MERA, symmetries and AdS/CFT

Sukhi Singh (EQuS, Macquarie)























Hamiltonian:

- A) Local
- B) Critical/Gapless $(N \to \infty)$ -> described by a CFT

e.g.
$$\hat{H} = \sum_{k=1}^{N} \left(\sigma_X^k \sigma_X^{k+1} + \sigma_Y^k \sigma_Y^{k+1} + \sigma_Z^k \sigma_Z^{k+1} \right)$$
 (Heisenberg model)

























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Ground state:
$$|\Psi_{gs}\rangle = \sum_{i_1i_2...i_N} \hat{\Psi}_{i_1i_2...i_N} |i_1\rangle |i_2\rangle...|i_N\rangle$$

Locality & criticality => Limited entanglement in GS

$$S_l = \frac{c}{3} \log l$$

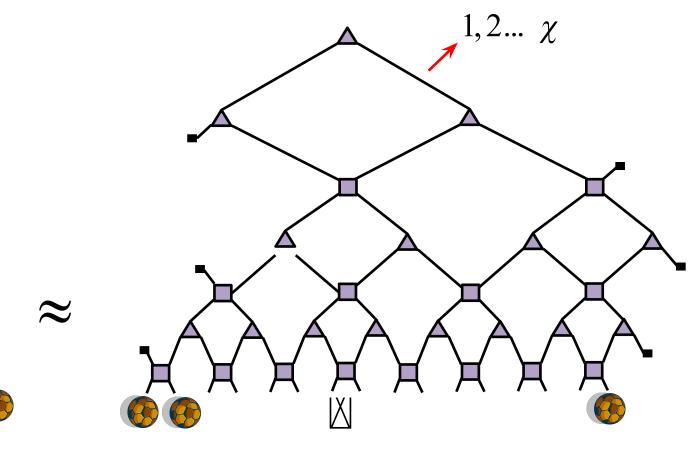
c: central charge of the underlying CFT

$$S_l = -\operatorname{tr}(\rho_l \log \rho_l)$$

$$\rho_l = \operatorname{tr}_{\bar{l}} \left(\left| \Psi_{gs} \right\rangle \left\langle \Psi_{gs} \right| \right)$$

MERA

$$\left|\Psi_{gs}\right\rangle = \sum_{i_1 i_2 \dots i_N} \hat{\Psi}_{i_1 i_2 \dots i_N} \left|i_1\right\rangle \left|i_2\right\rangle \dots \left|i_N\right\rangle$$



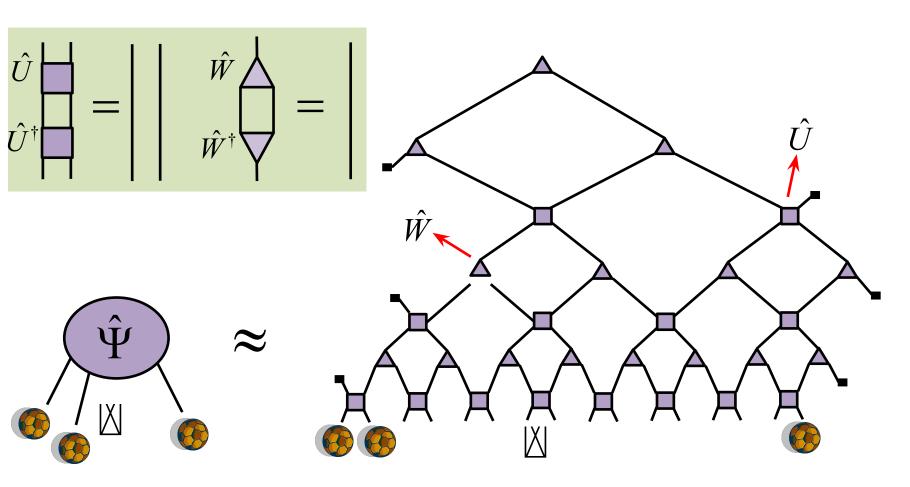
 d^L coefficients

 $\hat{\Psi}$

 $O(L\chi^4)$ coefficients

Disentanglers & isometries

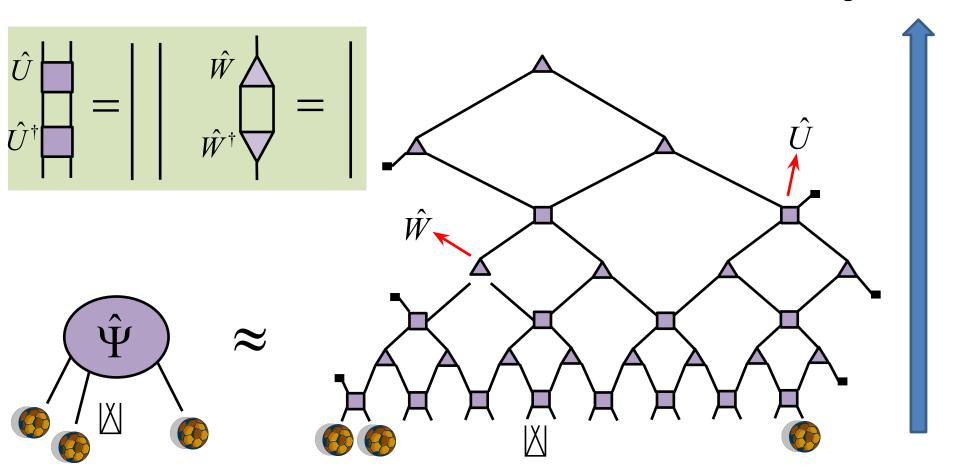
$$\left|\Psi_{gs}\right\rangle = \sum_{i_1 i_2 \dots i_N} \hat{\Psi}_{i_1 i_2 \dots i_N} \left|i_1\right\rangle \left|i_2\right\rangle \dots \left|i_N\right\rangle$$



