

Canonical Energy-Momentum Tensor of General Relativity

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Abstract

In this article, we provide the natural derivation of symmetrical, coordinate independent canonical energy-momentum tensor for gravitational field.

1 Introduction

The stress-energy-momentum pseudotensor is traditional method to describe the concept of energy-momentum of gravitational field. The Landau-Lifshitz pseudotensor $t_{LL}^{\mu\nu}$ is derived from the Einstein field equation $\mathbb{G}^{\mu\nu} = \kappa T^{\mu\nu}$ to satisfied the conservation law $(T^{\mu\nu} + t_{LL}^{\mu\nu})_{,\nu} = 0$, where $\mathbb{G}^{\mu\nu}$ is Einstein field tensor, $\kappa = \frac{8\pi G}{c^4}$, and $T^{\mu\nu}$ is the energy-momentum tensor of source.

$$t_{LL}^{\mu\nu} = -\frac{1}{2\kappa}\mathbb{G}^{\mu\nu} + \frac{1}{2\kappa(-g)}[(-g)(g^{\mu\nu}g^{\alpha\beta} - g^{\mu\alpha}g^{\nu\beta})]_{,\alpha\beta}$$

$t_{LL}^{\mu\nu}$ is symmetrical but depend explicit in the Christoffel symbols (i.e., coordinate dependent and vanish in specific coordinate). The Dirac pseudotensor $t_D^{\mu\nu}$ starting from the equivalent action \mathcal{L}^* and derived using standard Noether's theorem derivation. The original Einstein-Hilbert action $\mathcal{L} = R\sqrt{-g} = \mathcal{L}[g^{\mu\nu}, g^{\mu\nu}_{,\gamma}, g^{\mu\nu}_{,\gamma\eta}]$ depend on the second derivative of metric tensor, where R and g are Ricci scalar and $g = \det(g_{\mu\nu})$, respectively. The Ostrogradsky instability indicate the Lagrangian should not depend on higher order derivative more than 1st order. The equivalent action

$$\begin{aligned}\mathcal{L}^* &= \mathcal{L} - \partial_\mu(\sqrt{-g}g^{\mu\nu}\Gamma_{\nu\sigma}^\sigma - \sqrt{-g}g^{\sigma\nu}\Gamma_{\nu\sigma}^\mu) \\ &= \sqrt{-g}g^{\mu\nu}(\Gamma_{\mu\nu}^\tau\Gamma_{\tau\sigma}^\sigma - \Gamma_{\mu\sigma}^\tau\Gamma_{\tau\nu}^\sigma) \\ &= \mathcal{L}^*[g^{\mu\nu}, g^{\mu\nu}_{,\gamma}]\end{aligned}$$

have same EoM with advantage only depend on the 1st order derivative of metric tensor, which can apply standard Noether's theorem derivation. However, the equivalent action \mathcal{L}^* lost the scalar property. Also, the Dirac pseudotensor $t_D^{\mu\nu}$:

$$t_D^{\mu\nu} = \frac{1}{2\kappa(-g)}\left[g^{\mu\gamma}(g^{\alpha\beta}\sqrt{-g})_{,\gamma}(\Gamma_{\alpha\beta}^\nu - \delta_\beta^\nu\Gamma_{\alpha\sigma}^\sigma) - g^{\mu\nu}g^{\alpha\beta}(\Gamma_{\alpha\beta}^\rho\Gamma_{\rho\sigma}^\sigma - \Gamma_{\alpha\sigma}^\rho\Gamma_{\beta\rho}^\sigma)\right]$$

lost the symmetric property, and is coordinate dependent and vanishes in specific coordinate as $t_{LL}^{\mu\nu}$.

As we previously derive the symmetrical, gauge invariant canonical energy-momentum tensor for abelian and non-abelian gauge field, in this article we generalize to general relativity. To reserve the scalar property of the Lagrangian and avoid the higher-order derivatives, we will use Palatini variation for derivation the EoM. To further apply Noether's theorem, we will use the vielbeins technique for deriving the conservation law. A more detailed derivation can be found in the supplementary derivation.

2 Derivation of Palatini variation

The Palatini variation treat metric tensor $g^{\omega\sigma}$ and the connection $\Gamma_{\kappa\gamma}^\varepsilon$ as independent field. The curvature tensor $R_{\kappa\omega\sigma}^\varepsilon$:

$$R_{\kappa\omega\sigma}^\varepsilon = \Gamma_{\kappa\sigma,\omega}^\varepsilon - \Gamma_{\kappa\omega,\sigma}^\varepsilon + \Gamma_{\gamma\omega}^\varepsilon \Gamma_{\kappa\sigma}^\gamma - \Gamma_{\gamma\sigma}^\varepsilon \Gamma_{\kappa\omega}^\gamma \quad (1)$$

and the torsion tensor $T_{\beta\gamma}^\alpha$:

$$T_{\beta\gamma}^\alpha = \Gamma_{\beta\gamma}^\alpha - \Gamma_{\gamma\beta}^\alpha = -T_{\gamma\beta}^\alpha \quad (2)$$

In the following, we will NOT assume the connection to be torsion-free, i.e., $T_{\beta\gamma}^\alpha \neq 0 \leftrightarrow \Gamma_{\beta\gamma}^\alpha \neq \Gamma_{\gamma\beta}^\alpha$ in general. The Einstein-Hilbert action:

$$\mathcal{L} = \frac{1}{2\kappa} \sqrt{-g} g^{\mu\nu} \delta_\varepsilon^\omega R_{\kappa\omega\sigma}^\varepsilon = \mathcal{L}[g^{\mu\nu}, \Gamma_{\mu\nu}^\kappa, \Gamma_{\mu\nu,\gamma}^\kappa]$$

$$S = \frac{1}{2\kappa} \int d^4x \sqrt{-g} g^{\kappa\sigma} \delta_\varepsilon^\omega R_{\kappa\omega\sigma}^\varepsilon$$

Here we note that the Ricci tensor $R_{\kappa\sigma}$ and Ricci scalar R are:

$$R_{\kappa\sigma} = \delta_\varepsilon^\omega R_{\kappa\omega\sigma}^\varepsilon$$

and

$$R = g^{\kappa\sigma} R_{\kappa\sigma}$$

, so $\mathcal{L} = \frac{1}{2\kappa} \sqrt{-g} R$. Since the metric tensor $g^{\mu\nu}$ and the connection $\Gamma_{\mu\nu}^\alpha$ are independent field, the Lagrangian only depend on 1st order derivative of the connection. The Palatini variation:

$$\begin{aligned} \delta S &= \int \left(\frac{\partial \mathcal{L}}{\partial g^{\mu\nu}} \delta g^{\mu\nu} + \frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu}^\alpha} \delta \Gamma_{\mu\nu}^\alpha + \frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu,\gamma}^\alpha} \delta \Gamma_{\mu\nu,\gamma}^\alpha \right) d^4x \\ &= \int \underbrace{\frac{\partial \mathcal{L}}{\partial g^{\mu\nu}} \delta g^{\mu\nu} d^4x}_{\text{EoM\#1}} + \int \underbrace{\left[\frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu}^\alpha} - \left(\frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu,\gamma}^\alpha} \right)_{,\gamma} \right] \delta \Gamma_{\mu\nu}^\alpha d^4x}_{\text{EoM\#2}} + \int \left[\frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu,\gamma}^\alpha} \delta \Gamma_{\mu\nu}^\alpha \right]_{,\gamma} d^4x \end{aligned} \quad (3)$$

Two equations of motion (EoM) are

$$\begin{aligned} \frac{\partial \mathcal{L}}{\partial g^{\mu\nu}} &= 0 \\ \frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu}^\alpha} &= \left(\frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu,\gamma}^\alpha} \right)_{,\gamma} \end{aligned}$$

3 EoM#1 - Einstein Field Equation

Since the curvature tensor do not depend on metric tensor, the first EoM:

$$\frac{\partial \mathcal{L}}{\partial g^{\mu\nu}} = \frac{1}{2\kappa} \sqrt{-g} \left(R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R \right) = 0 \quad (4)$$

Define the Einstein field tensor $G_{\mu\nu}$:

$$\mathbb{G}_{\mu\nu} = R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R$$

For vacuum, $\mathbb{G}_{\mu\nu} = 0$ is the Einstein field equation in vacuum.

