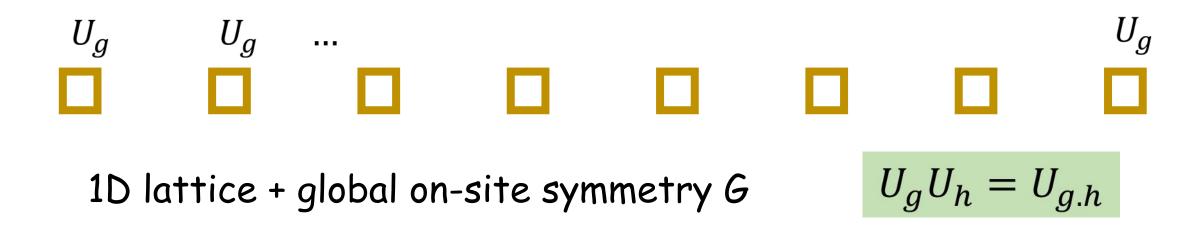
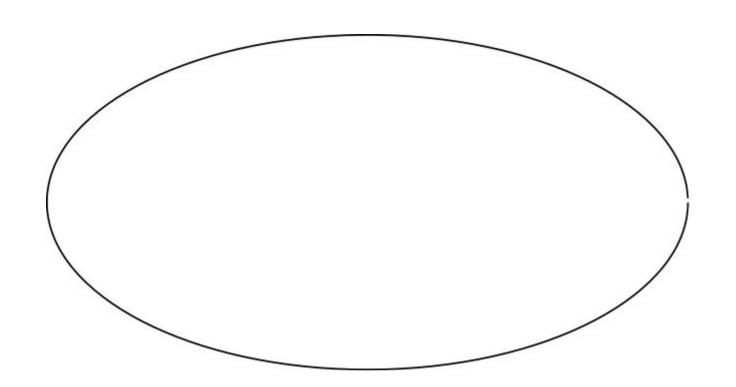
Tensor networks as tools for classifying 1D quantum phases of matter

Sukhbinder Singh

Max-Planck Institute for Gravitational Physics, Potsdam (Gravity, Quantum Fields, and Information Group)

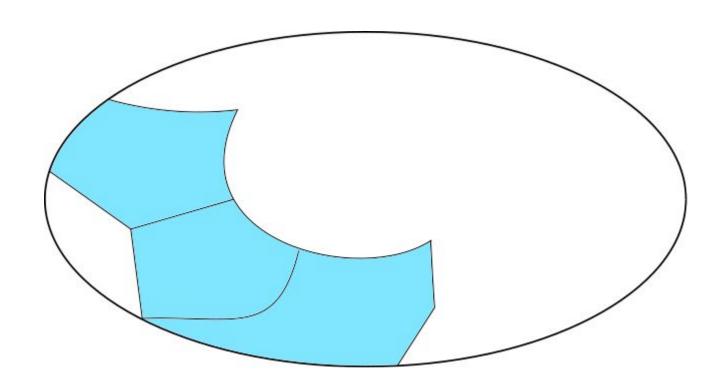
Based on: SS, N McMahon, G Brennen, 1812.08500 PRB 99 195139 (2019)





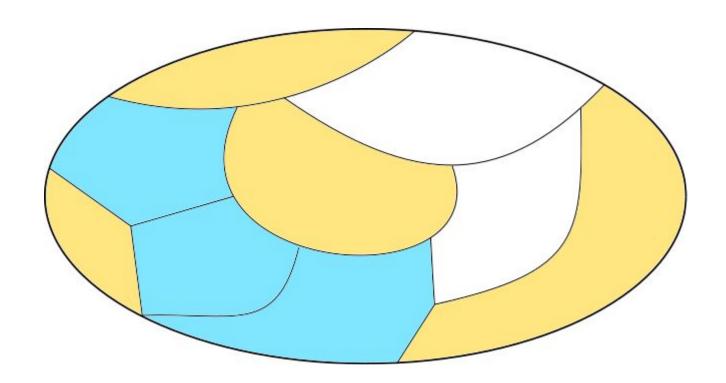
Space of local Hamiltonians with symmetry G

symmetry broken phase



Space of local Hamiltonians with symmetry G

- symmetry broken phase
- symmetry protected phase



Space of local Hamiltonians with symmetry G

Projective representations

$$V_g V_h = e^{i\omega(g,h)} V_{g,h}$$

Projective representations are defined only up to a phase factor

$$V_g \leftrightarrow e^{i\phi_g}V_g$$

Projective representations

$$V_g V_h = e^{i\omega(g,h)} V_{g,h}$$

Projective representations are defined only up to a phase factor

$$V_g \leftrightarrow e^{i\phi_g}V_g$$

Projective representations are equivalent under

$$\omega(g,h) = \omega(g,h) + \phi_g + \phi_h - \phi_{g,h} \bmod 2\pi$$

Projective representations

$$V_g V_h = e^{i\omega(g,h)} V_{g,h}$$

Projective representations are defined only up to a phase factor

$$V_g \leftrightarrow e^{i\phi_g}V_g$$

Projective representations are equivalent under

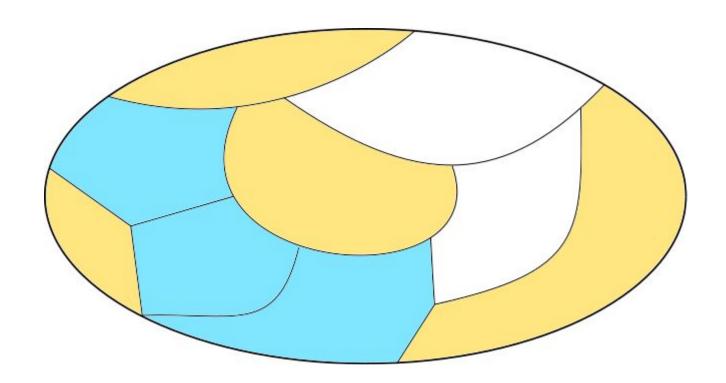
$$\omega(g,h) = \omega(g,h) + \phi_g + \phi_h - \phi_{g,h} \bmod 2\pi$$

1D Symmetry protected phases



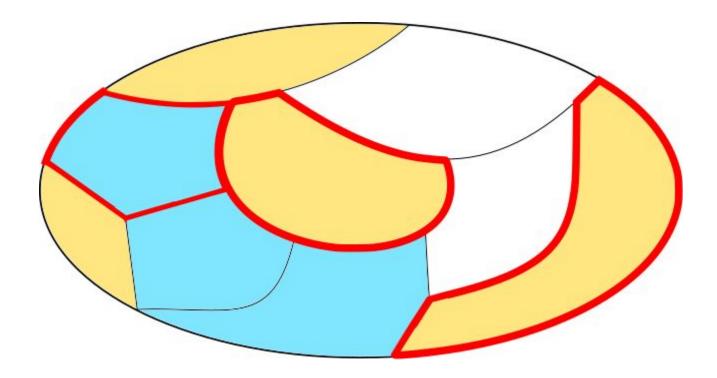
Equivalence classes of projective representations

- symmetry broken phase
- symmetry protected phase



Space of local Hamiltonians with symmetry G

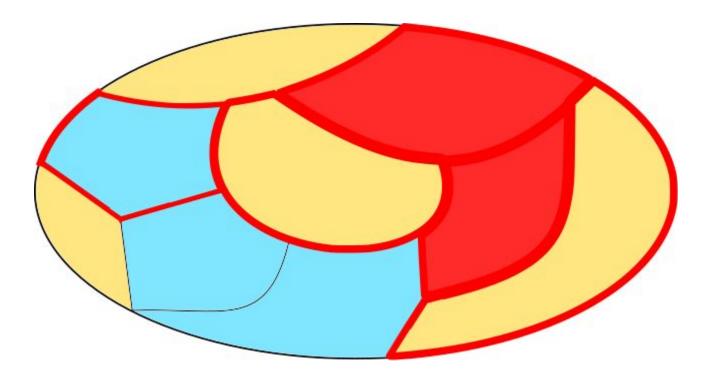
- symmetry broken phase
- symmetry protected phase
- critical



Space of local Hamiltonians with symmetry G

Furuya and Oshikawa PRL (2017)
Scaffidi, Parker, and Vasseur, PRX (2017)
Bridgeman and Williamson PRB (2017)
Verresen, Jones, and Pollmann PRL (2018)

- symmetry broken phase
- symmetry protected phase
- critical



Space of local Hamiltonians with symmetry G

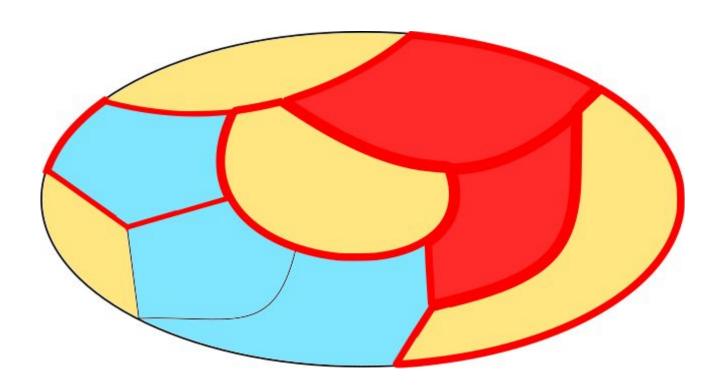
Furuya and Oshikawa PRL (2017)
Scaffidi, Parker, and Vasseur, PRX (2017)
Bridgeman and Williamson PRB (2017)
Verresen, Jones, and Pollmann PRL (2018)

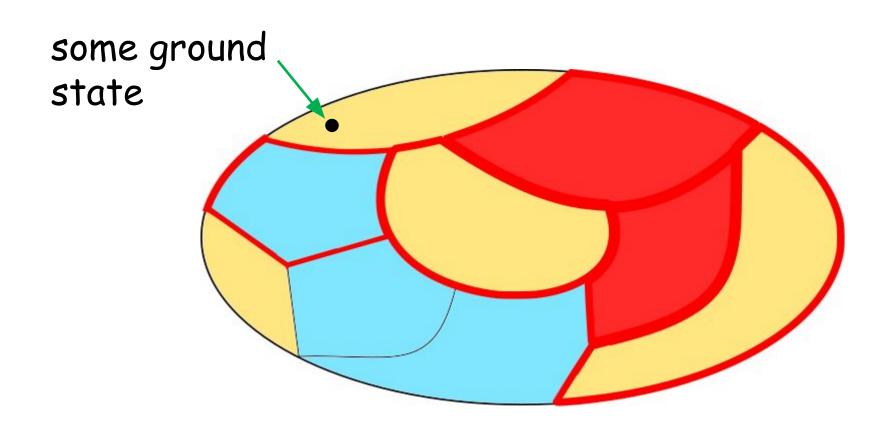
- symmetry broken phase
- symmetry protected phase
- critical

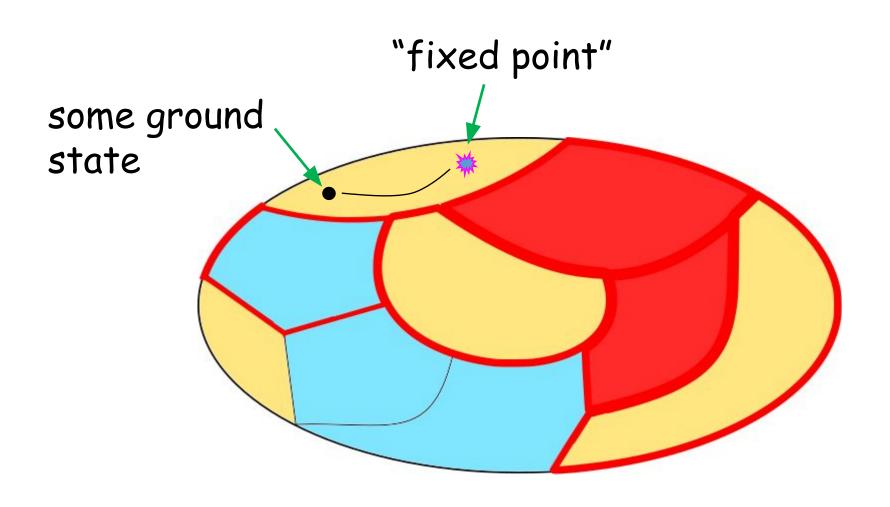
This talk --- approach classification of critical phases. Does symmetry protection play a role? But using the RG classification of phases

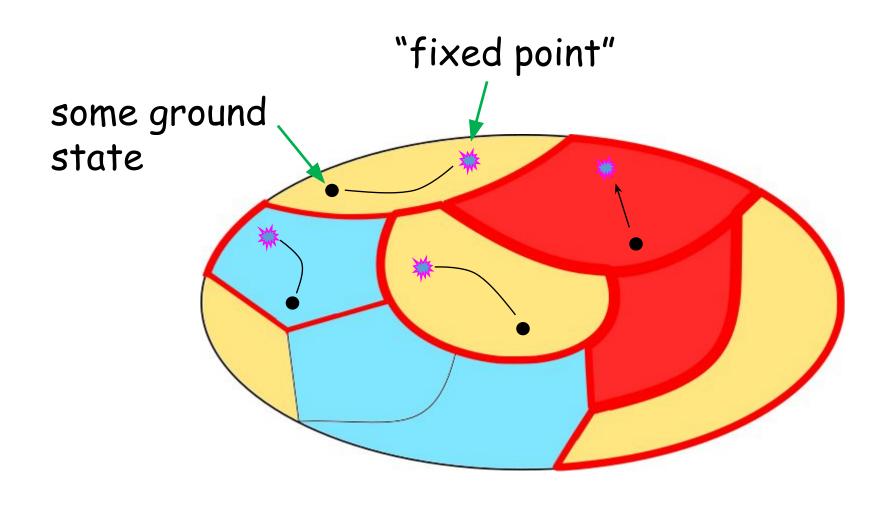
Space of local Hamiltonians with symmetry G

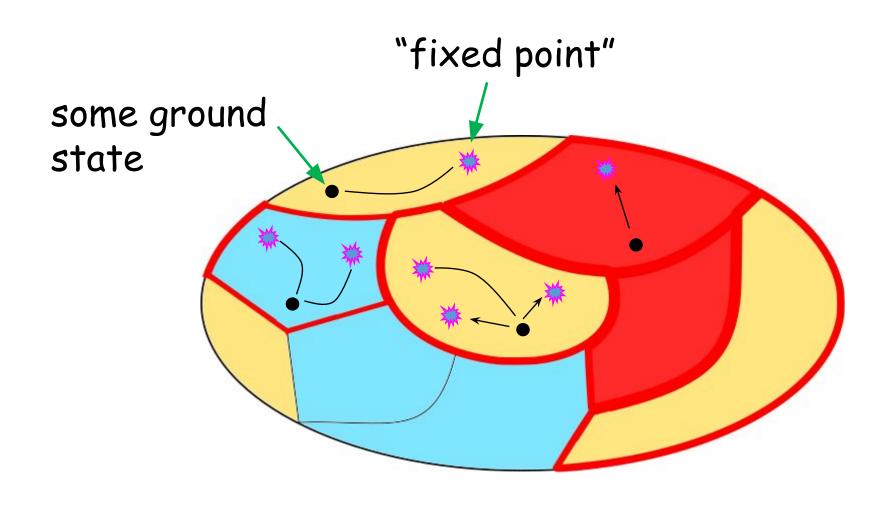
Furuya and Oshikawa PRL (2017)
Scaffidi, Parker, and Vasseur, PRX (2017)
Bridgeman and Williamson PRB (2017)
Verresen, Jones, and Pollmann PRL (2018)