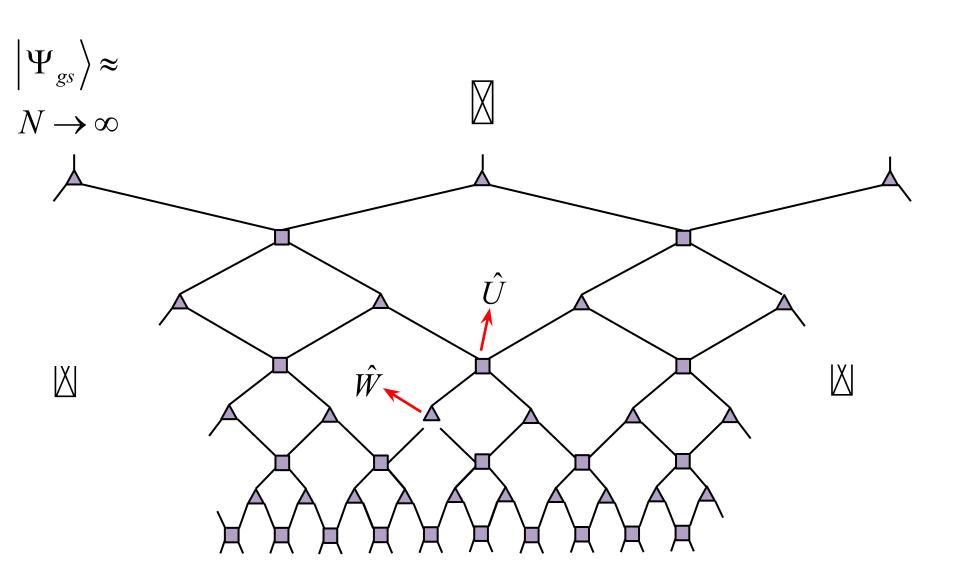
RG transformations

$$\left|\Psi_{gs}\right\rangle = \sum_{i_1 i_2 \dots i_N} \hat{\Psi}_{i_1 i_2 \dots i_N} \left|i_1\right\rangle \left|i_2\right\rangle \dots \left|i_N\right\rangle$$

