

# 1 Phase field modeling

Given is the following energy functional

$$F[\phi(x, y, t)] = \int_{-\infty}^{\infty} \underbrace{\left( \frac{U}{2} [a^2 ((\partial_x \phi)^2 + (\partial_y \phi)^2) + g(\phi)] + \mu_0 h(\phi) \right)}_{f(\phi, \partial_x \phi, \partial_y \phi)} dx,$$

where  $a$  and  $U$  are constants of the dimension length and energy respectively.  $f(\phi, \partial_x \phi, \partial_y \phi)$ <sup>1</sup> denotes the local free energy density and  $\mu$  the bulk free energy density difference between the two phases. Depending on the sign of  $\mu$ , this can either favor the growth of the one or the other phase.  $g(\phi) = \phi^2(1 - \phi)^2$  is the double well potential and  $h(\phi) = \phi^2(3 - 2\phi)$  is the interpolation function. Please note, that for the reason of phase-stability we have to demand  $|\mu| < U/6$ .

Variational principles provide the phase field equation

$$\begin{aligned} \frac{1}{M_\phi} \frac{\partial \phi}{\partial t} &= - \frac{\delta F}{\delta \phi} \\ &= \partial_x \frac{\partial f}{\partial (\partial_x \phi)} + \partial_y \frac{\partial f}{\partial (\partial_y \phi)} - \frac{\partial f}{\partial \phi} \\ &= U \left( a^2 (\partial_x^2 \phi + \partial_y^2 \phi) - \frac{1}{2} \frac{\partial g(\phi)}{\partial \phi} \right) - \mu \frac{\partial h(\phi)}{\partial \phi}. \end{aligned} \quad (1)$$

## 1.1 Stability of the homogenous and time independent solutions

We look for constant solutions of Eq. 1

$$0 = \frac{U}{2} \frac{\partial g(\phi)}{\partial \phi} + \mu \frac{\partial h(\phi)}{\partial \phi}$$

Calculation of the partial derivatives of the polynomial functions leads to

$$\begin{aligned} \frac{\partial g(\phi)}{\partial \phi} &= 2\phi(1 - \phi)(1 - 2\phi) \\ \frac{\partial h(\phi)}{\partial \phi} &= 6\phi(1 - \phi) \end{aligned}$$

---

<sup>1</sup> $\partial_x$  is an abbreviation for the partial derivative with respect to  $x$ , i.e.  $\partial_x \equiv \partial/\partial x$

Inserting in the equation above yields

$$\begin{aligned}
0 &= U\phi(1-\phi)(1-2\phi) + \mu 6\phi(1-\phi) \\
0 &= \phi(1-\phi)(1-2\phi) + \frac{6\mu}{U}\phi(1-\phi) \\
0 &= \phi(1-\phi) \left( 1-2\phi + \frac{6\mu}{U} \right) \\
\Rightarrow \phi_1 &= 0; \phi_2 = 1; \phi_3 = \frac{1}{2} + \frac{3\mu}{U};
\end{aligned}$$

Stability of the solutions:

- $\phi_1 = 0$  is a global (local) minimum if  $\mu_0 > 0$  ( $\mu_0 < 0$ ), i.e. stabile (meta stabile)
- $\phi_2 = 1$  is local (global) minimum if  $\mu_0 > 0$  ( $\mu_0 < 0$ ), d.h. meta stabile (stabile)
- $\phi_3 = \frac{1}{2} + 3\mu_0/U$  is for positive and negative  $\mu$  unstable (  $|\mu| < U/6$  )

## 1.2 Phase-field profile function

We show that

$$\phi_0(x, t) = \frac{1}{2} \left( 1 + \tanh \frac{(x - vt)}{2a} \right) \quad (2)$$

is a heterogeneous solution of the phase field equation (1) if  $v = 6M_\phi a \mu$ . Note that from  $\partial_x (\tanh(x)) = 1 - \tanh^2(x)$  we can deduct the following property of this solution  $\partial_x \phi_0 = \phi_0 (1 - \phi_0) / a$ .

