# **Process Modeling in the Pharmaceutical Industry using the Discrete Element Method**

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Received 26 February 2008; revised 7 May 2008; accepted 7 May 2008

Published online 18 June 2008 in Wiley InterScience (www.interscience.wiley.com). DOI 10.1002/jps.21466

ABSTRACT: The discrete element method (DEM) is widely used to model a range of processes across many industries. This paper reviews current DEM models for several common pharmaceutical processes including material transport and storage, blending, granulation, milling, compression, and film coating. The studies described in this review yielded interesting results that provided insight into the effects of various material properties and operating conditions on pharmaceutical processes. Additionally, some basic elements common to most DEM models are overviewed. A discussion of some common model extensions such as nonspherical particle shapes, noncontact forces, and interstitial fluids is also presented. While these more complex systems have been the focus of many recent studies, considerable work must still be completed to gain a better understanding of how they can affect the processing behavior of bulk solids. © 2008 Wiley-Liss, Inc. and the American Pharmacists Association J Pharm Sci 98:442–470, 2009

**Keywords:** pharmaceutical engineering; powder technology; discrete element method; dynamic simulation; molecular dynamics; blending; granulation; milling; tableting; coating

# **INTRODUCTION**

Processes involving particulate matter are prevalent throughout the pharmaceutical industry. In many cases, however, there are opportunities to improve the fundamental understanding of these processes to directly benefit equipment design, process efficiency, and scale-up. The use of various process modeling tools is becoming increasingly common and is playing a critical role in efforts to gain insight into these processes. Several computational process modeling tools are

used in the pharmaceutical industry and have been reviewed recently.<sup>1,2</sup> These models include computational fluid dynamics (CFD), the finite element method (FEM), and the discrete element method (DEM).

Computational fluid dynamics is an Eulerian method that treats the material as a continuum and numerically solves mass, momentum, and energy balances. This approach can be extremely accurate for systems consisting of fluids, and can be extended to systems consisting of dilute solids in either of two general ways. In the Eulerian–Eulerian approach, the solids are assumed to be a second continuous phase for which the aforementioned balance equations are also solved. In the Eulerian–Lagrangian approach, the solids are modeled discretely using DEM, and these results are coupled with the CFD fluid model. For systems

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Journal of Pharmaceutical Sciences, Vol. 98, 442–470 (2009) © 2008 Wiley-Liss, Inc. and the American Pharmacists Association



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of dense solids however, the material is usually assumed to be elastic or elasto-plastic, and other modeling approaches, such as FEM, are then required.

In contrast to CFD and FEM models, the discrete element method is a Lagrangian model and, as such, tracks the positions, velocities, and accelerations of each particle individually. This approach, reviewed by several authors, 3,4 is advantageous due to the high level of detail in the output describing the dynamic behavior of the particles. The particle positions and velocities can be used to calculate many particle-scale quantities of interest such as local concentrations and particle phase stresses, as well as examine particle-level phenomena such as segregation or agglomeration. While this approach does demand significant computational power, with the steadily increasing speed of computer hardware, the size of systems that can be modeled with DEM is continually increasing. Some recent DEM models simulate systems on the order of a couple hundred thousand particles.

The DEM approach is particularly beneficially since predictions can readily be made from simulation data that would be difficult, and possibly expensive, to obtain experimentally. Another strong point of DEM modeling is the ease at which parametric studies can be executed to study the effects of particle properties, process conditions, or equipment design. Further, since it is a discrete approach, particles can have a distribution of properties. Thus, the particles with in the system can have a distribution of sizes, for instance, or can even have varying material properties to model a mixture of various components. These details can generally not be addressed in continuum approaches.

This paper seeks to review the current status of DEM modeling of pharmaceutical processes. To this end, some of the fundamental aspects of DEM models are first reviewed, including the hard- and soft-particle approaches. Next, some common extensions of basic DEM models, such as the incorporation of cohesive forces, nonspherical particle shapes, and interstitial fluid effects are highlighted. This is followed by a review of recent DEM modeling work that focuses on some common pharmaceutical processes.

### The Discrete Element Method (DEM)

There are four main classes of discrete element models, each with their own relative advantages and state of development. These consist of cellular automata, Monte Carlo methods, hard-particle methods, and soft-particle methods. Cellular automata are simple, deterministic models that have been used to study flow in silos,<sup>5</sup> down inclined chutes,<sup>6</sup> and in rotating drums.<sup>7</sup> They have also shed light on phenomena such as segregation via percolation<sup>8</sup> and vibration,<sup>9</sup> as well as the effects of nonspherical particles.<sup>10</sup> In this approach, particles are constrained to a lattice, and their movement is governed by simple rules established through experimental observation. While these models are extremely fast, they generally only provide qualitative predictions.

The Monte Carlo method is another approach that has been applied to model granular materials. This approach calculates the probability of a random arrangement of particles based on the energy associated with that arrangement. Like the cellular automata, Monte Carlo simulations are rather simple to implement and fairly computationally efficient. Monte Carlo techniques have been used to study several aspects of granular behavior, including vibration-induced segregation, 11,12 the effect on nonspherical particle shapes on hopper flow, <sup>13</sup> and the distribution of coating in film coaters. <sup>14,15</sup> However, the hardand soft-particle approaches, as discussed in the next sections, are much more commonly used. These approaches are generally more flexible and typically permit a direct relation between material properties and model parameters, thereby increasing the likelihood of obtaining results that are in quantitative agreement with experiments.

### Hard-Particle Approach

The hard-particle approach 16-20 assumes that particles are rigid so that collisions are instantaneous and binary. As a result, hard-particle models are generally best suited for dilute, collisional flows where these assumptions are good approximations. Hard-particle models often are embedded in event-driven collision detection schemes that increment the simulation time from one collision to the next. Consequently, hard-particle, event-driven simulations can be computationally efficient when the time between particle collisions is large.

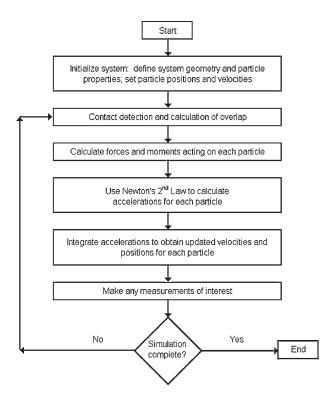
A hybrid, hard-particle-with-overlap model has also been developed by Hopkins. <sup>21</sup> This hybrid model uses the hard-particle collision model described above, but integrates the particle equations of motion between collisions at a specified time

step. Thus, this is referred to as a time-driven approach. This time-driven approach can be computationally more efficient than the event-driven approach for systems containing many particles.

## Soft-Particle Approach

In dense systems, the particle contacts are multiple and enduring. Therefore, the hardparticle interaction is not appropriate and the soft-particle model must be used. The soft-particle approach, originally developed by Cundall and Strack,<sup>22</sup> is not limited by the instantaneous contact time assumption of the hard-particle model and can therefore be used to investigate long-lasting and multiple particle contacts. This model allows for particle deformation which is modeled as an overlap of the particles. To prevent excessive errors from being introduced, the mean overlap value must be maintained to levels on the order of 0.1–1.0% of the particle size.<sup>23</sup> This is typically accomplished through appropriate selection of the spring stiffness. Using stiffness values that are too small will result in large overlap values, and will potentially lead to the introduction of significant error through inaccurate determination of contact forces and thus postcontact quantities such as accelerations and velocities. 24 The soft-particle model relies on a force-displacement (and/or force-displacement rate) relation to determine the interaction forces for each particle-particle and particle-wall contact. This approach proceeds via small time steps and is thus referred to as being time-driven. Accurate integration of the resulting particle equations of motion dictates a small simulation time step and, hence, long computation times. Soft-particle methods are generally utilized for dense, enduring-contact flows, such as those occurring in a hopper, blender, or pan coater.

