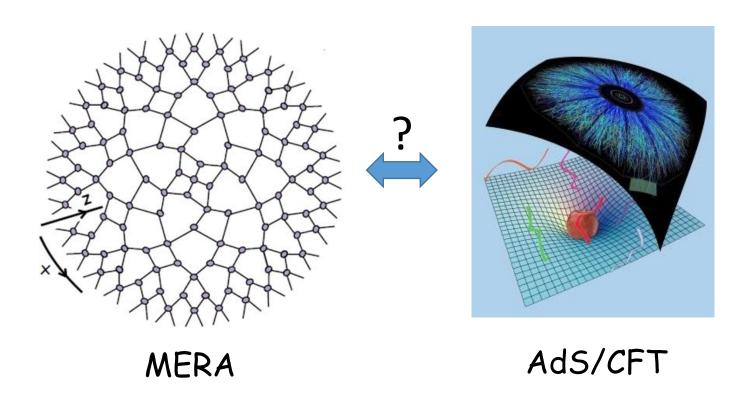
Tensor networks as toy models for holography?

Sukhi Singh

Institute for Quantum Optics and Quantum Information, Vienna



Contents

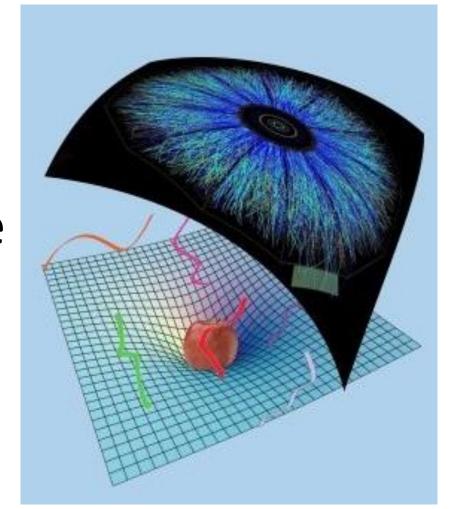
1) The holographic correspondence

2) Tensor network representations of ground states

3) Tensor networks (MERA) and holography

4) Open questions and Outlook

The holographic correspondence



Quantum gravity

- No universal agreement on what the problem is.
- The popular/conservative view is that gravity must be described by a quantum field theory.

o Challenges for a QFT description of gravity:

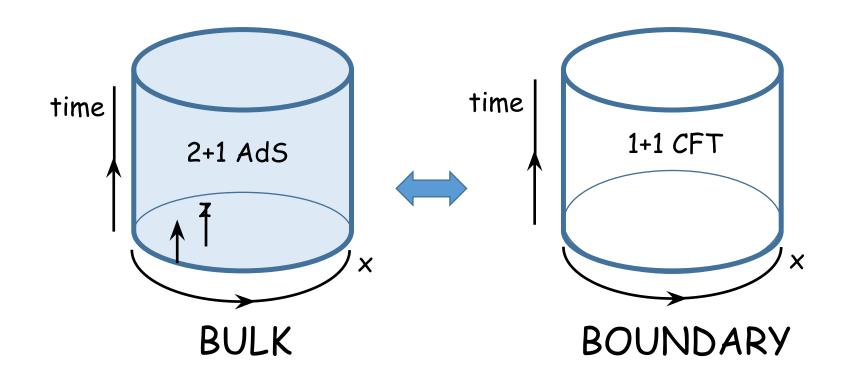
- A naïve QFT description is non-renormalizable and plagued with infinities.
- Gravity is different from other forces since gravity corresponds to dynamical space time is dynamical while QFT is usually formulated on a fixed space time.
- Gravity is not a fundamental force (entropic force)
- . .

The holographic principle

- Certain theories of gravity are equivalent to QFTs but in one less dimension!
- Conjecture from extrapolating insights from studying (quantum) black holes.
- QFT encodes the gravity system like a 2D hologram encodes a 3D image.
- This sounds like an intriguing idea but is there a concrete realization of this idea?
- The AdS/CFT correspondence

J. Maldacena, Adv. Theor. Math. Phys. 2, 231 (1998)

E. Witten, Adv. Theor. Math. Phys. 2, 253 (1998)



Gravity in anti-deSitter (AdS) spacetime



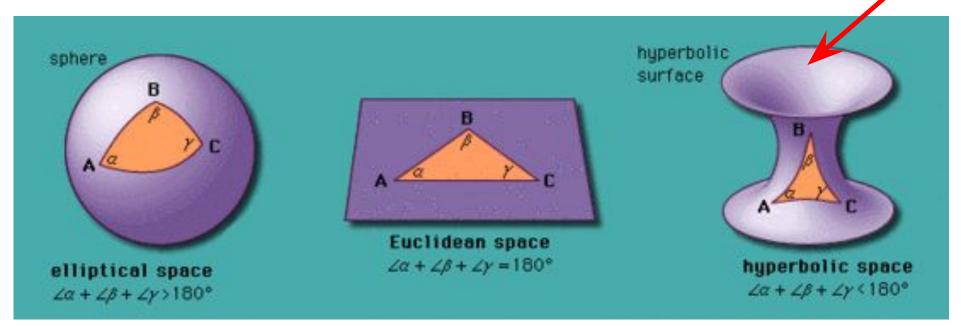
Conformal field theory (CFT)

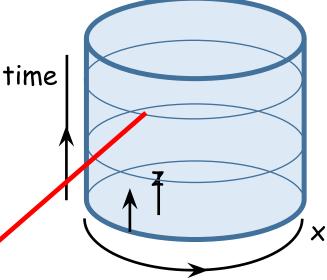
Anti-deSitter spacetime

 Solution to the vacuum Einstein equations with negative cosmological constant.

Space time with constant negative curvature.

Can be foliated into hyperbolic time-like surfaces.





Penrose diagram of 2+1 AdS spacetime as a solid cylinder.

Conformal field theory

- Action is invariant under the action of the conformal group
- In 1+1 dims, the action is invariant under global change of scale.

$$x^{\mu} \rightarrow \lambda x^{\mu}$$

- In 1+1 dims, the symmetry is very constraining. A CFT is completely specified by the following list of data:
- A list of primary fields
- Scaling dimension of each field
- Central charge
- OPE coefficients (which determine the two point correlators of the primary fields)

- The bulk and boundary theories are equivalent: Partition functions are equal.
- Concrete recipe to compute bulk correlation functions from the boundary theory and vice-versa.

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- Concrete recipe to compute bulk correlation functions from the boundary theory and vice-versa.
- AdS/CFT is a strong/weak type duality. Strongly coupled gravity (quantum gravity) is dual to a weakly coupled CFT, and vice-versa!

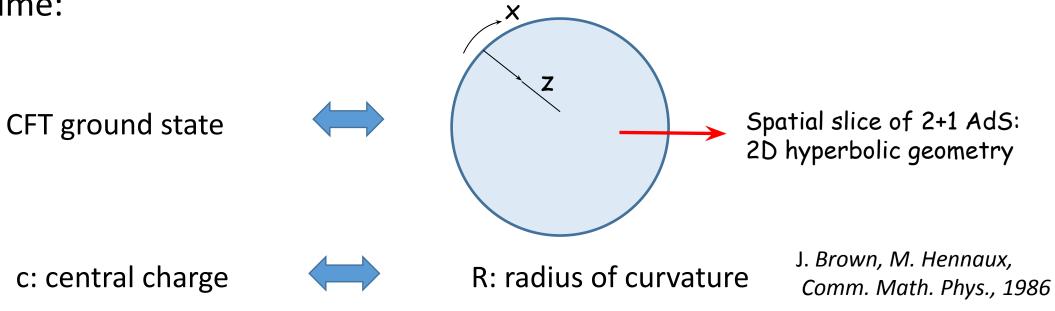
- The bulk and boundary theories are equivalent: Partition functions are equal.
- Concrete recipe to compute bulk correlation functions from the boundary theory and vice-versa.
- AdS/CFT is a strong/weak type duality. Strongly coupled gravity (quantum gravity) is dual to a weakly coupled CFT, and vice-versa!
- A formulation of quantum gravity as a (weakly coupled) CFT.
- Also, useful in condensed matter physics since CFTs appear commonly as effective description of certain strongly correlated matter (more on this later)

 Subir Sachdev et al

 Examples of how boundary and bulk properties are related in the weak gravity regime:

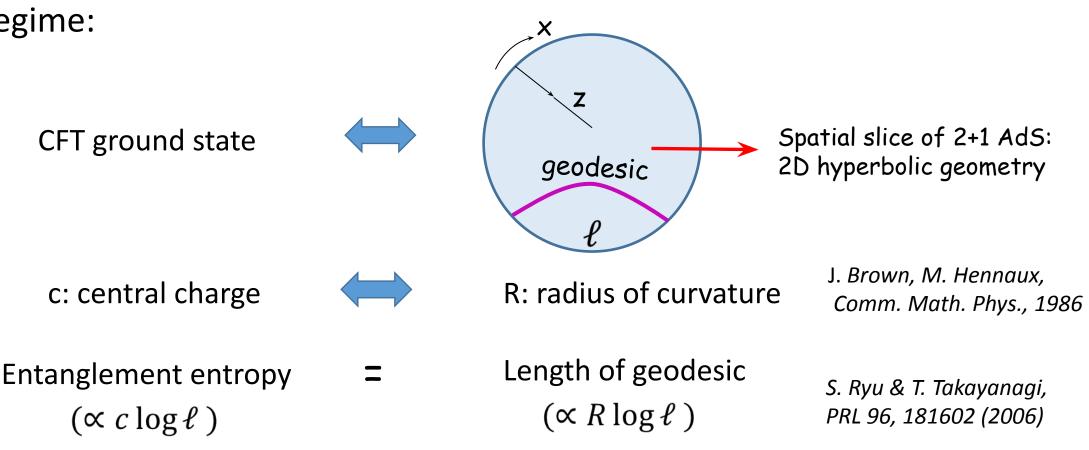
Examples of how boundary and bulk properties are related in the weak gravity

regime:



Examples of how boundary and bulk properties are related in the weak gravity

regime:

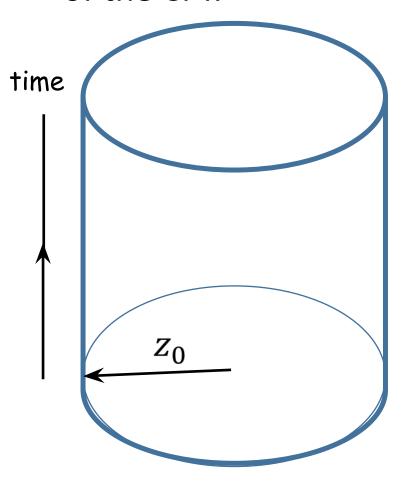


Examples of how boundary and bulk properties are related in the weak gravity

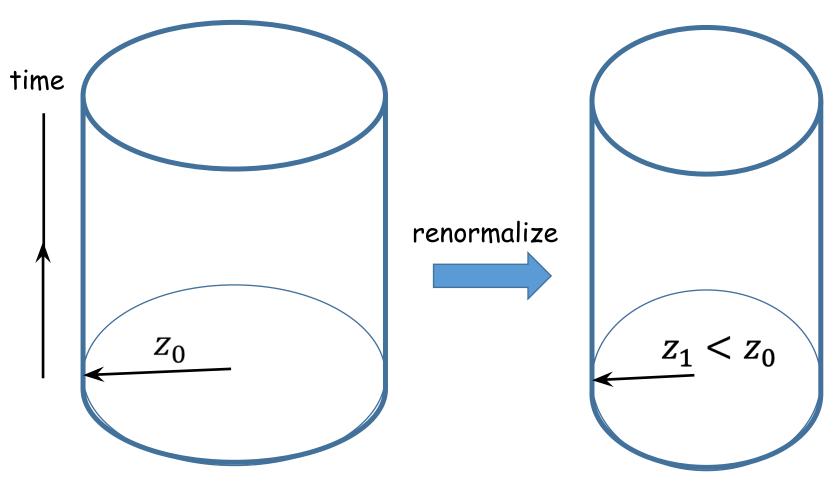
regime: CFT ground state Spatial slice of 2+1 AdS: 2D hyperbolic geometry geodesic J. Brown, M. Hennaux, R: radius of curvature c: central charge Comm. Math. Phys., 1986 Length of geodesic Entanglement entropy S. Ryu & T. Takayanagi, $(\propto R \log \ell)$ $(\propto c \log \ell)$ PRL 96, 181602 (2006)

A thermal state of the CFT corresponds to a AdS geometry with a black hole.