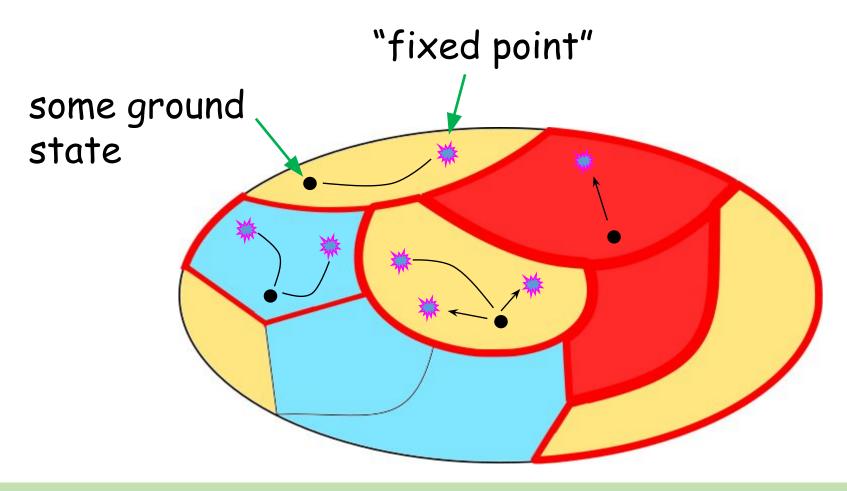




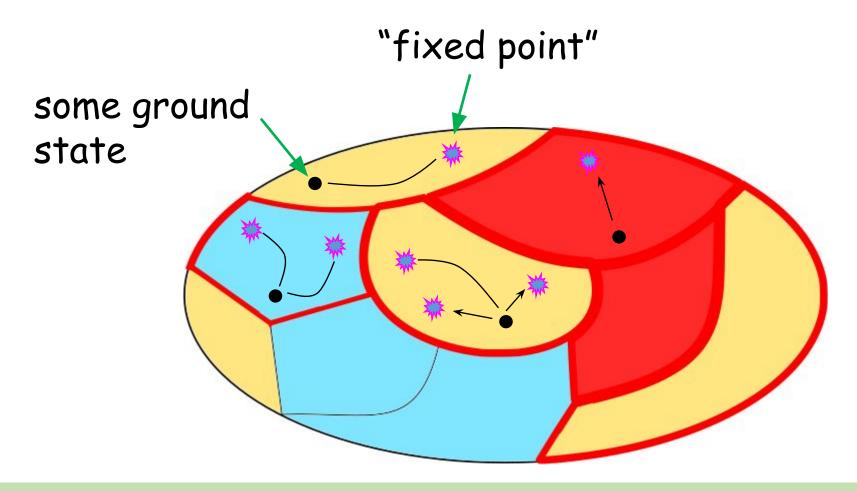




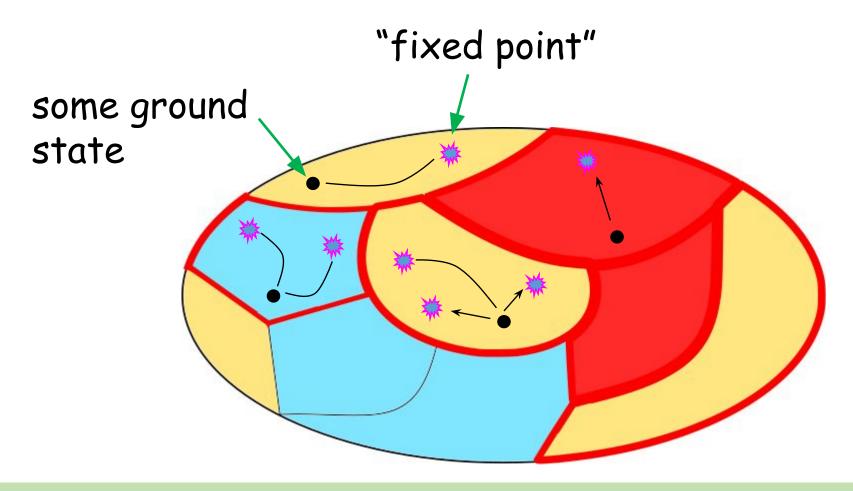
Different phases essentially differ in large length scale properties



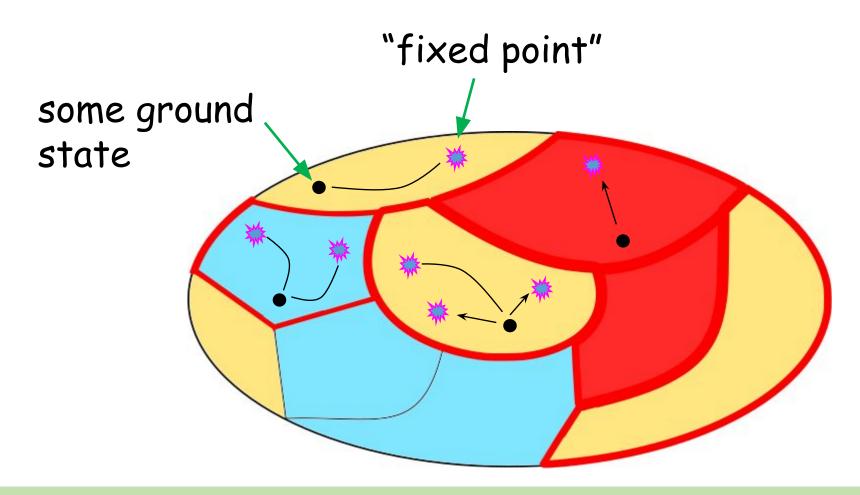
1. States in distinct phases must flow to distinct fixed points.



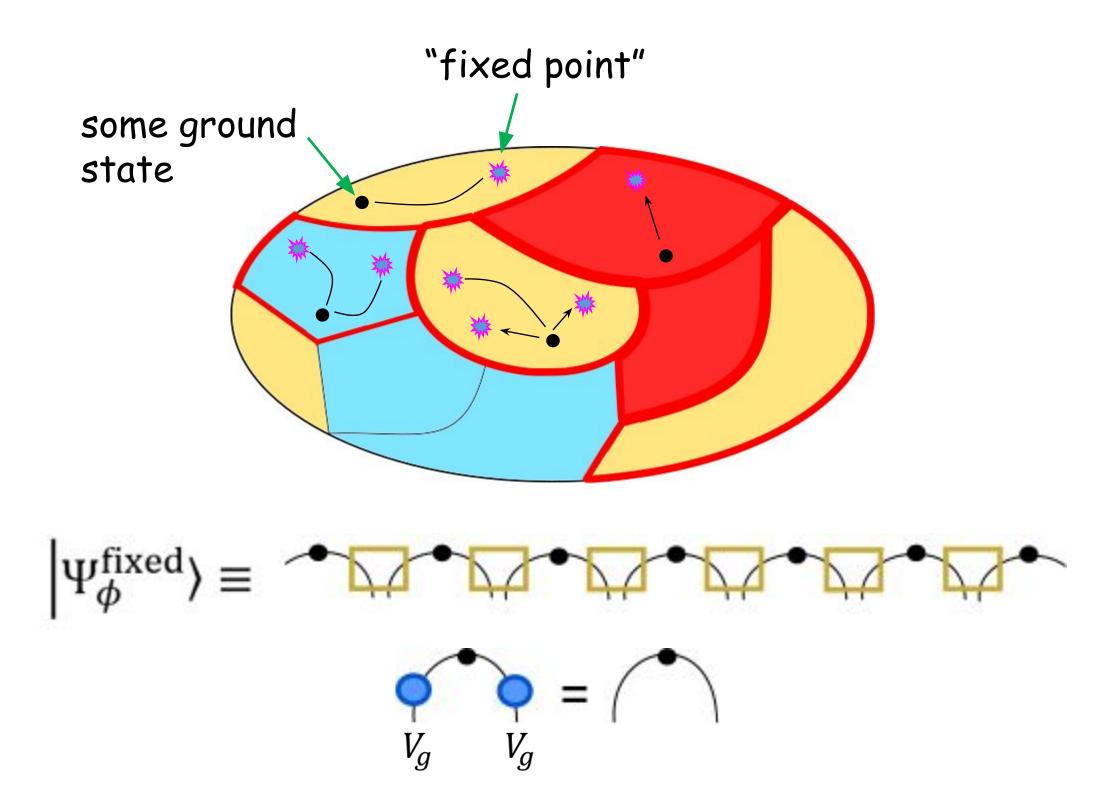
- 1. States in distinct phases must flow to distinct fixed points.
- 2. Gapped states must flow to short-range entangled fixed points.



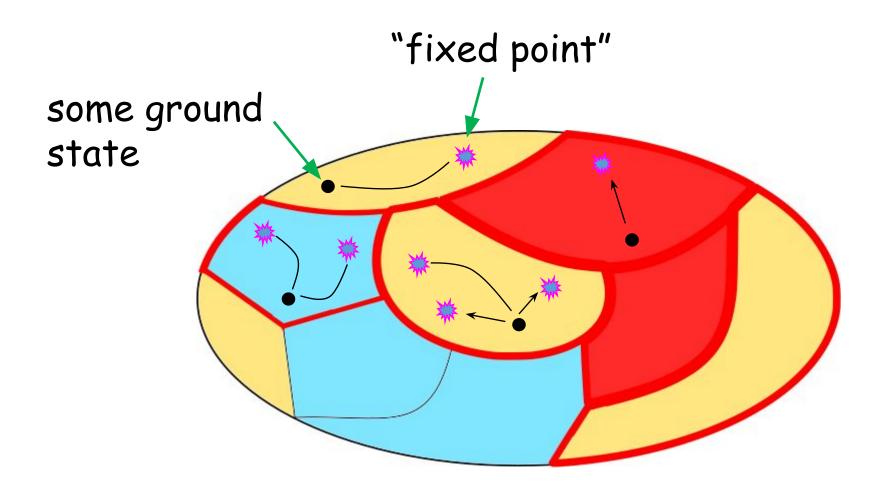
- 1. States in distinct phases must flow to distinct fixed points.
- 2. Gapped states must flow to short-range entangled fixed points.
- 3. Critical states must flow to long-range entangled fixed points.



- 1. States in distinct phases must flow to distinct fixed points.
- 2. Gapped states must flow to short-range entangled fixed points.
- 3. Critical states must flow to long-range entangled fixed points.
- 4. Symmetry protected phases flow to distinct fixed points only if the symmetry is protected along the RG.



Different phases essentially differ in large length scale properties



This talk --- approach classification of phases using specific symmetry protected RG flows, which are described by tensor networks

Outline

- 1. Tensor network representations of 1D ground states
 - 1.1 MPS
 - 1.2 MERA
- 2. MPS and 1D gapped symmetry protected phases
- 3. Symmetry protected MERA
 - 3.1 Gapped SP phases
 - 3.2 Critical SP phases

Outline

- 1. Tensor network representations of 1D ground states
 - 1.1 MPS
 - 1.2 MERA
- 2. MPS and 1D gapped symmetry protected phases
- 3. Symmetry protected MERA
 - 3.1 Gapped SP phases
 - 3.2 Critical SP phases

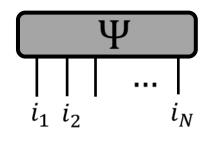
Matrix product states (MPS)

1D gapped ground state

$$|\Psi\rangle = \sum \Psi_{i_1 i_2 \dots i_N} \ |i_1\rangle \otimes |i_2\rangle \otimes \cdots |i_N\rangle$$

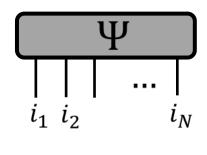
1D gapped ground state

$$|\Psi\rangle = \sum \Psi_{i_1 i_2 \dots i_N} \ |i_1\rangle \otimes |i_2\rangle \otimes \cdots |i_N\rangle$$



(tensor)

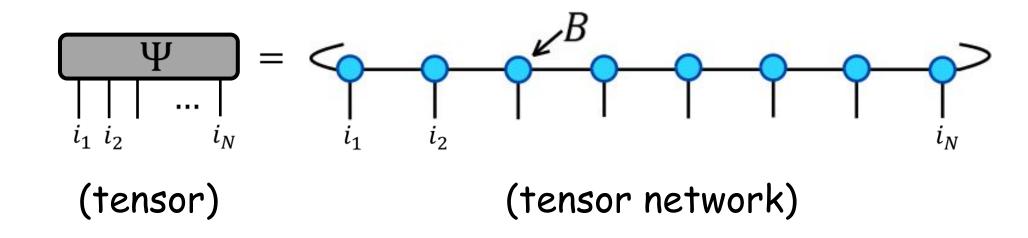
$$|\Psi\rangle = \sum \Psi_{i_1 i_2 \dots i_N} |i_1\rangle \otimes |i_2\rangle \otimes \dots |i_N\rangle$$



(tensor)

1D gapped ground state

$$|\Psi\rangle = \sum \Psi_{i_1 i_2 \dots i_N} |i_1\rangle \otimes |i_2\rangle \otimes \dots |i_N\rangle$$

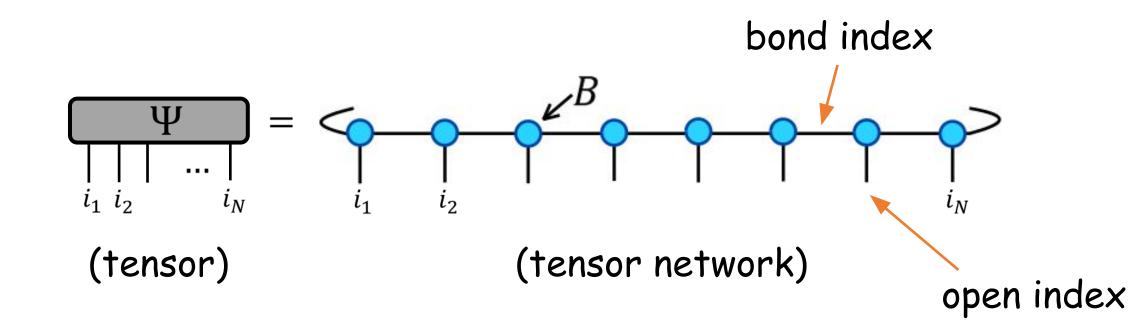


Fannes, Nachtergaele, and Werner, Commun. Math. Phys. (1992)

Verstraete and Cirac, PRB (2006)

1D gapped ground state

$$|\Psi\rangle = \sum \Psi_{i_1 i_2 \dots i_N} |i_1\rangle \otimes |i_2\rangle \otimes \dots |i_N\rangle$$

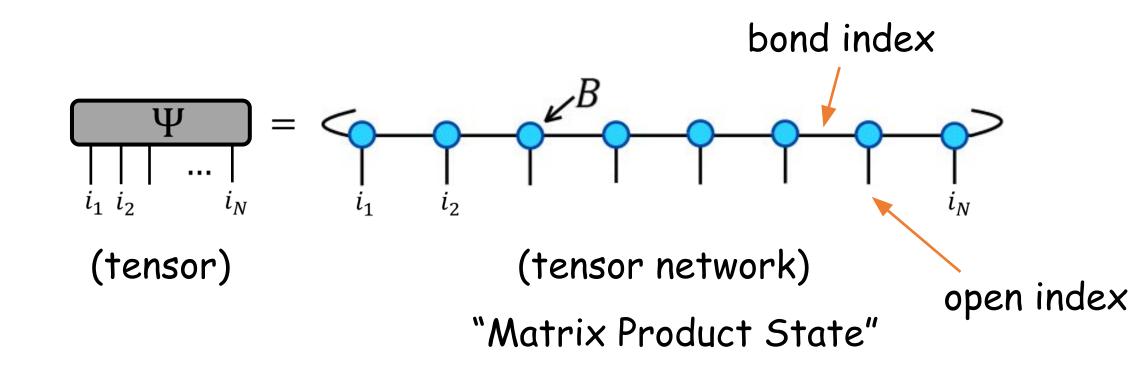


Fannes, Nachtergaele, and Werner, Commun. Math. Phys. (1992)

Verstraete and Cirac, PRB (2006)

1D gapped ground state

$$|\Psi\rangle = \sum \Psi_{i_1 i_2 \dots i_N} |i_1\rangle \otimes |i_2\rangle \otimes \dots |i_N\rangle$$



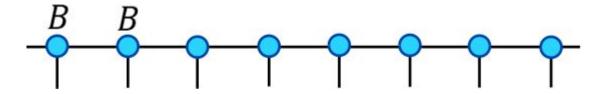
Fannes, Nachtergaele, and Werner, Commun. Math. Phys. (1992)

Verstraete and Cirac, PRB (2006)

