

# Honours Differential Equations

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# 1 ODE systems

## 1.1 Integrating factors

Consider linear DE of form

$$y' + P(x)y = Q(x)$$

Our integrating factor for this DE is:

$$I(x) = \exp\left(\int P(x)dx\right)$$

Then our solution is

$$y(x) = \frac{1}{I(x)} \int Q(x)I(x)dx + \frac{\alpha}{I(x)}$$

where here  $\alpha$  is a constant.

## 1.2 Change of variables

For higher order differential equations of form

$$y^{(n)} = F(y, y', \dots, y^{(n-1)}, t),$$

consider **change of variables**  $x_{i+1} = y^{(i)}$  for  $i \in \{0, 1, \dots, n-1\}$ . Then take derivative with respect to time to obtain a first order ODE system, which takes the form:

$$x'_j = F_j(t, x_1, \dots, x_n)$$

for  $j = 1, \dots, n$ . We either immediately write this as a matrix system or linearise near a critical point.

## 1.3 Existence and uniqueness for IVPs

An initial value problem (**IVP**) is defined as

$$\frac{dx}{dt} = f(x, t)$$

for **initial** condition  $x(t_0) = x_0$ . A solution  $x : I \rightarrow \mathbb{R}$  is a differentiable function that satisfies the IVP. Similarly for a first order system

$$x'_i = F_i(t, x_1, \dots, x_n)$$

to have a **unique** solution,  $F_i$  and  $\frac{\partial F_i}{\partial x_j}$  must be continuous in a region. Here  $i, j \in \{1, \dots, n\}$ .

## 1.4 Homogeneous systems

Now consider  $\mathbf{x}' = \mathbf{A}\mathbf{x}$ , where  $\mathbf{A}$  is a  $n \times n$  matrix.

Substituting  $\mathbf{x} = e^{rt}\boldsymbol{\xi}$  results in an eigenvector problem:

$$(\mathbf{A} - r_i \mathbf{I}_n)\boldsymbol{\xi}^{(i)} = \mathbf{0}.$$

If eigenvalues repeat we try  $\mathbf{x} = te^{rt}\boldsymbol{\xi} + e^{rt}\boldsymbol{\eta}$  which gives

$$(\mathbf{A} - r_i \mathbf{I}_n)\boldsymbol{\eta}^{(i)} = \boldsymbol{\xi}^{(i)}.$$

The **fundamental matrix** satisfies  $\boldsymbol{\Psi}' = \mathbf{A}\boldsymbol{\Psi}$  and contains the basis of solutions.

$$\boldsymbol{\Psi}(t) = [\mathbf{x}^{(1)}, \dots, \mathbf{x}^{(n)}]$$

for  $\mathbf{x}^{(i)} = \exp\{(r_i t)\}\boldsymbol{\xi}^{(i)}$  and hence our general solution is  $\mathbf{x} = \boldsymbol{\Psi}\mathbf{c}$ .

## 1.5 Non-homogeneous systems

Consider non-homogeneous ODE system:

$$\mathbf{x}' = \mathbf{A}\mathbf{x} + \mathbf{g}.$$

There are a couple of different approaches we can take to solve such a system.

- **Change of basis**

Let  $\mathbf{x} = \mathbf{T}\mathbf{y}$ , where  $\mathbf{T}$  is our eigenvector matrix from diagonalisation. So  $\mathbf{A} = \mathbf{T}\mathbf{D}\mathbf{T}^{-1}$ , and after some algebra we obtain:

$$\mathbf{y}' = \mathbf{D}\mathbf{y} + \mathbf{T}^{-1}\mathbf{g}$$

which can be solved by integrating factors. Finally revert back to  $\mathbf{x}$ .

- **Variation of parameters**

So  $\mathbf{x}_H = \mathbf{\Psi}\mathbf{c}$  solves the  $\mathbf{x}' = \mathbf{A}\mathbf{x}$ , where  $\mathbf{c}$  is a constant vector.

We then assume that the solution to our non-homogeneous system takes the form:

$$\mathbf{x} = \mathbf{\Psi}\mathbf{u}$$

for here  $\mathbf{u} = \mathbf{u}(t)$ . We then get  $\mathbf{\Psi}\mathbf{u}' = \mathbf{g}$ , which can be solved by eliminating variables and integrating.

