

QGS : Quantum groups seminar

Quantum affine algebras and spectral \hbar -matrices

1. The Yang-Baxter equation (YBE)

$$R(z) \in A \otimes A((z)) \text{ s.t. } \xleftarrow{\text{some algebra}}$$

$$R_{12}(z) R_{13}(zw) R_{23}(w) =$$

$$= R_{23}(w) R_{13}(zw) R_{12}(z)$$

$$\begin{array}{ccc} \curvearrowright & & \\ R_{12}(z) = R(z) \otimes I & & R_{23}(z) = I \otimes R(z) \end{array}$$

$$R_{13}(z) = (\text{flip} \otimes I)(R_{23}(z))$$

$R(z)$ = R -matrix with
spectral parameter z

2. Quantum integrable systems

$$\mathcal{A} = \text{end}(\mathbb{C}^2)$$

$q \in \mathbb{C}^\times$ — quantum parameter

$$R(z) := \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & \frac{\bar{q}^1(z-1)}{z - \bar{q}^{-2}} & \frac{1 - \bar{q}^{-2}}{z - \bar{q}^{-2}} & 0 \\ 0 & \frac{z(1 - \bar{q}^{-2})}{z - \bar{q}^{-2}} & \frac{\bar{q}^{-1}(z-1)}{z - \bar{q}^{-2}} & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

$$\begin{array}{c} \text{XXZ } \frac{1}{2}\text{-spin} \\ \text{Heisenberg model} \end{array} \quad \begin{array}{c} \text{6-vertex} \\ \text{model} \end{array} \quad \xrightarrow{\sim} \quad U_q \widehat{\mathfrak{sl}_2}$$



$$\mathbb{C}^2$$

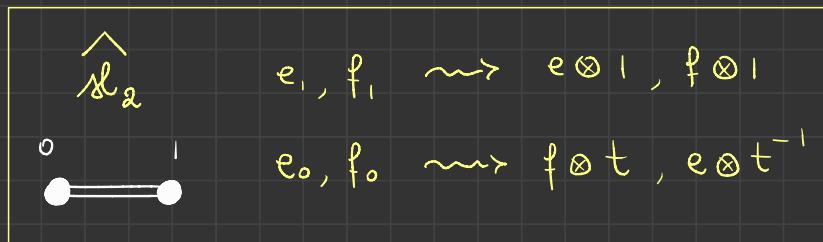
3. Quantum affine algebras

$$\widehat{\mathfrak{g}} = \text{simple f.d. Lie algebra } \mathfrak{g} / \mathbb{C} \quad \text{central}$$

↗ $\widehat{\mathfrak{g}}$ = $\underbrace{\mathfrak{g} \otimes \mathbb{C}[t, t^{-1}]}$ ↙
 ↘ $\mathcal{L}_{\mathfrak{g}}$ ↑
 ↘ derivation

affine Lie algebra
 $\mathfrak{g} \otimes \mathbb{C}[t, t^{-1}]$
loop algebra

(Kac-Moody algebra)



$$U_q \widehat{\mathfrak{g}} \supseteq (U_q \widehat{\mathfrak{g}})' \xrightarrow[q^c=1]{} (U_q \widehat{\mathfrak{g}})' \simeq U_q \mathcal{L}_{\mathfrak{g}}$$

Drinfeld - Jimbo
quantum group

Quantum loop algebra

(usual Kac - Moody)
presentation

(Drinfeld loop)
presentation

4. Representations

$\mathcal{U}_q \widehat{\mathfrak{g}}$



\mathcal{O}^{int}

$\mathcal{U}_q \mathfrak{Log}$



$\text{Rep}^{\text{fd}}(\mathcal{U}_q \mathfrak{Log})$

- $\dim(V) = \infty$
- semi simple
- $\underline{\lambda} \in \mathbb{N}^I$
- braided

- $\dim(V) < \infty$
- non-semi simple !
- $P(u) \in \mathbb{C}[u]^I$
- non-braided !