 The extra dimension on the gravity side corresponds to length or energy scale of the CFT.



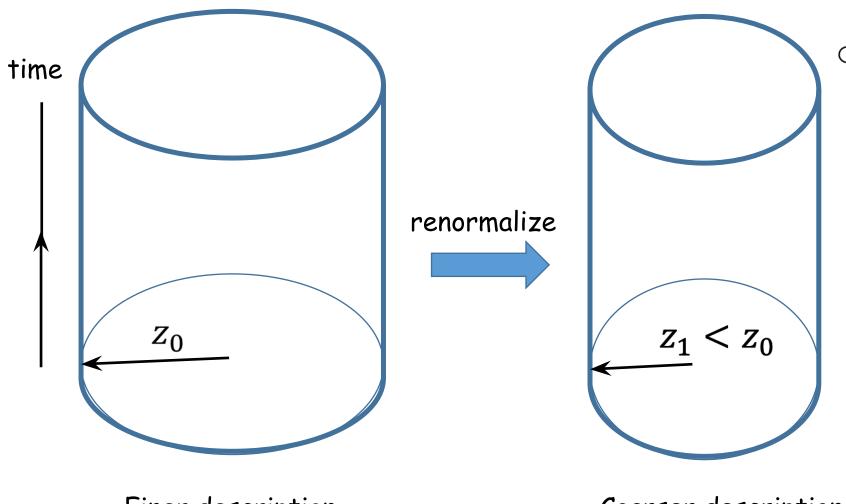
 The extra dimension on the gravity side corresponds to length or energy scale of the CFT.



Finer description

Coarser description

 The extra dimension on the gravity side corresponds to length or energy scale of the CFT.

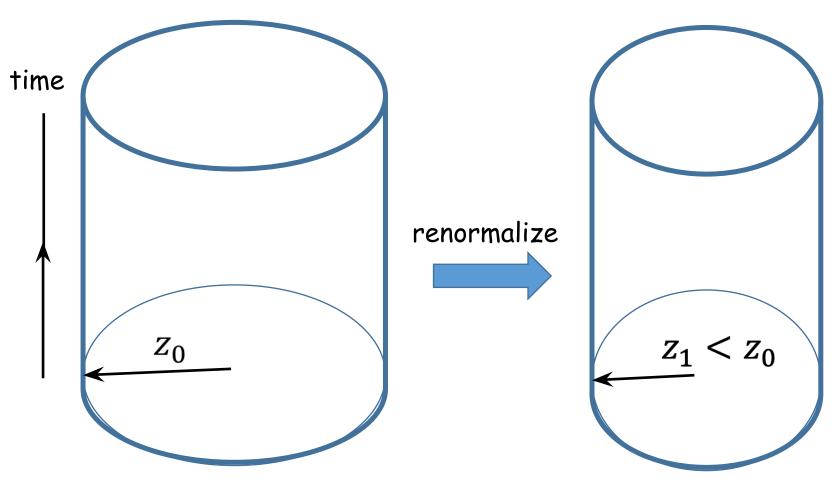


'Glue' effective descriptions of the CFT at different length scales $z_0 \rightarrow z_1 \rightarrow z_2 \rightarrow \dots$ into a solid cylinder.

Finer description

Coarser description

 The extra dimension on the gravity side corresponds to length or energy scale of the CFT.

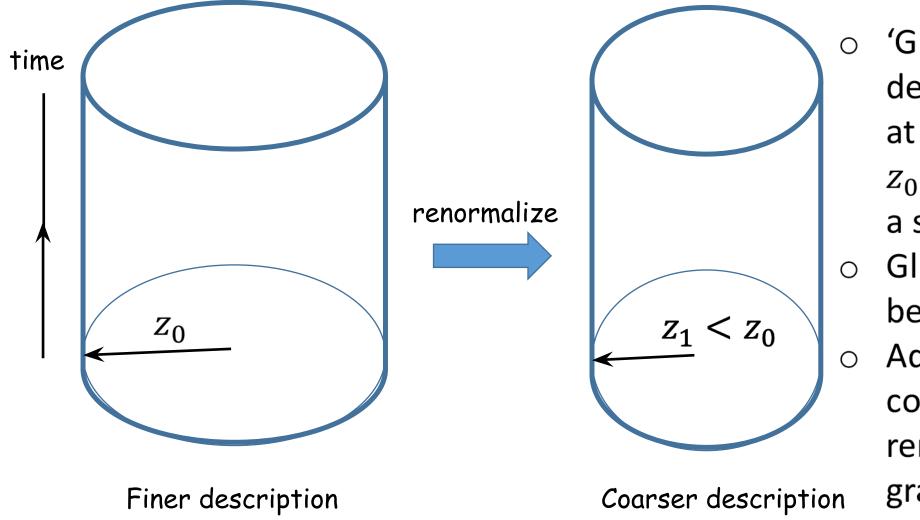


- 'Glue' effective descriptions of the CFT at different length scales $z_0 \rightarrow z_1 \rightarrow z_2 \rightarrow \dots$ into a solid cylinder.
- Glued descriptions can be recast as gravity.

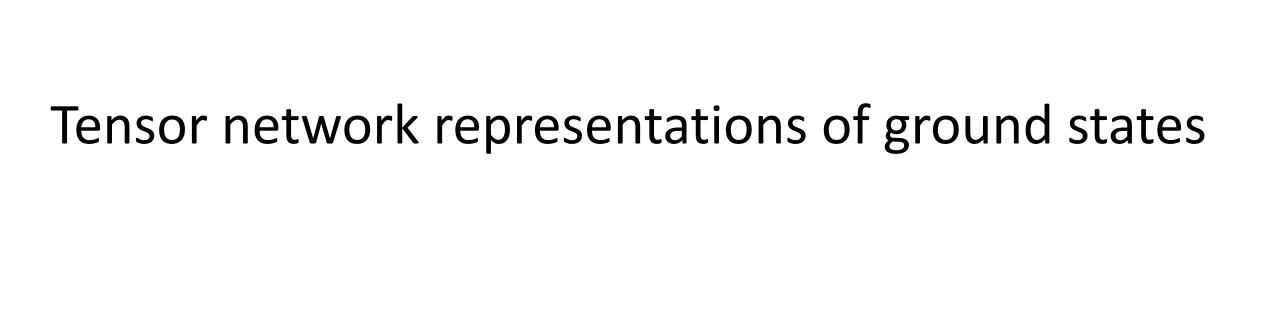
Finer description

Coarser description

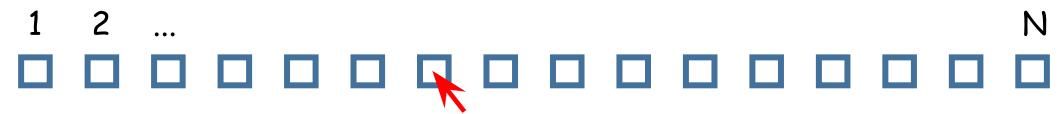
 The extra dimension on the gravity side corresponds to length or energy scale of the CFT.



- 'Glue' effective descriptions of the CFT at different length scales $z_0 \rightarrow z_1 \rightarrow z_2 \rightarrow \dots$ into a solid cylinder.
- Glued descriptions can be recast as gravity.
 - AdS/CFT hints at a connection between renormalization & gravity.



Typical problem in condensed matter physics



qubit, atom, spin, ... (Hilbert space of dimension d)

Local and translational invariant Hamiltonian

$$H(m) = \sum_{k} X_k X_{k+1} + m Z_k$$
 X, Z : Pauli matrices 1D quantum Ising model

Typical problem in condensed matter physics



qubit, atom, spin, ... (Hilbert space of dimension d)

Local and translational invariant Hamiltonian

$$H(m) = \sum_{k} X_k X_{k+1} + m Z_k$$
 X, Z : Pauli matrices 1D quantum Ising model

 \circ Find the phase diagram of the system at zero temperature \square Find the ground state for large N

$$m=0$$

Magnetic Disordered

 $m=1$

"critical point" (gapless)

(described by a CFT in the continuum limit!)

Key challenge

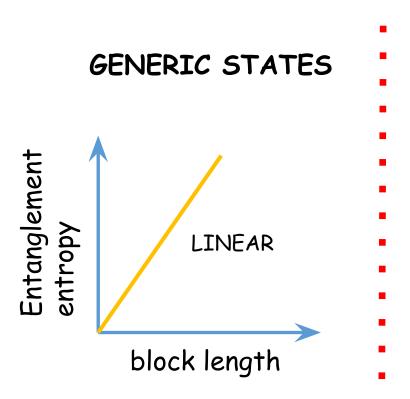
- Exponentially large Hilbert space.
- Have to diagonalize an exponentially large Hamiltonian
- \circ Ground state specified by d^N complex numbers

$$|\Psi\rangle = \sum \Psi_{i_1 i_2 \dots i_N} |i_1\rangle \otimes |i_2\rangle \otimes \dots |i_N\rangle$$

 $\{|i_k\rangle\}$: orthonormal basis on site k

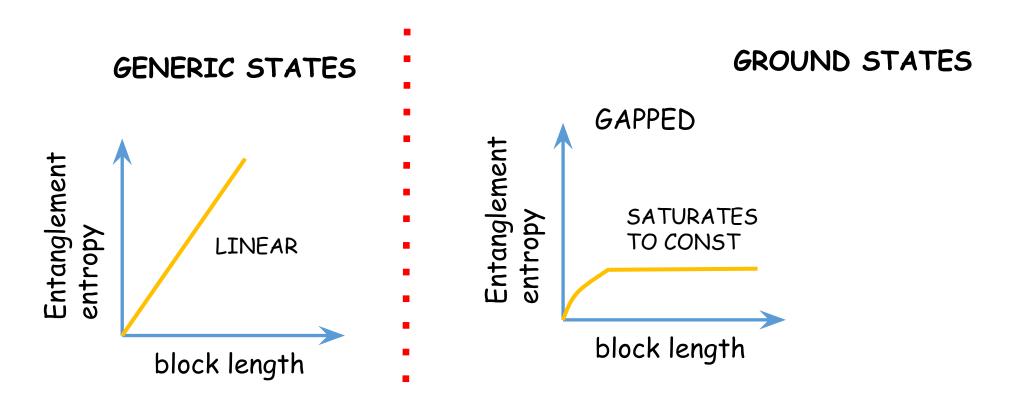
- (a) d=2, N=30 requires 17 GB memory
- (b) d=3, N=20 requires 55 GB memory

They have much less entanglement than generic states on the lattice.