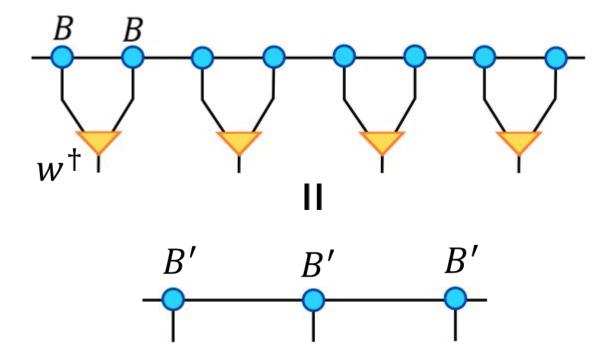
MERA

Based on a coarse-graining transformation

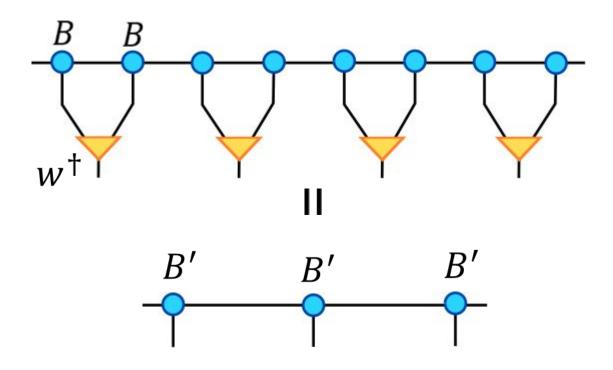
1D gapped ground state on an infinite lattice



1D gapped ground state on an infinite lattice

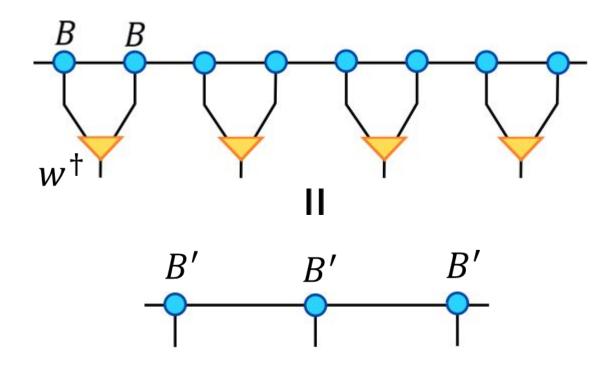


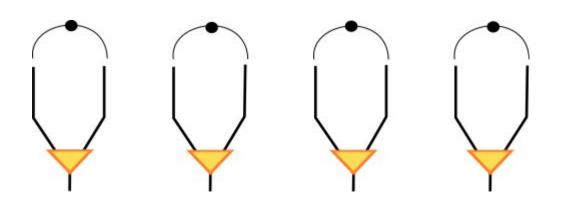
1D gapped ground state on an infinite lattice





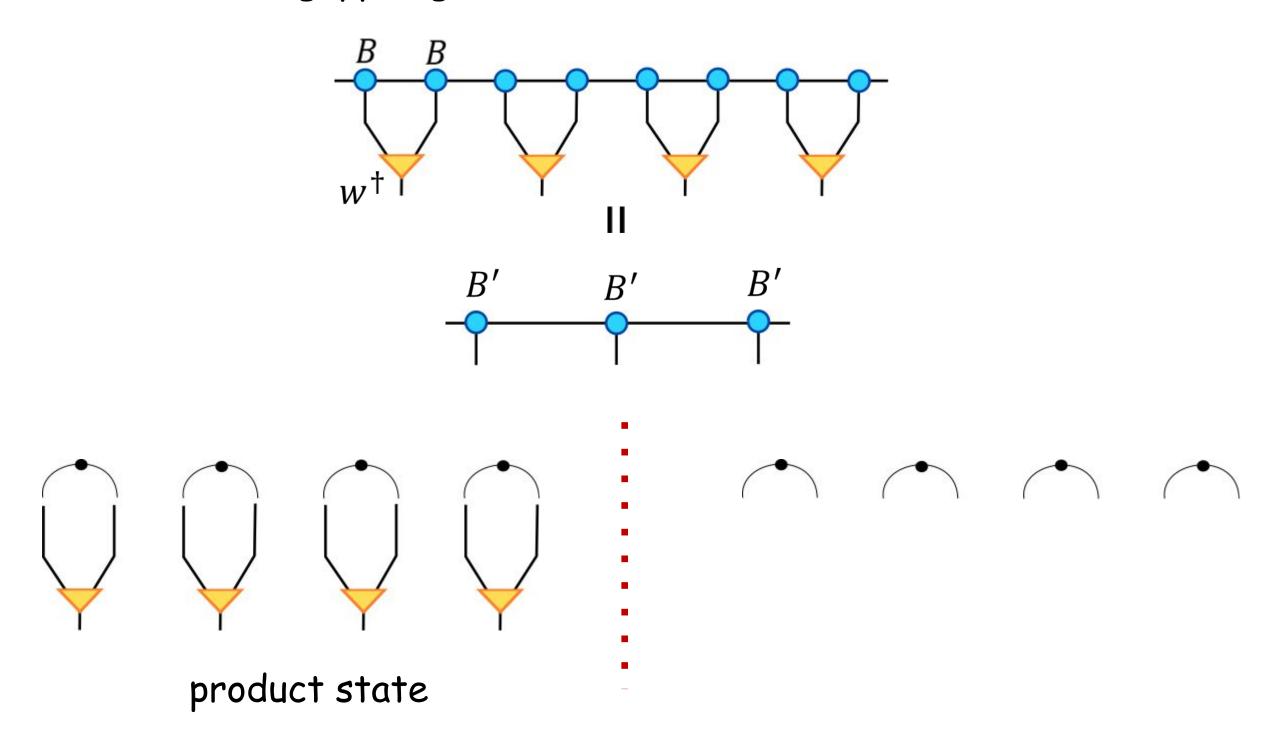
1D gapped ground state on an infinite lattice



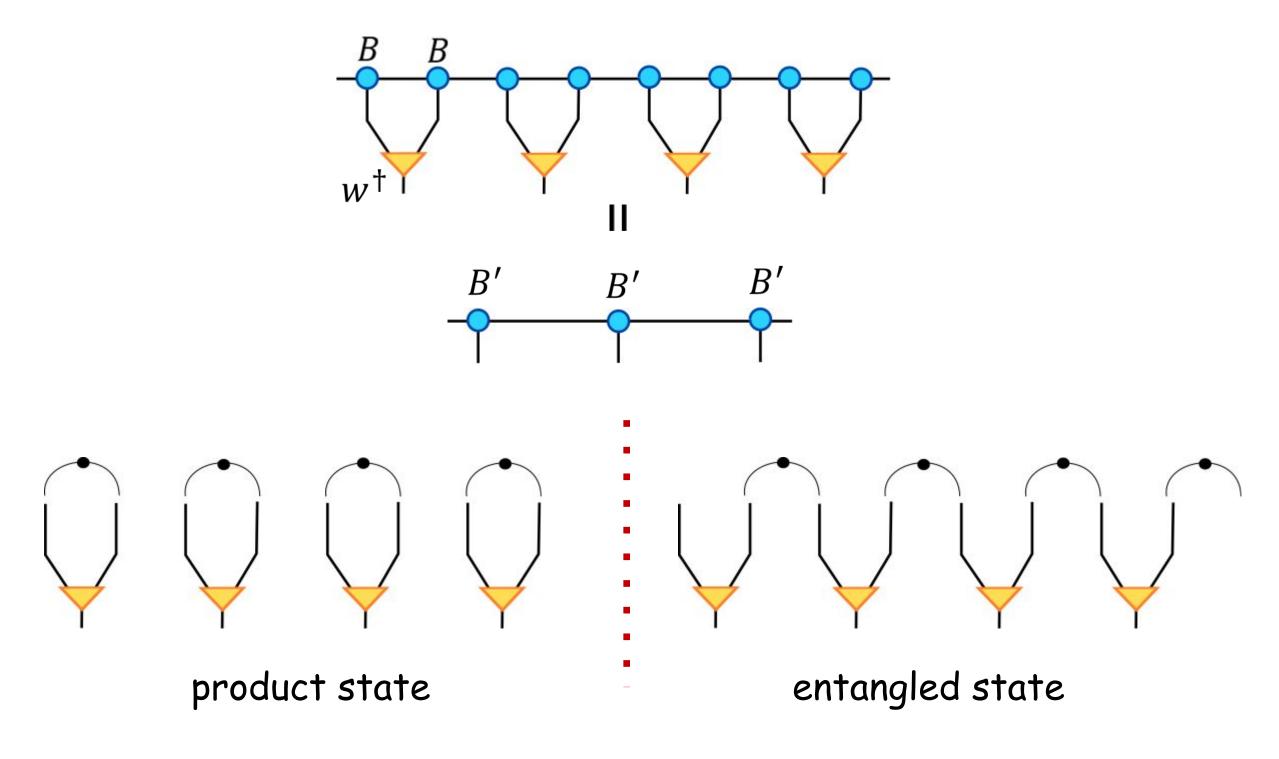


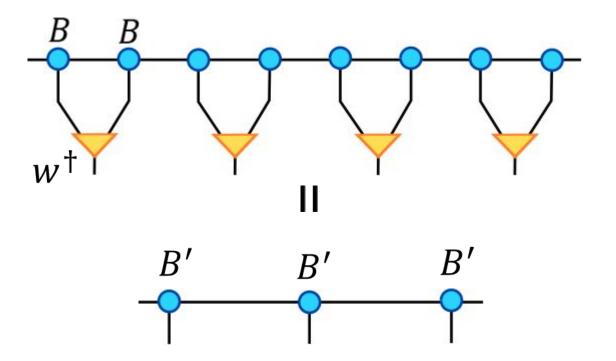
product state

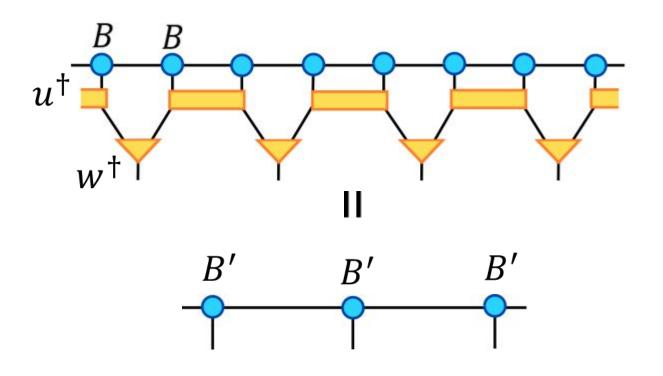
1D gapped ground state on an infinite lattice

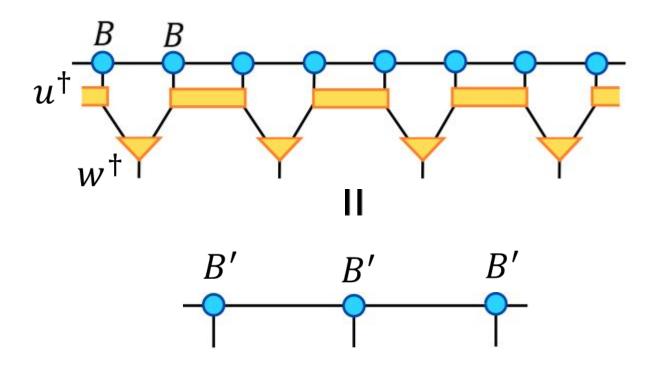


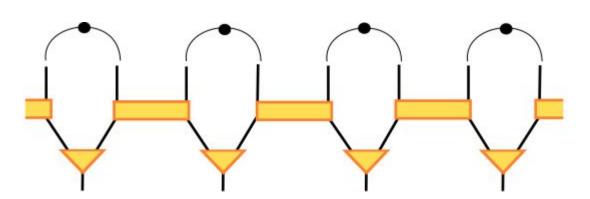
1D gapped ground state on an infinite lattice





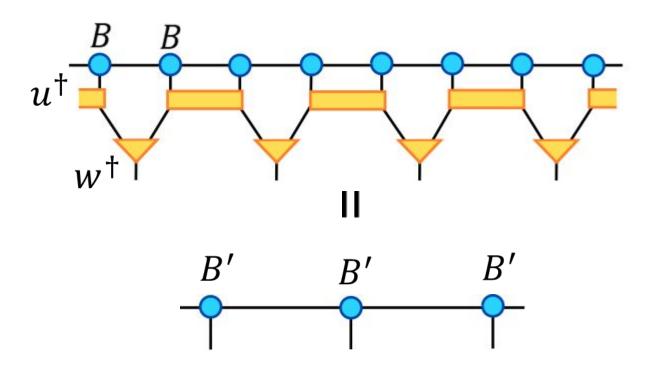


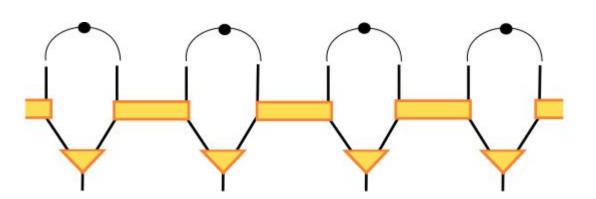




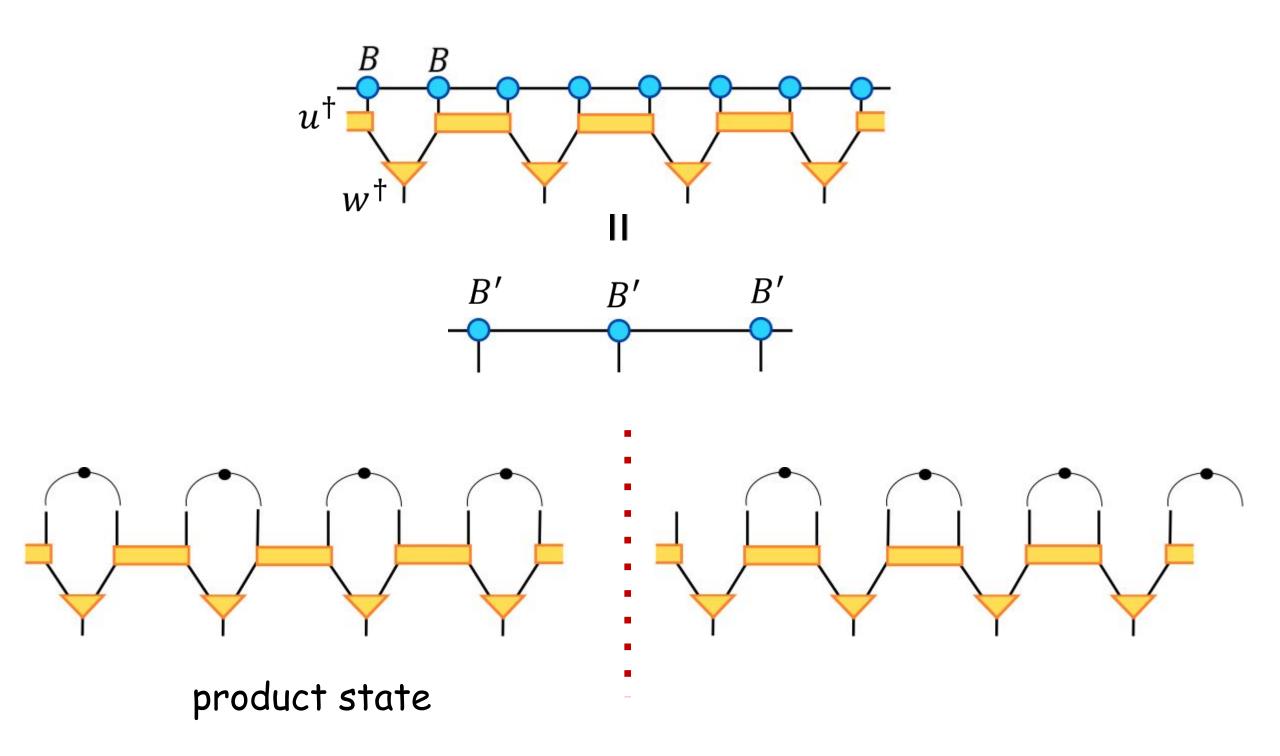
"Entanglement Renormalization"

