



High-Quality
Axions
in a Class of Chiral
 $U(1)$ Gauge
Theories

Yu-Cheng QIU

Strong CP and
Axion

Framework

QCD Axion

Fuzzy DM

Quintessence
Axion

Summary

High-Quality Axions in a Class of Chiral $U(1)$ Gauge Theories

Yu-Cheng QIU



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with Jin-Wei Wang, Tsutomu T. Yanagida



Who ordered the Axion?

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Strong CP problem

$$\mathcal{L} \supset \bar{\theta} G \tilde{G} , \quad \bar{\theta} < 10^{-10}$$

Why $\bar{\theta}$ is so small?

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Strong CP problem

$$\mathcal{L} \supset \bar{\theta} G \tilde{G}, \quad \bar{\theta} < 10^{-10}$$

Why $\bar{\theta}$ is so small?

► Peccei–Quinn mechanism:

$$\text{anomalous global } U(1) \xrightarrow{SSB} \text{shift symmetry } \frac{a}{F_a} \rightarrow \frac{a}{F_a} + \delta_{\text{PQ}}$$

$$\xrightarrow{\text{anomaly}} \mathcal{L} \supset \left(\bar{\theta} + \frac{a}{F_a} \right) G \tilde{G} \xrightarrow{\text{instanton}} V(a) \sim \cos \left(\bar{\theta} + \frac{a}{F_a} \right)$$



Why axion?

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► Observed Cosmological Constant

$$\Lambda_{\text{obs}} \simeq 10^{-120} M_{\text{Pl}}^4$$

- If it is the potential of a scalar, the mass must be extremely small, $\sim 10^{-33}$ eV.

► Fuzzy (Wave) Dark Matter

$$m \sim 10^{-22} - 10^{-19} \text{ eV}$$

All of them could be Nambu-Goldstone boson! Axion!

How good is the axion served?

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$U(1)_{\text{PQ}}$ is a global symmetry,
which could easily be explicitly broken **at least** by gravity (wormhole).

Suppose the axion belongs to $\Phi \sim e^{ia/F_a}$, then

$$g \frac{\Phi^n}{M_{\text{Pl}}^{n-4}} \xrightarrow{\langle \Phi \rangle \sim F_a} \Delta V(a) = M_{\text{Pl}}^4 \left(\frac{F_a}{M_{\text{Pl}}} \right)^n \cos \left(\frac{a}{F_a} + \arg[g] \right)$$

$$V = V(a) + \Delta V(a) \quad \bar{\theta} \xrightarrow{< 10^{-10}} \quad n > 10 \quad \text{for } F_a \sim 10^{16} \text{ GeV}$$

One shall suppress all terms at least up to order **10!**



Axion Quality Problem

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The goal is to construct a (field theory) framework
that can easily have high quality axions.

The recipe

- ▶ Suppose there are 2 complex scalar fields,
then one has global $U(1) \times U(1)$. **phase rotations**

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- ▶ Suppose there are 2 complex scalar fields,
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- ▶ One linear combination could be gauged, $U(1)_g$,
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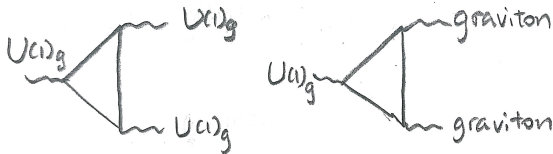
- ▶ Suppose there are 2 complex scalar fields,
then one has global $U(1) \times U(1)$. **phase rotations**
- ▶ One linear combination could be gauged, $U(1)_g$,
and another one is the global $U(1)_{\text{PQ}}$.
- ▶ To have an anomalous $U(1)_{\text{PQ}}$, chiral fermions shall be introduced.



Chiral $U(1)$ Gauge Theory

Introduce N pairs of chiral fermions charged under $U(1)_g$, $\{\psi_i, \bar{\psi}_i\}$, where $\psi_i \in (\mathbf{3}, 1, 0)$ and $\bar{\psi}_i \in (\mathbf{3}^*, 1, 0)$ under $SU(3)_c \times SU(2)_L \times U(1)_Y$. Anomaly cancellation condition requires that

$$\sum_i^N Q_i^3 = 0, \quad \sum_i^N Q_i = 0$$



e.g. $N = 5$, $\{Q_i\} = \{-9, -5, -1, 7, 8\} \times 2$

- ▶ Assume that 2 Higgs given all chiral fermion mass,
- ▶ and only through **Yukawa** interaction.

