oint \rho dV \quad \dots\dots\dots (1) \quad [\because Q = \oint \rho dV]$$

If J is current density, then

$$I = \oint \vec{J} \cdot d\vec{S} \quad \dots\dots\dots (2) \quad \left[\because J = \frac{I}{A} \right]$$

From eq 1 and 2,

$$\oint \vec{J} \cdot d\vec{S} = - \frac{\partial}{\partial t} \oint \rho dV \quad \dots\dots\dots (3)$$

Using gauss divergence theorem,

$$\oint \vec{J} \cdot d\vec{S} = \int_V (\vec{\nabla} \cdot \vec{J}) dV \quad \dots\dots\dots (4)$$

By eq 3 and 4,

$$\int_V (\vec{\nabla} \cdot \vec{J}) dV = -\frac{\partial}{\partial t} \oint \rho dV$$

$$\vec{\nabla} \cdot \vec{J} = -\frac{\partial \rho}{\partial t}$$

$$\boxed{\vec{\nabla} \cdot \vec{J} + \frac{\partial \rho}{\partial t} = 0}$$

As charge inside a conductor is 0, then $\rho = 0$ and $\vec{\nabla} \cdot \vec{J} = 0$.

DISPLACEMENT CURRENT

Displacement current (i_d) is the rate of change of an electric displacement field over a time.

According to equation of continuity,

$$\vec{\nabla} \cdot \vec{J} = -\frac{\partial \rho}{\partial t} \quad \dots\dots\dots (1)$$

If charge inside a conductor is 0, then

$$\vec{\nabla} \cdot \vec{J} = 0$$

Let us consider the total current density is $(\vec{J} + \vec{J}_D)$,

$$\vec{\nabla} \cdot (\vec{J} + \vec{J}_D) = 0$$

$$\vec{\nabla} \cdot \vec{J} + \vec{\nabla} \cdot \vec{J}_D = 0$$

$$\vec{\nabla} \cdot \vec{J}_D = -\vec{\nabla} \cdot \vec{J}$$

From eq 1,

$$\vec{\nabla} \cdot \vec{J}_D = -\left(-\frac{\partial \rho}{\partial t}\right) = \frac{\partial \rho}{\partial t}$$

According to Maxwell's 1st eq,

$$\vec{\nabla} \cdot \vec{J}_D = \frac{\partial (\vec{\nabla} \cdot \vec{D})}{\partial t}$$

$$[\because \rho = \vec{\nabla} \cdot \vec{D}]$$

$$\vec{\nabla} \cdot \vec{J}_D = \vec{\nabla} \cdot \frac{\partial \vec{D}}{\partial t}$$

[Where D = Electric Flux Density]

$$\boxed{\vec{J}_D = \frac{\partial \vec{D}}{\partial t}}$$

[Where J_D = Displacement Current Density]

According to Ampere's Law,

$$\vec{\nabla} \times \vec{B} = \mu_0 (\vec{J} + \vec{J}_D)$$

$$\vec{\nabla} \times \vec{B} = \mu_0 \left(\vec{J} + \frac{\partial \vec{D}}{\partial t} \right)$$

[Modify Form of Ampere's Law]

$$\boxed{\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}}$$

$$[\because i_d = \epsilon_0 \frac{\partial \phi_E}{\partial t}]$$

MAXWELL EQUATIONS

Maxwell 1 Equation

According to this law, the total electric flux passing through a closed surface is equal to the $\frac{1}{\epsilon_0}$ times of the total charge in closed in the surface.

$$\oint_s \vec{E} \cdot d\vec{S} = \frac{Q}{\epsilon_0}$$

$$\boxed{\oint_s \vec{E} \cdot d\vec{S} = \frac{1}{\epsilon_0} \int_V \rho dV} \quad \text{..... (1)} \quad [\text{Integral Form}]$$

Using gauss divergence theorem,

$$\oint_s \vec{E} \cdot d\vec{S} = \int_V (\vec{\nabla} \cdot \vec{E}) dV \quad \text{..... (2)}$$

