On Polarization Properties of Light, Gaussian Beams and Spin-Orbit Interaction of Light

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Submitted to

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Topics of discussion

Polarization properties of light

Gaussian Beam

Spin-orbit interaction of light

Polarization properties of light

- ▶ Jones Formalism
- ► Stokes-Muller formalism

Jones Vector

Electric field of $fully\ polarized\ {\rm EM}$ wave propagating along z-axis is given by

$$\boldsymbol{E}(\boldsymbol{r},t) = \begin{bmatrix} A_x(\boldsymbol{r})e^{i\delta_x} \\ A_y(\boldsymbol{r})e^{i\delta_y} \\ 0 \end{bmatrix} e^{-i(kz-\omega t)}$$

Define normalized **Jones vector** s.t. $J^* J = 1$ as

$$\boldsymbol{J}(\boldsymbol{r},t) = \frac{1}{\sqrt{A_x^2 + A_y^2}} \begin{bmatrix} A_x(\boldsymbol{r}) e^{i\delta_x} \\ A_y(\boldsymbol{r}) e^{i\delta_y} \end{bmatrix}$$

Note that intensity, $I = A_x^2 + A_y^2 = J^*J$

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Note that intensity, $I = A_x^2 + A_y^2 = J^*J$

Jones vector of usual polarization state

Polarization state	\boldsymbol{J}
H angle	$\begin{bmatrix} 1 \\ 0 \end{bmatrix}$
$ V\rangle$	0
$ P\rangle$	$\frac{1}{\sqrt{2}}\begin{bmatrix}1\\1\end{bmatrix}$
M angle	$\frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ -1 \end{bmatrix}$
$ L\rangle$	$\frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ i \end{bmatrix}$
$ R\rangle$	$\frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ -i \end{bmatrix}$

Jones Matrix & evolution of Jones vector

Jones matrix for an optical element be M s.t.

$$\boldsymbol{M} = \begin{bmatrix} m_{11} & m_{12} \\ m_{21} & m_{22} \end{bmatrix}$$

If a polarized light of Jones vector J_{in} passes through that optical element then the Jones vector of output light is given by

$$oldsymbol{J}_{out} = oldsymbol{M} oldsymbol{J}_{in}$$

- ightharpoonup Composition rule: $M = M_1 M_2 \dots M_n$
- Frame rotation by θ : $M_{\theta} = R(-\theta) M R(\theta)$ where $R(\theta)$ is passive rotation matrix.

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Jones matrix of usual optical element

Optical element

M

1	
Free space	$\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$
x-Polariser	$\begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$
Right circular polariser	$\frac{1}{2} \begin{bmatrix} 1 & i \\ -i & 1 \end{bmatrix}$
Linear di-attenuator	$\begin{bmatrix} a & 0 \\ 0 & b \end{bmatrix}$
Half-wave plate	$\begin{bmatrix} 1 & 0 \end{bmatrix}$
with fast axis horizontal	0 - 1
Quarter-wave plate	$\begin{bmatrix} 1 & 0 \end{bmatrix}$
with fast axis horizontal	$\begin{bmatrix} 0 & i \end{bmatrix}$
General phase retarder	$ \begin{bmatrix} e^{i\phi_x} & 0 \\ 0 & e^{i\phi_y} \end{bmatrix} $

Polarization properties of light

- ▶ Jones Formalism
- ➤ Stokes-Muller formalism

Coherency matrix

Coherency matrix, C defined as

$$oldsymbol{C} = \left\langle oldsymbol{E} \otimes oldsymbol{E}^\dagger
ight
angle = \left\langle oldsymbol{E} oldsymbol{E}^\dagger
ight
angle = \left[egin{array}{c} \langle E_x E_x^*
angle & \langle E_x E_y^*
angle \\ \langle E_y E_x^*
angle & \langle E_y E_y^*
angle \end{array}
ight]$$

only for fully polarized light

- $ightharpoonup C = C^{\dagger}$ (Hermitian).
- ightharpoonup Time averaged intensity, I = Tr(C)

Let basis set.

$$\left\{ \underbrace{\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}}_{\boldsymbol{V_0}}, \underbrace{\begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}}_{\boldsymbol{V_1}}, \underbrace{\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}}_{\boldsymbol{V_2}}, \underbrace{\begin{bmatrix} 0 & i \\ -i & 1 \end{bmatrix}}_{\boldsymbol{V_3}} \right\} s.t. \ \boldsymbol{C} = \frac{1}{2} \sum_{i=0}^{3} S_i \boldsymbol{V_i}$$

Wave Optics (2015), Polarization lecture- YouTube

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Stokes vector

$$m{C} = rac{1}{2} \sum_{i=0}^{3} S_i m{V_i} \longrightarrow egin{bmatrix} S_0 \ S_1 \ S_2 \ S_3 \end{bmatrix} = m{S} \; (Stokes \; vector)$$

