### Note

#### Feng-Yang Hsieh

# 1 Signal

We consider a simple extension of the standard model (SM) [1], which includes a vectorlike dark fermion ( $\overline{\chi}$ ,  $\chi$ ) and a complex singlet scalar S. A signature of CP violation could come from the Higgs-to-Higgs decays,  $h_3 \to h_2 h_1$ , where  $h_3/h_2/h_1$  are the heaviest scalar, second heaviest scalar, and the SM-like 125 GeV Higgs, respectively.

The signal process is the triple production of 125 GeV Higgs bosons via the gluon fusion:

$$gg \rightarrow h_3 \rightarrow h_2 h_1 \rightarrow h_1 h_1 h_1$$

The Higgs boson  $h_1$  would further decay to the  $b\bar{b}$  pair. We consider the banchmark point 1 (BP1), where  $m_{h_3}=450$  GeV,  $m_{h_2}=280$  GeV,  $m_{h_1}=125$  GeV. This process is generated at  $\sqrt{s}=13$  TeV. Following are the MadGraph scripts for generating signal samples:

```
import model cxSM_VLF_EFT
generate g g > h h h
output MG5/gghhh_bsm
launch MG5/gghhh_bsm
```

```
shower=Pythia8
detector=Delphes
analysis=OFF
madspin=ON
done
```

```
set param_card mh1 125
set param_card mh2 280
set param_card mh3 420
set param_card theta12 0.73
```

```
set param card theta13 1.67079632679
set param_card theta23 -0.73
set param card vs 200
set param card delta2 0
set param card Rdelta3 0
set param card Idelta3 -3.5
set param card b2 0
set param_card Rc1 0
set param_card Ic1 0
set param_card Rc2 0
set param card Ic2 0
set param card Rd3 0
set param card Id3 0
set param card msq -5033.406281907266
set param card lam 0.13850082540690806
set param card Rdelta1 -47.561525227572744
set param card Idelta1 853.05384671134
set param card Rb1 -70476.6380004269
set param_card Ib1 -30486.140015405872
set param card Rd1 -2.562109886826132
set param card Id1 2.257859679994403
set param_card d2 6.340799300844676
set param_card gh1ggr -0.00005478952893059635
set param_card gh1gagar -0.00003270447254456052
set param_card gh1Zgar -0.00005871986046374793
set param card gh2ggr -1.4279972541632635e-7
set param_card gh2gagar -8.237715486808595e-8
set param card gh2Zgar -1.3984990232267825e-7
set param card gh3ggr -6.031835872118092e-6
set param card gh3gagar -1.1377279177203616e-6
set param card gh3Zgar -2.2999597941282603e-6
set param_card decay 102 auto
set param_card decay 103 auto
set run card nevents 100000
```

```
set run_card ebeam1 6500.0
set run_card ebeam2 6500.0
set run_card ptb 24
set run_card etab 2.6
set spinmode none
decay h > b b~
```

## 2 SPANet pairing

done

We employ the novel neural network structure SPA-NET [2, 3, 4] to identify the correct pairings among the jets in the final states.

#### 2.1 Training dataset preparation

Preselection:  $\geq 6$  jets with transverse momentum  $p_T \geq 25$  GeV in range  $|\eta| < 2.5$ .

The input features for the SPA-NET are a list of jets, each represented by its 4-component vector  $(p_T, \eta, \phi, m)$  as well as a boolean b-tag. We only keep each event's 15 highest  $p_T$  jets. We define the correct jet assignments for each event by matching the jets to the simulated truth quarks within an angular distance of  $\Delta R < 0.4$ . Such an event will be dropped if a simulated truth quark is matched to more than one jet. Furthermore, some simulated truth quarks may not be matched to any jet, so the event will not be used in training either.

After the selection and matching, we could obtain the following results from 1M events:

• Total sample size: 522,899

• 1h sample size: 184,769

• 2h sample size: 161,476

• 3h sample size: 94,464

Here, the 1h sample is where we could define the correct jet assignments for 1 Higgs boson.

### 2.2 Training results

#### • Training sample:

- Total sample size: 470,609

- 1h sample size: 166,490

- 2h sample size: 145,309

- 3h sample size: 84,913

- 5% used on validation

#### • Testing sample:

- Total sample size: 52,290

- 1h sample size: 18,279

- 2h sample size: 16,167

- 3h sample size: 9,551

Some useful definitions for evaluating jet assignment performance:

#### • Event Efficiency

$$\epsilon^{\text{event}} \equiv \frac{\text{number of events with and all Higgs are correctly identified}}{\text{number of events}}$$
(1)

Higgs Efficiency

$$\epsilon^{\rm h} \equiv \frac{\text{number of correctly identified Higgs}}{\text{number of identifiable Higgs}}$$
(2)

The training results are shown in Table 1.

Table 1: Spa-Net pairing efficiencies on 3h events.

$N_{ m Jet}$	Event Fraction	Event Efficiency	Higgs Efficiency
= 6	0.077	0.532	0.650
=7	0.057	0.345	0.536
$\geq 8$	0.052	0.237	0.452
Total	0.186	0.375	0.548

# 3 $\chi^2$ pairing

 $\chi^2$  method considers all possible combinations of final jets and selects the configuration that minimizes the mass difference between Higgs candidates and SM Higgs, i.e., minimizes this:

$$\chi^2 = [m(j_1 j_2) - m_h]^2 + [m(j_3 j_4) - m_h]^2 + [m(j_5 j_6) - m_h]^2$$
(3)

where  $m(j_i j_j)$  is the invariant mass of jet i, j and  $m_h = 125$  GeV.

Table 2 is the performance of the  $\chi^2$  method.

Event Fraction Event Efficiency Higgs Efficiency 0.0770.4030.450=70.0570.1580.281 $\geq 8$ 0.0520.0000.077Total 0.1860.2150.294

Table 2:  $\chi^2$  pairing efficiencies on 3h events.

### 4 Estimate cross-section of background process

Besides the 6 b background, we need to consider the backgrounds that come from the mis-tagging of light jets or charm-jets to b-jets. We assume that the probability of a charm-jet being misidentified as b-jet is  $\mathcal{P}_{c\to b} = 0.1$  and that of light jets is  $\mathcal{P}_{j\to b} = 0.01$ . The b-tagging efficiency is assumed to be  $\mathcal{P}_{b\to b} = 0.7$ .