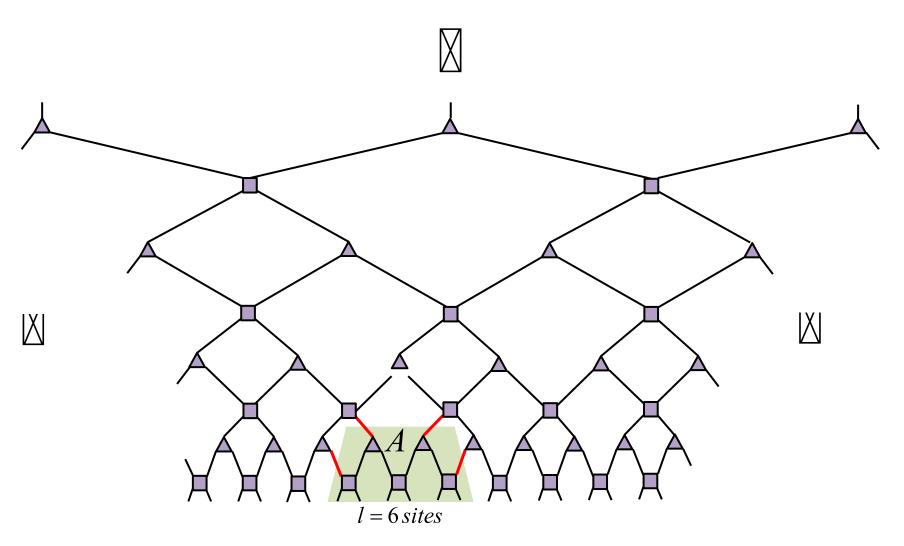
Length Scale



Scale-invariant MERA



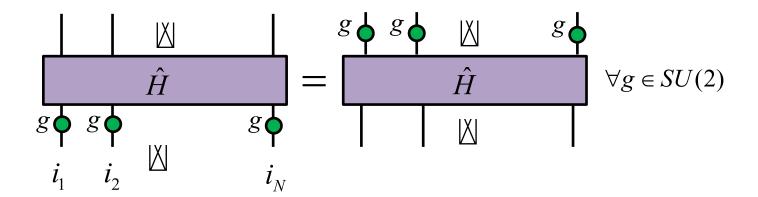
Entanglement from geometry in MERA



min number of connecting bonds $\approx \log l$

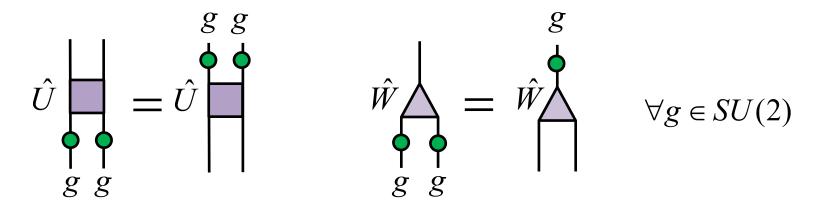
Global SU(2) symmetry

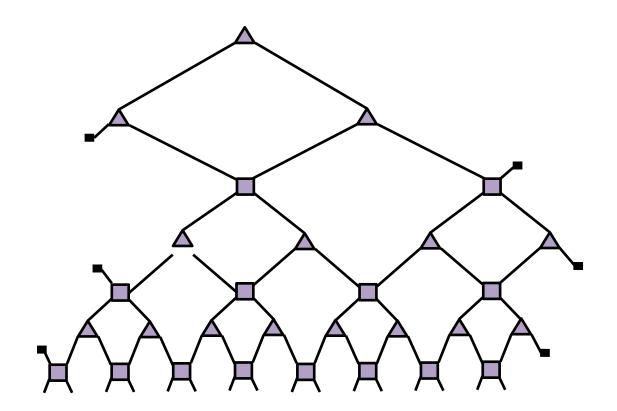
1) Hamiltonian:



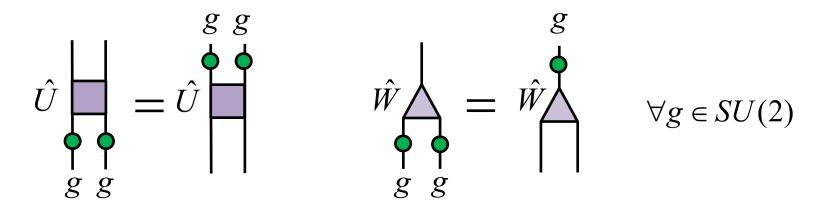
2) Ground state: total spin 0

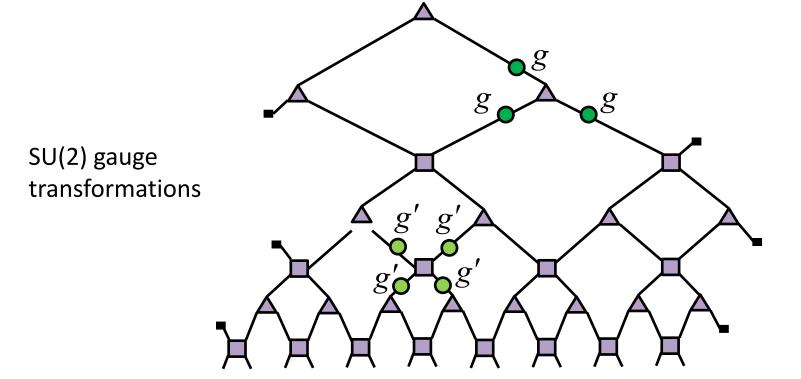
SU(2)-invariant MERA



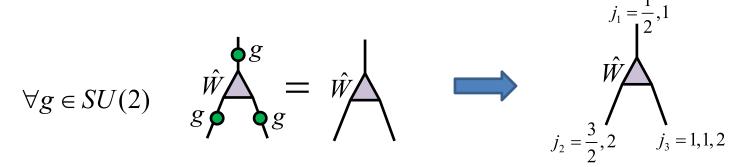


SU(2)-invariant MERA

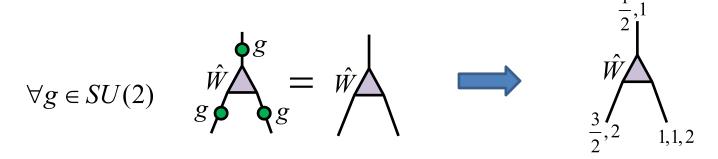




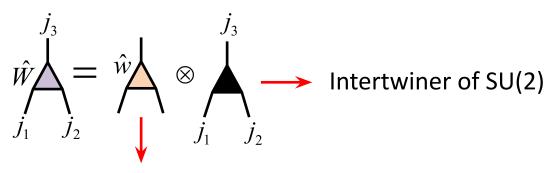
$$\forall g \in SU(2) \qquad \mathring{g} = \mathring{w}$$



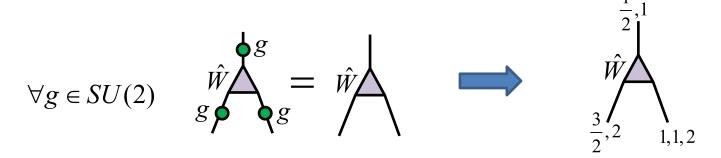
Spin conservation $|j_1 - j_2| \le j_3 \le j_1 + j_2$



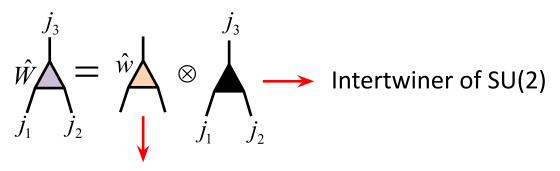
Spin conservation $|j_1 - j_2| \le j_3 \le j_1 + j_2$



