

# N-Body Simulation of High-Density Magneto-Optical Traps (MOT)

Documentation & Algorithm Analysis

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## Abstract

This document provides a comprehensive description of the physics, numerical methods, and usage of the `mot_simulation.py` suite. The simulation models the dynamics of large  $^{87}\text{Rb}$  atom clouds ( $N \sim 10^7 - 10^9$ ) in the high-density regime. It specifically addresses the transition from the temperature-limited regime to the density-limited regime dominated by collective effects: **Shadowing** (attenuation) and **Radiative Rescattering**. To make these  $O(N^2)$  interactions computationally feasible, the code employs a **Superparticle** approximation combined with **Particle-Mesh (FFT)** algorithms.

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# 1 Physics of the Magneto-Optical Trap

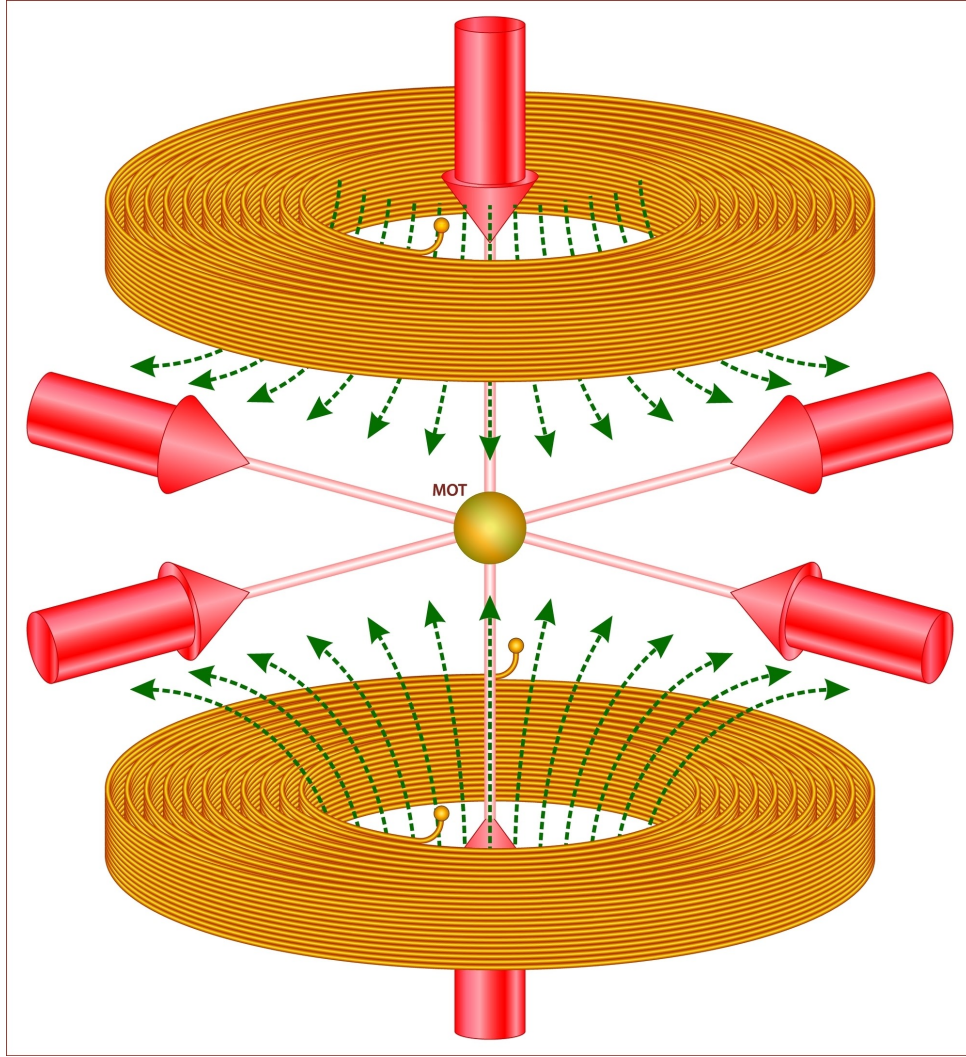


Figure 1: Visualization of a magneto-optical trap setup.

The Magneto-Optical Trap (MOT) relies on the interplay between dissipative optical forces (Doppler cooling) and position-dependent restoring forces arising from an inhomogeneous magnetic field (Zeeman effect).

## 1.1 Doppler Cooling (The Friction Force)

The simulation assumes a simplified two-level atomic structure (ground state  $J = 0$  to excited state  $J' = 1$ ). The atoms are illuminated by six counter-propagating laser beams along the Cartesian axes. The laser frequency  $\omega_L$  is detuned below the atomic resonance  $\omega_0$  (red detuning,  $\delta = \omega_L - \omega_0 < 0$ ).

