

Note

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1 Higgs Production

We aim to apply deep learning methods to distinguish vector boson fusion (VBF) from gluon-gluon fusion (GGF) and Higgs production at the Large Hadron Collider (LHC).

We want to apply the CWoLa method, so we can use the real data without knowing the true label.

2 Sample Preparation

2.1 Monte Carlo samples

We consider Standard Model (SM) di-photon Higgs events produced via GGF and VBF channels at a center-of-mass energy of $\sqrt{s} = 14$ TeV. The Higgs boson events are generated using **MadGraph** 3.3.1 [1] for both GGF and VBF production. The Higgs decays into the di-photon final state, and the parton showering and hadronization are simulated using **Pythia** 8.306 [2]. The detector simulation is conducted by **Delphes** 3.4.2 [3]. Jet reconstruction is performed using **FastJet** 3.3.2 [4] with the anti- k_t algorithm [5] and a jet radius of $R = 0.4$. These jets are required to have transverse momentum $p_T > 25$ GeV.

The following **MadGraph** scripts generate Monte Carlo samples for each production channel.

GGF Higgs Sample Generation

```
generate p p > h QCD<=99 [QCD]
output GGF_Higgs
launch GGF_Higgs
```

```
shower=Pythia8
detector=Delphes
```

```
analysis=OFF
madspin=OFF
done
```

```
set run_card nevents 100000
set run_card ebeam1 7000.0
set run_card ebeam2 7000.0
```

```
set run_card use_syst False
```

```
set pythia8_card 25:onMode = off
set pythia8_card 25:onIfMatch = 22 22
done
```

VBF Higgs Sample Generation

```
define v = w+ w- z
generate p p > h j j $$v
output VBF_Higgs
launch VBF_Higgs
```

```
shower=Pythia8
detector=Delphes
analysis=OFF
madspin=OFF
done
```

```
set run_card nevents 100000
set run_card ebeam1 7000.0
set run_card ebeam2 7000.0
```

```
set run_card use_syst False
```

```
set pythia8_card 25:onMode = off
set pythia8_card 25:onIfMatch = 22 22
done
```

2.2 Event selection

The selection cuts after the **Delphes** simulation:

- n_γ cut: The number of photons should be at least 2.
- n_j cut: The number of jets should be at least 2.
- $m_{\gamma\gamma}$ cut: The invariant mass of two leading photons $m_{\gamma\gamma}$ are required $120 \text{ GeV} \leq m_{\gamma\gamma} \leq 130 \text{ GeV}$.

Table 30 summarizes the cutflow number at different selection cuts.

Table 1: Number of passing events and passing rates for GGF and VBF Higgs production at different selection cuts.

Cut	GGF	pass rate	VBF	pass rate
Total	100000	1	100000	1
n_γ cut	48286	0.48	53087	0.53
n_j cut	9302	0.09	42860	0.43
$m_{\gamma\gamma}$ cut	8864	0.09	40694	0.41

Figure 1 shows the distributions of m_{jj} (the invariant mass of the two leading jets) and $\Delta\eta_{jj}$ (the pseudorapidity difference between the two leading jets). The scatter plot of m_{jj} versus $\Delta\eta_{jj}$ is presented in Figure 2.

2.3 Event image

The inputs for the neural networks are event images [6, 7, 8]. These images are constructed from events that pass the kinematic selection criteria described in section 2.2. Each event image has three channels corresponding to calorimeter towers, tracks, and photons. The following preprocessing steps are applied to all event constituents:

1. Translation: Compute the p_T -weighted center in the ϕ coordinates, then shift this point to the origin.
2. Flipping: Flip the highest p_T quadrant to the first quadrant.
3. Pixelation: Pixelate in a $\eta \in [-5, 5]$, $\phi \in [-\pi, \pi]$ box, with 40×40 pixels

Figure 3 shows the event images for GGF and VBF production modes.



Figure 1: Distributions of the invariant mass m_{jj} and pseudorapidity difference $\Delta\eta_{jj}$ of the two leading jets. Red dashed lines are selection cuts used to construct mixed datasets.



Figure 2: Scatter plot of m_{jj} versus $\Delta\eta_{jj}$. Red dashed lines are selection cuts used to construct mixed datasets.



(a) GGF: Calorimeter Tower



(b) VBF: Calorimeter Tower



(c) GGF: Track



(d) VBF: Track



(e) GGF: Photon



(f) VBF: Photon

Figure 3: Event images for GGF and VBF production, separately shown for calorimeter towers, tracks, and photons.

2.4 Mixed datasets

Based on figure 1, we set selection cuts of $m_{jj} > 300$ GeV and $\Delta\eta_{jj} > 3.1$. We consider three cases: applying each cut individually and simultaneously. These cuts define the signal region (SR), which is VBF-like, and the background region (BR), which is GGF-like. Table 2 summarizes the cutflow results for different selection criteria.

Table 2: Number of passing events and passing rates for GGF and VBF Higgs production under different selection cuts.

Cut	GGF	pass rate	VBF	pass rate
Total	100000	1.00	100000	1.00
n_γ cut	9302	0.09	42860	0.43
n_j cut	9302	0.09	42860	0.43
$m_{\gamma\gamma}$ cut	8864	0.09	40694	0.41
m_{jj} cut: SR	2695	0.03	29496	0.29
m_{jj} cut: BR	6169	0.06	11198	0.11
$\Delta\eta_{jj}$ cut: SR	2317	0.02	28160	0.28
$\Delta\eta_{jj}$ cut: BR	6547	0.07	12534	0.13
$m_{jj}, \Delta\eta_{jj}$ cuts: SR	1832	0.02	26446	0.26
$m_{jj}, \Delta\eta_{jj}$ cuts: BR	5684	0.06	9484	0.09

The total cross-section for VBF production is $\sigma_{\text{VBF}} = 4.278 \text{ pb}^{-1}$ at NNLO and for GGF production is $\sigma_{\text{GGF}} = 54.67 \text{ pb}^{-1}$ at N3LO, as referenced in [this link](#). The branching ratio for the di-photon decay channel is $\Gamma(h \rightarrow \gamma\gamma) = 2.270 \times 10^{-3}$, as given in [this link](#).