Table 3 shows the cross-section computed from MadGraph and the cross-section times the mis-tagging probabilities  $\mathcal{P}_{c\to b}$  and  $\mathcal{P}_{j\to b}$ . Table 4 shows the same results with kinetic cuts. We require the transverse momentum  $p_{\rm T}$  of each jet greater than 24 GeV and in the range  $|\eta| < 2.6$  at the MadGraph level. The 6b process contributes much more than the processes containing charm jets and light jets.

## 5 Compute pairing efficiency

To understand the pairing performance with different pairing methods, we compute how many events where 1h/2h/3h bosons are reconstructed correctly.

The pairing performance of SPA-NET are shown in Table 5. Table 6 is the performance of the  $\chi^2$  method. For both cases, we found the number of events where only two Higgs are

Table 3: The cross-sections of 6b and mis-tagging background processes. The cross-sections are computed from the MadGraph at  $\sqrt{s} = 13$  TeV.

process	$\sigma$ (pb)	$\sigma \times \mathcal{P}(\text{tagging efficiency}) \text{ (pb)}$
$(b\overline{b})(b\overline{b})(b\overline{b})$	$2.53 \times 10^3$	$2.97 \times 10^{2}$
$(b\overline{b})(b\overline{b})(c\overline{c})$	$2.72 \times 10^{2}$	$6.54 \times 10^{-1}$
$(b\overline{b})(c\overline{c})(c\overline{c})$	$3.73 \times 10^{1}$	$1.83 \times 10^{-3}$
$(b\overline{b})(b\overline{b})(jj)$	$7.44 \times 10^4$	1.79

Table 4: The cross-sections of 6b and mis-tagging background processes. The cross-sections are computed from the MadGraph at  $\sqrt{s} = 13$  TeV. We require the transverse momentum  $p_{\rm T}$  of each jets greater than 24 GeV in range  $|\eta| < 2.6$ .

process	$\sigma$ (fb)	$\sigma \times \mathcal{P}(\text{tagging efficiency}) \text{ (fb)}$
$(b\overline{b})(b\overline{b})(b\overline{b})$	$9.63 \times 10^{2}$	113.35
$(b\overline{b})(b\overline{b})(c\overline{c})$	$1.67 \times 10^{3}$	4.02
$(b\overline{b})(c\overline{c})(c\overline{c})$	$1.06 \times 10^{3}$	$5.19 \times 10^{-2}$
$(b\overline{b})(b\overline{b})(jj)$	$4.16 \times 10^{5}$	9.98
$(b\overline{b})(jj)(jj)$	$1.50 \times 10^{7}$	$7.73 \times 10^{-2}$

paired correctly is tiny, which means if we can pair two Higgs bosons correctly, then we have a high chance to pair the final Higgs correctly.

Note that Higgs Efficiencies of SPA-NET are inconsistent with Table 1. This issue needs more checking.

Table 5: Spa-Net pairing efficiencies on different categories.

	Correc	etly reco			
$N_{ m Jet}$	3h	2h	1h	0h	Higgs Efficiency
= 6	0.532	0.000	0.119	0.348	0.572
=7	0.345	0.021	0.166	0.469	0.414
$\geq 8$	0.237	0.022	0.186	0.554	0.314
Total	0.375	0.014	0.156	0.455	0.436

Table 6:  $\chi^2$  pairing efficiencies on different categories.

	Correc	etly reco			
$\overline{N_{ m Jet}}$	3h	2h	1h	0h	Higgs Efficiency
= 6	0.403	0.000	0.143	0.455	0.450
=7	0.158	0.070	0.228	0.544	0.281
$\geq 8$	0.000	0.000	0.231	0.769	0.077
Total	0.215	0.022	0.194	0.570	0.294

# 6 SPANet pairing and classification

We train a SPA-NET to identify the correct pairings and perform the signal/background classification at the same time.

#### 6.1 Training dataset

The selection and matching process for the jet pairing is the same as Section 2.1. We prepare the signal and background samples of the same size for classification.

For the jet assignment part,

#### • Training sample:

- Total sample size: 1,800,000

- 1h sample size: 318,053

- 2h sample size: 277,876

- 3h sample size: 162,444

- 5% used on validation

#### • Testing sample:

- Total sample size: 200,000

- 1h sample size: 35,372

- 2h sample size: 30,853

- 3h sample size: 18,004

For event classification,

#### • Training sample:

- Total sample size: 1,800,000

- Signal sample size: 900,000

- Background sample size: 900,000

- 5% used on validation

#### • Testing sample:

- Total sample size: 200,000

- Signal sample size: 100,000

- Background sample size: 100,000

This training takes around 10 hours on our server.

#### 6.2 Training results

The training results are presented in Table 7. This result is better than Table 5 since we use larger training datasets.

Table 7: SPA-NET training results on the tri-Higgs samples. SPA-NET is trained on jet pairing and event classification tasks at the same time.

	Correc	etly reco			
$N_{ m Jet}$	3h	2h	1h	0h	Higgs Efficiency
= 6	0.656	0.000	0.082	0.262	0.684
=7	0.436	0.017	0.168	0.379	0.504
$\geq 8$	0.341	0.018	0.173	0.468	0.411
Total	0.478	0.012	0.142	0.368	0.533

Table 8 presents the classification training results. We use the accuracy (ACC) and the area under the Receiver Operating Characteristic (ROC) curve (AUC) as two metrics.

Table 8: The SPA-NET classification training results with tri-Higgs sample.

#### 6.3 3h training dataset

We only consider 3h events in pairing tasks in this subsection. We prepare the signal and background samples of the same size for classification.

For the jet assignment part,

- Training sample:
  - Total sample size: 1,800,000
  - 1h sample size: 0
  - 2h sample size: 0
  - 3h sample size: 900,000
  - 5% used on validation
- Testing sample:
  - Total sample size: 200,000
  - 1h sample size: 0
  - 2h sample size: 0
  - 3h sample size: 100,000

For event classification,

- Training sample:
  - Total sample size: 1,800,000
  - Signal sample size: 900,000
  - Background sample size: 900,000
  - 5% used on validation
- Testing sample:
  - Total sample size: 200,000
  - Signal sample size: 100,000
  - Background sample size: 100,000

The training results are presented in Table 9. This result is similar to Table 7.

However, some issues should be resolved when we try to use all 3h events for combining training. The loss values are not reasonable.

Table 9: Spa-Net training results on the tri-Higgs samples, where we only consider 3h events. Spa-Net is trained on jet pairing and event classification tasks at the same time.

