

4) (a) choose a discrete set of functions  $\{\phi_i\}$

Such as that:

$$\hat{\psi}(r) = \sum_i \phi_i(r) \hat{a}_i$$

$$\int \phi_i^*(r) \phi_j(r) dr = \delta_{ij} \quad \text{and} \quad \int \phi_i(r) \phi_j^*(r) dr$$

From

$$\delta(r-r') = [\psi(r), \psi^\dagger(r')] = \sum_{ij} \phi_i(r) \phi_j(r') [a_i, a_j^\dagger]$$

Multiply by  $\phi_k(r')$  and  $\int \dots dr'$

$$\phi_k(r) = \sum_i \phi_i(r) [a_i, a_k^\dagger]$$

Multiply by  $\phi_l^*(r)$  and integrate:

$$\delta_{lk} = [a_l, a_k^\dagger]$$

And let us find coherent states in each  $a_i$ :

A coherent state  $^{\text{of}}$   $a_i$  is:

$$|\alpha_i\rangle = e^{-\frac{|\alpha_i|^2}{2} + \alpha_i \hat{a}_i^\dagger} |0\rangle.$$

← Maybe should  
show how to find  
it, but we saw  
it in class.

$$|\{\alpha_i\}\rangle = \otimes_i |\alpha_i\rangle.$$

$$\psi(r) |\{\alpha_i\}\rangle = \sum_j \phi_j(r) \alpha_j |\{\alpha_i\}\rangle = \sum_i \phi_i(r) \alpha_i |\{\alpha_i\}\rangle$$

So the ~~eigenfun~~ eigenstates of  $\psi(r)$  are  $|\{\alpha_i\}\rangle$

and the ~~eigenfunctions~~ are eigenvalues are  $\psi_r = \sum_i \phi_i(r) \alpha_i$ .

$$\langle \{\alpha_i\} | \mathcal{N} | \{\alpha_i\} \rangle = \int \langle \{\alpha_i\} | \psi^\dagger(r) \psi(r) | \{\alpha_i\} \rangle dr =$$

$$\begin{aligned} &= \sum_{ij} \int \phi_i^* \phi_j(r) \langle \{\alpha_i\} | a_i^\dagger a_j | \{\alpha_i\} \rangle dr = \sum_{ij} \int \langle \{\alpha_i\} | a_i^\dagger a_j | \{\alpha_i\} \rangle = \\ &= \sum_i |\alpha_i|^2. \end{aligned}$$

$$H = H_0 + V_{int}$$

$$H_0 = \int \psi^\dagger(r) \left( \underbrace{-\frac{\hbar^2}{2m} \nabla^2 + U(r)}_{h_0} \right) \psi(r) dr$$

$$V_{int} = \frac{1}{2} \int \psi^\dagger(r) \psi^\dagger(r') V \psi(r) \psi(r') dr dr'$$

If we assume that  $\{\phi_i\}$  are eigenfunctions of  $h_0$  then

$$H_0 = \sum_i \epsilon_i a_i^\dagger a_i \quad \text{and general interaction is } V_{int} = \sum_{ijkl} V_{ijkl} a_i^\dagger a_j^\dagger a_l a_k$$

$$\langle \{\alpha_i\} | H_0 | \{\alpha_i\} \rangle = \sum_i \epsilon_i |\alpha_i|^2$$

$$\langle \{\alpha_i\} | V_{int} | \{\alpha_i\} \rangle = \sum V_{ijkl} \alpha_i^* \alpha_j^* \alpha_l \alpha_k$$

$$V_{ijkl} = \langle ij | V | kl \rangle$$

If  $V$  is diagonalized ~~in~~ with in our basis then  $\langle \{\alpha_i\} | V_{int} | \{\alpha_i\} \rangle = \sum V_{ij} |\alpha_i|^2 |\alpha_j|^2$ .

Overall:

$$\langle \{\alpha_i\} | H | \{\alpha_i\} \rangle = \sum_i \epsilon_i |\alpha_i|^2 + \sum_{ijkl} V_{ijkl} \alpha_i^* \alpha_j^* \alpha_l \alpha_k$$

(b)  $a, a^\dagger$  are ladder operators,  $b, b^\dagger$  are creation & annihilation operators.

Both of them have the same ~~non~~ commutation relations.

On the other side  $a, a^\dagger$  are operators in a Hilbert space of ~~the~~ ~~an~~ functions while  $b, b^\dagger$  are operators in Fock-space.

Another difference is that  $b, b^\dagger$  generate ~~excitation~~ <sup>particles</sup> with certain energy while  $a, a^\dagger$  raise (or lower)

the energy of one particle. Therefore,  
when working with S.p. it is easier to  
use  $a, a^\dagger$  while when working with ~~some~~  
many body interactions  $b, b^\dagger$  are a lot more  
useful.