# ${\rm Cu_3Au}$ order-disorder entropy changes Thermodynamics Hand-in Project

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#### 1 Introduction

 $Cu_3Au$  is an alloy that undergoes an order-disorder (OD) transition from a  $L1_2$  ordered face-centred cubic (fcc) with Cu occupying the face centre to a disordered state in temperature range 560K to 675K[1], as shown in figure 1. This report aims to calculate the experimental and theoretical values of the entropy change  $\Delta S$  during this OD transition and attempt to account for any discrepancies. The entropy change of a material is important as other physical behavioural quantities like the isobaric expansivity and  $C_V$  may be predicted via the Maxwell's relations.

## 1.1 Background theory

#### 1.1.1 Entropy changes in OD transitions

The total change in entropy from our OD transition is given by [3]

$$\Delta S^{OD} = \Delta S^{vib} + \Delta S^{config} + \Delta S^{mag} + \Delta S^{elec}$$

where  $\Delta S^{vib}$  is the vibrational entropy change,  $\Delta S^{mag}$  is the magnetic entropy change and  $\Delta S^{elec}$  is the electronic entropy change from disordering. Finally  $\Delta S^{config}$  corresponds to the theoretical entropy of mixing and represents the entropy change when particles are rearranged from ordered fcc to a disordered state where there are still 4 atoms per lattice but their specific positions are now unknown. Figure 1 below shows the ordered fcc structure of Cu<sub>3</sub>Au where 1/4 of a Cu atom occupies the face centre and 1/8 of an Au atom occupies the vertices, summing up to give 4 atoms per lattice.



Figure 1: Ordered fcc Cu<sub>3</sub>Au[4]

#### 1.1.2 Measuring $C_P$

Here we outline the method for finding  $\Delta S$  given data  $C_P$  and temperature T. We define

$$C_P = \frac{\mathrm{d}Q_p}{\mathrm{d}T} = \left(\frac{\mathrm{d}H}{\mathrm{d}T}\right)_p$$

where state function H is the enthalpy. Since H = U + PV:

$$dH = TdS + VdP$$

and dividing by dT:

$$\frac{\mathrm{d}H}{\mathrm{d}T} = T\frac{\mathrm{d}S}{\mathrm{d}T} + V\frac{\mathrm{d}P}{\mathrm{d}T}.$$

If  $\Delta P = 0$  then:

$$C_P(T) = T \left(\frac{\mathrm{d}S}{\mathrm{d}T}\right)_P.$$

Rearranging and integrating gives us the change in entropy:

$$\Delta S^{OD} = \int_{(1)}^{(2)} \frac{C_P(T)}{T} dT.$$

#### 1.1.3 Ideal entropy of mixing

Consider a box with separator, containing  $n_A$  moles at  $V_A$  of gas A on one side and  $n_B$  moles at  $V_B$  of gas B on the other. We set  $\Delta T = 0$ ,  $\Delta P = 0$  and proceed to mix gases A and B.

The entropy of mixing is:

$$\begin{split} \Delta S &= \Delta S_A + \Delta S_B \\ &= n_i R \ln \left( \frac{V_f}{V_i} \right) \\ &= n_A R \ln \left( \frac{V_A + V_B}{V_A} \right) + n_B R \ln \left( \frac{V_A + V_B}{V_B} \right). \end{split}$$

Since P and T are fixed, the ideal state equation PV = nRT implies that:

$$\frac{V_A + V_B}{V_A} = \frac{n_A + n_B}{n_A}$$

and so we define **inverse mole fractions** as:

$$x_A = \frac{n_A}{n_A + n_B}$$

and

$$x_B = \frac{n_B}{n_A + n_B}.$$

After substituting and dividing through by  $n_A + n_B$  we find the **molar specific entropy of mixing**:

$$\Delta s_{mix} = -R(x_A \ln x_A + x_B \ln x_B).$$

In the context of Cu<sub>3</sub>Au we have that:

$$\Delta S_{mix} = \Delta S^{config}$$

$$= -4R(x_{Cu} \ln x_{Cu} + x_{Au} \ln x_{Au})$$

$$= 18.7 \text{Jmol}^{-1} K$$

where this is per 1 mole of Cu<sub>3</sub>Au and that  $x_{Cu} = \frac{3}{4}$  and  $x_{Au} = \frac{1}{4}$ .