For $N = k + l$ pairs,

$$\mathcal{L} \supset -\phi_1 \sum_{i=1}^k c_i \psi_i \bar{\psi}_i - \phi_2 \sum_{j=1}^l c_j \psi_j \bar{\psi}_j$$

TABLE I. Fermion charge assignment.

i	1	2	3	4	...				$k+1$	$k+2$...					
ψ_i	α_1	β_1	α_2	β_2	...	γ_1	γ_2	...	δ_1	η_1	δ_2	η_2	...	σ_1	σ_2	...
$\bar{\psi}_i$	β_1	α_1	β_2	α_2	...	γ_1	γ_2	...	η_1	δ_1	η_2	δ_2	...	σ_1	σ_2	...

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$$-q_1 = \alpha_1 + \beta_1 = \alpha_2 + \beta_2 = \cdots = 2\gamma_1 = 2\gamma_2 = \cdots$$

$$-q_2 = \delta_1 + \eta_1 = \delta_2 + \eta_2 = \cdots = 2\sigma_1 = 2\sigma_2 = \cdots$$

By $\sum Q_i = 0$, one could have

$$\sum_i (\alpha_i + \beta_i) + \sum_j \gamma_j + \sum_i (\delta_i + \eta_i) + \sum_j \sigma_j = 0$$

$$\implies -\frac{q_1}{q_2} = \frac{k}{l} = \frac{m}{n}$$

(n, m) are relatively prime numbers to each other.

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(n, m) are relatively prime numbers to each other.

The **lowest-order** non-renormalizable operator $(\phi_{1,2})$ that obey $U(1)_g$ is (?)

$$\mathcal{O} = \frac{1}{k!l!} \frac{\phi_1^k \phi_2^l}{M_{\text{Pl}}^{N-4}}$$

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Focus on the axionic mode.

$$\phi_1 \sim e^{i\tilde{a}/f_1}, \quad \phi_2 \sim e^{i\tilde{b}/f_2}$$

Spontaneous symmetry breaking, $\langle \phi_1 \rangle \sim f_1$ and $\langle \phi_2 \rangle \sim f_2$.

$$|D_\mu \phi_1|^2 + |D_\mu \phi_2|^2 \rightarrow \frac{1}{2}(\partial_\mu a)^2 + m_A^2 \left(A_\mu - \frac{1}{m_A} \partial_\mu b \right)^2$$

where

$$\begin{pmatrix} a \\ b \end{pmatrix} = \frac{-1}{\sqrt{q_1^2 f_1^2 + q_2^2 f_2^2}} \begin{pmatrix} q_2 f_2 & -q_1 f_1 \\ q_1 f_1 & q_2 f_2 \end{pmatrix} \begin{pmatrix} \tilde{a} \\ \tilde{b} \end{pmatrix}$$

$$\frac{a}{F_a} \in [0, 2\pi), \quad F_a = \frac{f_1 f_2}{\sqrt{m^2 f_1^2 + n^2 f_2^2}}$$

Breaking Operator

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$$\mathcal{O} = \frac{1}{k!l!} \frac{\phi_1^k \phi_2^l}{M_{\text{Pl}}^{N-4}} \propto \exp\left(i \frac{a N_{\text{DW}}}{F_a}\right), \quad N_{\text{DW}} = \text{GCD}[k, l]$$

Anomaly term is

$$\mathcal{L} \supset N_{\text{DW}} \frac{a}{F_a} G \tilde{G}$$

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Anomaly term is

$$\mathcal{L} \supset N_{\text{DW}} \frac{a}{F_a} G \tilde{G}$$

- ▶ Choosing large k and l , one could have highly suppressed \mathcal{O} .
- ▶ Given k and l , one needs to solve for fermion charges Q_i under

$$\sum_i Q_i^3 = 0$$

A general mathematical problem. Tricks Required!