Measurement of Stokes parameters by passing the light through

- 1. isotropic medium, $I_0 = S_0$
- 2. x-polariser, $I_1 = \frac{1}{2}(S_0 + S_1)$
- 3. 45° -polariser, $I_2 = \frac{1}{2}(S_0 + S_2)$
- 4. circular polariser, $I_3 = \frac{1}{2}(S_0 + S_3)$

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Stokes vector of usual polarization state

Polarization state	$oldsymbol{C}$	\boldsymbol{S}
$ H\rangle$	$\begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$	$\begin{bmatrix} 1 & 1 & 0 & 0 \end{bmatrix}^T$
$ V\rangle$	$\begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix}$	$\begin{bmatrix} 1 & -1 & 0 & 0 \end{bmatrix}^T$
$ P\rangle$	$\frac{1}{2} \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix}$	$\begin{bmatrix}1 & 0 & 1 & 0\end{bmatrix}^T$
M angle	$\frac{1}{2} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix}$	$\begin{bmatrix} 1 & 0 & -1 & 0 \end{bmatrix}^T$
$ L\rangle$	$\frac{1}{2} \begin{bmatrix} 1 & -i \\ i & 1 \end{bmatrix}$	$\begin{bmatrix}1 & 0 & 0 & 1\end{bmatrix}^T$
$ R\rangle$	$\frac{1}{2} \begin{bmatrix} 1 & i \\ -i & 1 \end{bmatrix}$	$\begin{bmatrix} 1 & 0 & 0 & -1 \end{bmatrix}^T$
Un-polarized	$\begin{array}{c c} 1 & 0 \\ \hline & 0 & 1 \end{array}$	$\boxed{\begin{bmatrix}1 & 0 & 0 & 0\end{bmatrix}^T}$

Muller Matrix & evolution of Stokes vector

Muller matrix for an optical element \mathfrak{M} s.t.

$$\mathfrak{M} = \begin{bmatrix} \mu_{11} & \cdots & \mu_{14} \\ \vdots & \ddots & \vdots \\ \mu_{41} & \cdots & \mu_{44} \end{bmatrix}$$

Evolution of Stokes vector as, $S_{out} = \mathfrak{M} S_{in}$

- Composition rule: $\mathfrak{M} = \mathfrak{M}_1 \mathfrak{M}_2 \dots \mathfrak{M}_n$
- Frame rotation by θ : $\mathfrak{M}_{\theta} = T^{-1}(\theta) \mathfrak{M} T(\theta)$ where

$$T(\theta) = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos 2\theta & \sin 2\theta & 0 \\ 0 & -\sin 2\theta & \cos 2\theta & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

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- ightharpoonup Composition rule: $\mathfrak{M} = \mathfrak{M}_1 \, \mathfrak{M}_2 \dots \mathfrak{M}_n$
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For elliptically polarized light, Stokes vector

$$\boldsymbol{S} = S_0 \begin{bmatrix} 1 \\ \cos 2\chi \cos 2\psi \\ \cos 2\chi \sin 2\psi \\ \sin 2\chi \end{bmatrix} \longrightarrow \underbrace{S_0(\cos 2\chi \cos 2\psi, \cos 2\chi \sin 2\psi, \sin 2\psi)}_{\text{On Poincare sphere of radius } S_0}$$

where azimuth (ψ) and ellipticity (χ) of polarization ellipse.

For elliptically polarized light, Stokes vector

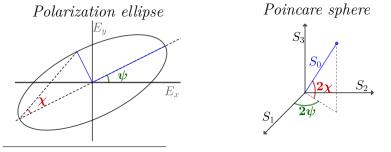
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Wave Optics (2015), Born, M; Wolf, E (1999)

For un-polarized light,

$$oldsymbol{S} = S_0 egin{bmatrix} 1 \ 0 \ 0 \ 0 \end{bmatrix} \longrightarrow \underbrace{(0,0,0)}_{ ext{At Origin}}$$

For partially polarized light,

$$\mathbf{S} = \begin{bmatrix} S_0 \\ S_1 \\ S_2 \\ S_3 \end{bmatrix} \longrightarrow \underbrace{(S_1, S_2, S_3)}_{\text{Inside sphere } s.t.}$$
$$0 < S_1^2 + S_2^2 + S_3^2 < S_0^2$$

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Gaussian Beam and properties

- ► Paraxial wave
- ► Gaussian beam solution and properties
- ► Modes of Gaussian beam

Maxwell's wave equation:

$$\nabla^2 \mathbf{E}(\mathbf{r}, t) - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \mathbf{E}(\mathbf{r}, t) = 0$$

Paraxial beam propagating predominantly in z-direction,

$$\boldsymbol{E}(x,y,z,t) = \boldsymbol{\psi}(x,y,z)e^{i(\omega t - kz)}$$

and taking slowly varying amplitude approx. i.e.

$$\left| \frac{\partial^2 \psi}{\partial z^2} \right| \ll k \left| \frac{\partial \psi}{\partial z} \right| \ll k^2 \left| \psi \right|$$

Paraxial wave equation

$$\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} - 2ik \frac{\partial \psi}{\partial r} = 0$$

One of the solutions is Gaussian beam.

