

QMII

Exercise 3

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November 13, 2022

Question 1

Given the transformation $U(\lambda) = e^{-i\lambda \frac{G}{\hbar}}$ and the generator $G = \frac{1}{2}(xp + px)$.

Part 1

Show how an infinitesimal transformation $U(\epsilon)$ transforms the position and momentum operators x , p .

Solution:

$$U(\epsilon) \approx 1 - i\epsilon \frac{G}{\hbar} = 1 - i\frac{\epsilon}{2\hbar}(xp + px).$$

$$\begin{aligned} UxU^\dagger &\approx \left[1 - i\frac{\epsilon}{2\hbar}(xp + px)\right] x \left[1 - i\frac{\epsilon}{2\hbar}(xp + px)\right] \\ &= x - i\frac{\epsilon}{2\hbar}(xp + px)x + i\frac{\epsilon}{2\hbar}x(xp + px) \\ &= x + i\frac{\epsilon}{2\hbar}[x, xp + px] \\ &= x + i\frac{\epsilon}{2\hbar}[x, 2xp - i\hbar] \\ &= x + i\frac{\epsilon}{\hbar}[x, xp] \\ &= x + i\frac{\epsilon}{\hbar}x[x, p] \\ &= x + i\frac{\epsilon}{\hbar}xi\hbar \\ &= \boxed{(1 - \epsilon)x}. \end{aligned}$$

$$\begin{aligned}
UpU^\dagger &\approx \left[1 - i\frac{\epsilon}{2\hbar}(xp + px)\right] p \left[1 - i\frac{\epsilon}{2\hbar}(xp + px)\right] \\
&= x - i\frac{\epsilon}{2\hbar}(xp + px)p + i\frac{\epsilon}{2\hbar}p(xp + px) \\
&= x + i\frac{\epsilon}{2\hbar}[p, xp + px] \\
&= x + i\frac{\epsilon}{2\hbar}[p, 2xp - i\hbar] \\
&= x + i\frac{\epsilon}{\hbar}[p, xp] \\
&= x + i\frac{\epsilon}{\hbar}p[p, x] \\
&= x - i\frac{\epsilon}{\hbar}pi\hbar \\
&= \boxed{(1 + \epsilon)p}.
\end{aligned}$$

Part 2

Given the results of the previous part, show how a finite transformation transforms x, p .

Solution:

Dividing λ into N equal sections and letting $N \rightarrow \infty$ we obtain:

$$\begin{aligned}
x' &= \left(1 - \frac{\lambda}{N}\right)^N x \stackrel{(N \rightarrow \infty)}{=} e^{-\lambda} x \\
p' &= \left(1 + \frac{\lambda}{N}\right)^N p \stackrel{(N \rightarrow \infty)}{=} e^{\lambda} p.
\end{aligned}$$

Part 3

Show what the transformation U does to $|x\rangle$.

Solution:

Firstly we'll note that U and x do not commute so we'll have to find these two expressions:

$$\begin{aligned}
Ux|x\rangle &\quad ; \quad xU|x\rangle. \\
Ux|x_0\rangle &= Ux_0|x_0\rangle = x_0U|x_0\rangle \\
UxU^\dagger U|x_0\rangle &= e^{-\lambda}xU|x_0\rangle \\
\Rightarrow Ux|x_0\rangle &= x_0U|x_0\rangle = e^{-\lambda}xU|x_0\rangle \\
\Rightarrow xU|x_0\rangle &= e^{\lambda}x_0U|x_0\rangle
\end{aligned}$$

Which means that $U|x_0\rangle$ is an eigenket of x with eigenvalue $e^{\lambda}x_0$:

$$U|x_0\rangle \propto |e^{\lambda}x_0\rangle.$$

Now we only need to see if U pulls out a constant:

$$\begin{aligned}\delta(x - x') &= \langle x' | x \rangle = \langle x' | U^\dagger U | x \rangle = |A|^2 \langle e^\lambda x' | e^\lambda x \rangle = |A|^2 \delta(e^\lambda(x - x')) = \frac{|A|^2}{e^\lambda} \delta(x - x') \\ &\Rightarrow \frac{|A|^2}{e^\lambda} = 1 \\ &\Rightarrow A = e^{\frac{\lambda}{2}} \\ &\Rightarrow \boxed{U | x \rangle = e^{\frac{\lambda}{2}} | e^\lambda x \rangle}\end{aligned}$$

Question 2

Show that:

$$R_{z'}(\gamma) R_{y'}(\beta) R_z(\alpha) = R_z(\alpha) R_y(\beta) R_z(\gamma).$$