By eq 1 and 2,

$$\int_V (\vec{\nabla} \cdot \vec{E}) dV = \frac{1}{\epsilon_0} \int_V \rho dV$$

$$\boxed{\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0}} \quad [\text{Differential Form}]$$

$$\vec{\nabla} \cdot (\vec{E} \epsilon_0) = \rho \quad [\text{Displacement current}(D) = \vec{E} \cdot \epsilon_0]$$

$$\boxed{\vec{\nabla} \cdot \vec{D} = \rho}$$

Maxwell 2 Equation

According to this law, the total magnetic field through closed surface is zero.

$$\boxed{\oint_s \vec{B} \cdot d\vec{S} = 0} \quad \text{..... (1)} \quad [\text{Integral Form}]$$

Using gauss divergence theorem,

$$\oint_s \vec{B} \cdot d\vec{S} = \int_V (\vec{\nabla} \cdot \vec{B}) dV \quad \text{..... (2)}$$

By eq 1 and 2,

$$\int_V (\vec{\nabla} \cdot \vec{B}) dV = 0$$

$$\boxed{\vec{\nabla} \cdot \vec{B} = 0} \quad [\text{Differential Form}]$$

Maxwell 3 Equation

According to faraday law of electromagnetic induction, emf induced in closed loop is equal to the rate of change of magnetic flux enclosed in loop.

$$\epsilon = -\frac{d\phi_B}{dt} \quad \text{..... (1)}$$

Since, the line integral of electric field also represent the emf.

$$\epsilon = \oint_l \vec{E} \cdot d\vec{l} \quad \text{..... (2)}$$

From eq 1 and 2,

$$\oint_l \vec{E} \cdot d\vec{l} = -\frac{d\phi_B}{dt} \quad \text{..... (3)}$$

As we know that,

$$\phi_B = \oint_s \vec{B} \cdot d\vec{S}$$

Put the value in eq 3,

$$\oint_l \vec{E} \cdot d\vec{l} = -\frac{d[\oint_s \vec{B} \cdot d\vec{S}]}{dt}$$

$$\boxed{\oint_l \vec{E} \cdot d\vec{l} = -\oint_s \frac{d\vec{B}}{dt} \cdot d\vec{S}} \quad \text{..... (4)} \quad [\text{Integral Form}]$$

Using stokes theorem,

$$\oint_l \vec{E} \cdot d\vec{l} = \oint_s (\vec{\nabla} \times \vec{E}) d\vec{S} \quad \dots\dots\dots (5)$$

From eq 4 and 5,

$$\oint_s (\vec{\nabla} \times \vec{E}) d\vec{S} = - \oint_s \frac{d\vec{B}}{dt} \cdot d\vec{S}$$

$$\boxed{\vec{\nabla} \times \vec{E} = - \frac{d\vec{B}}{dt}} \quad \text{[Differential Form]}$$

Maxwell 4 Equation

According to this law, the line integral of the magnetic field induction over a closed path is equal to μ times the total current passing through the loop.

$$\oint_l \vec{B} \cdot d\vec{l} = \mu_0 I_{(total)} \quad \dots\dots\dots (1)$$

If J is current density, then

$$I_{(total)} = \oint_s \vec{J} \cdot d\vec{S} \quad \dots\dots\dots (2)$$

From eq 1 and 2,

$$\oint_l \vec{B} \cdot d\vec{l} = \mu_0 \oint_s \vec{J} \cdot d\vec{S} \quad \dots\dots\dots (3)$$

According to maxwell total J is due to two types of current are:

- (a) Free current / conduction current density (\vec{J}), due to motion of free charges.
- (b) Displacement current density (\vec{J}_D), due to variation of electric field w.r.to time.

From eq 3,

$$\boxed{\oint_l \vec{B} \cdot d\vec{l} = \mu_0 \oint_s (\vec{J} + \vec{J}_D) \cdot d\vec{S}} \quad \dots\dots\dots (4) \quad \text{[Integral Form]}$$

Using stokes theorem,

$$\oint_l \vec{B} \cdot d\vec{l} = \oint_s (\vec{\nabla} \times \vec{B}) d\vec{S} \quad \dots\dots\dots (5)$$

From eq 4 and 5,

$$\oint_s (\vec{\nabla} \times \vec{B}) d\vec{S} = \mu_0 \oint_s (\vec{J} + \vec{J}_D) \cdot d\vec{S}$$

$$\vec{\nabla} \times \vec{B} = \mu_0 (\vec{J} + \vec{J}_D) \quad \dots\dots\dots (6)$$