Assuming the luminosity of $\mathcal{L} = 300 \text{ fb}^{-1}$, we can estimate the number of events belonging to the SR and BR. These results are summarized in table 3.

3 Training CNN

The total sample sizes are mentioned in section 2.4. We allocate 80% of the data for training and 20% for validation. The testing set consists of the SR's 10,000 VBF and 10,000 GGF events.

The convolutional neural network (CNN) model structure is summarized in figure 4. The internal node uses the rectified linear unit (ReLU) as the activation function. The loss function is the binary cross-entropy. The Adam optimizer minimizes the loss value. The learning rate is 10^{-4} , and the batch size is 512. We employ the early stopping technique to

Table 3: The number of events of mixed datasets under different selection cuts.

(a) $m_{jj} > 300$ GeV			(b) $\Delta\eta_{jj} > 3.1$		
	GGF	VBF		GGF	VBF
BR	2297	326	BR	2437	365
SR	1003	859	SR	863	820

(c) $m_{jj} > 300$ GeV, $\Delta\eta_{jj} > 3.1$		
	GGF	VBF
BR	2116	276
SR	682	770

prevent over-training issues with patience of 10.

The training results are summarized in table 4. The performance of the $\Delta\eta_{jj}$ cuts is better than the m_{jj} cut. Moreover, when both cuts are applied together, the performance is slightly worse than when applying either cut individually.

Table 4: The CNN training results. The ACC and AUC are evaluated based on 10 training. The selection cuts of $m_{jj} > 300$ GeV and $\Delta\eta_{jj} > 3.1$ are applied.

Cut	M_1/M_2		S/B	
	ACC	AUC	ACC	AUC
m_{jj}	0.712 ± 0.023	0.741 ± 0.041	0.576 ± 0.010	0.596 ± 0.014
$\Delta\eta_{jj}$	0.828 ± 0.043	0.889 ± 0.050	0.604 ± 0.014	0.630 ± 0.015
$m_{jj}, \Delta\eta_{jj}$	0.753 ± 0.022	0.792 ± 0.035	0.573 ± 0.007	0.596 ± 0.008

3.1 More events

This section assumes the luminosity of $\mathcal{L} = 3000 \text{ fb}^{-1}$. The number of events belonging to the SR and BR are summarized in table 5.

The training results are summarized in table 6. All datasets' performance is better than the results in table 4. The $\Delta\eta_{jj}$ cut performs better than the m_{jj} cut. Moreover, when both cuts are applied together, the performance is slightly worse than the $\Delta\eta_{jj}$ cut but better than m_{jj} . These results are similar to the previous one.

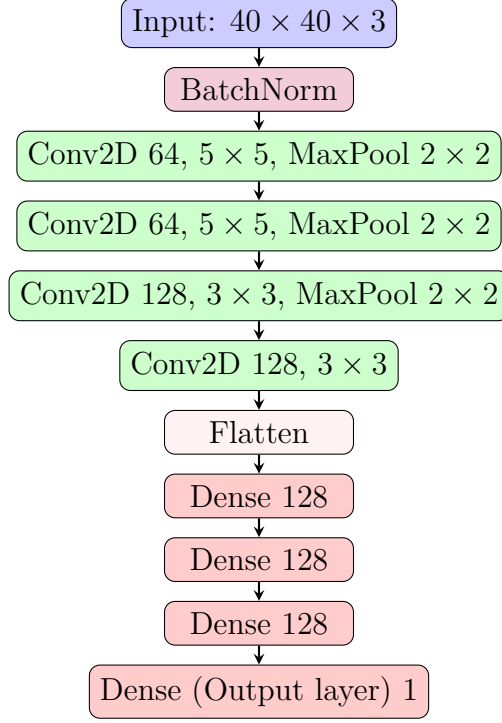


Figure 4: The architecture of the CNN model with key hyperparameters.

Table 5: The number of events of mixed datasets under different selection cuts.

(a) $m_{jj} > 300$ GeV			(b) $\Delta\eta_{jj} > 3.1$		
	GGF	VBF		GGF	VBF
BR	22967	3262	BR	24375	3652
SR	10034	8593	SR	8626	8204

(c) $m_{jj} > 300$ GeV, $\Delta\eta_{jj} > 3.1$		
	GGF	VBF
BR	21162	2763
SR	6821	7705

Table 6: The CNN training results. The ACC and AUC are evaluated based on 10 training. The selection cuts of $m_{jj} > 300$ GeV and $\Delta\eta_{jj} > 3.1$ are applied.

Cut	M_1/M_2		S/B	
	ACC	AUC	ACC	AUC
m_{jj}	0.907 ± 0.002	0.969 ± 0.002	0.598 ± 0.008	0.625 ± 0.009
$\Delta\eta_{jj}$	0.931 ± 0.004	0.979 ± 0.002	0.615 ± 0.005	0.648 ± 0.006
$m_{jj}, \Delta\eta_{jj}$	0.929 ± 0.003	0.978 ± 0.002	0.608 ± 0.004	0.638 ± 0.005

4 p_T normalization

To remove the potential dependence of the input samples on m_{jj} , we standardize the event images to remove the difference in input data distributions between the SR and BR. We calculate the mean and standard deviation of the event image transverse momentum and use these values to standardize each event image. We standardize each channel separately.

The number of events in the SR and BR are the same as previously in table 5.

The training results are summarized in table 7. The m_{jj} cut performs better than the previous one (table 6).