If the matter presents, $\mathcal{L} = \mathcal{L}_G + \mathcal{L}_M$, where $\mathcal{L}_G = \frac{1}{2\kappa} \sqrt{-g} R$ and \mathcal{L}_M is matters. The variation gives:

$$\begin{aligned} \frac{\partial (\mathcal{L}_G + \mathcal{L}_M)}{\partial g^{\mu\nu}} &= \frac{\partial \mathcal{L}_G}{\partial g^{\mu\nu}} + \frac{\partial \mathcal{L}_M}{\partial g^{\mu\nu}} = \frac{1}{2\kappa} \sqrt{-g} \mathbb{G}_{\mu\nu} + \frac{\partial \mathcal{L}_M}{\partial g^{\mu\nu}} = 0 \\ \rightarrow \mathbb{G}_{\mu\nu} &= \kappa \left(\frac{-2}{\sqrt{-g}} \frac{\partial \mathcal{L}_M}{\partial g^{\mu\nu}} \right) \equiv \kappa T_{\mu\nu} \end{aligned}$$

4 EoM#2 Relate to The Metric Compatible Condition

To show the second EoM relates to the metric compatible condition, we will use the following relations:

$$\frac{1}{\sqrt{-g}}(\sqrt{-g})_{,\gamma} = \Gamma_{\eta\gamma}^{\eta} \quad (5)$$

$$g^{\phi\psi}_{,\gamma} = \frac{1}{\sqrt{-g}}(\sqrt{-g}g^{\phi\psi})_{,\gamma} - g^{\phi\psi}\Gamma_{\eta\gamma}^{\eta} \quad (6)$$

We first calculate $\frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu}^{\alpha}}$ and $\frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu,\gamma}^{\alpha}}$ for later use:

$$\frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu}^{\alpha}} = \sqrt{-g}(g^{\kappa\sigma}\delta_{\alpha}^{\nu}\Gamma_{\kappa\sigma}^{\mu} + g^{\mu\nu}\Gamma_{\alpha\omega}^{\omega} - g^{\kappa\nu}\Gamma_{\kappa\alpha}^{\mu} - g^{\mu\sigma}\Gamma_{\alpha\sigma}^{\nu}) \quad (7)$$

$$\frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu,\gamma}^{\alpha}} = \sqrt{-g}(g^{\mu\nu}\delta_{\alpha}^{\gamma} - g^{\mu\gamma}\delta_{\alpha}^{\nu}) \quad (8)$$

$$\left(\frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu,\gamma}^{\alpha}}\right)_{,\gamma} = (\sqrt{-g}g^{\mu\nu})_{,\alpha} - (\sqrt{-g}g^{\mu\gamma})_{,\gamma}\delta_{\alpha}^{\nu} \quad (9)$$

Substituting Eq.(7) and Eq.(9) into EoM(more detail derivation, see supplement Eq.(10) below):

$$\frac{(\sqrt{-g}g^{\mu\nu})_{,\alpha} - (\sqrt{-g}g^{\mu\gamma})_{,\gamma}\delta_{\alpha}^{\nu}}{\left(\frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu,\gamma}^{\alpha}}\right)_{,\gamma}} = \frac{\sqrt{-g}g^{\kappa\sigma}\delta_{\alpha}^{\nu}\Gamma_{\kappa\sigma}^{\mu} + \sqrt{-g}g^{\mu\nu}\Gamma_{\alpha\omega}^{\omega} - \sqrt{-g}g^{\kappa\nu}\Gamma_{\kappa\alpha}^{\mu} - \sqrt{-g}g^{\mu\sigma}\Gamma_{\alpha\sigma}^{\nu}}{\frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu}^{\alpha}}} \quad (10)$$

$$g^{\mu\nu}_{,\alpha} + g^{\kappa\nu}\Gamma_{\kappa\alpha}^{\mu} + g^{\mu\kappa}\Gamma_{\kappa\alpha}^{\nu} = [g^{\mu\gamma}_{,\gamma} + g^{\kappa\sigma}\Gamma_{\kappa\sigma}^{\mu} + g^{\mu\eta}\Gamma_{\eta\gamma}^{\gamma} + g^{\mu\gamma}T_{\eta\gamma}^{\eta}]\delta_{\alpha}^{\nu} + g^{\mu\nu}T_{\alpha\gamma}^{\gamma} + g^{\mu\kappa}T_{\kappa\alpha}^{\nu}$$

Define the "covariant derivative"

$$g^{\mu\nu}_{;\alpha} = g^{\mu\nu}_{,\alpha} + g^{\kappa\nu}\Gamma_{\kappa\alpha}^{\mu} + g^{\mu\kappa}\Gamma_{\kappa\alpha}^{\nu}$$

Eq.(10) becomes to

$$g^{\mu\nu}_{;\alpha} = (g^{\mu\gamma}_{,\gamma} + g^{\mu\gamma}T_{\eta\gamma}^{\eta})\delta_{\alpha}^{\nu} + g^{\mu\nu}T_{\alpha\gamma}^{\gamma} + g^{\mu\kappa}T_{\kappa\alpha}^{\nu} \quad (11)$$

After some calculation (more detail derivation, see supplement Eq.(12) below), we arrive

$$g^{\mu\nu}_{;\alpha} = \frac{1}{3}g^{\mu\gamma}T_{\eta\gamma}^{\eta}\delta_{\alpha}^{\nu} + g^{\mu\nu}T_{\alpha\gamma}^{\gamma} + g^{\mu\kappa}T_{\kappa\alpha}^{\nu} \quad (12)$$

Eq.(12) indicate that:

Torsion-free then metric compatible	Metric compatible then torsion-free
If torsion-free: $T_{\beta\gamma}^{\alpha} = 0$, then $g^{\mu\nu}_{;\alpha} = 0$	If Metric compatible: $g^{\mu\nu}_{;\alpha} = 0$, then $T_{\beta\gamma}^{\alpha} = 0$

Table 1: Comparison of Torsion-free and Metric Compatibility Conditions

If torsion-free, then

$$g^{\mu\nu}_{;\alpha} = g^{\mu\nu}_{,\alpha} + g^{\kappa\nu}\Gamma_{\kappa\alpha}^{\mu} + g^{\mu\kappa}\Gamma_{\kappa\alpha}^{\nu} = 0$$

is the metric compatible condition.

5 Vielbeins Formalism

In previous tutorials, the variation of the Lie-algebra value gauge 1-form (gauge connection) \mathbf{B} on principal bundle is :

$$\Delta \mathbf{B} = \delta \mathbf{B} + \hat{\mathcal{L}}_{\delta x} \mathbf{B} \quad (13)$$

In local coordinate $p \rightarrow \{x^\mu\}$, the expressions of \mathbf{B} is

$$\mathbf{B}(p) \rightarrow B_\mu^a(x^\gamma) \hat{T}_a$$

, where \hat{T}_a is the generator of Lie algebra. For a given Lie algebra representation:

$$\hat{T}_a \rightarrow \left(\hat{T}_a \right)_c^b$$

We define the notation:

$$\hat{B}_{c\mu}^b \equiv B_\mu^a \left(\hat{T}_a \right)_c^b$$

. The curvature tensor:

$$\mathbf{G} \equiv d\mathbf{B} + [\mathbf{B} \wedge \mathbf{B}] \rightarrow \hat{G}_{b\mu\nu}^a = \hat{B}_{b\nu,\mu}^a - \hat{B}_{b\mu,\nu}^a + \hat{B}_{c\mu}^a \hat{B}_{b\nu}^c - \hat{B}_{c\nu}^a \hat{B}_{b\mu}^c \quad (14)$$

Eq.(13) in local coordinate with Lie algebra representation is:

$$\Delta \hat{B}_{c\mu}^b = \delta \hat{B}_{c\mu}^b + \underbrace{\delta x^\nu \partial_\nu \hat{B}_{c\mu}^b + \hat{B}_{c\mu}^b \partial_\mu \delta x^\nu}_{\hat{\mathcal{L}}_{\delta x} \mathbf{B}} \quad (15)$$

However, the connection on tangent bundle is more subtle. The curvature on tangent bundle is Eq.(1):

$$R_{\kappa\omega\sigma}^\epsilon = \Gamma_{\kappa\sigma,\omega}^\epsilon - \Gamma_{\kappa\omega,\sigma}^\epsilon + \Gamma_{\gamma\omega}^\epsilon \Gamma_{\kappa\sigma}^\gamma - \Gamma_{\gamma\sigma}^\epsilon \Gamma_{\kappa\omega}^\gamma \quad (1)$$

The principal curvature Eq.(14) is Lie-algebra value 2-form. In contrast, the tangent bundle curvature takes value in tangent vector. If we directly apply Lie derivative on tangent connection:

$$\left(\hat{\mathcal{L}}_{\delta x} \Gamma \right)_{\mu\nu}^\alpha = \delta x^\epsilon \Gamma_{\mu\nu,\epsilon}^\alpha + \Gamma_{\mu\epsilon}^\alpha \delta x_{,\nu}^\epsilon + \delta x_{,\mu\nu}^\alpha - \Gamma_{\mu\nu}^\epsilon \delta x_{,\epsilon}^\alpha + \Gamma_{\epsilon\nu}^\alpha \delta x_{,\mu}^\epsilon$$

due to tangent vector value

We can use Vielbeins formalism. In the vielbeins formalism, the tangent connection $\Gamma_{\kappa\sigma}^\epsilon$ becomes the spin connection $\omega_{b\sigma}^a$ in the frame bundle, which is \mathfrak{gl} -value 1-form. From the tangent formalism $\{\partial_\mu\}$ to the vielbeins formalism $\{\hat{e}_a\}$, define the transformation e_a^μ such that (more detail can be found in Supplementary: Vielbeins):

$$\begin{aligned} \partial_\mu &= e_\mu^a \hat{e}_a \\ g_{\mu\nu} &= e_\mu^a e_\nu^b \eta_{ab} \\ e_a^\mu e_\nu^\mu &= \delta_\nu^\mu \text{ and } e_a^\mu e_\mu^b = \delta_b^a \end{aligned}$$

