Honours Differential Equations Assignments

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Winter 2023

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1.

1.

1. Consider $f(x) = \frac{x^2}{2}$ where $0 \le x < L$.

Part (a) wants us to find its Fourier series with period 2L.

Define an extension of our function f as g:

$$g(x) = \begin{cases} f(x) & x \in [0, L) \\ f(-x) & x \in [-L, 0) \end{cases}$$

and therefore g has period 2L. Since g is also an even function its Fourier series can only consist of cosine terms.

$$\therefore g_{FS}(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos \frac{n\pi x}{L}$$

Its coefficients are:

$$a_{0} = \frac{1}{L} \int_{-L}^{L} g(x) dx$$

$$= \frac{1}{L} \int_{0}^{L} f(x) dx + \frac{1}{L} \int_{-L}^{0} f(-x) dx$$

$$= \frac{1}{L} \int_{0}^{L} f(x) dx - \frac{1}{L} \int_{0}^{-L} f(-x) dx$$

$$= \frac{1}{L} \int_{0}^{L} f(x) dx + \frac{1}{L} \int_{0}^{L} f(x^{*}) dx^{*}$$

$$= \frac{2}{L} \int_{0}^{L} f(x) dx$$

where we define $x = -x^*$ and similarly

$$a_n = \frac{1}{L} \int_{-L}^{L} \cos \frac{n\pi x}{L} g(x) dx$$

$$= \frac{1}{L} \int_{0}^{L} \cos \frac{n\pi x}{L} f(x) dx + \frac{1}{L} \int_{-L}^{0} \cos \frac{n\pi x}{L} f(-x) dx$$

$$= \frac{2}{L} \int_{0}^{L} \cos \frac{n\pi x}{L} f(x) dx.$$

Calculating this explicitly:

$$a_0 = \frac{2}{L} \int_0^L f(x) dx$$
$$= \frac{2}{L} \int_0^L \frac{1}{2} x^2 dx$$
$$= \frac{1}{L} \left[\frac{x^3}{3} \right]_0^L$$
$$= \frac{1}{3} L^2$$

and

$$a_n = \frac{2}{L} \int_0^L \cos \frac{n\pi x}{L} f(x) dx$$

$$= \frac{2}{L} \int_0^L \cos \frac{n\pi x}{L} \frac{1}{2} x^2 dx$$

$$= \frac{1}{L} \int_0^L x^2 \cos \frac{n\pi x}{L} dx$$

$$= \frac{1}{n\pi} \left[x^2 \sin \frac{n\pi x}{L} \right]_0^L - \frac{2}{n\pi} \int_0^L x \sin \frac{n\pi x}{L} dx$$

$$= \frac{L^2}{n\pi} \sin n\pi - \frac{2}{n\pi} \int_0^L x \sin \frac{n\pi x}{L} dx$$

$$= \frac{L^2}{n\pi} \sin n\pi - \frac{2}{n\pi} \left(-\frac{L^2}{n\pi} \cos n\pi + \left(\frac{L}{n\pi} \right)^2 \sin n\pi \right)$$

$$= \left(\frac{L^2}{n\pi} - \frac{2L^2}{(n\pi)^2} \right) \sin n\pi + \left(\frac{2L^2}{(n\pi)^2} \right) \cos n\pi$$

$$= (-1)^n \frac{2L^2}{(n\pi)^2}$$

where we have integrated by parts, and $\sin n\pi = 0$ for $\forall n \in \mathbb{N}$. Putting all this together we get our Fourier series for f:

$$\therefore f_{FS}(x) = \frac{L^2}{6} + \sum_{n=1}^{\infty} \left((-1)^n \frac{2L^2}{(n\pi)^2} \cos \frac{n\pi x}{L} \right)$$

and is valid for $\forall x \in [0, L)$.