Part not fixed by symmetry

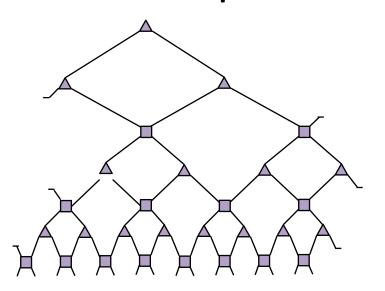


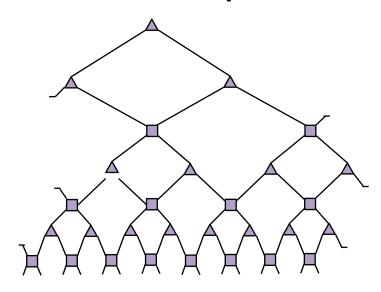
Spin conservation $|j_1 - j_2| \le j_3 \le j_1 + j_2$

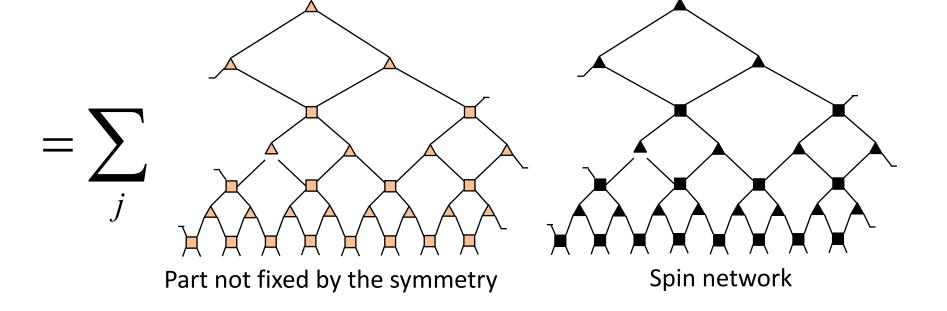


Part not fixed by symmetry

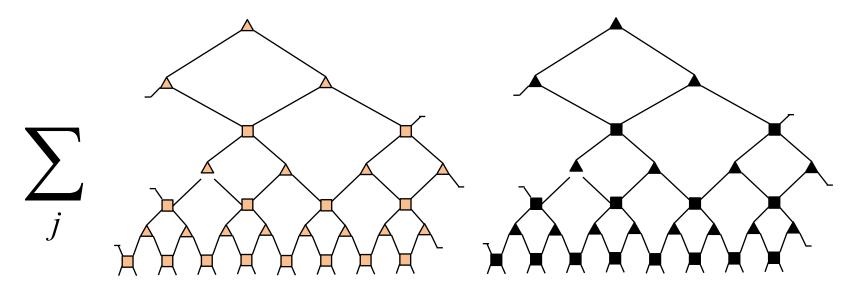
$$\hat{J}_{1} \qquad \hat{J}_{2} \qquad \hat{J}_{2} \qquad \hat{J}_{3} \qquad \hat{J}_{4} \qquad \hat{J}_{3} \qquad \hat{J}_{4} \qquad \hat{J}_{3} \qquad \hat{J}_{4}$$







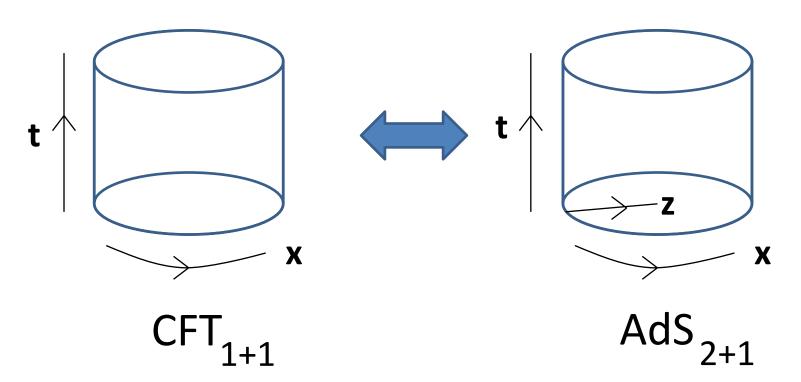
Computational advantages



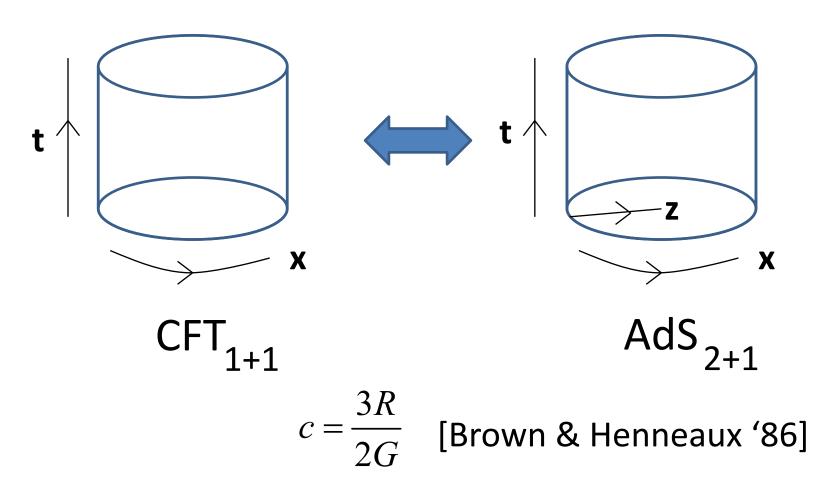
- 1) Compactified description
- 2) Faster tensor network manipulations
- 3) Can target specific symmetry sectors