The DEM soft-particle algorithm is relatively straightforward and is shown in a flowchart in Figure 1. At the onset of a simulation, the particle properties such as particle size distribution (PSD), density, mass, and moment of inertia are all defined. If process equipment is to be modeled, the geometry and dimensions are also defined. The particles are then inserted into the computational domain by defining a position and velocity for each, often by placing the particles on a lattice or by using a random number generator. The simulation then proceeds to a contact detection stage where all particle—particle and particle—



**Figure 1.** A flowchart of a general soft-particle DEM algorithm.

wall contacts are identified. For each contact, the soft-particle deformation, which is modeled as an overlap, is calculated. Using the overlap values, force-displacement relations are used to calculate the forces acting on each particle. The total forces and moments acting on each particle are then summed and Newton's second law is used to calculate the translational and rotational accelerations. The time step is incremented and accelerations are integrated over time to determine updated particle velocities and positions. After measurement of any quantities of interest, such as velocity profiles, solid phase stresses, or local concentrations, the simulation repeats the contact detection for the updated particle positions and the loop is repeated.

Contact models. At the heart of a DEM simulation is a particle interaction model. This particle interaction model is used to calculate the forces acting in either particle—particle or particle—wall contacts. Each of these two contact types can be resolved using the same contact model, and the material properties (e.g. coefficient of restitution, friction coefficient, etc.) for each contact type can differ so that dissimilar materials can be modeled. These forces are used to calculate accelerations,

which are then integrated in time for updated particle velocities and positions. The modeled forces in a soft-particle system are typically subdivided into normal and tangential components. A number of different normal force—displacement models have been proposed and these have been reviewed by several authors. <sup>3,25–27</sup>

One of the simplest and most commonly used normal force model is the linear spring and dashpot (LSD) model. <sup>22,23,28–33</sup> Here, the repulsive portion of the contact is modeled with a linear spring and the dissipative portion of the contact is modeled in parallel with a viscous dashpot. Together, these elements represent the normal portion of a contact fairly well. This commonly used model is advantageous in that the coefficient of restitution, contact time, and maximum overlap can be analytically determined. However, the resulting coefficient of restitution is not dependent on the impact velocity.

Others have modified the repulsive portion of the LSD model with a nonlinear spring following the Hertz theory of elastic contacts. <sup>34,35</sup> This Hertzian model is usually coupled with a dashpot to provide dissipation. <sup>36–38</sup> However, using this model yields a coefficient of restitution that increases with impact velocity, which is contrary to experimental data. <sup>39</sup> Thus, later work by Kuwabara and Kono<sup>40</sup> and Brilliantov et al. <sup>41</sup> assumed the material to be viscoelastic instead of elastic. With this modified dissipation term, the coefficient of restitution decreases with increasing impact velocity, in accordance with experimental results. <sup>40</sup> However, this Hertzian implementation

has no intrinsic time scale for non-zero damping coefficients.<sup>25,26</sup> Therefore, the simulation time step must be selected by using an estimate of the maximum expected impact velocity.

Another approach was proposed by Walton and Braun<sup>42</sup> who assumed that materials plastically deform during contact. They developed a partially-latching, hysteretic spring model shown schematically in Figure 2 that uses one linear spring to model the repulsive force during loading, and another, stiffer linear spring to model the force during the unloading portion of the contact. This model has been used by many researchers<sup>24,43–45</sup> due to several advantages. First, no viscous damping coefficient is required. This is beneficial as the damping coefficients can be difficult to relate to particle properties. Second, either a constant or impact velocity dependent coefficient of restitution can be specified.

A number of tangential force—displacement models of varying complexity have also been proposed and reviewed by several authors. <sup>3,25,26</sup> One of the simplest models was first introduced by Cundall and Strack<sup>22</sup> and subsequently used by several others. <sup>46–48</sup> Here, the tangential force is modeled with a linear spring where the associated displacement is integrated from the relative velocity at the contact point. This tangential force determined by the spring is limited by the Coulomb criterion, which specifies that the maximum frictional force is proportional to the normal force, where the coefficient of friction is the proportionality constant. This tangential force model is an attractive choice because it is relatively

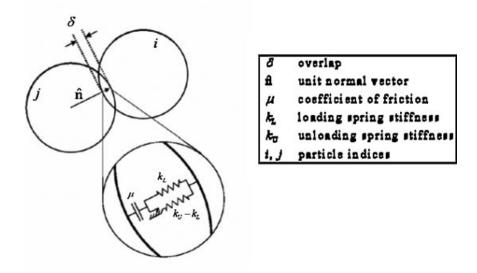


Figure 2. A schematic of the Walton and Braun<sup>42</sup> hysteretic spring model.

easy to implement and reproduces experimental observations fairly well.<sup>25</sup> It does however require a tangential displacement value to be stored in memory for each contact.

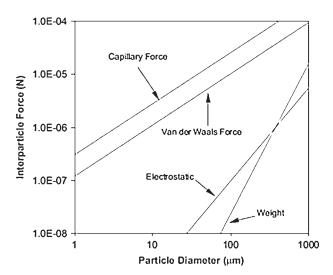
Other tangential force models have been based on the detailed analysis of frictional, elastic spheres by Mindlin and Deresiewicz<sup>49</sup> which specifies that the tangential force is dependent on the loading and unloading history of the contact. The full model considers eleven different loading/unloading possibilities, but generally, only two are considered in DEM implementations. The implementation of this model is significantly more complex than the Cundall and Strack model discussed previously, but compares only slightly better with experimental results, as reported by Schäfer et al.<sup>25</sup> Despite the additional complexity of this tangential model, it has been used frequently by several authors.<sup>42–45,50</sup>

### **Model Extension**

Historically, discrete element models have considered only cohesionless particles of circular (in 2D) or spherical (in 3D) shape, while neglecting the effects of interstitial fluids. These simplifications have been made primarily due to the complexities of contact detection and the extensive computational power required to model more complex systems. However, as the speed of computers has increased, the scope of systems than can be modeled in a reasonable amount of time has similarly increased. Today, it is not uncommon for researchers to use DEM to model systems with secondary forces such as cohesion, nonspherical particle shapes, or interstitial fluid effects. Each of these extensions will be discussed in the following paragraphs.

### Noncontact Forces

Noncontact forces due to capillary cohesion, electrostatics, or van der Waals interactions may significantly affect the behavior of bulk solids, especially for systems with small particle sizes or where moisture is present. Figure 3 shows the magnitudes of these forces and how they become quite significant relative to particle weight as the particle diameter decreases. The importance of these forces has been reviewed by Seville et al. <sup>51</sup> and the application of such forces in DEM models has been reviewed by Zhu et al. <sup>4</sup> Cohesive forces strongly affect the flowability of powders <sup>52</sup> and can lead to ratholing and



**Figure 3.** A comparison of the magnitude of interparticle forces for varying particle diameters. This figure is reprinted from Zhu et al.<sup>4</sup> with permission from Elsevier.

arching during hopper discharge, blockage of pneumatic conveying lines, as well as agglomeration and caking of powders. Additionally, cohesive forces can significantly affect fluidization processes by causing agglomeration and subsequent defluidization. While cohesive forces are significant in most powders, many DEM models have not included cohesive forces. Thus, further study of cohesive systems is clearly needed to gain a better understanding about these effects.