For part (b) we want $f_{FS}(0)$. By Fourier's convergence theorem:

$$f_{FS}(0) = \frac{L^2}{6} + \sum_{n=1}^{\infty} \left((-1)^n \frac{2L^2}{(n\pi)^2} \right).$$

For part (c):

$$f(0) = \frac{L^2}{6} + \sum_{n=1}^{\infty} \left((-1)^n \frac{2L^2}{(n\pi)^2} \cos \frac{n\pi x}{L} \right)$$

$$= 0$$

$$\therefore \frac{L^2}{6} + \frac{2L^2}{\pi^2} \sum_{n=1}^{\infty} \frac{(-1)^n}{n^2} = 0$$

$$\therefore \sum_{n=1}^{\infty} \frac{(-1)^n}{n^2} = -\frac{\pi^2}{12}$$

2. For part (a) we given the solution to Laplace's equation:

$$u(r,\theta) = \frac{a_0}{2} + \sum_{n=1}^{\infty} r^n (a_n \cos n\theta + b_n \sin n\theta)$$

find expressions for coefficients a_0 , a_n and b_n .

Because our solution has period 2π , using the Euler-Fourier formulae:

$$a_0 = \frac{1}{\pi} \int_{-\pi}^{\pi} f(\theta) d\theta$$

$$a_n = \frac{1}{\pi} \int_{-\pi}^{\pi} \cos(n\theta) f(\theta) d\theta$$

and

$$b_n = \frac{1}{\pi} \int_{-\pi}^{\pi} \sin(n\theta) f(\theta) d\theta.$$

For part (b) we begin with the following expression:

$$u(r,\theta) = \frac{1}{2\pi} \int_{-\pi}^{\pi} f(\psi)d\psi$$
$$+ \frac{1}{\pi} \sum_{n=1}^{\infty} r^n \left(\cos(n\theta) \int_{-\pi}^{\pi} \cos(n\psi) f(\psi)d\psi + \sin(n\theta) \int_{-\pi}^{\pi} \sin(n\psi) f(\psi)d\psi \right)$$

Now since $e^{in(\theta-\psi)} = \cos n(\theta-\psi) + i\sin n(\theta-\psi)$:

$$Re(e^{in(\theta-\psi)}) = \cos n(\theta-\psi)$$

= $\cos(n\theta)\cos(n\psi) + \sin(n\theta)\sin(n\psi)$

$$\therefore \int_{-\pi}^{\pi} Re(e^{in(\theta-\psi)}) f(\psi) d\psi = \int_{-\pi}^{\pi} (\cos(n\theta)\cos(n\psi) + \sin(n\theta)\sin(n\psi)) f(\psi) d\psi$$
$$= \cos(n\theta) \int_{-\pi}^{\pi} \cos(n\psi) f(\psi) d\psi$$
$$+ \sin(n\theta) \int_{-\pi}^{\pi} \sin(n\psi) f(\psi) d\psi$$

Therefore our original expression becomes:

$$u(r,\theta) = \frac{1}{2\pi} \int_{-\pi}^{\pi} f(\psi) d\psi + \frac{1}{\pi} \sum_{n=1}^{\infty} \int_{-\pi}^{\pi} r^n Re(e^{in(\theta - \psi)}) f(\psi) d\psi$$

Taking the real component of each element in a sum is equivalent to taking the real component of the overall sum:

$$\therefore u(r,\theta) = \frac{1}{\pi} Re \left(\frac{1}{2} \int_{-\pi}^{\pi} f(\psi) d\psi + \sum_{n=1}^{\infty} \int_{-\pi}^{\pi} r^n e^{in(\theta - \psi)} f(\psi) d\psi \right)$$

Then by the linearity of integrals:

$$\therefore u(r,\theta) = \frac{1}{\pi} Re \left(\int_{-\pi}^{\pi} \left[\frac{1}{2} + \sum_{n=1}^{\infty} r^n e^{in(\theta - \psi)} \right] f(\psi) d\psi \right)$$

Finally for part (c) we have that

$$\sum_{n=0}^{\infty} q^n = \frac{1}{1-q}$$

and let:

$$q = re^{i(\theta - \psi)}$$

= $r(\cos(\theta - \psi) + i\sin(\theta - \psi))$.