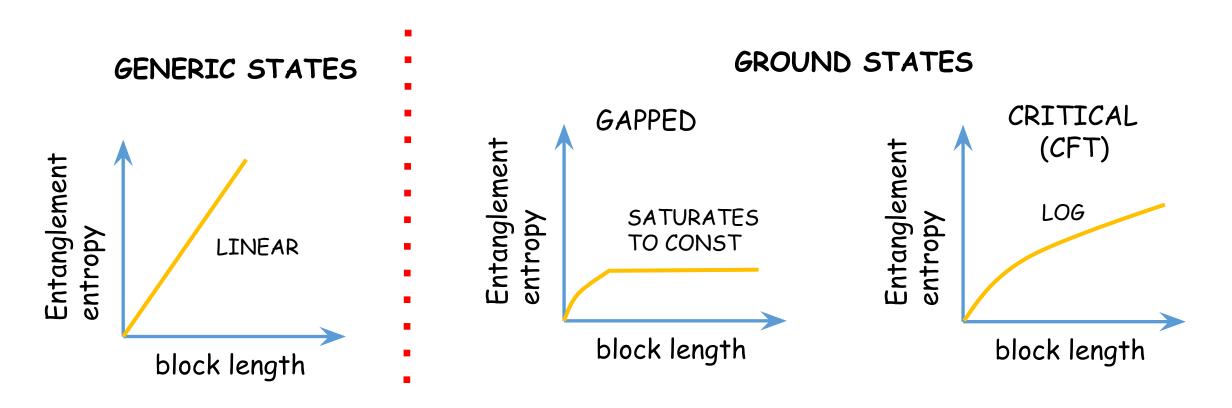
Entanglement entropy $(\rho) = -Tr(\rho \log \rho)$

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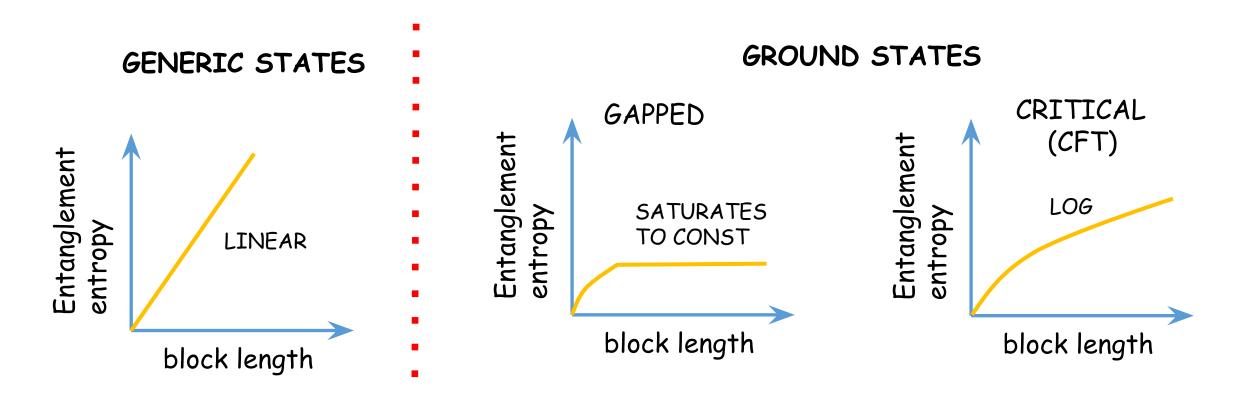
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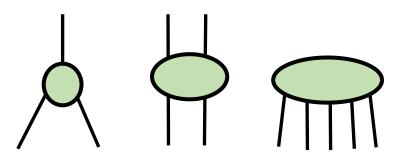
$$|c\rangle = \sum_{i=1}^{d} c_i |i\rangle$$

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$$M = \sum_{i} M_{ij} |i\rangle\langle j|$$

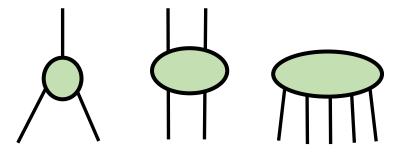
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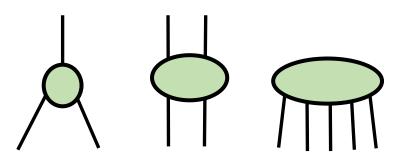
$$M = \sum M_{ij} |i\rangle\langle j| \qquad \bigcirc_{i}^{j}$$



$$C_{ij} = \sum_{k} A_{ik} B_{kj}$$

$$|c\rangle = \sum_{i=1}^{d} c_i |i\rangle$$

$$M = \sum_{i} M_{ij} |i\rangle\langle j| \qquad \bigcirc_{i}^{J}$$



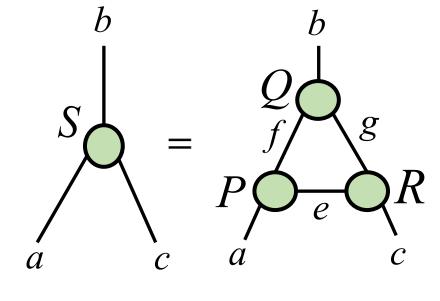
matrix product

$$C \bigcirc = \bigcirc_{j}^{i} A$$

$$B$$

$$C_{ij} = \sum_{k} A_{ik} B_{k}$$

tensor contraction



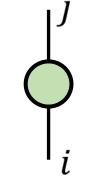
$$S_{abc} = \sum_{efg} P_{afe} Q_{fbg} R_{egc}$$

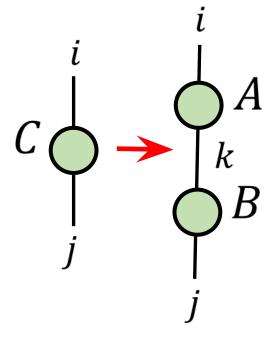
$$|\mathbf{c}\rangle = \sum_{i=1}^{d} \mathbf{c}_i |i\rangle$$

matrix decomposition

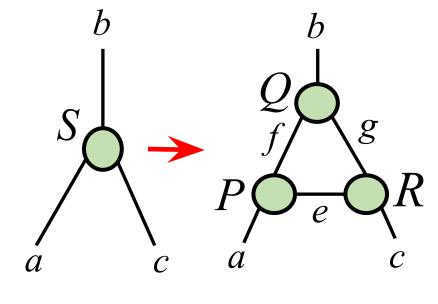
tensor decomposition

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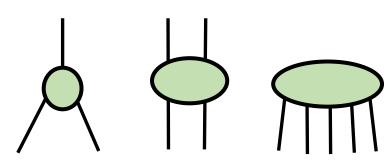




$$C_{ij} = \sum_{k} A_{ik} B_{k}$$

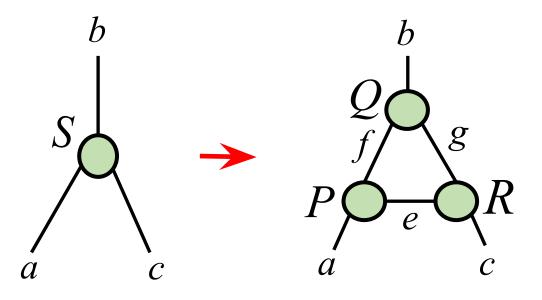


$$S_{abc} = \sum_{efg} P_{afe} Q_{fbg} R_{egc}$$



tensor

tensor network



Ground states as tensor networks

$$|\Psi\rangle = \sum \Psi_{i_1 i_2 \dots i_N} |i_1\rangle \otimes |i_2\rangle \otimes \dots |i_N\rangle$$

(tensor)

$$\begin{array}{|c|c|c|c|c|c|}\hline \Psi & = \\ \hline \vdots & \vdots & \vdots \\ i_1 & i_2 & & i_N \end{array}$$

 d^N coefficients



$$|\Psi\rangle = \sum \Psi_{i_1 i_2 \dots i_N} |i_1\rangle \otimes |i_2\rangle \otimes \dots |i_N\rangle$$

 $(tensor) \rightarrow (tensor network)$

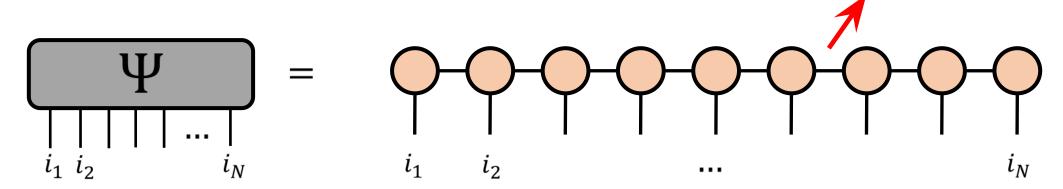
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$$(tensor) \rightarrow (tensor network)$$

Bond index $1, ..., \chi$



 d^N coefficients

$$|\Psi\rangle = \sum \Psi_{i_1 i_2 \dots i_N} |i_1\rangle \otimes |i_2\rangle \otimes \dots |i_N\rangle$$

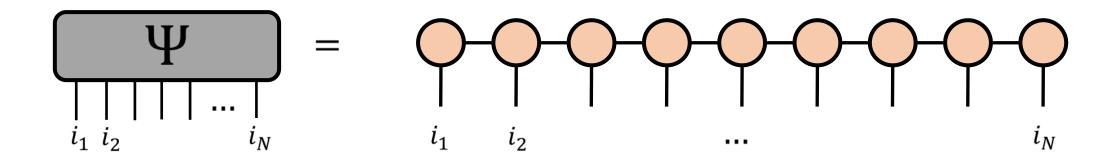
$$(tensor) \rightarrow (tensor network)$$

Bond index $1, ..., \chi$

 d^N coefficients

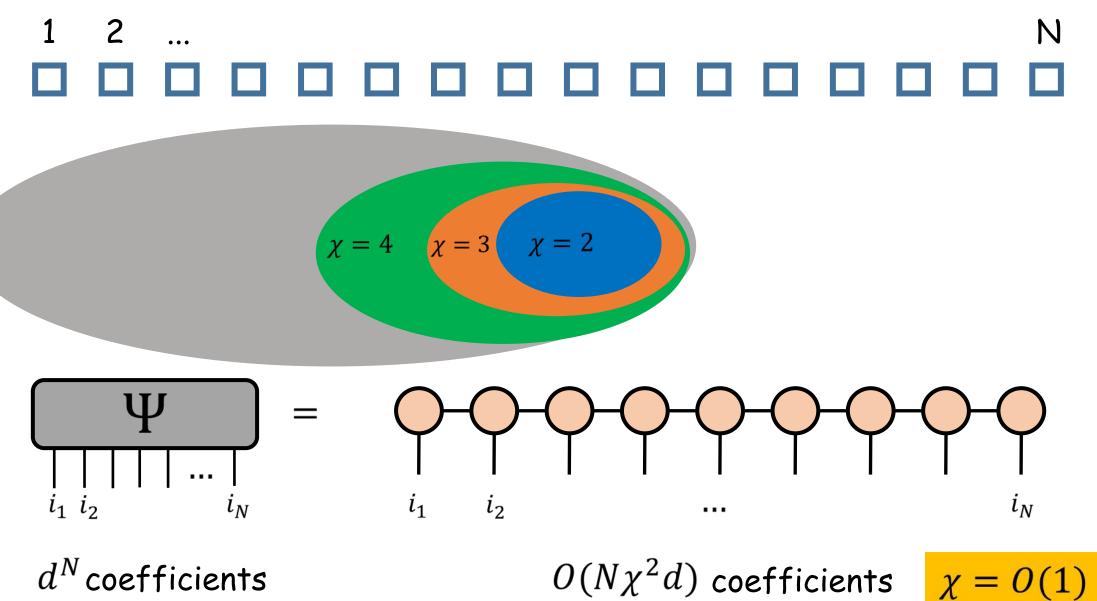
 $O(N\chi^2 d)$ coefficients $\chi = O(\exp N)$