Some of the simplest methods to model cohesion use a proportionality constant to predict the attractive force between particles. These methods predict the cohesive force by a spring constant, <sup>56,57</sup> a constant proportional to particle weight, <sup>58–60</sup> a square-well potential, <sup>53,61–63</sup> or by assuming the cohesive force is proportional to contact area. <sup>64</sup> Each approach is relatively simple to implement, but uses constants that may not be directly related to experimentally measurable values.

A more rigorous treatment models cohesion with pendular liquid bridges. The capillary force of a liquid bridge is specified by the Young–Laplace equation; however, this equation generally cannot be solved analytically. Thus, Fisher<sup>65</sup> and, more recently Lian et al.,<sup>66</sup> assumed a torodial liquid bridge shape, which allows for an estimation of the capillary force. Lian et al.<sup>66</sup> showed that the error associated with this torodial assumption is less than 10%. While the expressions based on the torodial assumption permit an estimation of the capillary force, they are not

explicit and so an iterative scheme must be used. This approach has been used in several DEM models to study the fundamental behaviors of cohesive materials in shear flow,52 during packing,<sup>67</sup> and while forming repose angles,<sup>55</sup> as well as in more complex systems such as a vibrated bed,<sup>68</sup> rotary drum,<sup>69</sup> and high-shear granulator.<sup>70</sup> Mikami et al.<sup>54</sup> developed regression expressions that provided an explicit calculation of the force as a function of the separation distance, and therefore provided a more amenable method for incorporating cohesive forces into DEM models. These expressions were later extended to the polydisperse case.<sup>71</sup> Due to the ease of implementation, these regression equations have been used to incorporate cohesive forces in studies of uniaxial compression, 71 shear flow, 72 fluidization, 54,73 and hopper discharge. 74

Other attractive forces such as van der Waals forces and surface adhesion have also been incorporated in to DEM models. van der Waals attractive forces have been added to study packing and rearrangement of fine particles<sup>75</sup> as well as fluidization. <sup>76–78</sup> Thornton and Yin developed a theory for contacts in the presence of adhesion, <sup>79</sup> and later applied this model in the study of agglomerate fracture.<sup>80</sup> A few researchers have recently applied discrete element models to colloidal systems as well. For these systems, a DLVO potential, based on the theory of Derjaguin, Landau, Verwey, and Overbeek<sup>81,82</sup> is often employed. This DLVO potential is based on the cumulative effects of electrostatic repulsion and van der Waals attractive forces. Additionally, an attractive force, based on the theory of Johnson-Kendall–Roberts (JKR), 83 is also included in some models for colloidal systems.84-86

# Particle Shape

Many studies of granular materials using DEM have modeled particles with spheres in 3D or disks (circular shapes) in 2D. The use of these simple shapes allows for straightforward contact detection and contact resolution. However, these sphere/disk systems are often too idealized to accurately model some phenomena exhibited by real granular materials. For instance, it has been shown that spherical particles have a smaller angle of repose <sup>87,88</sup> and a reduced strength <sup>89,90</sup> as compared to nonspherical particles. The underlying cause is that the rotation of spheres is only resisted by frictional contacts with neighboring particles whereas for nonspherical particles,

rotation tends to be inhibited by mechanical interlocking.  $^{91-93}$ 

Several different methods have been used in an attempt to reduce the rotational freedom of spheres/disks within DEM models. Ting and Corkum<sup>94</sup> artificially increased the particle moment of inertia thereby reducing particle rotation. Others have damped<sup>95</sup> or completely fixed<sup>96–98</sup> particle rotation. Another approach incorporates a rolling resistance.<sup>38,99–101</sup> This approach shows promise however, it also introduces additional parameters to the model via rolling resistance coefficients.

Due to the aforementioned differences between experimental particles and DEM models of spheres, a common extension of the basic, noncohesive, sphere-based DEM models is to incorporate nonspherical particle shapes. For systems of spheres/disks, contact detection and contact resolution are relatively straightforward. However, for nonspherical particles, these calculations can become quite complex. A wide range of nonspherical particles have been implemented in DEM models, and many of these have been reviewed by Džiugys and Peters<sup>102</sup> and Hogue. <sup>103</sup>

One of the simplest approaches to model nonspherical particles is by "gluing" two or more spheres or disks together in either a rigid or flexible manner. This clustering approach has been commonly used in 2D with disks<sup>87,88,104,105</sup> and in 3D with spheres. 50,106–113 This approach retains the relative simplicity of contact detection and contact calculation associated with spheres or disks. Additionally, arbitrary or nonsymmetric particle shapes can be modeled. However, a drawback is the difficulty in modeling shapes that have sharp edges. 114,115 Also, a potential disadvantage is the roughness of the resulting particles if they consist of only a small number of spheres/disks. It has been shown that the dynamic behavior of a sphere-cluster particle may be significantly different than that of the smooth particle it is intended to represent. 116

A similar method uses the idea of mathematical dilation<sup>117</sup> to create spherocylinders and spherodisks. <sup>13,118–123</sup> These so-called dilated shapes are created by placing a sphere with a fixed diameter at every point on some basic geometric shape. A spherocylinder is created by dilating a line segment with a sphere while a spherodisk (a 3D disk) is created by dilating a 2D disk with a sphere. Thus, each particle is modeled by an infinite collection of spheres, resulting in a smooth particle surface. However, the contact detection

for dilated shapes becomes somewhat more complicated than that for spherical particles.

Another class of shapes is those that can be mathematically defined. Most commonly, these include ellipses in  $2D^{89,90,102,124-127}$  and ellipsoids in  $3D^{92,128-130}$  as defined by

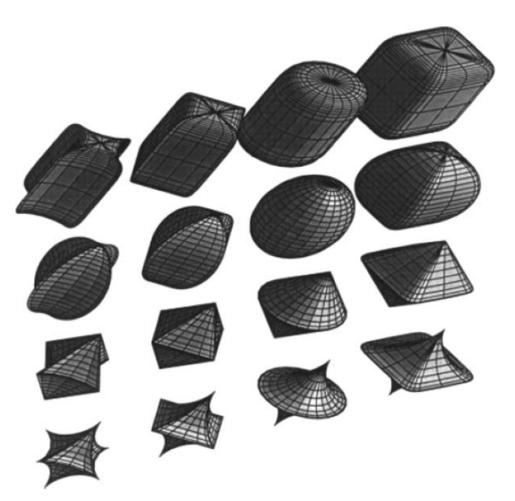
$$\left(\frac{x}{a}\right)^2 + \left(\frac{y}{b}\right)^2 + \left(\frac{z}{c}\right)^2 - 1 = 0 \tag{1}$$

where  $\alpha$ , b, and c define the lengths of the three principle axes which are aligned with the local x, y, and c coordinates, respectively. Superquadrics,  $c^{131,132}$  the general class of shapes to which ellipses belong, have also been used in  $c^{23,130,133,134}$  and  $c^{33}$  to represent nonspherical particles. The shapes of several superquadrics, shown in Figure 4, are mathematically given by

$$\left|\frac{x}{a}\right|^{\alpha} + \left|\frac{y}{b}\right|^{\alpha} + \left|\frac{z}{c}\right|^{\alpha} - 1 = 0 \tag{2}$$

where the exponent  $\alpha>0$  characterizes the blockiness of the particle. For  $\alpha=2$ , the equation of an ellipsoid is recovered. For increasing values of  $\alpha$ , the corners begin to sharpen and the particle becomes increasingly blocky. Superquadrics are advantageous since they can algebraically model a wide variety of particle shapes with various values of the exponent  $\alpha$  and the aspect ratios as determined by a, b, and c. Additionally, the normal vector is algebraically defined. However, for these shapes, contact detection is usually prohibitively slow as the intersection of two nonlinear functions must be calculated.

