

Stars and Planets Problem Set1

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Exercise I.1 Radiation Pressure

(a)

The dimension of P_{rad} is $[E][L]^{-3}$. The physical quantities that are related to pressure from the derivation of state equations are:

$$\begin{aligned}kT &\sim [E] \\ h &\sim [E][T]\end{aligned}$$

However, only with the two quantities above we cannot cancel out T and add in L . Given the pressure is produced by radiation, we can introduce the speed of light $c \sim [L][T]^{-1}$ to help us. After some calculations, we can get that the dimension of radiation pressure is:

$$P_{\text{rad}} \sim \frac{(kT)^4}{h^3 c^3}$$

(b)

From the energy density u_ν , we can get the distribution function for the energy $f(\nu)$:

$$f(\nu) = u_\nu / (h\nu) = \frac{8\pi\nu^2}{c^3} \frac{1}{e^{h\nu/kT} - 1}$$

And the relation between the distribution function for the energy and the momenta is:

$$\begin{aligned}f(\nu)d\nu &= f(p)dp \\ p &= \frac{h}{\lambda} = \frac{h\nu}{c} \\ dp &= \frac{h}{c}d\nu\end{aligned}$$

Therefore, the distribution function for the momenta is:

$$f(p) = \frac{8\pi}{h^3} p^2 \frac{1}{e^{cp/kT} - 1}$$

The radiation pressure is:

$$P_{\text{rad}} = \int_0^{+\infty} cp f(p) dp = \frac{8\pi^5}{45} \frac{(kT)^4}{h^3 c^3}$$

The relation between radiation energy density and radiation pressure is:

$$P_{\text{rad}} = \frac{u_{\text{rad}}}{3}$$

Exercise I.2 Physical structure of protoplanetary disks

(a)

The standard astrophysical “cosmic” abundances for the elements gives:

$$X = 0.75, Y = 0.25$$

Since the hydrogen and helium gas are in molecular form, the relative mean molecular weight is:

$$\begin{aligned}\mu_{\text{H}_2} &= 2 \\ \mu_{\text{He}} &= 4\end{aligned}$$

The mean molecular weight μ of the gas is:

$$\begin{aligned}\frac{1}{\mu} &= \frac{X}{\mu_{\text{H}_2}} + \frac{Y}{\mu_{\text{He}}} \\ \mu &= 2.29\end{aligned}$$

(b)

Assume the mass of the star is M , the gravitational constant is G and the positive direction of z axis is pointing upward. Write down Newton’s second law in the vertical direction:

$$g_{z,\star} = -\frac{GM}{r^2 + z^2} \frac{z}{\sqrt{r^2 + z^2}} = -\frac{GM}{r^2} \frac{z}{(1 + (\frac{z}{r})^2)^{\frac{3}{2}}}$$

For $z \ll r$:

$$\frac{z}{(1 + (\frac{z}{r})^2)^{\frac{3}{2}}} = z(1 - \frac{3}{2}(\frac{z}{r})^2) \sim z$$

Therefore:

$$g_{z,\star} = -\frac{GM}{r^2} z \approx -\Omega_K^2 z$$

(c)

The equation of state for the gas in the disk is:

$$P = nkT = \frac{\rho}{\mu m_u} kT$$

Since the $g_{z,\star}$ has included the minus sign, the hydrostatic balance is:

$$\frac{dP}{dz} = \rho g_{z,\star}$$

Combine the two equations and we get the ordinary differential equation for $\rho(z)$:

$$\frac{d\rho}{dz} = -\frac{\Omega_K^2 \mu m_u}{kT} z \rho$$

The solution is:

$$\rho(z) = C e^{-\frac{\Omega_K^2 \mu m_u}{2kT} z^2}$$

(d)

The standard deviation σ or the scaleheight H is:

$$\sigma = H = \sqrt{\frac{kT}{\mu m_u \Omega_K^2}}$$

And the $\rho(z)$ can be written as:

$$\rho(z) = C e^{-\frac{z^2}{2H^2}}$$

(e)

The surface density is:

$$\begin{aligned}\Sigma &= \int_{-\infty}^{+\infty} \rho(z) dz = \int_{-\infty}^{+\infty} C e^{-\frac{z^2}{2H^2}} dz \\ \Sigma &= CH\sqrt{2\pi} \\ C &= \frac{\Sigma}{H\sqrt{2\pi}}\end{aligned}$$

Therefore, $\rho(z)$ can be rewritten as:

$$\rho(z) = \frac{\Sigma}{H\sqrt{2\pi}} e^{-\frac{z^2}{2H^2}}$$

(f)

When $z \ll r$, use the small angle approximation:

$$\tan \theta = \frac{z}{r} \approx \theta$$

Since z only depends on r , take derivatives on both sides:

$$\frac{d\theta}{dr} = \frac{d}{dr}\left(\frac{z}{r}\right)$$

The fraction ϵ of the stellar luminosity L_\star that is being intercepted by the disk at radius interval $[r, r + \Delta r]$ is:

$$\begin{aligned}\epsilon &= \frac{1}{2\pi} \int \sin \theta d\theta d\phi = \int_{\theta}^{\theta+d\theta} \sin \theta d\theta = \int_r^{r+dr} \frac{z}{r} \frac{d}{dr}\left(\frac{z}{r}\right) dr \\ &= \frac{1}{2} \frac{d}{dr}\left(\frac{z}{r}\right)^2 \Delta r = \frac{\beta^2}{2} \frac{d}{dr}\left(\frac{H}{r}\right)^2 \Delta r = \beta^2 \frac{H}{r} \left(\frac{H}{r}\right)' \Delta r\end{aligned}$$