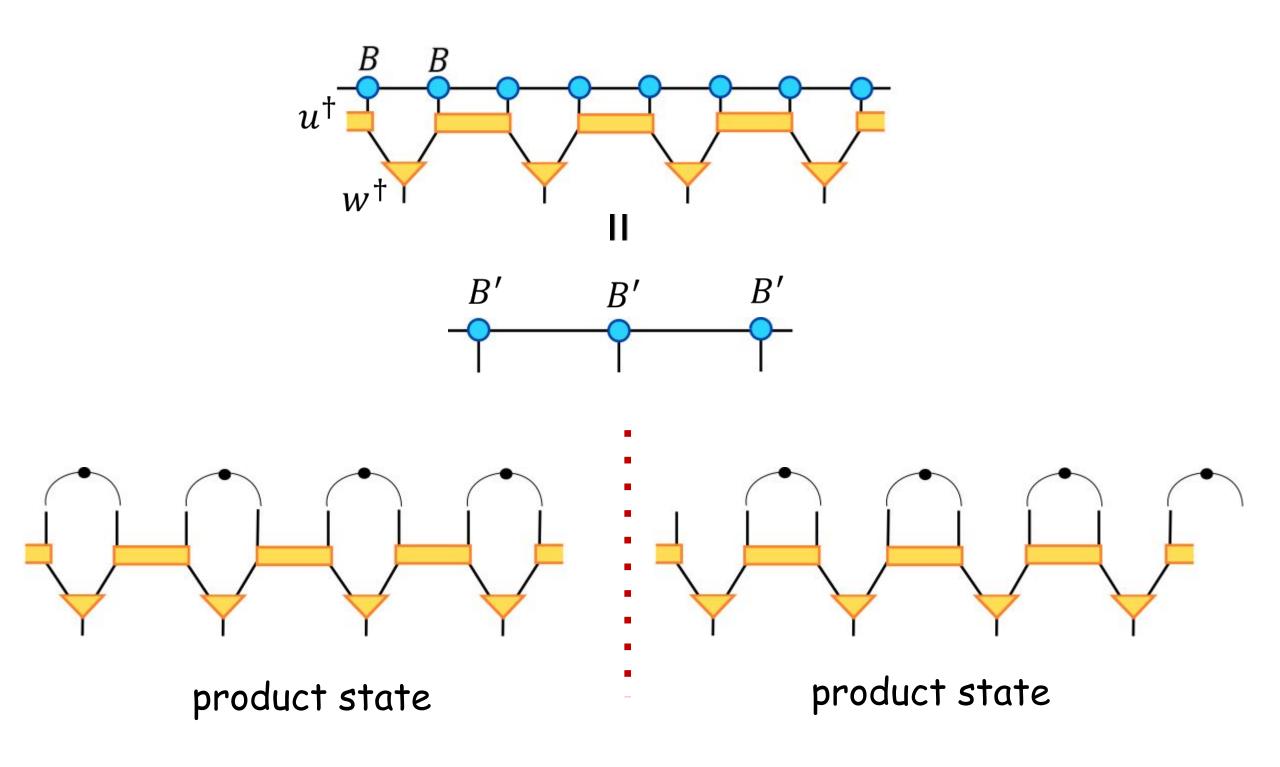
Vidal PRL (2007)

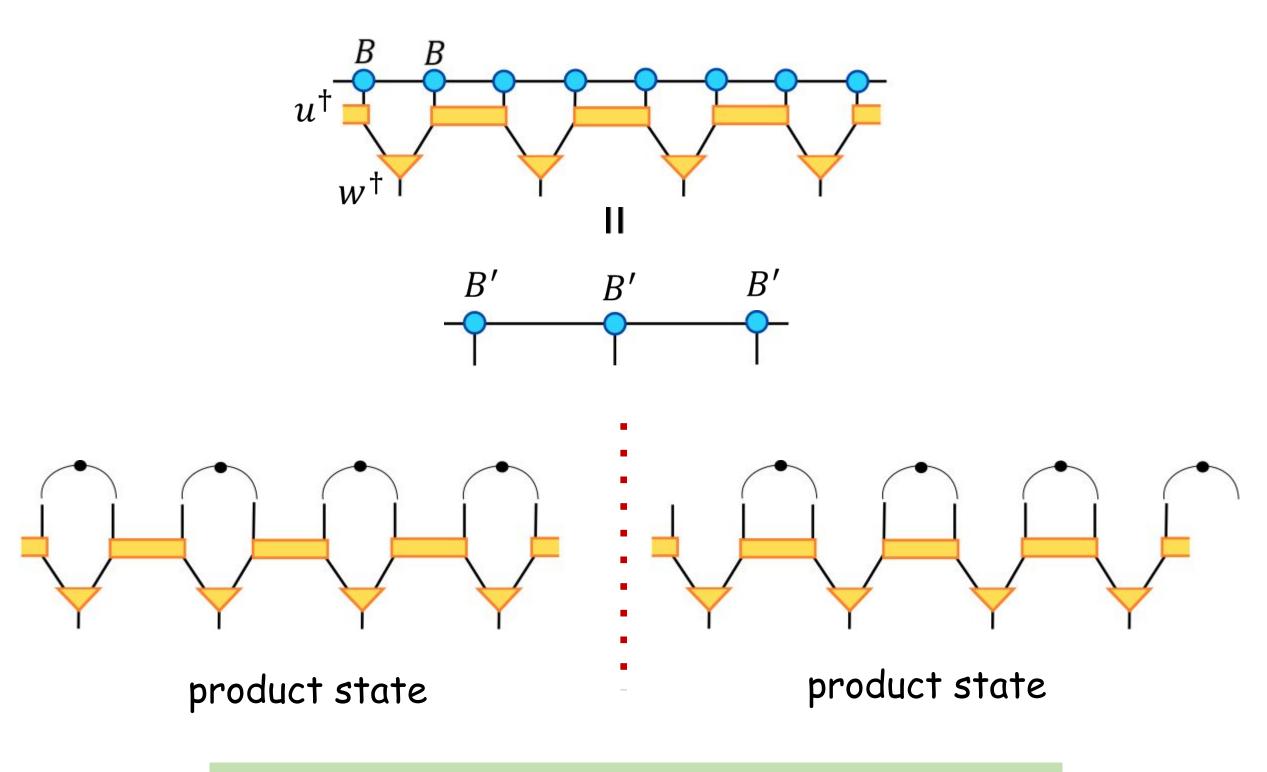




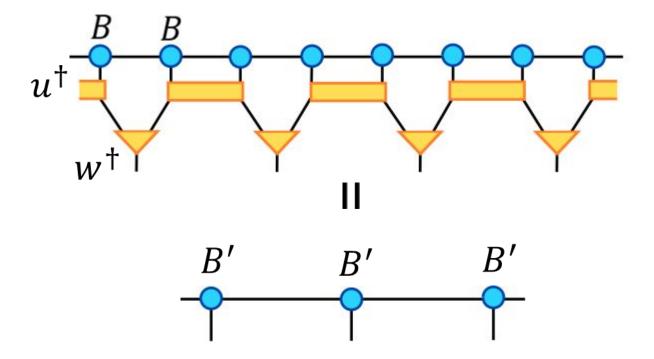
product state

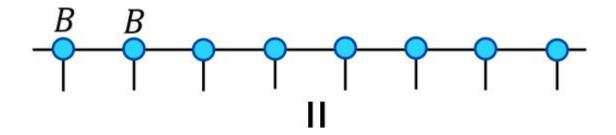


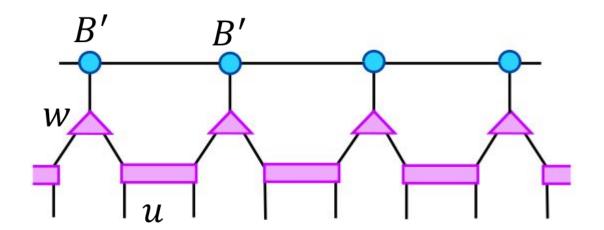


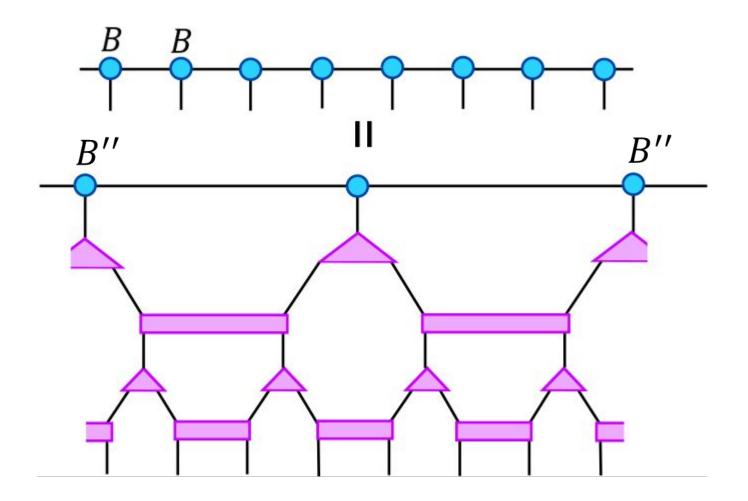


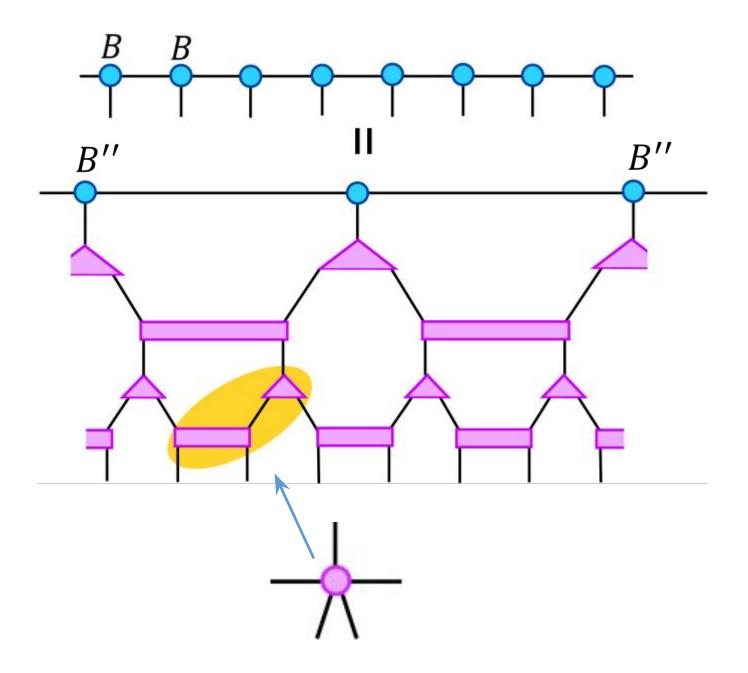
More effective at removing short range details



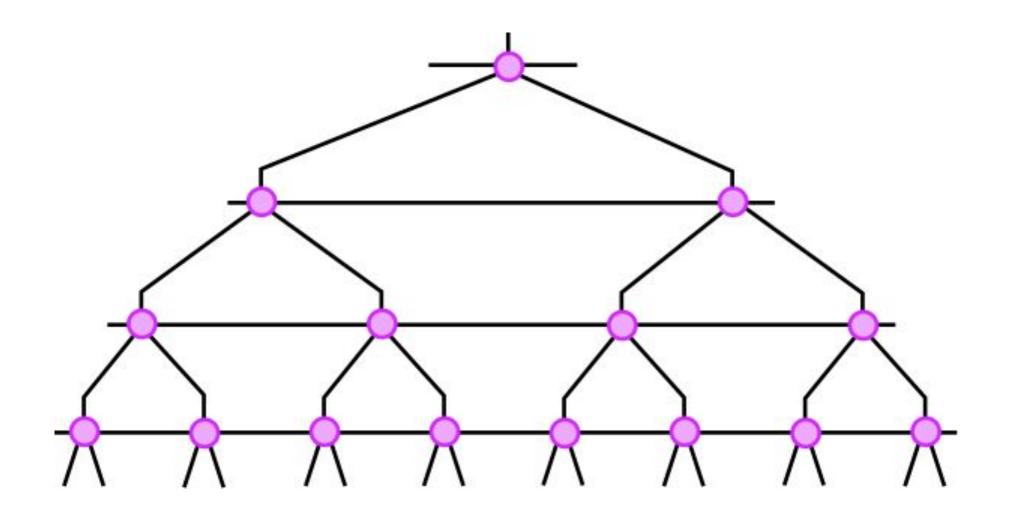






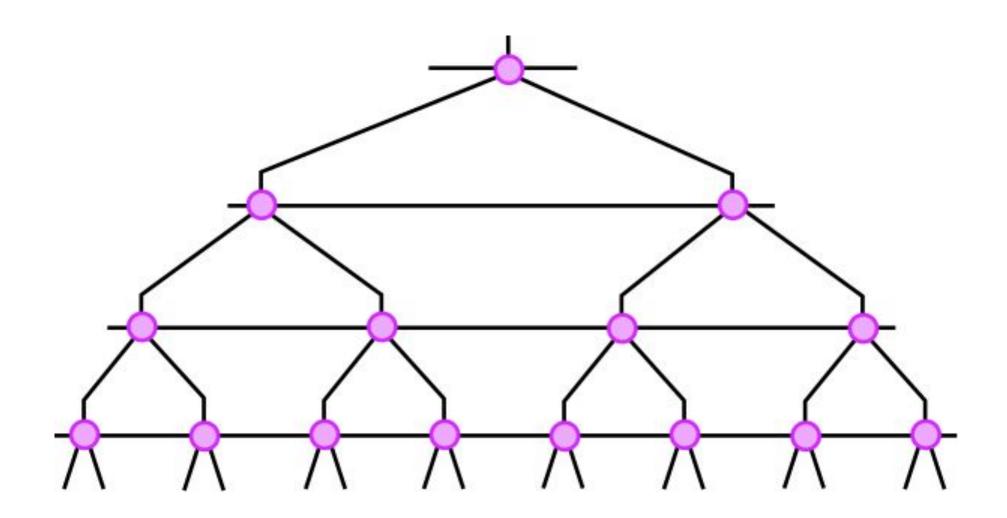


Multi-scale Entanglement Renormalization Ansatz (MERA)
Vidal PRL (2008)



Multi-scale Entanglement Renormalization Ansatz (MERA)
Vidal PRL (2008)

MERA representation of a state = RG flow + Fixed point

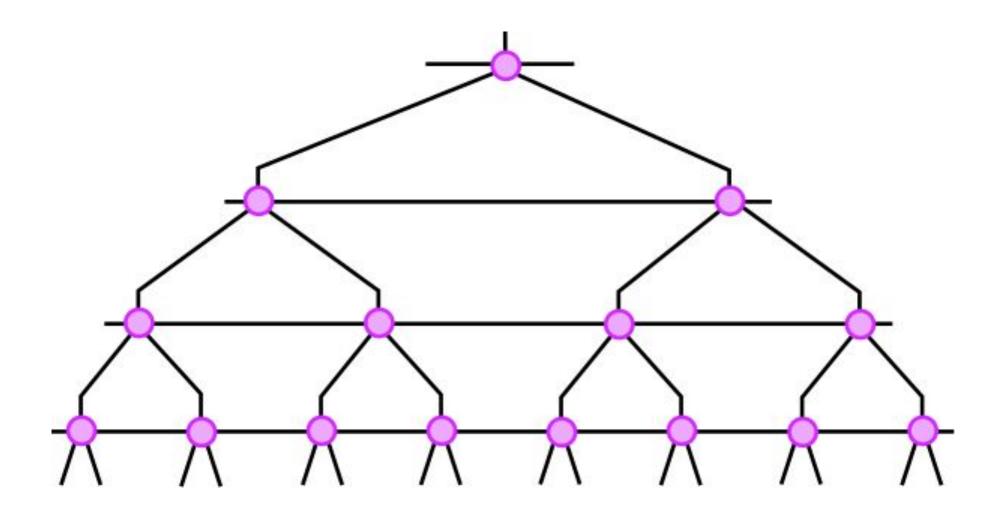


Multi-scale Entanglement Renormalization Ansatz (MERA)

Vidal PRL (2008)

MERA representation of a state = RG flow + Fixed point

MERA obtained from an MPS has finite correlation length.