Polygonal particles are another class of particle shapes popular in the field of geomechanics. Arbitrary or symmetric polygons have been modeled in  $2D^{48,115,134,136-140}$  and in  $3D.^{141-145}$  Additionally, polygons constructed of beams and triangular elements have also modeled in  $2D.^{114,146}$  Polygonal particles pose certain



**Figure 4.** Images showing some possible particle shapes using 3D superquadrics. This figure reprinted from Hogue $^{103}$  with permission from Emerald Publishing Group.

difficulties surrounding contact resolution involving corners and edges. Additionally, modeling round or smooth particles with this method is inefficient due to the large number of facets involved.

Potapov and Campbell<sup>147</sup> developed another shape representation for 2D systems. They connected four circular arcs in a continuous manner to create an oval shape. This representation allows for nonround particles to be modeled with only a small increase in the computational effort. Additionally, this approach of connected arcs can be extended to model any other 2D polygon (e.g. triangle, square, pentagon, *etc.*). A drawback of this method is that there is no simple implementation of this general scheme in 3D. One exception is the case of ovals, which have been extended to 3D ovoids<sup>148–150</sup> by rotating the oval around an axis of symmetry.

Another approach is the discrete function representation for 2D shapes, as defined by Williams and O'Connor. This method discretizes the surface of an arbitrarily shaped particle based on a single parameter resulting in a list of vertices that can easily be searched for contacts. One example is given by Hogue and Newland who discretized particles via a polar representation. Here, the surface of a given particle is specified by a polar angle and radius for several vertices on the surface. This approach was also used by Wu and Cocks in die filling investigations. An extension to 3D has also been proposed where two parameters are used to

discretize the particle surface. This discrete representation approach can represent a wide variety of arbitrary shapes that can be either convex or concave, smooth or angular, and elongated or compact.

A final particle shape worth mentioning is that of a round, bi-convex tablet. <sup>116,153</sup> This shape is defined mathematically by the intersection of one small sphere that defines the tablet band and two larger spheres that define each of the convex caps. This shape is advantageous in that it precisely models the shape of interest, and it does so in a computationally efficient manner.

While each of these nonspherical shape representations has its relative merits, it is important to consider the increase in computational time associated with each method. Some authors have reported the relative differences in run time for various shape representations 116,147 and several observed computational times are summarized in Table 1. While these results are not comprehensive, they show that sphere clusters, continuously joined arcs, and bi-convex tablet shapes all require only modest additional computational effort beyond that of spherical particles. In contrast, the computational cost for polygons and superquadrics may be on the order of two to ten times that for a system of spheres. The computational times presented in Table 1 are stated relative to the case of disks/spheres in an effort to eliminate any differences between number of particles, the length of time simulated, and the computational

**Table 1.** Computational Times for Various Particle Shape Representations as a Multiple of the Time Required for Disk-Shaped or Spherical Particles

Dimensionality	Shape	Increase in Computational Time Over that for Disks	References
2D	Quasi-triangle	1.7	147
2D	Quasi-square	1.8	147
2D	Superquadric	3	134
2D	Square	3.9	147
2D	Superquadric	12	147
Dimensionality	Shape	Increase in Computational Time Over that for Spheres	References
3D	10 sphere cluster	1.1	116
3D	26 sphere cluster	1.2	116
3D	Biconvex tablet	1.5	116
3D	Ovoid	2	149
3D	Superquadric	${\sim}23$	33
3D	66 sphere cluster	2.2	116
3D	178 sphere cluster	6.3	116

resources utilized. However, differences in the contact algorithm, the degree of optimization, or other factors may still result in widely varying computational times (e.g., the differences for 2D superquadrics). <sup>134,147</sup>

### Fluid Effects

The effects of interstitial fluids are neglected in many DEM models. For many systems, such as those with relatively large particles, typically greater than 500  $\mu$ m, this assumption is usually satisfactory. However, for systems with fine particles, it may be critical to include interstitial fluid effects to obtain accurate model predictions. Additionally, for processes such as fluid bed drying and pneumatic transport, the fluid phase must be modeled.

The inclusion of such effects is currently an emerging area of study where an Eulerian CFD model is coupled with a Lagrangian DEM model. Thus, the fluid phase is treated as a continuum, while the particulates are modeled as discrete elements. Typically, these studies are carried out with a CFD model coupled with a DEM model using either the hard-<sup>155</sup> or soft-particle<sup>156</sup> approach. Further details of these coupled models are described in several excellent reviews.<sup>4,157</sup>

Several pharmaceutical processes have been studied using these coupled models. For example, air effects have been included in the study of fluidized beds, <sup>155,156,158–162</sup> spouted beds, <sup>163,164</sup> and the effects of cohesion therein. <sup>54,58,165</sup> The presence of air has also been shown to affect the die filling process. <sup>140,166</sup> Finally, the incorporation of air effects is required in pneumatic conveying models. <sup>167–169</sup>

# Validation of DEM Models

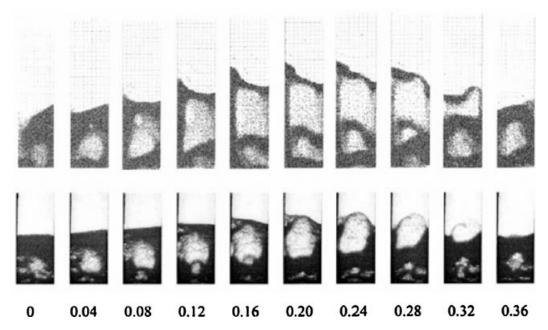
The verification and validation of computational models are of utmost importance. The term verification is used to describe the process of ensuring that the numerical solutions are in fact an accurate solution to the model. Thus, this issue addresses the computational implementation of the mathematical equations. The term validation describes the process by which the accuracy of a model's predictions are compared with experimental data. This area has been widely discussed in the CFD literature via, for example, textbooks, <sup>170</sup> policy statements, <sup>171,172</sup> and review articles. <sup>173</sup> However, the verification and validation of DEM models does not appear to be as

rigorous and the procedures certainly are not as well defined. This is attributable primarily to the difficulty in obtaining experimental data with which to compare. The extent and methods of verification and validation used by researchers vary widely, and often these are not fully described in the literature. A notable exception is the work of Asmar et al. <sup>56</sup> that outlined several fundamental, two-body contacts to verify their DEM model. It was also suggested that these results be used as a benchmark for verification by other researchers.

Discussions of validation studies have appeared considerably more frequently in the literature than those regarding verification. Most validation studies for DEM models have been qualitative and rely on visualization of experimental flows on the surface or at a transparent wall. For example, experimental observations have been used to qualitatively validate flow in a rotating cylinder, <sup>174–176</sup> fluidized bed as shown in Figure 5, <sup>162</sup> tumbling mill, <sup>177</sup> vibrated bed, <sup>178</sup> and during heap formation. <sup>179</sup> However, the opacity of bulk solids limits the applicability of these visualization techniques to quasi-2D systems (e.g. a thin slice of a cylindrical rotating drum). Thus, validation studies of 3D systems that rely on visualization of the material on a free surface or at a transparent wall in order to assess particle phenomena, such as the flow pattern, occurring within the interior may be tenuous.