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Gaussian Beam

- ► Paraxial wave
- ► Gaussian beam solution and properties
- ► Modes of Gaussian beam

Gaussian beam solution

Ansatz:
$$\psi(\mathbf{r}, z) = A \exp \left[-i \left(p(z) + \frac{kr^2}{2q(z)} \right) \right]$$

$$\psi(\mathbf{r}, z) = A\left(\frac{w_0}{w(z)}\right) \exp\left(\tan^{-1}\left(\frac{z}{z_0}\right) - i\frac{kr^2}{2R(z)} - \frac{r^2}{w^2(z)}\right)$$

Gaussian beam solution

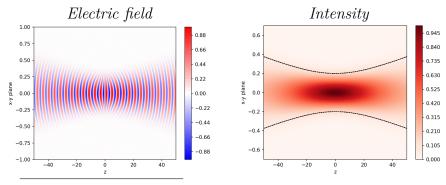
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Milonni, P.W (2010), Kogelnik, H. (1966), ECE 4300- Cornell

$$\psi(\boldsymbol{r},z) = A \underbrace{\left(\frac{\boldsymbol{w}_0}{w(z)}\right)}_{\text{term I}} \exp \left(i \tan^{-1} \left(\frac{z}{z_0}\right) - i \frac{kr^2}{2R(z)} - \frac{r^2}{w^2(z)}\right)$$

Term I related to **spreading of beam**.

 $w \to \text{Physical radius}$ $w_0 \to \text{Beam waist}$

$$w(z) = w_0 \sqrt{1 + \left(\frac{z}{z_0}\right)^2}$$

$$z_0 = \frac{\pi w_0^2}{\lambda}$$

$$\psi(\boldsymbol{r},z) = A \underbrace{\left(\frac{\boldsymbol{w}_0}{w(z)}\right)}_{\text{term I}} \exp\left(i \tan^{-1}\left(\frac{z}{z_0}\right) - i \frac{kr^2}{2R(z)} - \frac{r^2}{w^2(z)}\right)$$

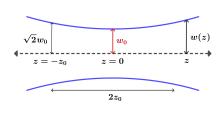
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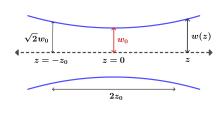
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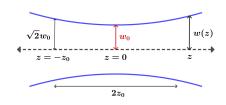
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$$z_0 = \frac{\pi w_0^2}{\lambda}$$



$$\psi(\boldsymbol{r},z) = A \underbrace{\left(\frac{\boldsymbol{w}_0}{w(z)}\right)}_{\text{term I}} \exp\left(i \tan^{-1}\left(\frac{z}{z_0}\right) - i \frac{kr^2}{2R(z)} - \frac{r^2}{w^2(z)}\right)$$

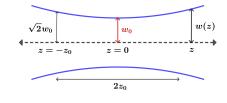
Term I related to **spreading of beam**.

 $w \to \text{Physical radius}$

 $w_0 \to \text{Beam waist}$

$$w(z) = w_0 \sqrt{1 + \left(\frac{z}{z_0}\right)^2}$$

 $z_0 \to \text{Rayleigh length}$



$$z_0 = \frac{\pi w_0^2}{\lambda}$$

Milonni, P.W (2010), Kogelnik, H. (1966), ECE 4300- Cornell

$$\psi(\mathbf{r}, z) = A\left(\frac{w_0}{w(z)}\right) \exp\left(\underbrace{i \tan^{-1}\left(\frac{z}{z_0}\right)}_{\text{term II}} - i\frac{kr^2}{2R(z)} - \frac{r^2}{w^2(z)}\right)$$

Term II related to Gouy phase (ϕ_G) .

$$\phi_G = \tan^{-1} \left(\frac{z}{z_0} \right)$$

$$\psi(\mathbf{r}, z) = A\left(\frac{w_0}{w(z)}\right) \exp\left(\underbrace{i \tan^{-1}\left(\frac{z}{z_0}\right)}_{\text{term II}} - i\frac{kr^2}{2R(z)} - \frac{r^2}{w^2(z)}\right)$$

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Term III related to radius of curvature (R) of beam wave-front.

$$R(z) = z \left[1 + \left(\frac{z_0}{z} \right)^{2^{-1}} \right]$$

$$\psi(\mathbf{r}, z) = A\left(\frac{w_0}{w(z)}\right) \exp\left(i \tan^{-1}\left(\frac{z}{z_0}\right) \underbrace{-i\frac{kr^2}{2R(z)}}_{\text{term III}} - \frac{r^2}{w^2(z)}\right)$$