Remark $\mathcal{U}_q \mathfrak{Log}$ is \mathbb{Z} -graded ($\deg(e_i/f_i) = 0$,
 $\deg(e_0/f_0) = \pm 1$) and endowed with

$(a \in \mathbb{C}) \quad \tau_a \subset \mathcal{U}_q \mathfrak{Log}$

$$\tau_a(x) = a^m x \quad \text{if } \deg(x) = m$$

$$V \in \text{Rep}^{\text{fd}}(\mathcal{U}_q \mathfrak{Log}) \iff V(a) := \tau_a^*(V)$$

Thm (Drinfeld)

$$\exists \quad R(z) \in (\mathcal{U}_q \widehat{\otimes} \mathcal{U}_q)^{\wedge} \text{ s.t. }$$

- $R_{vw}(z) \in \text{end}(V \otimes W)[[z]]$
- flip $\circ R_{vw}(z)$ intertwiner
- $\Delta \otimes 1(R(z)) = R_{13}(z) R_{23}(z)$
- $1 \otimes \Delta(R(z)) = R_{13}(z) R_{12}(z)$

$\rightsquigarrow R_{vw}(z)$ satisfies (YBE).

Idea: $\mathcal{U}_q \widehat{\otimes}$ is quasitriangular Hopf algebra

$$R = q^{\Omega_0} \cdot \sum_{\mu > 0} e_\mu \otimes f_\mu \in \mathcal{U}_q \widehat{\otimes} \mathcal{U}_q$$

$$\Omega_0 = \sum u_i \otimes u^i \in \mathfrak{h}$$

$$+ c \otimes d + d \otimes c$$

$$\rightsquigarrow R_{vw}(z) := \tau_z \otimes 1(R) G_{V \otimes W}$$

5. Reflection equation

From integrable models with boundary conditions (Sklyanin, Cherednik, ...)

$k(z) \in \mathcal{A}(z)$ s.t.

$$k_1(z) R_{21}^{\psi}(zw) k_2(w) R(w/z)$$

$$= R_{21}^{\psi\psi}(w/z) k_2(z) R(zw) k_1(w)$$

$\rightsquigarrow k(z) =$
 *k-matrix with
 spectral parameter z*

examples

- many computational examples for $U_q \mathfrak{so}_N$
(e.g. Delius - George , Regelski - Vlaar ...)
- universal constant solutions for
finite-type $U_q \mathfrak{so}_N$ (Bao - Wang ,
Balagovic - Kolb)

6. Quantum symmetric pairs

In both cases, the construction of the k -matrix is deeply related to Letzter-Kolb quantum symmetric pairs

$B \subseteq U_q \otimes \mathbb{C}^{\times}$ of Kac-Moody

coideal $\Delta(B) \subseteq B \otimes U_q \otimes$

($q=1$) $B \longrightarrow U(\mathfrak{g}^\theta)$

$\theta = \text{Ad}(w_X) \circ w \circ \tau$

finite-type Dynkin
subdiagram

diagram automorphism
with $\tau^2 = 1$ and

$\tau|_X =$ action induced

by w_X

longest element in W_X

$\beta \subseteq \mathcal{U}_g \otimes$ controls the intertwining property of k on $V \in \mathcal{O}^{\text{int}}$:

$\forall v \in V, \beta \in \beta$

$$k_V \cdot \beta \cdot v = \varphi(\beta) \cdot k_V \cdot v$$

↑ some diagram
automorphism

$$\Rightarrow k_V : V \xrightarrow{\sim} V^\varphi \text{ as } \beta\text{-modules}$$

In affine type, the intertwining equation involves the parameter inversion, i.e.