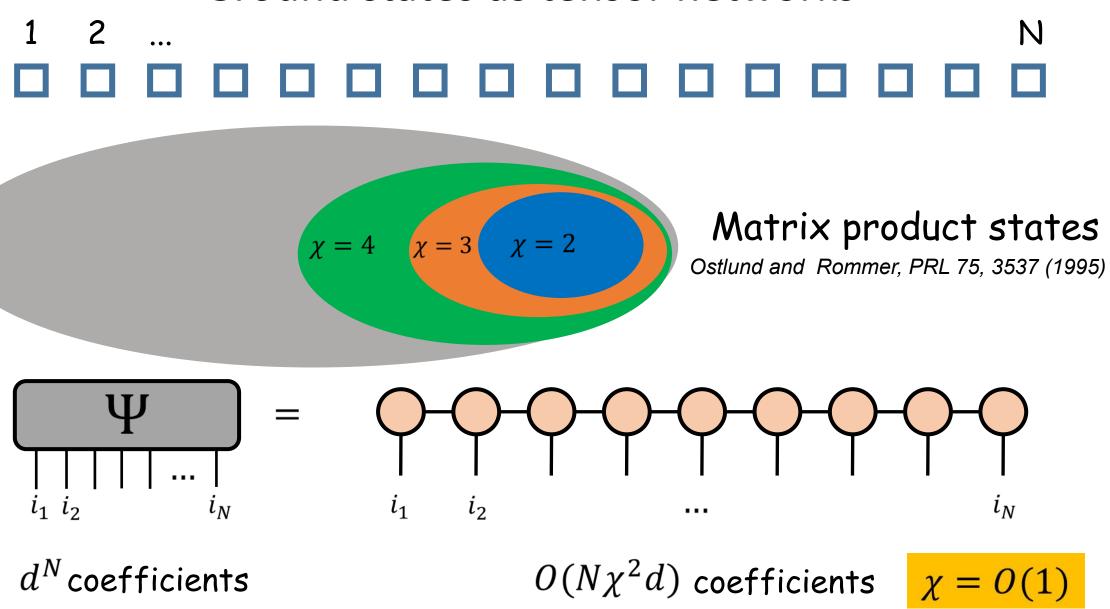


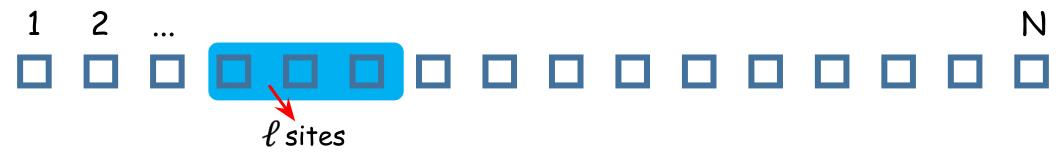


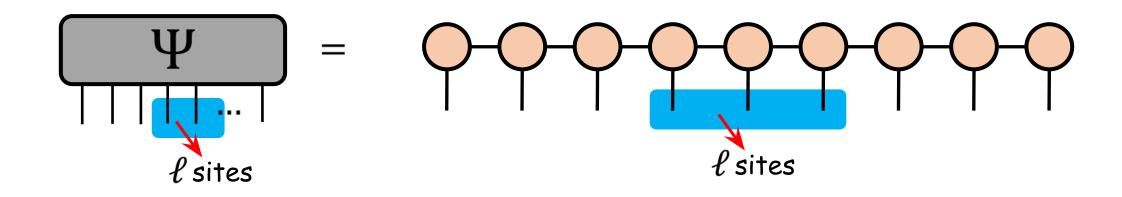
 d^N coefficients

 $O(N\chi^2 d)$ coefficients





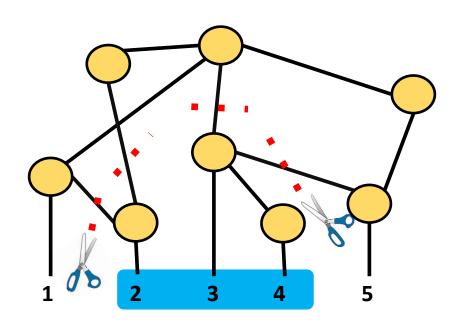


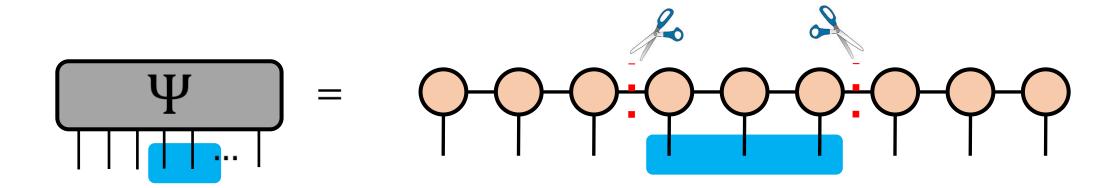


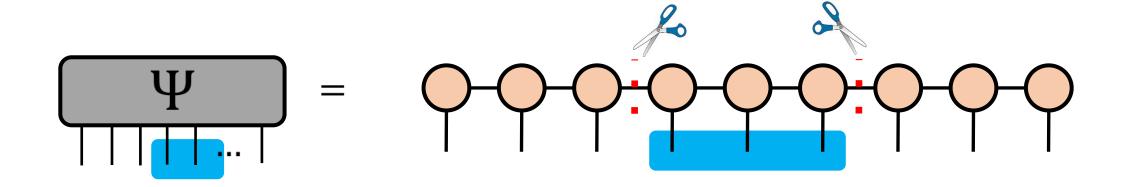
Rule for estimating upper bound of entanglement scaling from a TN

G. Evenbly and G. Vidal, J. Stat. Phys. 145, 891 (2011)

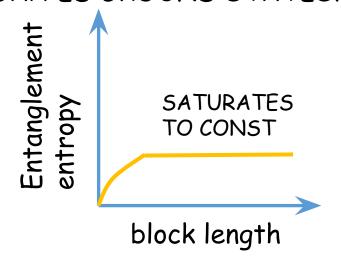
"Entanglement entropy & Minimal number of bond cuts"

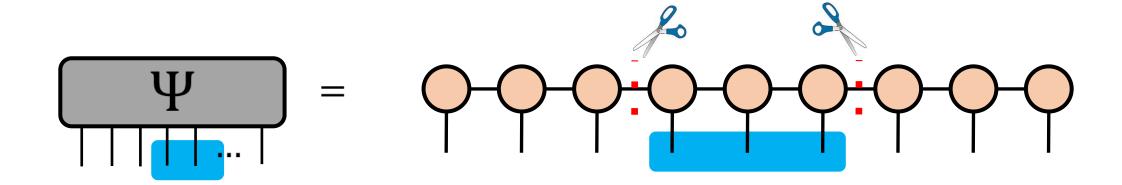




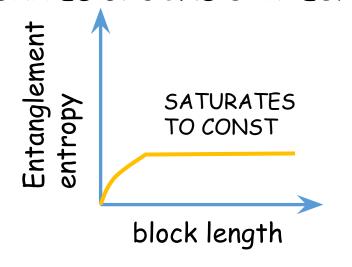


GAPPED GROUND STATES!



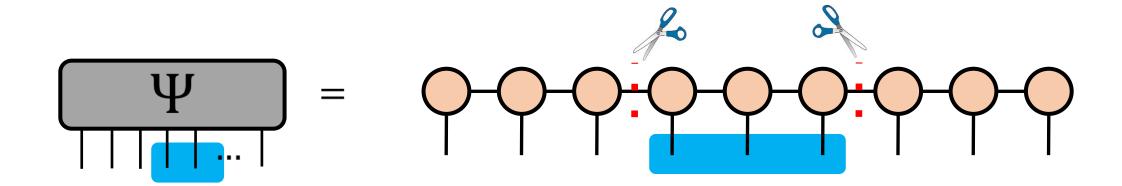


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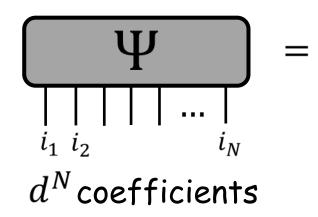


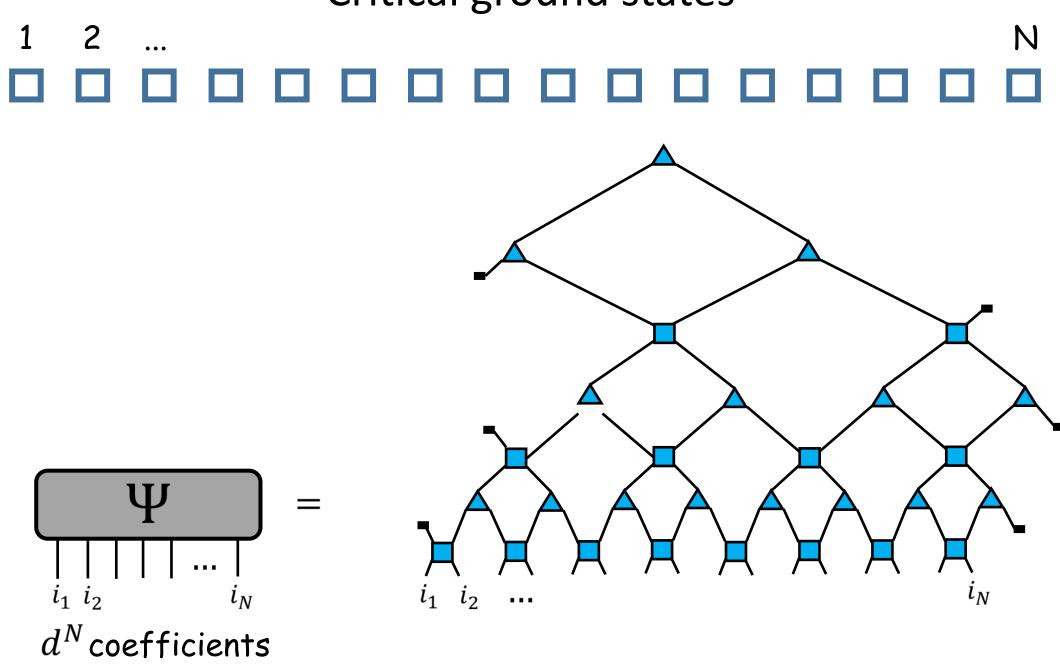
- MPS can accomodate the entanglement scaling of gapped ground states.
- Can use MPS as a variational ansatz for gapped ground states.
- Lots of supporting numerical evidence and theoretical results.

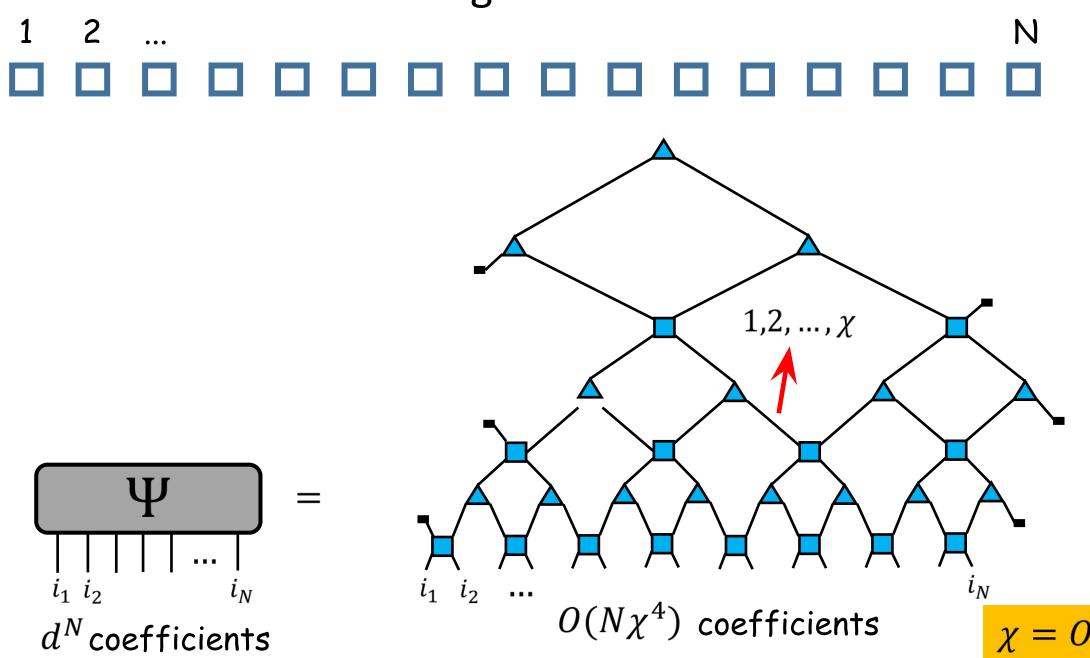
- We will be interested in critical ground states.
- Can we build a tensor network representation for critical ground states?
- o Lesson from the MPS: entanglement in a tensor network state depends on the geometry of the tensor network.