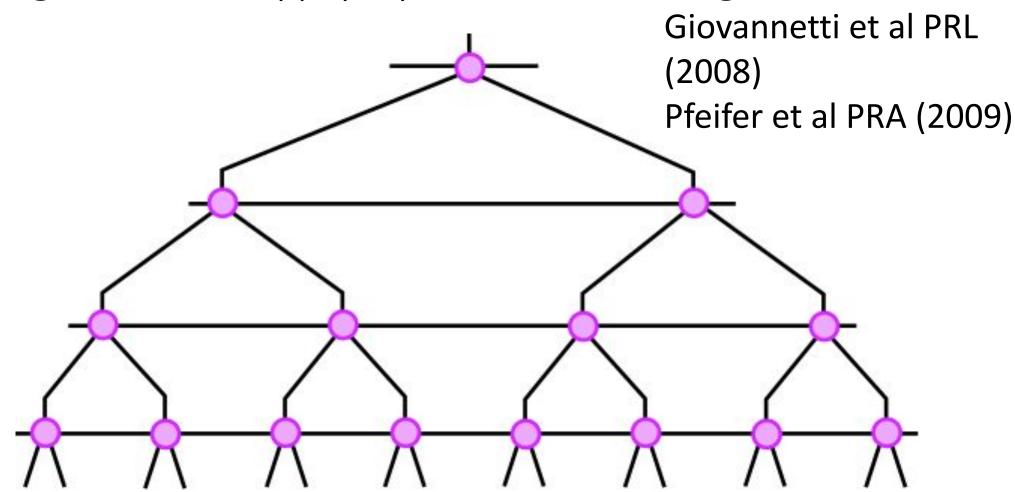
Multi-scale Entanglement Renormalization Ansatz (MERA)

Vidal PRL (2008)

MERA representation of a state = RG flow + Fixed point

MERA obtained from an MPS has finite correlation length.

But generic choice of tensors describes a state with polynomial decaying correlations and logarithmic growth of entanglement entropy (properties of critical ground states).



MPS & 1D gapped SP phases

Every MPS is the ground state of a local Hamiltonian

Every MPS is the ground state of a local Hamiltonian

Classifying 1D gapped SP phases



Classifying all symmetric MPSs

Chen, Gu, and Wen, PRB (2011) Schuch, Perez-Garcia, and Cirac, PRB (2011) Pollmann, Berg, Turner, and Oshikawa, PRB (2012)

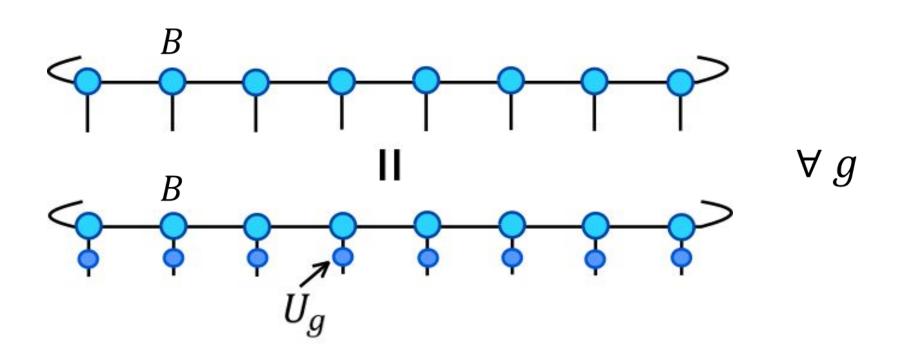
Every MPS is the ground state of a local Hamiltonian

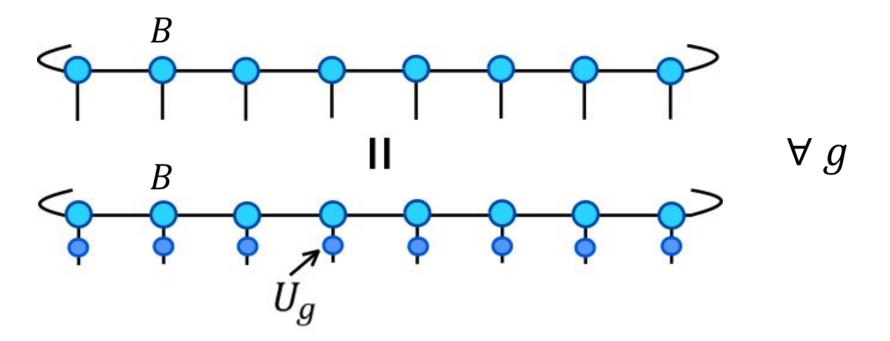
Classifying 1D gapped SP phases

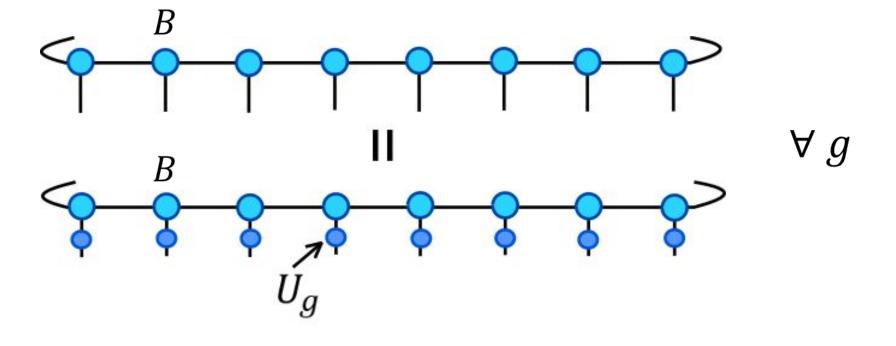
 \longleftrightarrow

Classifying all symmetric MPSs

Chen, Gu, and Wen, PRB (2011) Schuch, Perez-Garcia, and Cirac, PRB (2011) Pollmann, Berg, Turner, and Oshikawa, PRB (2012)

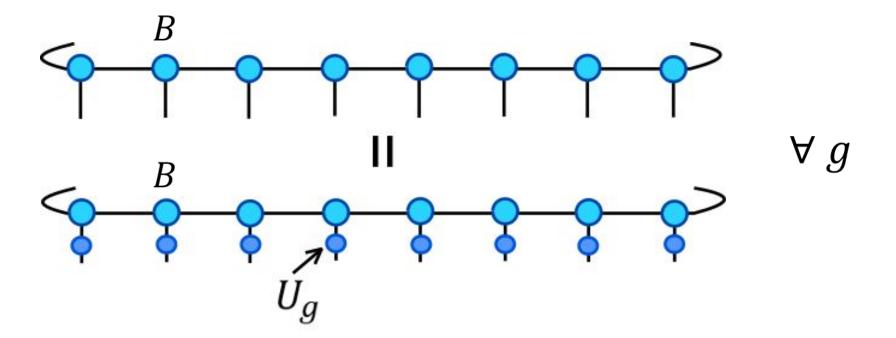






$$U_g = Y_g = Y_g^{B}$$

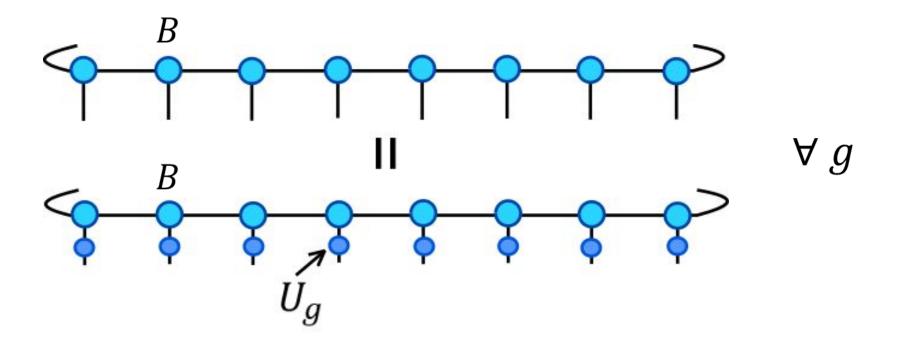
$$Y_g^{\dagger}$$



If MPS is normal & in the canonical form

$$U_g = Y_g = Y_g^B$$

$$Y_g^{\dagger}$$



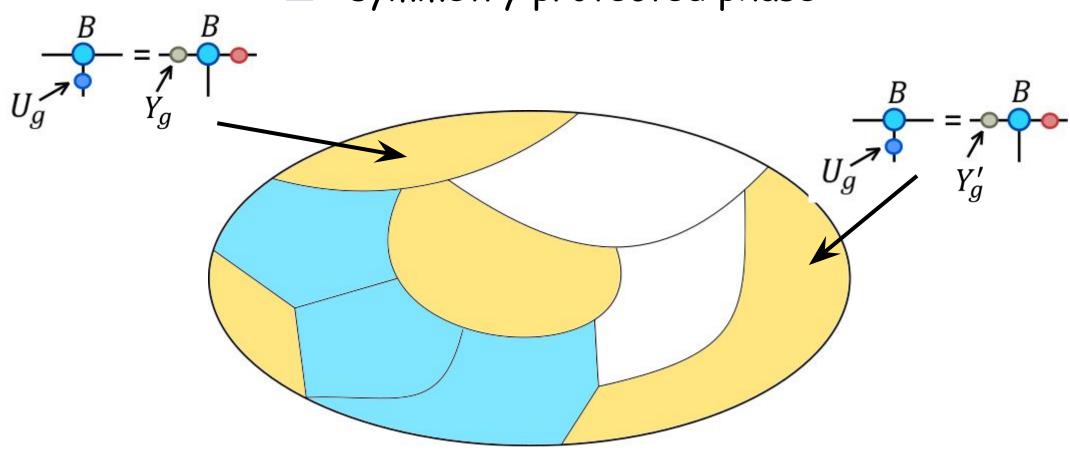
If MPS is normal & in the canonical form

$$U_g = V_g = V_g = V_g^{B}$$

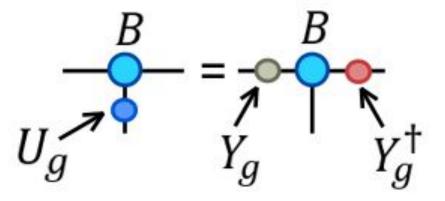
$$V_g = V_g^{\dagger}$$

$$Y_g$$
 can be projective $Y_g Y_h = e^{i\omega(g,h)} Y_{g,h}$

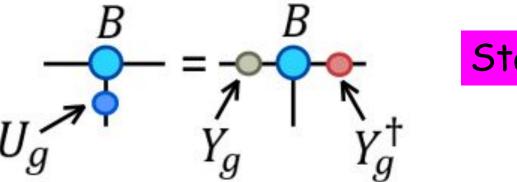
- symmetry broken phase
- symmetry protected phase



 Y_g and Y_g' are inequivalent projective representations



Star equation!

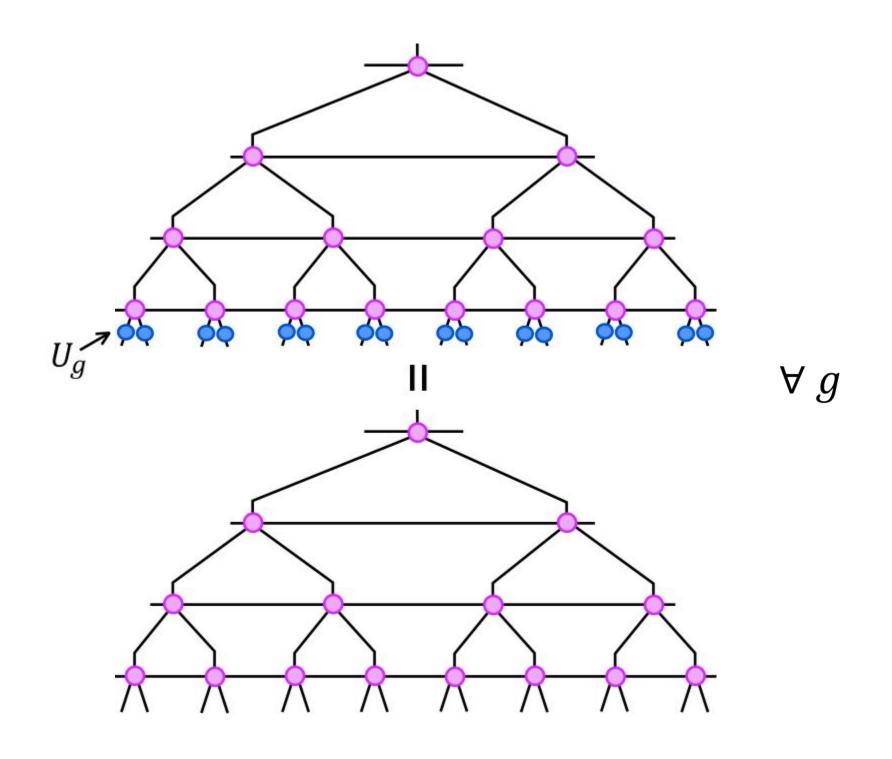


Star equation!

- 1. Conceptual tool for classifying phases
- 2. Detecting phases in numerical simulations
 Pollmann and Turner, PRB (2012)
 SS, PRB (2015)

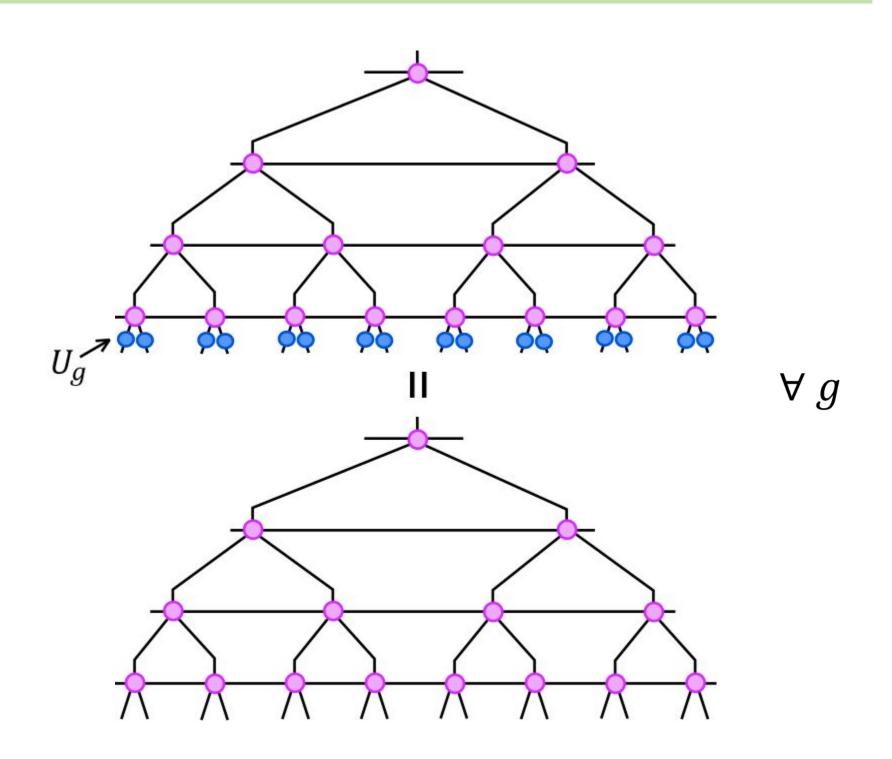
Symmetry protected MERA

Symmetry protected MERA

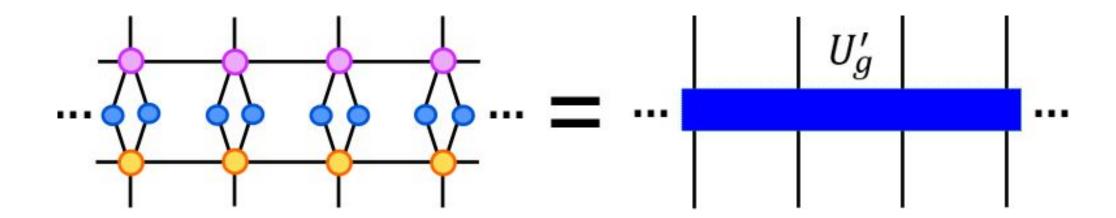


Symmetry protected MERA

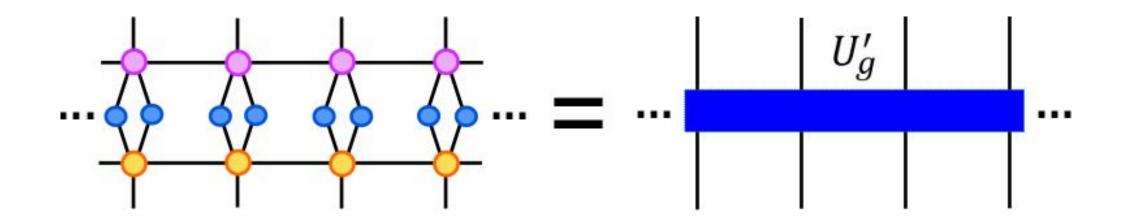
This should include both gapped and critical states



Symmetry protected entanglement renormalization



Symmetry protected entanglement renormalization



Assume U_g' is linear. But is it also on-site?