For an atom moving with velocity  $\vec{v}$ , the laser frequency seen by the atom is shifted by the Doppler effect:  $\omega' = \omega_L - \vec{k} \cdot \vec{v}$ .

- If the atom moves *towards* a laser beam, the light is Doppler-shifted closer to resonance, increasing the absorption rate.
- If the atom moves *away*, the light is shifted further from resonance, decreasing absorption.

The net result is a radiation pressure force that opposes velocity, acting as a viscous drag:  $\vec{F}_{cool} \approx -\beta\vec{v}$ .

## 1.2 Spatial Confinement (The Restoring Force)

To trap atoms, the force must be position-dependent. This is achieved using a quadrupole magnetic field generated by anti-Helmholtz coils:

$$\vec{B}(x, y, z) = B'(x\hat{x} + y\hat{y} - 2z\hat{z}) \quad (1)$$

where  $B'$  is the magnetic field gradient.

This field causes a spatial variation in the atomic energy levels due to the **Zeeman Effect**. The detuning becomes position-dependent:

$$\delta_{\text{eff}}(\vec{r}) = \delta - \vec{k} \cdot \vec{v} - \frac{\mu_B g_F}{\hbar} |\vec{B}(\vec{r})| \quad (2)$$

By choosing correct circular polarizations ( $\sigma^+$ ,  $\sigma^-$ ) for the laser beams, atoms further from the center are brought closer to resonance with the beam pushing them back towards the center. This creates a harmonic restoring force:  $\vec{F}_{trap} \approx -\kappa\vec{r}$ .

## 2 Collective Effects in High-Density Regimes

Standard MOT theory assumes atoms are independent. However, when the number of atoms  $N$  exceeds  $\sim 10^6$  or density  $n$  exceeds  $10^{10} \text{ cm}^{-3}$ , the cloud becomes optically thick. Two competing collective effects emerge.

### 2.1 The Shadow Effect (Attenuation)

As laser beams propagate through the cloud, photons are absorbed by the outer layers of atoms. This creates a "shadow" in the center of the trap.

- **Mechanism:** The intensity of each beam  $I_\alpha$  decays according to the Beer-Lambert law:

$$\frac{dI_\alpha}{dz} = -n(\vec{r})\sigma_{sc}I_\alpha \quad (3)$$

- **Consequence:** The restoring force from the "far side" beam is weaker than the "near side" beam. This imbalance creates a net compressive force that squeezes the cloud, potentially leading to instabilities.

### 2.2 Radiative Rescattering (The Repulsive Force)

The most critical effect for the density limit is **Rescattering**. Photons absorbed by one atom are not destroyed; they are re-emitted via spontaneous emission. Ideally, these photons escape the cloud. In a dense cloud, however, a re-emitted photon has a high probability of being re-absorbed by a neighboring atom.