The spin connection:

$$\omega_{b\sigma}^a = e_\epsilon^a e_b^\kappa \Gamma_{\kappa\sigma}^\epsilon + e_\mu^a e_{b,\sigma}^\mu$$

The frame bundle curvature $\mathcal{R}_{b\omega\sigma}^a$ is

$$\begin{aligned} \mathcal{R}_{b\omega\sigma}^a &= \omega_{b\sigma,\omega}^a - \omega_{b\omega,\sigma}^a + \omega_{c\omega}^a \omega_{b\sigma}^c - \omega_{c\sigma}^a \omega_{b\omega}^c \\ &= e_\epsilon^a e_b^\kappa R_{\kappa\omega\sigma}^\epsilon \end{aligned}$$

With this preparation, the Hilbert-Einstein action becomes as follows:

$$S = \frac{1}{2\kappa} \int d^4x \sqrt{-g} g^{\kappa\sigma} \delta_\varepsilon^\omega R_{\kappa\omega\sigma}^\varepsilon = \frac{1}{2\kappa} \int d^4x \eta^{ae} e_e^\sigma e_c^\omega \mathcal{R}_{a\omega\sigma}^c \quad (16)$$

, where $e = \det(e_\mu^a) = \sqrt{-g}$.

The variation is similar to Eq.(3):

$$\begin{aligned} \delta S &= \int \left(\frac{\partial \mathcal{L}}{\partial e_a^\mu} \delta e_a^\mu + \frac{\partial \mathcal{L}}{\partial \omega_{c\mu}^b} \delta \omega_{c\mu}^b + \frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \delta \omega_{c\mu,\gamma}^b \right) d^4x \\ &= \int \frac{\partial \mathcal{L}}{\partial e_a^\mu} \delta e_a^\mu d^4x + \int \underbrace{\left[\frac{\partial \mathcal{L}}{\partial \omega_{c\mu}^b} - \left(\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \right)_{,\gamma} \right]}_{EoM\#4} \delta \omega_{c\mu}^b d^4x + \int \left[\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \delta \omega_{c\mu}^b \right]_{,\gamma} d^4x \end{aligned} \quad (17)$$

Since $g_{\mu\nu} = \eta_{ab} e_\mu^a e_\nu^b$, the variation $\delta g_{\mu\nu} = 2\eta_{ab} e_\mu^a \delta e_\nu^b$, the *EoM*#3 is similar to *EoM*#1. *EoM*#4 is also similar to *EoM*#2. Now, we apply the Noether variation:

$$\Delta S = \int [EoM] d^4x + \int \left[\partial_\gamma \left(\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \delta \omega_{c\mu}^b \right) + (\mathcal{L} \delta x^\gamma)_{,\gamma} \right] d^4x$$

The spin connection $\omega_{c\mu}^b$ is \mathfrak{gl} -value 1-form, we apply similar variation as Eq.(15):

$$\Delta \omega_{c\mu}^b = \delta \omega_{c\mu}^b + \underbrace{\delta x^\varepsilon \partial_\varepsilon \omega_{c\mu}^b + \omega_{c\varepsilon}^b \partial_\mu \delta x^\varepsilon}_{\hat{\mathcal{L}}_{\delta x} \omega} \quad (18)$$

We have:

$$\begin{aligned} \Delta S &= \int [EoM] d^4x + \int \left[\partial_\gamma \left(\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} (\Delta \omega_{c\mu}^b - \delta x^\varepsilon \partial_\varepsilon \omega_{c\mu}^b - \omega_{c\varepsilon}^b \partial_\mu \delta x^\varepsilon) \right) + (\mathcal{L} \delta x^\gamma)_{,\gamma} \right] d^4x \\ &= \int [EoM] d^4x + \int \frac{\partial_\gamma}{(*)} \left[\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \Delta \omega_{c\mu}^b - \left(\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \omega_{c\mu,\varepsilon}^b \delta x^\varepsilon + \frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \omega_{c\varepsilon}^b \delta x_{,\mu}^\varepsilon - \delta_\varepsilon^\gamma \mathcal{L} \delta x^\varepsilon \right) \right] d^4x \end{aligned}$$

Evaluate the $(*)$ term:

$$\underbrace{\partial_\gamma \left[\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \omega_{c\varepsilon}^b \delta x_{,\mu}^\varepsilon \right]}_{(*)} = \underbrace{\partial_\gamma \left[\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \omega_{c\varepsilon}^b \delta x^\varepsilon \right]}_{(*1)}_{,\mu} - \underbrace{\partial_\gamma \left[\left(\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \right)_{,\mu} \omega_{c\varepsilon}^b \delta x^\varepsilon \right]}_{(*2)} - \underbrace{\partial_\gamma \left[\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \omega_{c\varepsilon,\mu}^b \delta x^\varepsilon \right]}_{(*3)}$$

We first calculate $\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b}$ for later use:

$$\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} = \frac{1}{2\kappa} e \eta^{ce} (e_e^\mu e_b^\gamma - e_e^\gamma e_b^\mu) \left(= -\frac{\partial \mathcal{L}}{\partial \omega_{c\gamma,\mu}^b} \right) \quad (19)$$

The $(*1)$ term:

$$\underbrace{\partial_\gamma \left[\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \omega_{c\varepsilon}^b \delta x^\varepsilon \right]}_{(*1)}_{,\mu} = \left[\frac{1}{2\kappa} e \eta^{ce} (e_e^\mu e_b^\gamma - e_e^\gamma e_b^\mu) \omega_{c\varepsilon}^b \delta x^\varepsilon \right]_{(19)}_{,\mu\gamma} = 0$$

The $(*2)$ term rely on *EoM*#4, using

$$\underbrace{-\partial_\gamma \left[\left(\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \right)_{,\mu} \omega_{c\varepsilon}^b \delta x^\varepsilon \right]}_{(*2)} \underbrace{=}_{(19)} -\partial_\gamma \left[\left(-\frac{\partial \mathcal{L}}{\partial \omega_{c\gamma,\mu}^b} \right)_{,\mu} \omega_{c\varepsilon}^b \delta x^\varepsilon \right] \underbrace{=}_{EoM\#4} \partial_\gamma \left(\frac{\partial \mathcal{L}}{\partial \omega_{c\gamma}^b} \omega_{c\varepsilon}^b \delta x^\varepsilon \right)$$

and

$$\frac{\partial \mathcal{L}}{\partial \omega_{c\gamma}^b} = \frac{1}{2\kappa} e \eta^{ae} (e_e^\gamma e_b^\mu - e_e^\mu e_b^\gamma) \omega_{a\mu}^c + \frac{1}{2\kappa} e \eta^{ce} (e_e^\mu e_f^\gamma - e_e^\gamma e_f^\mu) \omega_{b\mu}^f \quad (20)$$

We can calculate

$$\frac{\partial \mathcal{L}}{\partial \omega_{c\gamma}^b} \omega_{c\varepsilon}^b = \frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} (\omega_{d\varepsilon}^b \omega_{c\mu}^d - \omega_{d\mu}^b \omega_{c\varepsilon}^d) \quad (21)$$

The $(*)$ term than become:

$$\frac{\partial_\gamma \left[\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \omega_{c\varepsilon}^b \delta x_{,\mu}^\varepsilon \right]}{(*)} = \frac{0}{(*)1} - \frac{\partial_\gamma \left[\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \omega_{c\varepsilon,\mu}^b \delta x^\varepsilon \right]}{(*)3} + \frac{\partial_\gamma \left[\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} (\omega_{d\varepsilon}^b \omega_{c\mu}^d - \omega_{d\mu}^b \omega_{c\varepsilon}^d) \delta x^\varepsilon \right]}{(*)2}$$

We have:

$$\begin{aligned} \Delta S &= \int [EoM] d^4x + \int \partial_\gamma \left[\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \Delta \omega_{c\mu}^b - \left(\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \omega_{c\mu,\varepsilon}^b \delta x^\varepsilon - \frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \omega_{c\varepsilon,\mu}^b \delta x^\varepsilon + \frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} (\omega_{d\varepsilon}^b \omega_{c\mu}^d - \omega_{d\mu}^b \omega_{c\varepsilon}^d) \delta x^\varepsilon - \delta_\varepsilon^\gamma \mathcal{L} \delta x^\varepsilon \right) \right] d^4x \\ &= \int [EoM] d^4x + \int \partial_\gamma \left[\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \Delta \omega_{c\mu}^b - \left(\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} (\omega_{c\mu,\varepsilon}^b - \omega_{c\varepsilon,\mu}^b + \omega_{d\varepsilon}^b \omega_{c\mu}^d - \omega_{d\mu}^b \omega_{c\varepsilon}^d) \delta x^\varepsilon - \delta_\varepsilon^\gamma \mathcal{L} \delta x^\varepsilon \right) \right] d^4x \\ &= \int [EoM] d^4x + \int \partial_\gamma \left[\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \Delta \omega_{c\mu}^b - \left(\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \mathcal{R}_{c\varepsilon\mu}^b - \delta_\varepsilon^\gamma \mathcal{L} \right) \delta x^\varepsilon \right] d^4x \end{aligned}$$