Argument for symmetry remaining on-site

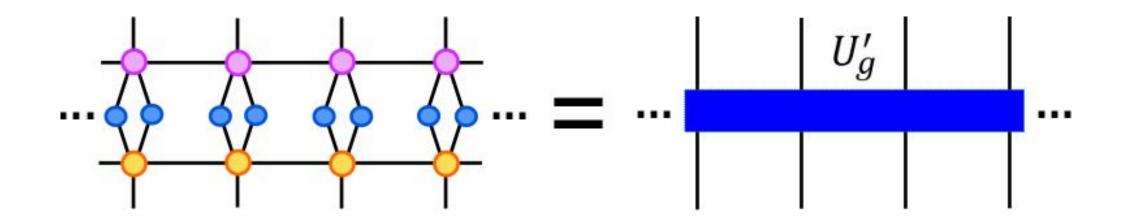
1. Global on-site symmetries can always be gauged

2. If a symmetry cannot be gauged it acts anomalously (acts in a non-onsite way)

3. At least in QFT legitimate RG flows must preserve any anomalies ('t Hooft anomaly matching condition)

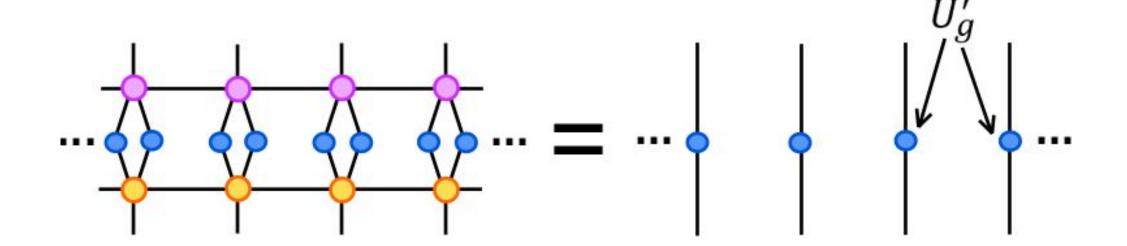
4. A trivial anomaly (on-site) must remain trivial (on-site) under a legitimate RG flow

Symmetry protected entanglement renormalization



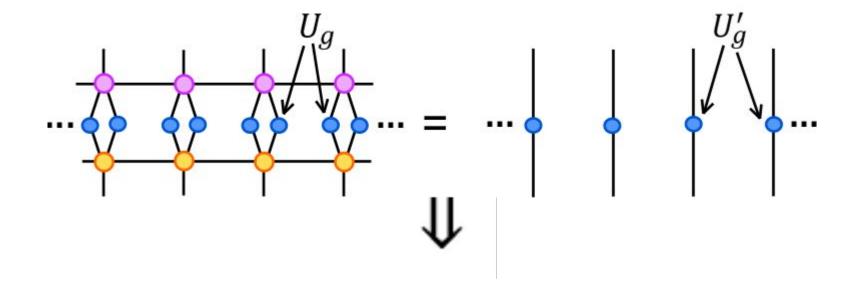
Assume U_g' is linear. But is it also on-site?

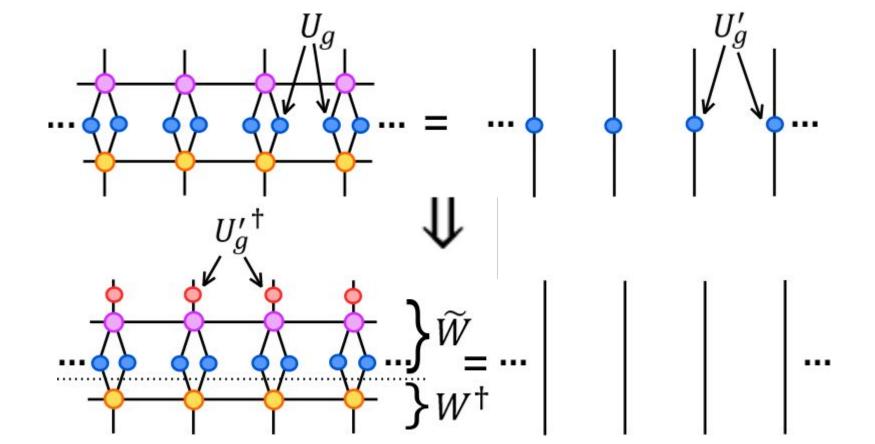
Symmetry protected entanglement renormalization

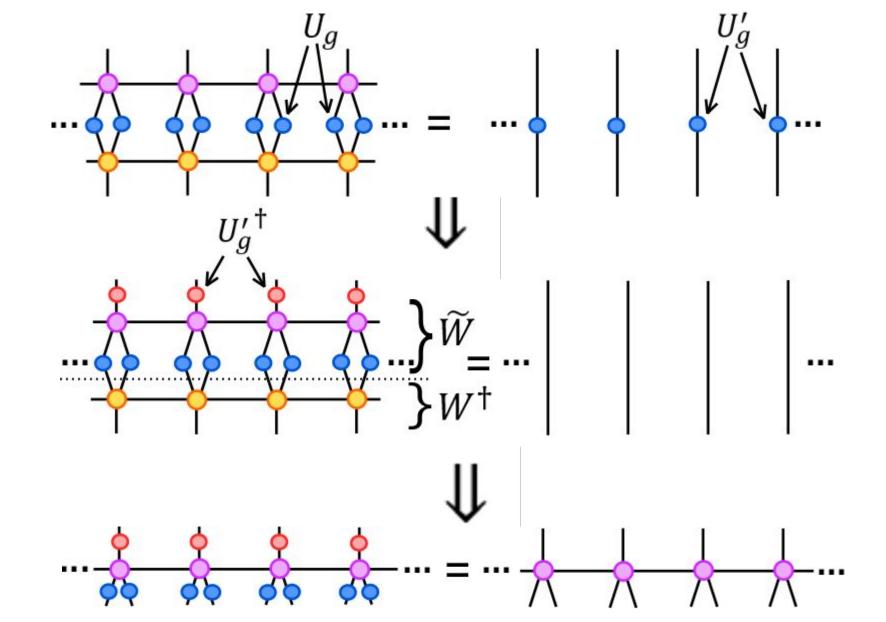


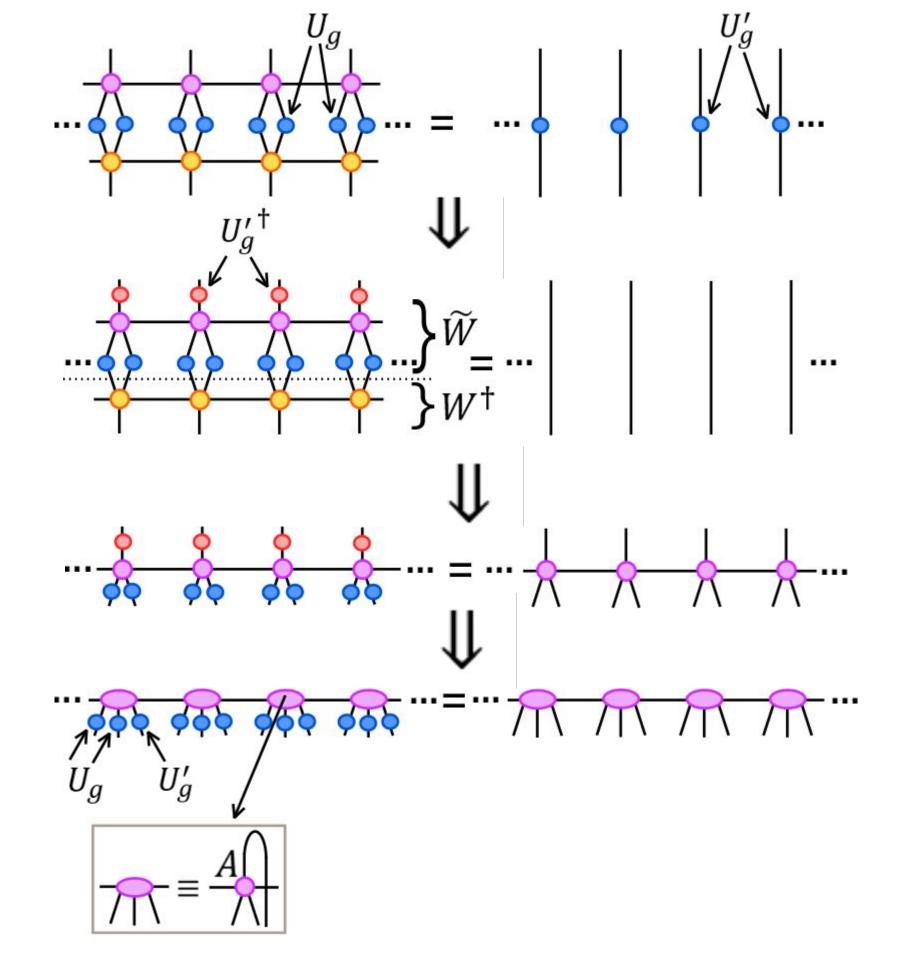
Assume U_g' is linear. But is it also on-site?

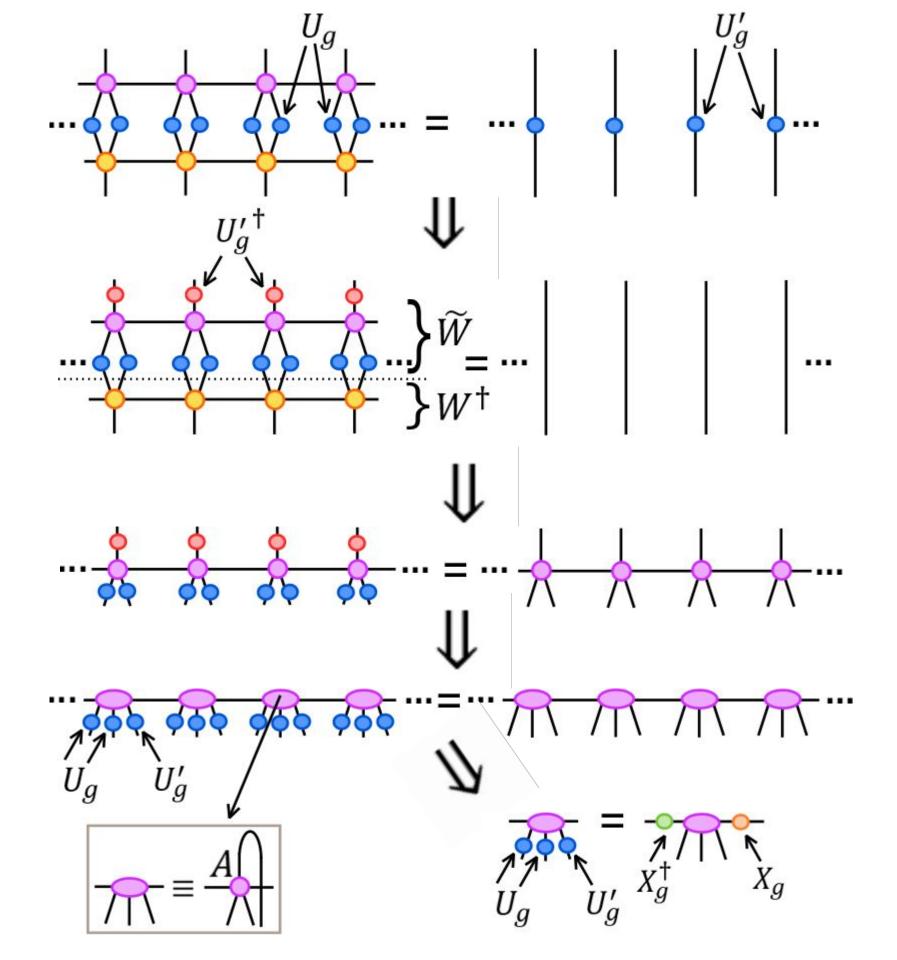
Yes (anomaly matching)