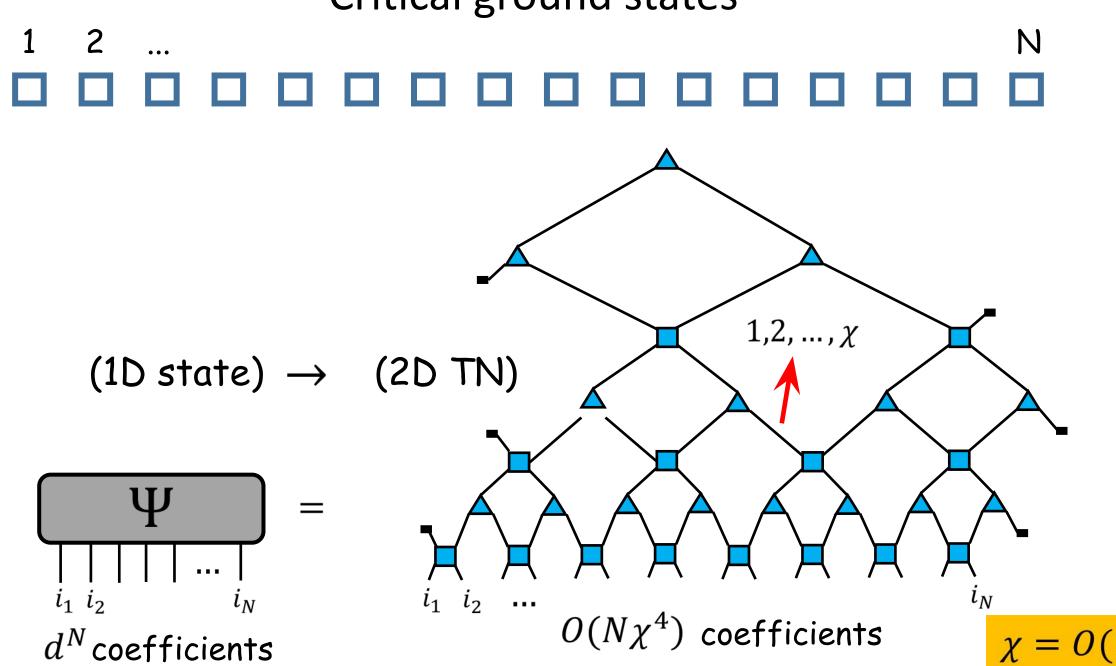


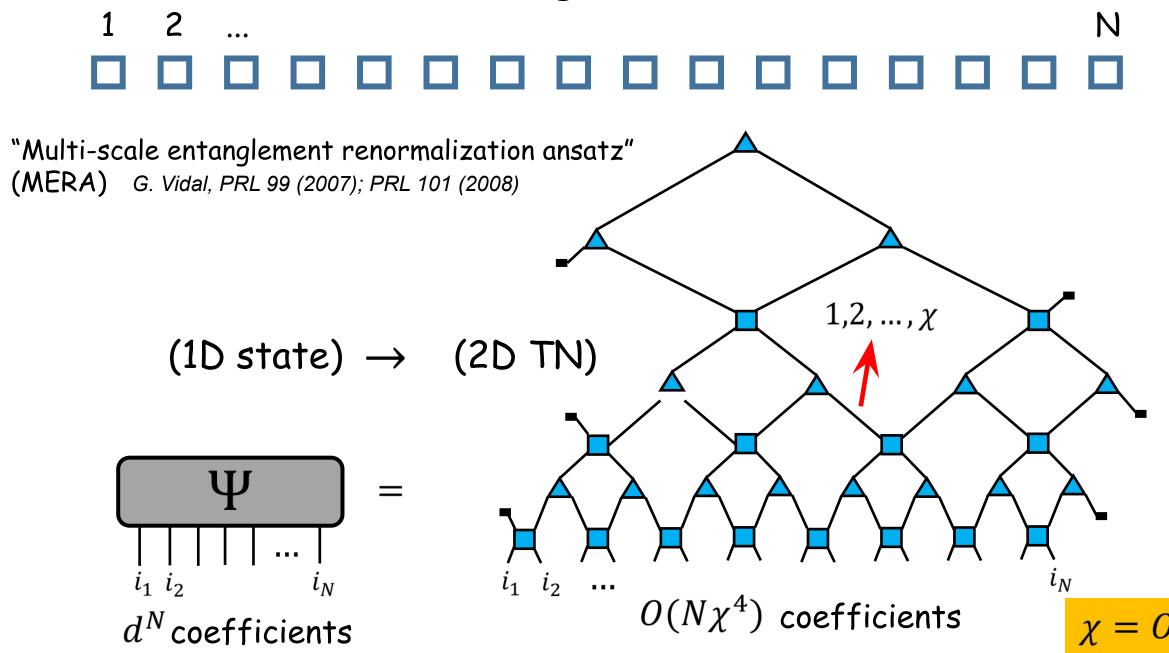




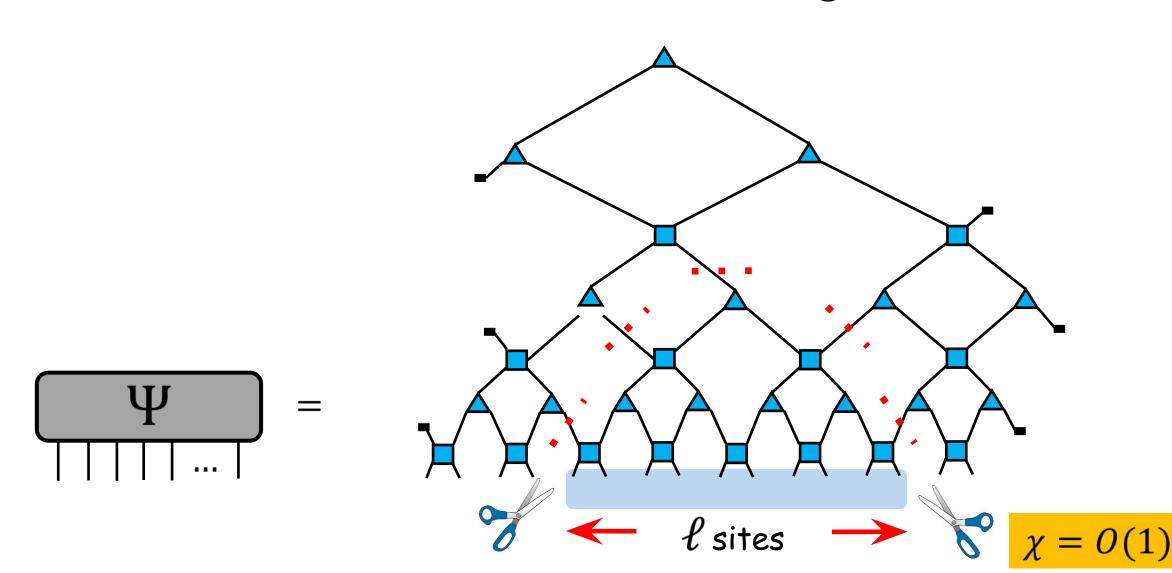






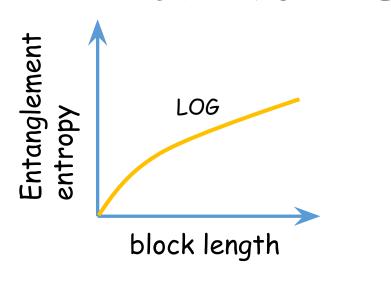


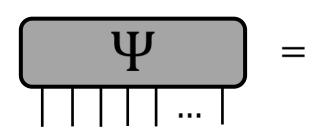
Minimal number of bond cuts $\propto \log \ell$