The canonical energy-momentum tensor:

$$\begin{aligned} t_\varepsilon^\gamma &= \frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} \mathcal{R}_{c\varepsilon\mu}^b - \delta_\varepsilon^\gamma \mathcal{L} \\ &= \frac{1}{2\kappa} e \eta^{ce} (e_e^\gamma e_b^\mu - e_e^\mu e_b^\gamma) \mathcal{R}_{c\varepsilon\mu}^b - \delta_\varepsilon^\gamma \mathcal{L} \\ &= \frac{1}{2\kappa} (g^{\beta\mu} R_{\beta\varepsilon\mu}^\gamma + g^{\beta\gamma} R_{\beta\varepsilon} - \delta_\varepsilon^\gamma R) \sqrt{-g} \end{aligned} \quad (22)$$

The canonical energy-momentum 2-form:

$$t_{\alpha\varepsilon} = g_{\alpha\gamma} t_\varepsilon^\gamma = \frac{1}{2\kappa} (g_{\alpha\gamma} g^{\beta\mu} R_{\beta\varepsilon\mu}^\gamma + R_{\alpha\varepsilon} - g_{\alpha\varepsilon} R) \sqrt{-g}$$

If metric compatible, we have:

$$t_{\alpha\varepsilon} = \frac{1}{\kappa} \mathbb{G}_{\alpha\varepsilon} \sqrt{-g} \quad (23)$$

This canonical energy-momentum 2-form has following properties:

1. symmetric
2. coordinate independent
3. vanish when vacuum, which does not depend on the choice of the connection.

References

1. Canonical Energy-Momentum Tensor of Abelian Fields, arXiv:2503.15031.

Supplementary Derivation

Eq.(4)

$$\begin{aligned}
\frac{\partial \mathcal{L}}{\partial g^{\mu\nu}} &= \frac{1}{2\kappa} \frac{\partial}{\partial g^{\mu\nu}} (\sqrt{-g} g^{\kappa\sigma} \delta_\epsilon^\omega R_{\kappa\omega\sigma}^\epsilon) = \frac{1}{2\kappa} \sqrt{-g} \left(\delta_\epsilon^\omega \delta_\mu^\kappa \delta_\nu^\sigma R_{\kappa\omega\sigma}^\epsilon - \frac{1}{2} g_{\mu\nu} g^{\kappa\sigma} \delta_\epsilon^\omega R_{\kappa\omega\sigma}^\epsilon \right) \\
&= \frac{1}{2\kappa} \sqrt{-g} \left(\delta_\epsilon^\omega \delta_\mu^\kappa \delta_\nu^\sigma R_{\kappa\omega\sigma}^\epsilon - \frac{1}{2} g_{\mu\nu} g^{\kappa\sigma} R_{\kappa\sigma} \right) \\
&= \frac{1}{2\kappa} \sqrt{-g} \left(R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R \right) \\
&= \frac{1}{2\kappa} \sqrt{-g} \left(R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R \right)
\end{aligned}$$

Eq.(5)

$$\begin{aligned}
(\sqrt{-g})_{,\gamma} &= -\frac{1}{2} \sqrt{-g} g_{\phi\psi} g^{\phi\psi}_{,\gamma} \\
&= -\frac{1}{2} \sqrt{-g} g_{\phi\psi} (-g^{\eta\psi} \Gamma_{\eta\gamma}^\phi - g^{\phi\eta} \Gamma_{\eta\gamma}^\psi) \\
&= \frac{1}{2} \sqrt{-g} (g_{\phi\psi} g^{\eta\psi} \Gamma_{\eta\gamma}^\phi + g_{\phi\psi} g^{\phi\eta} \Gamma_{\eta\gamma}^\psi) \\
&= \frac{1}{2} \sqrt{-g} (\delta_\phi^\eta \Gamma_{\eta\gamma}^\phi + \delta_\psi^\eta \Gamma_{\eta\gamma}^\psi) \\
&= \frac{1}{2} \sqrt{-g} (\Gamma_{\eta\gamma}^\eta + \Gamma_{\eta\gamma}^\eta) \\
&= \sqrt{-g} \Gamma_{\eta\gamma}^\eta
\end{aligned}$$

Eq.(6)

$$g^{\phi\psi}_{,\gamma} = \frac{\sqrt{-g}}{\sqrt{-g}} g^{\phi\psi}_{,\gamma} = \frac{1}{\sqrt{-g}} (\sqrt{-g} g^{\phi\psi})_{,\gamma} - \underbrace{\frac{1}{\sqrt{-g}} (\sqrt{-g})_{,\gamma}}_{(5)=\Gamma_{\eta\gamma}^\eta} g^{\phi\psi} = \frac{1}{\sqrt{-g}} (\sqrt{-g} g^{\phi\psi})_{,\gamma} - g^{\phi\psi} \Gamma_{\eta\gamma}^\eta$$

Eq.(7)

Evaluate $\frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu}^\alpha}$:

$$\begin{aligned}
\frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu}^\alpha} &= \frac{\partial}{\partial \Gamma_{\mu\nu}^\alpha} \left(\frac{1}{2\kappa} \sqrt{-g} g^{\kappa\sigma} \delta_\epsilon^\omega R_{\kappa\omega\sigma}^\epsilon \right) \\
&= \frac{1}{2\kappa} \sqrt{-g} g^{\kappa\sigma} \delta_\epsilon^\omega \frac{\partial}{\partial \Gamma_{\mu\nu}^\alpha} (\Gamma_{\kappa\sigma,\omega}^\epsilon - \Gamma_{\kappa\omega,\sigma}^\epsilon + \Gamma_{\eta\omega}^\epsilon \Gamma_{\kappa\sigma}^\eta - \Gamma_{\eta\sigma}^\epsilon \Gamma_{\kappa\omega}^\eta) \\
&= \frac{1}{2\kappa} \sqrt{-g} g^{\kappa\sigma} \delta_\epsilon^\omega (\delta_\alpha^\eta \delta_\mu^\epsilon \delta_\nu^\omega \Gamma_{\kappa\sigma}^\eta + \Gamma_{\eta\omega}^\epsilon \delta_\alpha^\eta \delta_\kappa^\mu \delta_\sigma^\nu - \delta_\alpha^\eta \delta_\mu^\epsilon \delta_\nu^\omega \Gamma_{\kappa\omega}^\eta - \Gamma_{\eta\sigma}^\epsilon \delta_\alpha^\eta \delta_\kappa^\mu \delta_\omega^\nu) \\
&= \frac{1}{2\kappa} (\sqrt{-g} g^{\kappa\sigma} \delta_\alpha^\nu \Gamma_{\kappa\sigma}^\mu + \sqrt{-g} g^{\mu\nu} \Gamma_{\alpha\omega}^\omega - \sqrt{-g} g^{\kappa\nu} \Gamma_{\kappa\alpha}^\mu - \sqrt{-g} g^{\mu\sigma} \Gamma_{\alpha\sigma}^\nu)
\end{aligned}$$

Eq.(8)

Evaluate $\frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu,\gamma}^\alpha}$:

$$\begin{aligned}
\frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu,\gamma}^\alpha} &= \frac{\partial}{\partial \Gamma_{\mu\nu,\gamma}^\alpha} \left(\frac{1}{2\kappa} \sqrt{-g} g^{\kappa\sigma} \delta_\epsilon^\omega R_{\kappa\omega\sigma}^\epsilon \right) \\
&= \frac{1}{2\kappa} \sqrt{-g} g^{\kappa\sigma} \delta_\epsilon^\omega \frac{\partial}{\partial \Gamma_{\mu\nu,\gamma}^\alpha} (\Gamma_{\kappa\sigma,\omega}^\epsilon - \Gamma_{\kappa\omega,\sigma}^\epsilon + \Gamma_{\eta\omega}^\epsilon \Gamma_{\kappa\sigma}^\eta - \Gamma_{\eta\sigma}^\epsilon \Gamma_{\kappa\omega}^\eta) \\
&= \frac{1}{2\kappa} \sqrt{-g} g^{\kappa\sigma} \delta_\epsilon^\omega (\delta_\alpha^\epsilon \delta_\kappa^\mu \delta_\sigma^\nu \delta_\omega^\gamma - \delta_\alpha^\epsilon \delta_\kappa^\mu \delta_\omega^\nu \delta_\sigma^\gamma) \\
&= \frac{1}{2\kappa} \sqrt{-g} (g^{\mu\nu} \delta_\alpha^\gamma - g^{\mu\gamma} \delta_\alpha^\nu)
\end{aligned}$$

