

For John, notes on TMDs in b space

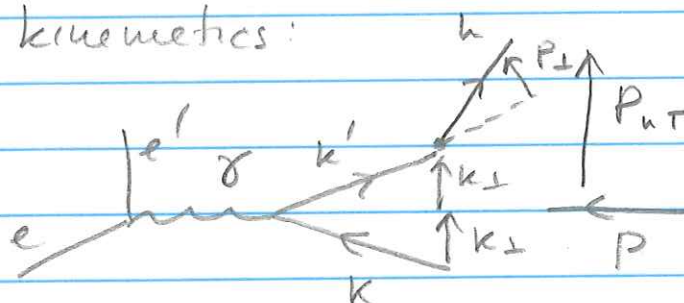
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In TMD formalism all structure functions can be expressed in terms of convolutions of TMDs. For instance, the unpolarised SIDIS structure function F_{UU} is

$$F_{UU} = C [f_1 D_1] \text{ where}$$

$$C[\] \equiv x \sum_a e_a^2 \int d^2 k_\perp d^2 p_\perp \delta^{(2)}(z \vec{k}_\perp + \vec{p}_\perp - \vec{P}_{hT}) f_1(x, k_\perp^2) D_1(z, p_\perp^2)$$

and kinematics:



$$P_h \approx z k' + p_\perp$$

$$\Rightarrow \vec{P}_{hT} \approx z \vec{k}_\perp + \vec{p}_\perp$$

now let us write F_{UU} in b space:

$$F_{UU} = x \sum_a e_a^2 \int d^2 k_\perp d^2 p_\perp \frac{1}{z^2} \delta^{(2)}\left(\vec{k}_\perp + \frac{\vec{p}_\perp}{z} - \frac{\vec{P}_{hT}}{z}\right) f_1(x, k_\perp^2) D_1(z, p_\perp^2)$$

$$\int \frac{d^2 b}{(2\pi)^2} e^{i \vec{b} \cdot \left(\vec{k}_\perp + \frac{\vec{p}_\perp}{z} - \frac{\vec{P}_{hT}}{z}\right)}$$

$$F_{UU} = x \sum_a e_a^2 \int d^2 b e^{-i \vec{b} \cdot \vec{P}_{hT}/z} \int \frac{d^2 k_\perp}{2\pi} e^{i \vec{b} \cdot \vec{k}_\perp} f_1(x, k_\perp^2)$$

$$\int \frac{d^2 p_\perp}{2\pi z^2} e^{i \vec{b} \cdot \vec{p}_\perp/z} D_1(z, p_\perp^2)$$

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Now let me define

$$\tilde{f}_1(x, b^2) \equiv \int \frac{d^2 k_\perp}{2\bar{u}} e^{i\vec{b}\vec{k}_\perp} f_1(x, k_\perp^2)$$

$$\tilde{D}_1(z, b^2) \equiv \int \frac{d^2 p_\perp}{2\pi z^2} e^{i\vec{b}\vec{p}_\perp/2} D_1(z, p_\perp^2)$$

notice that I have $1/2\bar{u}$ in the definition which is an unusual choice for F.T. but I do it in order to have symmetry in formulas for F.T. and anti-F.T

$$\begin{aligned} \tilde{f}_1(x, b^2) &= \int \frac{k_\perp dk_\perp}{2\pi} \underbrace{\int d\varphi e^{i b k_\perp \cos \varphi}}_{2\bar{u} J_0(k_\perp b)} f_1(x, k_\perp^2) = \\ &= \int k_\perp dk_\perp J_0(k_\perp b) f_1(x, k_\perp^2) \end{aligned}$$

$$\tilde{D}_1(z, b^2) = \int p_\perp dp_\perp J_0\left(\frac{p_\perp b}{2}\right) \frac{1}{z^2} D_1(z, p_\perp^2)$$

and we finally have

$$F_{un} = 2\bar{u} \times \sum_a e_a^2 \int b db J_0\left(\frac{b p_{qT}}{2}\right) \tilde{f}_1(x, b^2) \tilde{D}_1(z, b^2)$$

How do we parametrise f_1 and D_1 ?
Usual choice is gaussian parametrisation

$$f_1(x, k_\perp) = f_1(x) \frac{1}{\pi \langle k_\perp^2 \rangle} e^{-k_\perp^2 / \langle k_\perp^2 \rangle}$$

(notice $\int d^2 k_\perp f_1(x, k_\perp) = f_1(x) \leftarrow$ collinear PDF)

$$\text{so, } \tilde{f}(x, b^2) = \frac{f_1(x)}{2\pi} e^{-\frac{b^2 \langle k_\perp^2 \rangle}{4}}$$

$$D_1(z, p_\perp^2) = D_1(z) \frac{1}{\pi \langle p_\perp^2 \rangle} e^{-p_\perp^2 / \langle p_\perp^2 \rangle}$$

$$\text{so, } \tilde{D}_1(z, b^2) = \frac{D_1(z)}{2\pi z^2} e^{-\frac{b^2 \langle p_\perp^2 \rangle}{4z^2}}$$

It is known that \tilde{f}_1 and \tilde{D}_1 obey Collins-Soper evolution equations, so that the Sudakov form factor should be added to above formulae, i.e

$$\tilde{f}_1(x, b^2) = \frac{f_1(x)}{2\pi} e^{-\frac{b^2 \langle k_\perp^2 \rangle}{4}} e^{-S_{\text{sud}}(b)/2}$$

$$\tilde{D}_1(z, b^2) = \frac{D_1(z)}{2\pi z^2} e^{-\frac{b^2 \langle p_\perp^2 \rangle}{4z^2}} e^{-S_{\text{sud}}(b)/2}$$