	Correc	etly reco			
$N_{ m Jet}$	3h	2h	1h	0h	Higgs Efficiency
= 6	0.680	0.000	0.084	0.236	0.708
=7	0.477	0.014	0.150	0.359	0.536
$\geq 8$	0.311	0.027	0.184	0.477	0.391
Total	0.491	0.014	0.139	0.356	0.547

# 7 Review the pairing methods

In Refs. [5, 6]: We select the 6 b-tagged jets with the highest transverse momentum. The requirements for the transverse momentum and pseudo-rapidity are applied. We subsequently make use of the observable:

$$\chi^{2,(6)} = \sum_{qr \in J} (m_{qr} - m_h)^2 \tag{4}$$

where  $J = \{j_1j_2, j_3j_4, j_5j_6\}$  is the set of all possible 15 pairings of 6 b-tagged jets. Out of all the possible combinations, we pick the one with the smallest value  $\chi_{\min}^{2,(6)}$ . The pairing of b-jets defining  $\chi_{\min}^{2,(6)}$  is our best candidate for the reconstruction of the three Higgs bosons, h. No pairing efficiency is provided.

In Ref. [7]: We select the 6 b-tagged jets with the highest transverse momentum and form pairs in different combinations, with the aim of first reconstructing individual SM-like Higgs bosons,  $h_1$ , and subsequently the two scalars  $h_2$  and  $h_3$ . To this end, we introduce two observables:

$$\chi^{2,(4)} = \sum_{qr \in I} (m_{qr} - m_h)^2 \tag{5}$$

$$\chi^{2,(6)} = \sum_{qr \in J} (m_{qr} - m_h)^2 \tag{6}$$

where we have defined the sets  $I = \{i_1i_2, i_3i_4\}$  and  $J = \{j_1j_2, j_3j_4, j_5j_6\}$ , constructed from different pairings of 4 and 6 b-tagged jets, respectively, and where  $m_{qr}$  denotes the invariant mass of the respective pairing, qr. Note that the set I that defines  $\chi_{\min}^{2,(4)}$  should be a subset of the arrangement J.

We select the combinations of b-tagged jets entering in I and J based on the minimization of the sum

$$\chi^{2,(4)} + \chi^{2,(6)} \tag{7}$$

We then "identify" candidates for the scalars  $h_2$  and  $h_3$  with the pairing configurations  $I_{\min}$  and  $J_{\min}$  which minimise  $\chi^{2,(4)}$  and  $\chi^{2,(6)}$  respectively. Note that this procedure does not guarantee that  $I_{\min}$  indeed reconstructs to  $h_2$ ; in fact, we found this to be the case in about 40% on average for all benchmark samples, being slightly higher than a "blind guess" that would lead to a probability of 1/3.

## 8 6b requirement

They could utilize stronger kinetic and b-tagging requirements in experiments. However, we do not use the b-tag information for the previous  $\chi^2$  method. Thus, we change the preselection condition to make a fair comparison.

For testing samples, we only consider the event that contains at least 6 b-tagged jets. The Higgs Bosons are reconstructed by pairing the six leading b-jets for the  $\chi^2$  method. Thus, the best one is chosen from 15 possible pairing combinations.

Table 10 and Table 11 are the pairing performance on 6 b-tagged samples. The  $\chi^2$  method performs better than the previous results (Table 6) but still performs worse than the SPA-NET pairing.

Table 10:  $\chi^2$  pairing efficiencies on different categories. The  $\chi^2$  method only considers the six leading b-jets.

	Correc	etly reco			
$N_{ m Jet}$	3h	2h	1h	0h	Higgs Efficiency
= 6	0.500	0.000	0.124	0.376	0.541
=7	0.322	0.072	0.185	0.486	0.389
$\geq 8$	0.238	0.026	0.156	0.579	0.308
Total	0.329	0.013	0.159	0.499	0.391

# 9 Another pairing method

The pairing algorithm is defined to minimize

$$|m_{h_1} - 120| + |m_{h_2} - 115| + |m_{h_3} - 110| \tag{8}$$

where  $m_{h_i}$  is the mass of the *i*-th Higgs boson candidate (sorted by  $p_{\rm T}$ ) in units of GeV. The numbers in this definition are chosen based on the peaks of the  $m_{h_i}$  distributions in simulated signal events.

Table 11: SPA-NET training results on the tri-Higgs samples, where we only consider 3h events. SPA-NET is trained on  $\geq 0b$  datasets and tested on  $\geq 6b$  datasets.

	Correc	etly reco			
$\overline{N_{ m Jet}}$	3h	2h	1h	0h	Higgs Efficiency
			0.074		0.704
=7	0.522	0.013	0.127	0.337	0.574
$\geq 8$	0.354	0.019	0.180	0.447	0.427
Total	0.492	0.013	0.136	0.359	0.546

Figure 1 is the Higgs boson invariant mass  $m_h$  distributions. The peaks of the  $m_h$  distributions in our simulation are 121, 119, 116 GeV. We modify the numbers in Equation 8 based on these values. Table 12 is the pairing performance on 6 b-tagged samples. The

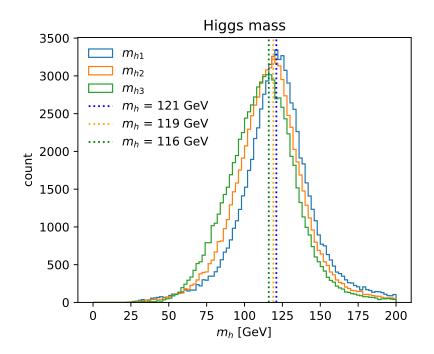


Figure 1: The Higgs boson mass distributions. Here, we only consider the event containing at least 6 b-tagged jets.

absolute value method performs worse than the  $\chi^2$  method (Table 10).

Table 12:  $\chi^2$  pairing efficiencies on different categories. The  $\chi^2$  method only considers the six leading b-jets.

	Correc	etly reco			
$\overline{N_{ m Jet}}$	3h	2h	1h	0h	Higgs Efficiency
= 6	0.476	0.000	0.135	0.388	0.522
=7	0.293	0.011	0.185	0.511	0.362
$\geq 8$	0.215	0.026	0.168	0.589	0.289
Total	0.303	0.015	0.167	0.515	0.367

# 10 New simulation setting

We apply the new setting to the event generation for a fair comparison. The following are the notes on the latest simulation setting.

The 6b analysis uses R = 0.4 anti- $k_{\rm T}$  jets. All jets are required to have  $p_{\rm T} > 20$  GeV and  $|\eta| < 2.5$ , to be within the tracker acceptance for b-tagging.

All events must pass the Preselection. They must have at least six jets, defined by the earlier selection criteria. Of these six jets, at least four jets must have  $p_T > 40$  GeV, and at least four jets must be b-tagged.