Now |q| < 1 since we defined r < 1 and $e^{i2\pi} = 1$. Returning to our sum:

$$\sum_{n=1}^{\infty} q^n = \frac{1}{1-q} - 1$$
$$= \frac{q}{1-q}.$$

Furthermore:

$$\frac{1}{2} + \sum_{n=1}^{\infty} q^n = \frac{1}{2} \frac{1+q}{1-q}$$

$$= \frac{1}{2} \frac{1+r(\cos(\theta-\psi)+i\sin(\theta-\psi))}{1-r(\cos(\theta-\psi)+i\sin(\theta-\psi))}$$

$$= \frac{1}{2} \left[\frac{(1+rc)+i(rs)}{(1-rc)+i(-rs)} \right] \times \frac{1-rc+i(rs)}{1-rc+i(rs)}$$

$$= \frac{1}{2} \frac{1-r^2+i(2rs)}{1+r^2-2rc}$$

$$\therefore Re(\frac{1}{2} + \sum_{n=1}^{\infty} q^n) = \frac{1-r^2}{1+r^2-2rc}$$

Finally we have that:

$$\begin{split} u(r,\theta) &= \frac{1}{\pi} Re \left(\int_{-\pi}^{\pi} \left[\frac{1}{2} + \sum_{n=1}^{\infty} r^n e^{in(\theta - \psi)} \right] f(\psi) d\psi \right) \\ &= \frac{1}{\pi} \left(\int_{-\pi}^{\pi} Re \left[\frac{1}{2} + \sum_{n=1}^{\infty} r^n e^{in(\theta - \psi)} \right] f(\psi) d\psi \right) \\ &= \frac{1}{2\pi} \left(\int_{-\pi}^{\pi} \frac{1 - r^2}{1 + r^2 - 2r \cos(\theta - \psi)} f(\psi) d\psi \right). \end{split}$$

1. For part (a), $\phi_n(x)$ are the orthonormal eigenfunctions and λ_n the real eigenvalues of the corresponding homogeneous regular S-L problem:

$$-\frac{\mathrm{d}}{\mathrm{d}x}\left(p(x)\frac{\mathrm{d}y}{\mathrm{d}x}\right) + q(x)y = \lambda r(x)y$$

with initial conditions:

- $\alpha_1 y(0) + \alpha_2 y'(0) = 0$
- $\beta_1 y(1) + \beta_2 y'(1) = 0.$

for $x \in [0, 1]$. The solution to the following:

$$-\frac{\mathrm{d}}{\mathrm{d}x}\left(p(x)\frac{\mathrm{d}y}{\mathrm{d}x}\right) + q(x)y = \mu r(x)y + f(x)$$

is then:

$$y(x) = \sum_{n=1}^{\infty} b_n \phi_n(x)$$

where:

$$b_n = \frac{c_n}{\lambda_n - \mu}$$

and

$$c_n = \int_0^1 \phi_n(x) f(x) dx.$$

For part (b) solve the following:

$$\frac{\mathrm{d}^2}{\mathrm{d}x^2}y(x) + 7y(x) = 2\sin 5x + 3\sin 7x$$

with boundary conditions $y(0) = y(\pi) = 0$ for $\forall x \in [0, \pi]$.

First define change of variables $x = \pi t$.

$$\therefore y(x) \iff y(t)$$

$$\therefore \frac{\mathrm{d}}{\mathrm{d}x} y(t = \frac{x}{\pi}) = \frac{1}{\pi} \frac{\mathrm{d}y}{\mathrm{d}t}$$

$$\therefore \frac{\mathrm{d}}{\mathrm{d}x} \left(\frac{\mathrm{d}}{\mathrm{d}x} y(t = \frac{x}{\pi}) \right) = \frac{1}{\pi^2} \frac{\mathrm{d}^2 y}{\mathrm{d}t^2}$$

Then our ODE becomes:

$$\frac{1}{\pi^2} \frac{\mathrm{d}^2 y}{\mathrm{d}t^2} + 7y(t) = 2\sin 5\pi t + 3\sin 7\pi t$$

with boundary conditions y(0) = y(1) = 0 for $\forall t \in [0,1]$. This is now of S-L form, and its corresponding homogeneous S-L system is:

$$y'' + \pi^2 \lambda y = 0.$$

We know that S-L problems have real valued eigenvalues, and so we consider the sign of λ separately. Now if $\lambda = 0$ we have linear solutions:

$$y = a_1 t + a_2$$

but since $y(0) = a_2 = 0$ and $y = a_1 = 0$, only trivial solutions remain.

