

# RQI with a Multigrid in Energy Preconditioner for Massively Parallel Neutron Transport

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M&C + SNA + MC 2015 Nashville, TN, April 23, 2015



#### **Outline**

- Background
- Block Krylov
- RQI
- MGE Preconditioner
- Results
- Conclusions

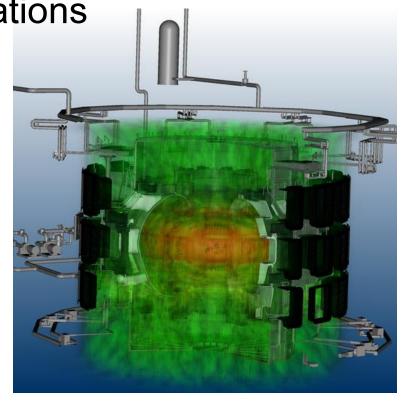




### **Large Problems and HPC**

 Systems of interest can be geometrically large and/or require fine discretizations

- We also want to do coupled multi-physics
- High performance computers (HPC) have big memories, many cores
- Enable solution of "grand challenge" problems



Major radius of 6.2m

http://www.efda.org/the\_iter\_project/iter\_ the machine.htm

# Transport Eqn. Operator Form

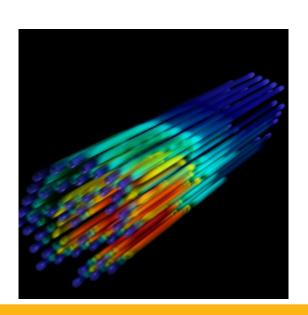
- After discretization, operator form of transport eqn:
- $\mathbf{L} = \hat{\Omega} \cdot \nabla + \Sigma$  is the transport operator,
- M converts harmonic moments into discrete angle sets,
- $\mathbf{D} = \mathbf{M}^T \mathbf{W} = \sum_{a=1}^n Y_{lm}^{e/o} w_a$  is the discrete-to-moment operator.
- f contains the fission source,  $\nu\Sigma_f$ ;
- $\mathbf{F} = \chi f^T$ ,
- **S** is the scattering matrix,

$$\mathbf{L}\psi = \mathbf{MS}\phi + \frac{1}{k}\mathbf{MF}\phi$$
$$\phi = \mathbf{D}\psi$$



# Denovo's Methods (used here)

- Within-group solvers:
  - Krylov (GMRES, BiCGStab), source iteration
- Multigroup Solvers:
  - Transport Two-Grid upscatter acceleration of Gauss Seidel
  - Block Krylov (GMRES, BiCGStab)
  - Multigrid in energy preconditioning
- Eigenvalue solvers:
  - Power iteration (rebalance, CMFD)
  - Rayleigh Quotient Iteration





## **Multigroup Solution**

- Gauss Seidel does energy iterations over the G space-angle inner iterations
- Space-angle done with Krylov on single group
- Fundamentally serial in energy

- Block Krylov puts a block of groups at the same iteration level
- Handles space-angle and energy all at once
- Krylov should converge more quickly than GS



### **Block Krylov: One Iteration Level**

Apply A to whole upscatter block at once

$$\mathbf{S} = \begin{pmatrix} \mathbf{S}_{11} & 0 & 0 & 0 & 0 \\ \mathbf{S}_{21} & \mathbf{S}_{22} & 0 & 0 & 0 \\ \mathbf{S}_{31} & \mathbf{S}_{32} & \mathbf{S}_{33} & \mathbf{S}_{34} & \mathbf{S}_{35} \\ \mathbf{S}_{41} & \mathbf{S}_{42} & \mathbf{S}_{43} & \mathbf{S}_{44} & \mathbf{S}_{45} \\ \mathbf{S}_{51} & \mathbf{S}_{52} & \mathbf{S}_{53} & \mathbf{S}_{54} & \mathbf{S}_{55} \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \\ v_3 \\ v_4 \end{pmatrix}$$
 Upscatter RHS Upscatter Block 
$$(\mathbf{I} - \mathbf{D} \mathbf{L}^{-1} \mathbf{M} \mathbf{S}_{\text{up-block}}) \phi = \mathbf{D} \mathbf{L}^{-1} \mathbf{M} (q_e^{\text{up}})$$

- Iteration vector includes all upscattering groups
- Enables parallelization in energy ("multisets")



