Decomposition of three qubit unitaries
Balint Pato
2020


#### Abstract

We would like to implement arbitrary three-qubit unitary decomposition into twoqubit and one-qubit gates in Cirq (https://github.com/quantumlib/Cirq/issues/451). This write-up explores the methodology described in [SBM06] which is the best known algorithm in terms of CNOT count, giving 20 CNOT decomposition for any three qubit unitaries.


## 1 Lower Bounds

Decomposing a unitary in the general case results in an exponential amount of single qubit and two-qubit gates. Based on [SMB04] the theoretical lower bound for CNOTs in three qubit circuits are $\frac{1}{4}\left(4^{n}-3 n-1\right)=54 / 4=13.5 \simeq 14$ in the gateset $\left\{R_{x}, R_{z}, C N O T\right\}$.

The following algorithm gives 20 CNOTs.

## 2 Quantum multiplexors

The idea of quantum multiplexors is described in [SBM06]:
A quantum multiplexor is a unitary that leaves the $S$, select qubits in their original state while changes the data qubits depending on the value of $S$. For each possible classic value of $S$ the multiplexor can act with a different unitary.

While this sounds like controlled gates - they are not. This is a generalization of the control notion, where based on $S$ different unitaries can be executed, in contrast controlled gates are only "choosing" between identity or a single unitary.

Notation: An n-qubit multiplexor $U$ in circuit diagrams is denoted by " $\square$ " on each select qubit, connected by a vertical line to a gate on the remaining data qubits, with the $U$ symbol on the rectangle over the data qubits. While technically $U$ denotes the whole unitary and not just the $n_{\text {data qubits }}$ qubit unitary that the actual "box" covers, this circuit notation allows to express the multiplexor structure.

Examples:

- $S$ is the most significant qubit - this case U is block diagonal: $U=\left(\begin{array}{ll}U_{0} & \\ & U_{1}\end{array}\right)$, which we can denote with $U=U_{0} \oplus U_{1}$

- $\mathrm{CNOT}=I \oplus \sigma^{x}$


## 3 Shannon decomposition

The classical Shannon expansion of boolean functions is an important result that describes an $N$ variable boolean function as the XOR of two $N-1$ variable Boolean functions that are both restricted in one variable to 0 or 1 respectively:

$$
f\left(x_{1}, x_{2}, . ., x_{N}\right)=\left(x_{1} \wedge f\left(x_{1}=0, x_{2}, x_{3}, \ldots, x_{N}\right)\right) \oplus\left(\neg x_{1} \wedge f_{1}\left(x_{1}=1, x_{2}, x_{3}, \ldots, x_{N}\right)\right)
$$

We can start to see how the language of quantum multiplexors will be helpful to express this kind of decomposition of step-by-step smaller qubit count operators.

### 3.1 The unoptimized algorithm

We start with an unoptimized algorithm, that gives 24 CNOTs. We will optimize it later to achieve the promised 20 CNOT count.

- Step 1: Cosine Sine decomposition (Theorem 10)


Where $U D=\left(\begin{array}{cc}u_{1} & 0 \\ 0 & u_{2}\end{array}\right) \in S U(8)$ is a multiplexor between $u_{1}$ and $u_{2}$ and similarly $V D H=\left(\begin{array}{cc}v_{1 h} & 0 \\ 0 & v_{2 h}\end{array}\right) \in S U(8)$ and $C S=\left(\begin{array}{cc}C & -S \\ S & C\end{array}\right) \in S U(8)$, where $C$ and $S$ are

4 x 4 diagonal matrices that satisfy $C^{2}+S^{2}=I$, finally $\theta_{i}, i \in\{0,1,2,3\}$ are the four possible rotations of the $R_{y}$ rotation that is being multiplexed by the four possible states of the two least significant bits.
Since SciPy version 1.5.0 the cossin method implements the Cosine Sine decomposition.

```
from scipy.linalg import cossin
(u1, u2), theta, (v1h, v2h) = cossin(U, 4, 4, separate=True)
```

- Step 2: demultiplex the multiplexed $R_{y}$ : it's easy to verify that the following circuit satisfies the required multiplexing logic, that is $\mid k>\rightarrow R_{y}\left(\theta_{k}\right)$ :

, where

$$
\begin{aligned}
\alpha & =\theta_{0}+\theta_{1}+\theta_{2}+\theta_{3} \\
\beta & =\theta_{0}+\theta_{1}-\theta_{2}-\theta_{3} \\
\gamma & =\theta_{0}-\theta_{1}-\theta_{2}+\theta_{3} \\
\delta & =\theta_{0}-\theta_{1}+\theta_{2}-\theta_{3}
\end{aligned}
$$