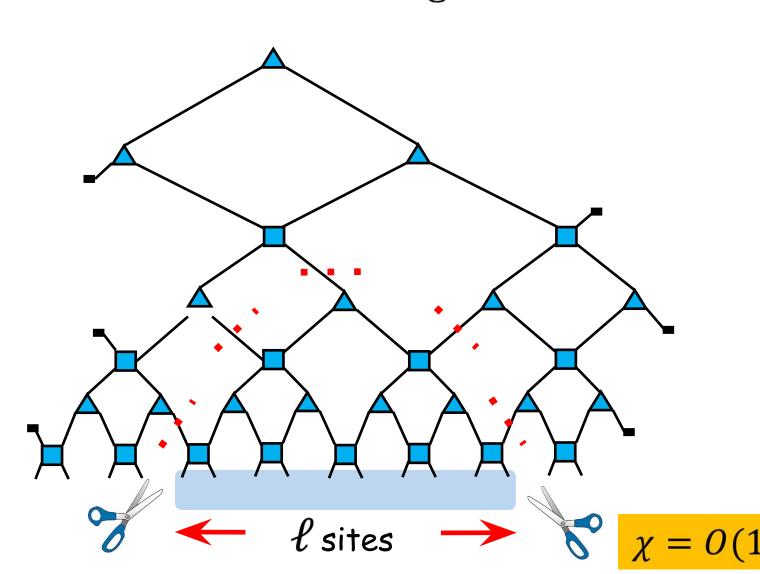


Minimal number of bond cuts $\propto \log \ell$

CRITICAL GROUND STATES

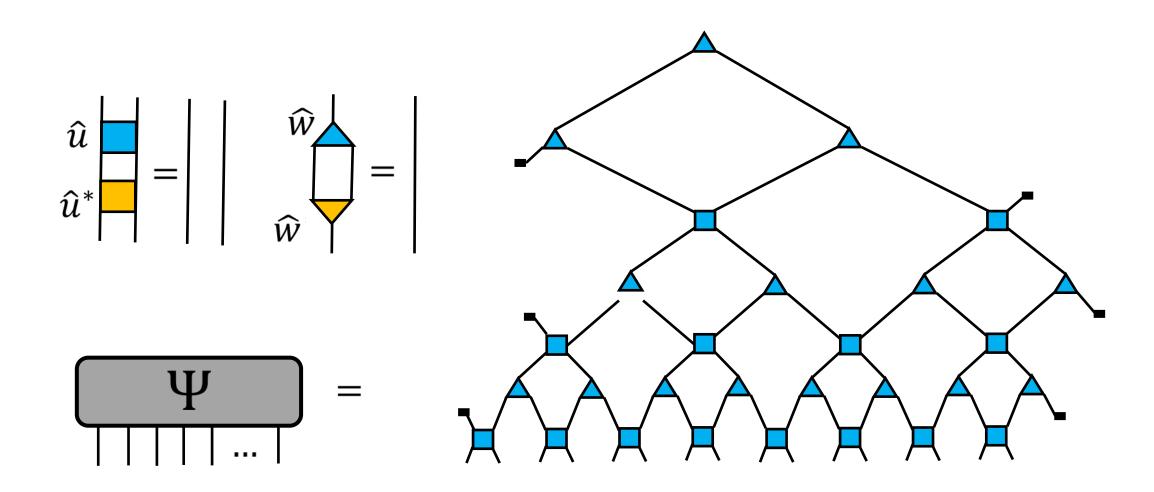


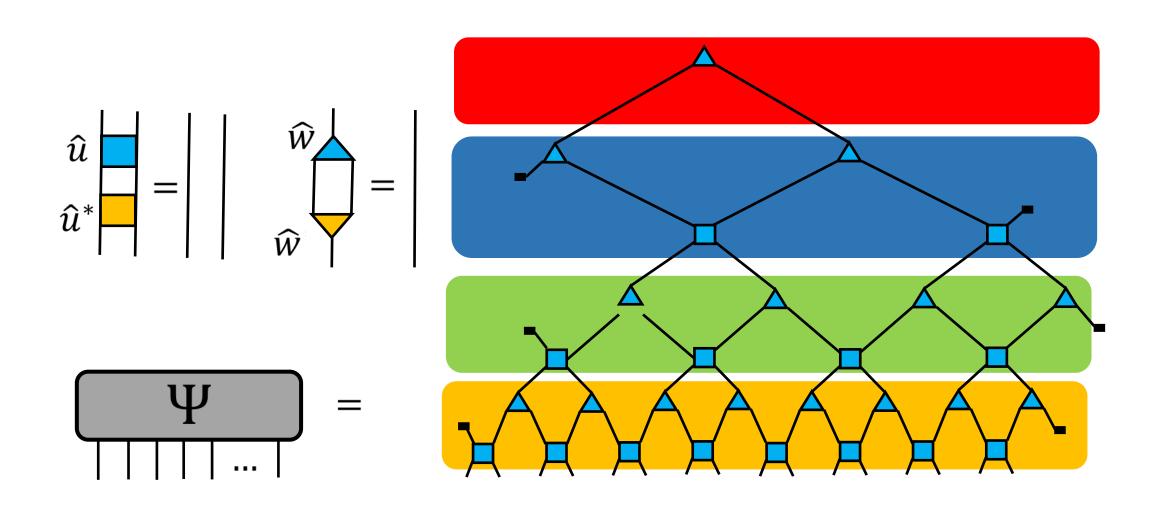


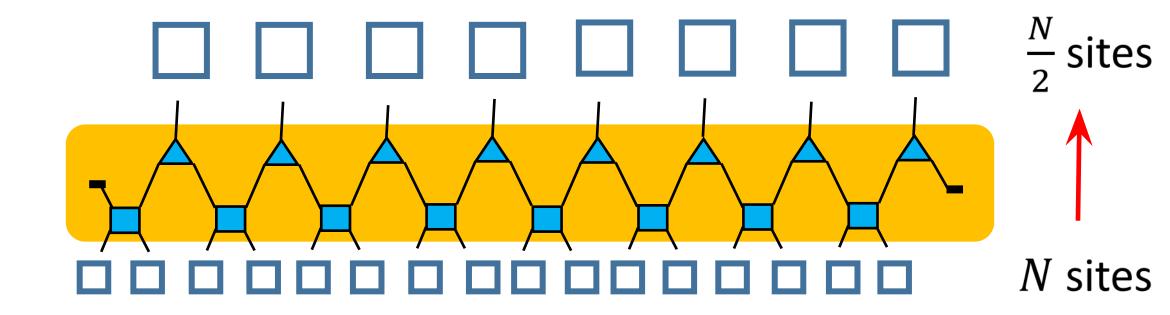


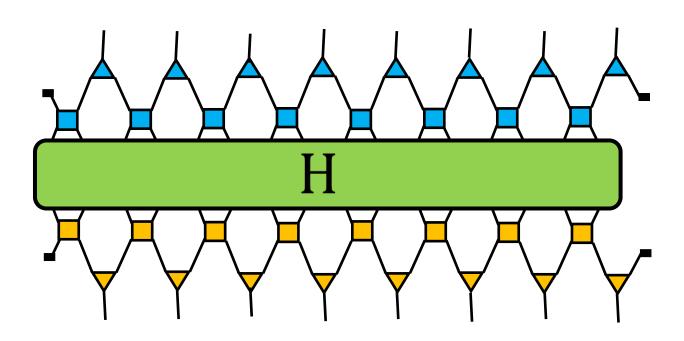
MERA as an efficient ansatz for critical ground states

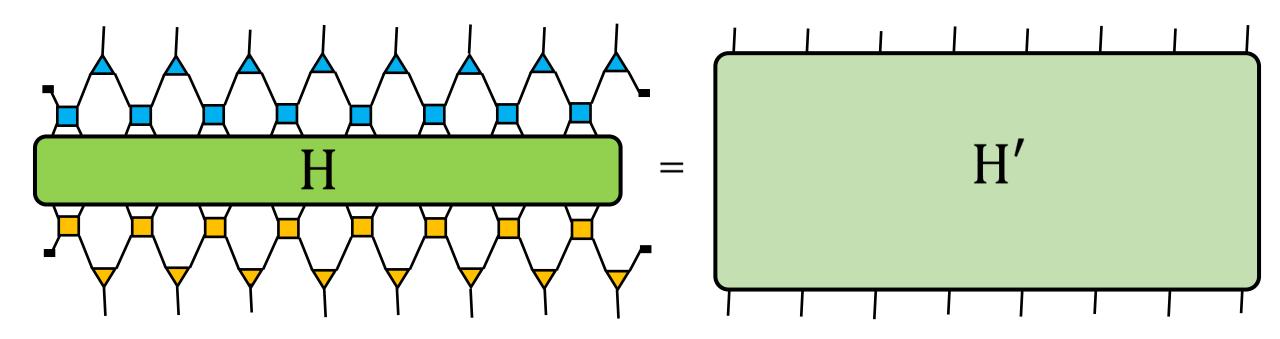
- Algorithms known for: critical $\widehat{H} \rightarrow$ ground state as MERA.
- Applied to several 1D quantum critical lattice models.
- Ground state properties can be efficiently computed.



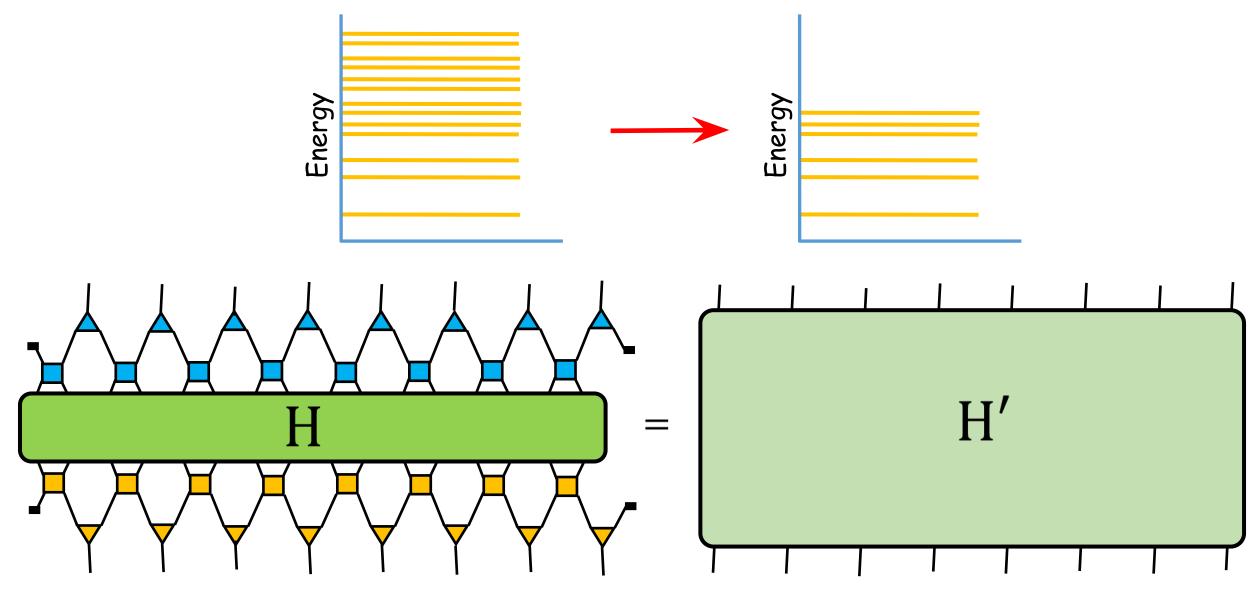




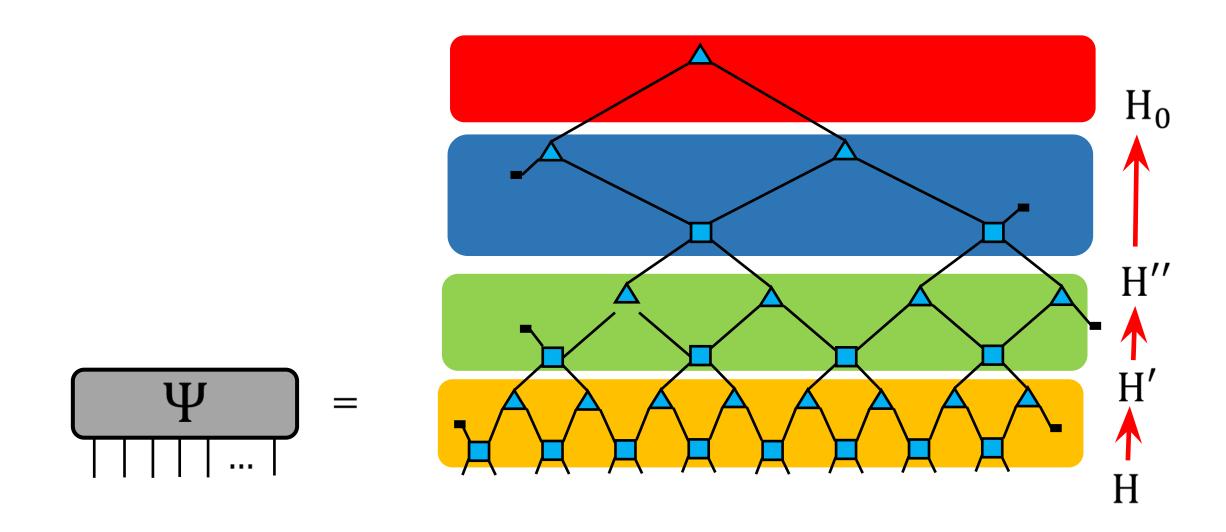


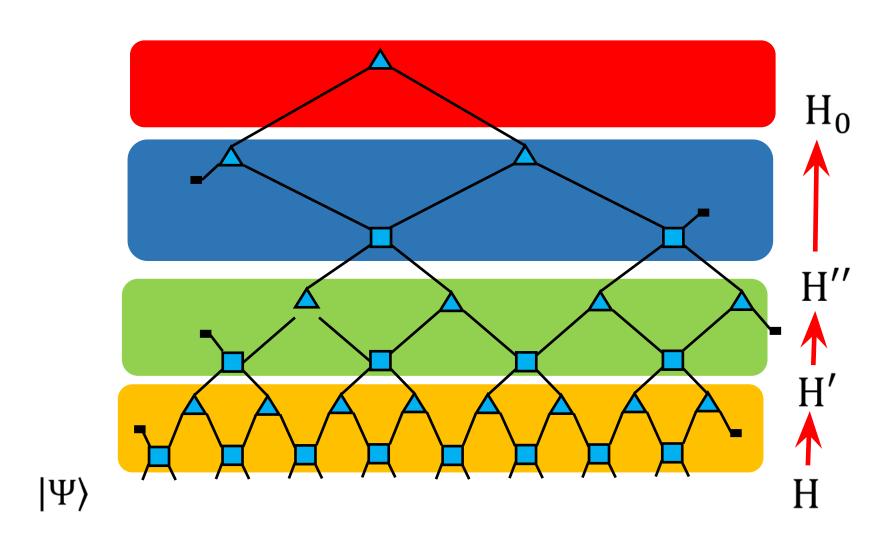


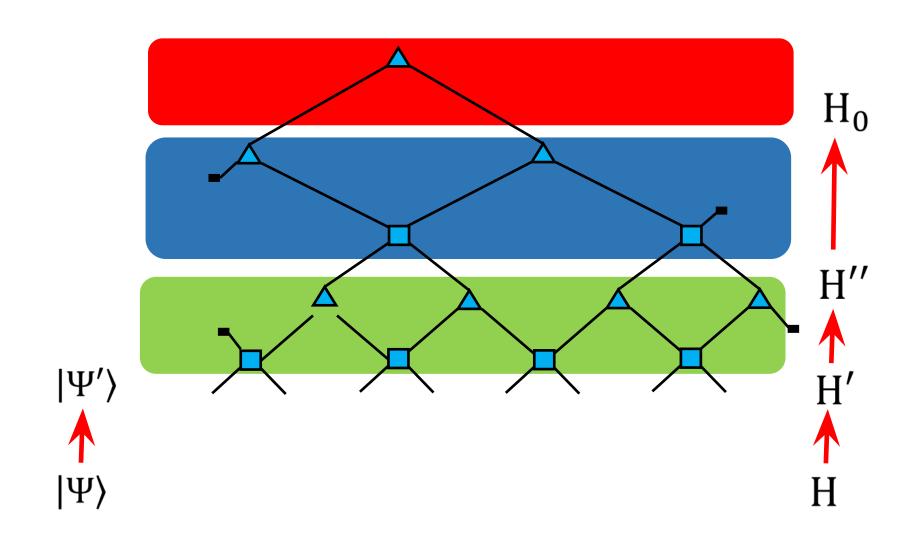
Single layer as a renormalization transformation