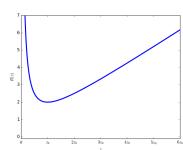
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Term IV related to Gaussian intensity profile.

$$I(r,z) \sim \exp\left(-\frac{2r^2}{w^2(z)}\right)$$

$$\psi(\mathbf{r}, z) = A\left(\frac{w_0}{w(z)}\right) \exp\left(i \tan^{-1}\left(\frac{z}{z_0}\right) - i \frac{kr^2}{2R(z)} \underbrace{-\frac{\mathbf{r}^2}{\mathbf{w}^2(z)}}_{\text{term IV}}\right)$$

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Term IV related to Gaussian intensity profile.

$$I(r,z) \sim \exp\left(-\frac{2r^2}{w^2(z)}\right)$$

$$\psi(\mathbf{r}, z) = A \exp \left[-i \left(p(z) + \frac{kr^2}{2 \mathbf{q}(z)} \right) \right]$$

 $q(z) \longrightarrow \text{characteristic of a beam if } \lambda \text{ known.}$

$$q(z) = z + i z_0$$

$$\frac{1}{q(z)} = \frac{1}{R(z)} - i \frac{\lambda}{\pi w^2(z)}$$

$$Q_{in} \longrightarrow \underbrace{\begin{bmatrix} A & B \\ C & D \end{bmatrix}}_{Optical} \longrightarrow q_{out} = \frac{Aq_{in} + B}{Cq_{in} + D}$$
Optical element

$$\psi(\mathbf{r}, z) = A \exp \left[-i \left(p(z) + \frac{kr^2}{2 \mathbf{q}(z)} \right) \right]$$

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$$\begin{array}{c}
A & B \\
C & D
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Optical element

Gaussian Beam and properties

- ► Paraxial wave
- ► Gaussian beam solution and properties
- ► Modes of Gaussian beam

Hermite-Gaussian mode

$$\psi_{m,n}(\mathbf{r},z) = A\left(\frac{w_0}{w(z)}\right) H_m\left(\frac{\sqrt{2}x}{w(z)}\right) H_n\left(\frac{\sqrt{2}y}{w(z)}\right) \exp\left(-\frac{r^2}{w^2(z)}\right) \cdot \exp\left(i\left(m+n+1\right)\tan^{-1}\left(\frac{z}{z_0}\right) - i\frac{kr^2}{2R(z)}\right)$$

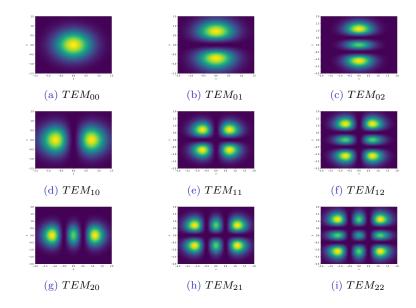
 $m, n = 0 \implies \psi = \text{Gaussian}$

Hermite-Gaussian mode

$$\psi_{m,n}(\mathbf{r},z) = A\left(\frac{w_0}{w(z)}\right) H_m\left(\frac{\sqrt{2}x}{w(z)}\right) H_n\left(\frac{\sqrt{2}y}{w(z)}\right) \exp\left(-\frac{r^2}{w^2(z)}\right) \cdot \exp\left(i\left(m+n+1\right)\tan^{-1}\left(\frac{z}{z_0}\right) - i\frac{kr^2}{2R(z)}\right)$$

 $m, n = 0 \implies \psi = \text{Gaussian}$

Hermite-Gaussian Intensity profile



Laguerre-Gaussian mode

$$\psi_{p,l}(r,\phi,z) = A \frac{w_0}{w(z)} \left[\frac{r\sqrt{2}}{w(z)} \right]^{|l|} L_p^{|l|} \left(\frac{2r^2}{w^2(z)} \right) \exp\left(-\frac{r^2}{w^2(z)} \right) \cdot \exp\left(-il\phi + i(2p+l+1)\tan^{-1}\left(\frac{z}{z_0}\right) - i\frac{kr^2}{2R(z)} \right)$$

 $l, p = 0 \implies \psi = \text{Gaussian}$

$$\exp(-il\phi) \longrightarrow \text{Helical phase}$$
(carries OAM)

Milonni, P.W (2010), Kogelnik, H. (1966), Bliokh, K. Y.; Rodríguez-Fortuño, F. J. (2015)

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$$l, p = 0 \Rightarrow \psi = Gaussian$$

$$\exp(-il\phi) \longrightarrow \text{Helical phase}$$
(carries OAM)

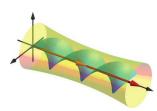
Milonni, P.W (2010), Kogelnik, H. (1966), Bliokh, K. Y.; Rodríguez-Fortuño, F. J. (2015)