If $\lambda < 0$ we have solutions:

$$y = b_1 \cosh \pi \sqrt{\lambda} t + b_2 \sinh \pi \sqrt{\lambda} t$$

and using boundary conditions, $y(0) = b_1 = 0$ and $y(1) = b_2 \sinh \pi \sqrt{\lambda} = 0$. This also yields only trivial solutions since:

$$\sinh \pi \sqrt{\lambda} = \frac{1}{2} \left(e^{\pi \sqrt{\lambda}} - e^{-\pi \sqrt{\lambda}} \right) = 0$$

implies $e^{2\pi\sqrt{\lambda}} = 1$ and the only eigenvalue satisfying this is $\lambda = 0$.

Finally if $\lambda > 0$ we have solutions of form:

$$y = c_1 \sin \pi \sqrt{\lambda} t + c_2 \cos \pi \sqrt{\lambda} t$$

and our boundary conditions yields $y(0) = c_2 = 0$ and:

$$y(1) = c_1 \sin \pi \sqrt{\lambda} = 0$$

implies $\lambda_n = n^2$ where $n \in \mathbb{N}$. So our eigenfunctions are:

$$\phi_n(t) = k_n \sin \pi \sqrt{\lambda_n} t$$
$$= k_n \sin n\pi t.$$

Since $\langle \phi_n, \phi_n \rangle = 1$ and $r(x) = \pi^2$ we then have that:

$$\int_0^1 \pi^2 k_n^2 (\sin^2 n\pi t) dt = 1.$$

$$\therefore k_n = \frac{1}{\pi} \left[1 - \frac{1}{2\pi n} \sin 2\pi n \right]^{-1/2}$$

Returning to our original ODE and rearranging it into S-L form:

$$-y'' = 7\pi^2 y - \pi^2 (2\sin 5\pi t + 3\sin 7\pi t)$$

where we have $\mu = 7$, $r(t) = \pi^2$ and

$$f(t) = -\pi^2 (2\sin 5\pi t + 3\sin 7\pi t).$$

Let solutions be of the following form:

$$y(t) = \sum_{n=1}^{\infty} b_n \phi_n(t)$$

where

$$b_n = \frac{c_n}{\lambda_n - \mu}$$

and

$$c_n = \int_0^1 \phi_n(x) f(x) dx$$

$$= -\pi^2 k_n \left[\int_0^1 \sin n\pi t \left(2\sin 5\pi t + 3\sin 7\pi t \right) dt \right]$$

$$= -\pi^2 k_n \left[2 \int_0^1 \sin(n\pi t) \sin(5\pi t) dt + 3 \int_0^1 \sin(n\pi t) \sin(7\pi t) dt \right]$$

$$= -\pi^2 k_n \left[\frac{\sin(n-5)\pi}{(n-5)\pi} - \frac{\sin(n+5)\pi}{(n+5)\pi} \right]$$

$$- \frac{3}{2} \pi^2 k_n \left[\frac{\sin(n-7)\pi}{(n-7)\pi} - \frac{\sin(n+7)\pi}{(n+7)\pi} \right]$$

for $n \neq 5, 7$. We then have that:

$$b_n = \frac{c_n}{n^2 - 7}$$

and

$$k_n = \frac{1}{\pi} \left[1 - \frac{1}{2\pi n} \sin 2\pi n \right]^{-1/2}.$$

These coefficients form the solutions to our ODE:

$$y(t) = \sum_{n=1}^{\infty} b_n \phi_n(t)$$

where

$$\phi_n(t) = k_n \sin n\pi t.$$

2. For part (a) find the Laplace transform of:

$$f(t) = \sinh at$$
.