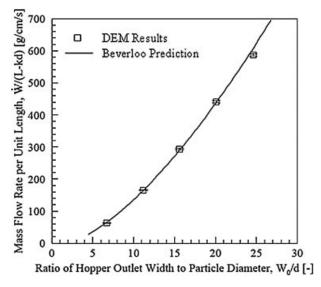
One approach to validation for 3D systems relies on making measurements of macroscopic quantities. For example, the measurement of hopper discharge rates 180,181 can be compared to the well established Beverloo correlation. 182 Figure 6 shows one such comparison with excellent agreement between the DEM and Beverloo predictions. This hopper discharge rate is an example of a simple benchmark test for model validation. Researchers have also measured other macroscopic quantities for validation including the extent of segregation of bidisperse material during hopper discharge 183 and the power draw in a tumbling mill. 177

Tomographic techniques have also been developed towards validation of 3D granular systems. These techniques are nonintrusive and are not hindered by the opacity of solids. Therefore, they are ideal methods to probe the internal microstructure and particle velocities within 3D systems. Nuclear magnetic resonance (NMR)<sup>184</sup> has been used for validation of a long rotating cylinder.<sup>109,185</sup> Also, the blend uniformity in a



**Figure 5.** Comparison of fluidized bed simulations (upper) and experiments (lower) at the given times in seconds. This figure is reprinted from van Wachem et al. <sup>162</sup> with permission from Elsevier.

hopper has been quantified using gamma-ray tomography<sup>186</sup> and the packing of particulates has been examined using X-ray microtomography.<sup>187,188</sup> The positron emission particle tracking



**Figure 6.** DEM predictions of the mass discharge rate per unit length of a monodisperse material in a flat-bottomed, wedge-shaped hopper as a function of the ratio of outlet width to particle size  $W_0/d$ .<sup>181</sup> The line indicates the discharge rate as predicted by the Beverloo correlation<sup>182</sup> modified for wedge-shaped hoppers.<sup>298</sup>

(PEPT) technique<sup>189</sup> has also been used to probe various equipment including high-shear mixers, <sup>190,191</sup> a V-blender, <sup>192</sup> a ploughshare mixer, <sup>193</sup> and fluidized beds. <sup>194,195</sup> This method tracks a single radioactively labeled particle for times on the order of several minutes to an hour to compile the internal flow velocities, which can then be compared with DEM predictions.

### **Limitations of DEM Models**

The discrete element method is an emerging area of computational modeling. The primary limitation of the technique is its inherent computational intensity. Since the DEM approach tracks each individual particle, any increase in the number of modeled particles, N, increases the computational time, which typically scales as  $N \log(N)$ . This effectively limits the number of particles that can be modeled. The number of particles in a system of volume V and relative density  $\phi$  is given by

$$N = \frac{6V\phi}{\pi d^3} \sim \frac{V}{d^3} \tag{3}$$

where d is the particle diameter. Thus, as the volume of a system increases or the particle diameter decreases, the number of particles increases significantly. This is illustrated in

**Table 2.** The Approximate Number of Particles in a 1 L Volume for the Given Particle Diameter and Assuming a Relative Density of 0.625

$d$ ( $\mu$ m)	N
10	$10^{12} \\ 10^{11} \\ 10^{10}$
20	$10^{11}$
50	$10^{10}$
100	$10^9$
200	$10^{8}$
500	$10^7$
1000	$10^{6}$
2000	$10^5$
5000	$10^{4}$

The dark shading indicates system sizes that are typically inaccessible with current DEM models and computational power. The unshaded areas are accessible, and the light shading indicates an intermediate regime.

Table 2 where the approximate number of particles in a 1 L volume is tabulated for the given particle diameter assuming a relative density of 0.625. Typical DEM models at present are capable of modeling on the order of  $10^5$  particles. Thus, to model a 1 L system on a one-to-one basis, the particle diameter could be no smaller than approximately 1-2 mm.

While the above example demonstrates that only select systems with a small number of particles can be modeled on a one-to-one basis, there is utility in modeling larger systems. For these systems with large numbers of particle required, DEM approaches can still be quite useful in elucidating qualitative trends and learning more about particulate behavior and how it is affected by certain particle properties, process parameters, and equipment design.

For those systems involving a large number of particles, several different approaches have been taken to ease the required computational effort. One frequent approach takes advantage of symmetry within a system of interest thereby allowing for the use of periodic boundary conditions. 24,70,101,196 The use of such boundaries reduces the size of the computational domain and thus, the number of modeled particles, without affecting the measured results. A second approach artificially increases the size of the modeled particles. For example, 10 µm particles could be approximated with 10 mm particles in a DEM model which would reduce the total number of particles within the system for a given solids volume. However, caution must be exercised using this approach as particle-level phenomena may differ between the two scales. Further, as in

the example cited above, effects such as cohesion are usually quite significant for 10  $\mu$ m particles but typically insignificant for 10 mm particles. Finally, a third approach partitions the computational domain into smaller regions, but still models the flow in each region. The flow in each region can then be modeled with DEM sequentially on a single computer<sup>197</sup> or distributed to several computers (i.e. parallelized). <sup>23</sup> In contrast, others have only used DEM to model certain regions of interest and have used other approaches such as FEM<sup>198</sup> or stochastic methods<sup>199</sup> to model the flow in the bulk of the domain.

In addition to the significant effect that the number of particles has on the computational time, certain complexities added to the model can also increase the computational time considerably. Such complexities include nonspherical particle shapes, as discussed earlier and summarized in Table 1. Additionally, the inclusion of secondary forces, such as cohesion, or the coupling of the DEM model with CFD to capture interstitial fluid effects will also necessitate a longer computational time.

There are also some other technical concerns regarding DEM models. One common concern is the difficulty associated with sufficiently validating the models. An advantage of DEM is that it can probe complex systems where experimental measurements are difficult or expensive to obtain. This becomes a detriment when it comes to validating the model experimentally, as the necessary measurements may not be available. However, as discussed in the preceding section, researchers are continually devising new technigues and using new experimental tools to obtain measurements for validation. Another common concern is the use of spherical shapes to represent nonspherical granular materials. While this is a typical assumption used to reduce the computational time requirements for a given simulation, results from experiments and DEM simulations have shown that nonspherical particles can significantly affect particulate flow behavior. This is also an ongoing area of research with a focus on developing fast contact detection algorithms for nonspherical particle shapes.

# APPLICATION TO PHARMACEUTICAL PROCESSES

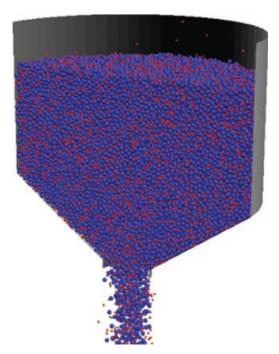
A wide variety of particle handling and processing equipment is typically used in the pharmaceutical industry. The discrete element method has been used with good success to model many of these processes resulting in a better understanding of these systems. In the following sections, highlights of that body of work are presented in the areas of transport and storage, blending, granulation, milling, compression, and film coating.

# **Transport and Storage**

The transport and storage of bulk solids has been comprehensively studied using DEM. Much of the research has focused on relatively simple geometries to permit examination of the underlying granular physics. For instance, numerous studies have examined the solid phase stresses in simple shear flow with periodic boundaries. <sup>24,32,42</sup> Additionally, others have studied the stresses and solids fractions within flows down an incline, <sup>38,50,200</sup> and during the formation of a heap. <sup>48,201,202</sup> In addition to these relatively simple geometries, more complex systems such as hoppers, pneumatic conveying lines, and screw conveyors have also been studied using DEM.