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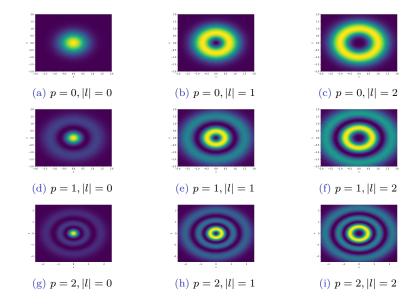
Helical phase front

$$\exp(-il\phi) \longrightarrow \text{Helical phase}$$
 (carries OAM)

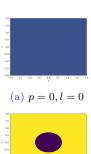


Milonni, P.W (2010), Kogelnik, H. (1966), Bliokh, K. Y.; Rodríguez-Fortuño, F. J. (2015)

Laguerre-Gaussian Intensity profile



Laguerre-Gaussian Phase profile

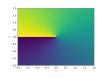








(g)
$$p = 2, l = 0$$



(b) p = 0, l = 1



(e)
$$p = 1, l = 1$$



(h)
$$p = 2, l = 1$$







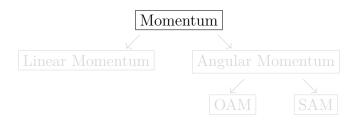
(f) p = 1, l = 2

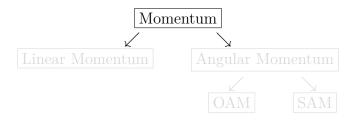


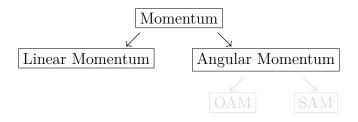
(i)
$$p = 2, l = 2$$

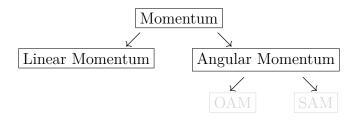
Spin-orbit interaction

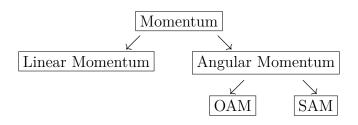
- ► Momentum of Light
- ► Geometric phase of light
- ► SOI in anisotropic medium











Linear momentum of light

Monochromatic beam with angular frequency ω propagating in z-direction :

$$\boldsymbol{\mathcal{E}}(\boldsymbol{r},t) = \operatorname{Re} \left\{ \boldsymbol{E}(\boldsymbol{r}) e^{-(\omega t - kz)} \right\}$$

$$\boldsymbol{\mathcal{B}}(\boldsymbol{r},t) = \mathrm{Re} \Big\{ \boldsymbol{B}(\boldsymbol{r}) e^{-i(\omega t - kz)} \Big\}$$

Maxwell-Faraday law:

$$\nabla \times \boldsymbol{E} = i\omega \boldsymbol{B}$$

Time-averaged linear momentum per length,

$$\mathcal{P}_z = \frac{1}{c^2} \int d\tau \ \langle \mathbf{S} \rangle_z = \frac{\epsilon_0}{2i\omega} \iint dx \, dy \ \langle \mathbf{E} \times (\nabla \times \mathbf{E}) \rangle_z$$

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Time-averaged AM per length:

$$\mathcal{J}_z = \frac{1}{c^2} \int d\tau [\mathbf{r} \times \langle \mathbf{S} \rangle]_z = \frac{\epsilon_0}{2i\omega} \iint dx dy [\mathbf{r} \times \langle \mathbf{E} \times (\nabla \times \mathbf{E}) \rangle]_z$$

For paraxial beam,

$$\mathcal{J}_{z} = \frac{\epsilon_{0}}{2i\omega} \iint dx \, dy \, \left[E_{\xi}^{*} \left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right) E_{\xi} \right]_{\xi=x,y} + \frac{\epsilon_{0}}{2i\omega} \iint dx \, dy \, (E_{x}^{*} E_{y} + E_{y}^{*} E_{x})$$

Time-averaged AM per length:

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Allen, L. (1999), Wave Optics (2015), S.J. van Enk; G. Nienhuis (1992), Cohen-Tannoudji, C., Dupont-Roc, J. (1989)

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Orbital AM, \mathcal{L}

Time-averaged AM per length:

$$\mathcal{J}_z = \frac{1}{c^2} \int d\tau [\boldsymbol{r} \times \langle \boldsymbol{S} \rangle]_z = \frac{\epsilon_0}{2i\omega} \iint dx dy \left[\boldsymbol{r} \times \langle \boldsymbol{E} \times (\nabla \times \boldsymbol{E}) \rangle \right]_z$$

For paraxial beam,

$$\mathcal{J}_{z} = \frac{\epsilon_{0}}{2i\omega} \iint dx \, dy \, \left[E_{\xi}^{*} \left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right) E_{\xi} \right]_{\xi=x,y} + \frac{\epsilon_{0}}{2i\omega} \iint dx \, dy \, \left(E_{x}^{*} E_{y} + E_{y}^{*} E_{x} \right)$$