Pairing is performed on all events passing the Preselection. In the case of 4b or 5b events, the six jets considered for pairing into Higgs candidates are the b-tagged jets, and the extra jets are selected as the highest  $p_{\rm T}$  of the remaining light-flavor jets. In the case of 6b events, the 6 jets are the 6 leading  $p_{\rm T}$  b-tagged jets.

### 11 Pairing for new simulated samples

We modified the **Delphes** card to apply the anti- $k_{\rm T}$  clustering algorithm with R=0.4 and change the b-tagging efficiency to the ATLAS DL1r 77% working point. At this working point, the light-jet (charm-jet) rejection is about 130 (4.9).

Figure 2 is the Higgs boson invariant mass  $m_h$  distributions for new simulated samples. All events passing the Preselection and matching are used to generate this plot. The peaks of the  $m_h$  distributions in our simulation are 119, 115, 111 GeV. Based on these values, we modify the numbers in Equation 8.

Table 13, 14 and 15 are the pairing performance on new simulated samples. The SPA-NET performs the best.

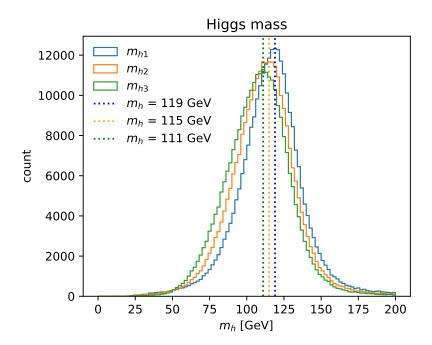


Figure 2: The Higgs boson mass distributions. Here, we consider all events passing the Preselection and the correct pairing can be obtained.

Table 13:  $\chi^2$  pairing efficiencies on different categories. We minimize the quantity defined in Equation 3. The  $\chi^2$  method considers the possible combinations of 6 jets.

		Correc	tly reco			
$N_{ m Jet}$	Fraction	3h	2h	1h	0h	Higgs Efficiency
= 6	0.241	0.519	0.000	0.112	0.368	0.557
=7	0.335	0.272	0.010	0.186	0.532	0.341
$\geq 8$	0.424	0.147	0.008	0.196	0.649	0.218
Total	1.000	0.279	0.007	0.172	0.542	0.341

Table 14: Absolute value pairing efficiencies on different categories. We minimize the quantity defined in Equation 8. The absolute value method considers the possible combination of 6 jets.

		Correc	etly reco			
$N_{ m Jet}$	Fraction	3h	2h	1h	0h	Higgs Efficiency
= 6	0.241	0.457	0.000	0.131	0.412	0.500
=7	0.335	0.239	0.009	0.184	0.568	0.306
$\geq 8$	0.424	0.136	0.007	0.197	0.661	0.206
Total	1.000	0.248	0.006	0.176	0.570	0.311

Table 15: SPA-NET training results on the tri-Higgs samples, where we only consider 3h events. The SPA-NET method considers all jets in the final state.

		Correc	etly reco			
$N_{ m Jet}$	Fraction	3h	2h	1h	0h	Higgs Efficiency
= 6	0.241	0.619	0.000	0.097	0.283	0.652
=7	0.335	0.481	0.008	0.150	0.361	0.536
$\geq 8$	0.424	0.350	0.011	0.171	0.469	0.414
Total	1.000	0.459	0.007	0.146	0.388	0.512

# 12 Modify the matching strategy

For the previous exercise, we define the correct jet assignments for each event by matching the jets to the simulated truth quarks within an angular distance of  $\Delta R < 0.4$ . Such an event will be dropped if a simulated truth quark is matched to more than one jet. Furthermore, some simulated truth quarks may not be matched to any jet, so the event will not be used in training either.

We modify our strategy for more than one jet case. If more than one jet can be matched to a simulated truth quark in the  $\Delta R = 0.4$  cone, we choose the nearest one by the  $\Delta R$  distance. This method is the same as the di-Higgs analysis [8].

Table 16 is the cutflow number at different selection cuts.

Table 16: The number of passing events, efficiencies, and passing rates for signal processes at different selection cuts.

	Count	Efficiency	Pass rate
Total	100000	1.00	1.00
$\geq 6 \text{ jets}$	61454	0.61	0.61
$\geq 4$ jets with $p_{\rm T} > 40$ GeV	50341	0.82	0.50
$\geq 4$ b-jets	32337	0.64	0.32
Matching 3h	9944	0.31	0.10
Matching 2h	11341	0.35	0.11
Matching 1h	8941	0.28	0.09

Table 17 and 18 are the pairing performance with the new matching strategy. The results are similar to the previous one (Table 13 and 14).

Table 17:  $\chi^2$  pairing efficiencies on different categories. We minimize the quantity defined in Equation 3. The  $\chi^2$  method considers the possible combinations of 6 jets.

		Correc	etly reco			
$N_{ m Jet}$	Fraction	3h	2h	1h	0h	Higgs Efficiency
= 6	0.242	0.520	0.000	0.112	0.368	0.557
=7	0.335	0.277	0.007	0.182	0.534	0.343
$\geq 8$	0.422	0.161	0.008	0.190	0.641	0.229
Total	1.000	0.287	0.006	0.168	0.539	0.347

Table 18: Absolute value pairing efficiencies on different categories. We minimize the quantity defined in Equation 8. The absolute value method considers the possible combination of 6 jets.

		Correc	Correctly reconstructed Higgs			
$N_{ m Jet}$	Fraction	3h	2h	1h	0h	Higgs Efficiency
= 6	0.242	0.450	0.000	0.132	0.418	0.494
=7	0.335	0.249	0.007	0.177	0.566	0.313
$\geq 8$	0.422	0.145	0.009	0.194	0.653	0.215
Total	1.000	0.254	0.006	0.173	0.567	0.316

# 13 SM tri-Higgs samples

We prepare the SM tri-Higgs samples to identify the reason for the matching issue.

The SM tri-Higgs process is generated at the centre-of-mass energy  $\sqrt{s} = 13$  TeV with the NNPDF30\_nlo\_as\_0119 PDF set [9]. In Delphes, we use the anti- $k_{\rm T}$  clustering algorithm with R=0.4 and set the b-tagging efficiency to the ATLAS DL1r 77% working point. At this working point, the light-jet (charm-jet) rejection is about 130 (4.9). Following are the MadGraph scripts for generating SM tri-Higgs samples:

```
generate p p > h h h [QCD] QED^2<=6</pre>
output MG5/pphhh_sm
launch MG5/pphhh_sm
shower=Pythia8
detector=Delphes
analysis=OFF
madspin=ON
done
Cards/delphes_card.dat
set run card nevents 100000
set run card ebeam1 6500.0
set run card ebeam2 6500.0
set run_card pdlabel lhapdf
set run_card lhaid 266000
set run card ptb 19
set run_card etab 2.6
set spinmode none
decay h > b b~
```

done

Table 19 is the cutflow number of the SM tri-Higgs process. The matching efficiency is similar to the resonant tri-Higgs case (Table 16).

