

Analysis Note of the Experiment

Supplementary note on diffractive analyses of 2015 proton-proton data:

- [1] Measurement of Central Exclusive Production of $h\bar{h}$ pairs ($h = \pi, K, p$) with Roman Pot detectors in diffractive proton-proton interactions at $\sqrt{s} = 200$ GeV
- [2] Measurement of particle production with Roman Pot detectors in diffractive proton-proton interactions at $\sqrt{s} = 200$ GeV

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Abstract

In this note we present supplementary material related to analyses of diffractive processes based on 2015 data from proton-proton collisions at $\sqrt{s} = 200$ GeV. This dataset was collected with newly installed Roman Pot detectors in Phase II* configuration which ensured efficient triggering and measuring diffractively scattered protons.

We focus on calculation of efficiencies, corrections to Monte Carlo simulations and derivation of systematic uncertainties of our measurements.

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1. TPC track quality cuts

2. Efficiencies

2.1 TPC track acceptance and reconstruction efficiency

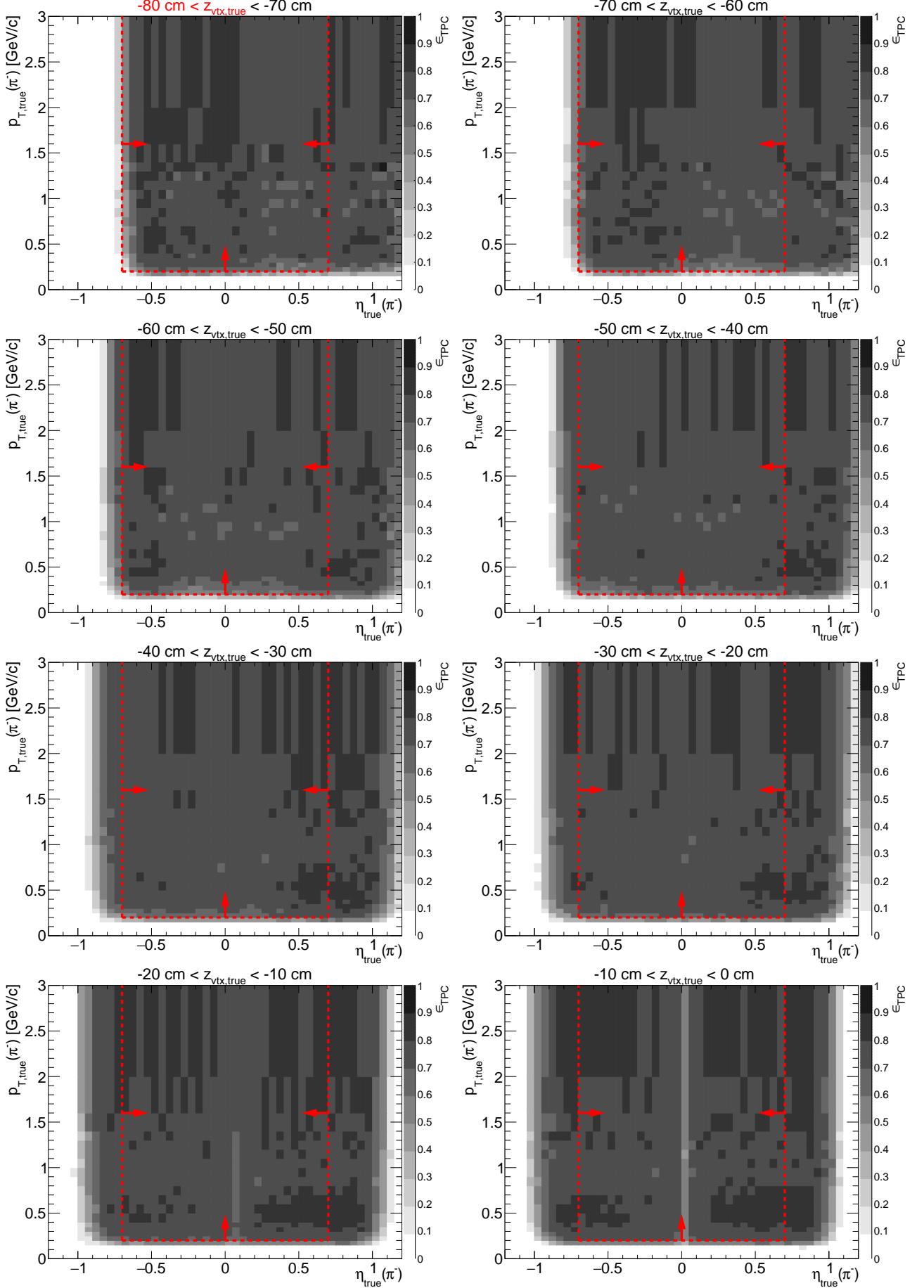
We define joint acceptance and efficiency of the reconstruction of a track in the TPC, ϵ_{TPC} , as the probability that particle from the primary interaction generates signal in the detector which is reconstructed as a track that satisfies all quality criteria and whose p_T and η are within the kinematic region of the measurement (cuts ?? and ??).

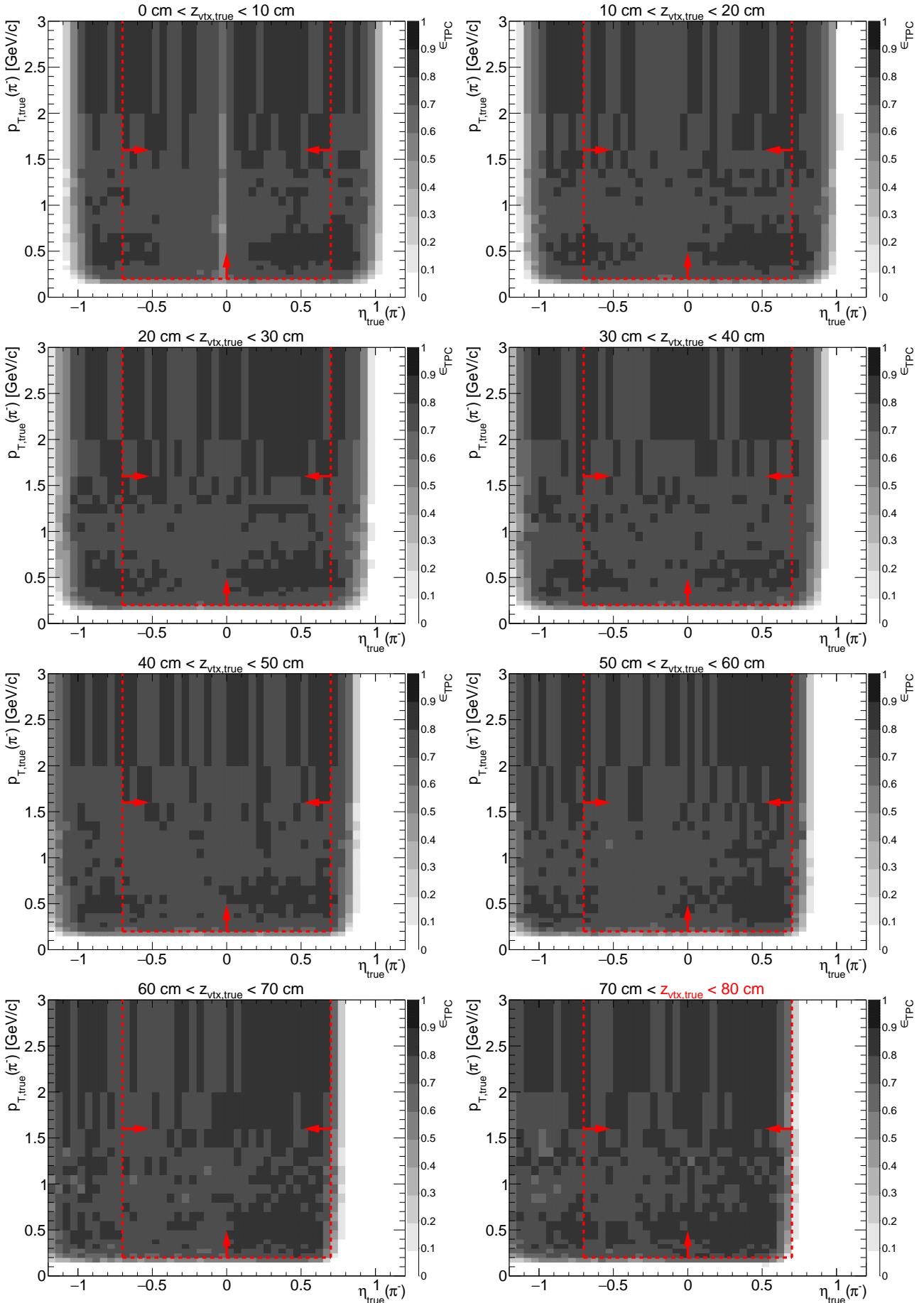
Thechnically this quantity is derived from STARsim MC embedded into zero-bias triggers in the following procedure:

1. True-level primary particles of given ID and charge, with all physics (p_T^{true} , η^{true}) and detector (z_{vx}) quantities within defined region of the measurement, are selected (*set A*).
2. Each particle from *set A* is checked if global TPC track with more than half of hit points generated by this particle, was reconstructed. All global tracks which are associated with true-level primary particles and satisfy kinematic and quality criteria (cuts ?? and ??), form *set B*.
3. The joint TPC acceptance and efficiency is calculated as the ratio of the histograms of true-level quantities (such as p_T , η , z_{vx}) for particles from *set B* and particles from *set A*:

$$\epsilon_{\text{TPC}}(p_T, \eta, z_{vx}; \text{sign, PID}) = \frac{(p_T, \eta, z_{vx}) \text{ histogram for particles of given sign and ID from } \textit{set B}}{(p_T, \eta, z_{vx}) \text{ histogram for particles of given sign and ID from } \textit{set A}}. \quad (2.1)$$

Figure 2.1: TPC acceptance and reconstruction efficiency of π^- . Each plot represents the TPC efficiency ϵ_{TPC} (z -axis) as a function of true particle pseudorapidity η (x -axis) and transverse momentum p_T (y -axis) in single z -vertex bin whose range is given at the top. Red lines and arrows indicate region accepted in analyses.





2.2 TOF acceptance, hit reconstruction and track-matching efficiency

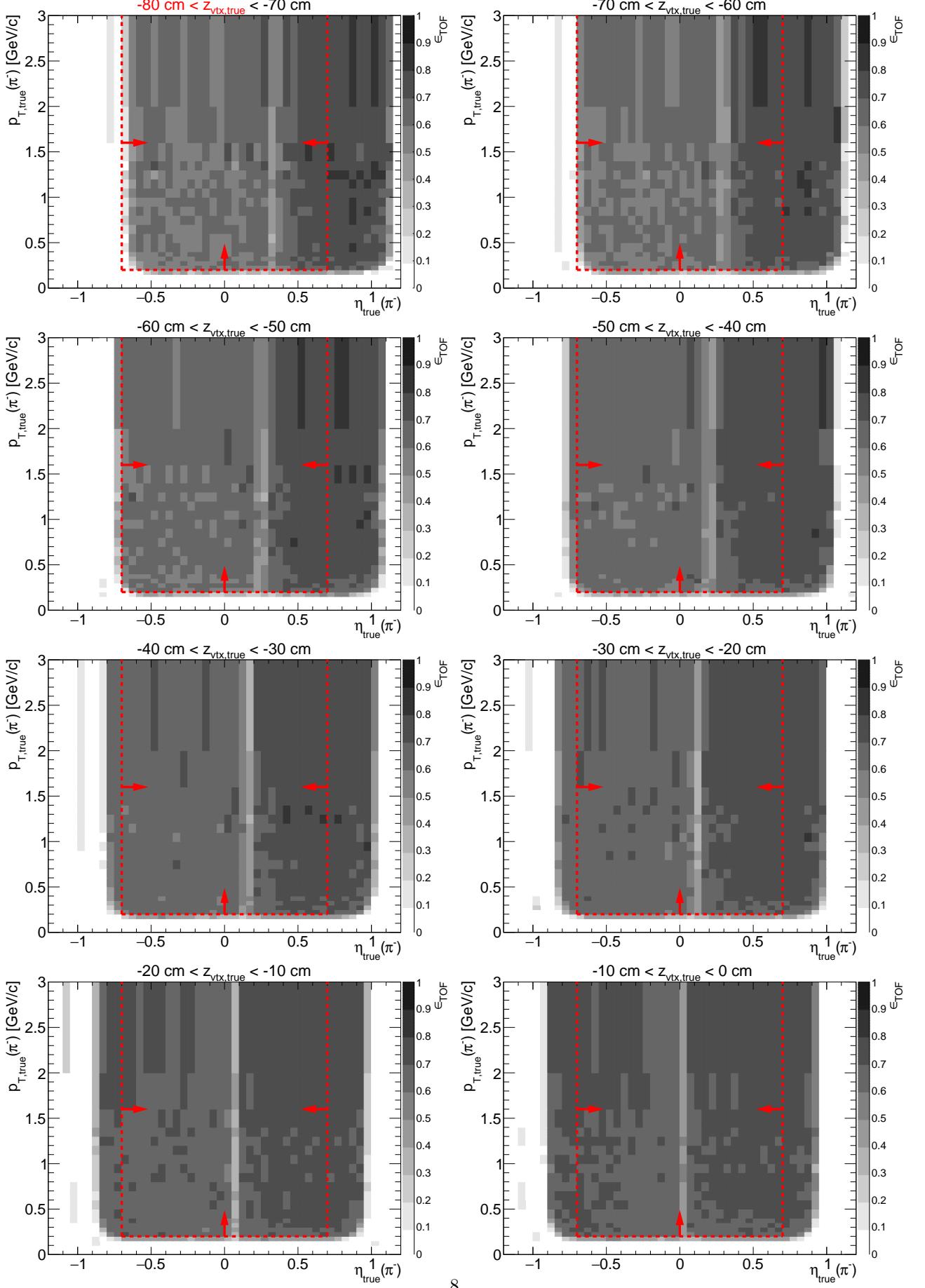
Combined TOF acceptance, hit reconstruction efficiency and matching efficiency with TPC tracks, ϵ_{TOF} , is defined as the probability that the global TPC track that satisfy kinematic and quality criteria (cuts ?? and ??) is matched with hit in TOF (matching flag of the track is different from 0). This quantity is generally referred as “TOF efficiency”.