# 2 Data analysis

#### 2.1 Visualisations

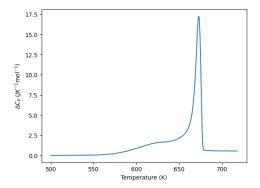
The temperature T and heat capacity  $C_P$  data are from Benisek and Dachs.[1] Using the following Python code we can plot figures 2 and 3.

```
import pandas as pd
import matplotlib.pyplot as plt

df = pd.read_csv("data.csv")

x = df["temp"]
y1 = df["cp"]
y2 = df["cp"] / x

plt.plot(x, y1)
plt.show()
```



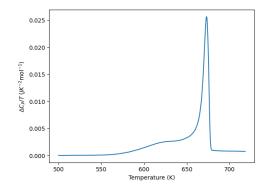


Figure 2:  $\Delta C_P (JK^{-1} \text{mol}^{-1})$  against T(K)

Figure 3:  $\frac{\Delta C_P}{T}$   $(JK^{-2}\text{mol}^{-1})$  against T (K)

It can be shown that there is a peak near T = 680K using the following code:

```
import numpy as np
i1 = np.argmax(y1)  # finds index of max value
print(x[i1])
i2 = np.argmax(y2)
print(x[i2])
```

which gives the following outputs:

```
> 672.96
> 672.96
```

The slight difference in temperatures (673K compared to 680K) are the result of either experimental precision or that our peak occurs in a range of values rather than a specific temperature.

## 2.2 Calculating $\Delta S^{OD}$

From before we have the following integral on endpoints [a, b]:

$$\Delta S^{OD} = \int_{a}^{b} \frac{C_{P}(T)}{T} dT$$

$$\approx \sum_{i} \frac{C_{P}(t_{i})}{t_{i}} \delta t_{i}$$

where given the *i*th datum of data we have that  $\delta t_i = t_{i+1} - t_i$  for  $t_i \in [a, b]$  and holds  $\forall i \in \{1, \ldots, n\}$ . Here *n* corresponds to the number of data entries. This sum may be found using the following code:

```
df_int = pd.concat([x,y2], axis = 1)
df_cut = df_int[df_int.temp.between(560, 675)]
df_cut.columns = ['temp', 'cpt']

temp = df_cut["temp"].tolist()
cpt = df_cut["cpt"].tolist()
l = len(temp)

total_area = 0
for i in list(range(1, 1)):
    base = temp[i]-temp[i-1]
    height = cpt[i]
    area = base*height
    total_area += area
print(total_area) # in units R
```

which gives the following output:

```
> 0.38718425519279304
```

This choice of endpoints [560K, 675K] is suggested by Benisek and Dachs:

$$\Delta S^{OD} = \int_{560}^{675} \frac{C_P(T)}{T} dT$$
$$\approx 0.39R$$
$$= 3.24JK^{-1} \text{mol}^{-1}$$

where  $R = 8.31 J K^{-1} \text{mol}^{-1}$  is the molar gas constant. Choosing another set of endpoints:

$$\Delta S^{OD} = \int_{550}^{700} \frac{C_P(T)}{T} dT$$
$$\approx 3.58 J K^{-1} \text{mol}^{-1}$$

and so after taking averages of the two we claim:

$$\Delta S^{OD} = 3.41 J K^{-1} \text{mol}^{-1}$$
.

### 2.3 Physical interpretations

Summarising our results we have that:

$$\Delta S^{OD} = \Delta S^{vib} + \Delta S^{config} + \Delta S^{mag} + \Delta S^{elec}$$

for

$$\Delta S^{config} = 18.7 \text{Jmol}^{-1} K \text{ and } \Delta S^{OD} = 3.41 J K^{-1} \text{mol}^{-1}.$$

This suggests that some of the other entropies (vibrational, magnetic and electronic) are negative valued, confirmed via Paras and Allanore.[3] They also note that  $Cu_3Au$  undergoes multiple stages of transitions before disordering.  $(I \to II \to D)$  Furthermore we have that  $Cu_3Au(I)$  is an orderd fcc and  $Cu_3Au(II)$  a tetragonal structure consisting of 18  $Cu_3Au(I)$  unit cells.[2] The disordered fcc state (D) for temperatures greater than T = 675K can be thought of as a solid "solution" of Cu and Cu and Cu atoms where entropy change is computed via ideal entropy of mixing.

## 3 Conclusions

This report calculates the ideal entropy of mixing and compares it to the numerically integrated entropy value from experimental data. It is then concluded that the discrepancies between these two values were due to other entropy values. We also remark that improved quenching techniques and experimental data with higher precision will result in a better data analysis.

#### References

- [1] Artur Benisek and Edgar Dachs. "The vibrational and configurational entropy of disordering in Cu3Au". In: *Journal of Alloys and Compounds* 2015 (May 2015), pp. 585–590. DOI: 10.1016/j.jallcom.2014.12.215.
- [2] H. Okamoto. "The Au-Cu (Gold-Copper) system". In: Journal of Phase Equilibria 8 (Oct. 1987), pp. 454-474. DOI: 10.1007/BF02893155.
- [3] J. Paras and Antoine Allanore. "Contribution of electronic entropy to the order-disorder transition of Cu3Au". In: *Physical Review Research* 3 (June 2021). DOI: 10.1103/PhysRevResearch. 3 023239
- [4] Mark Weller, Tina Overton, and Jonathan Rourke. *Inorganic Chemistry*. OUP Oxford, 2018.