Table 7: The CNN training results with p_T normalization technique. The ACC and AUC are evaluated based on 10 training. The selection cuts of $m_{jj} > 300$ GeV and $\Delta\eta_{jj} > 3.1$ are applied.

Cut	M_1/M_2		S/B	
	ACC	AUC	ACC	AUC
m_{jj}	0.874 ± 0.004	0.946 ± 0.003	0.624 ± 0.005	0.663 ± 0.006
$\Delta\eta_{jj}$	0.928 ± 0.005	0.979 ± 0.002	0.597 ± 0.005	0.630 ± 0.006
$m_{jj}, \Delta\eta_{jj}$	0.917 ± 0.003	0.973 ± 0.002	0.603 ± 0.004	0.636 ± 0.006

5 Different cut setting

We set selection cuts of $m_{jj} > 225$ GeV and $\Delta\eta_{jj} > 2.3$ to ensure the SR and BR datasets have similar sizes. Table 8 summarizes the cutflow results for different selection criteria.

Assuming the luminosity of $\mathcal{L} = 3000 \text{ fb}^{-1}$, we can estimate the number of events belonging to the SR and BR. These results are summarized in table 9

Table 8: Number of passing events and passing rates for GGF and VBF Higgs production under different selection cuts.

Cut	GGF	pass rate	VBF	pass rate
Total	100000	1.00	100000	1.00
n_γ cut	9302	0.09	42860	0.43
n_j cut	9302	0.09	42860	0.43
$m_{\gamma\gamma}$ cut	8864	0.09	40694	0.41
m_{jj} cut: SR	3638	0.04	32993	0.33
m_{jj} cut: BR	5226	0.05	7701	0.08
$\Delta\eta_{jj}$ cut: SR	3611	0.04	32914	0.33
$\Delta\eta_{jj}$ cut: BR	5253	0.05	7780	0.08
$m_{jj}, \Delta\eta_{jj}$ cuts: SR	2842	0.03	31113	0.31
$m_{jj}, \Delta\eta_{jj}$ cuts: BR	4457	0.04	5900	0.06

Table 9: The number of events of mixed datasets under different selection cuts.

(a) $m_{jj} > 225$ GeV			(b) $\Delta\eta_{jj} > 2.3$		
	GGF	VBF		GGF	VBF
BR	19457	2244	BR	19557	2267
SR	13544	9612	SR	13444	9589

(c) $m_{jj} > 225$ GeV, $\Delta\eta_{jj} > 2.3$		
	GGF	VBF
BR	16594	1719
SR	10581	9064

The training results are summarized in table 10. The results are better than the table 7 by 1%. Similarly, the m_{jj} cut performs best.

Table 10: The CNN training results with p_T normalization technique. The ACC and AUC are evaluated based on 10 training. The selection cuts of $m_{jj} > 225$ GeV and $\Delta\eta_{jj} > 2.3$ are applied.

Cut	M_1/M_2		S/B	
	ACC	AUC	ACC	AUC
m_{jj}	0.864 ± 0.004	0.940 ± 0.004	0.632 ± 0.006	0.673 ± 0.007
$\Delta\eta_{jj}$	0.913 ± 0.006	0.972 ± 0.003	0.605 ± 0.007	0.640 ± 0.009
$m_{jj}, \Delta\eta_{jj}$	0.896 ± 0.007	0.961 ± 0.004	0.616 ± 0.005	0.653 ± 0.006

6 Supervised training

This section tests the supervised training on CNN. The training, validation, and testing sample sizes are summarized in table 11. The events passing all selection requirements (section 2.2) are considered.

Table 11: Sizes of various samples used for supervised training.

	Training	Validation	Testing
GGF	100k	25k	25k
VBF	100k	25k	25k

The training results are summarized in table 12. These results demonstrate the upper limit of CNN training.

Table 12: The CNN training results with p_T normalization technique. The ACC and AUC are evaluated based on 10 training.

ACC	AUC
0.784 ± 0.001	0.861 ± 0.001

6.1 Testing sample in SR and BR

The testing events used to evaluate the table 12 are all events passing the selection and not restricted to the particular SR. Thus, to make a fair comparison with previous results,

we must evaluate the training performance on the events in SR and BR.

The new testing dataset consists of the 10,000 VBF and 10,000 GGF events from SR and BR. The numbers of SR and BR events are computed from table 8.

The training results of table 10 are re-evaluated on the new testing set and shown in table 13. The results are better than the table 10. It seems that the events in the BR can be distinguished better than those in the SR.

Table 13: The CNN training results with p_T normalization technique. The ACC and AUC are evaluated based on 10 training. The selection cuts of $m_{jj} > 225$ GeV and $\Delta\eta_{jj} > 2.3$ are applied.

Cut	M_1/M_2		S/B	
	ACC	AUC	ACC	AUC
m_{jj}	0.863 ± 0.004	0.940 ± 0.002	0.716 ± 0.003	0.780 ± 0.004
$\Delta\eta_{jj}$	0.914 ± 0.004	0.972 ± 0.003	0.702 ± 0.003	0.754 ± 0.003
$m_{jj}, \Delta\eta_{jj}$	0.896 ± 0.006	0.962 ± 0.004	0.723 ± 0.003	0.780 ± 0.002

7 Use jet tagging results to construct mixed datasets

This section uses the jet tagging results to construct the mixed datasets.

Assuming the luminosity of $\mathcal{L} = 3000 \text{ fb}^{-1}$, we can estimate the number of events belonging to the SR and BR. The SR and BR are defined based on the number of gluon jets n_g and quark jets n_q . The selection results are summarized in table 31.

Table 14: The number of events of mixed datasets under different selection cuts. Here, $agbq$ means that $n_g = a, n_q = b$.