Taking divergence both the sides in eq 6,

$$\vec{\nabla} \cdot (\vec{\nabla} \times \vec{B}) = \mu_0 (\vec{\nabla} \cdot \vec{J} + \vec{\nabla} \cdot \vec{J}_D)$$

According to vector identity,

$$\begin{aligned} \vec{\nabla} \cdot (\vec{\nabla} \times \vec{B}) &= 0 \\ \vec{\nabla} \cdot \vec{J} + \vec{\nabla} \cdot \vec{J}_D &= 0 \\ \vec{\nabla} \cdot \vec{J}_D &= -\vec{\nabla} \cdot \vec{J} \quad \dots\dots\dots (7) \end{aligned}$$

According to equation of continuity,

$$\begin{aligned} \vec{\nabla} \cdot \vec{J} + \frac{\partial \rho}{\partial t} &= 0 \\ \vec{\nabla} \cdot \vec{J} &= -\frac{\partial \rho}{\partial t} \end{aligned}$$

From eq 7,

$$\vec{\nabla} \cdot \vec{J}_D = \frac{\partial \rho}{\partial t}$$

$$\vec{\nabla} \cdot \vec{J}_D = \frac{\partial(\vec{\nabla} \cdot \vec{D})}{\partial t}$$

$$\vec{J}_D = \frac{\partial \vec{D}}{\partial t}$$

Put value of \vec{J}_D in eq 6, we get

$$\vec{\nabla} \times \vec{B} = \mu_0 \left(\vec{J} + \frac{\partial \vec{D}}{\partial t} \right)$$

$$\vec{\nabla} \times \mu_0 \vec{H} = \mu_0 \left(\vec{J} + \frac{\partial \vec{D}}{\partial t} \right) \quad [\because \vec{B} = \mu_0 \vec{H}]$$

$\vec{\nabla} \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$	Or	$\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$	[Differential Form]
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ELECTROMAGNETIC WAVE PROPAGATION IN FREE SPACE

Electromagnetic Waves are waves that consist of oscillating electric and magnetic fields, which are propagate through space without requiring a medium. These waves obey maxwell equations and travel at the speed of light in vacuum.

Maxwell's equations describe the behavior of EM Waves.

1. Gauss's Law for Electricity

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

2. Gauss's Law for Magnetism

$$\vec{\nabla} \cdot \vec{B} = 0$$

3. Faraday's Law of Electromagnetic Induction

$$\vec{\nabla} \times \vec{E} = -\frac{d\vec{B}}{dt}$$

4. Ampere's Law

$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

Using these equations, we derive the wave equation for EM Waves in free space.

Maxwell's equations in free space

Since, free space has no charges ($\rho = 0$) and no currents ($\vec{J} = 0$), Maxwell's equations simplify to:

$$1. \quad \vec{\nabla} \cdot \vec{E} = 0$$

$$2. \quad \vec{\nabla} \cdot \vec{B} = 0$$

$$3. \quad \vec{\nabla} \times \vec{E} = -\frac{d\vec{B}}{dt}$$

$$4. \quad \vec{\nabla} \times \vec{B} = \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

Wave Equation for \vec{E}

According to Faraday's Law,

$$\vec{\nabla} \times \vec{E} = -\frac{d\vec{B}}{dt}$$

Now, take the curl on both sides

$$\vec{\nabla} \times (\vec{\nabla} \times \vec{E}) = -\vec{\nabla} \times \frac{d\vec{B}}{dt} \quad \text{..... (1)}$$

Using vector identity,

$$\vec{\nabla} \times (\vec{\nabla} \times \vec{E}) = \vec{\nabla}(\vec{\nabla} \cdot \vec{E}) - \vec{\nabla}^2 \vec{E} \quad \text{..... (2)}$$

From eq 1 and 2,

$$\vec{\nabla}^2 \vec{E} = \vec{\nabla} \times \frac{d\vec{B}}{dt} \quad \text{..... (3)} \quad [\because \vec{\nabla} \cdot \vec{E} = 0]$$

From Ampere's Law,

$$\vec{\nabla} \times \vec{B} = \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

