

# Quantum Noise and error mitigation

## Quantum Computing Minicourse ICTP-SAIFR

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8 April 2024

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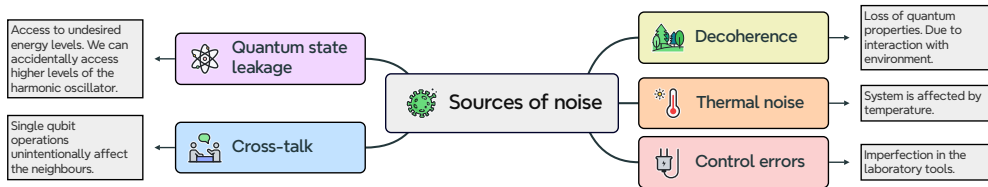


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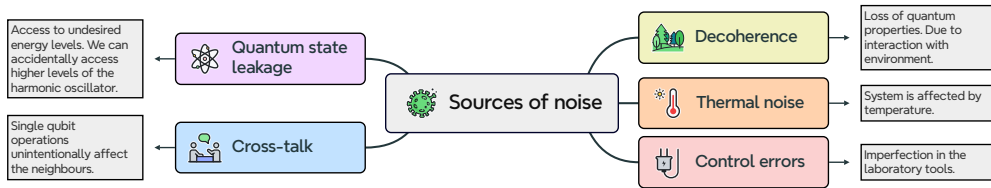
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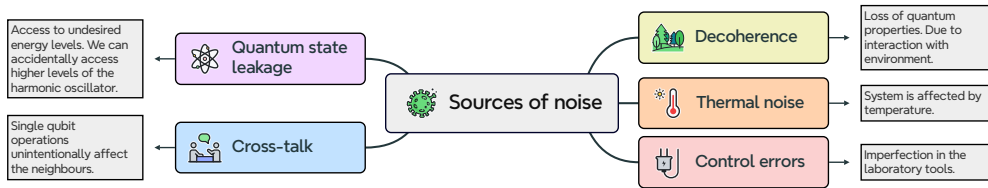


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What about the shot-noise? Is it similar to these noises?

## Without noise we can use state vectors

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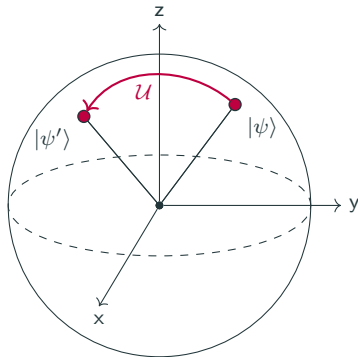
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### About pure states

- Pure states are preserved until the quantum system is isolated, thus doesn't interact with the environment.
- A pure state of a single qubit can be visualized as a point on the surface of the Bloch sphere.
- Using unitaries  $\mathcal{U}$  we can move any point of the sphere into another point in a reversible way.





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Noise corrupts the quantum states, building **mixed states**, namely classical mixture of pure states. An useful tool to represent both pure and mixed states is the **density matrix**, which in general is defined as:

$$\rho = \sum_{i=0}^n p_i |\psi_i\rangle \langle \psi_i|,$$

where  $p_i$  is the statistical weight of the pure state  $|\psi_i\rangle$  in the mix. If  $n = 1$ , then  $\rho = |\psi\rangle \langle \psi|$  and  $|\psi\rangle$  is pure.

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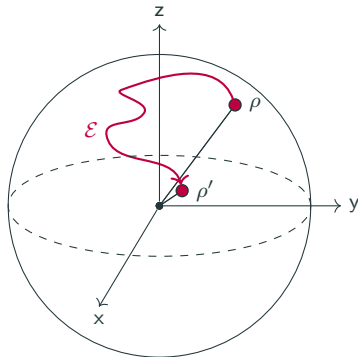
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## About mixed states

- The evolution of a density matrix is described in terms of superoperators  $\rho' = \mathcal{E}(\rho)$ , which is a generalization of the unitary case:  $\rho' = U^\dagger \rho U$ .
- A mixed state of a single qubit can be visualized as a point within the surface of the Bloch sphere.
- The effect of the noise can be seen as moving a point from the surface of the sphere to the inner region. This is not reversible!