(a) SR: $2q0g$; BR: $1q1g, 0q2g$			(b) SR: $2q0g, 1q1g$; BR: $0q2g$			(c) SR: $2q0g$; BR: $0q2g$		
	GGF	VBF		GGF	VBF		GGF	VBF
SR	16828	10229	SR	30752	11779	SR	16828	10229
BR	16865	1596	BR	2941	47	BR	2941	47

For now, we use the true information from `Delphes` and do not consider the mis-tagging case.

The training results are summarized in table 15. All different jet-tagging conditions produced similar performance. However, the results are worse than those of kinematic cuts (table 13).

Table 15: The CNN training results with p_T normalization technique. The ACC and AUC are evaluated based on 10 training.

Datasets	M_1/M_2		S/B	
	ACC	AUC	ACC	AUC
SR: $2q0g$; BR: $1q1g, 0q2g$	0.623 ± 0.005	0.642 ± 0.005	0.653 ± 0.008	0.706 ± 0.009
SR: $2q0g, 1q1g$; BR: $0q2g$	0.934 ± 0.000	0.689 ± 0.012	0.662 ± 0.006	0.719 ± 0.008
SR: $2q0g$; BR: $0q2g$	0.900 ± 0.000	0.740 ± 0.010	0.655 ± 0.008	0.710 ± 0.009

The training results without p_T normalization are summarized in table 16. All different jet-tagging conditions produced similar performance. However, the results are worse than the ones with p_T normalization (table 15) by 2%.

Table 16: The CNN training results without p_T normalization technique. The ACC and AUC are evaluated based on 10 training.

Datasets	M_1/M_2		S/B	
	ACC	AUC	ACC	AUC
SR: $2q0g$; BR: $1q1g, 0q2g$	0.614 ± 0.007	0.632 ± 0.011	0.646 ± 0.008	0.690 ± 0.011
SR: $2q0g, 1q1g$; BR: $0q2g$	0.934 ± 0.000	0.695 ± 0.015	0.643 ± 0.009	0.689 ± 0.011
SR: $2q0g$; BR: $0q2g$	0.900 ± 0.000	0.743 ± 0.011	0.632 ± 0.007	0.677 ± 0.008

7.1 Loss weighted

Since the sample sizes are unbalanced, we add the class weights. The weights are proportional to the reciprocal of the number of events.

The training results with class weights are summarized in table 17. All different jet-tagging conditions produced similar performance.

8 Total scaling of transverse momentum

The p_T normalization removes the magnitude information of the input datasets. Thus, we would expect the training performance of the p_T normalization datasets would be worse than the one without it. However, table 15 and 16 shows the opposite results.

Table 17: The CNN training results without p_T normalization technique. The ACC and AUC are evaluated based on 10 training.

Datasets	M_1/M_2		S/B	
	ACC	AUC	ACC	AUC
SR: $2q0g$; BR: $1q1g, 0q2g$	0.621 ± 0.006	0.635 ± 0.007	0.645 ± 0.009	0.688 ± 0.013
SR: $2q0g, 1q1g$; BR: $0q2g$	0.934 ± 0.000	0.679 ± 0.016	0.624 ± 0.005	0.662 ± 0.008
SR: $2q0g$; BR: $0q2g$	0.900 ± 0.000	0.730 ± 0.013	0.621 ± 0.005	0.658 ± 0.008

To explore the reason why the p_T normalization could improve the training performance, we try the total p_T scaling, which computes the mean and standard deviation of all input samples. Then, use these values to standardize the input datasets.

8.1 Results

The training results with p_T scaling are summarized in table 18. All different jet-tagging conditions produced similar performance. However, the results are worse than the ones with p_T normalization (table 15).

Table 18: The CNN training results with p_T scaling technique. The ACC and AUC are evaluated based on 10 training. The selection cuts on the number of gluon jets are applied.

Datasets	M_1/M_2		S/B	
	ACC	AUC	ACC	AUC
SR: $2q0g$; BR: $1q1g, 0q2g$	0.622 ± 0.004	0.637 ± 0.008	0.638 ± 0.009	0.678 ± 0.011
SR: $2q0g, 1q1g$; BR: $0q2g$	0.934 ± 0.000	0.673 ± 0.032	0.619 ± 0.019	0.652 ± 0.029
SR: $2q0g$; BR: $0q2g$	0.900 ± 0.000	0.733 ± 0.011	0.621 ± 0.006	0.657 ± 0.009

The training results with p_T normalization are summarized in table 19.

The training results without p_T normalization are summarized in table 20.

9 Data augmentation

To improve the training performance, we will consider various data augmentation methods.

Table 19: The CNN training results with p_T normalization technique. The ACC and AUC are evaluated based on 10 training. The selection cuts on the number of gluon jets are applied.

Datasets	M_1/M_2		S/B	
	ACC	AUC	ACC	AUC
SR: 2q0g; BR: 1q1g, 0q2g	0.615 ± 0.005	0.632 ± 0.007	0.650 ± 0.011	0.703 ± 0.015
SR: 2q0g, 1q1g; BR: 0q2g	0.934 ± 0.000	0.662 ± 0.014	0.630 ± 0.008	0.675 ± 0.011
SR: 2q0g; BR: 0q2g	0.900 ± 0.000	0.716 ± 0.012	0.640 ± 0.007	0.690 ± 0.009

Table 20: The CNN training results without p_T normalization technique. The ACC and AUC are evaluated based on 10 training. The selection cuts on the number of gluon jets are applied.