Take the time derivative on both sides

$$\vec{\nabla} \times \frac{\partial \vec{B}}{\partial t} = \mu_0 \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2}$$

Put the value in eq 3, we get

$$\boxed{\vec{\nabla}^2 \vec{E} = \mu_0 \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2}} \quad \text{OR} \quad \boxed{\vec{\nabla}^2 \vec{E} - \mu_0 \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2} = 0}$$

Wave Equation for \vec{B}

According to Ampere's Law,

$$\vec{\nabla} \times \vec{B} = \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

Now, take the curl on both sides

$$\vec{\nabla} \times (\vec{\nabla} \times \vec{B}) = \mu_0 \epsilon_0 \vec{\nabla} \times \frac{\partial \vec{E}}{\partial t} \quad \text{..... (1)}$$

Using vector identity,

$$\vec{\nabla} \times (\vec{\nabla} \times \vec{B}) = \vec{\nabla}(\vec{\nabla} \cdot \vec{B}) - \vec{\nabla}^2 \vec{B} \quad \text{..... (2)}$$

From eq 1 and 2,

$$-\vec{\nabla}^2 \vec{B} = \mu_0 \epsilon_0 \vec{\nabla} \times \frac{\partial \vec{E}}{\partial t} \quad \text{..... (3)} \quad [\because \vec{\nabla} \cdot \vec{B} = 0]$$

From Faraday's Law,

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

Take the time derivative on both sides

$$\vec{\nabla} \times \frac{\partial \vec{E}}{\partial t} = -\frac{\partial^2 \vec{B}}{\partial t^2}$$

Put the value in eq 3, we get

$$\boxed{\vec{\nabla}^2 \vec{B} = \mu_0 \epsilon_0 \frac{\partial^2 \vec{B}}{\partial t^2}} \quad \text{OR} \quad \boxed{\vec{\nabla}^2 \vec{B} - \mu_0 \epsilon_0 \frac{\partial^2 \vec{B}}{\partial t^2} = 0}$$

Compare this equation with Wave Standard Equation

$$\nabla^2 f = \frac{1}{V^2} \frac{\partial^2 f}{\partial t^2}$$

We can identify wave speed as

$$V = \frac{1}{\sqrt{\mu_0 \epsilon_0}}$$

In free space,

$$V = \frac{1}{\sqrt{(4\pi \times 10^{-7})(8.85 \times 10^{-12})}} = 3 \times 10^8 \text{ m/s} = c$$

Speed of light in vacuum.

Characteristics of EM Waves in Free Space

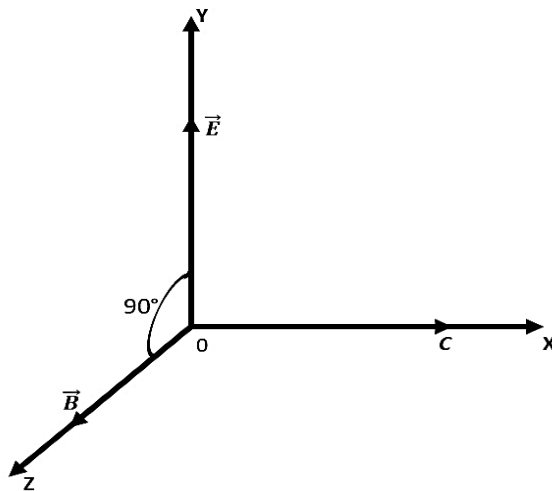
1. Transverse nature

The E and B are perpendicular to each other and to the direction of wave propagation(k).

If the wave propagates in the x-direction, then

$$E = E_0 \hat{y} \cos(kx - \omega t)$$

$$B = B_0 \hat{z} \cos(kx - \omega t)$$



2. Speed of Wave

The speed of the wave in free space is $V = \frac{1}{\sqrt{\mu_0 \epsilon_0}} = 3 \times 10^8 \text{ m/s} = c$.

3. Relation Between E and B

The magnitudes E and B are related as $c = \frac{E_0}{B_0}$.

ENERGY FLOW AND POYNTING VECTOR

The energy carried by EM Waves is characterized by the Poynting vector, which represents the power per unit area transported by the wave.

$$S = E \times B$$

Where,

S is Poynting vector (W/m^2)

E is Electric Field (V/m)

B is Magnetic Field (T)