MERA and AdS/CFT

AdS/CFT correspondence

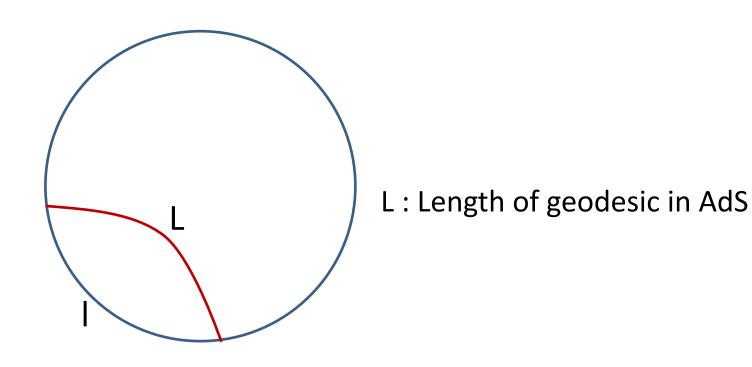


AdS/CFT correspondence



Small c => quantum geometry in the bulk?
Gravity dual of the critical ising model, c=1/2
[Castro et. al, PRD '12]

Gravitational dual of entanglement entropy



$$L \approx \frac{R}{2G} \log l$$
$$S_l \approx L$$

[Ryu & Takayanagi PRL '06]

MERA and AdS/CFT

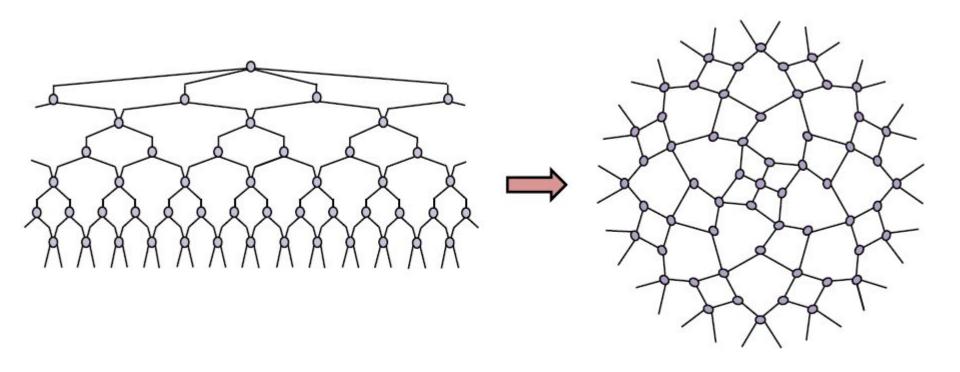


Fig. source: Evenbly, Vidal (2011)