- **Method of undetermined coefficients**

Our non-homogeneous ODE system has solutions of form:

$$\mathbf{x} = \mathbf{x}_H + \mathbf{x}_p$$

Solving the homogeneous ODE gives us  $\mathbf{x}_H$ .

On the other hand we just need to find a **particular solution**  $\mathbf{x}_p$  that satisfies our non-homogeneous ODE. Then our solution is complete.

Whilst the fastest, this method is not guaranteed to work.

## 1.6 Critical points & linearisation

Consider non-linear ODE system

$$x' = F(x, y),$$

$$y' = G(x, y).$$

We define  $\mathbf{x}^0 = \begin{bmatrix} x^0 \\ y^0 \end{bmatrix}$  as a **critical point** when  $F(\mathbf{x}^0) = G(\mathbf{x}^0) = 0$ .

Non-linear systems may then be linearised by Taylor expanding them around a critical point  $\mathbf{x}^0$ , and discarding higher order terms.

i.e. let  $\mathbf{u} = \begin{bmatrix} u_1 \\ u_2 \end{bmatrix}$  where  $u_1 = x - x^0$  and  $u_2 = y - y^0$ .

$$\therefore u'_1 = x'$$

$$\approx F(x^0, y^0) + \left( \frac{\partial F}{\partial x} \right)_{x^0} (x - x^0) + \left( \frac{\partial F}{\partial y} \right)_{y^0} (y - y^0)$$

$$\therefore u'_2 = y'$$

$$\approx G(x^0, y^0) + \left( \frac{\partial G}{\partial x} \right)_{x^0} (x - x^0) + \left( \frac{\partial G}{\partial y} \right)_{y^0} (y - y^0)$$

Then we end up with the following linear system:

$$\mathbf{u}' = \mathbf{A}\mathbf{u}$$

where  $\mathbf{A} = \begin{bmatrix} \partial F/\partial x & \partial F/\partial y \\ \partial G/\partial x & \partial G/\partial y \end{bmatrix}_{\mathbf{x}=\mathbf{x}^0}$  and is a  $2 \times 2$  Jacobian matrix.

Our critical points  $\mathbf{x}^0$  may also be classified:

Eigenvalues	Critical Points	Stability
$r_1 > r_2 > 0$	Node (source)	unstable
$r_1 < r_2 < 0$	Node (sink)	asympt. stable
$r_2 < 0 < r_1$	saddle	unstable
$r_1 = r_2 > 0$	Proper/Improper node	unstable
$r_1 = r_2 < 0$	Proper/Improper node	asympt. stable
$r_1, r_2 = \lambda \pm i\mu$ ( $\lambda > 0$ )	focus	unstable
$r_1, r_2 = \lambda \pm i\mu$ ( $\lambda < 0$ )	focus	asympt. stable
$r_1 = i\mu, r_2 = -i\mu$	center	stable

Linearisation preserves critical point behaviour **except** when eigenvalues are of  $r = \pm i\mu$  form, for which then classification is unknown.

## 1.7 Stability of critical points

**Stable** critical points  $\mathbf{x}^0$ : All solutions start and stay near  $\mathbf{x}^0$ .

$$\forall \epsilon > 0; \exists \delta > 0; \forall \mathbf{x}_{\text{solution}} \text{ to } \mathbf{x}' = \mathbf{F}(\mathbf{x}, t): \\ |\mathbf{x}(0) - \mathbf{x}^0| < \delta \implies |\mathbf{x}(t) - \mathbf{x}^0| < \epsilon \text{ for } \forall t \geq 0$$

**Attracting** critical points  $\mathbf{x}^0$ : All solutions tends to  $\mathbf{x}^0$ .

$$\forall \delta > 0 : |\mathbf{x}(0) - \mathbf{x}^0| < \delta \implies \lim_{t \rightarrow \infty} \mathbf{x}(t) = \mathbf{x}^0$$

**Asymptotically stable** critical points  $\mathbf{x}^0$ : Attracting **and** stable

## 1.8 Lyapunov's theory and limit cycles

In this section  $\dot{\mathbf{x}}$  means its first time derivative. So consider:

$$\dot{x} = F(x, y),$$

$$\dot{y} = G(x, y)$$

defined in  $\mathbb{R}^2$ . Let  $\mathbf{x}^0 \in D$  be a critical point.