Take $k = 3$ and $l = 2x$, and assume that $\text{GCD}[3, 2x] = 1$.

$$\delta\bar{\theta} \sim \frac{2}{3!(2x)!} \left(\frac{F_a}{M_{\text{Pl}}} \right)^N \left(\frac{9 + 4x^2}{2} \right)^{N/2} \frac{M_{\text{Pl}}^4}{m_\pi^2 F_\pi^2}$$

- $F_a = 10^{12} \text{ GeV}$, $x = 7$, $\delta\bar{\theta} \sim 10^{-26}$

$$\{-19, -9, -14 \mid -17, 23, -4, 10, -2, 8, -2, 8, -2, 8, -2, 8, 1, 5\}$$

- $F_a = 10^9 \text{ GeV}$, $x = 4$, $\delta\bar{\theta} \sim 10^{-23}$

$$\{-5, -3, -4 \mid -3, 6, 1, 2, 1, 2, 1, 2\}$$

Assume that $\psi_i \in (1, \mathbf{2}, 0)$ and $\bar{\psi}_i \in (1, \mathbf{2}^*, 0)$ under SM gauge group.

$$\mathcal{L} \supset N_{\text{DW}} \frac{a}{F_a} W \tilde{W}$$

- Electroweak instanton contribution negligible without SUSY.

Nomura, Watari, and Yanagida (2000)

- Potential mainly from \mathcal{O} .

$$V = \frac{\Lambda_a}{2} \left(1 - \cos \frac{aN_{\text{DW}}}{F_a} \right), \quad \Lambda_a = \frac{2^{2-N/2}}{k!} \frac{f_1^k f_2^l}{M_{\text{Pl}}^{N-4}}$$

The mass of axion is determined by $\partial_a^2 V|_{\min}$,

$$m_a = N_{\text{DW}} M_{\text{Pl}} \sqrt{\frac{2}{k!l!}} \left(\frac{m^2 + n^2}{2} \right)^{N/4} \left(\frac{F_a}{M_{\text{Pl}}} \right)^{-1+N/2}$$

- $k = 7, l = 42$ and $F_a = 10^{16} \text{ GeV} \longrightarrow m_a = 2.5 \times 10^{-20} \text{ eV}$.

$$\{-21|4, 3, 16, -9, 20, -13\} \times 7, \quad N_{\text{DW}} = 7$$

Same as the Fuzzy DM scenario, the potential is

$$V = \frac{\Lambda_a}{2} \left(1 - \cos \frac{aN_{\text{DW}}}{F_a} \right), \quad \Lambda_a = \frac{2^{2-N/2}}{k!l!} \frac{f_1^k f_2^l}{M_{\text{Pl}}^{N-4}}$$

- $k = 8$ and $l = 72$, $\Lambda_a = 1.88\Lambda_{\text{obs}}$ and $F_a/N_{\text{DW}} = 3.3 \times 10^{16} \text{ GeV}$.

$$\{-27|5, 1, 15, , -9, 19, -13, 29, -23, 3\} \times 8, \quad N_{\text{DW}} = 8$$

Higher order (non Yukawa) fermion mass term may generate

$$\mathcal{O} \stackrel{?}{\sim} \frac{1}{M_{\psi}^{N-4}} \phi_1^k \phi_2^l, \quad M_{\psi} < M_{\text{Pl}}$$

Is there a general way to avoid this possible loop?

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$$\mathcal{L} \supset -\phi_1 \sum_i c_i \psi_i \bar{\psi}_i - \phi_2 \sum_j c_j \psi_j \bar{\psi}_j$$

Interestingly, there is an anomaly-free \mathbf{Z}_{2N} symmetry, where

$$\psi_i(+1), \quad \bar{\psi}_i(+1), \quad \phi_{1,2}(-2).$$

We gauge it.

Now we have gauged $U(1)_g$ and \mathbf{Z}_{2N} .

The Z_{2N} and fermion loop.