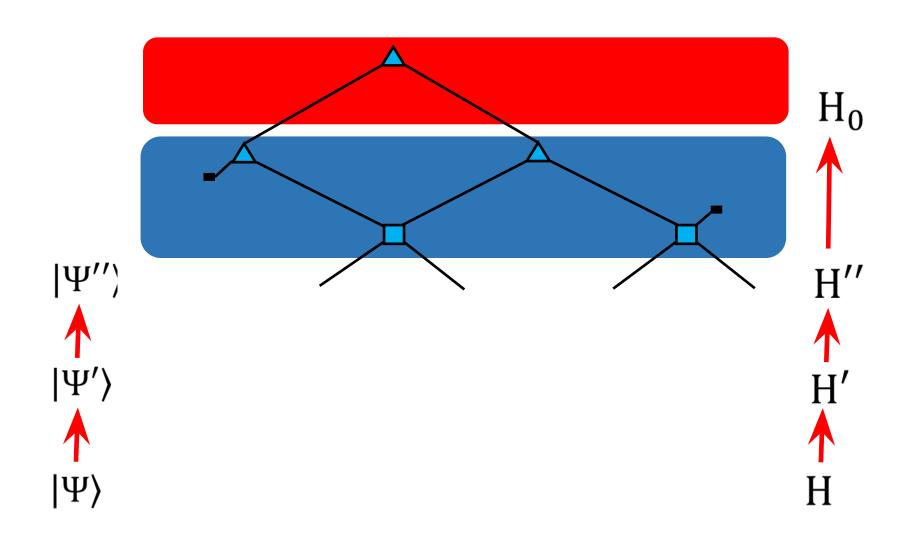


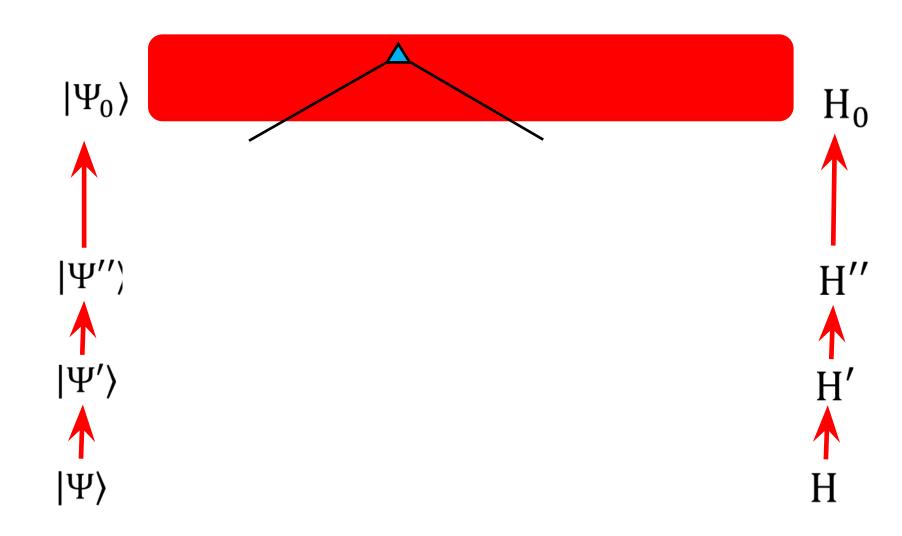
Renormalized Hamiltonian

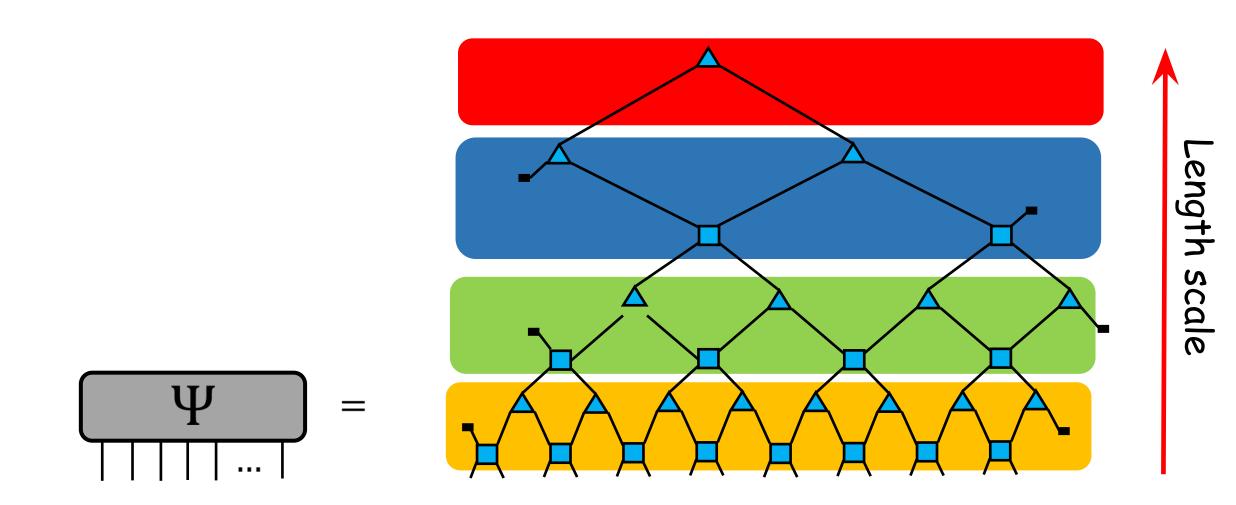




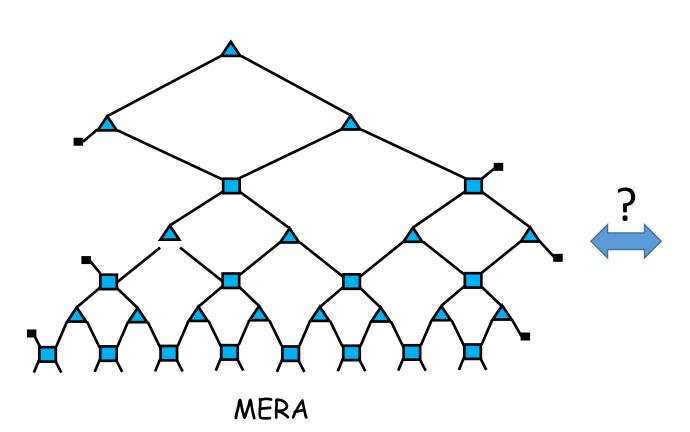


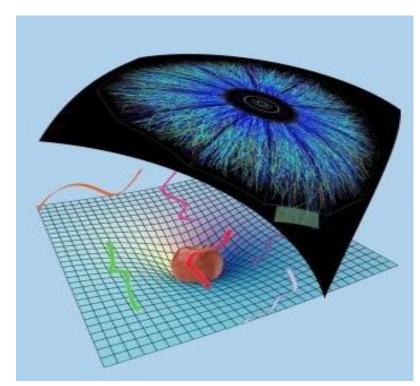






MERA and the AdS/CFT correspondence





Holographic correspondence

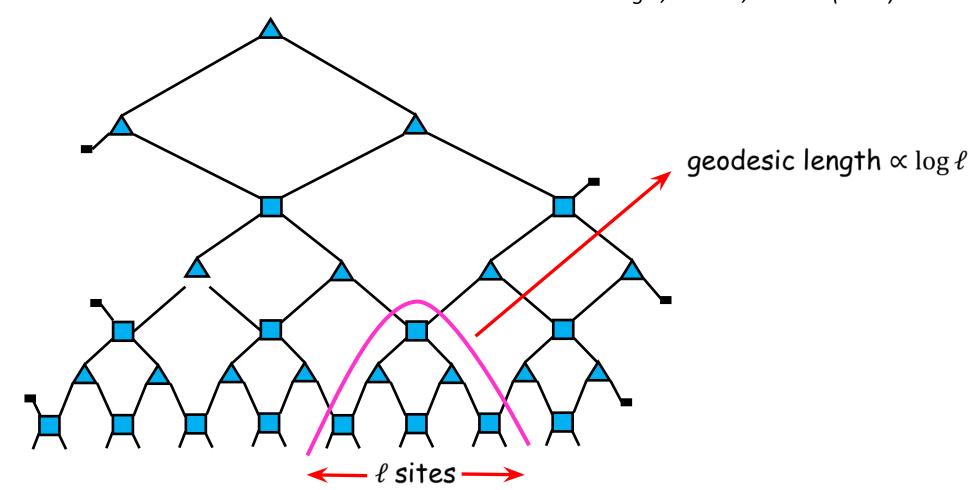
CFT side

- MERA is a "natural" ansatz for 1D critical ground states.
- Critical Hamiltonians are described by a CFT in the continuum

 At critical point, system becomes scale-invariant, critical ground state corresponds to the vacuum of the CFT.
- The properties of the CFT (central charge, primary fields, scaling dimensions etc.) can be extracted easily from the MERA tensors.
- Can incorporate scale-invariance exactly into the MERA tensor network.

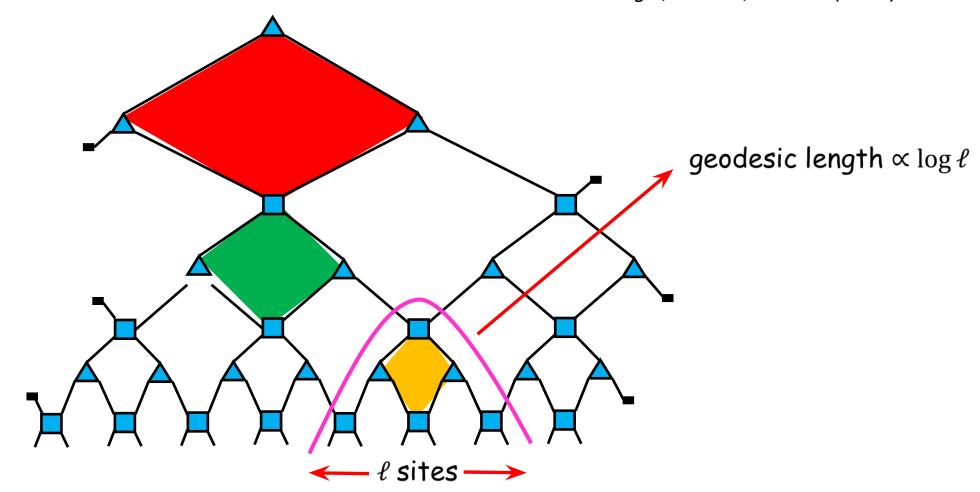
MERA tensor network has a hyperbolic geometry!

B. Swingle, PRD 86, 065007 (2012)

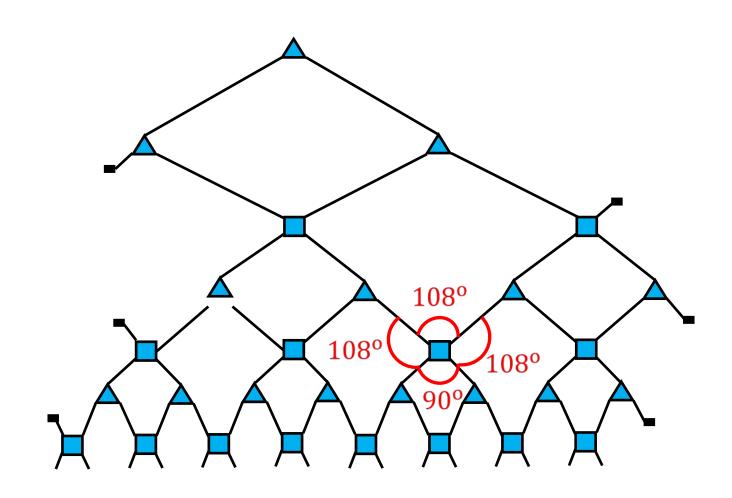


MERA tensor network has a hyperbolic geometry!