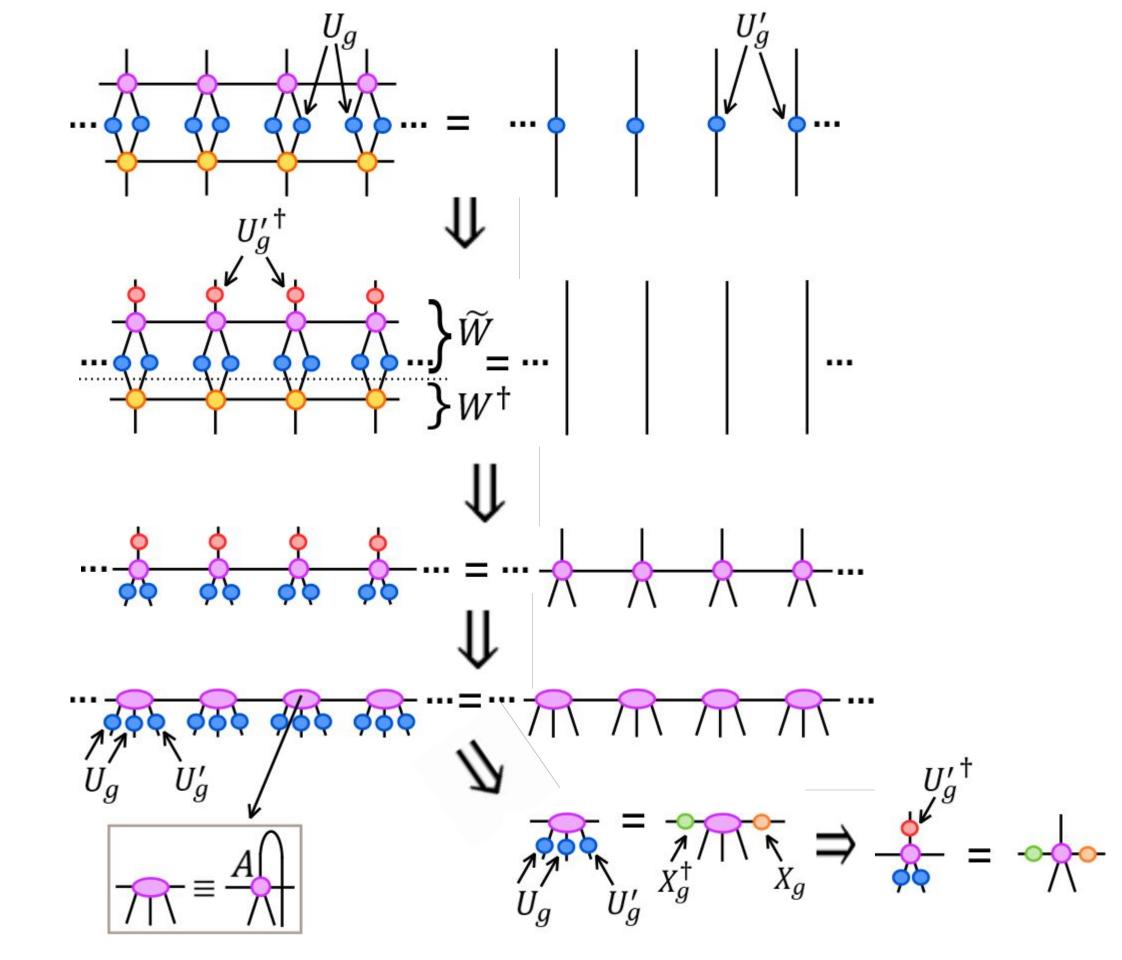


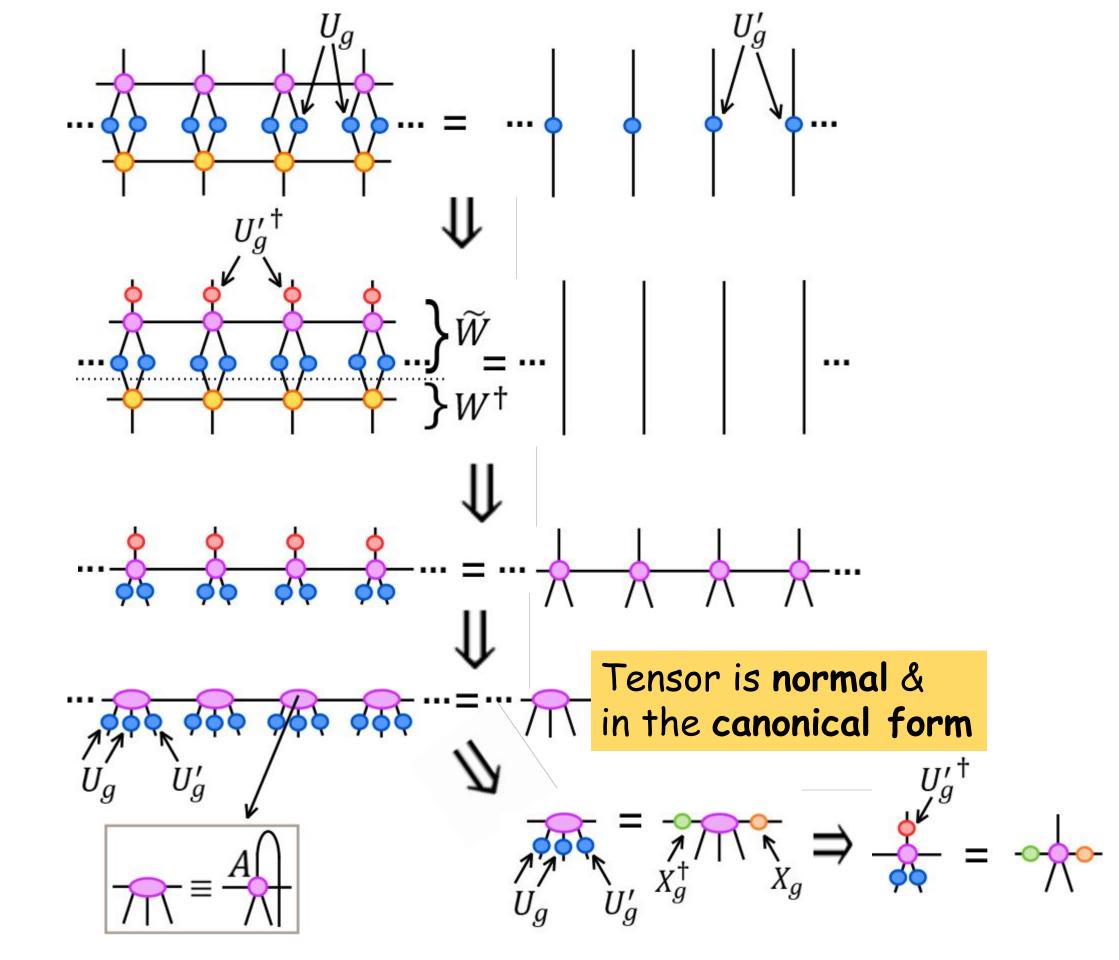


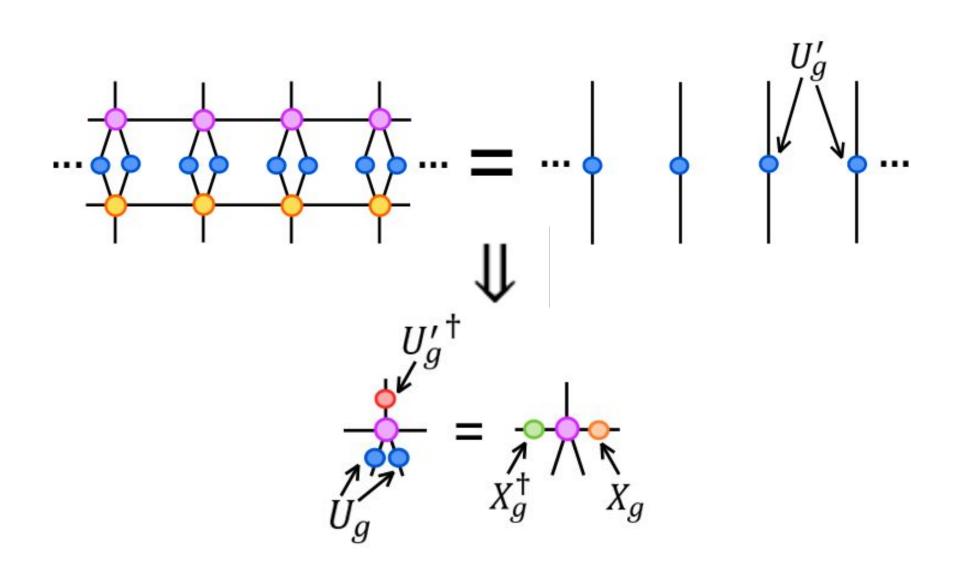


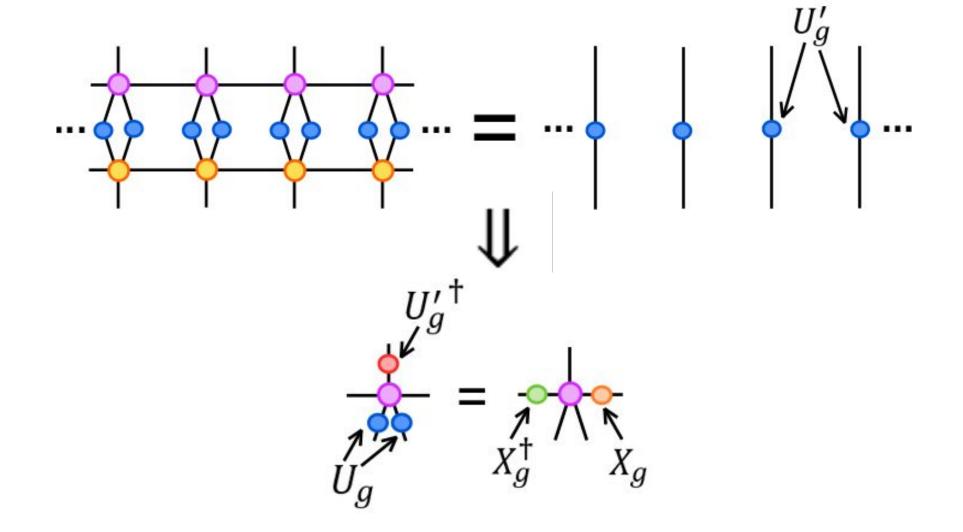












 X_g can be **projective**

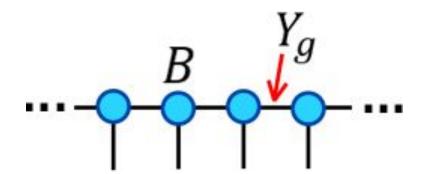
In a gapped phase? (No)

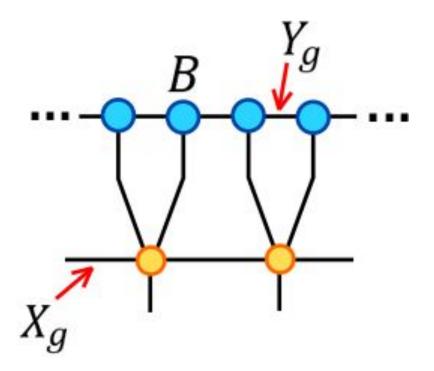
In a critical phase? (?)

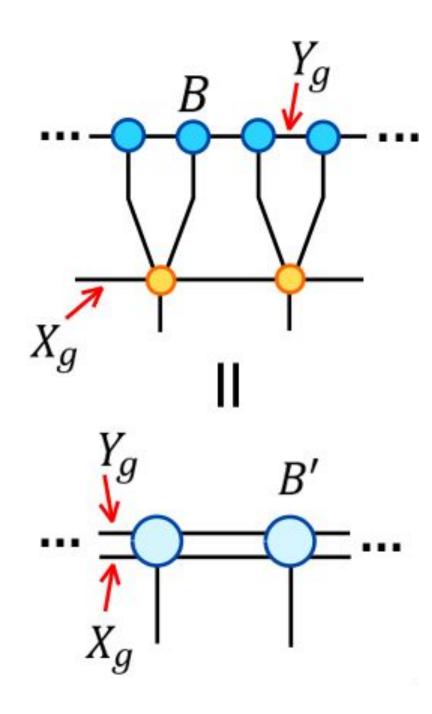
MERA and 1D gapped SP phases

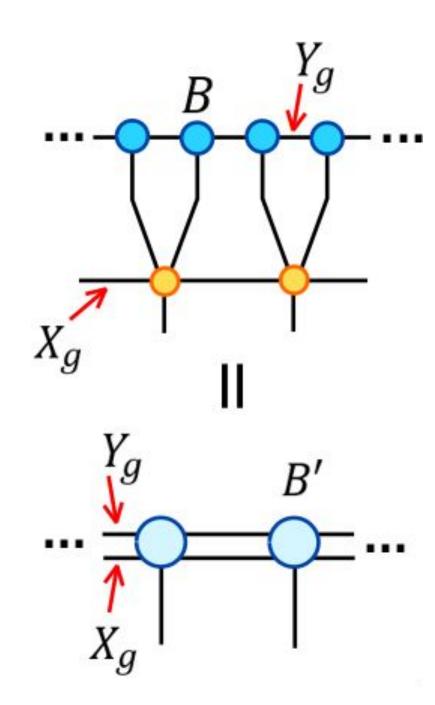
Obtain MERA by coarse-graining MPS

MERA representation of a state = RG flow + Fixed point

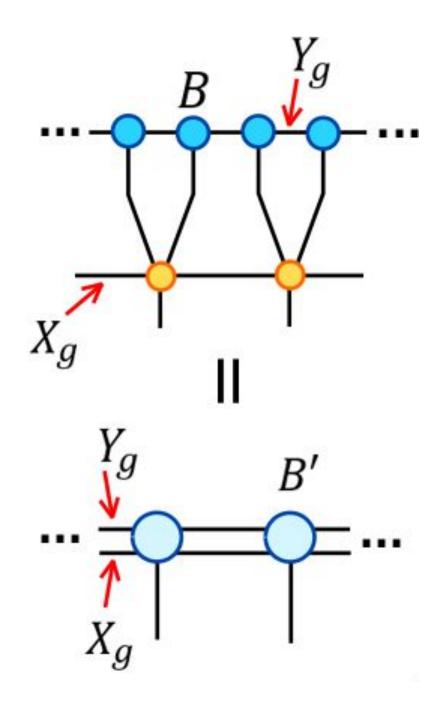






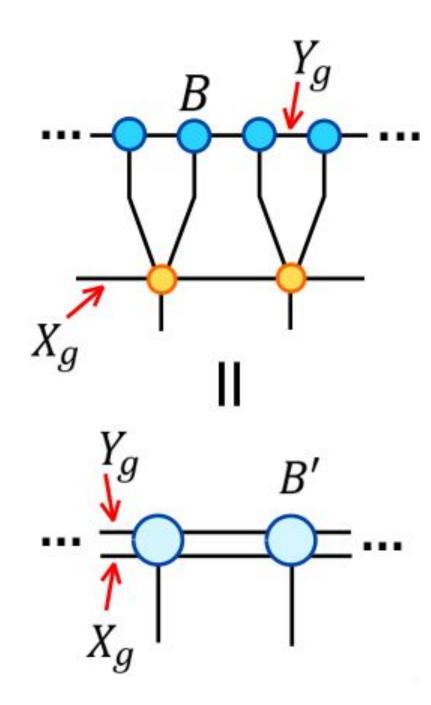


Since the coarse-grained MPS belongs to the same phase, rep $Y_g \otimes X_g$ must be equivalent to Y_g



Since the coarse-grained MPS belongs to the same phase, rep $Y_g \otimes X_g$ must be equivalent to Y_g

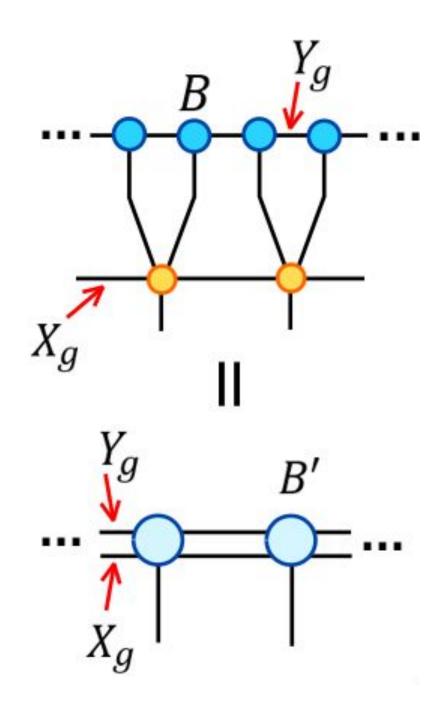
This implies X_g must be a linear rep.



Since the coarse-grained MPS belongs to the same phase, rep $Y_g \otimes X_g$ must be equivalent to Y_g

This implies X_g must be a linear rep.

This must be true for any projective rep Y_g



Since the coarse-grained MPS belongs to the same phase, rep $Y_g \otimes X_g$ must be equivalent to Y_g

This implies X_g must be a linear rep.

This must be true for any projective rep Y_g

MERA does not pick up a projective rep at least along the RG flow.

$$|\Psi_{\phi}^{\mathrm{fixed}}\rangle \equiv$$

$$v_{g} \quad v_{g}$$

$$u_{\phi}^{\mathrm{fixed}} \equiv$$

$$w_{\phi}^{\mathrm{fixed}} \equiv$$

No projective rep also at the fixed point

No projective rep also at the fixed point

MERA representation = RG flow + fixed point

No projective rep also at the fixed point

MERA representation = RG flow + fixed point

So no projective reps in the MERA representation of a gapped symmetry protected ground state.

Our argument does not apply to critical states.

Our argument does not apply to critical states.

Because we assumed an MPS representation of the ground states.

Our argument does not apply to critical states.

Because we assumed an MPS representation of the ground states.

Our argument does not apply to critical states.

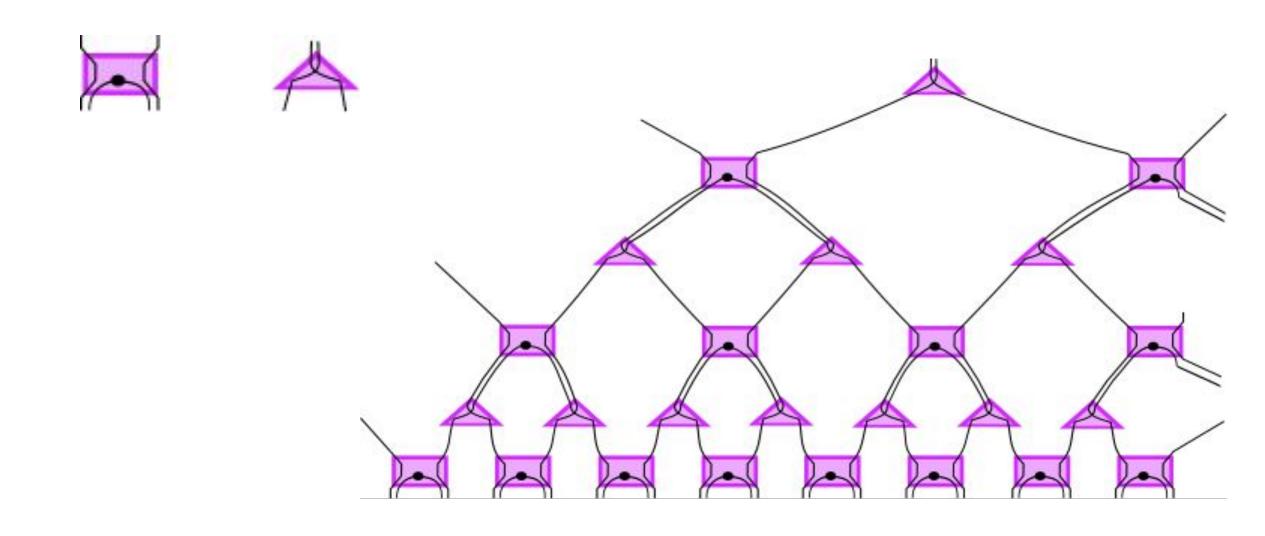
Because we assumed an MPS representation of the ground states.