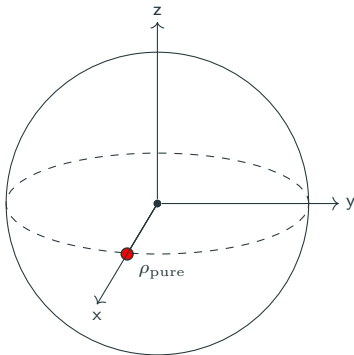


# Pure and mixed states in the Bloch sphere

## Example of simple pure state

The density matrix of  $|+\rangle$  is

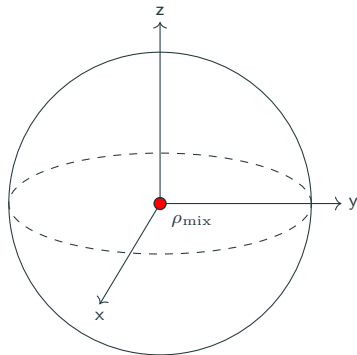
$$\begin{aligned}\rho_{\text{pure}} &= |+\rangle \langle +| = \\ &= \begin{bmatrix} 1/\sqrt{2} & 1/\sqrt{2} \end{bmatrix} \begin{bmatrix} 1/\sqrt{2} \\ 1/\sqrt{2} \end{bmatrix} = \begin{bmatrix} 1/2 & 1/2 \\ 1/2 & 1/2 \end{bmatrix}\end{aligned}$$



## Example of simple mixed state

The density matrix of  $\frac{1}{2}(|0\rangle \langle 0| + |1\rangle \langle 1|)$  is

$$\begin{aligned}\rho_{\text{mix}} &= \frac{1}{2}(|0\rangle \langle 0| + |1\rangle \langle 1|) = \\ &= \frac{1}{2} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} = \begin{bmatrix} 1/2 & 0 \\ 0 & 1/2 \end{bmatrix}\end{aligned}$$



# The Pauli representation of the noise

We can now define a superoperator  $\mathcal{N}$ , which represents the action of the noise on a state:  $\rho_{\text{noisy}} = \mathcal{N}(\rho)$ .

We can exploit the Pauli's operator to write a noise model:

- we use  $X$  to apply random bitflip  $|0\rangle = X|1\rangle$  and  $|1\rangle = X|0\rangle$ ;
- we use  $Z$  to apply random phase flip  $|0\rangle = Z|0\rangle$  and  $-|1\rangle = Z|1\rangle$ ;
- we use  $Y$  to apply more complex manipulations since  $Y = iXZ$ .

Then the combined effect of these Pauli components is:

$$\mathcal{N}(\rho) = p_I \rho + p_x X \rho X + p_y Y \rho Y + p_z Z \rho Z,$$

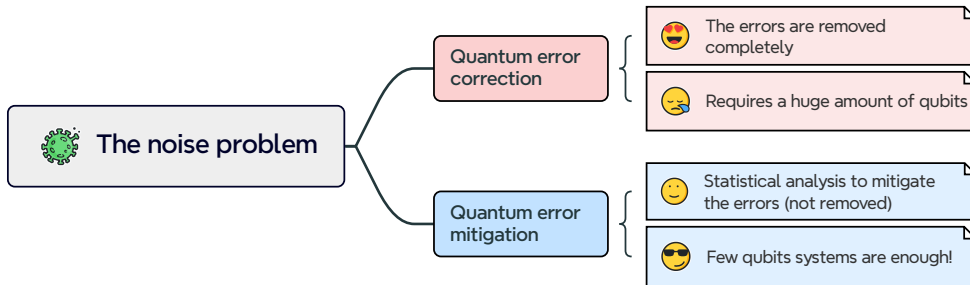
where  $p_I + p_x + p_y + p_z = 1$  and the **first term** of the sum represent the case in which  $\rho$  remains unchanged.

## Final comments

1. The effect of this noise model is to push any state  $\rho$  closer to the maximally mixed state  $\rho_{\text{mix}} = \frac{1}{2^N} I$ .
2. The expectation value of a Pauli observable ( $X, Y, Z$  or a combination  $P$ ) over  $\rho_{\text{mix}}$  is zero.
3. From 1. and 2., the more intense is the noise, the more  $\langle \rho | P | \rho \rangle \rightarrow 0$



# How can we face this problem?



The goal of quantum error mitigation is to use our knowledge about the noise to try to mitigate its effect.

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Suppose we are interested in computing the expectation value  $\langle P \rangle^0$  of an observable  $P$  over the state we get applying the quantum circuit  $\mathcal{U}$  to a state  $\rho_0$ :

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But our system is noisy! Thus what we get is a noisy expectation value  $\langle P \rangle_{\text{noisy}}^0$ .