The function  $E : D \subset \mathbb{R}^2 \rightarrow \mathbb{R}$  is a Lyapunov function where  $E(x^0, y^0) = 0$ , whenever it exists. Note that the time derivative of E is:

$$\frac{dE}{dt} = \frac{\partial E}{\partial x} F + \frac{\partial E}{\partial y} G$$

- Let  $E > 0$  for  $\forall \mathbf{x} \neq \mathbf{x}^0$ .  
If  $\frac{dE}{dt} \leq 0$  then  $\mathbf{x}^0$  is stable.  
If  $\frac{dE}{dt} < 0$  then  $\mathbf{x}^0$  is asymptotically stable.
- If every neighbourhood of  $\mathbf{x}^0$  contains  $\mathbf{x}^*$  such that  $E(\mathbf{x}^*) > 0$   
**and** if  $\frac{dE}{dt} > 0$  then  $\mathbf{x}^0$  is unstable.

Now **limit cycles** are defined as periodic solutions such that at least one other **non-closed trajectory** approaches the limit cycle as  $t \rightarrow \infty$ .

## 2 Fourier series

### 2.1 Real Fourier series

Let  $f(x)$  and  $f'(x)$  be **piecewise continuous** in  $[-L, L]$  with **period**  $2L$ .  
i.e.  $f(x) = f(x + 2L)$  for  $\forall x$ . Then the Fourier series for  $f(x)$  is

$$f_{FS}(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} \left( a_n \cos \frac{n\pi x}{L} + b_n \sin \frac{n\pi x}{L} \right).$$

The **convergence** of our Fourier series depends on the continuity of  $f(x)$ :

- If  $f(x)$  is continuous then  $f_{FS}(x) = f(x)$ .
- If  $f(\alpha)$  is discontinuous then at point  $\alpha$  we have

$$f_{FS}(\alpha) = \frac{f(\alpha^+) + f(\alpha^-)}{2}.$$

Note that  $f(x)$  is continuous at  $\alpha$  if  $f(\alpha) = \lim_{x \rightarrow \alpha} f(x)$  and we define:

$$f(\alpha^-) = \lim_{x \rightarrow \alpha^-} f(x)$$

and

$$f(\alpha^+) = \lim_{x \rightarrow \alpha^+} f(x),$$

i.e. limits from left and right respectively. It is important to also note that the derivative of a Fourier series is **not necessarily convergent**.

Now consider  $S_n = \sin \frac{n\pi x}{L}$  and  $C_n = \cos \frac{n\pi x}{L}$ . We then have the following **orthogonality relations**:

$$\langle S_n, S_m \rangle = \langle C_n, C_m \rangle = L\delta_{mn}$$

$$\langle S_n, C_m \rangle = 0$$

where we define the inner product as:

$$\langle u(x), v(x) \rangle = \int_{-L}^L u(x)v(x)dx$$

and use the following trigonometric identities:

$$\cos A \cos B = \frac{1}{2} [\cos(A - B) + \cos(A + B)]$$

$$\sin A \sin B = \frac{1}{2} [\cos(A - B) - \cos(A + B)]$$

$$\sin A \cos B = \frac{1}{2} [\sin(A + B) + \sin(A - B)].$$



Now integrating the following expression:

$$\int_{-L}^L \left( \frac{a_0}{2} + \sum_{n=1}^{\infty} \left( a_n \cos \frac{n\pi x}{L} + b_n \sin \frac{n\pi x}{L} \right) \right) dx = \int_{-L}^L f(x) dx$$

gives:

$$a_0 = \frac{1}{L} \int_{-L}^L f(x) dx.$$

Similarly:

$$a_n = \frac{1}{L} \int_{-L}^L \cos \frac{n\pi x}{L} f(x) dx$$

$$b_n = \frac{1}{L} \int_{-L}^L \sin \frac{n\pi x}{L} f(x) dx.$$

Note that  $\delta_{mn}$  is the **Kronecker delta** and is defined as:

$$\delta_{mn} = \begin{cases} 1 & m = n \\ 0 & m \neq n \end{cases}$$

Furthermore notice that:

- The Fourier series of **even** functions contains only **cosines**.
- The Fourier series of **odd** functions contains only **sines**.

