

Quantum Mechanics Study Guide

Lectures 6 – 9: Operators, Projections, and Uncertainty

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Winter 2026

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1 Lecture 6: Generalized Quantum Systems

1.1 Core Concepts

- **Postulate 2:** Physical observables are represented by **operators** acting on kets.
- **Postulate 3:** Measurement outcomes are the **eigenvalues** $\{a_n\}$ of that operator.
- **Hermitian Operators:** Observables must be Hermitian ($\hat{A} = \hat{A}^\dagger$) because:
 1. They have **real eigenvalues** (measurement results must be real).
 2. Eigenvectors with distinct eigenvalues are **orthogonal**.

1.2 Matrix Representation (How to Build Operators)

To represent an operator \hat{A} as a matrix in a specific basis $\{|a_n\rangle\}$, calculate the matrix elements A_{nm} :

$$A_{nm} = \langle a_n | \hat{A} | a_m \rangle \quad (1)$$

If the basis consists of eigenvectors of \hat{A} , the matrix is **diagonal** with eigenvalues on the diagonal.

Memorize: Spin Matrices (in Z-basis)

For a spin-1/2 particle, the operators in the S_z basis are:

$$\begin{aligned} S_z &= \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \\ S_x &= \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \\ S_y &= \frac{\hbar}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \end{aligned}$$

2 Lecture 7: Projection Operators

2.1 The Projection Operator

A projector isolates the component of a vector along a specific direction (eigenstate).

$$\hat{P}_n \equiv |a_n\rangle \langle a_n| \quad (2)$$

2.2 Important Identities

- **Completeness Relation:** Summing all projectors returns the identity matrix. Use this to insert "1" into equations to simplify them.

$$\sum_n \hat{P}_n = \sum_n |a_n\rangle \langle a_n| = \mathbb{I} \quad (3)$$

- **Spectral Decomposition:** Any operator can be written as a sum of its eigenvalues weighted by its projectors.

$$\hat{A} = \sum_n a_n \hat{P}_n = \sum_n a_n |a_n\rangle \langle a_n| \quad (4)$$

2.3 Postulate 5: Measurement Collapse

If you measure observable \hat{A} and get result a_n , the state collapses.

Collapse Formula

$$|\psi_f\rangle = \frac{\hat{P}_n |\psi_i\rangle}{\sqrt{\langle \psi_i | \hat{P}_n | \psi_i \rangle}} \quad (5)$$

2.4 Solving Stern-Gerlach (S-G) Problems

Problem Type: "What is the probability of exiting the final port?" or "What is the final state?"

Step-by-Step Method:

1. **Identify the Input:** Start with the initial state vector $|\psi_{in}\rangle$.
2. **Apply Projectors:** For every S-G analyzer the particle passes through, apply the corresponding projection operator.
 - Example: Passing through an S_z+ filter means applying $\hat{P}_{z+} = |+\rangle\langle +|$.
 - $|\psi_{after}\rangle = \hat{P}_{filter} |\psi_{before}\rangle$.
3. **Normalize (if asking for state):** If asking for the state after measurement, divide by the norm (length) of the vector.
4. **Calculate Probability (if asking for prob):**

$$P = |\langle \text{final state} | \text{current state} \rangle|^2 \quad \text{OR} \quad P = \langle \psi_{current} | \hat{P}_{final} | \psi_{current} \rangle \quad (6)$$

3 Lecture 8: Expectation Values

3.1 Definitions

The **Expectation Value** is the average result of a large number of measurements on identical systems.

$$\langle \hat{A} \rangle = \langle \psi | \hat{A} | \psi \rangle \quad (7)$$

The **RMS Deviation** (Uncertainty) is:

$$\Delta A = \sqrt{\langle A^2 \rangle - \langle A \rangle^2} \quad (8)$$

3.2 Calculation Strategy

To calculate ΔA :

1. Calculate $\langle A \rangle$:

$$\langle A \rangle = \langle \psi | \hat{A} | \psi \rangle$$

(Matrix multiplication: Row Vector \times Matrix \times Column Vector).

2. Calculate \hat{A}^2 : Square the matrix \hat{A} .

3. Calculate $\langle A^2 \rangle$:

$$\langle A^2 \rangle = \langle \psi | \hat{A}^2 | \psi \rangle$$

4. Plug into the RMS formula.

4 Lecture 9: The Uncertainty Principle

4.1 Commutators

The commutator measures if two operators commute (order doesn't matter) or not.

$$[\hat{A}, \hat{B}] \equiv \hat{A}\hat{B} - \hat{B}\hat{A} \quad (9)$$

- If $[\hat{A}, \hat{B}] = 0$: The observables are **compatible**. They share a common basis and can be measured simultaneously with arbitrary precision.
- If $[\hat{A}, \hat{B}] \neq 0$: The observables are **incompatible**.

4.2 Spin Commutation Relations

Remember the cyclic permutation ($x \rightarrow y \rightarrow z \rightarrow x$):

$$\begin{aligned} [S_x, S_y] &= i\hbar S_z \\ [S_y, S_z] &= i\hbar S_x \\ [S_z, S_x] &= i\hbar S_y \end{aligned}$$

4.3 Heisenberg Uncertainty Principle

The general form for any two operators:

General Uncertainty Relation

$$\Delta A \Delta B \geq \frac{1}{2} | \langle [\hat{A}, \hat{B}] \rangle | \quad (10)$$

4.4 Useful Identity: Magnitude of Spin

While individual components are uncertain, the *magnitude* of the spin vector is constant:

$$\hat{S}^2 = \hat{S}_x^2 + \hat{S}_y^2 + \hat{S}_z^2 = \frac{3\hbar^2}{4} \mathbb{I} \quad (11)$$