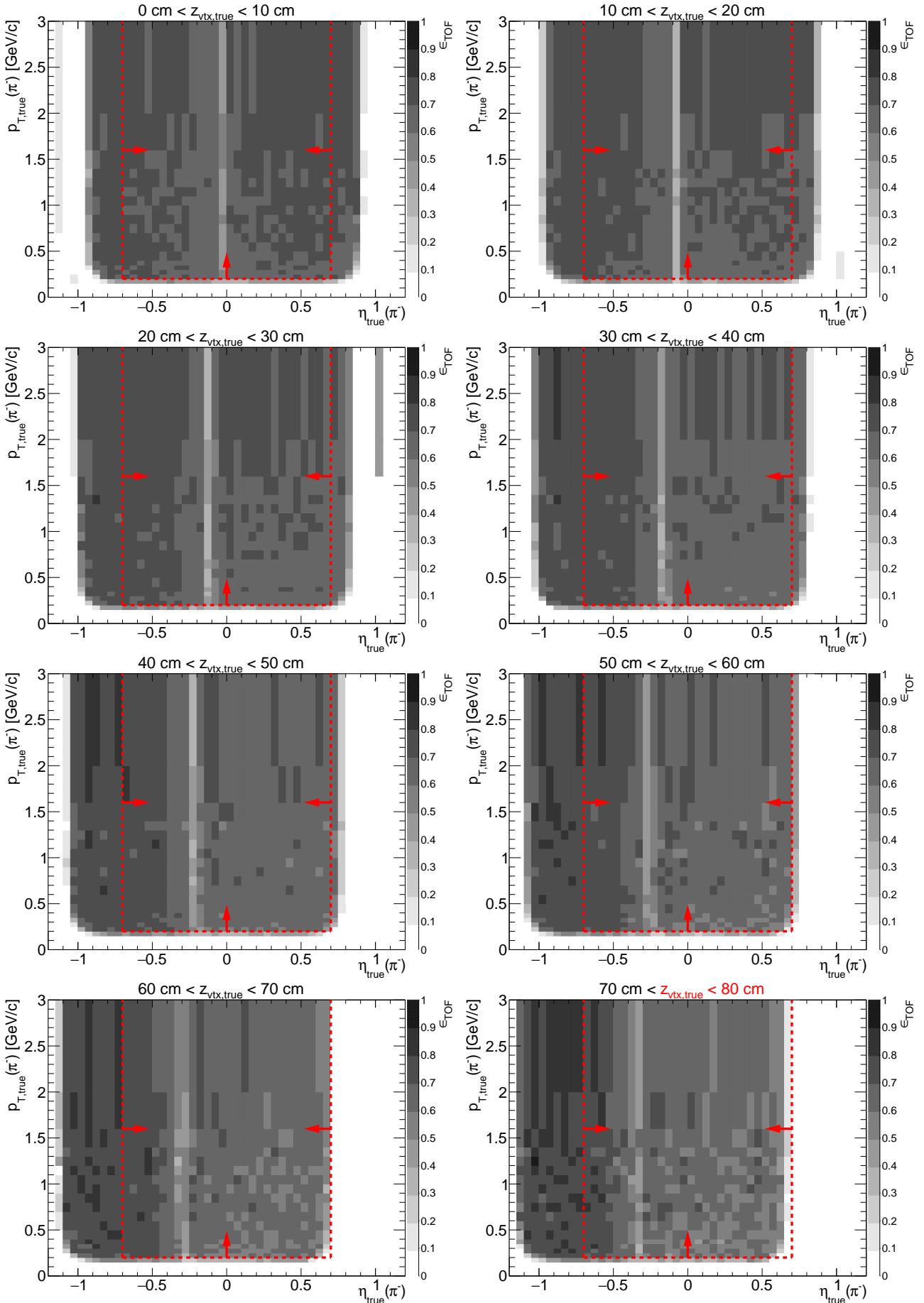
It is calculated in two ways. In the first approach the STARsim MC embedded into zero-bias triggers is used. Tracks belonging to *set B* from Sec. ?? are utilized. From these tracks a sub-sample of tracks with non-zero TOF matching flag is extracted (*set C*). The TOF efficiency is calculated as

$$\epsilon_{\text{TOF}}(p_T, \eta, z_{vx}; \text{sign, PID}) = \frac{(p_T, \eta, z_{vx}) \text{ histogram for particles of given sign and ID from set } C}{(p_T, \eta, z_{vx}) \text{ histogram for particles of given sign and ID from set } B}. \quad (2.2)$$

NEED TO ADD INFO ABOUT REJECTING TRACKS MATCHED WITH HITS THAT ARE IN DEAD MODULES BASED ON THE DATA

Figure 2.2: TOF acceptance, reconstruction and matching efficiency of π^- . Each plot represents the TOF efficiency ϵ_{TOF} (z -axis) as a function of true particle pseudorapidity η (x -axis) and transverse momentum p_T (y -axis) in single z -vertex bin whose range is given at the top. Red lines and arrows indicate region accepted in analyses.





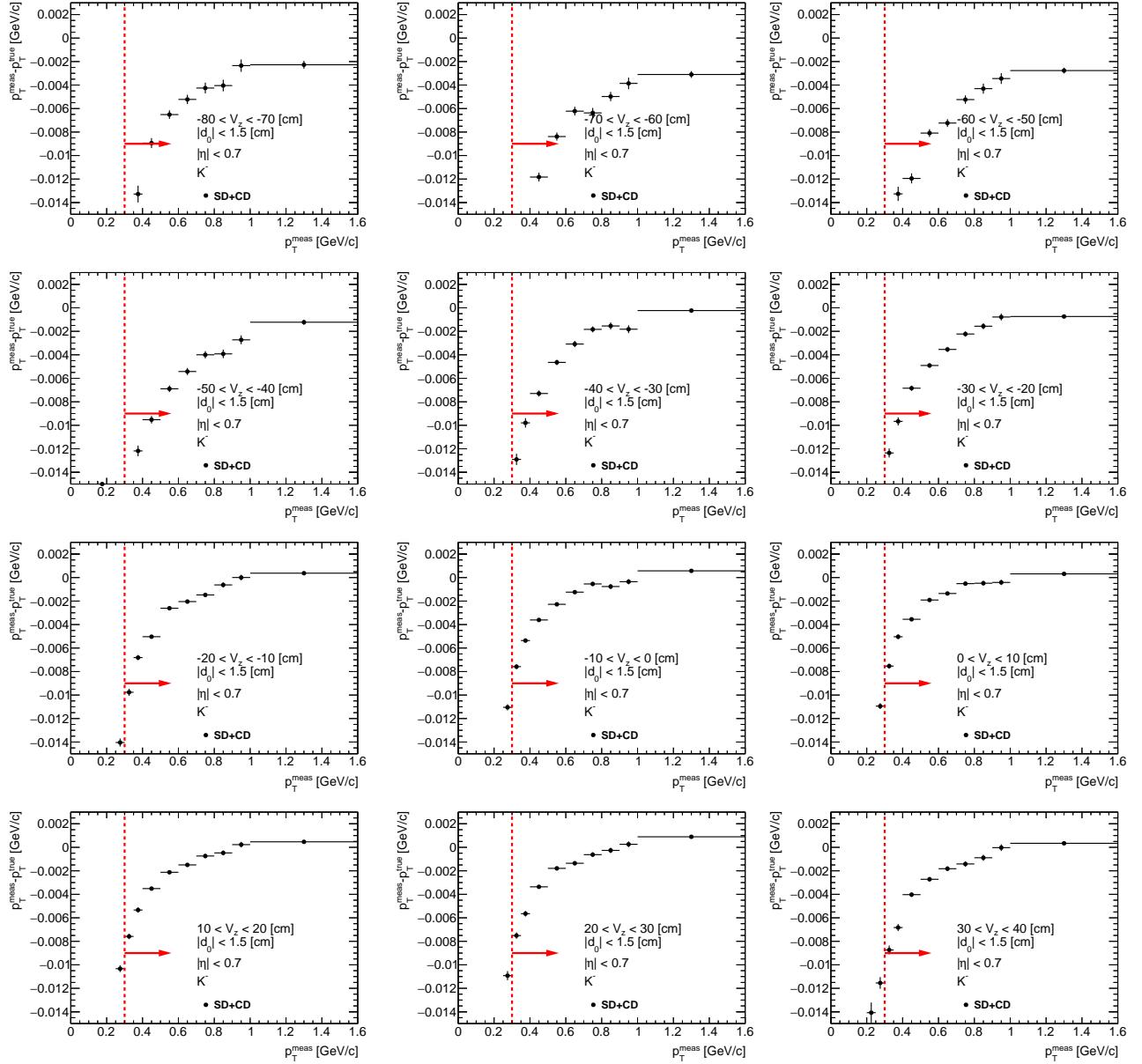
2.3 TPC vertex reconstruction efficiency

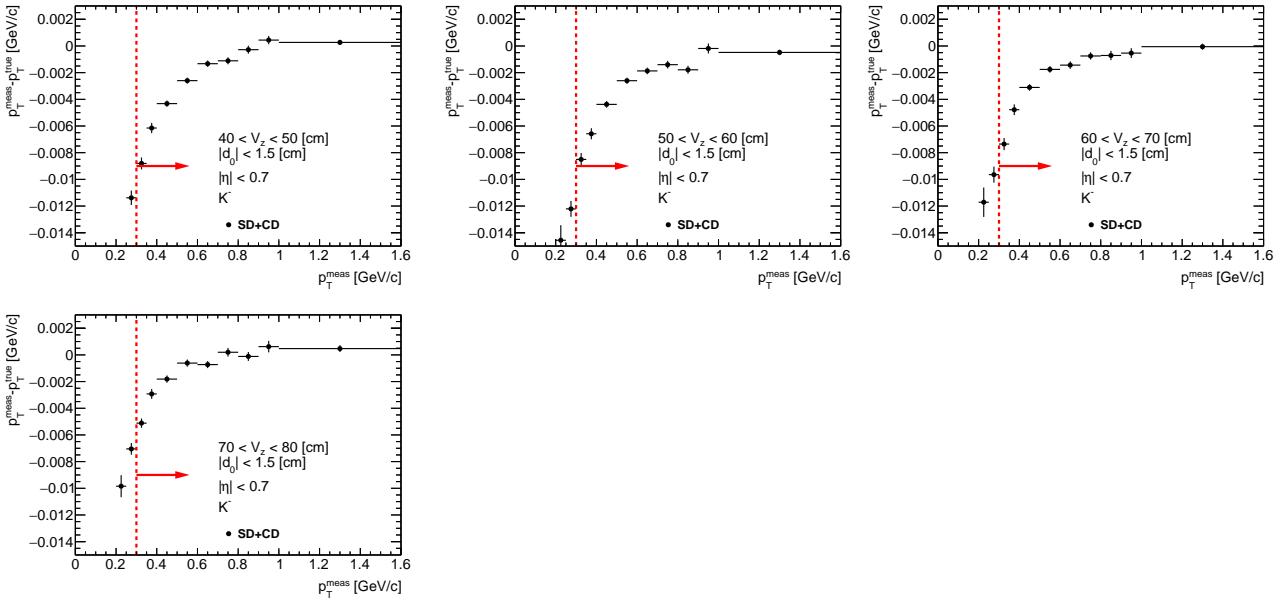
The definition of vertex reconstruction efficiency established in this analysis is the probability that two global tracks, both associated with true-level primary particles from the kinematic region of the measurement, both satisfying kinematic and quality criteria (cuts ?? and ??) and both matched with hits in TOF, form a vertex listed in the collection of reconstructed primary vertices and DCA(R) and DCA(z) of both global tracks calculated w.r.t. this vertex is contained within the limits of cut ??.