Table 19: The number of passing events, efficiencies, and passing rates for signal processes at different selection cuts.

	Count	Efficiency	Pass rate
Total	100000	1.00	1.00
$\geq 6 \text{ jets}$	74878	0.75	0.75
$\geq 4$ jets with $p_{\rm T} > 40$ GeV	70399	0.94	0.70
$\geq 4$ b-jets	48734	0.69	0.49
Matching 3h	17599	0.36	0.18
Matching 2h	18601	0.38	0.19
Matching 1h	10759	0.22	0.10

Table 20 and 21 are the pairing performance of the SM tri-Higgs process.

Table 20:  $\chi^2$  pairing efficiencies of the SM tri-Higgs samples. We minimize the quantity defined in Equation 3. The  $\chi^2$  method considers the possible combinations of 6 jets.

		Correc	etly reco			
$N_{ m Jet}$	Fraction	3h	2h	1h	0h	Higgs Efficiency
= 6	0.169	0.556	0.000	0.107	0.338	0.591
=7	0.311	0.334	0.008	0.154	0.505	0.390
$\geq 8$	0.520	0.193	0.007	0.204	0.596	0.266
Total	1.000	0.298	0.006	0.172	0.524	0.359

### 14 Use Pythia for Higgs decay

We use MadSpin to implement the Higgs decay in the previous exercise, while Pythia could also do it. Thus, we generate the samples with Pythia decay and compute the cutflow table.

Table 22 is the cutflow number of the Pythia decayed resonant tri-Higgs process. The matching efficiency is similar to the MadSpin case (Table 16).

# 15 Matching rate in different categories

Based on the *b*-jet multiplicity, we can categorize events into 4b, 5b, and 6b regions after Preselection. They are required to have exactly 4, exactly 5, or  $\geq 6$  *b*-tagged jets,

Table 21: Absolute value pairing efficiencies of the SM tri-Higgs samples. We minimize the quantity defined in Equation 8. The absolute value method considers the possible combination of 6 jets.

		Correc	etly reco			
$N_{ m Jet}$	Fraction	3h	2h	1h	0h	Higgs Efficiency
= 6	0.169	0.651	0.000	0.095	0.254	0.683
=7	0.311	0.401	0.014	0.158	0.427	0.463
$\geq 8$	0.520	0.226	0.009	0.213	0.552	0.303
Total	1.000	0.352	0.009	0.176	0.463	0.417

Table 22: The number of passing events, efficiencies, and passing rates for signal processes at different selection cuts.

	Count	Efficiency	Pass rate
Total	10000	1.00	1.00
$\geq 6 \text{ jets}$	5209	0.52	0.52
$\geq 4$ jets with $p_{\rm T} > 40$ GeV	4501	0.78	0.41
$\geq 4$ b-jets	2035	0.50	0.20
Matching 3h	705	0.35	0.07
Matching 2h	644	0.32	0.06
Matching 1h	575	0.28	0.06

respectively. We compute the matching efficiency and event fraction in each category to obtain more details about the samples. Similarly, we can compute the matching efficiency and event fraction in different  $N_{\text{Jet}}$  categories.

Table 23 and 24 are the matching rates in different  $N_{b\text{-Jet}}$  and  $N_{\text{Jet}}$  categories for resonant samples, respectively. Only the events passing the Preselection would be considered.

The matching rate in the 6b region is much higher than in the 4b and 5b regions. However, the event fraction of 6b categories is 12%. Therefore, the total matching efficiency is 31%. The matching rate in the 8j region is higher than in the 6j and 7j regions. If there is a jet not decaying from the b-parton, then the matching would fail in the 6j case. Thus, this result is satisfied our expectation.

Table 23: The matching rates on different  $N_{b\text{-Jet}}$  categories. The numerator is the number of events each b-parton can be matched to a jet. The denominator is the number of events passing the Preselection.

$N_{b ext{-Jet}}$	Fraction	Match Efficiency
=4	0.509	0.187
=5	0.368	0.354
$\geq 6$	0.123	0.669
Total	1.000	0.308

Table 24: The matching rates on different  $N_{\text{Jet}}$  categories. The numerator is the number of events each b-parton can be matched to a jet. The denominator is the number of events passing the Preselection.

$N_{ m Jet}$	Fraction	Match Efficiency
= 6	0.341	0.218
=7	0.320	0.322
$\geq 8$	0.338	0.384
Total	1.000	0.308

Table 25 and 26 are the matching rates in different  $N_{b\text{-Jet}}$  and  $N_{\text{Jet}}$  categories for SM samples, respectively.

Similarly, the matching rate in the 6b region is much higher than in the 4b and 5b regions. However, the event fraction of the 6b category is much lower than 4b and 5b regions. Thus, the total matching efficiency is 36%.

Table 25: The matching rates on different  $N_{b\text{-Jet}}$  categories. The numerator is the number of events each b-parton can be matched to a jet. The denominator is the number of events passing the Preselection.

$N_{b ext{-Jet}}$	Fraction	Match Efficiency
= 4	0.475	0.227
=5	0.378	0.395
$\geq 6$	0.147	0.704
Total	1.000	0.361

Table 26: The matching rates on different  $N_{\text{Jet}}$  categories. The numerator is the number of events each b-parton can be matched to a jet. The denominator is the number of events passing the Preselection.

$N_{ m Jet}$	Fraction	Match Efficiency
= 6	0.282	0.227
=7	0.307	0.352
$\geq 8$	0.411	0.460
Total	1.000	0.361

## 16 Pairing efficiency in different categories

Similarly, we can categorize events into 4b, 5b, and 6b regions after Preselection, which are required to have exactly 4, exactly 5, or  $\geq 6$  b-tagged jets, respectively. Then, we compute the pairing efficiency and event fraction in each category to obtain more details about the samples.

Table 27 is the pairing efficiency in different  $N_{b\text{-Jet}}$  categories for resonant samples. Only the events passing the Preselection and whose b-partons all can be matched would be considered. The pairing efficiency in the 6b region is higher than in the 4b and 5b regions. However, the event fraction of the 6b category only contributes 27%. Therefore, the total pairing efficiency is 25%. Note that these results are consistent with Table 18.