MERA and AdS/CFT

Brian Swingle 2009

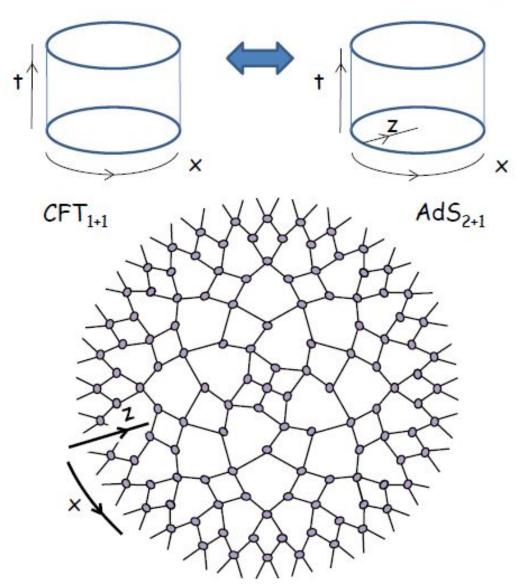
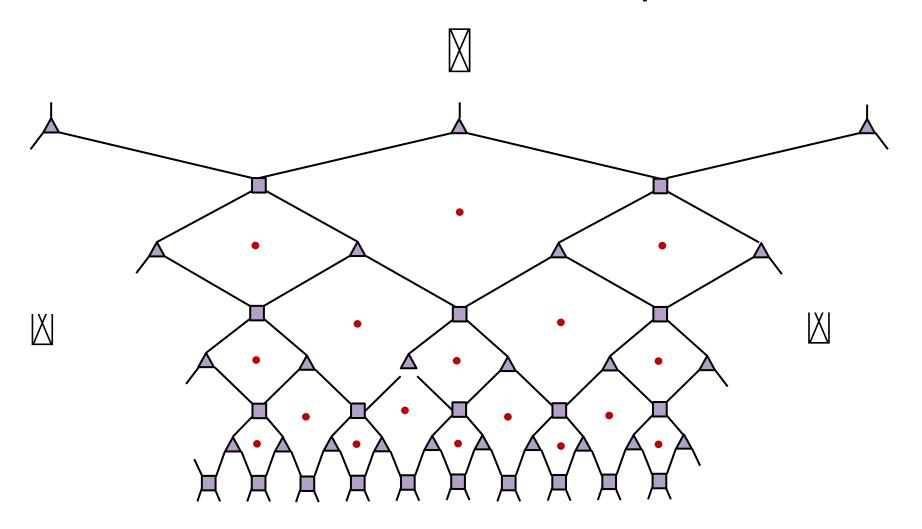
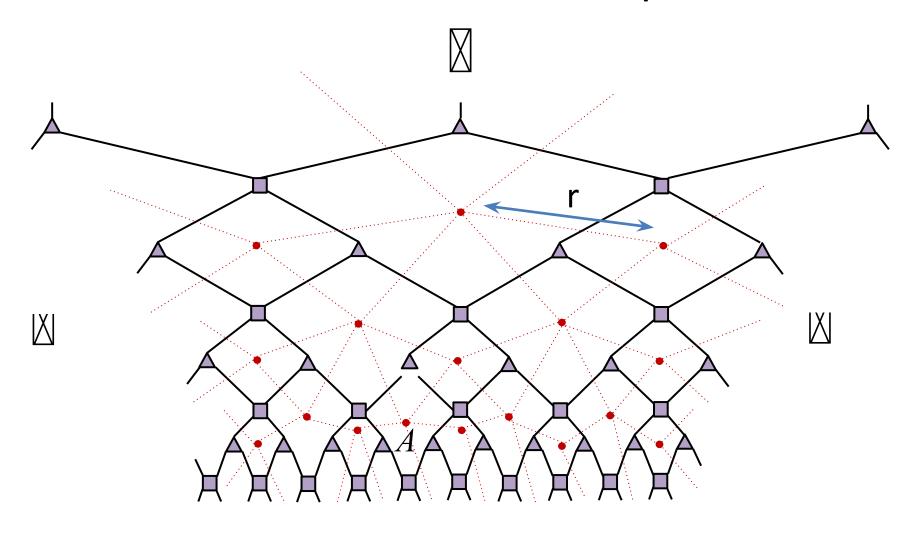


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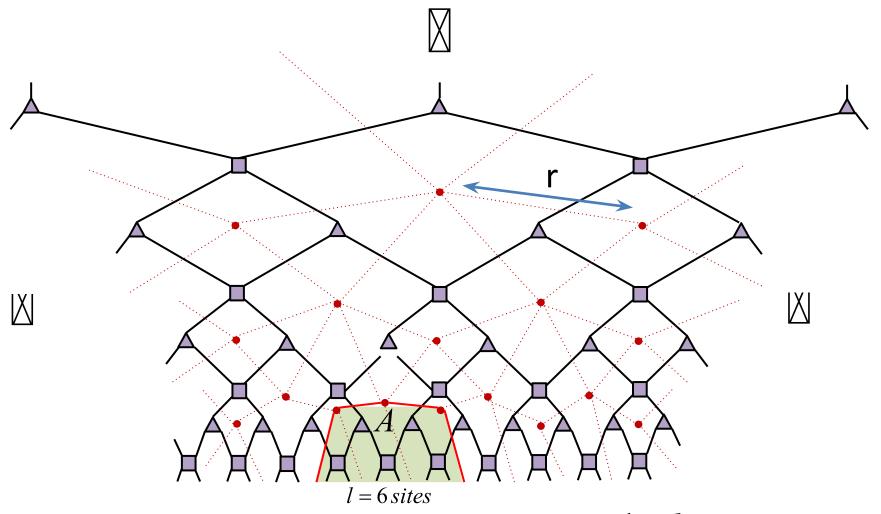
MERA as discrete AdS space



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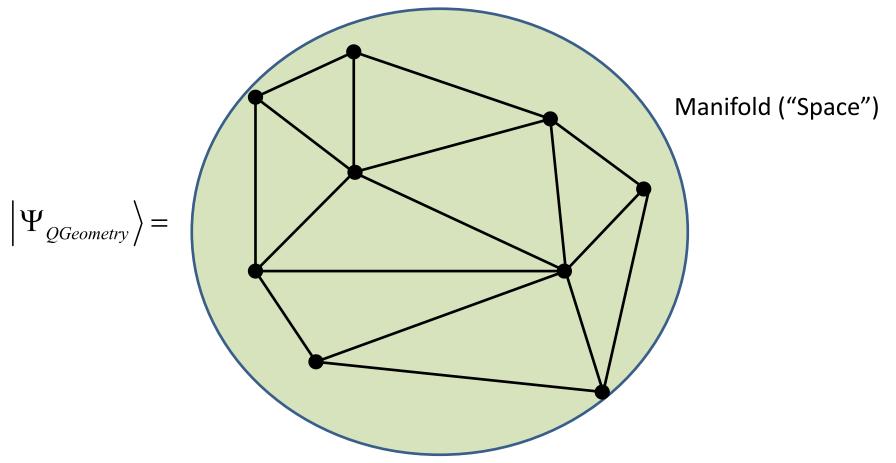
MERA as discrete AdS space