**Even** functions are defined  $f(-x) = f(x)$ , and:

$$\int_{-L}^L f_{\text{even}} dx = 2 \int_0^L f_{\text{even}} dx.$$

Similarly **odd** functions are defined  $f(-x) = -f(x)$ , and:

$$\int_{-L}^L f_{\text{odd}} dx = 0.$$

We can also extend a function defined in  $[0, L]$  in several ways:

1. Define even function:

$$g(x) = \begin{cases} f(x) & x \in [0, L] \\ f(-x) & x \in (-L, 0) \end{cases}$$

where its Fourier expansion is a cosine series.

2. Define odd function:

$$h(x) = \begin{cases} f(x) & x \in (0, L) \\ 0 & x = 0, L \\ -f(-x) & x \in (-L, 0) \end{cases}$$

where its Fourier expansion is a sine series.

## 2.2 Complex Fourier series

Expanding  $f(x)$  defined in  $[-L, L]$  with period  $2L$ :

$$f_{FS}(x) = \sum_{n=-\infty}^{\infty} c_n \exp\left(\frac{in\pi}{L}x\right)$$

using Euler's formula  $e^{i\theta} = \cos \theta + i \sin \theta$ . Its coefficients are:

$$c_n = \frac{1}{2L} \int_{-L}^L \exp\left(-\frac{in\pi}{L}x\right) f(x) dx$$

for  $\forall n \in \mathbb{Z}$  and:

$$c_n = \begin{cases} (a_n - ib_n)/2 & n > 0 \\ (a_0)/2 & n = 0 \\ (a_n + ib_n)/2 & n < 0. \end{cases}$$

Here we define the **inner product** for complex functions as

$$\langle f, g \rangle = \int_{\alpha}^{\beta} f^*(x) g(x) dx$$

where  $f^*(x)$  is the complex conjugate of  $f(x)$ . Then:

$$\begin{aligned} \left\langle \exp\left(\frac{im\pi}{L}x\right), \exp\left(\frac{in\pi}{L}x\right) \right\rangle &= \int_{-L}^L \exp\left(-\frac{im\pi}{L}x\right) \exp\left(\frac{in\pi}{L}x\right) dx \\ &= 2L \delta_{mn} \end{aligned}$$

and since  $f(x) = f_{FS}(x)$  we obtain our formula.

## 2.3 Parseval's theorem

Parseval's theorem states that given a periodic  $f(x)$  with convergent Fourier series we have that

$$\begin{aligned} \langle f, f \rangle &= \int_{-L}^L |f(x)|^2 dx \\ &= 2L \sum_{n=-\infty}^{\infty} |c_n|^2 \\ &= L \left[ \frac{|a_0|^2}{2} + \sum_{n=1}^{\infty} (|a_n|^2 + |b_n|^2) \right] \end{aligned}$$

and is derived by orthogonality.

### 3 PDEs

#### 3.1 Separation of variables

The only methodology considered is separation of variables. So for PDE:

$$\hat{D}[u(x_1, \dots, x_n)] = 0$$

where  $\hat{D}$  is our differential operator, we look for solutions of form:

$$u(x_1, \dots, x_n) = X_1(x_1) \cdots X_n(x_n)$$

subject to **initial** and **boundary** conditions.

#### 3.2 Heat equation

The heat equation is an equation of the following form:

$$\frac{\partial u}{\partial t} = \alpha^2 \frac{\partial^2 u}{\partial x^2}.$$

where  $\alpha^2$  is the thermal diffusivity constant.