Datasets	M_1/M_2		S/B	
	ACC	AUC	ACC	AUC
SR: 2q0g; BR: 1q1g, 0q2g	0.620 ± 0.004	0.636 ± 0.005	0.643 ± 0.006	0.686 ± 0.007
SR: 2q0g, 1q1g; BR: 0q2g	0.934 ± 0.000	0.680 ± 0.014	0.624 ± 0.010	0.660 ± 0.016
SR: 2q0g; BR: 0q2g	0.900 ± 0.000	0.727 ± 0.010	0.628 ± 0.008	0.666 ± 0.011

9.1 p_T smearing

The p_T smearing method simulates detector resolution effects on the transverse momentum of event constituents. This method resamples the transverse momentum p_T of event constituents according to the normal distribution:

$$p'_T \sim \mathcal{N}(p_T, f(p_T)), \quad f(p_T) = \sqrt{0.052p_T^2 + 1.502p_T}, \quad (1)$$

where p'_T is the augmented transverse momentum, and $f(p_T)$ is the energy smearing function applied by **Delphes** (the p_T 's are normalized in units of GeV). The preprocessing is applied after the p_T smearing augmentation.

The training results of the 2q0g datasets (Table 31 (a)) are summarized in table 21.

9.2 ϕ shifting

The ϕ shifting method shifts entire events by a random angle $\Delta\phi \in [-\pi, \pi]$ to enlarge the diversity of training datasets.

The training results of the 2q0g datasets are summarized in table 22.

Table 21: CNN training results with different augmentation sizes. The ACC and AUC are evaluated based on 10 training.

Datasets	M_1/M_2		S/B	
	ACC	AUC	ACC	AUC
Original	0.615 ± 0.005	0.632 ± 0.007	0.650 ± 0.011	0.703 ± 0.015
+5	0.625 ± 0.006	0.653 ± 0.009	0.661 ± 0.010	0.714 ± 0.012
+10	0.629 ± 0.005	0.658 ± 0.005	0.666 ± 0.008	0.721 ± 0.009
+15	0.629 ± 0.003	0.660 ± 0.003	0.661 ± 0.015	0.710 ± 0.018

Table 22: CNN training results with different augmentation sizes. The ACC and AUC are evaluated based on 10 training.

Datasets	M_1/M_2		S/B	
	ACC	AUC	ACC	AUC
Original	0.615 ± 0.005	0.632 ± 0.007	0.650 ± 0.011	0.703 ± 0.015
+5	0.641 ± 0.003	0.680 ± 0.004	0.683 ± 0.010	0.736 ± 0.013
+10	0.642 ± 0.006	0.684 ± 0.008	0.686 ± 0.008	0.739 ± 0.011
+15	0.643 ± 0.005	0.685 ± 0.006	0.687 ± 0.009	0.742 ± 0.010

9.3 $\eta - \phi$ smearing

We apply the $\eta - \phi$ smearing on the training samples. Specifically, the (η, ϕ) coordinates of constituents are resampled according to a normal distribution centered on the original coordinate and with a standard deviation inversely proportional to the p_T

$$\eta' \sim \mathcal{N}\left(\eta, \frac{\Lambda}{p_T}\right), \quad \phi' \sim \mathcal{N}\left(\phi, \frac{\Lambda}{p_T}\right) \quad (2)$$

where η', ϕ' are the augmented coordinates, p_T is the transverse momentum of the constituent, and the smearing scale is set to be $\Lambda = 100$ MeV.

The training results on the $2q0g$ datasets are summarized in Table 23. The +5 and +10 augmentation cases show performance comparable to the original dataset. However, applying +15 augmentations degrades the performance, suggesting that introducing too many augmented samples may lead the training in the wrong direction.

9.4 Without pre-processing

The ϕ shifting seems to cancel the ϕ translation in the pre-processing. Thus, we expect the model trained on the ϕ shifting dataset could perform similarly to the no pre-processing

Table 23: CNN training results with different augmentation sizes. The ACC and AUC are evaluated based on 10 training.

Datasets	M_1/M_2		S/B	
	ACC	AUC	ACC	AUC
Original	0.615 ± 0.005	0.632 ± 0.007	0.650 ± 0.011	0.703 ± 0.015
+5	0.618 ± 0.004	0.640 ± 0.006	0.658 ± 0.009	0.711 ± 0.013
+10	0.617 ± 0.004	0.641 ± 0.006	0.654 ± 0.010	0.705 ± 0.012
+15	0.612 ± 0.006	0.628 ± 0.008	0.635 ± 0.009	0.679 ± 0.013

datasets.

The testing results of the $2q0g$ datasets are summarized in table 24. The performance of pre-processing datasets is generally better than that without pre-processing. The reason may be that the original datasets are applied pre-processed. Thus, the samples have higher density for the ϕ center at 0. The model would prefer to learn these events first.

We can train the model on only the augmented datasets to ensure the effect of the original samples.

Table 24: CNN training results with different augmentation sizes. The ACC and AUC are evaluated based on 10 training.

Datasets	w/ pre-processing		w/o pre-processing	
	ACC	AUC	ACC	AUC
Original	0.637 ± 0.007	0.686 ± 0.008	0.625 ± 0.006	0.669 ± 0.008
+5	0.682 ± 0.011	0.735 ± 0.013	0.669 ± 0.011	0.720 ± 0.015
+10	0.685 ± 0.008	0.739 ± 0.010	0.673 ± 0.008	0.726 ± 0.010
+15	0.688 ± 0.007	0.743 ± 0.009	0.674 ± 0.008	0.726 ± 0.008

9.5 Only augmentation datasets

In section 9.4, we found that the performance of pre-processing datasets is generally better than without pre-processing. We train the model on only the augmented datasets to ensure the effect of the original samples.

The testing results of the only augmented sample are summarized in table 25. The performance without original samples is similar to that with original samples. It seems that the impact of the original datasets is limited. For 10, 15 augmentation cases, with and without original samples perform almost the same.

Table 25: CNN training results with different augmentation sizes. The ACC and AUC are evaluated based on 10 training. Here, $+x$ contains original and augmented samples; $=x$ contains only augmented samples.