3. Energy Loss Correction

Particles passing through the detector material loose energy as they travel. The track momentum p_T^{true} is reconstructed by fitting a helical path to the track points left in the detector. Fitting the track points to an ideal helical track tends to underestimate the momentum due to these energy loss effects. To minimize biases due to this effect, correction procedure is applied during standard track momentum reconstruction procedure for both data and MC simulation. For this procedure all particles are assumed to be pions and the reconstructed momentum p_T^{meas} is corrected by the amount of energy loss for a pion. For anything that is not a pion some rest bias is still present since on average energy loss is specific for each particle type. These biases can be determined from simulated tracks run through GEANT. The correction $p_T^{meas} - p_T^{true}$ was calculated for each particle species as a function of p_T^{meas} , η and z -vertex. The energy loss correction averaged over $|\eta| < 0.7$ for K^- is shown in Fig. 3.1. The energy loss corrections for other particle species are shown in Figs. C.1 to C.9 in Appendix C.

Figure 3.1: Energy loss correction $p_T^{meas} - p_T^{true}$ for K^- as a function of reconstructed transverse momentum p_T^{meas} ($|\eta| < 0.7$) in single z -vertex bin whose range is given on each plot. Red lines and arrows indicate region accepted in analyses.





4. Roman Pot simulation

5. dE/dx adjustment

Particle identification in our analyses is done using merged information from the TPC (specific energy loss of tracks dE/dx) and from the TOF (time of hit matched to TPC track). As can be seen in Fig. E.5, dE/dx information from the MC events simulated in STARsim (in red) poorly matches the data points (black). This results e.g. in large systematic error of estimate of particle identification efficiency.

This problem was discussed under ticket #3272 (Ref. [3]). There were trials to improve the TPC calibration in simulation, but the problem remained. It was finally concluded that the origin of the problem lies in the model of energy loss used in the STARsim, therefore any further action was postponed.

In order to tune simulated response of the TPC in terms of dE/dx , hence also reduce the systematic uncertainty related to particle identification, a correction method was developed based on proper transformation (recalculation) of simulated dE/dx to obtain new dE/dx whose distribution matches the data. We know that n_X^σ (where $X = \pi, K, p, \dots$) variable for particle X follows a gaussian distribution

$$n_X^\sigma = \left(\ln \frac{dE/dx}{\langle dE/dx \rangle_X} \right) / \sigma_{dE/dx}, \quad f(n_X^\sigma) = \mathcal{N}(n_X^\sigma; \mu = 0, \sigma = 1), \quad (5.1)$$

therefore dE/dx itself by definition follows log-normal distribution:

$$f(dE/dx) = \mathcal{LogN}(dE/dx; \mu = \langle dE/dx \rangle, \sigma = \sigma_{dE/dx}) = \frac{1}{\sqrt{2\pi} \cdot \sigma \cdot dE/dx} e^{-\frac{\ln^2 \frac{dE/dx}{\langle dE/dx \rangle}}{2\sigma^2}}. \quad (5.2)$$

The desired transformation should preserve the shape of dE/dx distribution so that it is still described by \mathcal{LogN} , however it should change μ and σ so that these values are equal to ones in the data. The transformation that satisfies above postulate is

$$dE/dx' = c \cdot (dE/dx)^a. \quad (5.3)$$

Parameters of the distribution $\mathcal{LogN}(dE/dx')$ are then

$$\mu' = c \cdot \mu^a, \quad \sigma' = a \cdot \sigma. \quad (5.4)$$

From above we get formulae for parameters of the transformation:

$$a = \sigma'/\sigma, \quad c = \frac{\mu'}{\mu^a}. \quad (5.5)$$

To sum up, one has to find the MPV and width parameter of the dE/dx spectrum of each particle in the data and MC, and use relations (5.5) in order to find parameters of the transformation introduced in Eq. (5.3).

The most challenging part of the task was extraction of the $\langle dE/dx \rangle$ and $\sigma_{dE/dx}$ from the data. In case of MC one can select tracks matched to true-level particles of given ID and thus separate dE/dx of different particles, which makes extraction of the distribution shape straightforward. Unfortunately, it is not possible to apply the same method to the data - here one has to deal with overlapping of the reconstructed dE/dx from different particles. Therefore fits of sum of $f(dE/dx)$ corresponding to different particles were performed to reconstructed track dE/dx in narrow momentum bins. The width of momentum bins was chosen to compromise statistics and validity of assumption of constant parameters of dE/dx distribution over bin range.

It was found during the fitting that log-normal distribution is not a perfect model of the reconstructed dE/dx . The problems with description of the data were mainly in the tail-part of the distribution from single particle. Precise model was necessary to obtain satisfactory quality of fits and trustworthy values of parameters. After some research the best model of dE/dx distribution from single particle was found to be

$$f(dE/dx) = \begin{cases} \frac{A}{\sqrt{2\pi} \cdot \sigma \cdot dE/dx} \exp \left(-\frac{1}{2} \left(\frac{\ln \frac{dE/dx}{\langle dE/dx \rangle}}{\sigma} \right)^2 \right) & \text{for } \frac{\ln \frac{dE/dx}{\langle dE/dx \rangle}}{\sigma} \leq k \\ \frac{A}{\sqrt{2\pi} \cdot \sigma \cdot dE/dx} \exp \left(-k \cdot \frac{\ln \frac{dE/dx}{\langle dE/dx \rangle}}{\sigma} + \frac{1}{2} k^2 \right) & \text{for } \frac{\ln \frac{dE/dx}{\langle dE/dx \rangle}}{\sigma} > k. \end{cases} \quad (5.6)$$

Such form was motivated by the function presented in Ref. [4], here adopted for the log-normal instead of normal distribution. Because the modification of the log-normal distribution is introduced only at high-end tail, the validity of the transformation discussed above still holds. To reduce fit complexity the k parameter was set the same for all particle species and fixed to value equal 2.2, which worked well for both data and embedded MC. Particles and their anti-particles were assumed to have the same dE/dx distributions for a given momentum and were analyzed together. The same track selection was used both for data and MC, with the quality criteria

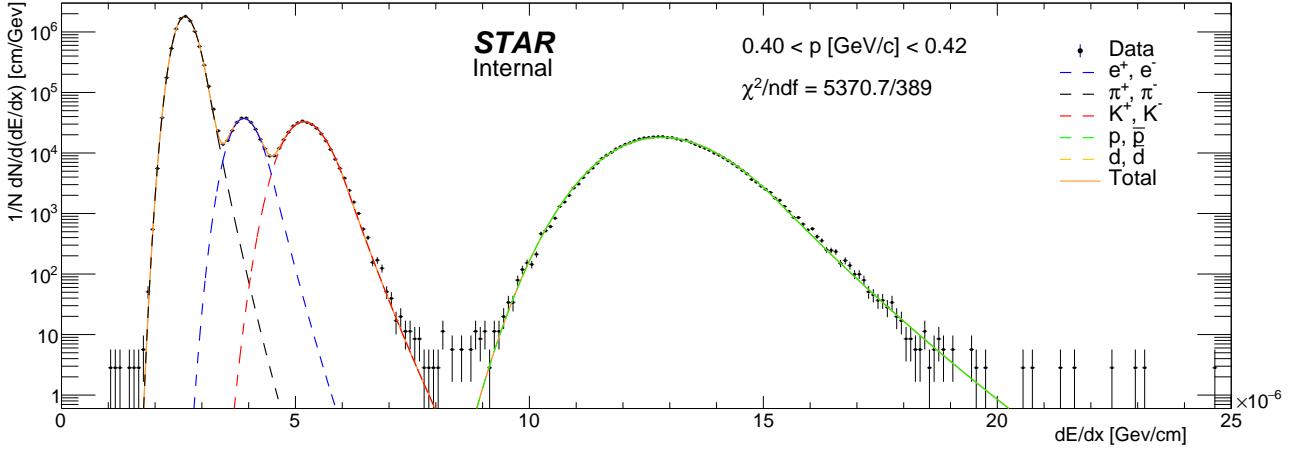


Figure 5.1: Sample fit of sum of functions from Eq. (5.6) corresponding to different particle species to dE/dx spectra in the data in a single momentum bin.

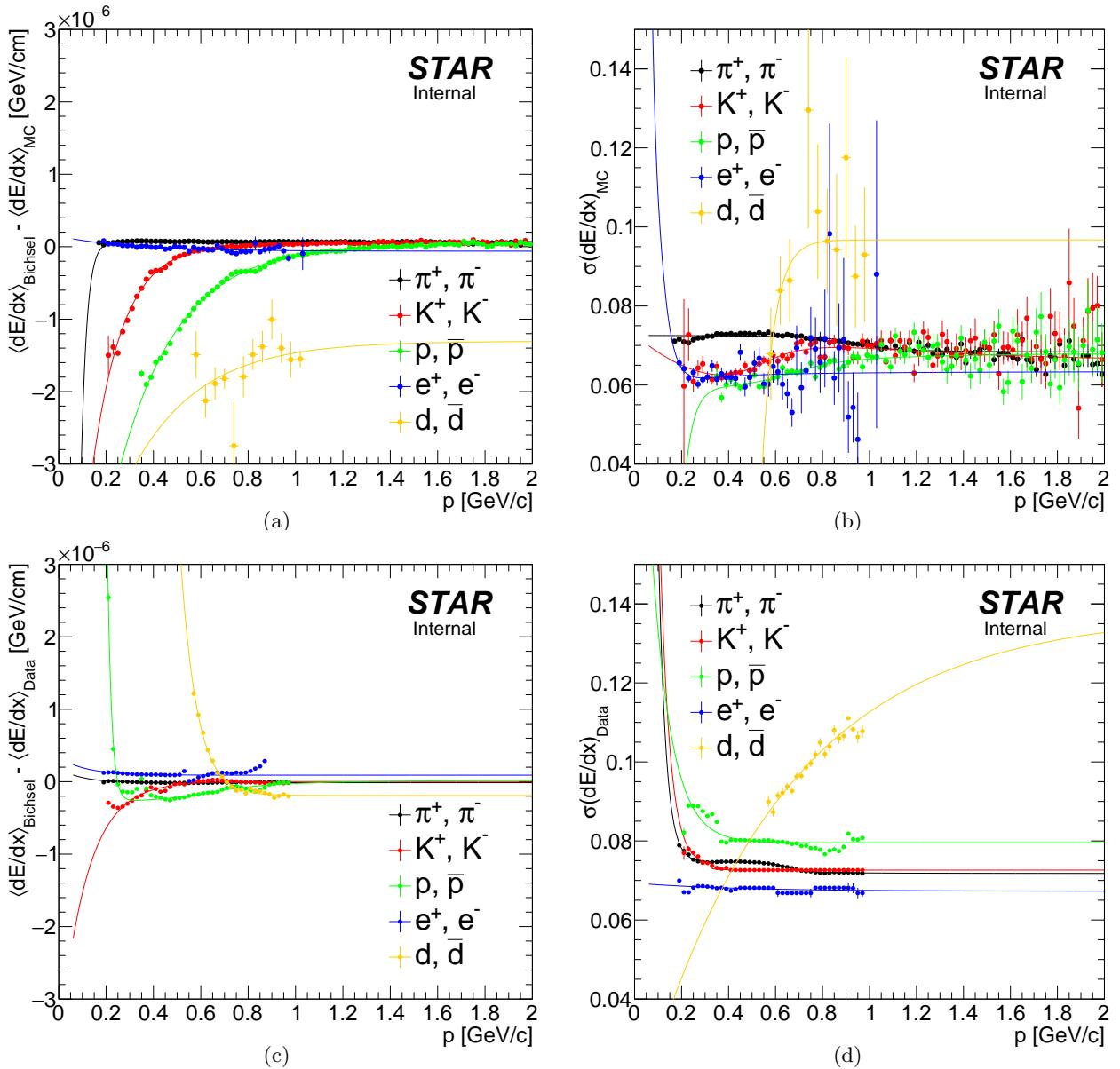


Figure 5.2: Difference between MPV of dE/dx predicted by Bichsel parametrization and obtained from the fit of Eq. (5.6) to dE/dx distribution in the data (5.2c) and MC sample (5.2a) and dE/dx width parameter in data (5.2d) and MC (5.2b) as a function of reconstructed particle momentum for a few particle species. Solid lines represent fits to points of corresponding color. Only statistical errors are shown.

described in Sec. 1. The sample fit in a single momentum bin can be found in Fig. 5.1. Fits in all momentum bins can be found in Appendix E.