##### 3.2.1 Standard boundary conditions

We firstly define:

- **Initial condition:**  $u(x, 0) = f(x)$  for  $0 \leq x \leq L$
- **Boundary condition:**  $u(0, t) = u(L, t) = 0$  for  $\forall t > 0$

Let solutions be of form:

$$\begin{aligned} u(x, t) &= X(x) \cdot T(t) \\ \therefore X(x) \cdot \dot{T}(t) &= \alpha^2 X''(x) \cdot T(t) \end{aligned}$$

Only a constant function may satisfy the first equality:

$$\frac{1}{\alpha^2} \frac{\dot{T}}{T} = \frac{X''}{X} = -\lambda.$$

Writing this as two ODEs:

$$\begin{aligned} \dot{T} + \alpha^2 \lambda T &= 0 \\ X'' + \lambda X &= 0. \end{aligned}$$

The first one we can directly integrate, yielding:

$$T(t) = a_1 \exp(-\alpha^2 \lambda t).$$

The second ODE is a spring system, hence it has solution of form:

$$X(x) = b_1 \cos \lambda^{1/2} x + b_2 \sin \lambda^{1/2} x.$$

However this time before proceeding we need to consider boundary conditions:

$$X(0) = X(L) = 0.$$

We find  $X(0) = b_1 = 0$  and  $X(L) = b_2 \sin \lambda^{1/2} L = 0$ .

The second equation implies that  $\lambda$  must of the following form:

$$\lambda_n = \left(\frac{n\pi}{L}\right)^2 \quad \text{for } \forall n \in \mathbb{N}$$

and so

$$X'' + \lambda X = 0 \implies X_n = b_2 \sin \lambda_n^{1/2} x.$$

Since  $\lambda$  is discretised:

$$\therefore T_n = a_1 \exp(-\alpha^2 \lambda_n t).$$

Our general solution must then be:

$$u(x, t) = \sum_{n=1}^{\infty} c_n \exp(-\alpha^2 \lambda_n t) \sin \lambda_n^{1/2} x$$

where  $\lambda_n = \left(\frac{n\pi}{L}\right)^2$ . Using initial condition  $u(x, 0) = f(x)$ :

$$f(x) = \sum_{n=1}^{\infty} c_n \sin \lambda_n^{1/2} x$$

and we recognise this as an odd Fourier series with period  $2L$ .

$$\begin{aligned} \therefore \int_{-L}^L \sin(\lambda_n^{1/2} x) f(x) dx &= \sum_{n=1}^{\infty} c_n \int_{-L}^L \left(\sin(\lambda_n^{1/2} x)\right)^2 dx \\ \therefore 2 \int_0^L \sin(\lambda_n^{1/2} x) f(x) dx &= c_n L \end{aligned}$$

The final step we split the integration range and use  $x = -x^*$ .

$$\therefore c_n = \frac{2}{L} \int_0^L \sin(\lambda_n^{1/2} x) f(x) dx$$

This is fine because we can extend  $u(x, t)$  via reflection for negative  $x$ .

### 3.2.2 Fixed boundary temperatures

We reconsider the heat equation:

$$\frac{\partial u}{\partial t} = \alpha^2 \frac{\partial^2 u}{\partial x^2}$$

but now with the following **non-homogeneous boundary conditions**:

- $u(0, t) = T_1$
- $u(L, t) = T_2$
- $u(x, 0) = f(x)$

Physically our rod has fixed boundary temperatures, namely  $T_1$  and  $T_2$ .

We approach this problem with a change of variables:

$$v(x) = \lim_{t \rightarrow \infty} u(x, t).$$

Using our boundary conditions  $v$  must be linear:

$$\therefore v(x) = \frac{T_2 - T_1}{L}x + T_1$$

since  $v'' = 0$ ,  $v(0) = T_1$  and  $v(L) = T_2$ . We then deduce that:

$$u(x, t) = v(x) + \omega(x, t)$$

for  $\omega(x, t)$  satisfies the same heat equation with initial conditions:

- $\omega(0, t) = \omega(L, t) = 0$
- $\omega(x, 0) = f(x) - v(x)$

Recognising this as our initial example:

$$\omega(x, t) = \sum_{n=1}^{\infty} c_n \exp(-\alpha^2 \lambda_n t) \sin \lambda_n^{1/2} x$$

where again  $\lambda_n = \left(\frac{n\pi}{L}\right)^2$  and because  $\omega(x, t)$  is a Fourier series with period  $2L$ :

$$c_n = \frac{2}{L} \int_0^L \sin(\lambda_n^{1/2} x) \{f(x) - v(x)\} dx.$$

**3.2.3 Insulated rod ends**

For the final example we consider:

$$\frac{\partial u}{\partial t} = \alpha^2 \frac{\partial^2 u}{\partial x^2}$$

and define the following conditions:

- $\frac{\partial}{\partial x} u(0, t) = \frac{\partial}{\partial x} u(L, t) = 0$
- $u(x, 0) = f(x)$

We begin again with a separation of variables:

### 3.3 Wave equation

### 3.4 Laplace's equation

Laplace's equation takes the form  $\nabla^2 u = 0$ . In two dimensions:

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0$$

and we only consider boundary conditions. (Dirichlet conditions)

#### 3.4.1 Rectangular boundary conditions

We open with the following example:

- **Boundary for y:**  $u(x, 0) = u(x, b) = 0$
- **Boundary for x:**  $u(0, y) = 0$  and  $u(a, y) = f(y)$

where  $x \in [0, a]$  and  $y \in [0, b]$ . Begin by separation of variables:

$$\begin{aligned} u(x, y) &= X(x) \cdot Y(y) \\ \therefore \frac{X''}{X} &= -\frac{Y''}{Y} = \lambda. \end{aligned}$$

Recognising the previous statement as two ODEs:

$$X'' - \lambda X = 0 \text{ for } X(0) = 0$$

$$Y'' + \lambda Y = 0 \text{ for } Y(0) = Y(b) = 0$$

The second ODE we have already solved in the heat equation. It has solution:

$$Y_n = a_1 \sin(\lambda_n^{1/2} y) \text{ for } \lambda_n = \frac{n^2 \pi^2}{b^2}.$$

The first ODE has solutions of form:

$$X_n = a_2 \cosh(\lambda_n^{1/2} x) + a_3 \sinh(\lambda_n^{1/2} x)$$

where these are the hyperbolic functions:

$$\sinh x = \frac{1}{2}(e^x - e^{-x})$$

$$\cosh x = \frac{1}{2}(e^x + e^{-x}).$$

Using our boundary condition  $X(0) = 0$  gives:

$$X_n = a_3 \sinh(\lambda_n^{1/2} x).$$

Now putting all of this together we get:

$$u(x, y) = \sum_{n=1}^{\infty} c_n \sinh(\lambda_n^{1/2} x) \sin(\lambda_n^{1/2} y)$$



To find coefficients  $c_n$  we use  $u(a, y) = f(y)$ .

$$\therefore f(y) = \sum_{n=1}^{\infty} c_n \sinh(\lambda_n^{1/2} a) \sin(\lambda_n^{1/2} y)$$

Since we have a Fourier series with period  $2b$ :

$$\begin{aligned} \int_{-b}^b \sin(\lambda_n^{1/2} y) f(y) dy &= \sum_{n=1}^{\infty} c_n \sinh(\lambda_n^{1/2} a) \int_{-b}^b \sin(\lambda_n^{1/2} y) dy \\ &= c_n \sinh(\lambda_n^{1/2} a) \cdot b \end{aligned}$$

We can split the first integral to give us:

$$c_n = \frac{2}{b \sinh(\lambda_n^{1/2} a)} \int_0^b \sin(\lambda_n^{1/2} y) f(y) dy$$

where  $\lambda_n = \left(\frac{n\pi}{b}\right)^2$  and our solution is complete.