Eq.(9)

Evaluate $\left(\frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu,\gamma}^\alpha} \right)_{,\gamma}$:

$$\begin{aligned}
\left(\frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu,\gamma}^\alpha} \right)_{,\gamma} &= \frac{1}{2\kappa} [\sqrt{-g} (g^{\mu\nu} \delta_\alpha^\gamma - g^{\mu\gamma} \delta_\alpha^\nu)]_{,\gamma} \\
&= \frac{1}{2\kappa} (\sqrt{-g} g^{\mu\nu})_{,\gamma} \delta_\alpha^\gamma - (\sqrt{-g} g^{\mu\gamma})_{,\gamma} \delta_\alpha^\nu \\
&= \frac{1}{2\kappa} (\sqrt{-g} g^{\mu\nu})_{,\alpha} - \frac{1}{2\kappa} (\sqrt{-g} g^{\mu\gamma})_{,\gamma} \delta_\alpha^\nu
\end{aligned}$$

Eq.(10)

$$\begin{aligned}
\left(\frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu,\gamma}^\alpha} \right)_{,\gamma} &= \frac{\partial \mathcal{L}}{\partial \Gamma_{\mu\nu}^\alpha} \\
\frac{(\sqrt{-g} g^{\mu\nu})_{,\alpha}}{(A)} - \frac{(\sqrt{-g} g^{\mu\gamma})_{,\gamma} \delta_\alpha^\nu}{(B)} &= \frac{\sqrt{-g} g^{\kappa\sigma} \delta_\alpha^\nu \Gamma_{\kappa\sigma}^\mu}{(C)} + \frac{\sqrt{-g} g^{\mu\nu} \Gamma_{\alpha\omega}^\omega}{(D)} - \frac{\sqrt{-g} g^{\kappa\nu} \Gamma_{\kappa\alpha}^\mu}{(E)} - \frac{\sqrt{-g} g^{\mu\sigma} \Gamma_{\alpha\sigma}^\nu}{(F)} \\
\frac{(\sqrt{-g} g^{\mu\nu})_{,\alpha}}{(A)} + \frac{\sqrt{-g} g^{\kappa\nu} \Gamma_{\kappa\alpha}^\mu}{(E)} &= \frac{(\sqrt{-g} g^{\mu\gamma})_{,\gamma} \delta_\alpha^\nu}{(B)} + \frac{\sqrt{-g} g^{\kappa\sigma} \delta_\alpha^\nu \Gamma_{\kappa\sigma}^\mu}{(C)} + \frac{\sqrt{-g} g^{\mu\nu} \Gamma_{\alpha\omega}^\omega}{(D)} - \frac{\sqrt{-g} g^{\mu\sigma} \Gamma_{\alpha\sigma}^\nu}{(F)} \\
\frac{1}{\sqrt{-g}} (\sqrt{-g} g^{\mu\nu})_{,\alpha} + g^{\kappa\nu} \Gamma_{\kappa\alpha}^\mu &= \left[\frac{1}{\sqrt{-g}} (\sqrt{-g} g^{\mu\gamma})_{,\gamma} + g^{\kappa\sigma} \Gamma_{\kappa\sigma}^\mu \right] \delta_\alpha^\nu + g^{\mu\nu} \Gamma_{\alpha\gamma}^\gamma - g^{\mu\kappa} \Gamma_{\alpha\kappa}^\nu \\
&\quad (B)+(C) \\
\underbrace{\frac{1}{\sqrt{-g}} (\sqrt{-g} g^{\mu\nu})_{,\alpha}}_{(6)=g^{\mu\nu},\alpha} - \underbrace{g^{\mu\nu} \Gamma_{\gamma\alpha}^\gamma}_{(6)=g^{\mu\nu},\alpha} + g^{\kappa\nu} \Gamma_{\kappa\alpha}^\mu + \underbrace{g^{\mu\kappa} \Gamma_{\kappa\alpha}^\nu}_{(6)=g^{\mu\kappa} \Gamma_{\kappa\alpha}^\nu} &= \underbrace{\left[\frac{1}{\sqrt{-g}} (\sqrt{-g} g^{\mu\gamma})_{,\gamma} + g^{\eta\gamma} \Gamma_{\eta\gamma}^\mu \right]}_{(\#)} \delta_\alpha^\nu + \underbrace{g^{\mu\nu} \Gamma_{\alpha\gamma}^\gamma - g^{\mu\nu} \Gamma_{\gamma\alpha}^\gamma}_{g^{\mu\nu} T_{\alpha\gamma}^\gamma} + \underbrace{g^{\mu\kappa} \Gamma_{\kappa\alpha}^\nu - g^{\mu\kappa} \Gamma_{\alpha\kappa}^\nu}_{g^{\mu\kappa} T_{\kappa\alpha}^\nu} \quad (24)
\end{aligned}$$

Calculate (#) term:

$$\begin{aligned}
\underbrace{\left[\frac{1}{\sqrt{-g}} (\sqrt{-g} g^{\mu\gamma})_{,\gamma} + g^{\eta\gamma} \Gamma_{\eta\gamma}^\mu \right]}_{(\#)} &= \underbrace{\frac{1}{\sqrt{-g}} (\sqrt{-g} g^{\mu\gamma})_{,\gamma}}_{(6)=g^{\mu\gamma},\gamma} - \underbrace{g^{\mu\gamma} \Gamma_{\eta\gamma}^\eta}_{(6)=g^{\mu\gamma},\gamma} + g^{\eta\gamma} \Gamma_{\eta\gamma}^\mu + \underbrace{g^{\mu\eta} \Gamma_{\eta\gamma}^\gamma + g^{\mu\gamma} \Gamma_{\eta\gamma}^\eta - g^{\mu\eta} \Gamma_{\eta\gamma}^\gamma}_{g^{\mu\gamma} T_{\eta\gamma}^\eta} \\
&= g^{\mu\gamma},\gamma + g^{\eta\gamma} \Gamma_{\eta\gamma}^\mu + g^{\mu\eta} \Gamma_{\eta\gamma}^\gamma + g^{\mu\gamma} T_{\eta\gamma}^\eta
\end{aligned}$$

Eq.(24) become

$$g^{\mu\nu},\alpha + g^{\kappa\nu} \Gamma_{\kappa\alpha}^\mu + g^{\mu\kappa} \Gamma_{\kappa\alpha}^\nu = \underbrace{\left[g^{\mu\gamma},\gamma + g^{\kappa\sigma} \Gamma_{\kappa\sigma}^\mu + g^{\mu\eta} \Gamma_{\eta\gamma}^\gamma + g^{\mu\gamma} T_{\eta\gamma}^\eta \right]}_{(\#)} \delta_\alpha^\nu + g^{\mu\nu} T_{\alpha\gamma}^\gamma + g^{\mu\kappa} T_{\kappa\alpha}^\nu \quad (25)$$

Eq.(12)

Starting from Eq.(11)

$$\begin{aligned}
g^{\mu\nu}{}_{;\alpha} &= [g^{\mu\gamma}{}_{;\gamma} + g^{\mu\gamma} T_{\eta\gamma}^\eta] \delta_\alpha^\nu + g^{\mu\nu} T_{\alpha\gamma}^\gamma + g^{\mu\kappa} T_{\kappa\alpha}^\nu \\
g^{\mu\nu}{}_{;\alpha} - \delta_\alpha^\nu g^{\mu\gamma}{}_{;\gamma} &= \delta_\alpha^\nu g^{\mu\gamma} T_{\eta\gamma}^\eta + g^{\mu\nu} T_{\alpha\gamma}^\gamma + g^{\mu\kappa} T_{\kappa\alpha}^\nu \\
\delta_\nu^\alpha g^{\mu\nu}{}_{;\alpha} - \delta_\nu^\alpha \delta_\alpha^\nu g^{\mu\gamma}{}_{;\gamma} &= \delta_\nu^\alpha \delta_\alpha^\nu g^{\mu\gamma} T_{\eta\gamma}^\eta + \delta_\nu^\alpha g^{\mu\nu} T_{\alpha\gamma}^\gamma + \delta_\nu^\alpha g^{\mu\kappa} T_{\kappa\alpha}^\nu \\
g^{\mu\nu}{}_{;\nu} - 4 g^{\mu\gamma}{}_{;\gamma} &= 4 g^{\mu\gamma} T_{\eta\gamma}^\eta + g^{\mu\nu} T_{\nu\gamma}^\gamma + g^{\mu\kappa} T_{\kappa\nu}^\nu \\
-3 g^{\mu\nu}{}_{;\nu} &= 2 g^{\mu\gamma} T_{\eta\gamma}^\eta \\
g^{\mu\nu}{}_{;\nu} &= -\frac{2}{3} g^{\mu\gamma} T_{\eta\gamma}^\eta
\end{aligned}$$

Substituting back into Eq.(11)

$$\begin{aligned}
g^{\mu\nu}{}_{;\alpha} &= \left(-\frac{2}{3} g^{\mu\gamma} T_{\eta\gamma}^\eta + g^{\mu\gamma} T_{\eta\gamma}^\eta \right) \delta_\alpha^\nu + g^{\mu\nu} T_{\alpha\gamma}^\gamma + g^{\mu\kappa} T_{\kappa\alpha}^\nu \\
&= \frac{1}{3} g^{\mu\gamma} T_{\eta\gamma}^\eta \delta_\alpha^\nu + g^{\mu\nu} T_{\alpha\gamma}^\gamma + g^{\mu\kappa} T_{\kappa\alpha}^\nu
\end{aligned}$$

If torsion-free, i.e., $T_{\beta\gamma}^\alpha = 0$, hence:

$$g^{\mu\nu}{}_{;\alpha} = 0$$

is metric compatible.