Results of the fits for all considered particle species (pions, kaons, protons, electrons and deuterons) are commonly presented in Fig. 5.2 with color markers. Figures 5.2c and 5.2a show the offset of the MPV of reconstructed dE/dx relative to the Bichsel parametrization in the data and embedded MC, respectively, and Fig. 5.2d and 5.2b show the width of reconstructed dE/dx (in the same order). Function able to qualitatively describe dependence of the parameters as a function of track momentum was empirically found to be given by Eq. (5.7):

$$g(p) = P_1 + P_2 \cdot \exp(-P_3 \cdot p) + P_4 \cdot \arctan(P_5 \cdot (p - P_6)) \quad (5.7)$$

This function was fitted to points corresponding to each particle type and fit result is shown in Fig. 5.2 with lines colored in accordance to markers. Values of parameters of above function are tabulated in Tab. 5.1.

PID	$\langle dE/dx \rangle_{\text{Bichsel}} - \langle dE/dx \rangle_{\text{MC}}$						$\sigma(dE/dx)_{\text{MC}}$					
	P_1	P_2	P_3	P_4	P_5	P_6	P_1	P_2	P_3	P_4	P_5	P_6
π^\pm	7.183e-8	-1.647e-4	41.68				0.0705			-1.42e-3	9.860	0.951
K^\pm	4.359e-8	-9.285e-6	7.697				0.0511	0.034	1.675	1.01e-2	4.934	0.528
p, \bar{p}	3.556e-8	-8.621e-6	3.980				0.0630	-7.725	27.17	3.37e-3	5.245	0.670
e^\pm	-6.219e-8	2.065e-7	3.241				0.0354	0.982	26.58	1.79e-2	41.515	0.095
d, \bar{d}	-1.305e-6	-5.268e-6	3.486				0.0967	-1526	18.75			

PID	$\langle dE/dx \rangle_{\text{Bichsel}} - \langle dE/dx \rangle_{\text{Data}}$						$\sigma(dE/dx)_{\text{Data}}$					
	P_1	P_2	P_3	P_4	P_5	P_6	P_1	P_2	P_3	P_4	P_5	P_6
π^\pm	-1.399e-8	2.012e-7	10.39				0.0734	1.907	31.86	-8.20e-4	22.788	0.653
K^\pm	2.325e-9	-3.690e-6	8.712				0.0808	-0.040	7.951	5.62e-3	-17.08	0.269
p, \bar{p}	-1.458e-7	0.6655	59.06	1.171e-7	4.660	0.644	0.0795	0.181	12.12			
e^\pm	9.005e-8	2.494e-7	8.834				0.0680	8.8e-4	1.549			
d, \bar{d}	-1.910e-7	5.637e-3	14.48				0.1161	-0.147	2.890			

(a)

(b)

Table 5.1: Parameters of functions from Fig. 5.2 describing reconstructed track dE/dx as a function of reconstructed momentum for a few particle species in the data (5.1b) and embedded MC (5.1a). Blank cells denote parameters equal 0. Units of parameters P_i are such that if one provides momentum in Eq. (5.7) in GeV/c the resultant offset of dE/dx MPV with respect to Bichsel parametrization is in GeV/cm , and the resultant σ parameter is unitless.

The correctness of the entire procedure described in this section was verified by comparing the reconstructed track dE/dx between the data and embedded MC without and with the dE/dx transformed using Eq. (5.3) and parameters from Tab. 5.1. Some difficulty arised in this comparison due to inconsistent relative content of different particle species in the data and embedded MC sample. Problem was ressolved by separating dE/dx histograms of different particle species in MC (in the same way as it was done for extraction of dE/dx MPV and σ for each particle ID) and fitting the sum of histograms from different particle types to the data histogram (in momentum bins). The only free parameters in the fit were relative contents of histogram from singe particle type in the data histogram. A sample comparison between the dE/dx in data and embedded MC is presented in Fig. 5.3. Comparison in all other momentum bins is contained in Appendix E (Fig. E.5). Fits were done for adjusted dE/dx (filled green). Histograms for unadjusted dE/dx (hashed red) were composed using the same relative content of particles as obtained from the fit of adjusted dE/dx . The ratio of the MC to the data shown in the lower pad of Fig. 5.3 and Fig. E.5 clearly demonstrates better agreement of the MC and the data in terms of position and width of peaks in dE/dx spectrum after the adjustment.

Exactly the same cross-check was done for n_X^σ variables, which are directly used to identify particles in our analyses. They carry information about the distance (in width units) of reconstructed dE/dx from the MPV of dE/dx according to Bichsel parametrization - they are kind of a pull variables. For every track n_X^σ was recalculated using its definition as shown in Eq. (5.8):

$$n_X^\sigma = \ln \left(\frac{dE/dx'}{\langle dE/dx \rangle_{\text{Bichsel}}^X} \right) / \delta(dE/dx), \quad (5.8)$$

where instead of original dE/dx the adjusted one (dE/dx') was used. Uncertainty of reconstructed dE/dx , $\delta(dE/dx)$, was left unchanged. The sample comparison of n_{pion}^σ , n_{kaon}^σ and n_{proton}^σ variables can be found in Fig. 5.4. Comparison in full set of momentum bins is placed in Appendix E. As in case of dE/dx comparison,

in the lower pads of Fig. 5.4 the ratio of the MC to the data also demonstrates better agreement of the MC and the data, which was the goal of the adjustment.

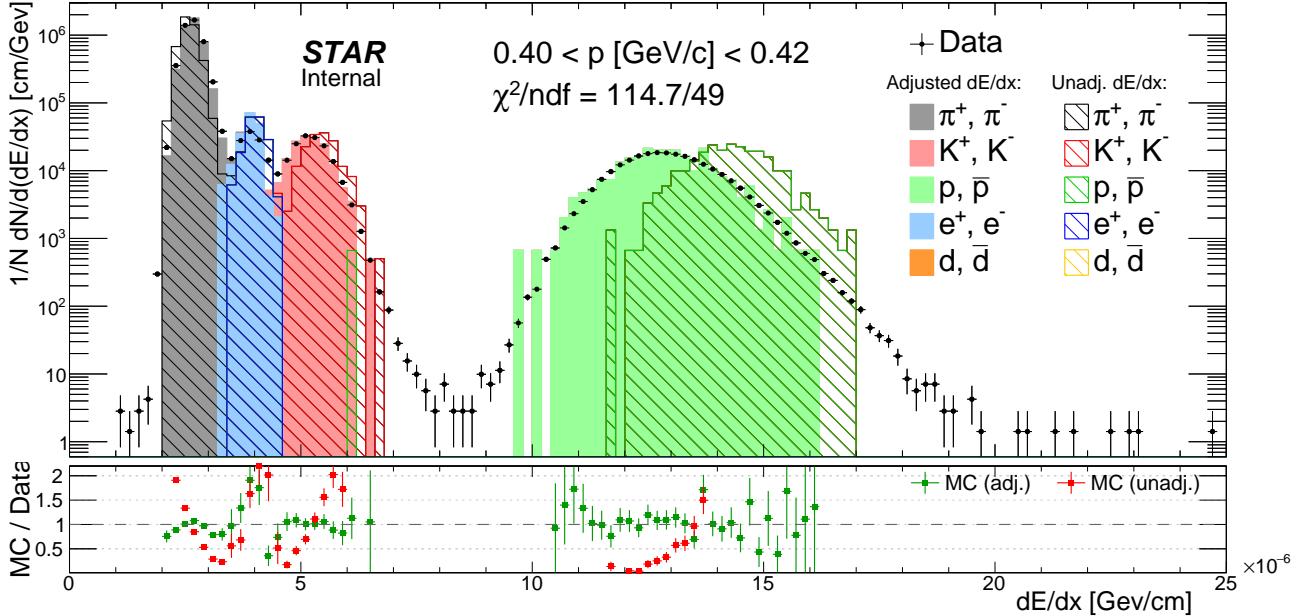


Figure 5.3: Sample comparison of dE/dx spectrum between data and embedded MC (before and after dE/dx adjustment) in a single momentum bin. Lower pad shows the ratio between embedded MC and data before and after dE/dx adjustment. In both upper and lower pads the same color code is used. Only statistical errors are shown. Due to limited statistics of embedded MC some data points do not have corresponding entries in MC.

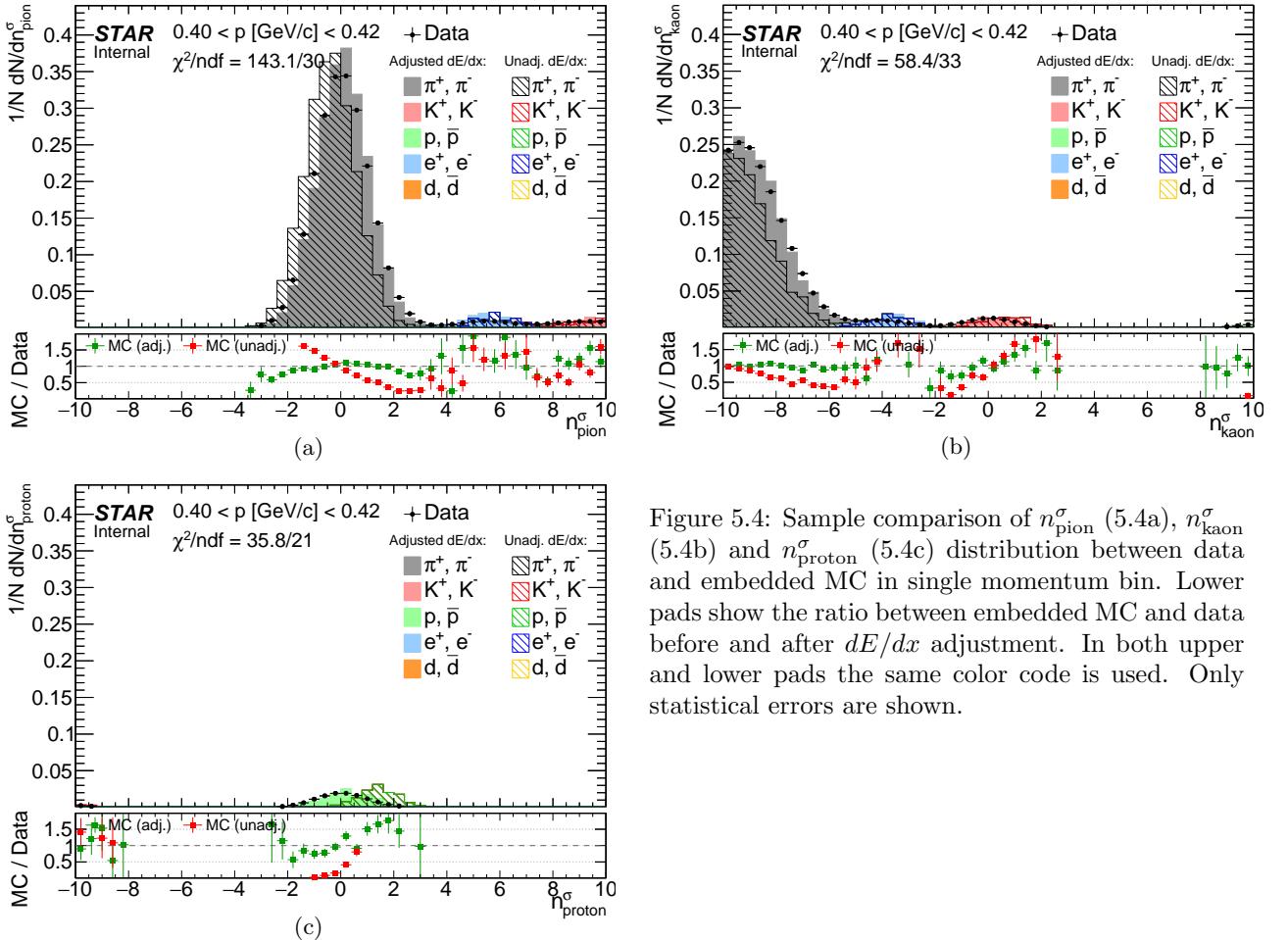


Figure 5.4: Sample comparison of n_{pion}^σ (5.4a), n_{kaon}^σ (5.4b) and n_{proton}^σ (5.4c) distribution between data and embedded MC in single momentum bin. Lower pads show the ratio between embedded MC and data before and after dE/dx adjustment. In both upper and lower pads the same color code is used. Only statistical errors are shown.