- Step 3: demultiplex the two-qubit multiplexor $U D$ :a block diagonal matrix can be diagonalized in a way that creates two two-qubit gates and a multiplexed $R_{z}$ in the middle multiplexed on qubits 0 and 1, acting on qubit 2, i.e $u_{1} \oplus u_{2}=(I \otimes V)(D \oplus$ $\left.D^{\dagger}\right)(I \otimes W):$

$$
\left(\begin{array}{ll}
u_{1} & \\
& u_{2}
\end{array}\right)=\left(\begin{array}{ll}
V & \\
& V
\end{array}\right)\left(\begin{array}{cc}
D & \\
& D^{\dagger}
\end{array}\right)\left(\begin{array}{ll}
W & \\
& W
\end{array}\right)
$$

The calculation of $V, D$ and $W$ comes from:

$$
\begin{aligned}
u_{1} & =V D W \\
u_{2} & =V D^{\dagger} W \\
u_{2}^{\dagger} & =W^{\dagger} D V^{\dagger} \\
u_{1} u_{2}^{\dagger} & =V D W W^{\dagger} D V^{\dagger}=V D^{2} V^{\dagger}
\end{aligned}
$$

Where D is diagonal. We implemented cirq. unitary_eig to ensure that the resulting eigenvectors are orthogonal - i.e the resulting $V$ is unitary:

```
u1u2 = u1 @ u2.conj().T
eigvals, V = cirq.unitary_eig(u1u2)
d = np.diag(np.sqrt(eigvals))
```

$W$ can be easily expressed as $W=D V^{\dagger} u_{2}$.

- Step 4: Implementing the Diagonal D is very similar to CS , using a 4 -way $R_{z}$ gate. At this point we have $C S$ implemented as a four-way $R_{y}$ gate, $W$ and $V$ two-qubit unitaries, $D$ implemented as a four-way $R_{z}$ gate.

- Step 5: similarly decompose $V D H$ - giving 4 CNOTs for the 4 -way multiplexed $R_{z}$
- Step 6: decompose the four two qubit operators (using the KAK decomposition) that gives 3 CNOTs for each operator

This gives 24 CNOTs $=4 \times$ two-qubit operators x 3 CNOTs (KAK) +2 x 4 -way multiplexed $R_{z}$ gates x 4 CNOTs ( $V D H$ and $U D$ diagonals) $+1 \times 4$-way multiplexed $R_{y}$ gate $\times 4$ CNOTs in the middle ( $C S$ )

### 3.2 Optimizations

### 3.2.1 $\mathrm{CNOT} \rightarrow \mathrm{CZ}$ in $R_{y}$

Appendix A. 1 in [SBM06] explains how replacing the CNOTs with CZs works equivalently in the multiplexed $R_{y}$ implementation. The terminal CZ, as the CZ gate is diagonal, can be merged with the neighboring generic two-qubit multiplexer (UD).

Now, we are down to 3 CZs and 20CNOTs.

### 3.2.2 Eager diagonals

Based on Theorem 14 in Appendix A. 2 by [SBM06] there exists a diagonal gate $\Delta$ that we can extract from any two-qubit unitary that will leave a two-qubit gate that can be decomposed with only two CNOT gates:


This diagonal commutes through the controls of the multiplexed $R_{y}$ and can be merged with the generic two-qubit multiplexers on the left of the circuit. As the KAK decomposition implemented in Cirq recognizes the two-CNOT circuits (is this true in general?), as long as we extract this diagonal, we can win 3 more CNOTs.

In order to understand this decomposition we need to look at [MBS03] for more details.

## 4 Extracting a diagonal from two-qubit circuits

### 4.1 Invariants of two-qubit unitaries

[MBS03] describes equivalence classes of two qubit special unitaries depending on whether they require zero, one, two or three CNOTs to implement them.
$U(4)$ is the group of two-qubit unitaries, $S U(4)$ is the group of determinant one two-qubit unitaries, the special unitary group.

Def: $\gamma: U(4) \rightarrow U(4), \gamma(u)=u\left(\sigma^{y}\right)^{\otimes 2} u^{T}\left(\sigma^{y}\right)^{\otimes 2}$
Def: $\chi(u)(x)=p(x)=\operatorname{det}(x I-u)$ the characteristic polynomial.
Def: $u$ and $v$ two are equivalent up to local unitaries if there exist one-qubit operators that, when pre- and post-composing with $u$, up to a global phase we get $v$. Denoted by $u \equiv v$.
[MBS03], Proposition II.1: $\forall u, v \in S U(4) u \equiv v \Longleftrightarrow \chi(\gamma(u))=\chi( \pm \gamma(v))$
[MBS03], Proposition III.1: An operator $u \in S U(4)$ can be simulated with 0 CNOT gates, if $\chi(\gamma(u))=(x+1)^{4}$ or $(x-1)^{4}$
[MBS03], Proposition III.2: An operator $u \in S U(4)$ can be simulated with 1 CNOT gate, if $\chi(\gamma(u))=(x+i)^{2}(x-i)^{2}$
[MBS03], Proposition III.3: An operator $u \in S U(4)$ can be simulated with 2 CNOT gates, if $\chi(\gamma(u))$ has real coefficients, which occurs if $\operatorname{tr}[\gamma(u)] \in \mathcal{R}$

### 4.2 Extracting the diagonal

The diagonal extraction is presented by in another paper by the same authors in [SMB04]. $C_{2}^{1}$ represents a CNOT gate with control on the 1st qubit, target on the 2nd qubit:
[SMB04], Proposition V.2: For any $u \in S U(4)$ one can find $\theta, \phi, \psi$ so that $\chi\left[\gamma\left(u C_{2}^{1}(I \otimes\right.\right.$ $\left.\left.R_{z}(\psi)\right) C_{2}^{1}\right]=\chi\left[\gamma\left(C_{2}^{1}\left(R_{x}(\theta) \otimes R_{z}(\phi)\right) C_{2}^{1}\right]\right.$.