On the otherhand, if metric compatible, i.e., $g^{\mu\nu}{}_{;\alpha} = 0$, hence:

$$\begin{aligned}
0 &= \frac{1}{3} g^{\mu\gamma} T_{\eta\gamma}^\eta \delta_\alpha^\nu + g^{\mu\nu} T_{\alpha\gamma}^\gamma + g^{\mu\kappa} T_{\kappa\alpha}^\nu \\
0 &= \delta_\nu^\alpha \frac{1}{3} g^{\mu\gamma} T_{\eta\gamma}^\eta \delta_\alpha^\nu + \delta_\nu^\alpha g^{\mu\nu} T_{\alpha\gamma}^\gamma + \delta_\nu^\alpha g^{\mu\kappa} T_{\kappa\alpha}^\nu \\
0 &= \frac{4}{3} g^{\mu\gamma} T_{\eta\gamma}^\eta + g^{\mu\nu} T_{\nu\gamma}^\gamma + g^{\mu\kappa} T_{\kappa\nu}^\nu \\
0 &= \frac{4}{3} g^{\mu\gamma} T_{\eta\gamma}^\eta - 2 g^{\mu\nu} T_{\gamma\nu}^\gamma \\
0 &= T_{\kappa\nu}^\nu
\end{aligned} \tag{26}$$

Put back into Eq.(26):

$$\begin{aligned}
0 &= 0 + 0 + g^{\mu\kappa} T_{\kappa\alpha}^\nu \\
0 &= T_{\kappa\alpha}^\nu
\end{aligned}$$

imply torsion-free.

Eq.(16)

$$\begin{aligned}
R &= g^{\kappa\sigma} \delta_\epsilon^\omega R_{\kappa\omega\sigma}^\epsilon \\
&= (\eta^{de} e_d^\kappa e_e^\sigma) \delta_\epsilon^\omega (e_c^\epsilon e_\kappa^a \mathcal{R}_{a\omega\sigma}^c) \\
&= \eta^{ae} e_e^\sigma e_c^\omega \mathcal{R}_{a\omega\sigma}^c
\end{aligned}$$

Eq.(19)

$$\begin{aligned}
\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} &= \frac{\partial}{\partial \omega_{c\mu,\gamma}^b} \left(\frac{1}{2\kappa} e \eta^{ae} e_e^\sigma e_f^\omega \mathcal{R}_{a\omega\sigma}^f \right) \\
&= \frac{1}{2\kappa} e \eta^{ae} e_e^\sigma e_f^\omega \frac{\partial}{\partial \omega_{c\mu,\gamma}^b} \left(\omega_{a\sigma,\omega}^f - \omega_{a\omega,\sigma}^f + \omega_{d\omega}^f \omega_{a\sigma}^d - \omega_{d\sigma}^f \omega_{a\omega}^d \right) \\
&= \frac{1}{2\kappa} e \eta^{ae} e_e^\sigma e_f^\omega \left(\delta_b^f \delta_a^c \delta_\sigma^\mu \delta_\omega^\gamma - \delta_b^f \delta_a^c \delta_\omega^\mu \delta_\sigma^\gamma \right) \\
&= \frac{1}{2\kappa} e \eta^{ae} \delta_a^c \left(e_e^\sigma e_f^\omega \delta_b^f \delta_\sigma^\mu \delta_\omega^\gamma - e_e^\sigma e_f^\omega \delta_b^f \delta_\omega^\mu \delta_\sigma^\gamma \right) \\
&= \frac{1}{2\kappa} e \eta^{ce} (e_e^\mu e_b^\gamma - e_e^\gamma e_b^\mu)
\end{aligned}$$

Similarly,

$$\frac{\partial \mathcal{L}}{\partial \omega_{c\gamma,\mu}^b} = \frac{1}{2\kappa} e \eta^{ce} (e_e^\gamma e_b^\mu - e_e^\mu e_b^\gamma) = -\frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b}$$

Eq.(20)

$$\begin{aligned}
\frac{\partial \mathcal{L}}{\partial \omega_{c\gamma}^b} &= \frac{\partial}{\partial \omega_{c\gamma}^b} \left(\frac{1}{2\kappa} e \eta^{ae} e_e^\sigma e_f^\omega \mathcal{R}_{a\omega\sigma}^f \right) \\
&= \frac{1}{2\kappa} e \eta^{ae} e_e^\sigma e_f^\omega \frac{\partial}{\partial \omega_{c\gamma}^b} \left(\omega_{a\sigma,\omega}^f - \omega_{a\omega,\sigma}^f + \omega_{d\omega}^f \omega_{a\sigma}^d - \omega_{d\sigma}^f \omega_{a\omega}^d \right) \\
&= \frac{1}{2\kappa} e \eta^{ae} e_e^\sigma e_f^\omega \left(\delta_b^f \delta_a^c \delta_\omega^\gamma \omega_{a\sigma}^d + \omega_{d\omega}^f \delta_b^d \delta_a^c \delta_\sigma^\gamma - \delta_b^f \delta_a^c \delta_\sigma^\gamma \omega_{a\omega}^d - \omega_{d\sigma}^f \delta_b^d \delta_a^c \delta_\omega^\gamma \right) \\
&= \frac{1}{2\kappa} e \eta^{ae} e_e^\sigma e_f^\omega \left(\delta_b^f \delta_\omega^\gamma \omega_{a\sigma}^c + \omega_{b\omega}^f \delta_a^c \delta_\sigma^\gamma - \delta_b^f \delta_\sigma^\gamma \omega_{a\omega}^c - \omega_{b\sigma}^f \delta_a^c \delta_\omega^\gamma \right) \\
&= \frac{1}{2\kappa} e \left(\eta^{ae} e_e^\sigma e_f^\omega \delta_b^f \delta_\omega^\gamma \omega_{a\sigma}^c + \eta^{ae} e_e^\sigma e_f^\omega \omega_{b\omega}^f \delta_a^c \delta_\sigma^\gamma - \eta^{ae} e_e^\sigma e_f^\omega \delta_b^f \delta_\sigma^\gamma \omega_{a\omega}^c - \eta^{ae} e_e^\sigma e_f^\omega \omega_{b\sigma}^f \delta_a^c \delta_\omega^\gamma \right) \\
&= \frac{1}{2\kappa} e \left(\eta^{ae} e_e^\sigma e_b^\gamma \omega_{a\sigma}^c + \eta^{ce} e_e^\gamma e_f^\omega \omega_{b\omega}^f - \eta^{ae} e_e^\gamma e_b^\omega \omega_{a\omega}^c - \eta^{ce} e_e^\sigma e_f^\gamma \omega_{b\sigma}^f \right) \\
&= \frac{1}{2\kappa} e \eta^{ae} e_e^\mu e_b^\gamma \omega_{a\mu}^c + \frac{1}{2\kappa} e \eta^{ce} e_e^\gamma e_f^\mu \omega_{b\mu}^f - \frac{1}{2\kappa} e \eta^{ae} e_e^\gamma e_b^\mu \omega_{a\mu}^c - \frac{1}{2\kappa} e \eta^{ce} e_e^\mu e_f^\gamma \omega_{b\mu}^f \\
&= \frac{1}{2\kappa} e \eta^{ae} (e_e^\mu e_b^\gamma - e_e^\gamma e_b^\mu) \omega_{a\mu}^c - \frac{1}{2\kappa} e \eta^{ce} (e_e^\mu e_f^\gamma - e_e^\gamma e_f^\mu) \omega_{b\mu}^f
\end{aligned}$$