$$k_V(z) : V(z) \longrightarrow V(\frac{1}{z})^\varphi$$

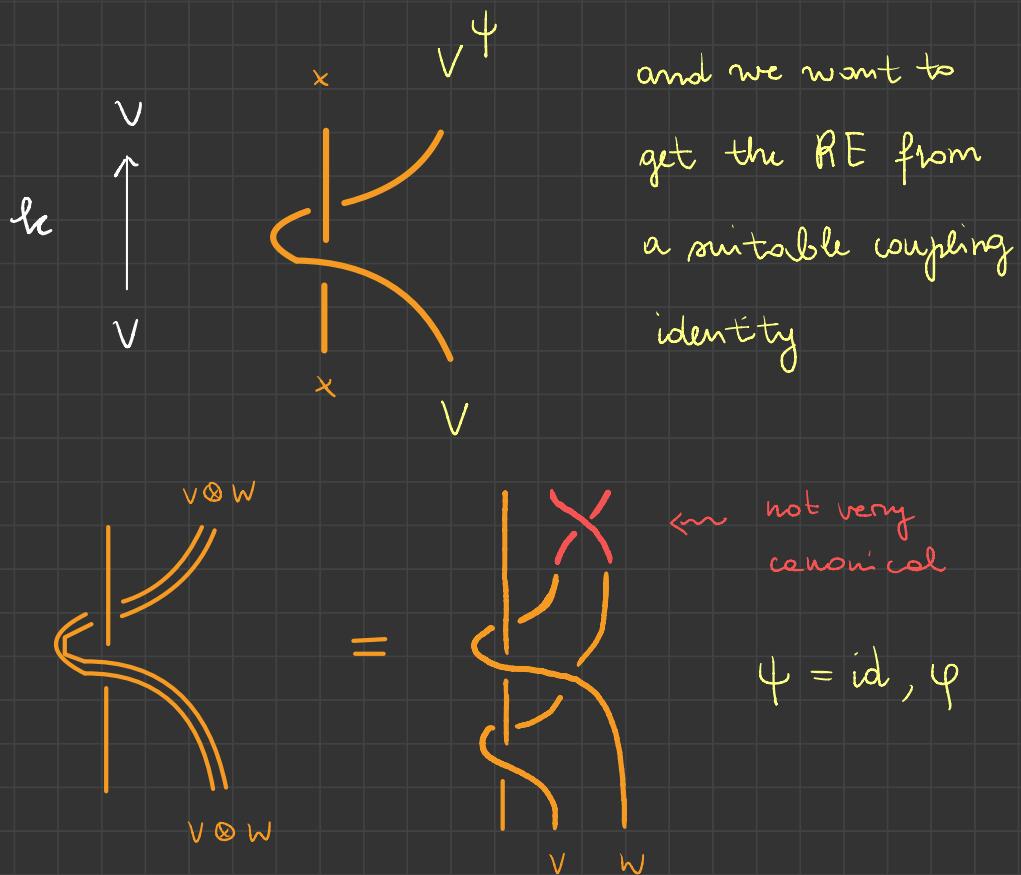
Goal: obtain spectral k -matrices from constant (Kac-Moody) k -matrices.

Big problem: φ cannot perform the parameter inversion

f. A new approach (A.-Vlaar)

- $(H, R) = \text{qt Hopf algebra}$
- $\psi: H \rightarrow H = \text{algebra automorphism}$
- $V \in \text{Rep}(H) =: \mathcal{C}$

We would like an operator $\ell_k: V \rightarrow V^\psi$



Assume instead we have a functor

$\bar{\Phi} : \mathcal{C} \rightarrow \mathcal{C}$ codifying the twisting so that

$k : V \rightarrow \bar{\Phi}(V)$ (say $\bar{\Phi} = \varphi^*$) and

$$\begin{array}{ccc} \bar{\Phi}(v \otimes w) & = & \bar{\Phi}(w) \otimes \bar{\Phi}(v) \\ \text{Diagram: } \begin{array}{c} \text{A vertical line } v \otimes w \text{ with two strands crossing.} \\ \text{The top strand is labeled } \bar{\Phi}(v \otimes w). \\ \text{The bottom strand is labeled } v \otimes w. \end{array} & & \begin{array}{c} \text{A vertical line } \bar{\Phi}(w) \otimes \bar{\Phi}(v) \text{ with two strands crossing.} \\ \text{The top strand is labeled } \bar{\Phi}(v \otimes w). \\ \text{The bottom strands are labeled } v \text{ and } w. \end{array} \end{array}$$

Our choice becomes almost canonical once we

assume $\bar{\Phi}$ is a braided functor

$$(\bar{\Phi}, \gamma) : \mathcal{C} \rightarrow \mathcal{C}^{\text{op}}$$

opposite braided
monoidal category

$$\boxed{\gamma : \bar{\Phi}(v \otimes w) \xrightarrow{\sim} \bar{\Phi}(w) \otimes \bar{\Phi}(v)}$$

precisely what we need !

Algebraically, our new setting is given by the following data:

- (H, R) = quasitriangular Hopf algebra
- $\psi: H \rightarrow H$ = algebra automorphism
- $J \in H \otimes H$ = Drinfeld twist

satisfying

$$(H^{\text{cop}})^{\psi} = H_J \quad R_{21}^{\psi\psi} = R_J$$

where

$(H^{\psi}, R^{\psi\psi})$ is the ψ -twisted qt HA

$$\Delta^{\psi}(x) := \psi \otimes \psi \circ \Delta \circ \psi^{-1}(x)$$

$$R^{\psi\psi} := \psi \otimes \psi(R)$$

(H_J, R_J) is the J -twisted qt HA

$$\Delta_J(x) := J \cdot \Delta(x) \cdot J^{-1}$$

$$R_J := J_{21} \cdot R \cdot J^{-1}$$

We assume (H, R) is fixed and we call (ψ, τ) a twist pair. Note that this is equivalent to have a braided equivalence

$$(\psi^*, \tau) : \text{Rep}(H, R) \longrightarrow \text{Rep}(H, R)^{\text{op}} \\ \parallel \\ \text{Rep}(H^{\text{cop}}, R_{21})$$

Finally we get the k -matrix.

