

Supplementary Materials: Inverse energy transfer in three-dimensional quantum vortex flows

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NUMERICAL METHOD

Using Schwarz mesoscopic model [1], vortex lines can be described as space curves $\mathbf{s}(\xi, t)$ of infinitesimal thickness, with a single quantum of circulation $\kappa = h/m_4 = 9.97 \times 10^{-8} \text{ m}^2/\text{s}$, where h is Planck's constant, $m_4 = 6.65 \times 10^{-27} \text{ kg}$ is the mass of one helium atom, ξ is the natural parameterization, arclength, and t is time. These conditions are a good approximation, since the vortex core radius of superfluid ${}^4\text{He}$ ($a_0 = 10^{-10} \text{ m}$) is much smaller than any of the length scales of interest in turbulent flows. The equation of motion is

$$\dot{\mathbf{s}}(\xi, t) = \mathbf{v}_s + \frac{\beta}{1 + \beta} [\mathbf{v}_{ns} \cdot \mathbf{s}'] \mathbf{s}' + \beta \mathbf{s}' \times \mathbf{v}_{ns} + \beta' \mathbf{s}' \times [\mathbf{s}' \times \mathbf{v}_{ns}], \quad (1)$$

where $\dot{\mathbf{s}} = \partial \mathbf{s} / \partial t$, $\mathbf{s}' = \partial \mathbf{s} / \partial \xi$ is the unit tangent vector, $\mathbf{v}_{ns} = \mathbf{v}_n - \mathbf{v}_s$, \mathbf{v}_n and \mathbf{v}_s are the normal fluid and superfluid velocities at \mathbf{s} and β, β' are temperature and Reynolds number dependent mutual friction coefficients [2]. The superfluid velocity \mathbf{v}_s at a point \mathbf{x} is determined by the Biot-Savart law

$$\mathbf{v}_s(\mathbf{x}, t) = \frac{\kappa}{4\pi} \oint_{\mathcal{T}} \frac{\mathbf{s}'(\xi, t) \times [\mathbf{x} - \mathbf{s}(\xi, t)]}{|\mathbf{x} - \mathbf{s}(\xi, t)|} d\xi, \quad (2)$$

where \mathcal{T} represents the entire vortex configuration. There is currently a lack of a well-defined theory of vortex reconnections in superfluid helium, like for the Gross-Pitaevskii equation [3–5]. An *ad hoc* vortex reconnection algorithm is employed to resolve the collisions of vortex lines [6].

A *two-way model* is crucial to understand the accurately interpret the back-reaction effect of the normal fluid on the vortex line and vice-versa [7]. We self-consistently evolve the normal fluid \mathbf{v}_n with a modified Navier-Stokes equation

$$\frac{\partial \mathbf{v}_n}{\partial t} + (\mathbf{v}_n \cdot \nabla) \mathbf{v}_n = -\nabla \frac{p}{\rho} + \nu_n \nabla^2 \mathbf{v}_n + \frac{\mathbf{F}_{ns}}{\rho_n}, \quad (3)$$

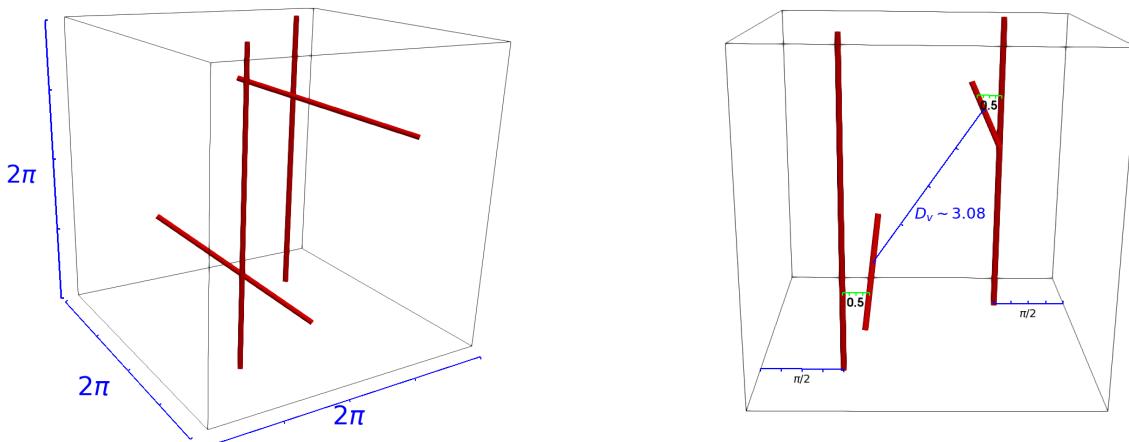


FIG. 1: Schematic diagram of the orthogonal vortex configuration.