(notice that Sudakov is partitioned $1/2$ & $1/2$ between \tilde{f}_1 and \tilde{D}_1)

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$S_{\text{rad}}(b)$ can be calculated perturbatively

$$S_{\text{rad}}(b) \equiv \int_{\mu_b}^Q \frac{d\mu^2}{\mu^2} \left(A \ln \frac{Q^2}{\mu^2} + B \right)$$

$$\downarrow \qquad \qquad \downarrow$$

$$\sum \left(\frac{d_s}{2\pi} \right)^n A^{(n)} \qquad \sum \left(\frac{d_s}{2\pi} \right)^n B^{(n)}$$

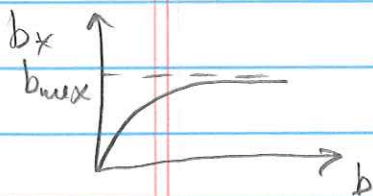
$$\mu_b = \frac{b_0}{b}, \quad b_0 = 2e^{-\gamma_E}$$

(notice that it also means that $f_2(x)$ and $D_1(z)$ should be evaluated at the scale of μ_b)

The problem is now with $b \rightarrow \infty$ as $\mu_b \rightarrow 0$ and we will have to deal with the Landau pole in d_s or \equiv non-perturbative physics.

A usual solution is to introduce a b_* prescription

$b_*(b)$ is a smooth function that maps all b onto $[0, b_{\text{max}}]$ region. A possible form



$$b_*(b) = \frac{b}{\sqrt{1 + b^2/b_{\text{max}}^2}}$$

$$b_*(0) = 0, \quad b_*(\infty) = b_{\text{max}}$$

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One can now write

$$e^{-S_{\text{sud}}(b)} = e^{-S_{\text{sud}}(b_*)} \underbrace{e^{-S_{\text{NP}}(b)}}_{\text{non perturbative Sudakov formfactor}}$$

We know that $S_{\text{NP}}(0) = 0$ and a possible parametrisation will be

$$S_{\text{NP}}(b) = g_2 \ln(b/b_*) \ln Q/Q_0$$

\uparrow a parameter $\uparrow \sqrt{1+b^2/b_{\text{max}}^2}$

$$S_{\text{sud}}(b_*) = \int_{\mu_{b_*}}^Q \frac{d\mu^2}{\mu^2} (A \ln Q^2/\mu^2 + B)$$

$$\mu_{b_*} = \frac{b_0}{b_*}$$

So that

$$\tilde{f}_2(x, b^2) = \frac{f_1(x, \mu_{b_*}^2)}{2u} e^{-\frac{b^2 \langle k_t^2 \rangle}{4} - \frac{g_2}{2} \ln \frac{b}{b_*} \ln \frac{Q}{Q_0} - \frac{S_{\text{sud}}(b_*)}{2}}$$

$$\tilde{D}_2(z, b^2) = \frac{D_2(z, \mu_{b_*}^2)}{2u \, z^2} e^{-\frac{b^2 \langle p_t^2 \rangle}{4z^2} - \frac{g_2}{2} \ln \frac{b}{b_*} \ln \frac{Q}{Q_0} - \frac{S_{\text{sud}}(b_*)}{2}}$$

$\langle k_t^2 \rangle$, $\langle p_t^2 \rangle$, g_2 are free parameters

Now lets talk of the data

We have multiplicities from Hermes & COMPASS

if we define $F_2 \equiv x \sum_q e_q^2 f(x, Q^2)$

Hermes measures $M_n^h(x, Q^2, z, P_{nT}) \equiv \frac{1}{\frac{dGPDs}{dx dQ^2}} \frac{d^4 G}{dx dQ^2 dz dP_{nT}}$

and

$$M_n^h(x, Q^2, z, P_{nT}) = 2\pi P_{nT} \frac{F_{nn}}{F_2}$$

COMPASS measures $\frac{d^2 n^h(x, Q^2, z, P_{nT})}{dz dP_{nT}} \equiv \pi \frac{F_{nn}}{F_2}$

The measurement is at low Q thus there should be little phase space for $S_{nd}(bx)$ and we will drop it assuming $S_{nd}(bx) \approx 0$ (at this energy)

We have a program that fits the dots, what we will need is a program that will perform numerically the b integration, see next page

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$$F_{un}(P_{\perp T}) = 2\bar{n} \times \sum_a e_a^2 \int b db J_0\left(\frac{b P_{\perp T}}{z}\right) \tilde{f}_1(x, b^*) \tilde{D}_1(z, b^*)$$

where

$$\tilde{f}_1(x, b^*) = \frac{f_1(x, \mu_{b^*})}{2\bar{n}} e^{-\frac{b^2 \langle k_{\perp}^2 \rangle}{4} - \frac{g_2}{2} \ln \frac{b}{b_*} \ln \frac{Q}{Q_0}}$$

$$\tilde{D}_1(z, b^*) = \frac{D_1(z, \mu_{b^*})}{2\bar{n} z^2} e^{-\frac{b^2 \langle p_{\perp}^2 \rangle}{4 z^2} - \frac{g_2}{2} \ln \frac{b}{b_*} \ln \frac{Q}{Q_0}}$$

$$b_* = \frac{b}{\sqrt{1 + b^2/b_{\text{max}}^2}}, \quad \underbrace{b_{\text{max}} = 1 \text{ (GeV}^{-1}\text{)}}_{\text{parameter}}$$

$$\mu_{b^*} = \frac{b_0}{b_*}, \quad b_0 = 2e^{-\delta_E}$$

In total we will have

$\langle k_{\perp}^2 \rangle, \langle p_{\perp}^2 \rangle, b_{\text{max}}, g_2 \rightarrow 4 \text{ parameters}$

↑ ↑
can be flavor dependent