Eq.(21)

$$\begin{aligned}
\frac{\partial \mathcal{L}}{\partial \omega_{c\gamma}^b} \omega_{c\varepsilon}^b &= \left[\frac{1}{2\kappa} e \eta^{ae} (e_e^\mu e_b^\gamma - e_e^\gamma e_b^\mu) \omega_{a\mu}^c - \frac{1}{2\kappa} e \eta^{ce} (e_e^\mu e_f^\gamma - e_e^\gamma e_f^\mu) \omega_{b\mu}^f \right] \omega_{c\varepsilon}^b \\
&= \frac{1}{2\kappa} e \eta^{ae} (e_e^\mu e_b^\gamma - e_e^\gamma e_b^\mu) \omega_{c\varepsilon}^b \omega_{a\mu}^c - \frac{1}{2\kappa} e \eta^{ce} (e_e^\mu e_f^\gamma - e_e^\gamma e_f^\mu) \omega_{b\mu}^f \omega_{c\varepsilon}^b \\
&= \frac{1}{2\kappa} e \eta^{ce} (e_e^\mu e_b^\gamma - e_e^\gamma e_b^\mu) \omega_{d\varepsilon}^b \omega_{c\mu}^d - \frac{1}{2\kappa} e \eta^{ce} (e_e^\mu e_b^\gamma - e_e^\gamma e_b^\mu) \omega_{d\mu}^b \omega_{c\varepsilon}^d \\
&= \frac{1}{2\kappa} e \eta^{ce} (e_e^\mu e_b^\gamma - e_e^\gamma e_b^\mu) (\omega_{d\varepsilon}^b \omega_{c\mu}^d - \omega_{d\mu}^b \omega_{c\varepsilon}^d) \\
&\stackrel{(19)}{=} \frac{\partial \mathcal{L}}{\partial \omega_{c\mu,\gamma}^b} (\omega_{d\varepsilon}^b \omega_{c\mu}^d - \omega_{d\mu}^b \omega_{c\varepsilon}^d)
\end{aligned}$$

Eq.(22)

$$\begin{aligned}
t_\varepsilon^\gamma &= \frac{\partial \mathcal{L}}{\partial \omega_{c\mu, \gamma}^b} \mathcal{R}_{c\varepsilon\mu}^b - \delta_\varepsilon^\gamma \mathcal{L} \\
&= \frac{1}{2\kappa} e \eta^{ce} (e_e^\mu e_b^\gamma - e_e^\gamma e_b^\mu) \mathcal{R}_{c\varepsilon\mu}^b - \delta_\varepsilon^\gamma \mathcal{L} \\
&= \frac{1}{2\kappa} \sqrt{-g} \eta^{ce} (e_e^\mu e_b^\gamma - e_e^\gamma e_b^\mu) e_\alpha^b e_c^\beta R_{\beta\varepsilon\mu}^\alpha - \delta_\varepsilon^\gamma \mathcal{L} \\
&= \frac{1}{2\kappa} \sqrt{-g} \eta^{ce} (e_e^\mu e_b^\gamma e_\alpha^b e_c^\beta - e_e^\gamma e_b^\mu e_\alpha^b e_c^\beta) R_{\beta\varepsilon\mu}^\alpha - \delta_\varepsilon^\gamma \mathcal{L} \\
&= \frac{1}{2\kappa} \sqrt{-g} \eta^{ce} (e_e^\mu \delta_\alpha^\gamma e_c^\beta - e_e^\gamma \delta_\alpha^\mu e_c^\beta) R_{\beta\varepsilon\mu}^\alpha - \delta_\varepsilon^\gamma \mathcal{L} \\
&= \frac{1}{2\kappa} \sqrt{-g} (\eta^{ce} e_e^\mu e_c^\beta \delta_\alpha^\gamma R_{\beta\varepsilon\mu}^\alpha - \eta^{ce} e_e^\gamma e_c^\beta \delta_\alpha^\mu R_{\beta\varepsilon\mu}^\alpha) - \delta_\varepsilon^\gamma \mathcal{L} \\
&= \frac{1}{2\kappa} \sqrt{-g} \left(g^{\beta\mu} R_{\beta\varepsilon\mu}^\gamma + g^{\beta\gamma} \delta_\alpha^\mu R_{\beta\mu\varepsilon}^\alpha \right) - \delta_\varepsilon^\gamma \frac{1}{2\kappa} \sqrt{-g} R \\
&= \left(g^{\beta\mu} R_{\beta\varepsilon\mu}^\gamma + g^{\beta\gamma} R_{\beta\varepsilon} - \delta_\varepsilon^\gamma R \right) \frac{1}{2\kappa} \sqrt{-g}
\end{aligned}$$

Eq.(23)

If metric compatible, the following is valid:

$$g_{\alpha\gamma} R_{\beta\varepsilon\mu}^\gamma = R_{\alpha\beta\varepsilon\mu} = -R_{\beta\alpha\varepsilon\mu} = -R_{\alpha\beta\mu\varepsilon} = R_{\varepsilon\mu\alpha\beta}$$

$$\begin{aligned}
t_{\alpha\varepsilon} &= \left(g_{\alpha\gamma} g^{\beta\mu} R_{\beta\varepsilon\mu}^\gamma + R_{\alpha\varepsilon} - g_{\alpha\varepsilon} R \right) \frac{1}{2\kappa} \sqrt{-g} \\
&= \left(g^{\beta\mu} R_{\alpha\beta\varepsilon\mu} + R_{\alpha\varepsilon} - g_{\alpha\varepsilon} R \right) \frac{1}{2\kappa} \sqrt{-g} \\
&= \left(g^{\beta\mu} R_{\beta\alpha\mu\varepsilon} + R_{\alpha\varepsilon} - g_{\alpha\varepsilon} R \right) \frac{1}{2\kappa} \sqrt{-g} \\
&= \left(R_{\alpha\mu\varepsilon}^\mu + R_{\alpha\varepsilon} - g_{\alpha\varepsilon} R \right) \frac{1}{2\kappa} \sqrt{-g} \\
&= \left(R_{\alpha\varepsilon} + R_{\alpha\varepsilon} - g_{\alpha\varepsilon} R \right) \frac{1}{2\kappa} \sqrt{-g} \\
&= \left(2R_{\alpha\varepsilon} - g_{\alpha\varepsilon} R \right) \frac{1}{2\kappa} \sqrt{-g} \\
&= \frac{1}{\kappa} \mathbb{G}_{\alpha\varepsilon} \sqrt{-g}
\end{aligned}$$

Supplementary: Vielbeins

The vielbein formalism chooses a new basis frame $\{\hat{e}_a\}$ rather than a natural basis $\{\partial_\alpha\}$,

$$\partial_\mu = e_\mu^a \hat{e}_a \quad (27)$$

, where e_μ^a is transformation called frame field (or vierbein field). Clearly, each point of frame field is an element in $GL(n, \mathbb{R})$ group. We choose the frame field satisfies:

$$g_{\mu\nu} = e_\mu^a e_\nu^b \eta_{ab} \quad (28)$$

The inverse transformation e_a^μ :

$$e_a^\mu e_\nu^\mu = \delta_\nu^\mu \text{ and } e_a^\mu e_\mu^b = \delta_b^a$$

Define $e = \det(e_\mu^a)$, we have:

$$\begin{aligned} g = \det(g_{\mu\nu}) &\stackrel{(28)}{=} \det(e_\mu^a e_\nu^b \eta_{ab}) = \det(e_\mu^a) \det(e_\nu^b) \det(\eta_{ab}) = \frac{1}{2\kappa} e \frac{1}{2\kappa} e (-1) \\ &\rightarrow e = \sqrt{-g} \end{aligned}$$

Recall the definition of tangent connection $\Gamma_{\alpha\mu}^\beta$:

$$\nabla_\mu \partial_\alpha = \Gamma_{\alpha\mu}^\beta \partial_\beta$$

We define the spin connection $\omega_{a\mu}^b$ in similar way:

$$\nabla_\mu \hat{e}_a = \omega_{a\mu}^b \hat{e}_b$$

The relation between $\Gamma_{\alpha\mu}^\beta$ and $\omega_{a\mu}^b$:

$$\begin{aligned} \nabla_\mu \partial_\alpha &\stackrel{(27)}{=} \nabla_\mu (e_\alpha^a \hat{e}_a) = e_{\alpha,\mu}^a \hat{e}_a + e_\alpha^a \nabla_\mu \hat{e}_a = e_{\alpha,\mu}^a \hat{e}_a + e_\alpha^a \omega_{a\mu}^b \hat{e}_b \\ &= \Gamma_{\alpha\mu}^\beta \partial_\beta = \Gamma_{\alpha\mu}^\beta e_\beta^b \hat{e}_b \\ &\rightarrow (e_{\alpha,\mu}^b + e_\alpha^a \omega_{a\mu}^b - \Gamma_{\alpha\mu}^\beta e_\beta^b) \hat{e}_b = 0 \\ &\rightarrow e_{\alpha,\mu}^b + e_\alpha^a \omega_{a\mu}^b - \Gamma_{\alpha\mu}^\beta e_\beta^b = 0 \end{aligned}$$

We can derive 3 useful formula:

$$e_{\alpha,\mu}^b = e_\beta^b \Gamma_{\alpha\mu}^\beta - e_\alpha^a \omega_{a\mu}^b \quad (29)$$

$$\Gamma_{\alpha\mu}^\beta = e_b^\beta e_\alpha^a \omega_{a\mu}^b + e_b^\beta e_{\alpha,\mu}^b \quad (30)$$

$$\omega_{a\mu}^b = e_a^\alpha e_\beta^b \Gamma_{\alpha\mu}^\beta - e_a^\alpha e_{\alpha,\mu}^b \quad (31)$$