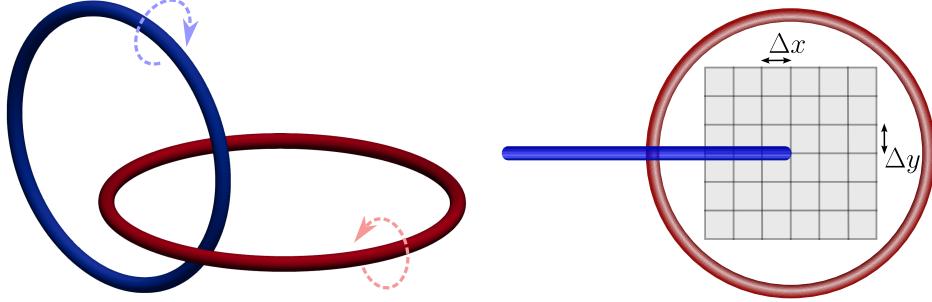


FIG. 2: Schematic diagram of the Hopf link vortex configuration.

$$\mathbf{F}_{ns} = \oint_{\mathcal{T}} \mathbf{f}_{ns} \delta(\mathbf{x} - \mathbf{x}) d\xi, \quad \nabla \cdot \mathbf{v}_n = 0, \quad (4)$$

where $\rho = \rho_n + \rho_s$ is the total density, ρ_n and ρ_s are the normal fluid and superfluid densities, p is the pressure, ν_n is the kinematic viscosity of the normal fluid and \mathbf{f}_{ns} is the local friction per unit length [8]

$$\mathbf{f}_{ns} = -\mathcal{D}\mathbf{s}' \times [\mathbf{s}' \times (\dot{\mathbf{s}} - \mathbf{v}_n)] - \rho_n \kappa \mathbf{s}' \times (\mathbf{v}_n - \dot{\mathbf{s}}), \quad (5)$$

where \mathcal{D} is a coefficient dependent on the vortex Reynolds number and intrinsic properties of the normal fluid. The regularization of the mutual friction force onto the normal fluid grid is physically motivated by the strongly localized injection of vorticity during the momentum exchange of point-like particles and viscous flow in classical fluid dynamics [9, 10]. In short, the localized vorticity induced by the relative motion between the vortex lines and the normal fluid is diffused to discretization of the grid spacing Δx in a time interval ϵ_R . In this way, the delta-forced friction as defined in Eq. 5 is regularized by a Gaussian function, the fundamental solution of the diffusion equation. Further details of the method for classical fluids are contained in [9, 10] and for FOUCAULT in [2].

In this Letter, we report all results using dimensionless units, where the characteristic length scale is $\tilde{\lambda} = D/D_0$, where $D^3 = (1 \times 10^{-3}\text{m})^3$ is the dimensional cube size, $D_0^3 = (2\pi)^3$ is the non-dimensional cubic computational domain. The time scale is given by $\tilde{\tau} = \tilde{\lambda}^2 \nu_n^0 / \nu_n$, where the non-dimensional viscosity ν_n^0 resolves the small scales of the normal fluid. In this work, we consider two vortex configurations - initially orthogonal vortices and Hopf links.

Orthogonal reconnection: The characteristic quantities are $\tilde{\lambda} = 1.59 \times 10^{-4}\text{m}$, $\nu_n^0 = 0.32$ and $\tilde{\tau} = 0.366\text{s}$ at $T = 1.9K$ and $\tilde{\tau} = 0.485\text{s}$ at $T = 2.1K$. The vortices are initialized as two pairs of orthogonal vortices, as shown in the schematic of Fig. 1. The separation between vortices in each pair d is set to be $d_v = 0.5$ in dimensionless units, and the shortest distance between pairs is $D_v = \sqrt{(\pi - d_v/2)^2 + \pi^2} \approx 3.08$, so that $d_v \ll D_v$. The Lagrangian discretization of vortex lines is $\Delta\xi = 0.025$ (a total of 1340 discretization points across the 4 vortex lines), using a timestep of $\Delta t_{VF} = 5.56 \times 10^{-6}$. For the normal fluid, a total of $N = 256^3$ mesh points were used, with a timestep of $\Delta t_{NS} = 45\Delta t_{VF}$.

Hopf link: The characteristic quantities are $\tilde{\lambda} = 1.59 \times 10^{-4}\text{m}$, $\nu_n^0 = 0.16$ and $\tilde{\tau} = 0.1836\text{s}$ at $T = 1.9K$ and $\tilde{\tau} = 0.2439\text{s}$ at $T = 2.1K$. Vortices are initialized as shown in Fig. 2, where the blue vortex ring is chosen at an initial offset $n_x \Delta x$ and $n_y \Delta y$ where $n_x, n_y \in \{-3, \dots, 3\}$ and $\Delta x = \Delta y = 0.125$ in units of the code. This gives a total of 49 individual reconnections for each temperature. Both of the rings have radius $R \approx 1$ also in units of the code. Furthermore, each reconnection is supplemented with a normal fluid ring around the superfluid vortex ring, which is generated by superimposing a normal fluid ring generated by a propagating ring of the same radius. In this way, we eliminate a transient phase of generating normal fluid structures. The Lagrangian discretization of the vortex lines is $\Delta\xi = 0.025$ (a total of 668 discretization points across both of the rings), using a timestep of $\Delta t_{VF} = 1.25 \times 10^{-5}$. For the normal fluid, a total of $N = 256^3$ mesh points were used, with a timestep of $\Delta t_{NS} = 40\Delta t_{VF}$.

HELICAL DECOMPOSITION OF THE HOPF LINK

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