Spin AM, \mathcal{S}

$$\mathcal{L} = \frac{\epsilon_0}{2i\omega} \iint dx \, dy \, \left[E_{\xi}^* \left(\boldsymbol{r} \times \nabla \right)_z E_{\xi} \right]_{\xi = x, y}$$

$$m{E}(r,\phi) = u(r) \; \exp(-il\phi) \; \hat{m{g}}$$
 $m{\mathcal{W}}_z = rac{\epsilon_0}{2} \iint dx \; dy \; m{E}^* \cdot m{E}$ $m{\frac{\mathcal{L}}{\mathcal{W}_z}} = rac{ ext{OAM}}{ ext{Total energy}} = rac{l}{\omega}$

Allen, L. (1999), Wave Optics (2015), S.J. van Enk (1992), Berry, M. V.; Soskin, M. S. (1998)

$$\mathcal{L} = \frac{\epsilon_0}{2i\omega} \iint dx \, dy \, \left[E_{\xi}^* \left(\boldsymbol{r} \times \nabla \right)_z E_{\xi} \right]_{\xi = x, y}$$

$$\boldsymbol{E}(r,\phi) = u(r) \, \exp(-il\phi) \, \hat{\boldsymbol{p}}$$

$$\mathcal{W}_z = \frac{\epsilon_0}{2} \iint dx \, dy \, \boldsymbol{E}^* \cdot \boldsymbol{E}$$

$$\frac{\mathcal{L}}{\mathcal{W}_z} = \frac{\text{OAM}}{\text{Total energy}} = \frac{l}{\omega}$$

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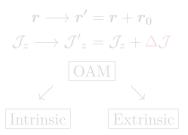
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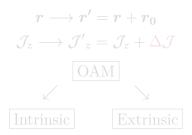
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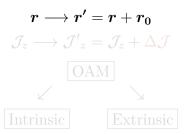


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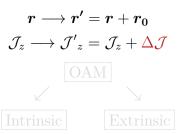
Allen, L. (1999), Wave Optics (2015), S.J. van Enk (1992), Berry, M. V.; Soskin, M. S. (1998)

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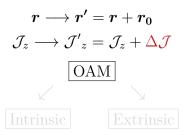
Allen, L. (1999), Wave Optics (2015), S.J. van Enk (1992), Berry, M. V.; Soskin, M. S. (1998)

$$\mathcal{L} = \frac{\epsilon_0}{2i\omega} \iint dx \, dy \, \left[E_{\xi}^* \left(\boldsymbol{r} \times \nabla \right)_z E_{\xi} \right]_{\xi = x, y}$$



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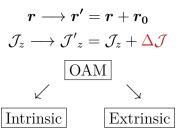
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$$egin{aligned} oldsymbol{r} & oldsymbol{r} & oldsymbol{r}' = oldsymbol{r} + oldsymbol{r}_0 \ \mathcal{J}_z & \longrightarrow \mathcal{J}'_z = \mathcal{J}_z + \Delta \mathcal{J} \ \hline & OAM \ & \swarrow & \searrow \ \hline & Intrinsic & Extrinsic \end{aligned}$$

Allen, L. (1999), Wave Optics (2015), S.J. van Enk (1992), Berry, M. V.; Soskin, M. S. (1998)

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 $oldsymbol{\checkmark}$
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$$\Delta \mathcal{J} \neq 0$$

Allen, L. (1999), Wave Optics (2015), S.J. van Enk (1992), Berry, M. V.; Soskin, M. S. (1998)

$$S = \frac{\epsilon_0}{2i\omega} \iint dx \, dy \, (E_x^* E_y + E_y^* E_x)$$

SAM is intrinsic.

For vortex beam,

$$\boldsymbol{E}(r,\phi) = u(r) \exp(-il\phi) \,\hat{\boldsymbol{p}}$$

$$\frac{\mathcal{S}}{\mathcal{W}_z} = \frac{\text{SAM}}{\text{Total energy}} = \frac{\sigma}{\omega}$$

$$\frac{\mathcal{J}_z}{\mathcal{W}_z} = \frac{\mathcal{L} + \mathcal{S}}{\mathcal{W}_z} = \frac{\text{Total AM}}{\text{Total energy}} = \frac{l + \sigma}{\omega}$$

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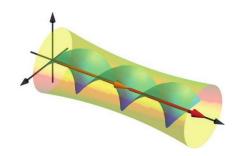
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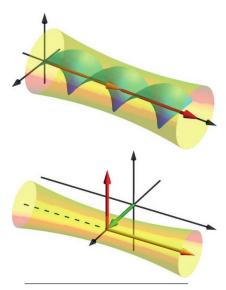
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Visualisation of AM