Since $e_a^\alpha e_\alpha^b = \delta_a^b$, then $-e_a^\alpha e_{\alpha,\mu}^b = e_{a,\mu}^\alpha e_\alpha^b$. The Eq.(31) can be rewritten:

$$\begin{aligned} \omega_{a\mu}^b &= e_a^\alpha e_\beta^b \Gamma_{\alpha\mu}^\beta + e_{a,\mu}^\alpha e_\alpha^b \\ &\rightarrow e_b^\gamma \omega_{a\mu}^b = e_b^\gamma e_\alpha^a e_\beta^b \Gamma_{\alpha\mu}^\beta + e_{a,\mu}^\alpha e_\alpha^b e_b^\gamma \\ &\rightarrow e_b^\gamma \omega_{a\mu}^b = e_a^\alpha \Gamma_{\alpha\mu}^\gamma + e_{a,\mu}^\gamma \\ &\rightarrow e_{a,\mu}^\gamma = e_b^\gamma \omega_{a\mu}^b - e_a^\alpha \Gamma_{\alpha\mu}^\gamma \end{aligned} \quad (32)$$

Since frame field takes value in GL group, the spin connection $\omega_{a\mu}^b$ is \mathfrak{gl} -value 1-form. Next, we will derive the relation between the tangent curvature $R_{\lambda\omega\sigma}^\kappa$ and the spin curvature $\mathcal{R}_{b\omega\sigma}^a$. The definition of these curvature are:

$$\begin{aligned}
R_{\lambda\omega\sigma}^\kappa &= \Gamma_{\lambda\sigma,\omega}^\kappa - \Gamma_{\lambda\omega,\sigma}^\kappa + \Gamma_{\nu\omega}^\kappa \Gamma_{\lambda\sigma}^\nu - \Gamma_{\nu\sigma}^\kappa \Gamma_{\lambda\omega}^\nu \\
\mathcal{R}_{b\omega\sigma}^a &= \omega_{b\sigma,\omega}^a - \omega_{b\omega,\sigma}^a + \omega_{c\omega}^a \omega_{b\sigma}^c - \omega_{c\sigma}^a \omega_{b\omega}^c
\end{aligned}$$

First, we take partial derivative on Eq.(30):

$$\begin{aligned}
\Gamma_{\lambda\sigma,\omega}^\kappa &= (e_a^\kappa e_\lambda^b \omega_{b\sigma}^a + e_b^\kappa e_{\lambda,\sigma}^b)_{,\omega} \\
&= e_{a,\omega}^\kappa e_\lambda^b \omega_{b\sigma}^a + e_a^\kappa e_{\lambda,\omega}^b \omega_{b\sigma}^a + e_a^\kappa e_\lambda^b \omega_{b\sigma,\omega}^a + e_{b,\omega}^\kappa e_{\lambda,\sigma}^b + e_b^\kappa e_{\lambda,\sigma\omega}^b \\
&= \underbrace{(e_c^\kappa \omega_{a\omega}^c - e_a^\eta \Gamma_{\eta\omega}^\kappa)}_{(32)} e_\lambda^b \omega_{b\sigma}^a + e_a^\kappa \underbrace{(e_\eta^b \Gamma_{\lambda\omega}^\eta - e_\lambda^c \omega_{c\omega}^b)}_{(29)} \omega_{b\sigma}^a + e_a^\kappa e_\lambda^b \omega_{b\sigma,\omega}^a + \underbrace{(e_c^\kappa \omega_{b\omega}^c - e_b^\eta \Gamma_{\eta\omega}^\kappa)}_{(32)} \underbrace{(e_\gamma^b \Gamma_{\lambda\sigma}^\gamma - e_\lambda^d \omega_{d\sigma}^b)}_{(29)} + e_b^\kappa e_{\lambda,\sigma\omega}^b
\end{aligned}$$

Expand and rearrange:

$$\begin{aligned}
\Gamma_{\lambda\sigma,\omega}^\kappa &= (e_c^\kappa \omega_{a\omega}^c e_\lambda^b \omega_{b\sigma}^a - e_a^\eta \Gamma_{\eta\omega}^\kappa e_\lambda^b \omega_{b\sigma}^a) + (e_a^\kappa e_\eta^b \Gamma_{\lambda\omega}^\eta \omega_{b\sigma}^a - e_a^\kappa e_\lambda^c \omega_{c\omega}^b \omega_{b\sigma}^a) + e_a^\kappa e_\lambda^b \omega_{b\sigma,\omega}^a \\
&\quad + (e_c^\kappa \omega_{b\omega}^c e_\gamma^b \Gamma_{\lambda\sigma}^\gamma - e_c^\kappa \omega_{b\omega}^c e_\lambda^d \omega_{d\sigma}^b - e_b^\eta \Gamma_{\eta\omega}^\kappa e_\gamma^b \Gamma_{\lambda\sigma}^\gamma + e_b^\eta \Gamma_{\eta\omega}^\kappa e_\lambda^d \omega_{d\sigma}^b) + e_b^\kappa e_{\lambda,\sigma\omega}^b \\
&= \left(\underbrace{\frac{e_c^\kappa e_\lambda^b \omega_{a\omega}^c \omega_{b\sigma}^a}{(*a)}}_{(32)} - \underbrace{\frac{e_a^\eta e_\lambda^b \Gamma_{\eta\omega}^\kappa \omega_{b\sigma}^a}{(*b)}}_{(29)} \right) + (e_a^\kappa e_\eta^b \Gamma_{\lambda\omega}^\eta \omega_{b\sigma}^a - e_a^\kappa e_\lambda^b \omega_{c\omega}^b \omega_{b\sigma}^a) + e_a^\kappa e_\lambda^b \omega_{b\sigma,\omega}^a \\
&\quad + \left(e_c^\kappa e_\gamma^b \Gamma_{\lambda\sigma}^\gamma \omega_{b\omega}^c - \underbrace{\frac{e_c^\kappa e_\lambda^b \omega_{a\omega}^c \omega_{b\sigma}^a}{(*a)}}_{(32)} \underbrace{\frac{-\Gamma_{\eta\omega}^\kappa \Gamma_{\lambda\sigma}^\eta}{\text{Move to left}}}_{(32)} + \underbrace{\frac{e_a^\eta e_\lambda^b \Gamma_{\eta\omega}^\kappa \omega_{b\sigma}^a}{(*b)}}_{(29)} \right) + e_b^\kappa e_{\lambda,\sigma\omega}^b
\end{aligned}$$

Remove (*a) and (*b), we have:

$$\Gamma_{\lambda\sigma,\omega}^\kappa + \Gamma_{\eta\omega}^\kappa \Gamma_{\lambda\sigma}^\eta = \left(\frac{e_a^\kappa e_\eta^b \Gamma_{\lambda\omega}^\eta \omega_{b\sigma}^a - e_a^\kappa e_\lambda^b \omega_{c\omega}^b \omega_{b\sigma}^a}{(\#a)} \right) + e_a^\kappa e_\lambda^b \omega_{b\sigma,\omega}^a + \left(\frac{e_c^\kappa e_\gamma^b \Gamma_{\lambda\sigma}^\gamma \omega_{b\omega}^c}{(\#b)} \right) + \frac{e_b^\kappa e_{\lambda,\sigma\omega}^b}{(\#c)} \quad (33)$$

Similarly, swap $\sigma \Leftrightarrow \omega$:

$$\Gamma_{\lambda\omega,\sigma}^\kappa + \Gamma_{\eta\sigma}^\kappa \Gamma_{\lambda\omega}^\eta = \left(\frac{e_a^\kappa e_\gamma^b \Gamma_{\lambda\sigma}^\gamma \omega_{b\omega}^c - e_a^\kappa e_\lambda^b \omega_{c\omega}^b \omega_{b\sigma}^a}{(\#b)} \right) + e_a^\kappa e_\lambda^b \omega_{b\omega,\sigma}^a + \left(\frac{e_c^\kappa e_\eta^b \Gamma_{\lambda\omega}^\eta \omega_{b\sigma}^a}{(\#a)} \right) + \frac{e_b^\kappa e_{\lambda,\omega\sigma}^b}{(\#c)} \quad (34)$$

Eq.(33) - Eq.(34):

$$\begin{aligned}
\Gamma_{\lambda\sigma,\omega}^\kappa - \Gamma_{\lambda\omega,\sigma}^\kappa + \Gamma_{\eta\omega}^\kappa \Gamma_{\lambda\sigma}^\eta - \Gamma_{\eta\sigma}^\kappa \Gamma_{\lambda\omega}^\eta &= e_a^\kappa e_\lambda^b (\omega_{b\sigma,\omega}^a - \omega_{b\omega,\sigma}^a + \omega_{c\omega}^a \omega_{b\sigma}^c - \omega_{c\sigma}^a \omega_{b\omega}^c) \\
R_{\lambda\omega\sigma}^\kappa &= e_a^\kappa e_\lambda^b \mathcal{R}_{b\omega\sigma}^a
\end{aligned}$$