Solution:

Firstly we'll note that y' after the z rotation is:

$$R_{y'}(\beta) = R_z(\alpha) R_y(\beta) R_z^{-1}(\alpha).$$

And that z' after the y rotation is:

$$R_{z'}(\gamma) = R_{y'}(\beta) R_z(\gamma) R_{y'}^{-1}(\beta).$$

So:

$$R_{z'}(\gamma) R_{y'}(\beta) R_z(\alpha) = R_{y'}(\beta) R_z(\gamma) \cancel{R_{y'}^{-1}(\beta) R_{y'}(\beta)} R_z(\alpha).$$

If we replace now $R_{y'}(\beta)$ and remember that rotations around the same axis commute:

$$R_{z'}(\gamma) R_{y'}(\beta) R_z(\alpha) = R_z(\alpha) R_y(\beta) \cancel{R_z^{-1}(\alpha) R_z(\alpha)} R_z(\gamma) = R_z(\alpha) R_y(\beta) R_z(\gamma) \checkmark.$$

Question 3

Given the following rotation in $j = \frac{1}{2}$:

$$\mathcal{D}(\hat{n}, \theta) = \frac{i}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}.$$

Find \hat{n} and θ .

Solution:

$$\begin{aligned}\mathcal{D} &= \frac{i}{\sqrt{2}} (\sigma_x + \sigma_z) = i\boldsymbol{\sigma} \cdot \left(\frac{\hat{x} + \hat{z}}{\sqrt{2}} \right). \\ \Rightarrow \begin{cases} \cos \frac{\theta}{2} = 0 \\ \sin \frac{\theta}{2} = -1 \end{cases} &\Rightarrow \boxed{\hat{n} = \frac{\hat{x} + \hat{z}}{\sqrt{2}}, \theta = 3\pi}.\end{aligned}$$

Question 4

Use the identity for $j = 1$:

$$e^{-i\theta J_y/\hbar} = 1 - i \sin \theta \frac{J_y}{\hbar} + (\cos \theta - 1) \left(\frac{J_y}{\hbar} \right)^2.$$

To find the elements of $d_{m'm}^{(j=1)}$.

Solution:

Firstly we'll note that in $j = 1$:

$$\frac{L_y}{\hbar} = \frac{1}{i\sqrt{2}} \begin{pmatrix} 0 & 1 & 0 \\ -1 & 0 & 1 \\ 0 & -1 & 0 \end{pmatrix} \Rightarrow \left(\frac{L_y}{\hbar} \right)^2 = \begin{pmatrix} \frac{1}{2} & 0 & -\frac{1}{2} \\ 0 & 1 & 0 \\ -\frac{1}{2} & 0 & \frac{1}{2} \end{pmatrix}.$$

Hence:

$$\begin{aligned} e^{-i\theta J_y/\hbar} &= \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} - i \sin \theta \frac{1}{i\sqrt{2}} \begin{pmatrix} 0 & 1 & 0 \\ -1 & 0 & 1 \\ 0 & -1 & 0 \end{pmatrix} + (\cos \theta - 1) \begin{pmatrix} \frac{1}{2} & 0 & -\frac{1}{2} \\ 0 & 1 & 0 \\ -\frac{1}{2} & 0 & \frac{1}{2} \end{pmatrix}. \\ &\Rightarrow d_{m'm}^{(j=1)} = \begin{pmatrix} \frac{1+\cos \theta}{2} & -\frac{\sin \theta}{\sqrt{2}} & \frac{1-\cos \theta}{2} \\ \frac{\sin \theta}{\sqrt{2}} & \cos \theta & -\frac{\sin \theta}{\sqrt{2}} \\ \frac{1-\cos \theta}{2} & \frac{\sin \theta}{\sqrt{2}} & \frac{1+\cos \theta}{2} \end{pmatrix}. \end{aligned}$$

Question 5

Given a state:

$$|\psi\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix}.$$

- Find $|\psi'\rangle$ after a rotation of $\frac{\pi}{2}$ around the \hat{z} axis.
- Find $|\psi'\rangle$ after a rotation of $\frac{\pi}{2}$ around the \hat{y} axis.