#### **Power Iteration**

• Ordinary eigenvalue equation traditionally solved with Power Iteration (PI) [define  $T = DL^{-1}$ ]:

$$\mathbf{A} = (\mathbf{I} - \mathbf{TMS})^{-1} \mathbf{TMF}$$

$$\sigma(\mathbf{A}) \equiv \{ \lambda \in \mathbb{C} : rank(\mathbf{A} - \lambda \mathbf{I}) \leq n \}$$

$$\mathbf{A}^{j} x_{i} = \lambda_{i}^{j} x_{i}$$

$$\phi^{j+1} = \frac{1}{k^{j}} \mathbf{A} \phi^{j} \qquad k^{j+1} = k^{j} \frac{f^{T} \phi^{j+1}}{f^{T} \phi^{j}}$$

Can be quite slow for dominance ratios close to 1

$$e^{j+1} \approx C |\frac{\lambda_2}{\lambda_1}| e^{j}$$



#### **Shifted Inverse Iteration is Faster**

- For some shift,  $\mu$ ,  $(\mathbf{A} \mu \mathbf{I})$  has same eigenvectors as  $\mathbf{A}$
- If invertible,  $\sigma([\mathbf{A} \mu \mathbf{I}]^{-1}) = \{1/(\lambda \mu) : \lambda \in \sigma(\mathbf{A})\}$
- Eigenvalues near shift become separated

$$\kappa_1 = \frac{1}{\lambda_1 - \mu}, \, \kappa_2 = \frac{1}{\lambda_2 - \mu}, \, \dots, \, \kappa_n = \frac{1}{\lambda_n - \mu}$$

• As  $\mu \to \lambda_1$ ,  $\kappa_1 \to \infty$ ; other terms stay finite

$$e^{i+1} \approx C \left| \frac{\lambda_1 - \mu}{\lambda_2 - \mu} \right| e^{i}$$

• Like power iteration on  $(\mathbf{A} - \mu \mathbf{I})^{-1}$ 



#### **RQI** is Faster Still

• For problems that look like  $\mathbf{A}x = \lambda \mathbf{B}x$ , the RQ is:

$$\rho(x) = \frac{x^T \mathbf{A} x}{x^T \mathbf{B} x}$$

- RQI: shifted inverse iteration with an optimal shift
- Better convergence properties than PI

$$(\mathbf{I} - \mathbf{TM\tilde{S}})\phi = (\gamma - \rho)\mathbf{TMF}\phi$$
  
 $\tilde{\mathbf{S}} \equiv \mathbf{S} + \rho\mathbf{F}$ 

Shifted matrix looks like an energy-block dense scattering



# **Block Krylov and RQI**

- Shifted system would be difficult for GS
- Systems that are block-dense in energy are still sparse in energy-space-angle
- Structurally ideal for Krylov methods
- However, RQI creates poorly conditioned systems:

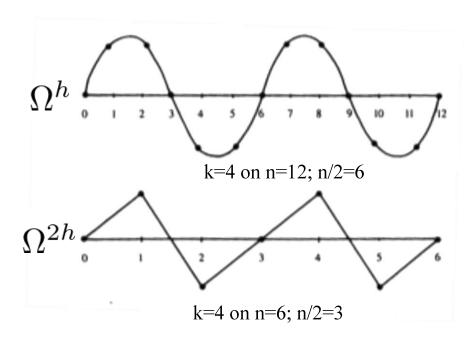
$$\kappa(\mathbf{A}) = ||\mathbf{A}|| ||\mathbf{A}^{-1}||$$

- Krylov methods can have difficulty
- Result: cannot get a valid RQ; RQI doesn't converge



# **Making Error Oscillatory**

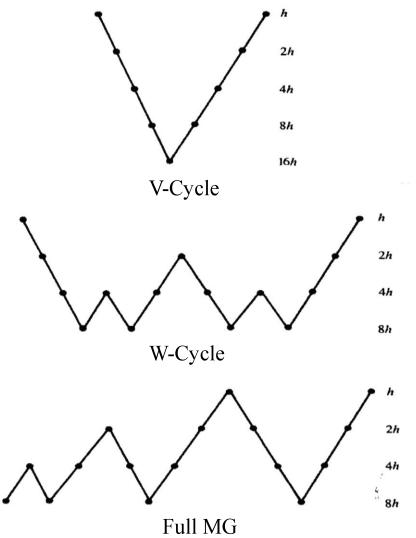
- Key to multigrid: <u>smooth modes look oscillatory on</u> <u>coarser grids</u>; remove them
- Map error from fine grid to coarse grid
- Smooth error is now relatively oscillatory
- Relax on coarse grid to remove this error