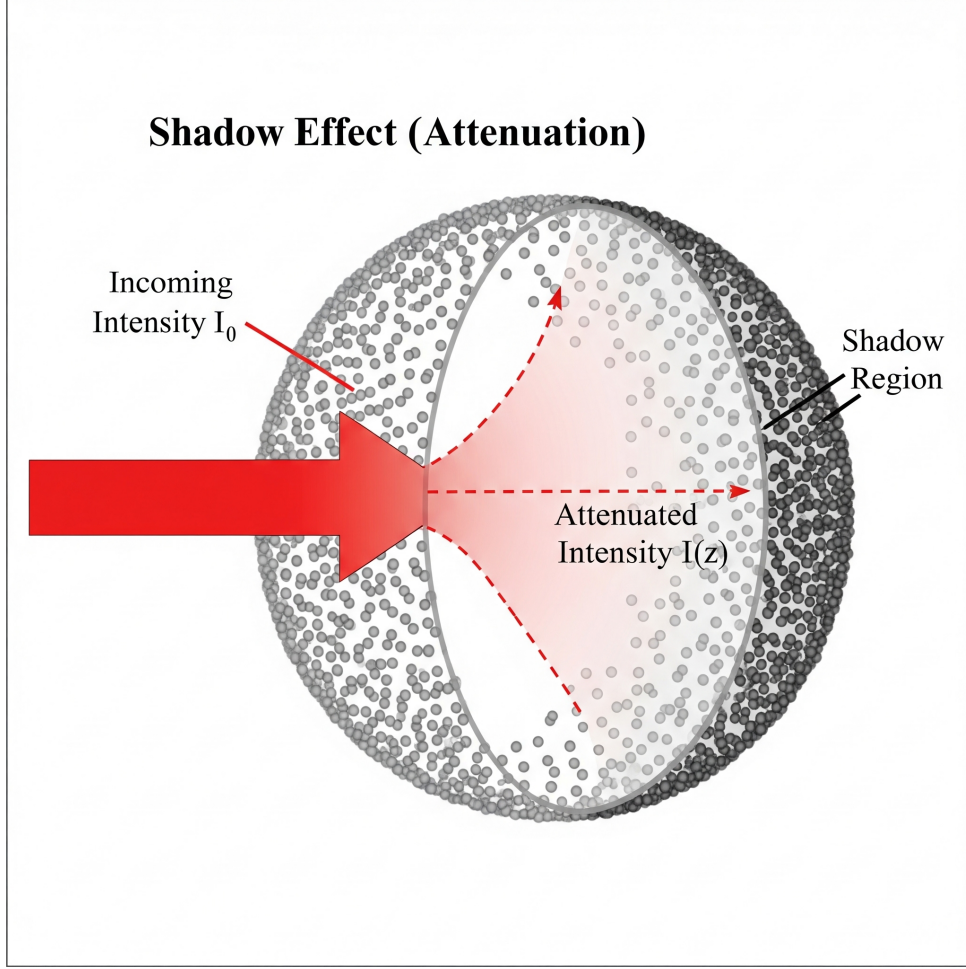


Figure 2: Visualization of shadowing.

### 2.2.1 Derivation of the Force

Consider an atom at position  $\vec{r}_j$  with a total scattering rate  $\Gamma_{sc}$ . This atom acts as a point source of radiation, re-emitting power  $P_j$ :

$$P_j = \hbar\omega_L \times \Gamma_{sc}(\vec{r}_j) \quad (4)$$

This power propagates outwards as a spherical wave. The intensity  $I_{resc}$  at a distance  $r$  from the source atom is given by the inverse-square law:

$$I_{resc}(r) = \frac{P_j}{4\pi r^2} \quad (5)$$

A second atom at position  $\vec{r}_i$  will absorb this scattered radiation. The force exerted on atom  $i$  is the momentum transfer rate from this rescattered field:

$$\vec{F}_{ij} = \sigma_R \frac{I_{resc}(|\vec{r}_i - \vec{r}_j|)}{c} \hat{n}_{ij} \quad (6)$$

where  $\sigma_R$  is the *rescattering cross-section* (which may differ from the laser cross-section due to spectral broadening) and  $c$  is the speed of light.

Combining these terms yields the full pairwise force equation:

$$\vec{F}_{ij} = \frac{\sigma_R P_j}{4\pi c} \frac{\vec{r}_i - \vec{r}_j}{|\vec{r}_i - \vec{r}_j|^3} \quad (7)$$

# Radiative Rescattering Force

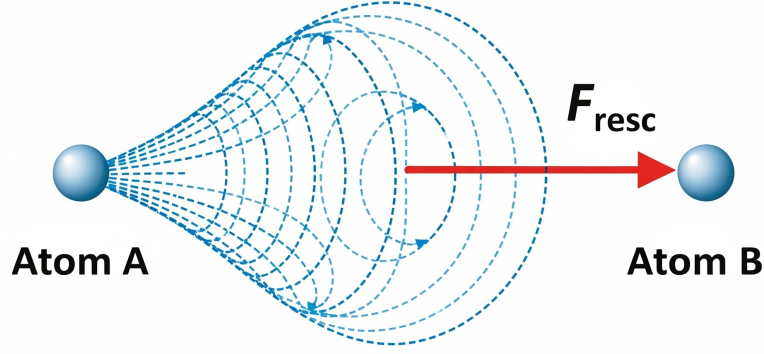


Figure 3: Visualization of re-scattering force.

## 2.2.2 Electrostatic Analogy

This equation is mathematically isomorphic to **Coulomb's Law** for electrostatics:

$$\vec{F}_{el} = \frac{1}{4\pi\epsilon_0} \frac{q_i q_j}{r^2} \hat{r} \quad (8)$$

By comparing the constants, we can define an "effective charge" for the atoms. If we assume the rescattering cross-section  $\sigma_R$  is uniform, the effective charge  $q_{\text{eff}}$  of an atom is proportional to the square root of its scattered power:

$$q_{\text{eff}} \sim \sqrt{\frac{\sigma_R P}{c}} \quad (9)$$

Consequently, a MOT in the multiple-scattering regime behaves like a uniformly charged gas of ions (a non-neutral plasma), subject to an external trapping potential.

## 2.2.3 The Density Limit

This repulsive force scales as  $1/r^2$ , effectively creating an outward "radiation pressure." The MOT cloud stops shrinking when this outward pressure balances the inward magneto-optical trapping pressure.

- **Low N:** Size is constant (Temperature limited). Density increases with  $N$ .
- **High N:** Density is constant (Density limited). The cloud size  $R$  expands as  $N^{1/3}$  to maintain constant density.

## 3 Numerical Algorithms

### 3.1 1. Time Integration: Velocity Verlet

**Complexity:**  $O(N)$

The simulation integrates the equations of motion using the Velocity Verlet algorithm, a symplectic integrator known for stability in conservative systems.