6. TPC track pointing resolution adjustment

It was found during the analysis that distributions of quantities which describe the pointing resolution of the TPC tracks do not agree well between the data and embedded MC. Namely, the resolutions of the global helices associated with the tracks were found to be significantly better in the STAR simulation than in the data, what manifests as narrower DCA and d_0 distribution in the embedded MC, comparing to corresponding distribution in the data (Fig. 6.5). This issue was discussed under ticket #3332 (Ref. [5]).

This problem could affect the momentum resolution and thus all other resolutions and response matrices used in data unfolding. Therefore the resolution adjustment procedure was performed to find appropriate parameters of the “artificial” helix deterioration and finally obtain agreement between DCA and d_0 distributions (and all related resolutions) in the data and embedded MC.

In order to reduce pointing resolution in the MC an additional smearing of the helix radius $\sigma(R)$ was introduced. Based on d_0 comparison in Fig. 6.5a it was decided to account also for the systematic bias of the helix radius $\Delta\mu(R)$ ¹, which may be present e.g. due to differences in the material budget used the simulation and reconstruction. Both smearing and bias of the helix radius were introduced only for MC tracks which were matched with the true-level particles since only simulated tracks require adjustment (tracks from zero-bias event used in embedding already contain all detector effects).

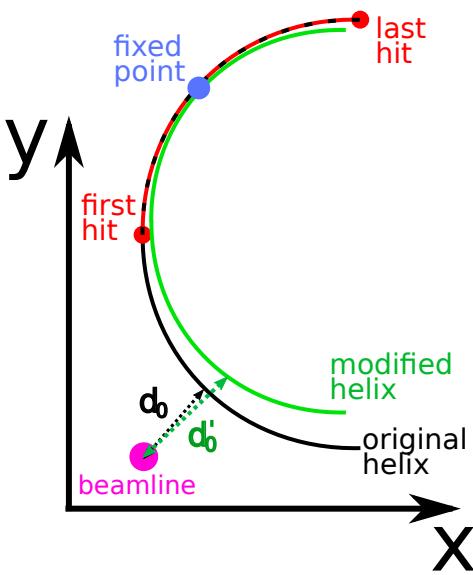


Figure 6.1: Sketch of helix modification procedure and d_0 calculation.

Extraction of $\Delta\mu(R)$ and $\sigma(R)$ parameter required to achieve agreement of pointing resolution between embedded MC and the data involved a few steps, as listed below:

1. Series of d_0 histograms in bins of p_T (100 MeV/c wide) was prepared, each for different size of distortion (different $\Delta\mu(R)$ and $\sigma(R)$) of global helix of the TPC tracks matched with true-level particles (example plot in single p_T bin is shown in Fig. 6.2):

- (a) for each set of parameters $\Delta\mu(R)$ and $\sigma(R)$ the helix radius R was recalculated independently for each track following the Eq. (6.1):

$$R' = R \times \mathcal{N}(1 + \Delta\mu(R), \sigma(R)), \quad (6.1)$$

- (b) new helix of a radius R' was assigned to a track and used to calculate d_0 . The modified helix was obtained by changing the radius of original helix from R to R' with a fixed middle point between the first and last TPC hit of a global track represented by the helix (Fig. 6.1). The momentum of the track was also recalculated:

$$p'_T = p_T \times \frac{R'}{R}, \quad \eta' = \eta \times \frac{R'}{R}. \quad (6.2)$$

2. In each p_T bin the χ^2/NDF was calculated between the data and MC d_0 histogram in a range $-1.5 \text{ cm} < d_0 < 1.5 \text{ cm}$ (corresponding to d_0 cut used in analyses), for every point in parameter space of radius distortion (for every set of $\Delta\mu(R)$ and $\sigma(R)$). An example (single p_T bin) of map of $-\chi^2/\text{NDF}$ in a parameter space is presented in Fig 6.3.
3. In each bin of recalculated p_T the 2-dim parabola $z(x, y; a, b, x_0, y_0, z_0)$ given in Eq. (6.3) ($z = \chi^2/\text{NDF}$, $x = \Delta\mu(R)$, $y = \sigma(R)$) was fitted to $-\chi^2/\text{NDF}$ in the global minimum region to obtain the best-fit distortion parameters.

$$z = z_0 - a(x - x_0)^2 - b(y - y_0)^2. \quad (6.3)$$

4. The best-fit smearing $\sigma(R)$ (equal to parabola parameter y_0) and best-fit bias $\Delta\mu(R)$ (x_0) from individual p_T bins was plotted as a function of global track p_T (Fig. 6.4). Each point was assigned with an error being a quadratic sum of two components: the error on x_0 (y_0) resulting from the parabola fit to $-\chi^2/\text{NDF}$, and length of corresponding semi-axis of ellipsis formed by the intersection of fitted parabola with the xy -plane at $z = z_0 - 1/\text{NDF}$ (from definition of the parameter uncertainty given by the change of overall χ^2 by 1 unit). Resultant formulae for the error of each individual point in Fig. 6.4 are

¹Transverse impact parameter d_0 takes positive value if the beamline is contained inside the helix (in the yz -plane projection), otherwise it is negative. Any asymmetry in the d_0 distribution in the MC with respect to the data indicates presence of systematic difference in reconstructed d_0 , hence also in reconstructed R .

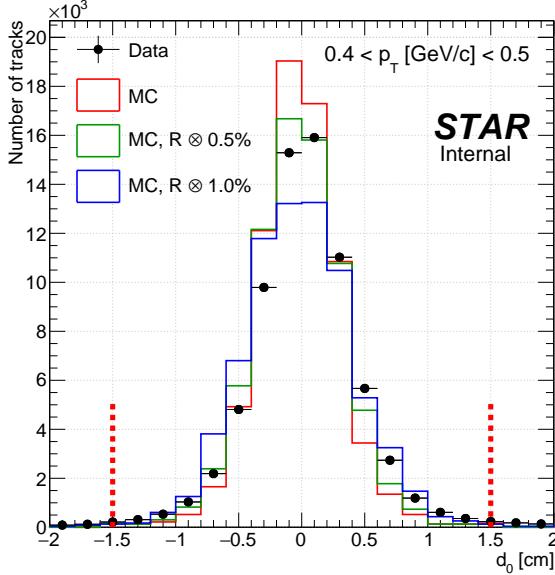


Figure 6.2: Example of comparison of d_0 histograms in single p_T bin in the data (black points) and embedded MC (colored lines) in the procedure of TPC pointing resolution adjustment. Only MC histograms for $\Delta\mu(R) = 0$ and $\sigma(R) = 0, 5 \times 10^{-3}$ and 10^{-2} were shown for explanatory purposes.

$$\delta(\Delta\mu(R)) = \sqrt{\delta_{\text{fit}}^2(x_0) + \frac{1}{2a\text{NDF}}}, \quad (6.4)$$

From Fig. 6.2 one can read that $\text{NDF} = 14$. In calculation of uncertainties correlation of $\Delta\mu(R)$ and $\sigma(R)$ have not been accounted.

5. The empirically determined functions were fitted to points representing $\Delta\mu(R)$ and $\sigma(R)$ dependence on the global track p_T . Their form and values of parameters are given in Fig. 6.4.

Helices of global TPC tracks were deteriorated according to Eq. (6.1) and the parametrizations of global track p_T -dependence of $\Delta\mu(R)$ and $\sigma(R)$ from Fig. 6.4, to verify if better agreement between the data and embedded MC is found after the adjustment. Filled histograms in Fig. 6.5 show d_0 and DCA distributions after the described adjustment, and filled circles in the bottom pad show their ratio to the data points. Clearly, there is much better agreement between embedded MC and the data after the pointing resolution adjustment. Remaining differences may arise from incomplete theoretical model of the CEP process implemented in GenEx leading to different p_T spectra of the data and the model (e.g. model does not contain resonant $\pi^+\pi^-$ production).

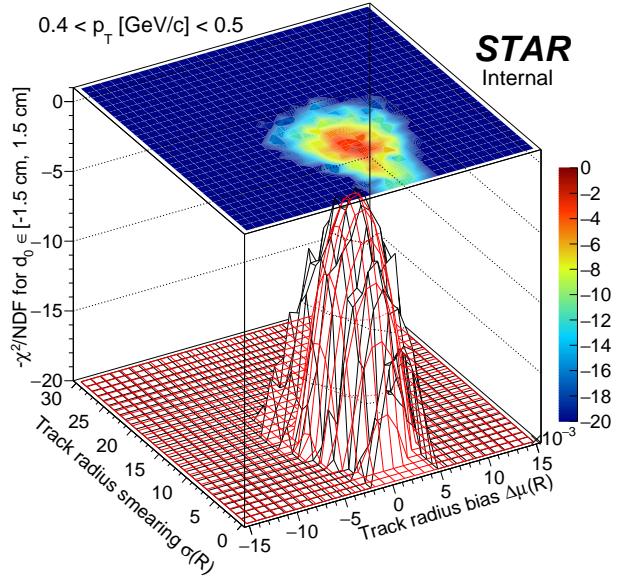


Figure 6.3: Example of $-\chi^2/\text{NDF}$ map in a parameter space in the procedure of TPC pointing resolution adjustment. The red surface represents parabola fitted in the vicinity of the global minimum.

$$\delta(\sigma(R)) = \sqrt{\delta_{\text{fit}}^2(y_0) + \frac{1}{2b\text{NDF}}}. \quad (6.5)$$

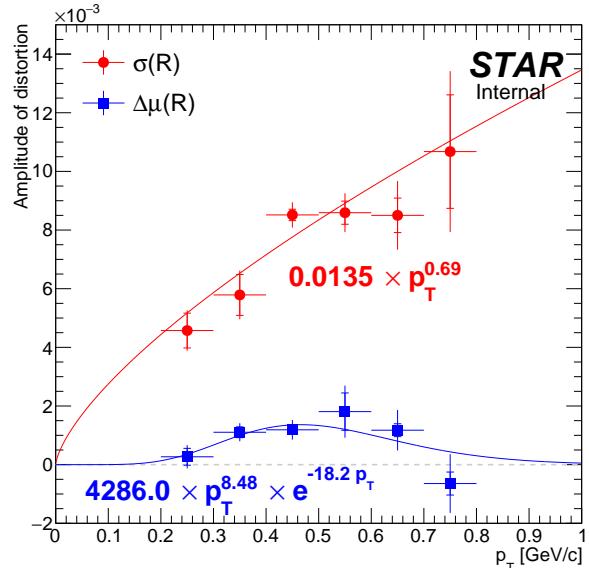
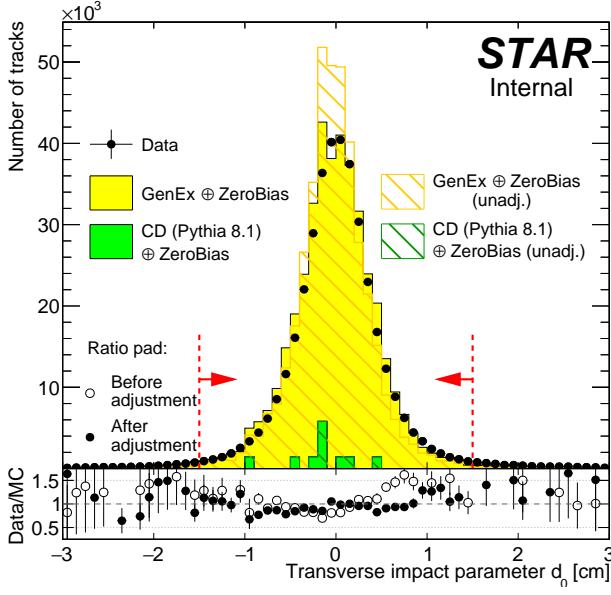
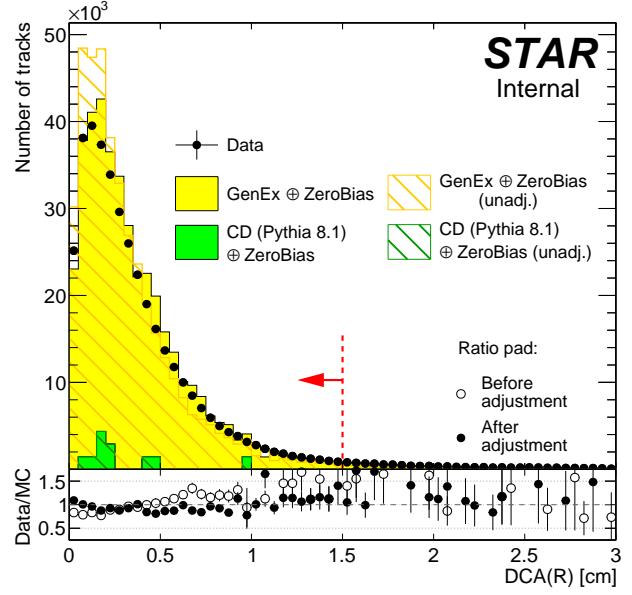