Map solution back to fine grid and correct



## Multigrid in Energy



- Multigrid in Energy (MGE) as right preconditioner
- Relaxation method is weighted Richardson

$$x^{i+1} = \omega \mathbf{TMS} x^i + \omega b + (1 - \omega) x^i$$

- Relaxations per level
- Number of levels
- omega



#### **Parallelization of MGE**

With one set, number of grid levels is

$$floor(log_2(G-1)) + 2$$

- Multisets: each set does its own "mini" V-cycle
- No communication beyond upscattering required
- Grid depth becomes

$$num_{-}g_{min} = floor\left(\frac{num_{-}groups}{num_{-}sets}\right)$$

$$num\_grids = floor(log_2(num\_g_{min} - 1)) + 2$$



# **Highlights From the Past**

- RQI doesn't really work without preconditioning
- Parameter selection: start with w1r2v2
- Use reduced angle set in preconditioner
- Shallow grid depth is sufficient
- Preconditioner scales very well in energy
- Preconditioned PI not faster than PI alone for
  - R. N. Slaybaugh, T. M. Evans, G. G. Davidson, and P. P. H. Wilson, "Rayleigh Quotient Iteration in 3D, Deterministic Neutron Transport," *PHYSOR 2012 Advances in Reactor Physics Linking Research, Industry, and Education,* Knoxville, TN, 2012, American Nuclear Society.
  - R. N. Slaybaugh, T. M. Evans, G. G. Davidson, and P. P. H. Wilson, "Multigrid in energy preconditioner for Krylov solvers," *Journal of Computational Physics*, **242**, pp. 405–419 (2013).



#### **Questions**

1. Will preconditioning with MGE facilitate the use of RQI?

2. Will the combination of RQI, MGE, and the block Krylov solver be advantageous for at least some problems?



### **MGE Does Not Help Pl**

- 2D C5G7 Benchmark
- 16 cores, tolerance 10<sup>-3</sup>, k tolerance 10<sup>-5</sup>

Solver	Precond	d Krylov	Eigen	Time(s)
PI	none	3,129	32	$8.54 \times 10^3$
PI	w1.4r2v2	438	31	1.77 x 10 <sup>4</sup>
PI	w1r3v3	253	31	2.28 x 10 <sup>4</sup>
RQI	none	n/a	n/a	n/a
RQI	w1r3v3	299	19	2.57 x 10 <sup>4</sup>



### MGE Helps RQI

- 2D C5G7 Benchmark
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#### **But PI Is Still Better**

- 2D C5G7 Benchmark
- 16 cores, tolerance 10<sup>-3</sup>, k tolerance 10<sup>-5</sup>

Solver	Precond	Krylov	Eigen	Time(s)
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### Problems Are Too Easy; Use Pl

- 3D C5G7 Benchmark, DR 0.7709
- 720 cores, tolerance 10<sup>-4</sup>, k tolerance 10<sup>-5</sup>

Solver	Precond	Krylov	Eigen	Time(s)
PI	none	3,129	32	$4.46 \times 10^3$
RQI	w1.3r2v2	302	19	2.32 x 10 <sup>4</sup>
PI	w1.3r2v2	288	32	2.84 x 10 <sup>4</sup>
RQI*	w1r3v3	103	9	3.02 x 10 <sup>4</sup>
PI*	w1r3v3	126	14	4.04 x 10 <sup>4</sup>
RQI	w1.5r3v3	187	19	3.24 x 10 <sup>4</sup>
PI	w1.5r3v3	192	32	$3.73 \times 10^4$

\*different tolerances



#### **Full PWR-900 Details**

2 x 2 spatial cells/pin

17 x 17 pins/assembly

289 assemblies (132 reflector, 159 fuel of varying enrichment)