### 3.2 2. Beam Attenuation: Vectorized Sorting (Grid/Tube)

**Complexity:**  $O(N \log N)$

Calculating the shadow for every atom is naively  $O(N^2)$ . We approximate the cloud as a bundle of parallel 1D tubes aligned with the laser axes.

1. **Binning:** Atoms are binned into transverse grid cells  $(u, v)$ .
2. **Sorting:** Within each tube, atoms are sorted by their longitudinal position  $(w)$ .
3. **CumSum:** The optical depth is computed via a cumulative sum along the sorted array, reducing the complexity to that of the sorting algorithm.

### 3.3 3. Rescattering: Particle-Mesh (FFT)

**Complexity:**  $O(N + M \log M)$

Directly summing the pairwise Coulomb-like rescattering force is  $O(N^2)$ . For  $N = 10^5$ , this requires  $10^{10}$  operations per step, which is computationally prohibitive. We exploit the electrostatic analogy to solve the problem using a grid-based field solver (Particle-Mesh method).

#### 3.3.1 Theoretical Basis: The Poisson Equation

Since the force follows an inverse-square law, the divergence of the force field is proportional to the source density (Gauss's Law). We can describe the rescattering field via a scalar potential  $\Phi$ :

$$\vec{F}_{\text{resc}}(\vec{r}) = -\sigma_R \nabla \Phi(\vec{r}) \quad (10)$$

The potential  $\Phi$  satisfies the **Poisson Equation**:

$$\nabla^2 \Phi(\vec{r}) = -\frac{S(\vec{r})}{c} \quad (11)$$

where  $S(\vec{r})$  is the continuous density of scattered power (Watts per  $\text{m}^3$ ).

#### 3.3.2 Algorithm Implementation Steps

The Particle-Mesh (PM) algorithm solves this differential equation in four steps:

##### Step 1: Charge Assignment (Binning)

The continuous atom positions are discretized onto a regular 3D grid of size  $M \times M \times M$ . The scattered power  $P_j$  of each superparticle is distributed to the nearest grid nodes using a weighting scheme (e.g., Nearest Grid Point or Cloud-in-Cell). This yields the source term grid  $S_{ijk}$ .

### Step 2: Poisson Solver in Fourier Space

The Poisson equation is a convolution, which becomes a simple multiplication in Fourier space. We apply the Fast Fourier Transform (FFT) to the source grid:

$$\hat{S}(\vec{k}) = \text{FFT}(S_{ijk}) \quad (12)$$

In Fourier space, the Laplacian operator  $\nabla^2$  becomes  $-k^2$ . The potential is found by multiplying with the Green's function:

$$\hat{\Phi}(\vec{k}) = \frac{\hat{S}(\vec{k})}{k^2} \quad (\text{with } \hat{\Phi}(0) = 0) \quad (13)$$

### Step 3: Field Calculation

The force is the gradient of the potential. In Fourier space, the gradient operator  $\nabla$  corresponds to multiplication by  $i\vec{k}$ . We compute the three components of the force field vector  $\vec{E}$  directly in k-space:

$$\hat{\vec{E}}(\vec{k}) = -i\vec{k} \hat{\Phi}(\vec{k}) \quad (14)$$

We then apply the Inverse FFT (IFFT) to transform these field components back to real space, obtaining the force field on the grid  $\vec{F}_{grid}$ .

### Step 4: Force Interpolation

The force acting on a specific atom at position  $\vec{r}_i$  is computed by interpolating the values from the surrounding grid nodes in  $\vec{F}_{grid}$  (e.g., Trilinear Interpolation).

### 3.3.3 Complexity Analysis

The complexity is dominated by the FFT step, which scales as  $O(M^3 \log M)$ , where  $M$  is the grid resolution. Crucially, this is independent of the number of atoms  $N$  (once  $N$  is large enough).

- **Pairwise Summation:**  $O(N^2)$
- **Particle-Mesh:**  $O(N) + O(M^3 \log M)$

For a typical simulation with  $N = 10^5$  atoms and a  $32^3$  grid, the Particle-Mesh method is several orders of magnitude faster.

## 4 Usage Guide

### 4.1 Running the Simulation

```
python mot_simulation.py [N_atoms] [B_gradient] [T_init_uK] [R_init_m]
```

Example: `python mot_simulation.py 10000 0.1 1000 5e-4`

### 4.2 Post-Processing

Use the included `mot_postprocessing.py` script to generate graphs and animations.

```
python mot_postprocessing.py [N_atoms] [B_gradient] [T_init_uK] [R_init_m]
```

It generates:

- **mot\_evolution.gif:** Animation of cloud dynamics and temperature.
- **Virial.png:** Stability analysis comparing Kinetic Energy vs. Trapping Potential.