### 3.4.2 Circular boundary conditions

Now we solve Laplace's equation

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0$$

but with a circular boundary. In polar coordinates  $(r, \theta)$ :

- $u(a, \theta) = f(\theta)$
- $u(r, \theta)$  is bounded

where  $a$  is the radius of our circle and  $\theta \in [0, 2\pi]$ . Since  $u = u(x, y)$ :

$$u'_\theta = u'_x x'_\theta + u'_y y'_\theta$$

$$u''_{\theta\theta} = (u''_{xx} x'_\theta + u''_{xy} y'_\theta) x'_\theta + u'_x x''_{\theta\theta} + (u''_{yy} y'_\theta + u''_{xy} x'_\theta) y'_\theta + u'_y y''_{\theta\theta}$$

$$u'_r = u'_x x'_r + u'_y y'_r$$

$$u''_{rr} = (u''_{xx} x'_r + u''_{xy} y'_r) x'_r + u'_x x''_{rr} + (u''_{xy} x'_r + u''_{yy} y'_r) y'_r + u'_y y''_{rr}$$

and here we have used the chain rule.

Applying these derivatives we obtain the following equation:

$$\frac{\partial^2 u}{\partial r^2} + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2} + \frac{1}{r} \frac{\partial u}{\partial r} = 0.$$

Using separation of variables:

$$u(r, \theta) = R(r)\Theta(\theta)$$

$$\therefore r^2 \frac{R''}{R} + r \frac{R'}{R} = -\frac{\ddot{\Theta}}{\Theta} = \lambda$$

where  $\lambda$  is our separation constant.

$$\therefore \ddot{\Theta} + \lambda\Theta = 0$$

$$\therefore r^2 R'' + rR' = \lambda R$$

## 4 Sturm-Liouville theory

### 4.1 Regular S-L problems

Sturm-Liouville theory is a general theory for 2nd order ODEs.

Consider the following eigenvalue ODE:

$$-\frac{d}{dx} \left( p(x) \frac{dy}{dx} \right) + q(x)y = \lambda r(x)y$$

where  $r(x)$  is our weight function. We define the following boundary conditions:

1.  $a_1 y(0) + a_2 y'(0) = 0$
2.  $b_1 y(1) + b_2 y'(1) = 0$ .

This is a **regular Sturm-Liouville** problem, where  $p(x)$ ,  $p'(x)$ ,  $q(x)$ ,  $r(x)$  are continuous functions and  $p(x)$ ,  $r(x)$  are strictly positive functions for  $\forall x \in [0, 1]$ .

**Eigenvalues**  $\lambda_n$  yield **eigenfunctions**  $\phi_n(x)$  which are nontrivial solutions to our S-L problem. Important consequences include:

- Eigenvalues  $\lambda_n$  of a S-L problem are **real**.

Furthermore each eigenvalue corresponds to one eigenfunction.

- Eigenfunctions  $\phi_n(x)$  are orthogonal:

$$\langle \phi_m, \phi_n \rangle = \int_0^1 r(x) \phi_m(x) \phi_n(x) dx = \delta_{mn}$$

in Hilbert space  $L^2([0, 1], r(x)dx)$ .

Note that our eigenfunctions are orthonormal and so we define:

$$\phi_n(x) = k_n y_n(x)$$

where  $k_n$  is our scale factor. Since  $\langle \phi_n, \phi_n \rangle = 1$ :

$$\therefore \int_0^1 r(x) k_n^2 y_n^2(x) dx = 1$$

and so we have that:

$$\begin{aligned} k_n &= \frac{1}{\sqrt{\langle y_n, y_n \rangle}} \\ &= \left( \int_0^1 r(x) y_n^2(x) dx \right)^{-1/2}. \end{aligned}$$

#### 4.1.1 General 2nd order ODEs

Consider the following general 2nd order eigenvalue ODE:

$$-P(x)\frac{d^2y}{dx^2} - \omega(x)\frac{dy}{dx} + q(x)y = \lambda r(x)y.$$

Multiply this by the following integrating factor:

$$F(x) = \exp\left[\int_0^x \frac{\omega(s) - p'(s)}{p(s)} ds\right]$$

yields an ODE of S-L form:

$$-\frac{d}{dx}\left[F(x)P(x)\frac{dy}{dx}\right] + F(x)q(x)y = \lambda F(x)r(x)y.$$