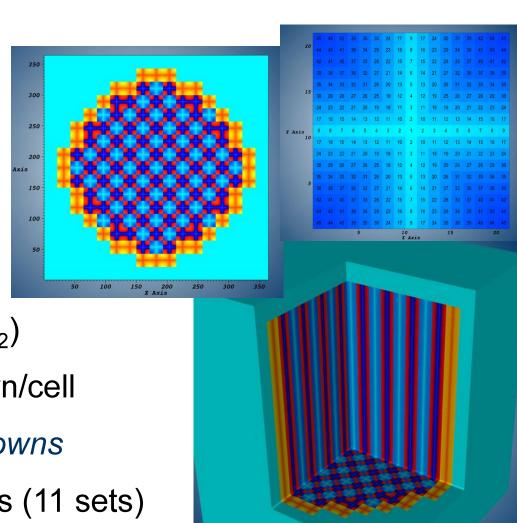
 $P_0$ : 1 moment

 $S_{12}$ : 168 angle sets (MGE:  $S_2$ )

233,858,800 cells; 1 unknown/cell

44 groups: 1.73 trillion unknowns

12,544 blocks; 137,984 cores (11 sets)





#### **RQI** Can Beat PI

RQI needed less time and fewer iterations

Solver	Precond	Krylov	Eigen	Time (m)
PI	none	5.602	149	612.2
PI	w1r2v2	946	86	720*
PI	w1r3v3	111	11	480*,+
RQI	w1r2v2	70	5	54.8
RQI	w1r3v3	76	6	330.4+

\*Exceeded wall time limit
+S<sub>12</sub> in MGE; different tolerances and decomposition

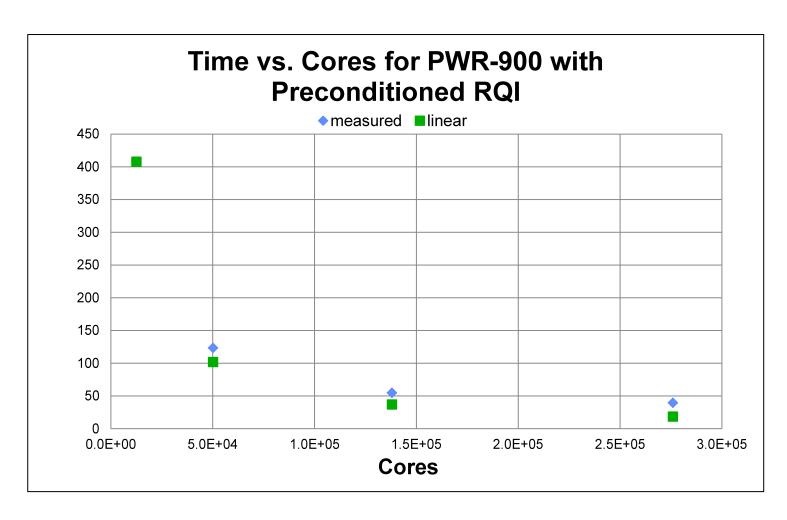


# **RQI+MGE Strong Scaling**

Sets	Domains	Time (m)	Tperfect	Efficiency
1	12,544	407.8	407.8	1.000
4	50,176	123.4	102.0	0.826
11	137,984	54.8	37.1	0.676
22	275,968	39.6	18.5	0.468



# **RQI+MGE Strong Scaling**





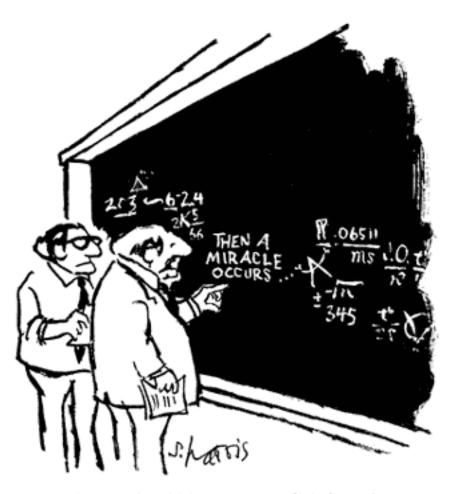
#### **Conclusions**

- RQI is an optimal eigenvalue method
  - Desirable for loosely coupled systems
  - Needs preconditioning for large problems
  - Can be a win for challenging problems

- Is enabled by multigroup Krylov solver
- Multigrid in Energy preconditioner can converge flux so RQI can converge challenging problems
- The methods used in concert performed well on large machines for real problems



## **Questions?**



"I think you should be more explicit here in step two."