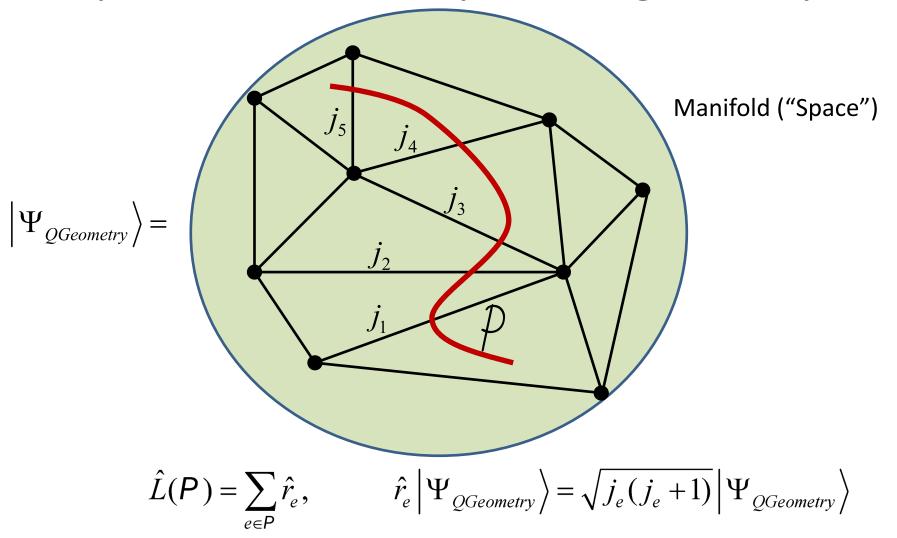
Length of geodesic in the dual graph = $r \log l$