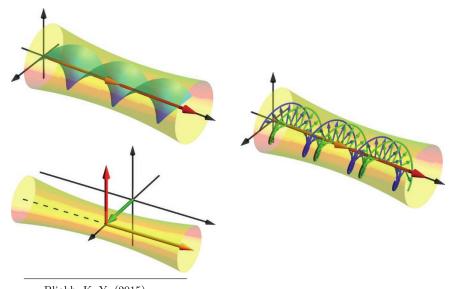


Visualisation of AM



Bliokh, K. Y. (2015)

Visualisation of AM

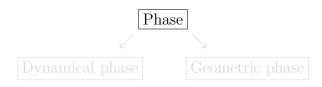


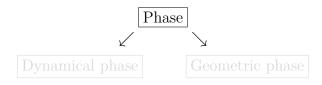
Bliokh, K. Y. (2015)

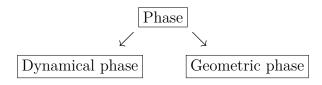
Spin-orbit interaction

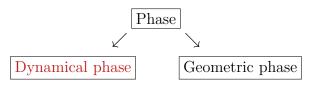
- ► Momentum of Light
- ► Geometric phase of light
- ► SOI in anisotropic medium

Geometric phase

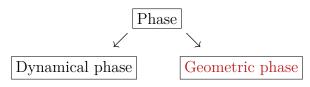




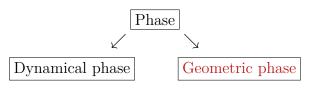




Associated with optical path length.



Associated with geometry of evolution.



Associated with geometry of evolution.

- ► Spin-redirection Berry phase
- ▶ Pancharatnam-Berry Phase

Associated with adiabatic evolution of wave-vector.

e.g., Polarized light through a helical optic fibre

$$\boldsymbol{J} \longrightarrow \boldsymbol{J'} = \boldsymbol{J} \exp(i \sigma \boldsymbol{\Theta})$$

$$\Theta = 2\pi(1 - \cos\theta)$$

 $\Theta \to$ solid angle obtained at the apex of the cone.

$$|L\rangle \longrightarrow e^{i\Theta} |L\rangle$$

 $|R\rangle \longrightarrow e^{-i\Theta} |R\rangle$

Associated with adiabatic evolution of wave-vector. e.g., Polarized light through a helical optic fibre.

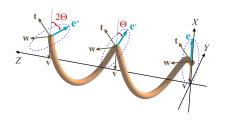
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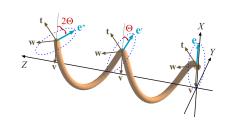
Associated with adiabatic evolution of wave-vector. e.g., Polarized light through a helical optic fibre.

$$J \longrightarrow J' = J \exp(i\sigma\Theta)$$

Helicity-dependant geometric phase

$$\Theta = 2\pi (1 - \cos \theta)$$

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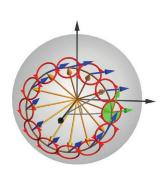
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Parallel transport

Associated with adiabatic evolution of wave-vector. e.g., Polarized light through a helical optic fibre.

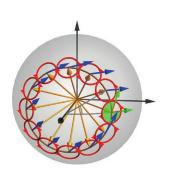
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Parallel transport

Associated with cyclic evolution in Poincare sphere keeping wave-vector fixed.

e.g., Michelson interferometer.

QP2 \rightarrow movable (aligned at β)

$$J_A = |x\rangle$$

$${m J}_A \longrightarrow {m J}_A'$$

$$\mathbf{J}_A' = |x\rangle \exp(i\phi_d) \exp(-i 2\varphi)$$

$$\varphi = \beta + \pi/4$$

Pancharatnam-Berry Phase

 $(\varphi \text{ depends on QP2})$

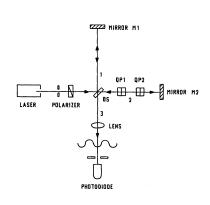
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QP1 \rightarrow fixed (aligned at $\pi/4$) QP2 \rightarrow movable (aligned at β)

$$egin{aligned} oldsymbol{J}_A &= |x\rangle \ oldsymbol{J}_A &\longrightarrow oldsymbol{J}_A' \ oldsymbol{J}_A' &= |x\rangle \exp(i\phi_d) \, \exp(-i\,2\varphi) \ &arphi &= \beta + \pi/4 \end{aligned}$$

Pancharatnam-Berry Pha $(\varphi \text{ depends on QP2})$



Associated with cyclic evolution in Poincare sphere keeping wave-vector fixed.

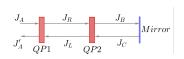
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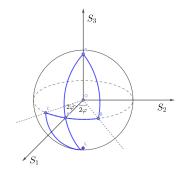
QP1 \rightarrow fixed (aligned at $\pi/4$) QP2 \rightarrow movable (aligned at β)

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Pancharatnam-Berry Phase (φ depends on QP2 alignment)





Associated with cyclic evolution in Poincare sphere keeping wave-vector fixed.