The flow from a hopper is frequently studied due to the prevalence of hoppers in all solids handling industries. Bulk quantities such as the mass discharge rate<sup>47,203-206</sup> are often measured and compare well with empirical correlations such as that proposed by Beverloo. 182 Cleary and Sawley extended this work to show how nonspherical particles affect the hopper flow rate. 23 Other work, shown in Figure 7, reported how bidisperse materials segregate during discharge due to differences in particle size. 101,183 Due to the discrete nature of DEM models, the positions and velocities of each particle are known at every time step. This information has been used to examine features within the granular bed that may be difficult to obtain experimentally. For example, the granular microstructure  $^{186,206}$  and the internal stresses and/or flow fields  $^{31,46,207-209}$ have been obtained for various hopper flow systems. Finally, other work has examined unusual hopper geometries, <sup>44</sup> the stresses acting on the hopper walls, <sup>47,203,204,208</sup> and the effects of hopper vibration, <sup>210,211</sup> flow promoting inserts, <sup>197</sup> and interstitial air on hopper discharge. 154

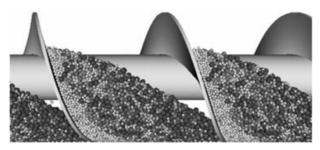
Discrete element models have also been coupled with computational fluid dynamics (CFD) to study the effects of interstitial fluids. One system for which these coupled models are often applied is pneumatic conveying. In some of the earliest



**Figure 7.** Snapshot of a hopper discharge simulation examining the extent of segregation of bidisperse material during discharge. <sup>183</sup>

work, Tsuji et al. <sup>167</sup> developed a model for the pneumatic transport of cohesionless material in a horizontal pipe. Later, other work included extensions for heat transfer, <sup>212</sup> attrition, <sup>213</sup> and electrostatics. <sup>214,215</sup> Other researchers have studied the saltation of solids that occurs in low-velocity flows when the particles fall out of suspension, <sup>216</sup> and developed phase diagrams to describe the various flow regimes observed in vertical and horizontal pipes. <sup>168</sup> Further work has examined flow regimes in inclined pipes, and found good agreement with results from an experimental system. <sup>169</sup>

Another solids transport system modeled via DEM is the screw conveyor. Such models have studied the conveyance of spherical<sup>33,217</sup> and nonspherical particles.<sup>135</sup> Results have shown good agreement with established empirical correlations for transfer velocity and power.<sup>217</sup> Figure 8 shows an image from a screw conveyor simulation where segregation of polydisperse material is observed to occur after only two revolutions of the screw. Other related work has examined solids feeding and conveying within a single-screw extruder<sup>218</sup> and mixing within a twin-screw extruder.<sup>219</sup>



**Figure 8.** A snapshot from a simulation of a screw conveyor where the particles are shaded according to their diameter. Size segregation is apparent here after only two screw revolutions. This figure is reprinted from Cleary<sup>33</sup> with permission from Emerald Group Publishing.

# **Blending**

Blending operations are one of the most common applications of DEM models and much of the recent work is highlighted in a review by Bertrand et al. 180 A number of different mixing systems have been modeled including V-blenders, tote blenders, and bladed mixers. However, one of the most frequently studied mixing systems is that of a rotating drum. For example, studies of rotating cylindrical drums, reviewed by McCarthy et al., 220 typically focus on mixing and segregation in the radial 13,69,196,221–224 and axial 175,196 directions. Other studies have reported measurements of the velocity field, 109,175,224,225 angle of repose, 109 and surface shape. 226 This simple geometry produces a steady flow where the mixing rate and mechanisms can be readily studied.

The mixing in other rotating devices has also been studied. The study of twin shell or V-blenders has focused on the flow and mixing mechanisms, <sup>45</sup> the sensitivity to simulation parameters, <sup>192</sup> different fill patterns, <sup>227</sup> and the extent of mixing over long times. <sup>228</sup> The effectiveness of a V-blender has also been compared to a double cone blender. 45 Other work has focused on various tote blenders. Computational mixing results in a pyramidal tote blender have compared favorably to 1:1 scale experiments with glass beads. 229 The mixing mechanisms in a conical tote blender were examined as a function of percent loading.230 Other work compared mixing in a conical tote blender with that in a V-blender.<sup>227</sup> Further, a study of baffles in a cylindrical tumbler rotating on an axis perpendicular to that of the cylinder has shown the importance of baffle design.<sup>231</sup>

Other studies have examined various mixing equipment with an impeller, such as the vertical

axis cylindrical mixer with a rotating blade shown in Figure 9. Stewart et al. 190 simulated the mixing in such a device and compared the results with experimental measurements obtained via positron emission particle tracking. The overall flow was successfully predicted by DEM, but the model did not closely predict the velocities in all regions of the bed. Later work by the same authors examined the granular microstructure during mixing. 232 Others have determined an optimal fill level 233 and examined the effect of different impeller shapes. 191

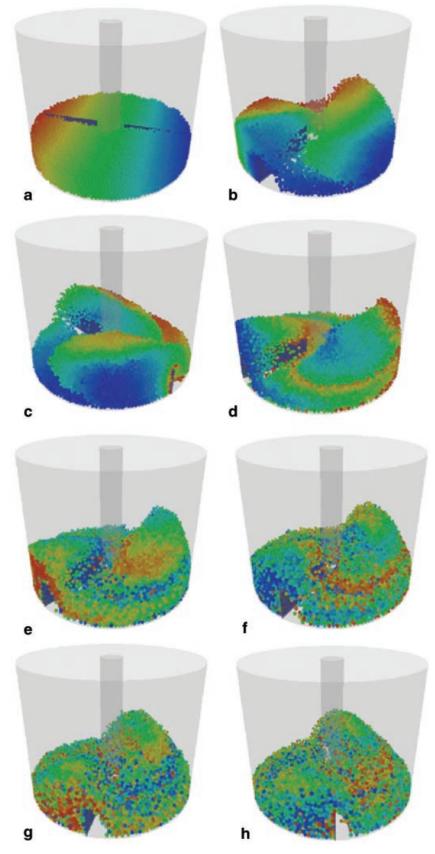
Terashita et al.<sup>234</sup> studied the particle segregation due to density differences while Zhou et al.<sup>235</sup> extended this work and examined segregation due to both density and size differences. Finally, a hybrid model was proposed and implemented in a horizontal convective mixer with an impeller.<sup>199</sup> This hybrid model employs DEM to model regions of high shear and a stochastic compartment model for the remainder of the domain.

Two other mixing devices have also been investigated with DEM. A helical ribbon blender has been analyzed, <sup>180,236</sup> and results show how the flow patterns and extent of mixing are affected for various rotational speeds and fill levels. <sup>236</sup> Further, the suitability of static mixers for bulk solids has been reviewed <sup>237</sup> and studied using DEM. <sup>238</sup> The observed flow regimes from the DEM study compared well with previous experimental results.

#### Granulation

The granulation process and the many approaches to modeling it have been summarized in a review by Cameron et al.<sup>239</sup> While a population balance models have played a key role in granulation process modeling, discrete element models have also played a small, yet significant part of this modeling effort thus far. Much of this DEM work has occurred only in the past few years.

One approach used in DEM modeling of granulation was taken by Gantt and Gatzke. 240 In this work, a high shear granulator was modeled using a DEM model that incorporated three key mechanisms of granulation: coalescence, consolidation, and breakage. The rates of each mechanism are directly simulated and incorporated to model a dynamic particle size distribution (PSD). The results from this DEM model were in good agreement with previous approaches using



**Figure 9.** Images from a DEM simulation of a bladed mixer after (a)  $0.1 \, \mathrm{s}$ , (b)  $0.4 \, \mathrm{s}$ , (c)  $0.7 \, \mathrm{s}$ , (d)  $1 \, \mathrm{s}$ , (e)  $2 \, \mathrm{s}$ , (f)  $3 \, \mathrm{s}$ , (g)  $4 \, \mathrm{s}$ , and (h)  $5 \, \mathrm{s}$ . This figure is reprinted from Bertrand et al.  $^{180}$  with permission from Elsevier.

population balances;<sup>241</sup> however, the DEM approach also has the ability to model dynamic operating conditions.