#### 4.1.2 Lagrange's identity

Our previous definition is motivated by the **Lagrange's identity**:

$$\begin{aligned}\langle \mathcal{L}[u], v \rangle - \langle u, \mathcal{L}[v] \rangle &= -\left[p\left(u'v^* - u(v^*)'\right)\right]_0^1 \\ &= -\left[p(x)\left(\frac{du}{dx} \cdot v^* - u \cdot \frac{dv^*}{dx}\right)\right]_0^1\end{aligned}$$

where  $u = u(x)$ ,  $v = v(x)$  are complex functions and

$$\mathcal{L}[u] = -\frac{d}{dx}\left[p(x)\frac{du}{dx}\right] + q(x)u.$$

Here the inner product is defined as

$$\langle u, v \rangle = \int_0^1 uv^* dx$$

and we have integrated by parts using the following identities:

$$[pu'v^*]' = (pu')'v^* + pu'(v^*)'$$

$$[pu(v^*)']' = (p(v^*)')'u + pu'(v^*)'.$$

For a S-L problem its properties follow from:

$$\langle \mathcal{L}[u], v \rangle = \langle u, \mathcal{L}[v] \rangle$$

where functions  $u$  and  $v$  satisfy its boundary conditions.

### 4.1.3 Series expansion

Now the set of orthonormal eigenfunctions  $\{\phi_n(x)\}$  from a S-L problem with boundary conditions may be used to expand function  $f(x)$ :

$$f_\phi(x) = \sum_{n=1}^{\infty} c_n \phi_n(x)$$

for  $\forall x \in [0, 1]$ . Integrating this on both sides:

$$\begin{aligned} \int_0^1 r(x) \phi_m(x) f(x) dx &= \int_0^1 r(x) \phi_m(x) \sum_{n=1}^{\infty} c_n \phi_n(x) dx \\ &= \sum_{n=1}^{\infty} c_n \int_0^1 r(x) \phi_m(x) \phi_n(x) dx \\ &= \sum_{n=1}^{\infty} c_n \delta_{mn} \\ &= c_m \end{aligned}$$

and so we get a general formula for our coefficients:

$$c_n = \int_0^1 r(x) \phi_n(x) f(x) dx.$$

If  $f(x)$  and  $f'(x)$  are piecewise continuous on  $x \in [0, 1]$  then:

$$\forall x \in (0, 1); \sum_{n=1}^{\infty} c_n \phi_n(x) = \frac{f(x^-) + f(x^+)}{2}$$

where these are the left and right handed limits.

## 4.2 Non-homogeneous S-L problems

Consider:

$$\mathcal{L}[y] = \mu r(x)y + f(x)$$

where  $f(x)$  is our non-homogeneous term, and that we define

$$\mathcal{L}[y] = -\frac{d}{dx} \left[ P(x) \frac{dy}{dx} \right] + q(x)y.$$

Given that this ODE satisfies the S-L boundary conditions:

$$\mathcal{L}[y] = \lambda r(x)y$$

is solved by a sum of eigenfunctions

$$y(x) = \sum_{n=1}^{\infty} b_n \phi_n(x).$$

Substituting this into our original ODE gives:

$$\begin{aligned} \sum_{n=1}^{\infty} b_n \mathcal{L}[\phi_n(x)] &= r(x) \sum_{n=1}^{\infty} b_n \lambda_n \phi_n(x) \\ &= \mu r(x) \sum_{n=1}^{\infty} b_n \phi_n(x) + f(x). \end{aligned}$$

Rearranging then gives:

$$\frac{f(x)}{r(x)} = \sum_{n=1}^{\infty} b_n (\lambda_n - \mu) \phi_n(x).$$

Now we can also expand  $\frac{f(x)}{r(x)}$  in terms of our eigenfunctions:

$$\frac{f(x)}{r(x)} = \sum_{n=1}^{\infty} c_n \phi_n(x)$$

with

$$\begin{aligned} c_n &= \int_0^1 r(x) \phi_n(x) \frac{f(x)}{r(x)} dx \\ &= \int_0^1 \phi_n(x) f(x) dx \end{aligned}$$

and so equating yields the following relation

$$b_n = \frac{c_n}{\lambda_n - \mu}.$$

### 4.3 Singular S-L problems

quite a lot of theory not covered in this course

general definition of singular s-l problems

bessel's equation (order 0 example)

conditions for singular problems

#### 4.4 S-L Series solutions



## 5 Laplace transforms