Datasets	w/ pre-processing		w/o pre-processing	
	ACC	AUC	ACC	AUC
+5	0.682 ± 0.011	0.735 ± 0.013	0.669 ± 0.011	0.720 ± 0.015
=5	0.682 ± 0.007	0.736 ± 0.009	0.668 ± 0.007	0.718 ± 0.010
+10	0.685 ± 0.008	0.739 ± 0.010	0.673 ± 0.008	0.726 ± 0.010
=10	0.687 ± 0.010	0.740 ± 0.012	0.675 ± 0.009	0.726 ± 0.011
+15	0.688 ± 0.007	0.743 ± 0.009	0.674 ± 0.008	0.726 ± 0.008
=15	0.687 ± 0.007	0.741 ± 0.010	0.672 ± 0.009	0.725 ± 0.012

10 Removing photon information

To investigate the role of photon information in model training, we conduct two exercises:

- **Case 1:** Remove both the photon channel and photon features in the Tower channel.
- **Case 2:** Remove the photon channel.
- **Case 3:** Remove the photon features in the Tower channel.

We consider the $2q0g$ datasets. The training results are summarized in Tables 26, 27, and 28.

Table 26: CNN training results with both the photon channel and Tower photon features removed (Case 1). The ACC and AUC are evaluated based on 10 training.

Datasets	M_1/M_2		S/B	
	ACC	AUC	ACC	AUC
Original	0.633 ± 0.005	0.664 ± 0.008	0.690 ± 0.005	0.750 ± 0.008
+5	0.644 ± 0.004	0.687 ± 0.005	0.693 ± 0.007	0.746 ± 0.006
+10	0.645 ± 0.004	0.689 ± 0.005	0.697 ± 0.009	0.751 ± 0.011
+15	0.645 ± 0.004	0.689 ± 0.005	0.698 ± 0.009	0.753 ± 0.010

Case 1 leads the performance improvement compared to the full-feature baseline (Table 22). Case 2 shows a little better performance of the all-feature input case. Case 3 demonstrates more improvement than case 2 but is still worse than case 1.

Table 27: CNN training results with only the photon channel removed (Case 2). The ACC and AUC are evaluated based on 10 training.

Datasets	M_1/M_2		S/B	
	ACC	AUC	ACC	AUC
Original	0.621 ± 0.005	0.640 ± 0.008	0.661 ± 0.007	0.715 ± 0.010
+5	0.636 ± 0.004	0.673 ± 0.006	0.673 ± 0.009	0.727 ± 0.011
+10	0.639 ± 0.004	0.677 ± 0.005	0.677 ± 0.007	0.728 ± 0.010
+15	0.640 ± 0.005	0.679 ± 0.007	0.678 ± 0.007	0.731 ± 0.010

Table 28: CNN training results with the Tower photon features removed (Case 3). The ACC and AUC are evaluated based on 10 training.

Datasets	M_1/M_2		S/B	
	ACC	AUC	ACC	AUC
Original	0.629 ± 0.004	0.653 ± 0.006	0.670 ± 0.008	0.726 ± 0.009
+5	0.642 ± 0.005	0.682 ± 0.008	0.690 ± 0.010	0.742 ± 0.011
+10	0.644 ± 0.003	0.686 ± 0.004	0.692 ± 0.006	0.746 ± 0.008
+15	0.645 ± 0.005	0.687 ± 0.007	0.690 ± 0.007	0.745 ± 0.007

11 Various decay channels

In this section, we consider $H \rightarrow W^+W^-$, $H \rightarrow ZZ$, and $H \rightarrow \tau^+\tau^-$.

11.1 Final state

Unlike the $H \rightarrow \gamma\gamma$ analysis, the final states of $H \rightarrow WW^*$, ZZ^* , and $\tau^+\tau^-$ are more diverse. Each decay mode offers several final-state configurations, depending on whether the intermediate particles decay leptonically or hadronically.

11.1.1 $H \rightarrow WW^*$

- **Fully leptonic:** $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$
- **Semi-leptonic:** $H \rightarrow WW^* \rightarrow \ell\nu jj$
- **Fully hadronic:** $H \rightarrow WW^* \rightarrow jjjj$

11.1.2 $H \rightarrow ZZ^*$

- **Fully leptonic:** $H \rightarrow ZZ^* \rightarrow 4\ell$
- **Semi-leptonic:** $H \rightarrow ZZ^* \rightarrow 2\ell 2j$
- **Invisible + leptonic:** $H \rightarrow ZZ^* \rightarrow 2\ell 2\nu$
- **Fully hadronic:** $H \rightarrow ZZ^* \rightarrow 4j$

11.1.3 $H \rightarrow \tau^+\tau^-$

- **Leptonic–Leptonic:** $H \rightarrow \tau^+\tau^- \rightarrow \ell\nu\bar{\nu} \ell\nu\bar{\nu}$
- **Leptonic–Hadronic:** $H \rightarrow \tau^+\tau^- \rightarrow \ell\nu\bar{\nu} \tau_{\text{had}}$
- **Hadronic–Hadronic:** $H \rightarrow \tau^+\tau^- \rightarrow \tau_{\text{had}}\tau_{\text{had}}$

11.2 Pre-selection cuts

Each decay mode requires tailored pre-selection cuts to suppress background while retaining a reasonable signal efficiency. Some considerations:

- How should pre-selection cuts be defined for each decay channel?
- Should we unify pre-selection criteria to allow CNN to transfer across channels?

Answers:

- Consider the **Fully leptonic:** $H \rightarrow ZZ^* \rightarrow 4\ell$ channel. Only require the number of leptons and jets.
- The invariant mass cuts can be removed.

12 ϕ shifting with fixed angle

In section 9.2, we implemented ϕ -shifting augmentation using a random rotation angle for each event. In this section, we consider shifting the azimuthal angle ϕ by fixed values and investigate its effect on training performance.

The training results on the $2q0g$ dataset are summarized in table 29. We find that using fixed-angle ϕ shifts yields performance comparable to the random-angle augmentation results shown previously in table 22.

