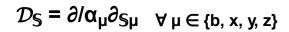
Innes D. Anderson-Morrison 2017

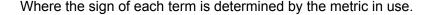
\mathcal{D}_S : the 4Set differential

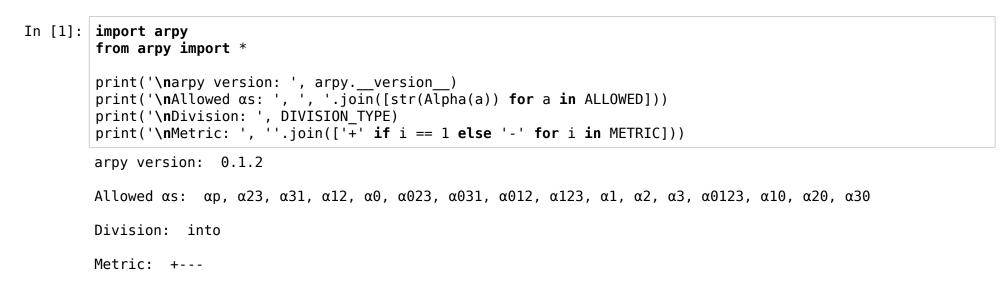
The 16 components of the algebra can be grouped into what I am calling 4-sets that consist of a 3-vector $\{x,y,z\}$ triplet (\mathbb{B} , \mathbb{T} , \mathbb{A} and \mathbb{E}) paired with an additional element: α_p , α_0 , α_{123} and α_{0123} respectively. From here on I will be referring to these four elements as \mathbf{p} , \mathbf{t} , \mathbf{s} and \mathbf{q} and the elements of each 3-vector as subscripts on their respective vector name: α_{23} as $\mathbf{B}_{\mathbf{x}}$ for example, with the paired additional element with the subscript \mathbf{b} (for paired blade).

Under this notation $\mathbb{B} = \{B_b, B_x, B_y, B_z\} = \{\alpha_p, \alpha_{23}, \alpha_{31}, \alpha_{12}\}$ and the general 4-set is denoted $S = \{S_b, S_x, S_y, S_z\}$ in boldface with S denoting the usual 3-vector elements.

A generalised 4-set differential can be defined for Cartesian coordinate as follows:







Defining and using the operators within arpy

As covered in the docs, the way to define a new operator fitting the pattern described above is to use the **differential_operator** function. The list of α indices provided acts as the set of values that μ is allowed to take in the above definition.

Here we define the four 4-set differentials: v_B , \mathcal{D}_T , \mathcal{D}_A and \mathcal{D}_E .

```
In [2]: DB = differential_operator(['p', '23', '31', '12'])
DT = differential_operator(['0', '023', '031', '012'])
DA = differential_operator(['123', '1', '2', '3'])
DE = differential_operator(['0123', '10', '20', '30'])
```

The action of such a differential is to always produce 16 components that can be grouped into 5 vector calculus elements.

```
In [3]: DB(XiT)

Out[3]: {

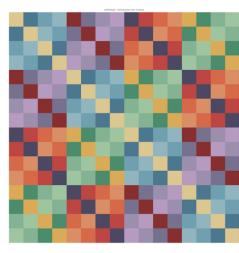
\alpha\theta = (\partial_{p}\xi_{0}, \partial_{23}\xi_{023}, \partial_{31}\xi_{031}, \partial_{12}\xi_{012})
\alpha\theta = (\partial_{p}\xi_{0}, \partial_{p}\xi_{023}, \partial_{12}\xi_{031}, -\partial_{31}\xi_{012})
\alpha\theta = (\partial_{p}\xi_{0}, \partial_{p}\xi_{023}, \partial_{p}\xi_{031}, \partial_{23}\xi_{012})
\alpha\theta = (\partial_{p}\xi_{0}, \partial_{12}\xi_{023}, \partial_{p}\xi_{031}, \partial_{23}\xi_{012})
\alpha\theta = (\partial_{p}\xi_{0}, \partial_{12}\xi_{023}, \partial_{p}\xi_{031}, \partial_{23}\xi_{012})
\alpha\theta = (\partial_{p}\xi_{0}, \partial_{12}\xi_{023}, \partial_{p}\xi_{031}, \partial_{p}\xi_{012})
\theta = (\partial_{p}\xi_{0}, \partial_{12}\xi_{023}, \partial_{12}\xi_{031}, \partial_{12}\xi_{012})
\theta = (\partial_{p}\xi_{0}, \partial_{12}\xi_{023}, \partial_{12}\xi_{031}, \partial_{12}\xi_{012})
```