#### **Parameter Effects?**

- User input parameters:
  - w = Richardson relaxation weight
  - r = number of relaxations per grid
  - v = number of concatenated V-cycles
- w = 1 always safe; up to 1.3 or 1.4 often beneficial
- Increasing r and v always reduced Krylov count
- r1v2 vs. r2v1
  - same impact on iteration count
  - may result in different times



#### **Small Test Success 1**

3 x 3 x 3, 0.1 grid size, vacuum BCs, 2 materials,
 4 downscattering-only groups, S<sub>2</sub>, P<sub>0</sub>

 $k_{ref} = 0.11752$ , dominance ratio = 1.396e-1

	PI	RQI
k	0.11752	0.11752
group iters	2 or 3 / group	4 to 8
eigen iters	7	6
total Krylov	73	39



#### **Small Test Success 2**

3 x 3 x 3, 0.1 grid size, reflecting BCs, 1 material,
 4 downscattering-only groups, S<sub>2</sub>, P<sub>0</sub>

 $k_{ref}$  = 2, dominance ratio = 1.630e-15

	PI	RQI
k	2	2
group iters	9 / group	17, 18
eigen iters	2	2
total Krylov	72	35



### **Transport Equation**

$$\begin{split} [\hat{\Omega} \cdot \nabla + \Sigma(\vec{r}, E)] \psi(\vec{r}, \hat{\Omega}, E) &= \\ \int dE' \int d\hat{\Omega}' \; \Sigma_s(\vec{r}, E' \to E, \hat{\Omega}' \cdot \hat{\Omega}) \psi(\vec{r}, \hat{\Omega}', E') \\ &+ \frac{\chi(E)}{k} \int dE' \; \nu \Sigma_f(\vec{r}, E') \int d\hat{\Omega}' \; \psi(\vec{r}, \hat{\Omega}', E') \end{split}$$

- $\psi(\vec{r}, \hat{\Omega}, E)$  is the angular neutron flux in neutrons per unit length squared per steradian,
- $\chi(E)$  is the fission spectrum,
- $\bullet$  k is the eigenvalue, which can be thought of as the asymptotic ratio of the number of neutrons in one generation to the number in the next,
- $\Sigma$ s are probabilities of interaction with units of inverse length,
- $\nu$  is the average number of neutrons released per fission.



# Krylov Methods are Powerful

 Solutions to Ax = b formed iteratively from Krylov subspace:

$$\mathcal{K}_k(\mathbf{A}, v_1) \equiv span\{v_1, \mathbf{A}v_1, \mathbf{A}^2v_1, ..., \mathbf{A}^{k-1}v_1\}$$

#### Useful:

- because robust, often converge quickly, easy to precondition, only matrix-vector products
- when A is large, not explicitly formed, sparse
- Drawbacks: can create large subspaces
- Restart methods available



#### **Gauss Seidel for Outer Iterations**

- Scattering to own group, lower group, higher group
- GS does the outer (energy) iteration:
- Inner (space-angle) iteration for 1 to G:

Update right hand side, repeat loop for upscattering groups until convergence



# **Krylov for Inner Iterations**

The Krylov solver applies A to a group-sized iteration vector, v, which represents :

The action of  $\mathbf{A}$  is implemented by doing the following for a group g:

$$(\mathbf{I} - \mathbf{D} \mathbf{L}^{-1} [\mathbf{M}] [\mathbf{S}]_{gg}) [\phi]_g^* =$$

- 1. matrix-vector multiply:  $y_g = [\mathbf{M}][\mathbf{S}]_{gg}v_g$ ,
- 2. sweep:  $z_g = \mathbf{D} \mathbf{L}^{-1} y_g$ ,
- 3. return:  $v_g \leftarrow v_g z_g$ .