Table 28 is the pairing efficiency in different  $N_{b\text{-Jet}}$  categories for SM samples. Similarly, the pairing efficiency in the 6b region is higher than in the 4b and 5b regions. However, the event fraction of the 6b category is lower than the 4b and 5b regions. Thus, the total matching efficiency is 35%. These results are consistent with Table 21.

Table 27: Absolute value pairing efficiencies of the resonant tri-Higgs samples on different  $N_{b\text{-Jet}}$  categories. We minimize the quantity defined in Equation 8. The absolute value method considers the possible combination of 6 jets.

		Correc	etly reco			
$N_{b ext{-Jet}}$	Fraction	3h	2h	1h	0h	Higgs Efficiency
=4	0.309	0.190	0.008	0.189	0.614	0.258
=5	0.424	0.225	0.005	0.179	0.591	0.288
$\geq 6$	0.267	0.373	0.005	0.146	0.475	0.425
Total	1.000	0.254	0.006	0.173	0.567	0.316

Table 28: Absolute value pairing efficiencies of the SM tri-Higgs samples on different  $N_{b\text{-Jet}}$  categories. We minimize the quantity defined in Equation 8. The absolute value method considers the possible combination of 6 jets.

		Correc	etly reco			
$N_{b ext{-Jet}}$	Fraction	3h	2h	1h	0h	Higgs Efficiency
=4	0.301	0.295	0.011	0.181	0.514	0.362
=5	0.419	0.310	0.011	0.193	0.486	0.381
$\geq 6$	0.280	0.477	0.005	0.145	0.373	0.529
Total	1.000	0.352	0.009	0.176	0.463	0.417

## 17 Matching rate in more categories

Similarly to Section 15, we compute the matching efficiency and event fraction in more categories to obtain more details about the samples.

Table 29 is the matching rates in different  $N_{b\text{-Jet}}$  categories for resonant samples. Only the events passing the Preselection would be considered. The matching rate in the 6b region is similar to the 7b and 8b regions. However, the event fraction of the 6b category is much higher than the 7b and 8b regions. Therefore, the matching efficiency of the  $\geq 6b$  category is dominated by the 6b region.

Table 29: The matching rates on different  $N_{b\text{-Jet}}$  categories. The numerator is the number of events each b-parton can be matched to a jet. The denominator is the number of events passing the Preselection. No event belongs to the  $\geq 9b$  category.

$N_{b ext{-}\mathrm{Jet}}$	Fraction	Match Efficiency
= 4	0.509	0.187
=5	0.368	0.354
= 6	0.113	0.669
=7	0.009	0.674
= 8	0.001	0.650
$\geq 9$	0.000	nan
Total	1.000	0.308

Table 30 is the matching rates in different  $N_{b\text{-Jet}}$  categories for SM samples. Similarly, the matching rate in the 6b region is similar to the 7b and 8b regions. However, the event fraction of 6b categories is much higher than the 7b  $\sim$  10b regions. Therefore, the matching efficiency of the  $\geq$  6b category is dominated by the 6b region.

### 18 High resonant samples

We generate samples with high resonant mass. Here we choose three benchmarks:  $(m_{h_3}, m_{h_2}) = (1500 \text{ GeV}, 1000 \text{ GeV}), (1500 \text{ GeV}, 850 \text{ GeV}), (1300 \text{ GeV}, 1000 \text{ GeV}).$ 

Table 31 is the matching rates in different  $N_{b\text{-Jet}}$  categories for resonant samples. Only the events passing the Preselection would be considered. These three benchmark points have similar match rates in the 6b region. This value is also similar to the previous one (Table 23).

Table 30: The matching rates on different  $N_{b\text{-Jet}}$  categories. The numerator is the number of events each b-parton can be matched to a jet. The denominator is the number of events passing the Preselection. No event belongs to the  $\geq 11b$  category.

$N_{b ext{-Jet}}$	Fraction	Match Efficiency
=4	0.475	0.227
=5	0.378	0.395
=6	0.132	0.702
=7	0.014	0.723
= 8	0.001	0.742
=9	0.000	1.000
= 10	0.000	1.000
≥ 11	0.000	nan
Total	1.000	0.361

Table 31: The matching rates on different  $N_{b\text{-Jet}}$  categories. The numerator is the number of events each b-parton can be matched to a jet. The denominator is the number of events passing the Preselection.

(a) $m_{h_3} = 1500 \text{ GeV}, m_{h_2} = 1000 \text{ GeV}$			(b) $m_{h_3} = 1500 \text{ GeV}, m_{h_2} = 850 \text{ GeV}$			
$N_{b ext{-Jet}}$	Fraction	Match Efficiency	$N_{b ext{-}\mathrm{Jet}}$	Fraction	Match Efficiency	
=4	0.603	0.260	=4	0.588	0.279	
=5	0.316	0.408	=5	0.320	0.466	
$\geq 6$	0.082	0.662	$\geq 6$	0.092	0.672	
Total	1.000	0.340	Total	1.000	0.375	

(c) 
$$m_{h_3} = 1300 \text{ GeV}, m_{h_2} = 1000 \text{ GeV}$$

$N_{b ext{-Jet}}$	Fraction	Match Efficiency
=4	0.578	0.294
=5	0.324	0.447
$\geq 6$	0.098	0.654
Total	1.000	0.379

Table 32 is the pairing efficiency in different  $N_{b\text{-Jet}}$  categories for resonant samples. Only the events passing the Preselection and whose b-partons all can be matched would be considered. Compared with the low resonant case (Table 27), the high resonant samples have higher pairing efficiencies.

Table 32: Absolute value pairing efficiencies of the resonant tri-Higgs samples on different  $N_{b\text{-Jet}}$  categories. We minimize the quantity defined in Equation 8. The absolute value method considers the possible combination of 6 jets.