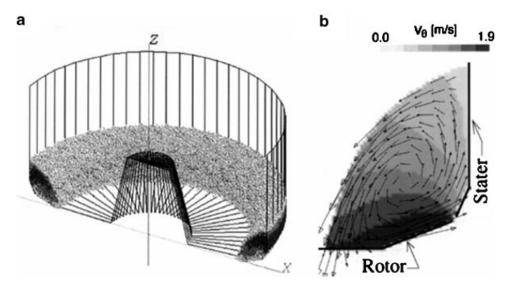
Gantt et al.<sup>242</sup> used another approach where a DEM model with periodic boundary conditions was used to represent flow in a high-shear granulator. The particle collision statistics compiled by the DEM simulation were used to develop a coalescence kernel. This kernel was then used with a Monte Carlo method to solve a multi-dimensional population balance.<sup>243</sup> Good agreement with experiments was observed, but could be improved with better incorporation of material properties such as wet granule yield strength, Young's modulus, and asperity size.

The agglomeration of particles is typically modeled using one of two methods. In the first method, the discrete elements always represent the primary particles.<sup>244</sup> Thus, the agglomerates many consist of many discrete elements. This method retains information on agglomerate porosity and morphology, and subsequent breakage is readily modeled. 245 This approach has also been used to study the impact of two agglomerates at different impact velocities and binder fluid properties. <sup>246</sup> In the second method, agglomerates are modeled as a single discrete element, the mass and volume of which are determined from those of the constituent primary particles. 247 This method allows for the simulation of larger systems since fewer discrete elements are required. The major disadvantage or limitation is that agglomerate

breakage cannot be readily modeled using this method.

A wide variety of different granulation equipment has been modeled with the DEM approach. Talu et al.<sup>248</sup> modeled agglomeration and breakage in 2D shear flow of a mixture of "wet" and "dry" particles showing the effect of the amount of liquid binder, the Stokes number, and the Capillary number on the PSD. Muguruma et al.<sup>70</sup> modeled the centrifugal tumbling granulator shown in Figure 10 where the liquid was uniformly distributed. The resulting velocity profiles were in agreement with experimental data using glass beads of same size. Mishra et al. 249 examined the agglomeration of particulates in a rotary drum with a model that included a spray zone and also considered the drying of particles. Granulation in fluidized beds has also been modeled. Goldschmidt et al.<sup>247</sup> modeled granulation in a top-spray fluidized bed. In that work, both the particles and droplets were modeled using the hard-particle approach. Results showed the development of the PSD over time as a function of the fluidization velocity and the spray pattern. In a similar study, Link et al. 250 examined granulation in a spout fluidized bed.

Dry granulation techniques have also been explored by a few authors with mixed success thus far. Odagi et al.<sup>251</sup> modeled a roller compaction process using DEM, but the pressures in the nip region were several orders of magnitude smaller than experimental values. In contrast,



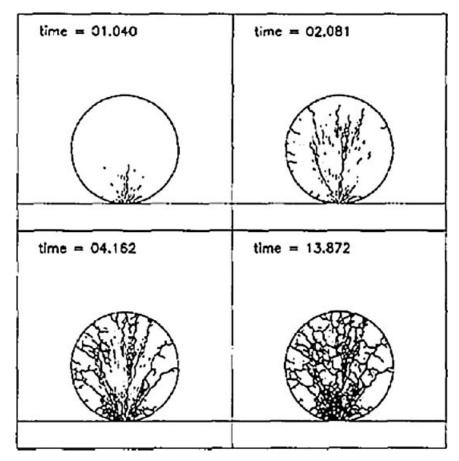
**Figure 10.** A DEM model of a centrifugal tumbling granulator with (a) particle positions and (b) particle velocity profiles. This figure is reprinted from Muguruma et al.<sup>70</sup> with permission from Elsevier.

Dec et al.<sup>252</sup> modeled roller compaction using the finite element method with good success. Thus, in cases where large pressures are present, such as the nip region of a roller compacter, a continuum model may be superior to DEM approaches. Additionally, Horio<sup>253</sup> examined dry granulation in a fluidized bed. This so-called pressure swing granulation relies on the intrinsic cohesivity of particles to form relatively weak agglomerates.

# Milling

The discrete element method approach to modeling comminution devices has provided insight into particle fracture patterns and the change in PSD. This work is usually focused on one of two processes. The first is the study of particle fracture and/or attrition while the second involves the dynamics of materials within the milling device. Several researchers examining particle

fracture have developed models exploring the effect of impact velocity and bond strength on fracture. The approaches taken to model breakage are varied. One approach, originally developed by Potapov and Campbell, <sup>254</sup> creates a nonporous, 2D agglomerate of Delaunay triangles<sup>254–256</sup> or arbitrary polygons via a Voroni tessellation. 255–257 Agglomerate fracture, as shown in Figure 11, occurs when the tensile strength of the glued joints is exceeded. In another approach, agglomerates are created with randomly arranged spherical particles and held together with surface adhesive forces. 80,258-261 Breakage of the agglomerate results when tensile forces between spheres exceed the surface adhesive forces. The mode by which agglomerates break is dependent on both the agglomerate strength and the impact velocity where large velocities may cause shattering, but smaller velocities may only lead to a plastic deformation of the agglomerate. A DEM and experimental study of weakly agglomerated



**Figure 11.** Images from a DEM model showing the impact and fracture of a single particle consisting of 2043 triangular elements. This figure is reprinted from Potapov and Campbell<sup>254</sup> with permission from Elsevier.

lactose particles has shown that they tend to deform plastically rather than fracture. <sup>258</sup>

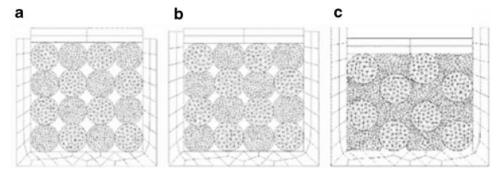
A different approach estimates the particle breakage using an analytic model. For example, Neil and Bridgwater, 262 following the work of Gwyn, 263 attempted to characterize attrition with a single parameter irrespective of the equipment causing the attrition. Another model proposed by Ghadiri and Zhang<sup>264</sup> models particle damage via a chipping mechanism rather than fracture as in the previous approaches. The chipping results in surface damage that is assumed to be a function of impact velocity and the particle size, hardness, and critical stress intensity. Ghadiri et al. 265,266 also model both the fracture and chipping of particulates. It is found that chipping occurs at low impact velocities while fracture occurs at higher values.