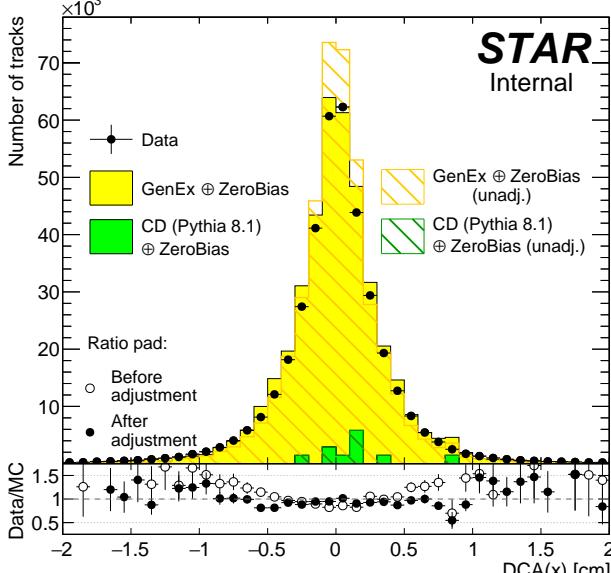
Figure 6.4: Best-fit parameters obtained in the procedure of the TPC track pointing resolution adjustment. Uncertainties on parameters resulting solely from the fit of Eq. (6.3) to $-\chi^2/\text{NDF}$ are represented by the lines with perpendicular endings. Total uncertainties (Eqs. (6.4), (6.5)) extend beyond. The empirical functions fitted to points are drawn with corresponding colors, and formula of each is written aside.



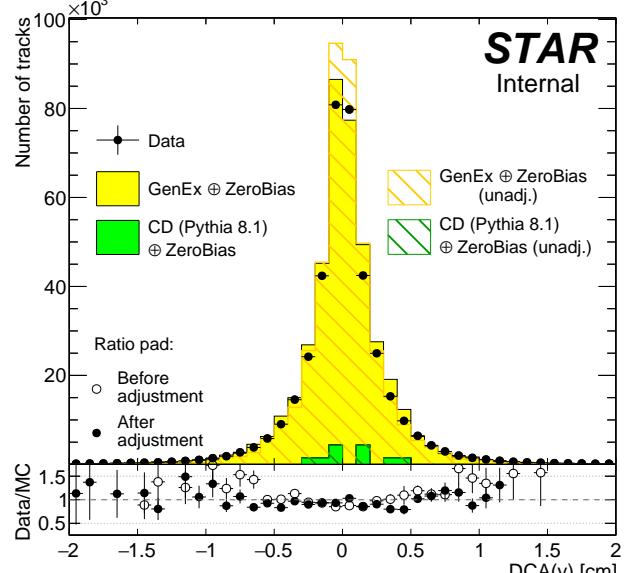
(a)



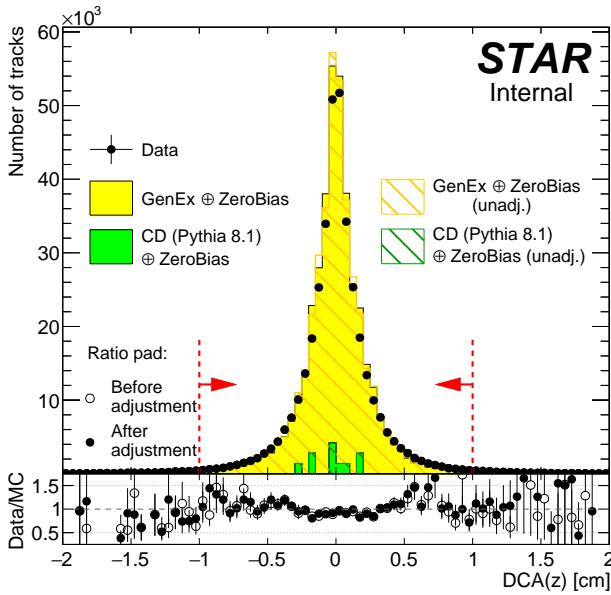
(b)



(c)



(d)



(e)

Figure 6.5: Comparison of distribution of pion transverse impact parameter d_0 (6.5a) and transverse (6.5b), x - (6.5c), y - (6.5d) and z -component (6.5e) of the DCA vector between the global helix and primary vertex in the data (CEP) and embedded MC (GenEx). Distributions for unadjusted helices are drawn as hashed histograms, while filled histograms are for adjusted helices. Normalizations of the signal and backgrounds were established from the comparison of p_T^{miss} and $\Delta\theta$ distributions after full selection (without cut on the presented quantity and without exclusivity cut), as described in Sec. XXX of Ref. [1]. Red dashed lines and red arrows indicate the range of each quantity which is accepted in analyses.

7. Dead material in front of TPC

Particle detected and reconstructed in the TPC must first pass through the detector material standing in between the accelerator vacuum and TPC gas. This affects track reconstruction efficiency, as the particle may interact with that material - in worst case inelastically, and induce secondary particles thus lower reconstruction efficiency. Accuracy of modeling of the detector material in the STAR simulation, especially in run 15 with the HFT installed, influences systematic error e.g. on the TPC track reconstruction efficiency. In this section the density of secondary vertices is compared between the data and embedded MC. The density of secondary vertices is directly proportional to the amount of the material in given volume, hence any discrepancy between secondary vertex distribution in the data and MC can be a hint for inaccuracies of the STAR simulation which should be accordingly covered by the systematic uncertainties. It should be stressed that this analysis is not aimed to tune the material budget in the STAR simulation, as there are much better data for this than high-luminosity proton-proton collisions from run 15. The aim of presented study is to obtain reasonable estimate of the component of systematic uncertainty of the TPC track reconstruction efficiency related to the error on the amount and distribution of simulated material.

Analysis of the distribution of secondary vertices was performed using both zero-bias (ZB) data and minimum-bias MC (Pythia) embedded into zero-bias triggers. Because of insufficient statistics of the ZB data, for the purpose of analysis presented in this section both standard ZB data sample (from ZB triggers in st_rp stream) and the subsample of RP_CP triggers (see Ref. [6] for trigger details) with identified elastic proton-proton scat-

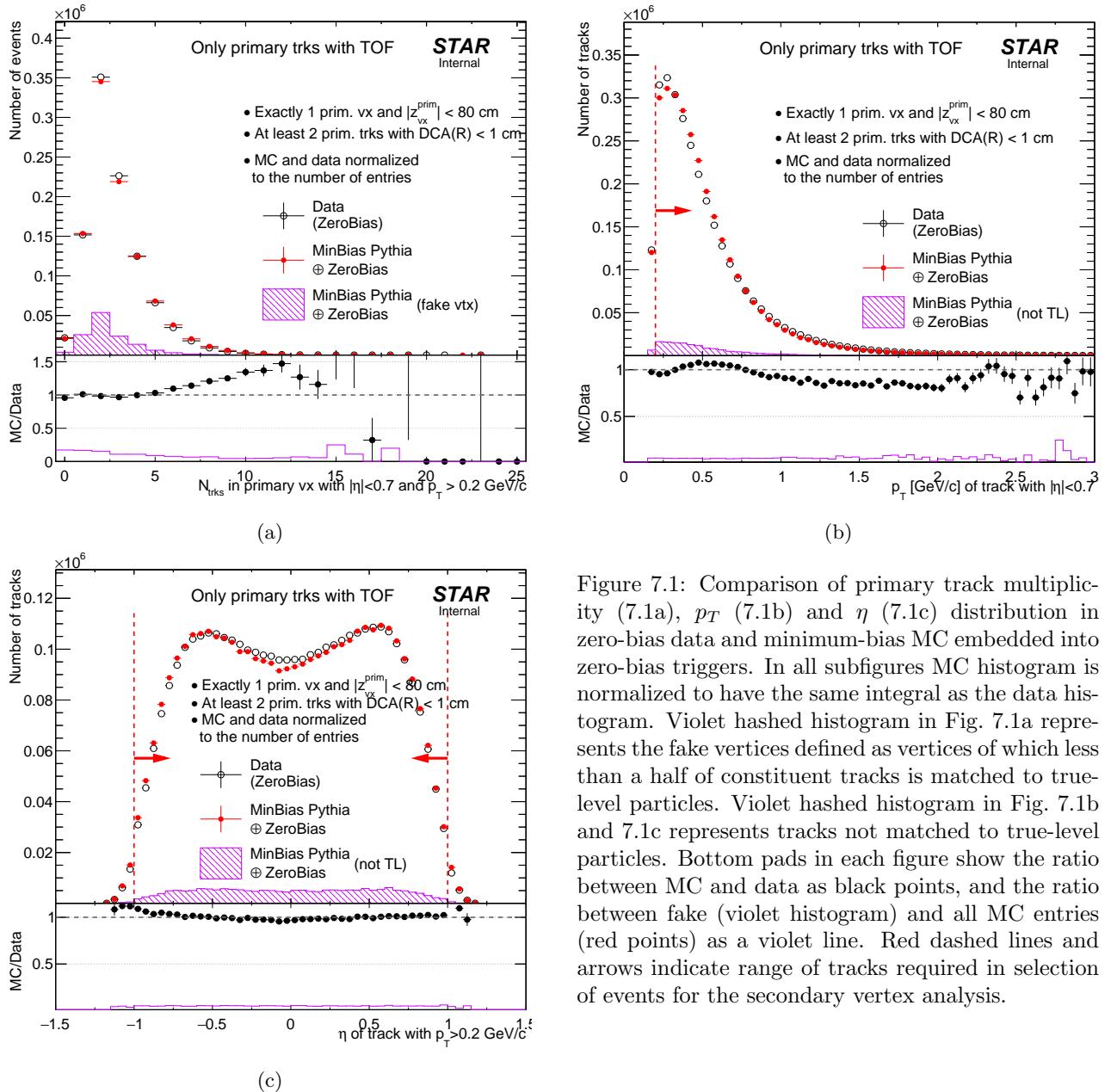


Figure 7.1: Comparison of primary track multiplicity (7.1a), p_T (7.1b) and η (7.1c) distribution in zero-bias data and minimum-bias MC embedded into zero-bias triggers. In all subfigures MC histogram is normalized to have the same integral as the data histogram. Violet hashed histogram in Fig. 7.1a represents the fake vertices defined as vertices of which less than a half of constituent tracks is matched to true-level particles. Violet hashed histogram in Fig. 7.1b and 7.1c represents tracks not matched to true-level particles. Bottom pads in each figure show the ratio between MC and data as black points, and the ratio between fake (violet histogram) and all MC entries (red points) as a violet line. Red dashed lines and arrows indicate range of tracks required in selection of events for the secondary vertex analysis.

tering events using loose RP track selection were used. The latter subsample is in good approximation a ZB sample in terms of central detector, as it was triggered only by the east and west coincidence of Roman Pots - any particles present in the TPC and TOF must be product of pile-up interaction. In all plots and later in the text we refer to this merged sample as ZB data sample.

Analysis started with the following selection of events:

1. Exactly 1 reconstructed primary vertex (with tracks matched to hits in TOF),
2. $|z_{vx}| < 80$ cm,
3. ≥ 2 prim. TOF tracks with: $DCA(R) < 1$ cm, $|\eta| < 1$, $p_T > 0.2$ GeV/c, $N_{\text{hits}}^{\text{fit}} \geq 25$, $N_{\text{hits}}^{\text{dE/dx}} \geq 15$, $N_{\text{hits}}^{\text{fit}}/N_{\text{hits}}^{\text{poss}} \geq 0.52$.

The aim of above criteria was to select pile-up-free events with well defined vertex. Cut on z -vertex is identical to one used in physics analyses. Figure 7.1 shows comparison of quantities characterizing an event. In general a moderate agreement between MC and data can be observed, considered sufficient for trustworthy result of described analysis.