(a	) n	$n_{h_3}$	=	1500	GeV	$^{\prime},m_{h_2}$	=	1000	GeV
----	-----	-----------	---	------	-----	---------------------	---	------	-----

		Correc	etly reco			
$N_{b ext{-Jet}}$	Fraction	3h	2h	1h	0h	Higgs Efficiency
=4	0.435	0.467	0.037	0.185	0.310	0.554
=5	0.399	0.568	0.029	0.174	0.229	0.645
$\geq 6$	0.166	0.671	0.012	0.124	0.193	0.720
Total	1.000	0.541	0.030	0.171	0.258	0.618

(b) 
$$m_{h_3} = 1500 \text{ GeV}, m_{h_2} = 850 \text{ GeV}$$

		Correc	etly reco			
$\overline{N_{b ext{-}\mathrm{Jet}}}$	Fraction	3h	2h	1h	0h	Higgs Efficiency
= 4	0.432	0.474	0.040	0.199	0.288	0.566
=5	0.387	0.584	0.028	0.165	0.223	0.657
$\geq 6$	0.181	0.678	0.019	0.115	0.189	0.729
Total	1.000	0.553	0.032	0.170	0.245	0.631

(c) 
$$m_{h_3} = 1300 \text{ GeV}, m_{h_2} = 1000 \text{ GeV}$$

		Correc	etly reco			
$N_{b ext{-Jet}}$	Fraction	3h	2h	1h	0h	Higgs Efficiency
=4	0.447	0.448	0.033	0.157	0.361	0.523
=5	0.379	0.525	0.017	0.124	0.334	0.578
$\geq 6$	0.174	0.644	0.026	0.131	0.199	0.705
Total	1.000	0.511	0.026	0.140	0.323	0.575

### 19 Matching failed events

Even when considering the 6b category, the matching efficiencies remain around 66%. To better understand this, we must investigate the reasons for the failed matching cases.

Matching algorithm: For each simulated truth quark, we compute the  $\Delta R$  between the quark and all jets and then match the quark to the jet within  $\Delta R = 0.4$  cone. If more than one jet can be matched to a simulated truth quark in the  $\Delta R = 0.4$  cone, we choose the nearest one by the  $\Delta R$  distance. If multiple quarks are matched to the same jet, then the matching is failed.

Here, we only consider the 6b category. First, we compute the number of quarks that can be matched to jets, meaning that at least one jet is within the  $\Delta R = 0.4$  cone of the quark. Events with  $N_{\rm q} < 6$  are the failed matching cases. Table 33 presents the quark matching results.

For the  $N_{\rm q}=6$  case, if multiple quarks are matched to the same jet, then the matching fails. Two types of  $N_{\rm q}=6$  failed cases exist. One is for each quark that has only one jet within the  $\Delta R=0.4$  cone. Another is there is more than one jet within the  $\Delta R=0.4$  cone of quark. Table 34 summarizes the vairous cases in the  $N_{\rm q}=6$  category.

Table 33: The quark matching table.  $N_{\rm q}$  is the number of quarks that can be matched to jets in an event.

(a) m	(a) $m_{h_3} = 420 \text{ GeV},$										
$m_{h_2}$ =	= 280	0 GeV		(b) SM							
	$V_{ m q}$	Count	Fraction	$N_{ m q}$	Count	Fraction					
	3	8	0.002	3	4	0.001					
۷	4	89	0.022	4	156	0.022					
!	5	871	0.220	5	1505	0.210					
(	6	2998	0.756	6	5510	0.768					
То	tal	3966	1.000	Total	7175	1.000					

#### 19.1 $\Delta R$ distance

Table 33 shows that only 76% event can match all quarks to at least one jet. We try to modify the  $\Delta R$  requirement to study more details about matching failed events.

Table 35 presents the quark matching results, where  $\Delta R = 0.5$ . By increasing the  $\Delta R$  distance, the fraction of event with  $N_{\rm q} = 6$  is increased while the values are still lower than

Table 34: The various cases of quark matching results in the  $N_{\rm q}=6$  case.

(a) 
$$m_{h_3} = 420 \text{ GeV}, m_{h_2} = 280 \text{ GeV}$$

$N_{\rm q}=6$	Count	Fraction
Successful matching	2655	0.886
Only one matching jet for each quark	285	0.095
More than one jet for some quark	58	0.019
Total	2998	1.000

(b) SM

$N_{ m q}=6$	Count	Fraction
Successful matching	5054	0.917
Only one matching jet for each quark	376	0.068
More than one jet for some quark	80	0.015
Total	5510	1.000

80%.

Table 35: The quark matching table.  $N_{\rm q}$  is the number of quarks that can be matched to jets in an event.

(a) $m_{h_3} = 420 \text{ GeV},$								
$m_{h_2} = 280 \text{ GeV}$					(b) SM			
	$N_{ m q}$	Count	Fraction	$N_{ m q}$	Count	Fraction		
	3	7	0.002	3	2	0.000		
	4	73	0.018	4	113	0.016		
	5	768	0.194	5	1321	0.184		
	6	3118	0.786	6	5739	0.800		

Total

7175

1.000

# 20 Train SPANet with the samples in 6b region

1.000

In this section, we only consider the signal events in the 6b region, i.e., there  $\geq 6$  b-tagged jets. We prepare the samples in the 6b region for the SPA-NET training.

For the jet assignment part,

Total

3966

#### • Training sample:

- Total sample size: 360,000

- 1h sample size: 38,595

- 2h sample size: 73,036

- 3h sample size: 243,762

- 5% used on validation

#### • Testing sample:

- Total sample size: 40,000

- 1h sample size: 4,360

- 2h sample size: 8,070

- 3h sample size: 27,087

Table 36, 37 and 38 are the pairing performance on the 6b samples. The SPA-NET performs the worst. One reason is the sample is not enough for training. We need to generate more training samples to improve the training performance.

Table 36:  $\chi^2$  pairing efficiencies of the resonant tri-Higgs samples on different categories. The testing set is the 6b samples. We minimize the quantity defined in Equation 3. The  $\chi^2$  method considers the possible combinations of 6 highest  $p_T$  b-jets.

		Correc	ctly reco			
$N_{ m Jet}$	Fraction	3h	2h	1h	0h	Higgs Efficiency
= 6	0.211	0.509	0.000	0.124	0.367	0.551
=7	0.318	0.429	0.005	0.135	0.431	0.477
$\geq 8$	0.471	0.311	0.011	0.148	0.530	0.367
Total	1.000	0.390	0.007	0.139	0.464	0.441

Another test is that we prepare the training dataset by different selection criteria. We only require at least 4 b-jets for this dataset. Thus, the sample sizes are the following:

#### • Training sample:

- Total sample size: 900,000

- 1h sample size: 246,462

Table 37: Absolute value pairing efficiencies of the resonant tri-Higgs samples on different categories. The testing set is the 6b samples. We minimize the quantity defined in Equation 8. The absolute value method considers the possible combination of 6 highest  $p_{\rm T}$  b-jets.