Def A cylindrical bialgebra is a tuple (H, R, ψ, τ, k) as before with $k \in H$ s.t.

$$\Delta(k) = \tau^{-1} \cdot 1 \otimes k \cdot R^+ \cdot k \otimes 1$$

new coupling identity

Cylindrical bialgebras provide the right framework to study a more general reflection equation:

$$\begin{aligned}
 (1) \quad & R_{21}^{\psi\psi} \cdot 1 \otimes k \cdot R^\psi \cdot k \otimes 1 = \\
 & = J_{21} \cdot R \cdot \underbrace{J^{-1} \cdot 1 \otimes k \cdot R^\psi \cdot k \otimes 1}_{\Delta(k)} = \\
 & = k \otimes 1 \cdot R_{21}^\psi \cdot 1 \otimes k \cdot R
 \end{aligned}$$

(2) k, R induce an external action of the cylindrical braid group on $V^{\otimes N}$

(3) Set $\vartheta(x) := k^{-1} \cdot \psi(x) \cdot k$ and let $B_k \subseteq H^\partial$ be the maximal coideal subalgebra. Then, $\text{Rep}(B_k)$ is a braided module category over $\text{Rep}(H)$.

8. New k -matrices for $\mathcal{U}_{q^{\text{of}}}$

Let $\mathcal{U}_{q^{\text{of}}}$ be a quantum Kac-Moody algebra and $\beta \subseteq \mathcal{U}_{q^{\text{of}}}$ a QSP corresponding to $\theta = \text{Ad}(w_x) \circ \omega \circ \tau$.

The construction of a k -matrix for $\mathcal{U}_{q^{\text{of}}}$ supported on β requires the use of quantum Weyl group operators:

T_i = q -analogues of triple exponentials

$[L, KR, LS]$ action of β_W on $\mathcal{U}_{q^{\text{of}}}$ and \mathcal{O}^{int} representations

[Kamnitzer-Tingley] $\exists t_X := (\text{Cartan}) T_{W_X}$

s.t.

$$R_X := t_X^{-1} \otimes t_X^{-1} \cdot \Delta(t_X)$$

t_X is a half-balance for $\mathcal{U}_{q^{\text{of}}} X$

Modified quantum Weyl group operators

Given Satake diagrams (X, τ) and (Y, η)
define suitable corrections

$$T_{X, \tau} := (\text{Cartan}) \cdot T_{W_X}$$

$$R_{X, \tau} := (\text{Cartan}) \cdot R_X$$

and we set

$$\vartheta_q(X, \tau) := \text{Ad}(T_{X, \tau}) \circ \omega \circ \tau \in U_q \otimes$$

$$\psi_{Y, \eta} := \text{Ad}(T_{Y, \eta}) \circ \omega \circ \tau$$

such that $(\psi_{Y, \eta}, R_{Y, \eta})$ is a twist pair.

Rem in finite type, we choose $Y = I$ and
 $\psi = \text{op}_I \circ \tau$ (cf. Balagovic - Kolb).

Finally we can state our

Theorem (A. - Vlaar)

There exists a canonical series $\mathfrak{X}_{X,\tau} = 1 + \sum_{\mu} \mathfrak{X}_{\mu}$
with $\mathfrak{X}_{\mu} \in (\mathcal{U}_q \mathbb{N}_+)_\mu$ s.t.

$$k_{Y,\eta} := \left(T_{Y,\eta}^{-1} \cdot T_{X,\tau} \right) \circ (\text{Cartan}) \circ \mathfrak{X}_{X,\tau}$$

is a k -matrix for the twist pair $(\psi_{Y,\eta}, R_{Y,\eta})$
with support $B_{X,\tau}$.

Affine type $Y = I \setminus \{o, \tau(o)\}$ and $\eta(o) = \tau(o)$

$(o \notin X)$ and $\eta|_Y = \sigma p_Y$. For any

$V \in \text{Rep}^{\text{fd}} \mathcal{U}_q \mathbb{L} \mathfrak{g}$, we get a formal series

$$k_{Y,\eta}(z) \in \text{end}(V)[[z]]$$

satisfying the spectral reflection equation
in $\psi_{Y,\eta}$.