### **Block Krylov: One Iteration Level**

Apply A to whole upscatter block at once

$$\mathbf{S} = \begin{pmatrix} [\mathbf{S}]_{11} & 0 & 0 & 0 & 0 \\ [\mathbf{S}]_{21} & [\mathbf{S}]_{22} & 0 & 0 & 0 \\ [\mathbf{S}]_{31} & [\mathbf{S}]_{32} & [\mathbf{S}]_{33} & [\mathbf{S}]_{34} & [\mathbf{S}]_{35} \\ [\mathbf{S}]_{41} & [\mathbf{S}]_{42} & [\mathbf{S}]_{43} & [\mathbf{S}]_{44} & [\mathbf{S}]_{45} \\ [\mathbf{S}]_{51} & [\mathbf{S}]_{52} & [\mathbf{S}]_{53} & [\mathbf{S}]_{54} & [\mathbf{S}]_{55} \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \\ v_3 \\ v_4 \\ v_5 \end{pmatrix}$$
 Upscatter RHS Upscatter Block 
$$(\mathbf{I} - \mathbf{D} \mathbf{L}^{-1} \mathbf{M} \mathbf{S}_{\text{up\_block}}) \phi^* = \mathbf{D} \mathbf{L}^{-1} \mathbf{M} (q_e^{\text{up}})$$

Iteration vector includes all upscattering groups



# **Enables Parallelization in Energy**

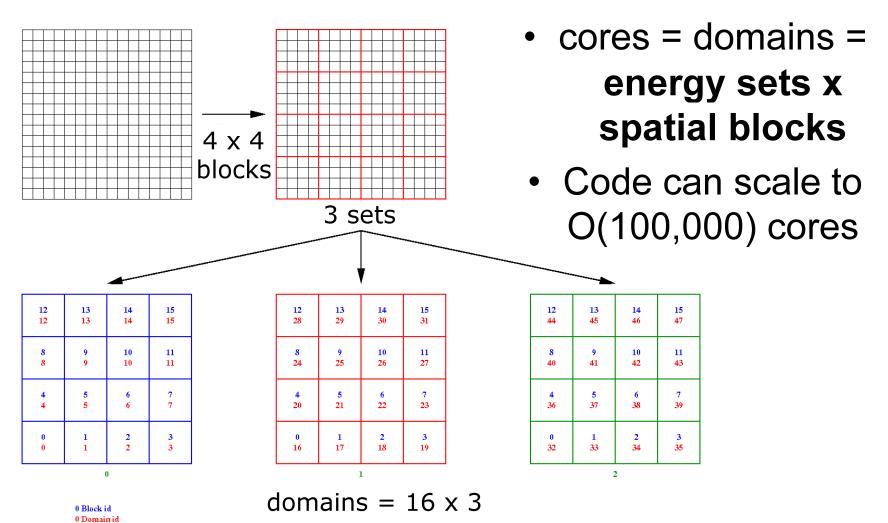
- Matrix-vector multiply for each upscatter group can be done in parallel
- Communicate result after each multiply
- If each upscatter group computed simultaneously:

$$\begin{aligned}
s_g &= \mathbf{S}_{g1} v_1 + \mathbf{S}_{g2} v_{g2} + \dots + \mathbf{S}_{g5} v_5 \\
\begin{pmatrix} s_1 \\ s_2 \\ s_3 \\ s_4 \\ s_5 \end{pmatrix} &= \begin{pmatrix} [\mathbf{S}]_{11} & 0 & 0 & 0 & 0 \\ [\mathbf{S}]_{21} & [\mathbf{S}]_{22} & 0 & 0 & 0 \\ [\mathbf{S}]_{31} & [\mathbf{S}]_{32} & [[\mathbf{S}]_{33}] & [\mathbf{S}]_{34} & [[\mathbf{S}]_{35}] \\
[\mathbf{S}]_{41} & [\mathbf{S}]_{42} & [[\mathbf{S}]_{43}] & [\mathbf{S}]_{44} & [[\mathbf{S}]_{45}] \\
[\mathbf{S}]_{51} & [\mathbf{S}]_{52} & [[\mathbf{S}]_{53}] & [\mathbf{S}]_{54} & [[\mathbf{S}]_{55}] & \\
\end{bmatrix}$$



0 Set id

## **Energy Set Decomposition**





### Large Test Failure

- 2-D C5G7 MOX benchmark problem
- External mesh file, 2 reflecting BCs, 10 materials,
   7 groups w/ 4 upscattering, S<sub>2</sub>, P<sub>0</sub>

 $k_{ref} = 1.18655 \pm 0.008$ , dominance ratio = 0.7709

	PI + TTG GS	PI + MG Krylov	RQI
k	1.18538	1.18702	~0.95
mg iters	157 GS	~100	1000
eigen iters	32	32	120*
total Krylov	21,365	3,129	119,006

<sup>\*</sup>terminated manually



# **RQI** Convergence Issues

The multigroup iterations do not converge