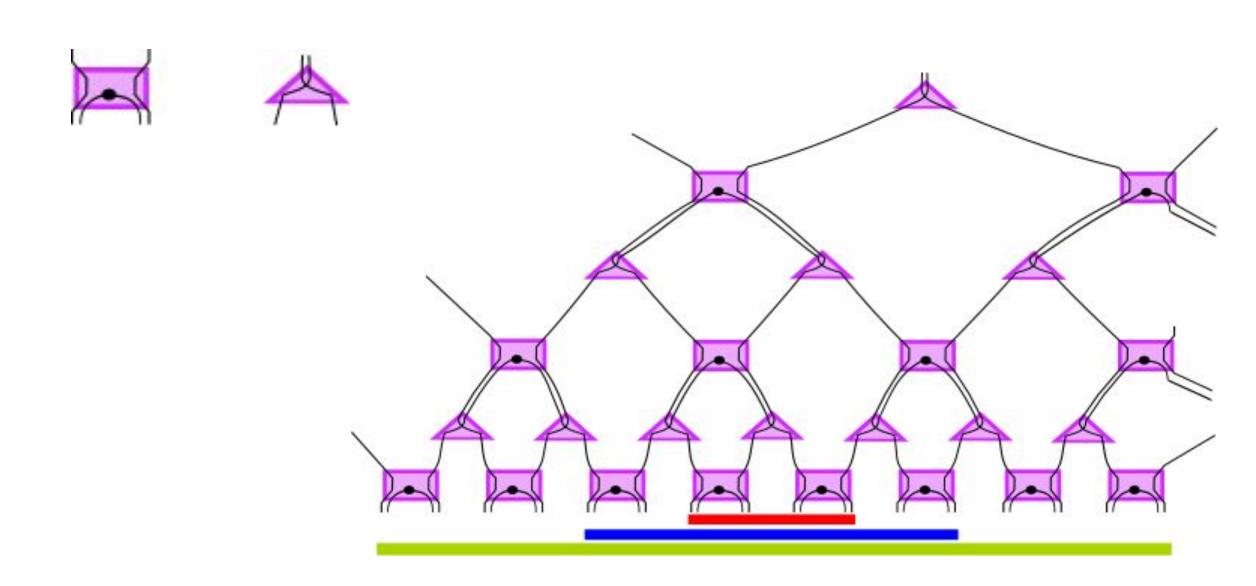
Our argument does not apply to critical states.

Because we assumed an MPS representation of the ground states.



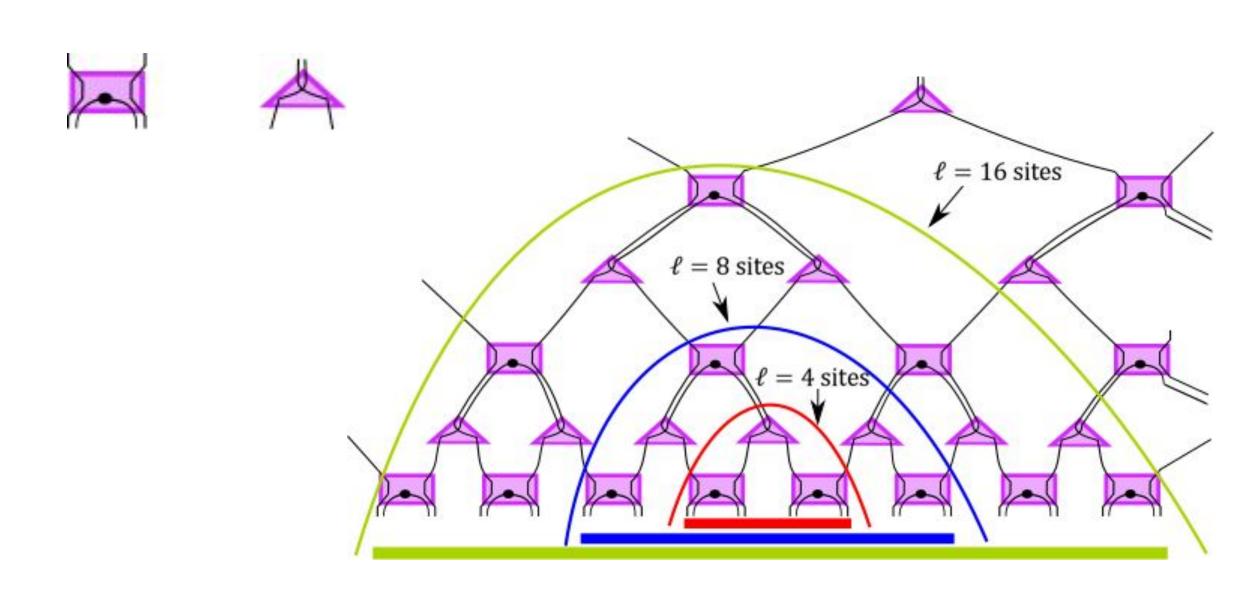
Our argument does not apply to critical states.

Because we assumed an MPS representation of the ground states.



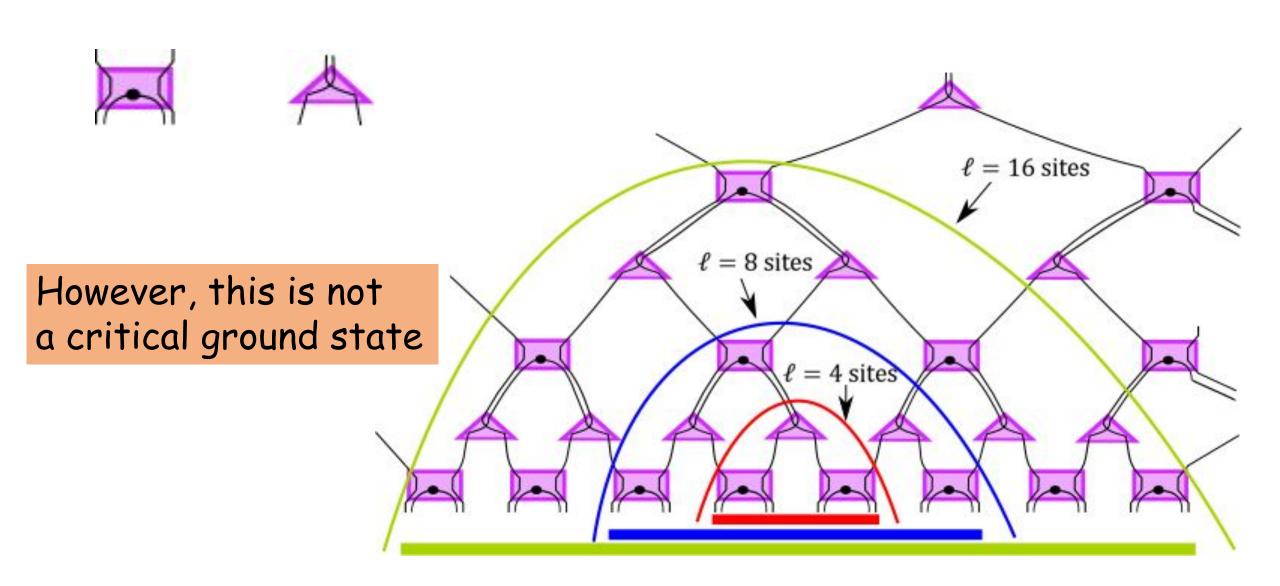
Our argument does not apply to critical states.

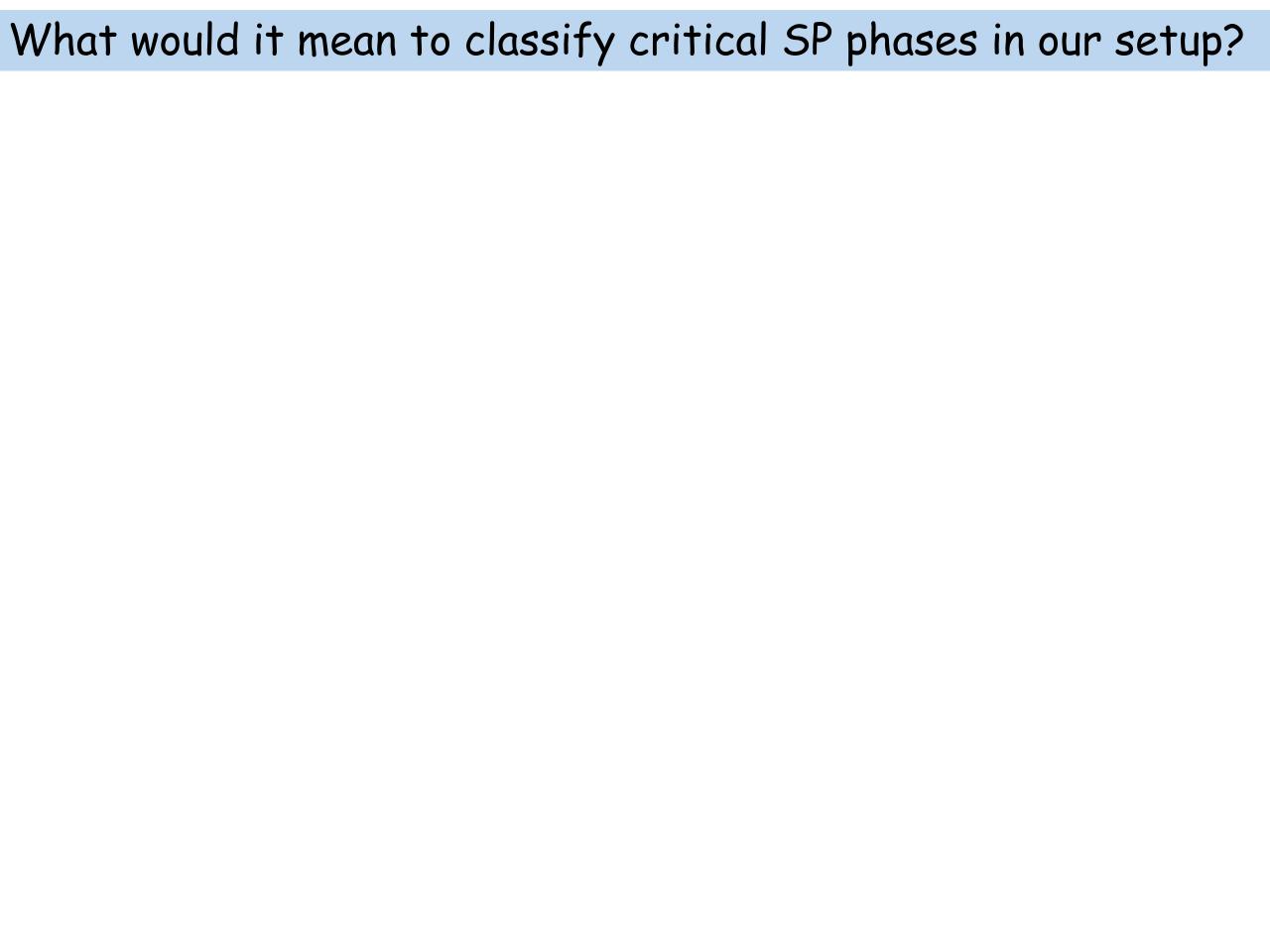
Because we assumed an MPS representation of the ground states.



Our argument does not apply to critical states.

Because we assumed an MPS representation of the ground states.





Entanglement renormalization is capable of reproducing critical fixed points (numerical evidence)

Entanglement renormalization is capable of reproducing critical fixed points (numerical evidence)

But there is no tight correspondence between MERA and critical ground states (unlike between MPS and gapped ground states)

Entanglement renormalization is capable of reproducing critical fixed points (numerical evidence)

But there is no tight correspondence between MERA and critical ground states (unlike between MPS and gapped ground states)

In fact, most states described by the MERA are unlikely to be ground states of local critical Hamiltonian.

Entanglement renormalization is capable of reproducing critical fixed points (numerical evidence)

But there is no tight correspondence between MERA and critical ground states (unlike between MPS and gapped ground states)

In fact, most states described by the MERA are unlikely to be ground states of local critical Hamiltonians.

So the problem would be to determine and classify all conformally invariant fixed points of symmetry protected entanglement renormalization.

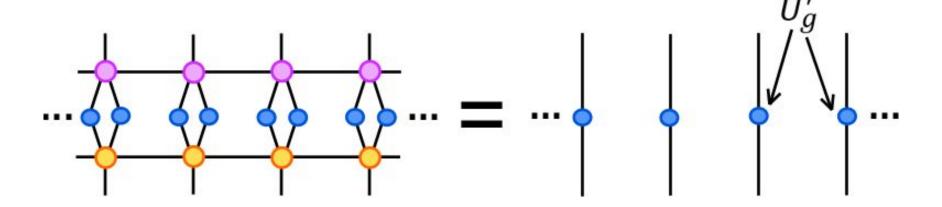
Entanglement renormalization is capable of reproducing critical fixed points (numerical evidence)

But there is no tight correspondence between MERA and critical ground states (unlike between MPS and gapped ground states)

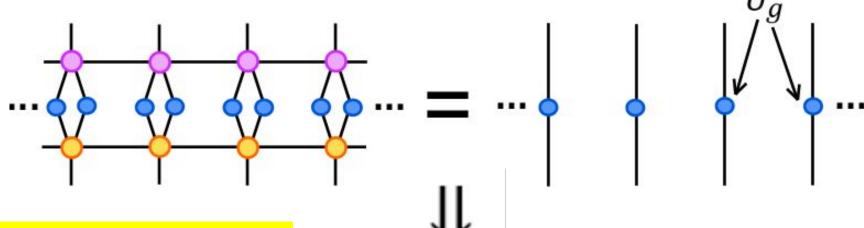
In fact, most states described by the MERA are unlikely to be ground states of local critical Hamiltonians.

So the problem would be to determine and classify all conformally invariant fixed points of symmetry protected entanglement renormalization.

$$\cdots \nearrow \qquad \bigvee_{U_g}^{U_g'^{\dagger}} = \bigvee_{X_g^{\dagger}}^{U_g'^{\dagger}} X_g$$



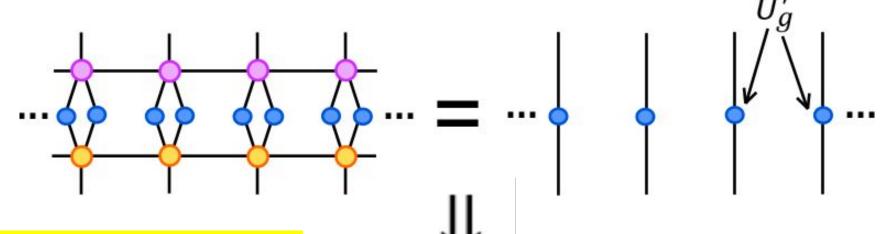
Symmetry protected Entanglement renormalization



Symmetry protected Entanglement renormalization

$$U_g^{\prime\dagger} = X_g^{\dagger} X_g^{\dagger}$$

 X_g can be **projective**

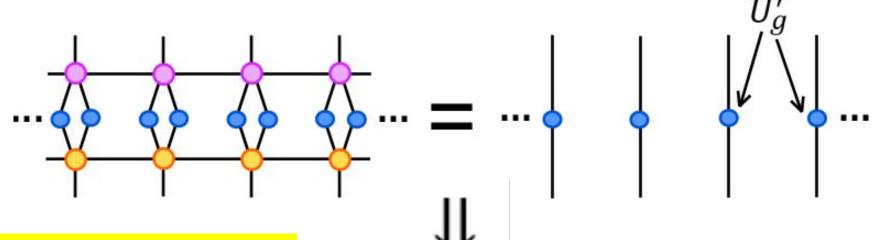


Symmetry protected Entanglement renormalization

$$U_g^{\prime\dagger} = X_g^{\dagger} X_g^{\dagger} X_g$$

 X_g can be **projective**

But X_g cannot be projective in a gapped phase.



Symmetry protected Entanglement renormalization

$$U_g^{\prime\dagger} = X_g^{\dagger} X_g^{\dagger}$$

 X_g can be projective

But X_g cannot be projective in a gapped phase.

What's next?

Determine and classify critical fixed points <-> distinct critical SP phases?