		Correc	tly reco			
$N_{b ext{-Jet}}$	Fraction	3h	2h	1h	0h	Higgs Efficiency
= 4	0.211	0.456	0.000	0.126	0.418	0.498
=5	0.318	0.388	0.004	0.141	0.466	0.438
$\geq 6$	0.471	0.280	0.010	0.153	0.558	0.337
Total	1.000	0.352	0.006	0.143	0.499	0.403

Table 38: The pairing performance of SPA-NET trained on 6b datasets. The testing set is the 6b samples. The SPA-NET method considers all jets in the final state.

		Correctly reconstructed Higgs				
$N_{ m Jet}$	Fraction	3h	2h	1h	0h	Higgs Efficiency
= 6	0.211	0.417	0.000	0.137	0.446	0.463
=7	0.318	0.318	0.003	0.156	0.523	0.372
$\geq 8$	0.471	0.219	0.010	0.172	0.599	0.283
Total	1.000	0.292	0.006	0.159	0.543	0.349

- 2h sample size: 318,057

- 3h sample size: 280,788

- 5% used on validation

#### • Testing sample:

- Total sample size: 100,000

- 1h sample size: 27,243

- 2h sample size: 35,050

- 3h sample size: 31,499

Table 39 is the pairing performance of SPA-NET training on the 4b samples and testing on 6b events. The SPA-NET performs better than the previous one (Table 38). Even if we only test on 6b samples, the event with 4 or 5 b-jets can improve the training performance.

Table 39: The pairing performance of SPA-NET trained on the datasets with the at least 4 b-jets requirement. The testing set is the 6b samples. The SPA-NET method considers all jets in the final state.

		Correctly reconstructed Higgs				
$N_{ m Jet}$	Fraction	3h	2h	1h	0h	Higgs Efficiency
= 6	0.211	0.648	0.000	0.088	0.264	0.677
=7	0.318	0.539	0.005	0.120	0.336	0.582
$\geq 8$	0.471	0.390	0.013	0.160	0.437	0.452
Total	1.000	0.492	0.008	0.132	0.368	0.541

### 21 DNN classifier

To distinguish the signal and background events, we use the deep neural networks (DNNs) to classify events as signal-like or background-like.

### 21.1 Input features

Following are the input features:

•  $\Delta R_{H_1}$ ,  $\Delta R_{H_2}$ ,  $\Delta R_{H_3}$ : The angular distance between the 2 jets form the leading, subleading and least-leading Higgs Boson candidate, respectively.

- rms  $\Delta R_{\text{dijet}}$ : The root mean squre of the angular distance between all possible dijet that can form a Higgs Boson candidate.
- Skewness  $\Delta A_{\text{dijet}}$ : The skewness of  $\cosh(\Delta \eta_{ij}) \cos(\Delta \phi_{ij})$ , where i, j are all possible dijet that can form a Higgs Boson candidate.
- $H_{T, 6jets}$ : Scalar sum of the  $p_T$  of the 6 jets that can form 3 Higgs Boson candidates.
- $m_H \cos \theta$ :  $\theta$  is the angular between the vector of the reconstructed mass of the Higgs Boson candidates  $(m_{H_1}, m_{H_2}, m_{H_3})$  and (120, 115, 110) GeV and  $m_H = 125$  GeV.
- $\eta m_{HHH}$  fraction: Defiend as

$$\frac{\sum_{i,j} 2p_{\mathrm{T},i}p_{\mathrm{T},j} \left(\cosh\left(\Delta\eta_{ij}\right) - 1\right)}{m_{HHH}^2} \tag{9}$$

where i, j are all possible dijet that can form a Higgs Boson candidate and  $m_{HHH}$  is the reconstructed tri-Higgs invariant mass.

• Shpericity and Aplanarity of 6 jets. Compute the momentum tensor  $M_{xyz}$ 

$$M_{xyz} = \sum_{i} \begin{pmatrix} p_{xi}^{2} & p_{xi}p_{yi} & p_{xi}p_{zi} \\ p_{xi}p_{yi} & p_{yi}^{2} & p_{yi}p_{zi} \\ p_{xi}p_{zi} & p_{yi}p_{zi} & p_{zi}^{2} \end{pmatrix}$$
(10)

where i is the jet index. Its eigenvalues are ordered and normalised such that  $\lambda_1 > \lambda_2 > \lambda_3$ . Sphericity is dened as

$$S = \frac{3}{2} \left( \lambda_2 + \lambda_3 \right) \tag{11}$$

The aplanarity is defined as

$$A = \frac{3}{2}\lambda_3\tag{12}$$

# References

- [1] T.-K. Chen, C.-W. Chiang, and I. Low, "Simple model of dark matter and CP violation," *Phys. Rev. D*, vol. 105, no. 7, p. 075025, 2022.
- [2] M. J. Fenton, A. Shmakov, T.-W. Ho, S.-C. Hsu, D. Whiteson, and P. Baldi, "Permutationless many-jet event reconstruction with symmetry preserving attention networks," *Phys. Rev. D*, vol. 105, p. 112008, Jun 2022.

- [3] M. J. Fenton, A. Shmakov, H. Okawa, Y. Li, K.-Y. Hsiao, S.-C. Hsu, D. Whiteson, and P. Baldi, "Extended Symmetry Preserving Attention Networks for LHC Analysis," 9 2023.
- [4] A. Shmakov, M. J. Fenton, T.-W. Ho, S.-C. Hsu, D. Whiteson, and P. Baldi, "SPANet: Generalized permutationless set assignment for particle physics using symmetry preserving attention," *SciPost Phys.*, vol. 12, p. 178, 2022.
- [5] A. Papaefstathiou, G. Tetlalmatzi-Xolocotzi, and M. Zaro, "Triple Higgs boson production to six *b*-jets at a 100 TeV proton collider," *Eur. Phys. J. C*, vol. 79, no. 11, p. 947, 2019.
- [6] A. Papaefstathiou and G. Tetlalmatzi-Xolocotzi, "Multi-Higgs Boson Production with Anomalous Interactions at Current and Future Proton Colliders," 12 2023.
- [7] A. Papaefstathiou, T. Robens, and G. Tetlalmatzi-Xolocotzi, "Triple Higgs Boson Production at the Large Hadron Collider with Two Real Singlet Scalars," *JHEP*, vol. 05, p. 193, 2021.
- [8] B. Stanislaus, Searching for Beyond the Standard Model resonances in the  $HH \to b\bar{b}b\bar{b}$  final state using the ATLAS detector. PhD thesis, Oxford U., 2020.
- [9] R. D. Ball et al., "Parton distributions for the LHC Run II," JHEP, vol. 04, p. 040, 2015.