As a next step the TPC tracks were selected for the search and reconstruction of secondary vertices. The requirements were as follows:

1. Global TPC tracks matched with TOF not associated with any primary TPC track,
2. $|\eta| < 0.7$, $p_T > 0.2$ GeV/c, $N_{\text{hits}}^{\text{fit}} \geq 25$, $N_{\text{hits}}^{\text{dE/dx}} \geq 15$, $N_{\text{hits}}^{\text{fit}}/N_{\text{hits}}^{\text{poss}} \geq 0.52$,
3. Distance of closest approach to the STAR z -axis ($x, y = (0, 0)$, $d_0^{(0,0)}$, larger than inner radius of the beampipe: $d_0^{(0,0)} > 2$ cm.

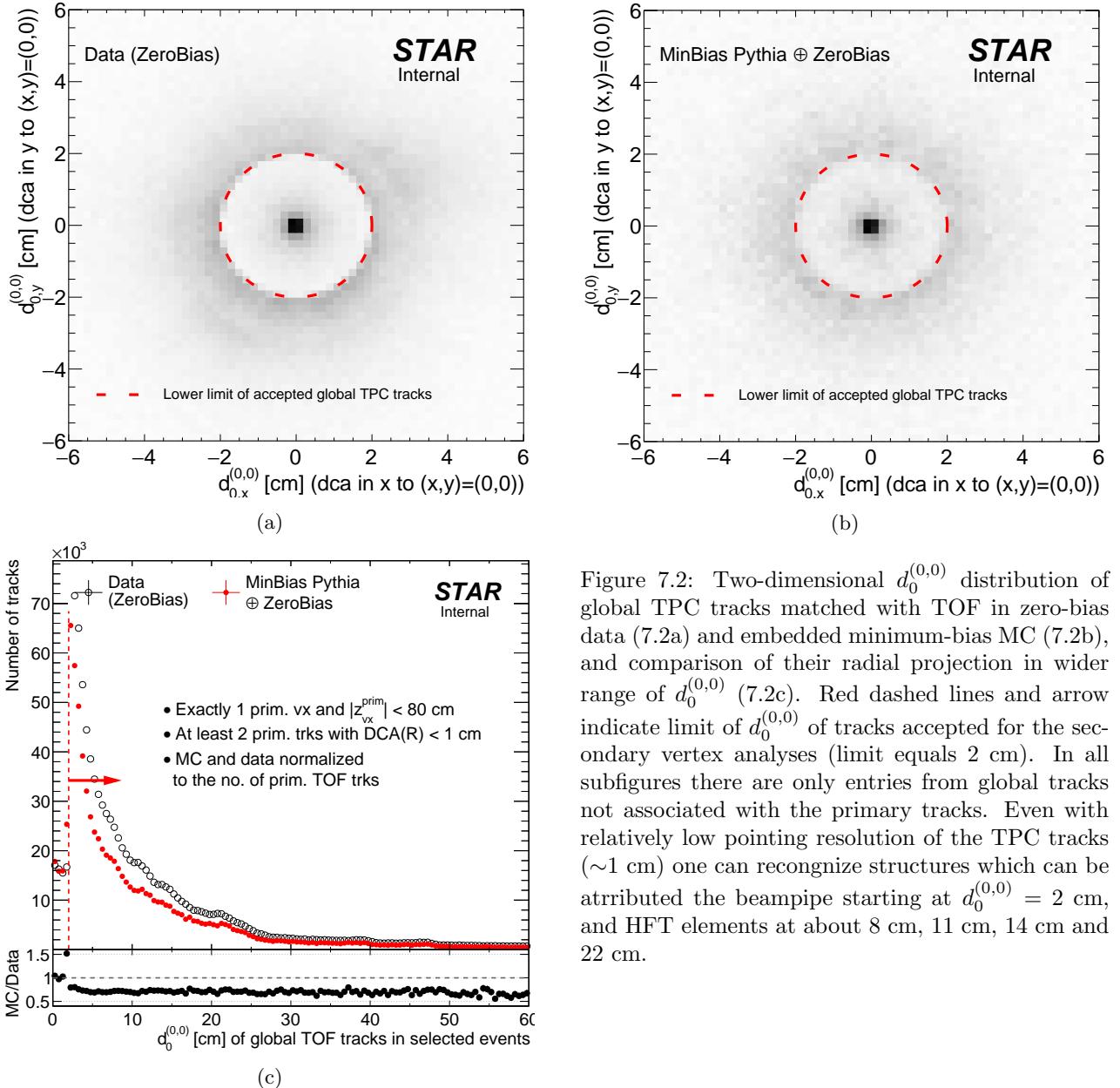


Figure 7.2: Two-dimensional $d_0^{(0,0)}$ distribution of global TPC tracks matched with TOF in zero-bias data (7.2a) and embedded minimum-bias MC (7.2b), and comparison of their radial projection in wider range of $d_0^{(0,0)}$ (7.2c). Red dashed lines and arrow indicate limit of $d_0^{(0,0)}$ of tracks accepted for the secondary vertex analyses (limit equals 2 cm). In all subfigures there are only entries from global tracks not associated with the primary tracks. Even with relatively low pointing resolution of the TPC tracks (~ 1 cm) one can recognize structures which can be attributed the beampipe starting at $d_0^{(0,0)} = 2$ cm, and HFT elements at about 8 cm, 11 cm, 14 cm and 22 cm.

These cuts were intended to select in-time TPC tracks with high chance of being a product of secondary interaction of primary particle with the detector material. The higher limit of accepted $d_0^{(0,0)}$ was set in analysis, the less background was found in the secondary vertex distribution for a price of limited access to secondary vertices of low radial distance from STAR z -axis. Cut of 2 cm was found a good compromise. In Fig. 7.2 we present comparison of $d_0^{(0,0)}$ distribution of selected global TOF-matched TPC tracks in the data and embedded MC (without cut on $d_0^{(0,0)}$). To some extent this distribution reflects the material density (secondary vertex density) in the radial direction, therefore we present it with the MC distribution normalized to the same total number of primary tracks as in the data. Number of secondary vertices is proportional to the number of primary particles, so we use such normalization to allow direct comparison of the distributions:

$$\text{MC normalization factor} = \frac{\langle N_{\text{trks/evt}}^{\text{DATA}} \rangle \times N_{\text{evts}}^{\text{DATA}}}{\langle N_{\text{trks/evt}}^{\text{MC}} \rangle \times N_{\text{evts}}^{\text{MC}}} = \frac{N_{\text{trks}}^{\text{DATA}}}{N_{\text{trks}}^{\text{MC}}} \quad (7.1)$$

Especially in Fig. 7.2c one can find structures/peaks that might be attributed to subdetectors (PXL, IST, SST) of the HFT. Notable is different yield of histograms which could indicate different amount of simulated dead material with respect to real conditions. The reason of this inconsistency was found in imperfect simulation of the pointing resolution of the TPC tracks - because the resolution is higher in the simulation, more true primary tracks are reconstructed as primary tracks hence less such tracks are accepted in the global track selection (comparing to data). This effect is accounted later in the background subtraction procedure.

After secondary track candidates were selected, the described algorithm for secondary vertex reconstruction was used:

1. Loop over all pairs of secondary track candidates, store pairs whose DCA is less than 0.5 cm (nearby tracks passing a proximity cut),
2. Link pairs of nearby tracks into sets of tracks connected by the common nearby tracks,
3. Loop over all sets defined in 2., in each set loop over all pairs from given set, reject worst-matching tracks (these with largest DCA to others) until all pairs of tracks have DCA less than 0.5 cm,
4. Based on number of tracks in secondary vertex, total charge, specific energy loss, dE/dx , cosine of the opening angle of two tracks $\cos(\theta)$ and invariant mass of two tracks m_{inv} determine if the vertex is from resonance decay or photoconversion (see Ref. [7]); if none of the two then assume hadronic vertex,
5. Calculate the vertex position as the average DCA point of all track pairs in the vertex.

As a result secondary vertices were reconstructed, whose multiplicity distribution is depicted in Fig. 7.3. Analysis was continued only with vertices of multiplicity equal 2. The first reason was that most of vertices consist of just a pair of tracks. Another reason was the background subtraction method developed only for vertices made of two tracks. In addition to this, only vertices representing primary particles in the pseudorapidity range $-0.7 < \eta < 0.7$ were analyzed. To enable such selection a variable η_{vtx} was defined, as shown in Fig. 7.4.

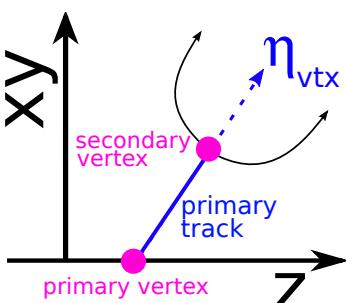


Figure 7.4: η_{vtx} definition (sketch).

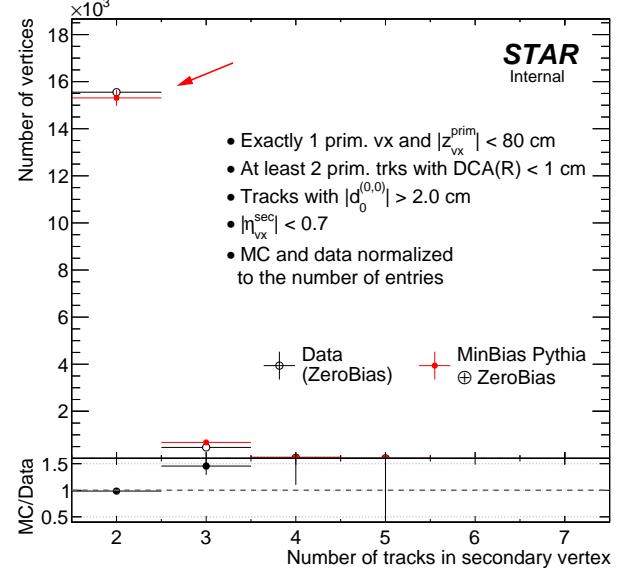
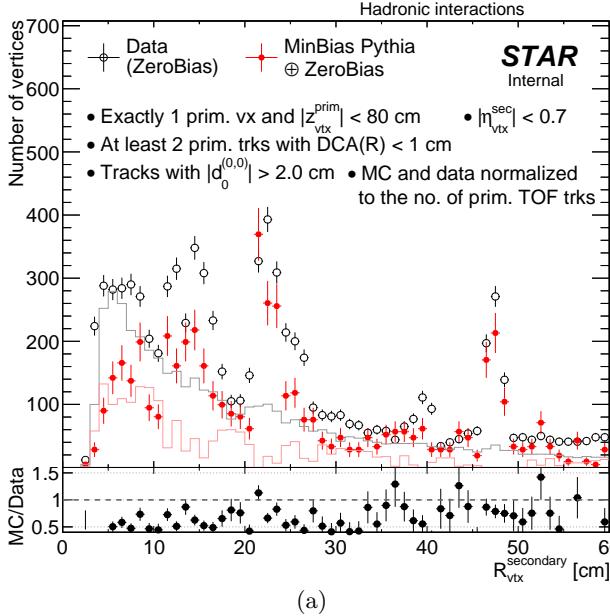


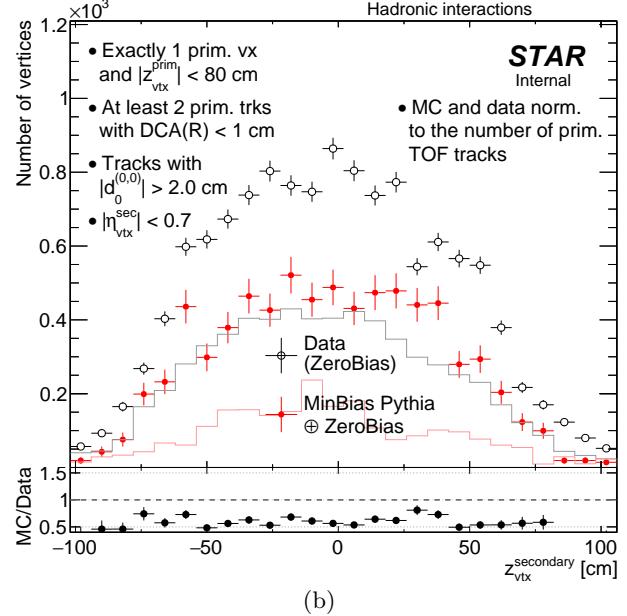
Figure 7.3: Multiplicity of tracks in reconstructed secondary vertices. Red arrow points to bin with vertices used in final analyses of vertex distribution.