Discrete element methods have also been used to model the dynamics of granular material and/or grinding media in comminution devices. Several studies have focused only on such dynamics and have not considered breakage of the particulate material. <sup>267–270</sup> Other work has incorporated a fragmentation or chipping model discussed above into a DEM simulation of particulate motion in a comminution device. These simulations have modeled breakage of nonporous agglomerates in various devices including a shear cell, <sup>271</sup> a semiautogenous (SAG) mill, <sup>272</sup> and a Hicom mill. <sup>273</sup> Similar work has incorporated Ghadiri's models for particle breakage and chipping in studies of shear deformation <sup>274</sup> and jet milling. <sup>275</sup>

### Compression

Various aspects relating to tablet compression have been thoroughly studied using DEM techniques. The process of die filling 140,152,166,276-278 is one aspect that has been studied in both the pharmaceutical and powder metallurgy industries. These studies typically describe the fill rate and particle flow field into a die and make comparisons with experimental results. Other results show the effects of powder characteristics such as friction, <sup>140,152,278</sup> particle-to-die size ratio, <sup>152</sup> effect of shoe speed, <sup>140</sup> and the effect of air pressure. 140,166 The compression of materials within a die has been studied 139,279-285 while others have investigated compression in a periodic system to eliminate die wall effects. 286-288 Rearrangement of particles during the compaction process has also been shown to be important. 75,283,289,290 In much of the compression work, the finite element method (FEM) has been coupled with DEM models so that systems with large particle deformations and large relative densities can be modeled. <sup>139,280,281,283,285,291</sup> This approach permits the modeling of materials with a wide variety of mechanical properties, such as the hardness and ductility, as shown in Figure 12. However, this method is more computationally intensive as individual particles must be meshed and analyzed via FEM.

A few other aspects related to compression have also been approached with DEM modeling. Couroyer et al.<sup>292</sup> used a previously mentioned particle fracture model<sup>265</sup> to examine crushing during compression for various material properties. Martin<sup>293</sup> investigated the springback during the unloading process. Springback was shown to be affected by the material properties, compaction history, and the die wall. In addition to elasticity, the decohesion of particles previously in contact was also shown to be an important factor during unloading. Sweeney and Martin<sup>294</sup> used DEM to determine the microstructure of a high-density



**Figure 12.** Modeling compression of a mixture of hard and soft ductile materials using a coupled DEM-FEM model. This figure is reprinted from Ransing et al.<sup>285</sup> with permission from The Royal Society.

compact. From this microstructure, they calculated a pore size distribution for varying relative densities that was in good agreement with experimental measurements. Hassanpour and Ghadiri<sup>295</sup> examined the Heckel analysis used for powder compression and determined that it may not be valid for all conditions. Rather, results show it is valid only in a limited range of Young's moduli and yield stresses.

# **Film Coating**

There has been relatively little work published on film coating studies using DEM. Certainly, the results from much of the work with rotating drums discussed in the "Blending" Section can be directly related to the mixing in pan coaters. However, the work by Yamane et al. 296 was one of the first studies to specifically study tablet motion in a pan coater. Inter-tablet uniformity was assessed via the standard deviation of the time each tablet was on the surface as a function of pan loading, particle size, and rotational speed. While the model assumed spherical particles, the results were in qualitative agreement with experimental results. Later DEM work 109 modeled the flow of spheres glued together into two-sphere clusters and performed a detailed comparison with experimental data from a system of mustard seeds obtained with MRI. The particle velocities compared well with those measured experimentally, but only after the coefficient of friction and particle shape (elongation) were adjusted. Further, the model predicted a slightly flatter free surface than was observed experimentally. A similar study by Pandey et al. 176 reported the effect of pan loading and speed on the particle velocities and the dynamic angle of repose. There was good agreement of the particle velocity distributions in the spray zone, and the angles of repose also showed good agreement despite difficulties in measurement of due to the nonlinear surface of the cascading layer.

Additional work related to film coating has looked at different geometries and the assumption of spherical tablets. Nakamura et al.<sup>297</sup> used a coupled DEM-CFD model to assess the intertablet coating uniformity in a rotating fluidized bed (RFB) as shown in Figure 13. The RFB coating results were compared with numerical results from a conventional fluidized bed. The RFB was shown to produce significantly less variability in applied coating. Other work has focused on

incorporating round convex tablet shapes into DEM models of pan coating equipment. 116

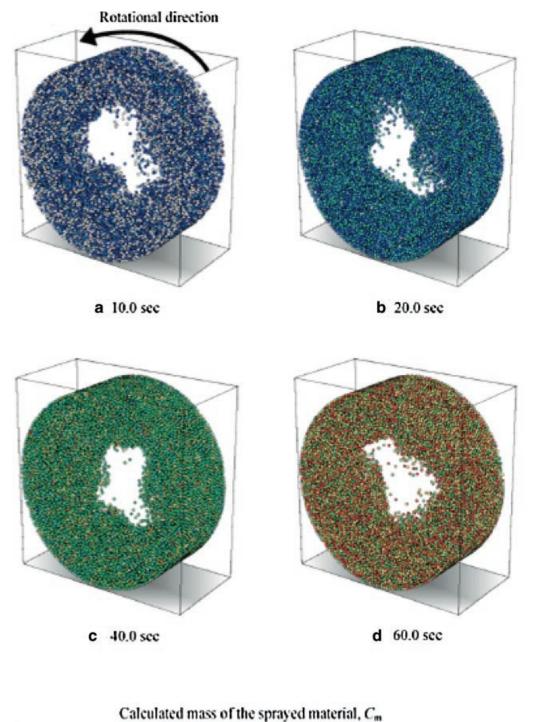
A different approach has applied the Monte Carlo method to study film coating. Nakamura et al. 14 analyzed the coating uniformity in a tumbling fluidized bed as a function of process conditions using Monte Carlo techniques. Similar work explored the coating uniformity in a Wurster fluidized bed coater. 15 This work showed that an upward cylindrical spray shape better matched experimental results than a solid conical spray shape. It also showed that coating uniformity improved with larger gas velocity and gap height.

### **CONCLUSIONS**

The discrete element method has been used to gain a better understanding of many important processes applicable to the pharmaceutical industry. Several aspects of key operations such as transport, blending, granulation, milling, compression, and film coating have been studied. Work has ranged from fundamental studies in simple geometries that attempt to gain insight into the underlying physics to studies of large process equipment that attempt to elucidate the effects of particle properties and operational parameters on the process.

Historically, DEM models have been limited to a modest number of spherical particles with no interstitial fluid effects. Recent DEM studies, however, are increasingly incorporating nonspherical particles, complex system geometries, noncontact cohesive forces, or interstitial fluid effects via a coupled model with CFD. Additionally, system sizes continue to increase into the hundreds of thousands of particles. Despite the increasing focus on these advanced systems, the effects of nonspherical particles, cohesion, and interstitial air effects are still not well understood. These effects are particularly important for typical pharmaceutical materials, and further study is warranted.

The future of DEM modeling looks very promising. This technique can be used to glean a wide variety of information from many diverse and complex systems comprised of particulate matter as highlighted above. The foremost limitation however, is the vast computational power inherently required which, in turn, limits the number of particles that can be modeled. With the ever increasing power and speed of computers



0 4,000 8,000 12,000 16,000 20,000 24,000 28,000 32,000 36,000 40,000 [-]

**Figure 13.** Visualization of the coating mass at four different times in a DEM model of a fluidized bed coater. This figure is reprinted from Nakamura et al.  $^{297}$  with permission from the Pharmaceutical Society of Japan.

however, this limitation will become less restrictive. Other concerns, including those that center on the typical assumptions of spherical particles and noncohesive interactions, are the focus of many recent and ongoing studies aimed at developing and refining models that are both accurate and computationally efficient. Incorporation of the results from these efforts will greatly expand the scope of systems that can be quantitatively modeled using DEM, thereby increasing its utility as a predictive design tool.

### **ACKNOWLEDGMENTS**

The authors wish to acknowledge various Pfizer colleagues for their helpful comments and suggestions which have been a great aid during the preparation of this report. Valuable discussions were held with C. Bentham, D. Kremer, M. Mullarney, V. Swaminathan, and N. Turnbull.

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