Note that when grouping into vector calculus <u>del notation (https://en.wikipedia.org/wiki/Del)</u> I am using the symbol ∇^S to denote the standard 3-vector ∇ operator with $\{x, y, z\}$ components being drawn from the 4-set S.

```
For \mathcal{D}_S S' we obtain: \partial_b S_b' \partial_b S' \nabla^S \cdot S_v' \nabla^S \cdot S'_b \nabla^S \cdot X \cdot S_v'
```

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The 5 ∂_b and ∇^S terms



This diagram (also shown above) is a visualisation of the Cayley table for the algebra that does not take sign into account. As such, it applies to all product based operations between αs and by extension, between 4-sets.

The colour of each cell denotes the resultant α value of the composition of the row α with the column α with α s in B, T, A, E order: α p, α 23, α 31, α 12, α 0, α 023, α 031, α 012, α 123, α 1, α 2, α 3, α 0123, α 10, α 20, α 30.

When we compose one full 4-set with another we *always* obtain a result that is shown in the diagram as a single colour 4x4 block with tone denoting **{b, x, y, z}**.

As an example below, we take the full product of the MultiVectors **A** and **T** (red and purple respectively in the diagram) and obtain α s from **E** (green in the diagram).

In fact, the four 4-sets form a group under *full-product-like* composition (identifying divison as a full-product of one term with the inverse of another) with the Magnetic 4-set **B** as the identity:

```
X | B T A E
B | B T A E
T | T B E A
A | A E B T
E | E A T B
```

The group formed is isomorphic to the <u>Klein four-group</u> (https://en.wikipedia.org/wiki/Klein_four-group) which has the <u>group presentation</u> (https://en.wikipedia.org/wiki/Presentation_of_a_group): $< a, b \mid a^2 = b^2 = (ab)^2 = I >$

(It should also be noted that the internal structor of each 4-set is also isomorphic to the Klein four-group.)

The general sign distribution of the algebra of course changes with respect to the operation being used, the metric and the choice of ordering for the elements of grade 2 and above. In practice, only the components of the electric field may be varied between **i0** and **0i**.

The following table shows the sign of each resulting component under full-product.

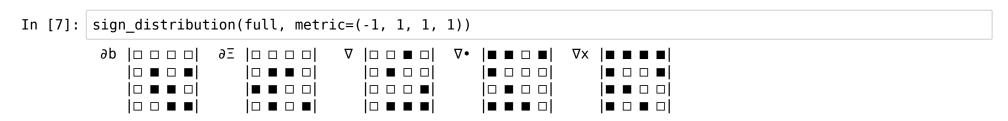
full product i0 +---

In [6]: sign_cayley(full) Ε αρ α23 α31 α12 α0 α023 α031 α2 α3 α0123 | - - - - | - - - - - - - - - - - -α10 α20 α30

Alternatively, the **sign_distribution** function can be used to look at what happens to each of the 5 groups of components under different setups. This also allows for specification of different metrics and allowed as.