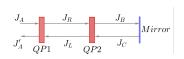
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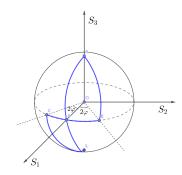
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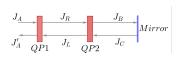
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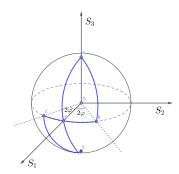
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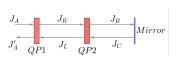
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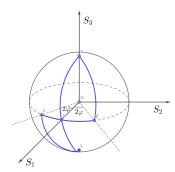
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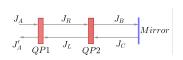
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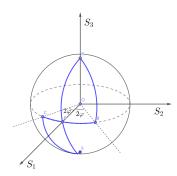
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Pancharatnam-Berry Phase (φ depends on QP2 alignment)





Wave Optics (2015), Chyba, T. H.; Wang, L. J. (1988)

Associated with cyclic evolution in Poincare sphere keeping wave-vector fixed.

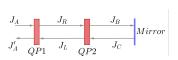
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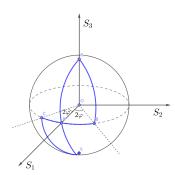
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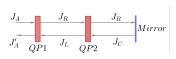
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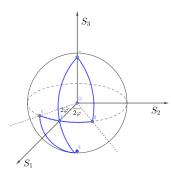
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alignment)





Wave Optics (2015), Chyba, T. H.; Wang, L. J. (1988)

Spin-orbit interaction

- ► Momentum of Light
- ► Geometric phase of light
- ► SOI in anisotropic medium

Three types of AM:

- ► IOAM
- ► EOAM
- ► SAM

Inter-conversion between AM in a process represents spin-orbit interaction of light

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Inter-conversion between AM in a process represents spin-orbit interaction of light

e.g. Quarter wave-plate

 $\pm\hbar$ SAM per photon is transferred to the wave plate. e.g. Half wave-plate

Circularly polarized
$$\longrightarrow$$
 W \longrightarrow polarized $(\sigma = \pm 1)$ P \longrightarrow $(\sigma = \mp 1)$

 $\pm 2\hbar$ SAM per photon is transferred to the wave plate and spin is flipped.

Wave Optics (2015)

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e.g. **q-plate**

Inhomogeneous orientation of the fast axis varying with azimuth (ϕ) .

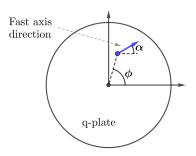
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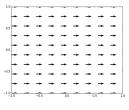
e.g. **q-plate**

Inhomogeneous orientation of the fast axis varying with azimuth (ϕ) .

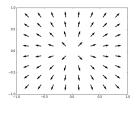
$$\alpha(\phi) = q\phi + \alpha_0$$



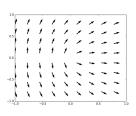
q-plate

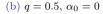


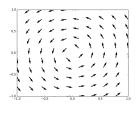




(a)
$$q = 1$$
, $\alpha_0 = 0$







(b)
$$q = 1$$
, $\alpha_0 = \pi/2$

Q-plate of phase retardation of π

$$\begin{bmatrix} 1 \\ i\sigma \end{bmatrix} \longrightarrow \mathbf{Q}_{\lambda/2} \end{bmatrix} \longrightarrow \begin{bmatrix} 1 \\ -i\sigma \end{bmatrix} \underbrace{\exp(i2\sigma q\phi)}_{\text{Vortex}} \exp(i2\sigma\alpha_0)$$

$$= \pm 1, l = 0) \longrightarrow (\sigma = \pm 1, l = \pm 2a)$$

 $q = 1 \rightarrow$ Angular momentum per photon is conserved **Q-plate of phase retardation of** $\pi/2$

$$\begin{bmatrix} 1 \\ i\sigma \end{bmatrix} \longrightarrow \boxed{\mathbf{Q}_{\lambda/4}} \longrightarrow \begin{bmatrix} \cos(\alpha - \sigma\pi/4) \\ \sin(\alpha - \sigma\pi/4) \end{bmatrix} \underbrace{\exp(i\sigma q\phi)}_{\text{Vortex}} \exp(i\sigma\alpha_0)$$

 $\alpha = 1 \rightarrow \text{Angular momentum per photon is conse$

Marrucci, L. (2006), Wave Optics (2015)

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$$(\sigma = \pm 1, l = 0) \longrightarrow (\sigma = \mp 1, l = \pm 2q)$$

$$(0 = \pm 1, t = 0) \longrightarrow (0 = \pm 1, t = \pm 2q)$$

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 $q=1 \to \text{Angular momentum per photon is conserved}.$

THANKS

TO ALL