$$S_l \approx r \log l?$$
 $r = \frac{c}{3}?$

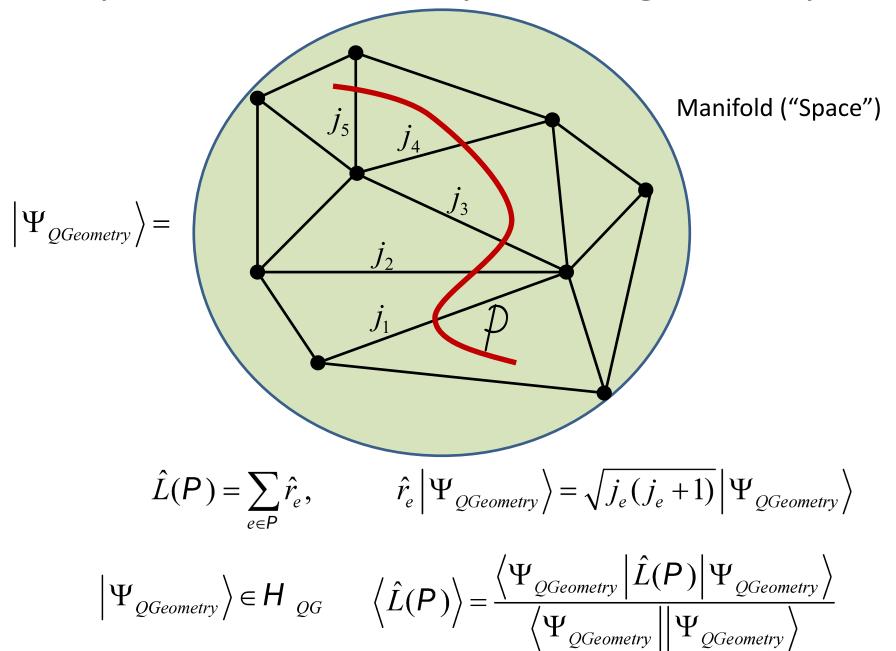
Spin networks and quantum geometry

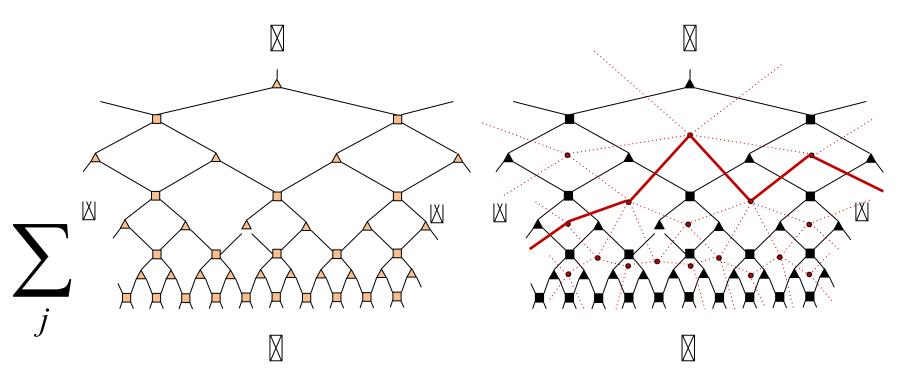


Spin networks and quantum geometry

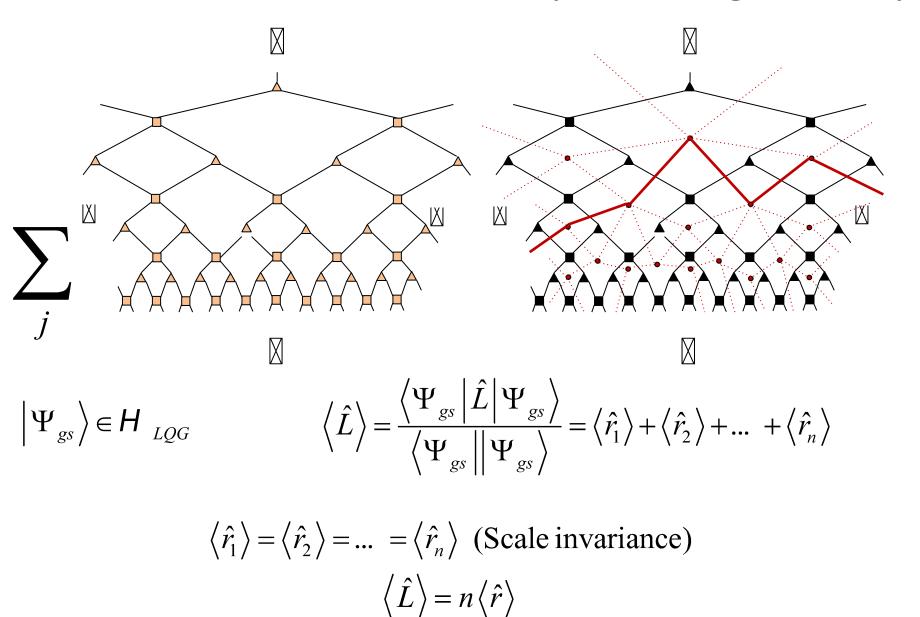


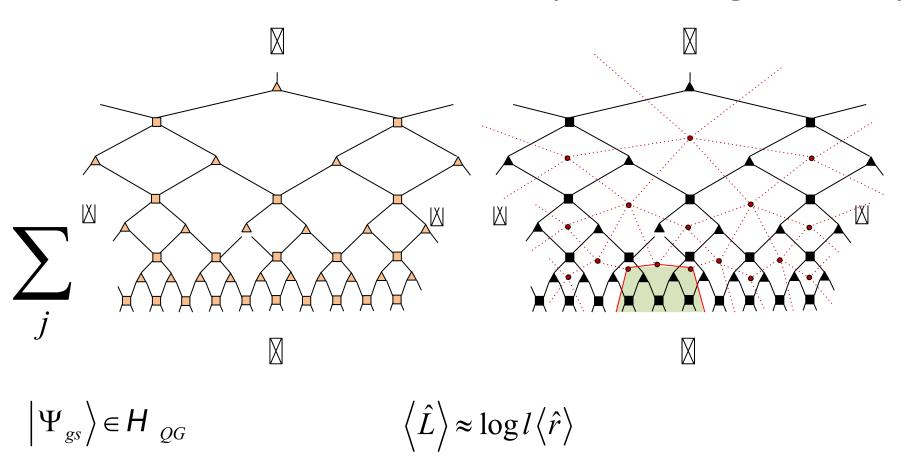
Spin networks and quantum geometry

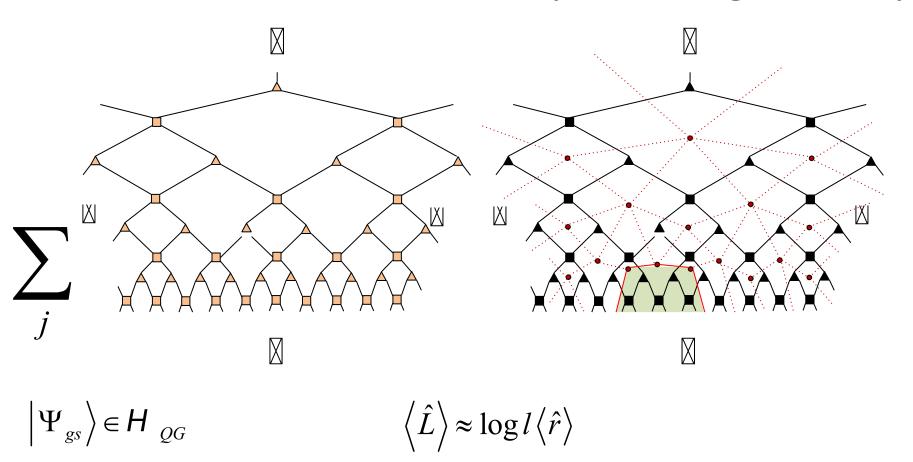




$$|\Psi_{gs}\rangle\in H_{LQG}$$







$$S_l \approx \langle \hat{L} \rangle$$
? $\frac{\langle \Psi_{gs} | \hat{r} | \Psi_{gs} \rangle}{\langle \Psi_{gs} | | \Psi_{gs} \rangle} \approx c$?

Summary

SU(2)-invariant MERA as a realization of AdS/CFT?

$$\left|\Psi_{gs}\right\rangle \in V^{\otimes N} \iff \left|\Psi_{gs}\right\rangle \in \mathcal{H}_{LQG}$$

- A) Have a description of bulk geometry
- B) Quantum geometry, expected for small c?
- C) Global symmetry Gauge symmetry

Can explore:

- A) AdS/CFT dictionary?
- B) other geometric quantities? holonomy, curvature
- C) different gauge groups? e.g. SU(2)_k

Thanks!