

$$\partial_t u = \nabla \cdot (D_V \nabla u) \quad x \in \Omega_1 \quad (1)$$

$$D_V \nabla u \cdot \nu = k(u - \bar{u}) \quad x \in \Gamma \quad (2)$$

We assume D_V to be constant in Ω_1 .

$$D(x) = \chi_\varepsilon(x) D_V$$

Ansatz:

$$u_\varepsilon(x, y) = u_0(x, y) + \varepsilon u_1(x, y) + \varepsilon^2 u_2(x, y) + \mathcal{O}(\varepsilon^3) \quad (3)$$

Chain rule for the gradient:

$$\nabla = (\nabla_x + \frac{1}{\varepsilon} \nabla_y) \quad (4)$$

The left hand side of (1) becomes:

$$\partial_t u_\varepsilon = \partial_t (u_0(x, y) + \varepsilon u_1(x, y) + \varepsilon^2 u_2(x, y)) \quad (5)$$

The right hand side of (1) becomes:

$$\begin{aligned} \nabla \cdot (D_V \nabla u_\varepsilon) &= \nabla \cdot (D_V \nabla (u_0 + \varepsilon u_1 + \varepsilon^2 u_2)) \\ &= (\nabla_x + \frac{1}{\varepsilon} \nabla_y) \cdot (D_V (\nabla_x + \frac{1}{\varepsilon} \nabla_y) (u_0 + \varepsilon u_1 + \varepsilon^2 u_2)) \\ &= D_V (\varepsilon^{-2} (\nabla_y \cdot \nabla_y u_0) \\ &\quad + \varepsilon^{-1} (\nabla_x \cdot \nabla_y u_0 + \nabla_y \cdot \nabla_x u_0 + \nabla_y \cdot \nabla_y u_1) \\ &\quad + \varepsilon^0 (\nabla_x \cdot \nabla_x u_0 + \nabla_x \cdot \nabla_y u_1 + \nabla_y \cdot \nabla_x u_1 + \nabla_y \cdot \nabla_y u_2) \\ &\quad + \varepsilon^1 (\nabla_x \cdot \nabla_x u_1 + \nabla_x \cdot \nabla_y u_2 + \nabla_y \cdot \nabla_x u_2) \\ &\quad + \varepsilon^2 (\nabla_x \cdot \nabla_x u_2)) \end{aligned}$$

Comparing the orders of ε we get for ε^{-2} :

$$\nabla_y \cdot \nabla_y u_0(x, y) = 0 \quad (6)$$

Because of the periodicity in y it is $u_0(x, y) = u_0(x)$.

Therefore we get from the ε^{-1} :

$$\nabla_x \cdot \nabla_y u_0(x) = \nabla_y \cdot \nabla_x u_0(x) = 0 \quad (7)$$

It follows:

$$\nabla_y \cdot \nabla_y u_1(x, y) = 0 \quad (8)$$

We write u_1 as linear combination of u_0

$$u_1(x, y) = \sum_{j=1}^n w_j(y) \partial_{x_j} u_0(x) \quad (9)$$

Kann man das wirklich machen, wenn $D_V = \text{const}$?

From the ε^0 term we get:

$$\partial_t u_0 = D_V(\nabla_x \cdot \nabla_x u_0 + \nabla_x \cdot \nabla_y u_1 + \nabla_y \cdot \nabla_x u_1 + \nabla_y \cdot \nabla_y u_2)$$

We integrate over the cell:

$$\begin{aligned} \int_{Y_S} \partial_t u_0 \, dy &= D_V \left(\int_{Y_S} \nabla_x \cdot \nabla_x u_0 \, dy \right. \\ &\quad + \int_{Y_S} \nabla_x \cdot \nabla_y u_1 \, dy \\ &\quad \left. + \int_{Y_S} \nabla_y \cdot (\nabla_x u_1 + \nabla_y u_2) \, dy \right) \end{aligned}$$

Consider the last integral and integrate by parts:

$$D_V \int_{Y_S} \nabla_y \cdot (\nabla_x u_1 + \nabla_y u_2) \, dy = D_V \int_{\Gamma} (\nabla_x u_1 + \nabla_y u_2) \cdot \nu \, d\Gamma(y)$$

Applying the boundary condition we get:

$$D_V \int_{Y_S} \nabla_y \cdot (\nabla_x u_1 + \nabla_y u_2) \, dy = k \int_{\Gamma} u_1 \cdot \nu \, d\Gamma(y)$$

The first two integrals together with the representation of u_1 as linear combination in (9) become:

$$\begin{aligned} &D_V \left(\int_{Y_S} \nabla_x \cdot \nabla_x u_0 \, dy + \int_{Y_S} \nabla_x \cdot \nabla_y u_1 \, dy \right) \\ &= D_V \left(\nabla_x \cdot \nabla_x u_0 \int_{Y_S} dy + \int_{Y_S} \nabla_x \cdot \nabla_y \sum_{j=1}^n w_j(y) \partial_{x_j} u_0(x) \, dy \right) \end{aligned}$$

From the boundary condition we get:

$$\begin{aligned} k(u_0 + \varepsilon u_1 + \varepsilon^2 u_2 - \bar{u}) &= D_V(\varepsilon^{-1} \nabla_y u_0 \\ &\quad + \varepsilon^0 (\nabla_x u_0 + \nabla_y u_1) \\ &\quad + \varepsilon (\nabla_x u_1 + \nabla_y u_2) \\ &\quad + \varepsilon^2 \nabla_x u_2) \cdot \nu \end{aligned}$$

$$0 = D_V \varepsilon^{-1} \nabla_y u_0(x, y) \cdot \nu \tag{10}$$

$$\Rightarrow u_0(x, y) = u_0(x) \text{ or } \nabla_y u_0(x, y) \perp \nu \tag{11}$$