Calculation of the derivatives:

$$\begin{aligned}
\frac{\partial \phi_0}{\partial x} &= \frac{1}{2} \frac{\partial}{\partial x} \left( 1 + \tanh \frac{(x - vt)}{2a} \right) = \frac{1}{4a} \left( 1 - \tanh^2 \frac{(x - vt)}{2a} \right) \\
&= \frac{1}{4a} \left( 1 + \tanh \frac{(x - vt)}{2a} \right) \left( 1 + 1 - 1 - \tanh \frac{(x - vt)}{2a} \right) \\
&= \frac{1}{4a} \left( 1 + \tanh \frac{(x - vt)}{2a} \right) \left( 2 - \left( 1 + \tanh \frac{(x - vt)}{2a} \right) \right) \\
&= \frac{1}{a} \frac{1}{2} \left( 1 + \tanh \frac{(x - vt)}{2a} \right) \left( 1 - \frac{1}{2} \left( 1 + \tanh \frac{(x - vt)}{2a} \right) \right) \\
&= \frac{1}{a} \phi_0 (1 - \phi_0)
\end{aligned} \quad (3)$$

$$\begin{aligned}
\frac{\partial^2 \phi_0}{\partial x^2} &= \frac{1}{a} \frac{\partial}{\partial x} [\phi_0 (1 - \phi_0)] = \frac{1}{a} \frac{\partial}{\partial \phi_0} [\phi_0 (1 - \phi_0)] \frac{\partial \phi_0}{\partial x} \\
&= \frac{1}{a^2} \phi_0 (1 - \phi_0) (1 - 2\phi_0),
\end{aligned} \quad (4)$$

$$\frac{\partial \phi_0}{\partial t} = -v \frac{\partial \phi_0}{\partial x} = -\frac{v}{a} \phi_0 (1 - \phi_0). \quad (5)$$

Where the second derivative has been calculated using the chain rule  $\frac{\partial f(\varphi(x))}{\partial x} = \frac{\partial f}{\partial \varphi} \frac{\partial \varphi}{\partial x}$ . Inserting in the nonlinear partial differential equation

$$\begin{aligned} -\frac{v}{a}\phi_0(1-\phi_0) &= M_\phi U \underbrace{\left[ a^2 \frac{1}{a^2} \phi_0(1-\phi_0)(1-2\phi_0) - \phi(1-\phi)(1-2\phi) \right]}_{=0} \\ &\quad - M_\phi \mu 6\phi_0(1-\phi_0) \\ \Leftrightarrow \quad -v &= -6M_\phi a\mu. \end{aligned}$$

### 1.3 Interface energy density

The interface energy density  $\gamma$  in the phase-field model corresponds to the total free energy of the heterogeneous solution  $\gamma = F[\phi_0(x, t)]$ :

$$\begin{aligned} F &= \frac{U}{2} \int_{-\infty}^{\infty} \left( a^2 \left( \frac{\partial \phi_0}{\partial x} \right)^2 + \phi_0^2(1-\phi_0)^2 \right) dx \\ &= \frac{U}{2} \int_{-\infty}^{\infty} \left( a^2 \left( \frac{1}{a} \phi_0(1-\phi_0) \right)^2 + \phi_0^2(1-\phi_0)^2 \right) dx \\ \left[ dx = \frac{a}{\phi_0(1-\phi_0)} d\phi_0 \right] &= aU \int_0^1 \frac{\phi_0^2(1-\phi_0)^2}{\phi_0(1-\phi_0)} d\phi_0 \\ &= aU \int_0^1 \phi_0(1-\phi_0) d\phi_0 \\ &= aU \left( \frac{1}{2} \phi_0^2 - \frac{1}{3} \phi_0^3 \right) \Big|_0^1 = \frac{aU}{6} \end{aligned}$$

### 1.4 Calibration of the field model

We calibrate the phase field model according to the 1D considerations above, i.e. we switch from the parameters  $a, U, M_\phi$  to the parameters  $\xi = 2a$  for the phase-field width,  $\Gamma = aU/6$  for interface energy density and  $M = M_\phi a^2 U$  for the kinetic coefficient  $M[\text{m}^2/\text{s}]$ . The calibrated phase-field model provide the following relation between the driving force  $\mu$  and the resulting stationary interface velocity  $v$

$$v = \frac{M}{\Gamma} \mu_0 = K \mu_0.$$

With these parameters we obtain the following phase-field equation

$$\frac{1}{M} \partial_t \phi = \underbrace{\partial_x^2 \phi + \partial_y^2 \phi}_{\text{Laplace-Operator}} - \frac{2}{\xi^2} \partial_\phi g(\phi) - \frac{\mu}{3\Gamma\xi} \partial_\phi h(\phi).$$