Table 29: CNN training results using fixed-angle ϕ -shifting augmentation. Here, $360/\theta$ denotes that events were augmented using rotation angles of θ , 2θ , ..., up to 360° . The ACC and AUC are evaluated based on 10 training.

Datasets	M_1/M_2		S/B	
	ACC	AUC	ACC	AUC
Original	0.615 ± 0.005	0.632 ± 0.007	0.650 ± 0.011	0.703 ± 0.015
360/90	0.637 ± 0.005	0.675 ± 0.007	0.684 ± 0.010	0.739 ± 0.012
360/60	0.642 ± 0.003	0.682 ± 0.005	0.689 ± 0.007	0.745 ± 0.010
360/45	0.643 ± 0.003	0.685 ± 0.004	0.689 ± 0.008	0.742 ± 0.012
360/30	0.643 ± 0.006	0.684 ± 0.009	0.688 ± 0.007	0.744 ± 0.008

13 $H \rightarrow ZZ^* \rightarrow 4\ell$ channel

13.1 Sample preparation

We consider SM Higgs decay into ZZ via GGF and VBF channels at a center-of-mass energy of $\sqrt{s} = 14$ TeV. We focus on the fully leptonic mode: $H \rightarrow ZZ^* \rightarrow 4\ell$. The Higgs boson events are generated using **MadGraph** 3.3.1 [1] for both GGF and VBF production. The parton showering and hadronization are simulated using **Pythia** 8.306 [2]. The detector simulation is conducted by **Delphes** 3.4.2 [3]. Jet reconstruction is performed using **FastJet** 3.3.2 [4] with the anti- k_t algorithm [5] and a jet radius of $R = 0.4$. These jets are required to have transverse momentum $p_T > 25$ GeV.

The following **MadGraph** scripts generate Monte Carlo samples for each production channel.

GGF Higgs Sample Generation

```
import model loop_sm
generate p p > h > l+ l- l+ l- QCD=0 QED<=4 [noborn=QCD]
output GGF_Higgs_ZZ_4l
launch GGF_Higgs_ZZ_4l

shower=Pythia8
detector=Delphes
analysis=OFF
madspin=OFF
done
```

Cards/delphes_card.dat

```
set run_card nevents 10000
set run_card ebeam1 7000.0
set run_card ebeam2 7000.0
```

```
set run_card use_syst False
```

done

VBF Higgs Sample Generation

```
define v = w+ w- z
generate p p > h j j $$v, (h > z z , z > l+ l- , z > l+ l-) QCD<=99
output VBF_Higgs_ZZ_4l
launch VBF_Higgs_ZZ_4l
```

```
shower=Pythia8
detector=Delphes
analysis=OFF
madspin=OFF
done
```

Cards/delphes_card.dat

```
set run_card nevents 10000
set run_card ebeam1 7000.0
set run_card ebeam2 7000.0
```

```
set run_card use_syst False
```

done

The selection cuts after the **Delphes** simulation:

- n_l cut: The number of leptons should be at least 4.

- n_j cut: The number of jets should be at least 2.

Table 30 summarizes the cutflow number at different selection cuts.

Table 30: Number of passing events and passing rates for GGF and VBF Higgs production at different selection cuts.

Cut	GGF	pass rate	VBF	pass rate
Total	10000	1	10000	1
n_l cut	2731	0.27	1902	0.19
n_j cut	687	0.07	1650	0.17

The branching ratio for the 4-lepton channel is $\Gamma(h \rightarrow 4\ell, \ell = e, \mu) = 1.240 \times 10^{-4}$, as given in [this link](#). Assuming the luminosity of $\mathcal{L} = 3000 \text{ fb}^{-1}$, we can estimate the number of events belonging to the SR and BR. The SR and BR are defined based on the number of gluon jets n_g and quark jets n_q . The selection results are summarized in table 31.

Table 31: The number of events of mixed datasets under different selection cuts. Here, $agbq$ means that $n_g = a, n_q = b$.

(a) SR: $2q0g$; BR: $1q1g, 0q2g$			(b) SR: $2q0g, 1q1g$; BR: $0q2g$		
	GGF	VBF		GGF	VBF
SR	722	228	SR	1287	261
BR	704	34	BR	138	1

13.2 Event image of $H \rightarrow ZZ^* \rightarrow 4\ell$ mode

We follow the same procedure described in section 2.3 to construct the event image. Each event image contains only two channels: calorimeter towers and tracks. To enable transferability of the trained CNN model across different decay modes, we remove decay product (lepton) information from the event images.

Figure 5 shows representative event images for the $H \rightarrow ZZ^* \rightarrow 4\ell$ mode, separated by production mechanism (GGF and VBF) and by feature type (calorimeter tower or track).

13.3 Testing results of di-photon classifier

We evaluate the performance of the $H \rightarrow \gamma\gamma$ classifier trained in section 10 on a different decay mode: $H \rightarrow ZZ^* \rightarrow 4\ell$. We focus on Case 1, where both the photon channel and