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full product i0 -+++



full product 0i +---

```
In [8]: oi = ['p', '23', '31', '12', '0', '023', '031', '012', '123', '1', '2', '3', '0123', '01', '02', '03']
      sign_distribution(full, allowed=oi)
                                    ∇• |■ ■ □ □|
                ∂Ξ |□ □ □ □|
                                              ∇x | □ □ □ □ |
      ∂b |□ □ □ □|
                           ▽ |□ □ ■ ■|
```

Generalised Force equations: SD_SS

As discussed in <u>Dr Williamson's paper (http://eprints.gla.ac.uk/110966/)</u>, the generalised <u>Lorentz Force (https://en.wikipedia.org/wiki/Lorentz_force)</u> may be obtained by computing the full-product of F (the even subgroup of the $\mathbb G$ which is equal to B + E) with its space-time derivative $\mathcal D_u$ F:

Force = $F\mathcal{D}_{u}F$

I say *may* as we are currently unsure as to whether this is the correct equation (which is the entire point of this notebook!) so we are looking into alternatives such as $\mathcal{D}_{\mu}\mathsf{FF}$, $\mathcal{D}_{\mu}(\mathsf{FF}^{\dagger})$, $\mathcal{D}_{\mu}(\mathsf{GG}^{\dagger})$ etc. I feel that it may be worthwhile looking at a generalisation of $\mathsf{F}\mathcal{D}_{\mu}\mathsf{F}$ namely $\mathsf{S}\mathcal{D}_{\mathsf{S}}\mathsf{S}$ where each of the three S terms may be any 4-set.

The simplest case to consider is that of $\mathbb{B}\mathcal{D}_{\mathbb{R}}\mathbb{B}$

As $\mathbb B$ is the group identity it is self contained under composition with all elements of $\mathcal D_{\mathbb B}\mathbb B$ remaining within $\mathbb B$ and all elements of the product $\mathbb B\mathcal D_{\mathbb B}\mathbb B$ also remaining within $\mathbb B$. (This can be seen from the leading diagonal of the Cayley table shown before: the composition of any 4-set with itself results in $\mathbb B$ -terms.)

Below I show the result of computing this product along with a partial simplification that is possible under the current version of arpy:

```
 \begin{array}{l} \text{In [9]: } & \text{DBB = DB(XiB)} \\ \text{DB(XiB, as\_del=True)} \\ \\ \text{Out[9]: } \\ & \begin{array}{l} \alpha p & (\nabla^B \bullet B, \ \partial_P \xi_P) \\ \alpha j k & (\partial_P B, \ -\nabla^B \Xi_P, \ -\nabla^B x B) \\ \\ \end{array} \\ \\ \text{In [10]: } & \begin{array}{l} \text{BDBB = ar('XiB $^$ DBB')} \\ \text{print('{$}$ terms:'.format(len([term \ \textbf{for} \ term \ \textbf{in} \ BDBB])))} \\ \\ \text{64 terms:} \end{array}
```

```
α<sub>p</sub>:
          \xi_{p} \left( \partial_{p} \xi_{p} + \partial_{23} \xi_{23} + \partial_{31} \xi_{31} + \partial_{12} \xi_{12} \right)
          + \xi_{23}( \partial_{23}\xi_{p} - \partial_{p}\xi_{23} - \partial_{12}\xi_{31} + \partial_{31}\xi_{12} )
          + \xi_{31}( \partial_{31}\xi_{p} + \partial_{12}\xi_{23} - \partial_{p}\xi_{31} - \partial_{23}\xi_{12} )
          + \xi_{12}( \partial_{12}\xi_p - \partial_{31}\xi_{23} + \partial_{23}\xi_{31} - \partial_p\xi_{12} )
α23:
          \xi_p(\partial_2 3\xi_p + \partial_p \xi_2 3 + \partial_{12}\xi_{31} + \partial_{31}\xi_{12})
          + \xi_{23}( \partial_{p}\xi_{p} + \partial_{23}\xi_{23} + \partial_{31}\xi_{31} + \partial_{12}\xi_{12} )
          + \xi_{31}( \partial_{12}\xi_p + \partial_{31}\xi_{23} + \partial_{23}\xi_{31} + \partial_p\xi_{12} )
          + \xi_{12} (\partial_{31} \xi_p + \partial_{12} \xi_{23} - \partial_p \xi_{31} - \partial_{23} \xi_{12})
αз1:
          \xi_p \left( \partial_{31} \xi_p + \partial_{12} \xi_{23} + \partial_p \xi_{31} + \partial_{23} \xi_{12} \right)
          + \xi_{23}( \partial_{12}\xi_{p} - \partial_{31}\xi_{23} + \partial_{23}\xi_{31} - \partial_{p}\xi_{12} )
          + \xi_{31}( \partial_{p}\xi_{p} + \partial_{23}\xi_{23} + \partial_{31}\xi_{31} + \partial_{12}\xi_{12} )
          + \xi_{12} ( \partial_{23} \xi_p + \partial_p \xi_{23} + \partial_{12} \xi_{31} + \partial_{31} \xi_{12} )
α12:
          \xi_p(\partial_{12}\xi_p + \partial_{31}\xi_{23} + \partial_{23}\xi_{31} + \partial_p\xi_{12})
          + \xi_{23} ( \partial_{31} \xi_p + \partial_{12} \xi_{23} + \partial_p \xi_{31} + \partial_{23} \xi_{12} )
          + \xi_{31} ( \partial_{23} \xi_p - \partial_p \xi_{23} - \partial_{12} \xi_{31} + \partial_{31} \xi_{12} )
          + \xi_{12}( \partial_{p}\xi_{p} + \partial_{23}\xi_{23} + \partial_{31}\xi_{31} + \partial_{12}\xi_{12} )
```