Which is exactly what we want: compose our two-qubit unitary $(u)$ from the left (in the circuit diagram) with a diagonal $\left(C_{2}^{1}\left(I \otimes R_{z}(\psi)\right) C_{2}^{1}\right.$ is diagonal) to get a unitary that can be implemented with only two CNOTs.

The proof is constructive and contains the algorithm to find $\psi$, however it does have a typo/bug in the formulae:

$$
\begin{aligned}
\Delta & :=C_{2}^{1}\left(I \otimes R_{z}(\psi)\right) C_{2}^{1} \\
\operatorname{tr}[\gamma(u \Delta)] & =\left(t_{1}+t_{4}\right) e^{-i \psi}+\left(t_{2}+t_{3}\right) e^{i \psi}
\end{aligned}
$$

, where $t_{1}, t_{2}, t_{3}, t_{4}$ are the diagonal entries of $\gamma\left(u^{T}\right)^{T}$.
Now, the paper claims that "We may ensure that this number is real by setting $\tan (\psi)=$ $\frac{I m\left(t_{1}+t_{2}+t_{3}+t_{4}\right)}{R e\left(t_{1}+t_{2}-t_{3}-t_{4}\right)}$.

Which is incorrect, it is relatively easy to deduce that $\tan (\psi)=\frac{I m\left(t_{1}+t_{2}+t_{3}+t_{4}\right)}{\operatorname{Re}\left(t_{1}+t_{4}-t_{2}-t_{3}\right)}$ is the right formula.

Also, there is one missing case mentioned in the paper:

- when $\operatorname{Re}\left(t_{1}+t_{4}-t_{2}-t_{3}\right)=0$ and will mean that:

$$
\begin{aligned}
\left(t_{1}+t_{4}\right) e^{-i \psi} & +\left(t_{2}+t_{3}\right) e^{i \psi} \in \mathbb{R} \Longrightarrow \\
e^{i \psi} & =-e^{-i \psi} \\
\psi & =3 \pi / 2 \text { or } \pi / 2
\end{aligned}
$$

The python code:

```
def special(u):
    return u / (np.linalg.det(u) ** (1 / 4))
```

```
def g(u):
    yy = np.kron(cirq.Y._unitary_(), cirq.Y._unitary_())
    return u @ yy @ u.T @ yy
def extract_right_diag(U):
    u = special(U)
    t = _gamma(_to_special(U).T).T.diagonal()
    k = np.real(t[0] + t[3] - t[1] - t[2])
    if k == 0:
        # in the end we have to pick a psi that makes sure that
        # exp(-i*psi) (t[0]+t[3]) + exp(i*psi) (t[1]+t[2]) is real
        # both pi/2 or 3pi/2 can work
        psi = np.pi/2
    else:
        psi = np.arctan(np.imag(np.sum(t)) / k)
    a, b = cirq.LineQubit.range(2)
    c_d = cirq.Circuit([cirq.CNOT(a, b), cirq.rz(psi)(b), cirq.CNOT(a, b)])
    return c_d._unitary_()
V = circuit._unitary_()
dV = extract_right_diag(V)
V = V @ dV
print(cirq.Circuit(
        cirq.optimizers.two_qubit_matrix_to_operations(
        a,b,V,allow_partial_czs=False
        )))
np.trace(g(special(V)))
```

$(-2.618033988749896-2.7755575615628914 \mathrm{e}-16 \mathrm{j})$


## References

[MBS03] Igor L Markov, Stephen S Bullock, and Vivek V Shende. "Recognizing SmallCircuit Structure in Two-Qubit Operators and Timing Hamiltonians to Compute Controlled-Not Gates". In: Quant-Ph/0308045 (2003), pp. 3-6. Doi: doi : 10 . 1103/PhysRevA.70.012310. arXiv: 0308045 [quant-ph]. URL: http://arxiv. org/abs / quant - ph / 0308045\% 7B\% 5C\%\%7D5Cnhttp : / / www . arxiv . org/pdf / quant-ph/0308045.pdf.
[SMB04] Vivek V Shende, Igor L Markov, and Stephen S Bullock. Minimal Universal TwoQubit CNOT-based Circuits. Tech. rep. 2004.
[SBM06] Vivek V Shende, Stephen S Bullock, and Igor L Markov. Synthesis of Quantum Logic Circuits. Tech. rep. 2006.