Solution:

a)

$$\begin{aligned} R_z \left(\frac{\pi}{2} \right) &= e^{-i\frac{\pi}{2} J_z/\hbar} = \begin{pmatrix} e^{-i\frac{\pi}{2}} & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -e^{-i\frac{\pi}{2}} \end{pmatrix}. \\ \Rightarrow |\psi'\rangle &= \frac{1}{\sqrt{2}} \begin{pmatrix} e^{-i\frac{\pi}{2}} & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -e^{-i\frac{\pi}{2}} \end{pmatrix} \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} = \boxed{\frac{1}{\sqrt{2}} \begin{pmatrix} -i \\ 1 \\ 0 \end{pmatrix}}. \end{aligned}$$

b)

$$\begin{aligned} R_y \left(\frac{\pi}{2} \right) &= e^{-i\frac{\pi}{2} J_y/\hbar} = d_{m'm}^{j=1} \left(\theta = \frac{\pi}{2} \right) = \begin{pmatrix} \frac{1}{2} & -\frac{1}{\sqrt{2}} & \frac{1}{2} \\ \frac{1}{\sqrt{2}} & 0 & -\frac{1}{\sqrt{2}} \\ \frac{1}{2} & \frac{1}{\sqrt{2}} & \frac{1}{2} \end{pmatrix}. \\ \Rightarrow |\psi'\rangle &= \frac{1}{\sqrt{2}} \begin{pmatrix} \frac{1}{2} & -\frac{1}{\sqrt{2}} & \frac{1}{2} \\ \frac{1}{\sqrt{2}} & 0 & -\frac{1}{\sqrt{2}} \\ \frac{1}{2} & \frac{1}{\sqrt{2}} & \frac{1}{2} \end{pmatrix} \begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} = \boxed{\frac{1}{2} \begin{pmatrix} \frac{1}{\sqrt{2}} - 1 \\ 1 \\ \frac{1}{\sqrt{2}} + 1 \end{pmatrix}}. \end{aligned}$$

Question 6

The dynamics of a particle with $j = 1$ is given by the hamiltonian:

$$\mathcal{H} = \epsilon \begin{pmatrix} 2 & \frac{1-i}{2} & 0 \\ \frac{1+i}{2} & 2 & \frac{1-i}{2} \\ 0 & \frac{1+i}{2} & 2 \end{pmatrix}.$$

a) Write \mathcal{H} as a sum of the elements of \mathbf{J} and J^2 i.e.:

$$\mathcal{H} = aJ^2 + b\mathbf{J} \cdot \hat{\mathbf{a}}.$$

b) Is \mathcal{H} symmetrical under an arbitrary rotation?

c) With which angle θ and around which axis $\hat{\mathbf{n}}$, $\hat{\mathbf{a}}$ must be turned such that it will point in $\hat{\mathbf{z}}$?

d) What are the euler angles α, β, γ that correspond to the rotation from the previous section?

e) Find the rotation matrix $\mathcal{D}(\hat{\mathbf{n}}, \theta)$.

f) Show that the rotation matrix diagonalize \mathcal{H} .

Solution:

a)

$$\mathcal{H} = \frac{\epsilon}{\hbar^2} J^2 + \frac{\epsilon}{\hbar} \mathbf{J} \cdot \left(\frac{\hat{\mathbf{x}} + \hat{\mathbf{y}}}{\sqrt{2}} \right).$$

$$\Rightarrow E_{lm} = \epsilon(j(j+1) + m) = (2+m)\epsilon.$$

b) No, the system is only symmetrical to rotations around the $\hat{\mathbf{a}} = \frac{\hat{\mathbf{x}} + \hat{\mathbf{y}}}{\sqrt{2}}$ axis.

c) Around $\hat{\mathbf{n}} = \frac{\hat{\mathbf{x}} - \hat{\mathbf{y}}}{\sqrt{2}}$ with $\theta = \frac{\pi}{2}$.

d)

$$\gamma = \frac{\pi}{4} \quad \beta = \frac{\pi}{2} \quad \alpha = 0.$$

e) With $d_{m'm}^{(j=1)}$ which we already calculated:

$$\begin{aligned} \mathcal{D} &= e^{-i\frac{\pi}{2}J_y/\hbar} e^{-i\frac{\pi}{4}J_z/\hbar} \\ &= \begin{pmatrix} \frac{1}{2} & -\frac{1}{\sqrt{2}} & \frac{1}{2} \\ \frac{1}{\sqrt{2}} & 0 & -\frac{1}{\sqrt{2}} \\ \frac{1}{2} & \frac{1}{\sqrt{2}} & \frac{1}{2} \end{pmatrix} \begin{pmatrix} e^{-i\frac{\pi}{4}} & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & e^{i\frac{\pi}{4}} \end{pmatrix} \\ &= \begin{pmatrix} e^{i\frac{\pi}{4}} & \sqrt{2} & e \end{pmatrix}. \end{aligned}$$