Firstly we have that:

$$e^{-st} \sinh at = \frac{1}{2}e^{-st} \left[e^{at} + e^{-at} \right]$$
$$= \frac{1}{2} \left[e^{-(s-a)t} - e^{-(s+a)t} \right].$$

Therefore:

$$\begin{split} \mathcal{L}[f(t)] &= \int_0^\infty e^{-st} \sinh(at) \mathrm{d}t \\ &= \int_0^\infty \frac{1}{2} \Big[e^{-(s-a)t} - e^{-(s+a)t} \Big] \mathrm{d}t \\ &= \lim_{T \to \infty} \Big[-\frac{1}{s-a} e^{-(s-a)t} + \frac{1}{s+a} e^{-(s+a)t} \Big]_0^T \\ &= \frac{a}{s^2 - a^2}. \end{split}$$

For part (b) show that:

$$\mathcal{L}[t^n] = \frac{n!}{s^{n+1}}$$

for all non-negative n.

We proceed via induction. Let n = 1:

$$\mathcal{L}[t] = \int_0^\infty t e^{-st} dt$$

$$= \left[-\frac{1}{s} t e^{-st} - \frac{1}{s^2} e^{-st} \right]_0^\infty$$

$$= \frac{1}{s^2}$$

since an exponential grows faster than a linear one.

We then assume that the following is true:

$$\mathcal{L}[t^k] = \frac{k!}{s^{k+1}}$$

and we want to show:

$$\mathcal{L}[t^{k+1}] = \frac{(k+1)!}{s^{k+2}}.$$

So integrating this by parts we get:

$$\mathcal{L}[t^{k+1}] = \int_0^\infty t^{k+1} e^{-st} dt$$

$$= \left[-\frac{1}{s} t^{k+1} e^{-st} \right]_0^\infty + \frac{k+1}{s} \int_0^\infty t^k e^{-st} dt$$

$$= \frac{k+1}{s} \mathcal{L}[t^k]$$

$$= \frac{(k+1)!}{s^{k+2}}.$$

Now verifying this for k = 0:

$$\mathcal{L}[t^0] = \int_0^\infty e^{-st} dt$$
$$= \left[-\frac{1}{s} e^{-st} \right]_0^\infty$$
$$= \frac{1}{s}.$$

For part (c) solve the following ODE:

$$y^{(4)}(t) = 3\sinh 2t$$

with initial conditions:

$$y(0) = y^{(1)}(0) = y^{(2)}(0) = y^{(3)}(0) = 1.$$

First begin by taking the Laplace transforms of both sides:

$$\mathcal{L}[y^{(4)}] = \mathcal{L}[3\sinh 2t] = \frac{6}{s^2 - 4}.$$

And our left hand side becomes:

$$\begin{split} \mathcal{L}[y^{(4)}] &= s^4 \mathcal{L}[y(t)] - s^3 y^{(0)}(0) - s^2 y^{(1)}(0) - s y^{(2)}(0) - y^{(3)}(0) \\ &= s^4 \mathcal{L}[y(t)] - s^3 - s^2 - s - 1. \end{split}$$

Equating and rearranging:

$$\begin{split} \mathcal{L}[y(t)] &= \frac{6}{s^4(s^2-4)} + \frac{1}{s} + \frac{1}{s^2} + \frac{1}{s^3} + \frac{1}{s^4} \\ &= \frac{1}{s} + \frac{5}{8} \frac{1}{s^2} + \frac{1}{s^3} + -\frac{1}{2} \frac{1}{s^4} + \frac{3}{16} \frac{2}{s^2-4}. \end{split}$$

Since we have the following standard transforms:

$$\frac{1}{s} = \mathcal{L}[1], \quad \frac{1}{s^2} = \mathcal{L}[t], \quad \frac{1}{s^3} = \frac{1}{2}\mathcal{L}[t^2], \quad \frac{1}{s^4} = \frac{1}{6}\mathcal{L}[t^3]$$

and

$$\frac{2}{s^2 - 4} = \mathcal{L}[\sinh 2t]$$

we have that by inspection:

$$y(t) = 1 + \frac{5}{8}t + \frac{1}{2}t^2 - \frac{1}{12}t^3 + \frac{3}{16}\sinh 2t.$$