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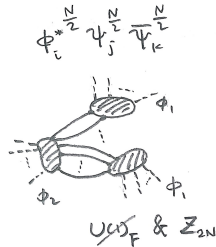
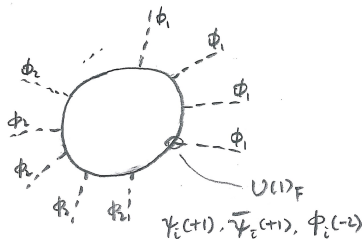
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$$\mathcal{L}_0 = -\phi_1 \sum_i c_i \psi_i \bar{\psi}_i - \phi_2 \sum_j c_j \psi_j \bar{\psi}_j$$

Accidentally, it contains a global $U(1)_F$.



- ▶ Converging diagram does not contribute to \mathcal{O} .
- ▶ Diverging diagrams are more suppressed due to the loop factors.

Cosmic evolution

$$\ddot{a} + 3H(t)\dot{a} + \partial_a V = 0$$

Suppose a_{ini} locates around the bottom, $\partial_a V \sim m_a^2 a$,

$$m_a \sim \sqrt{\frac{\Lambda_a N_{\text{DW}}}{F_a^2}} > \sqrt{N_{\text{DW}}} \times 10^{-33} \text{ eV} \quad \text{for} \quad F_a \lesssim M_{\text{Pl}}$$

One has to put a_{ini} around the hilltop.

One needs large enough $\frac{F_a}{N_{\text{DW}}}$ to stabilize it.

Previously,

- $k = 8$ and $l = 72$, $\Lambda_a = 1.88\Lambda_{\text{obs}}$ and $\frac{F_a}{N_{\text{DW}}} = 3.3 \times 10^{16} \text{ GeV}$.

Too complicated!

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Summary

Gauged \mathbf{Z}_{2N} and $N = 80$ pairs of fermions are needed!
Instability!

ugly!

How about something simple and beautiful?

Extra Dimension!

New Scenario on Quintessence Axion

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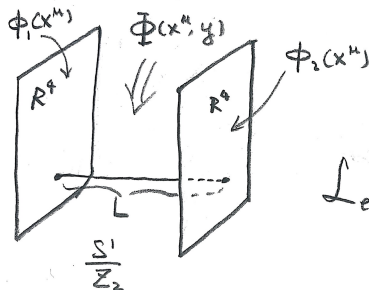
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Flat extra dimension. Fundamental scale is M_* .



$$M_* L = \left(\frac{M_{\text{Pl}}}{M_*} \right)^2$$

$$\mathcal{L}_{\text{eff}} \supset M_*^2 \phi_1 \phi_2 e^{-M_* L}$$

2303.02852 (PRD)

with Sudhakantha Girmohanta, Jin-Wei Wang, Tsutomu T. Yanagida

$$V = \frac{\Lambda_a}{2} \left(1 - \cos \frac{a}{F_a} \right)$$

$$\Lambda_a \approx \Lambda_{\text{obs}} \left(\frac{M_*}{1.47 \times 10^{17} \text{ GeV}} \right)^{515}, \quad F_a = \frac{M_*}{\sqrt{2}} \approx 1.04 \times 10^{17} \text{ GeV}$$

- ▶ Gauged \mathbf{Z}_{2N} is not needed.
- ▶ Only 2 pairs of fermions are needed.
- ▶ No instability problem. $F_a > 10^{17} \text{ GeV}$.
- ▶ Simple $U(1)_g$ charge assignment.

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- ▶ High quality axion(s) could be easily constructed in this framework.
- ▶ The quality is protected by gauged $U(1)_g$ and \mathbf{Z}_{2N} .
- ▶ The QCD axion constructed under this framework has $N_{\text{DW}} = 1$.
- ▶ $N_{\text{DW}} \frac{a}{F_a} W \tilde{W} \longrightarrow c_\gamma \frac{a}{F_a} F \tilde{F}$, quintessence axion could explain the isotropic cosmic birefringence. [Lin & Yanagida \(2022\)](#)
- ▶ If extra dimension is introduced, model become simpler and quintessence axion becomes more attractive.
- ▶ Lots of massive fermions may leave traces in the Cosmological Collider.
- ▶ ...

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Thank you!