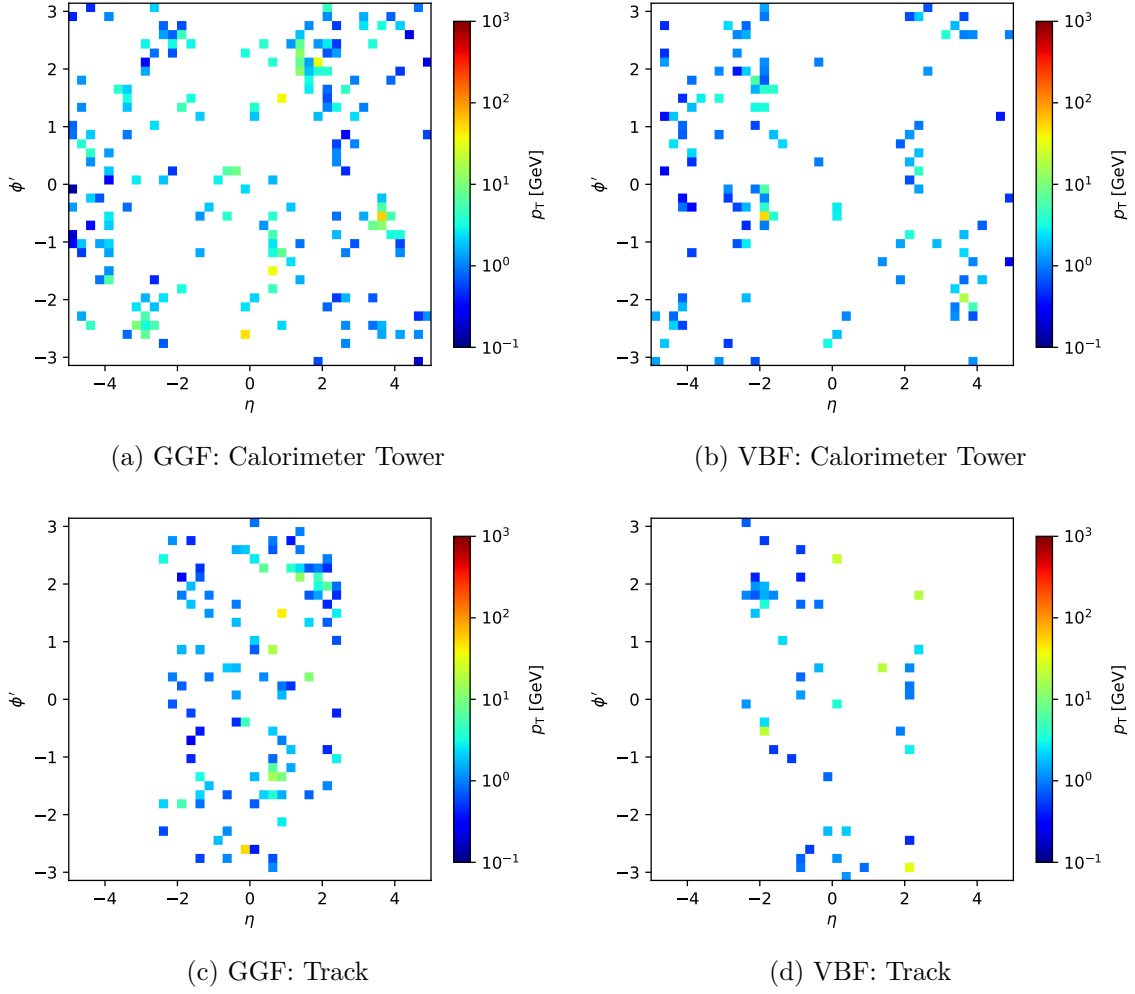


Figure 5: Event images for $H \rightarrow ZZ^* \rightarrow 4\ell$ events produced via GGF and VBF. Images are shown separately for the calorimeter tower and track channels.

photon-related features in the tower channel are removed from the input. This setting is designed to test whether a model trained without explicit decay production information can still extract meaningful patterns applicable to other decay modes.

The evaluation is performed using the $2q0g$ dataset for both decay channels. Table 32 summarizes the results. These results indicate that, although the classifier performs well on

Table 32: CNN training results for Case 1: both the photon channel and Tower photon-related features are removed. The classifier is trained on $H \rightarrow \gamma\gamma$ and evaluated on both $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^* \rightarrow 4\ell$. The ACC and AUC are evaluated based on 10 training.

Datasets	$H \rightarrow \gamma\gamma$		$H \rightarrow ZZ^* \rightarrow 4\ell$	
	ACC	AUC	ACC	AUC
Original	0.690 ± 0.005	0.750 ± 0.008	0.621 ± 0.022	0.665 ± 0.027
+5	0.695 ± 0.008	0.747 ± 0.005	0.592 ± 0.012	0.624 ± 0.016
+10	0.698 ± 0.008	0.752 ± 0.010	0.593 ± 0.013	0.627 ± 0.020
+15	0.700 ± 0.010	0.755 ± 0.010	0.589 ± 0.012	0.621 ± 0.019

its original $H \rightarrow \gamma\gamma$ dataset, its performance degrades when applied to the $H \rightarrow ZZ^* \rightarrow 4\ell$ events. This suggests that even after removing photon-related features, the learned representation may still carry decay-mode-specific biases that affect cross-channel transferability.

13.4 Testing results of supervised classifier

We train supervised CNN models separately on two Higgs decay channels. The training, validation, and testing dataset sizes are identical to those listed in table 11. Only events that pass all selection requirements are used. Importantly, all input features directly associated with the decay products are removed.

Table 33 summarizes the classification results. Each model is trained on one decay mode and evaluated on both $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^* \rightarrow 4\ell$ datasets. Compared to the results in table 12, the di-photon classifier shows slightly worse performance due to the exclusion of photon-specific information.

Although both models are trained without decay product features, the CNN still captures decay-mode-specific differences. Moreover, the model trained on $H \rightarrow ZZ^* \rightarrow 4\ell$ generalizes to the $H \rightarrow \gamma\gamma$ dataset worse than the reverse case.

Table 33: CNN classification results with decay product information removed. Each model is trained on one decay channel and tested on both $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^* \rightarrow 4\ell$. Results are averaged over 10 training runs.

Training channel	$H \rightarrow \gamma\gamma$		$H \rightarrow ZZ^* \rightarrow 4\ell$	
	ACC	AUC	ACC	AUC
$H \rightarrow \gamma\gamma$	0.775 ± 0.001	0.852 ± 0.001	0.752 ± 0.002	0.827 ± 0.003
$H \rightarrow ZZ^* \rightarrow 4\ell$	0.738 ± 0.002	0.806 ± 0.003	0.793 ± 0.001	0.872 ± 0.001

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