Or, in *partially* simplified notation we have this:

Note: further automated simplification is planned but this requires simplification of product terms accross common factors in addition to settling on the notation to use for some of the new grouped elements that have been identified. (See the next section for more details)

Simplifying further by hand I obtain the following:

Note: in the following expressions $\mu\nu\lambda$ run over $\mathbb{B}_{x,y,z}$ (α_{23} , α_{31} , α_{12}) in cyclic left to right order. There are some terms using this notation that I think may be grouped onto α_{ik} in a similar way to Curl.

```
\begin{split} \mathbb{B}\mathcal{D}_{\mathbb{B}}\mathbb{B} &= \\ \{ \\ &\alpha_{p} \colon P[\ \partial_{p}P + \nabla^{B}B\ ] + B[\ \partial_{jk}P - B \bullet \partial_{p}B + B_{\mu}(\partial_{\lambda}B_{\mu} - \partial_{\nu}B_{\lambda})\ ] \\ &\alpha_{jk} \colon P[\ \partial_{p}B + \nabla^{B}P\ ] + P.\partial_{p}P + B \bullet \nabla^{B}B \\ &\alpha_{\mu} \colon P(\partial_{jk\lambda}B_{\nu} + \partial_{jk\nu}B_{\lambda}) + B_{\lambda}\partial_{\nu}P + B_{\nu}\partial_{\nu}B_{\mu} + (B_{\nu}\partial_{\mu}B_{\nu} - B_{\lambda}\partial_{\mu}B_{\lambda}) + (B_{\nu}\partial_{p}B_{\lambda} - B_{\lambda}\partial_{p}B_{\nu}) \\ \} \end{split}
```

Each of the α_{μ} terms could be included in the α_{jk} group if we can come up with a name for what that operation *means*. I'm not 100% confident that each of those groups will ammount to a Physical description of something (in the way that Div, Grad and Curl do) but it may be worth further investigation. At the very least, it allows for a more compact representation of the result.

It is easy to see that we have a similar distribution of grouped elements (four 6-component Curl like terms for example) but that the original grouped term structure is not preserved. I think that this collection of terms can be generalised to the $\mathbb{S}\mathcal{D}_{\mathbb{S}}\mathbb{S}$ case modulo sign which will depend on the metric in use.