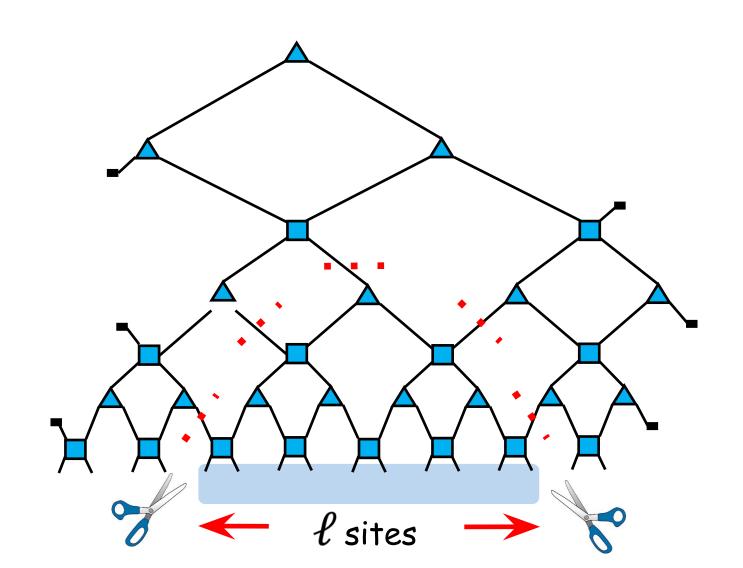
B. Swingle, PRD 86, 065007 (2012)



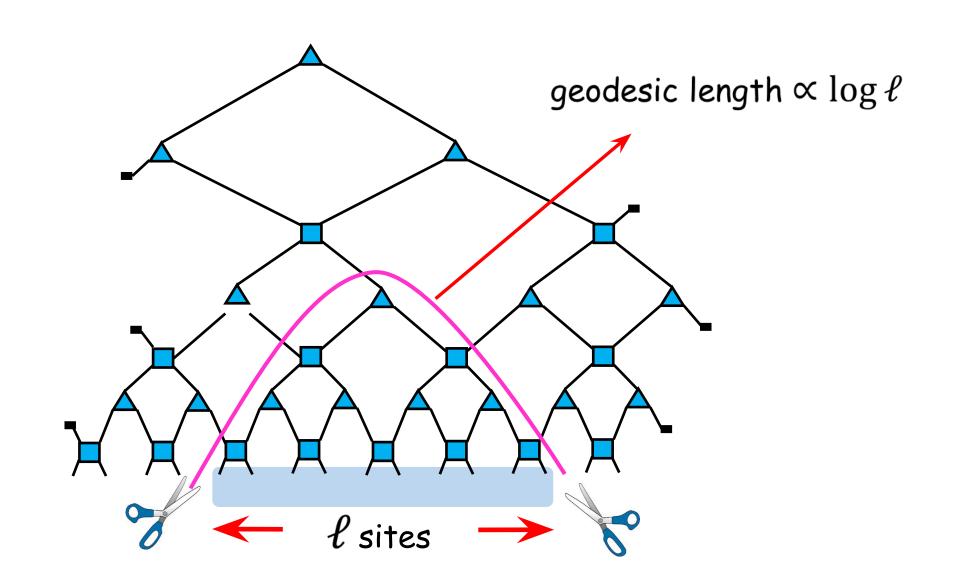
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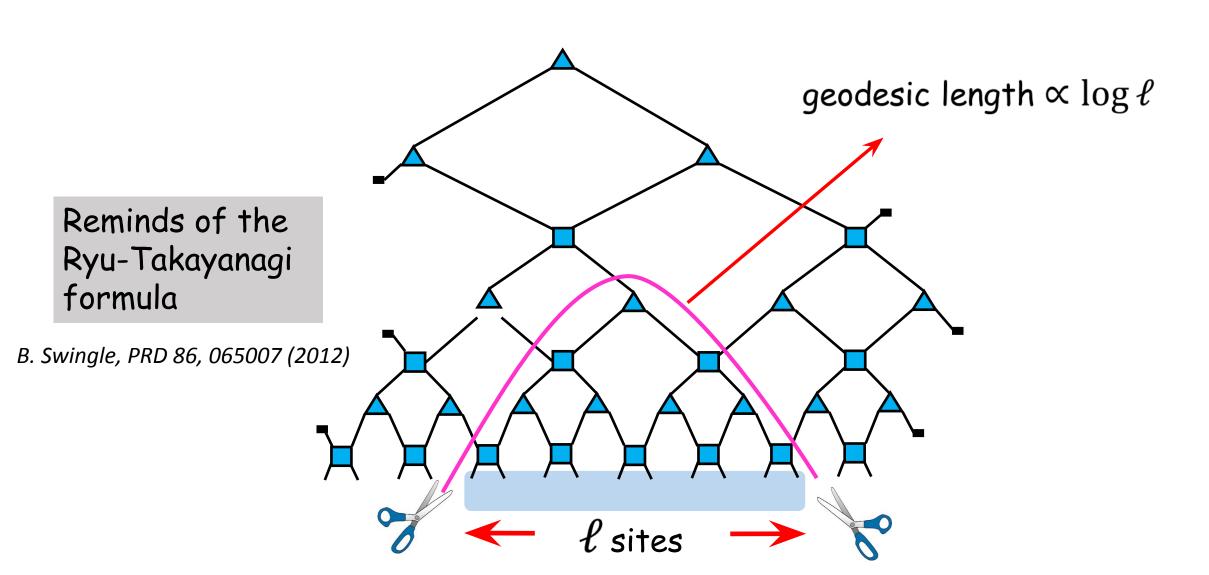
Minimal number of bond cuts $\propto \log \ell$

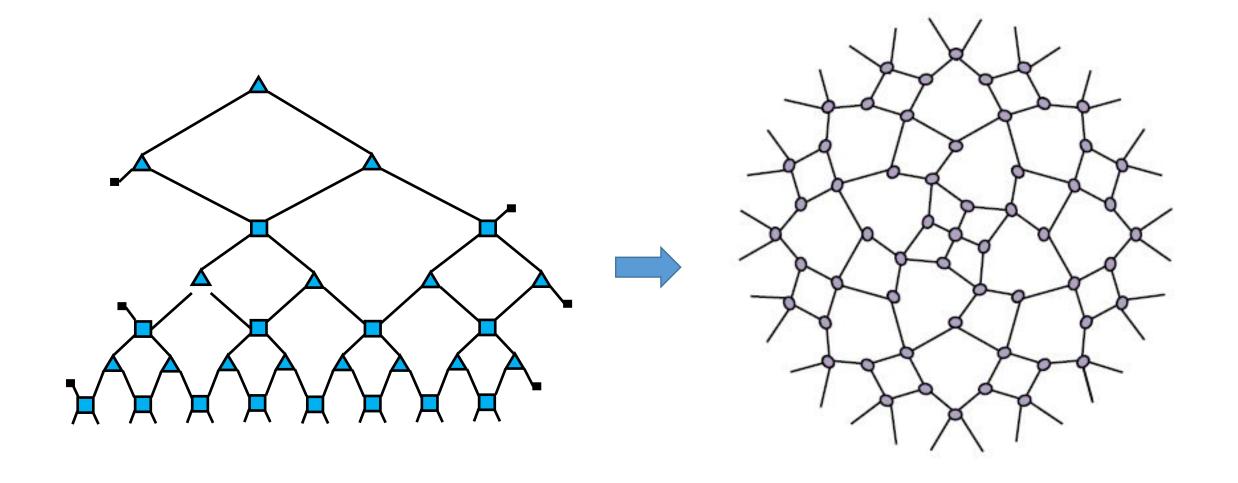


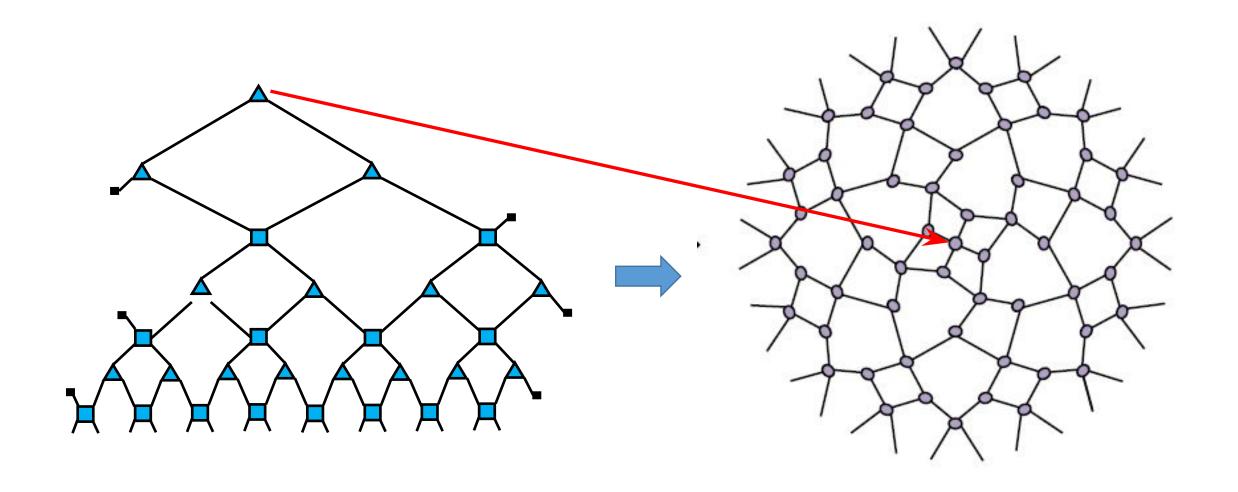
Block entropy identified with geodesics in the tensor network



Block entropy identified with geodesics in the tensor network



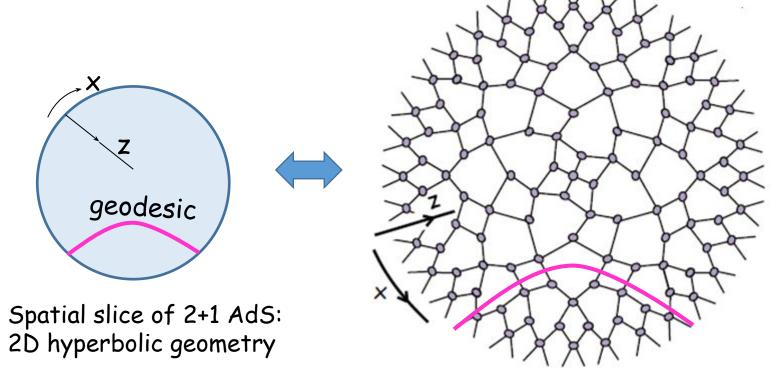




MERA/AdS conjecture B. Swingle, PRD 86, 065007 (2012)

MERA representation of a critical ground state can be regarded as a spatial slice of

discretised 2+1 AdS spacetime.



- Open indices correspond to sites of the lattice where the CFT lives.
- Extra dimension (z) corresponds to length scale.
- The MERA tensors that describe a critical ground state also encode a 2D gravity description.

MERA/AdS conjecture: what we have so far

- O Basic similarity: MERA is a 2D representation of a 1D ground state and has a hyperbolic geometry, and the extra dimension in the tensor network corresponds to length scale of the boundary system.
- Some non-trivial features of the AdS/CFT correspondence are present in the MERA:
- Analog of the Ryu-Takayanagi formula, B. Swingle, PRD 86, 065007 (2012)
- Thermal state/black hole correspondence, T. Hartman and J. Maldacena arXiv:1303.1080
- Bulk/boundary symmetries etc. s. singh and G. Vidal PRB (R) 88 121108 (2013)
 - B. Swingle, arXiv:1209.3304.
 - J. Molina-Vilaplana and P. Sodano, JFEP 10, 11 (2011), arXiv:1108.1277.
 - H. Matsueda, M. Ishihara, and Y. Hashizume, Phys. Rev. D 87, 066002 (2013), arXiv:1208.0206.
 - M. Nozaki, S. Ryu, and T. Takayanagi, JHEP 10, 193 (2012), arXiv:1208.3469.
 - C. Beny, New J. Phys. 15, 023020 (2013), arXiv:1110.4872.
 - B. Czech, L. Lamprou, S. McCandlish, J. Sully, arXiv:1512.01548, SU-ITP-15/18, SLAC-PUB-16292.

Open questions

- o Is the MERA/AdS conjecture true?
- Mostly qualitative observations. Can we make the conjecture more precise?
- We understand how MERA encodes the boundary critical ground state. How does the tensor network encode the dual bulk description?

One potential approach SS, arXiv:1701.04778

SS, N. McMahon, G. Brennen, arXiv:1702.00392

 Does the MERA encode a more general bulk/boundary correspondence, since in principle, one can represent the ground state of any CFT as a MERA (possibly including CFTs that are not known/expected to have a gravity dual).

Outlook

- If true, the MERA/AdS conjecture implies that a physical quantum gravity state is encoded in the MERA somehow, without even knowing the quantum gravity theory (but this in fact the promise of holography --- the bulk gravity is just a QFT in one lower dimension).
- Thus, MERA may be a toy model to explore the holographic correspondence.
- In reverse, insights from holography have already been applied to design better tensor networks for simulations of strongly correlated many-body systems e.g., the thermal state MERA, branching MERA etc.

THANKS!