Raw distributions of $R_{\text{vtx}}^{\text{secondary}}$ and $z_{\text{vtx}}^{\text{secondary}}$ are shown in Fig. 7.5a and Fig. 7.5b, respectively. In $R_{\text{vtx}}^{\text{secondary}}$ spectrum one can find peaks in the regions where the HFT subdetectors are placed. Peaks seem to lie on top of a tail whose origin has been identified with the secondary vertices made of pairs containing true primary tracks which were not associated with any primary vertex and unfortunately passed selection of global tracks for the secondary vertex reconstruction. Without this background subtraction, the ratio of MC to data varies mostly between 0.5 and 0.7. For this reason a method of estimation of the background was invented, as described in the next paragraph.

Background estimation makes use of different content of fake secondary vertices depending on the proximity cut used in the secondary vertex reconstruction. Figure 7.6 shows the percentage of background (fake pairs) distributed over the distance of closest approach between two tracks. A comment needs to be



(a)



(b)

Figure 7.5: Comparison of raw $R_{\text{vtx}}^{\text{secondary}}$ (7.5a) and $z_{\text{vtx}}^{\text{secondary}}$ (7.5b) distribution in the data (opened black circles) and embedded MC (filled red circles). Only vertices recognized as products of hadronic interactions are shown in the figure. Solid lines denote estimated background content in the distribution of corresponding color.

made that the agreement of the shape of the tails in data and MC distributions was achieved only after the adjustment of the TPC resolution in MC, as described in Sec. 6. This agreement allows to believe in proper description of the data by MC in terms of background distribution over DCA of two global tracks, which is used in the background estimation.

It agrees with intuition that the most optimal cut to select pairs from the secondary vertices is as low as about 0.5 cm, however one can select sample with slightly different ratio of signal to background if the proximity cut is changed to accept tracks of DCA within some higher limits. In Fig. 7.6 the nominal proximity cut is

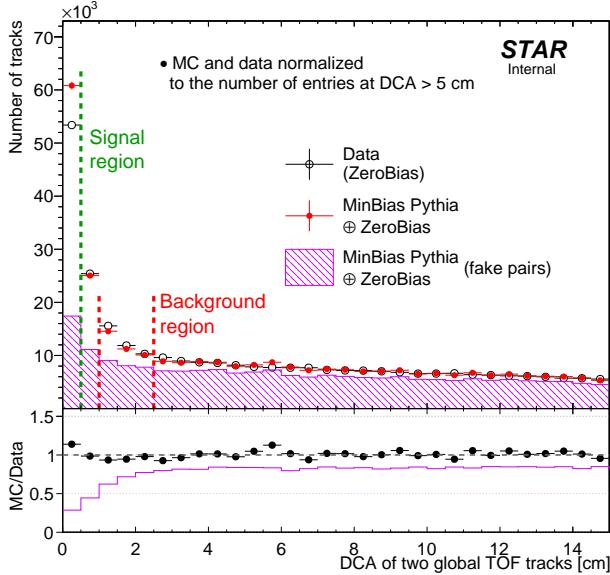


Figure 7.6: Comparison of DCA between all pairs of secondary track candidates selected for the secondary vertex reconstruction in the data and embedded MC. MC histogram is normalized to the data at $\text{DCA} > 5 \text{ cm}$. Violet hashed histogram depicts pairs contained in MC histogram and not originating from the same vertex. Solid violet line in the lower pad denotes ratio of violet and red histogram.

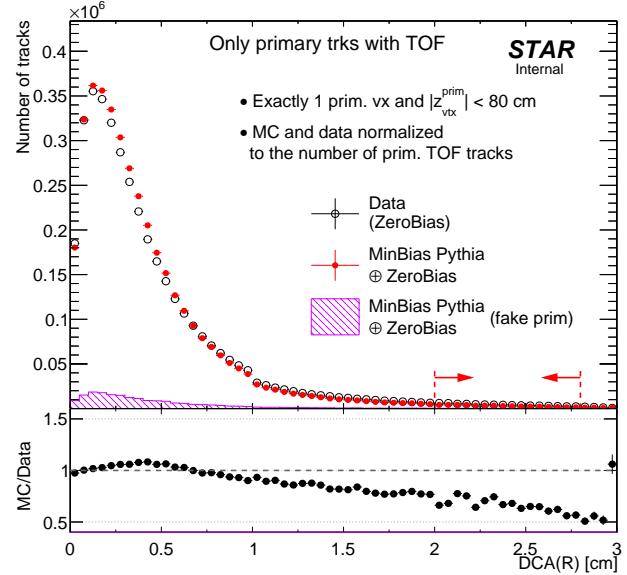


Figure 7.7: Comparison of radial DCA of all primary tracks matched with TOF and passing quality criteria in events selected for secondary vertex analysis, between the data and embedded MC. Violet hashed histogram represents tracks not matched to true-level particles. Red dashed lines and arrows limit region used to find normalization that compensates different background yield in reconstructed secondary vertex distributions in data and embedded MC.

marked with the green line (signal region), while the modified proximity cut is marked with red lines (background region). With such two versions of cuts used in vertexing the two independent distributions of secondary vertices can be obtained: one with the standard proximity cut - \mathcal{H}_1 , the other with modified proximity cut, in our case $1.0 \text{ cm} < \text{DCA} < 2.5 \text{ cm}$ - \mathcal{H}_2 . Limits in modified proximity cut were set to such values in order to ensure enough statistics as well as provide satisfactory resolution of secondary vertex position calculated as a middle point between DCA points on helices associated with the tracks. One can note that the content of histograms can be described by the set of equations given below:

$$\left\{ \begin{array}{l} \mathcal{H}_1 = (1 - B) \times \text{signal} + B \times \text{background}, \\ \mathcal{H}_2 = (1 - B') \times \text{signal} + B' \times \text{background}, \end{array} \right. \quad (7.2)$$

$$\left\{ \begin{array}{l} \text{signal} = \frac{B' \times \mathcal{H}_1 - B \times \mathcal{H}_2}{B' - B} \\ \text{background} = \frac{(1 - B) \times \mathcal{H}_2 - (1 - B') \times \mathcal{H}_1}{B' - B} \end{array} \right. \quad (7.4)$$

in which parameters B and B' denote the background fraction in the distribution resultant from analysis utilizing nominal and modified proximity cut, respectively. The solution to set of Eqs. (7.2), (7.4) is the following:

$$\left\{ \begin{array}{l} \text{signal} = \frac{B' \times \mathcal{H}_1 - B \times \mathcal{H}_2}{B' - B} \\ \text{background} = \frac{(1 - B) \times \mathcal{H}_2 - (1 - B') \times \mathcal{H}_1}{B' - B} \end{array} \right. \quad (7.5)$$

An important remark here is that the background fraction extracted from the ratio of violet and red histograms in Fig. 7.6 can be used directly in Eqs. (7.2)-(7.5) only for background estimation in MC. In case of background estimation in data parameters B and B' have to be corrected for the leakage of true primary tracks to set of selected secondary track candidates, as it was described in one of preceding paragraphs. The correction factor κ is extracted from the ratio of the radial DCA of the primary TPC tracks in events selected for the secondary vertex study (Fig. 7.7). Histogram range selected for calculation of the ratio was set to $2.0 \text{ cm} < \text{DCA}(R) < 2.6 \text{ cm}$, as this range coincides with the $d_0^{(0,0)}$ of global tracks accepted for the analysis. κ calculated in this range equals 1.48. Variation of value of κ with changed limits of $\text{DCA}(R)$ selected for the ratio calculation do not influence significantly the final result. The correction is done by multiplying fraction B and B' by κ only when estimating the background in the data.

Background determined with the described method is shown in Fig. 7.1 with the solid lines colored according to corresponding markers. This background was subtracted and final, background-free distributions of the secondary vertex positions in the transverse and longitudinal direction are presented in Fig. 7.8. Most relevant region - the HFT detector extending between $\sim 2 \text{ cm}$ and $\sim 30 \text{ cm}$ is satisfactorily well described by MC. Also, the inner wall of the TPC at $\sim 48 \text{ cm}$ well matches between data and MC.

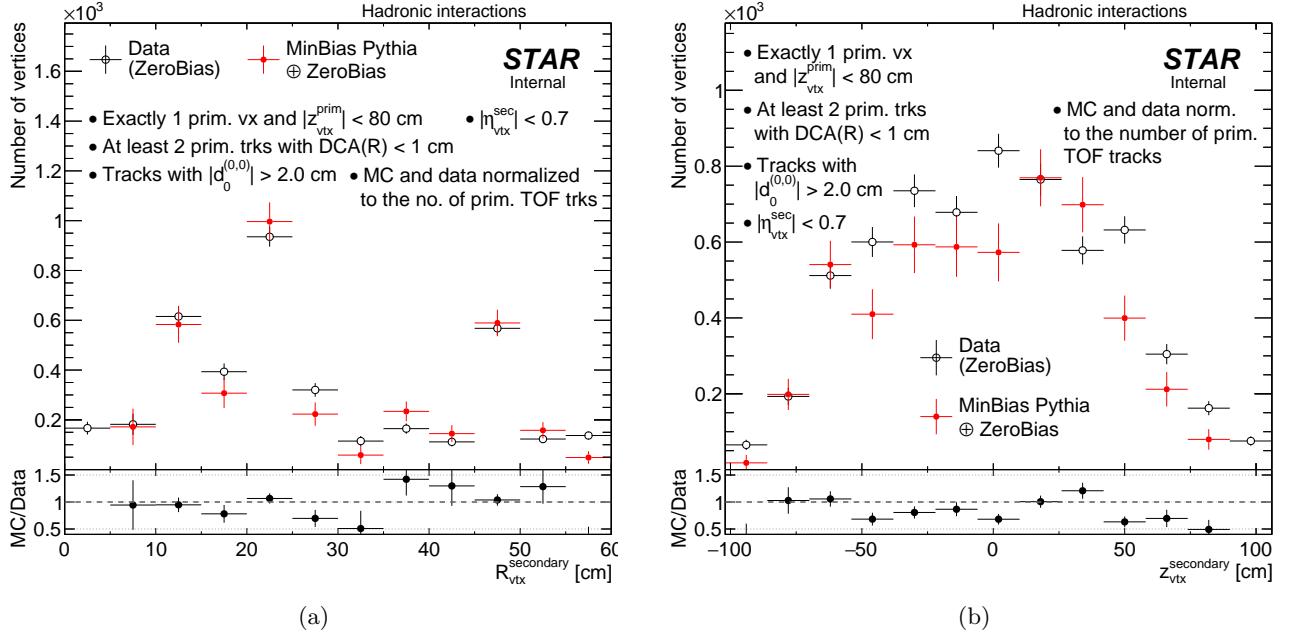


Figure 7.8: Comparison of background-subtracted $R_{\text{vtx}}^{\text{secondary}}$ (7.8a) and $z_{\text{vtx}}^{\text{secondary}}$ (7.8b) distribution in the data (open black circles) and embedded MC (filled red circles). Only vertices recognized as products of hadronic interactions are shown in the figure.

8. Systematic errors

8.1 TPC track reconstruction efficiency

8.1.1 Embedding (pile-up) effect

One major difference between simulation and real data is the presence of pile-up events. The average number of pile-up tracks in a triggering event is proportional to the BBC coincidence rate. It is expected that the difference between simulation and real data drops at lower BBC rates, and the effects of pile-up tracks could be much reduced by fitting the tracking efficiency as a function of BBC rate and using the extrapolated value at zero luminosity to compare with simulation.

The embedded MC was divided into two samples due to mean BBC_AND rate: $\langle \text{BBC_AND} \rangle = 700 \text{ kHz}$ and $\langle \text{BBC_AND} \rangle = 1400 \text{ kHz}$. Next, the track reconstruction efficiency was calculated for those two samples and no-pile-up MC corresponding to them. The difference between TPC track reconstruction efficiencies for pile-up and no-pile-up MCs was calculated as:

$$\Delta\epsilon_{TPC}^{1400/700 \text{ kHz}} = \frac{N_{reco}^{no-pile-up} - N_{reco}^{pile-up}}{N_{gen}} \quad (8.1)$$

where:

N_{gen} -number of MC tracks,

$N_{reco}^{no-pile-up}$ - number of reconstructed tracks, matched with MC tracks in no-pile-up MC,

$N_{reco}^{pile-up}$ - number of reconstructed tracks, matched with MC tracks in pile-up MC.

The difference between high and low pile-up runs is given by:

$$\Delta\epsilon_{TPC} = \Delta\epsilon_{TPC}^{1400 \text{ kHz}} - 2 \cdot \Delta\epsilon_{TPC}^{700 \text{ kHz}} \quad (8.2)$$

Finally, above difference, shown in Figs. 8.2 and 8.3 for π^\pm , varies between 2 – 3% and was taken as systematic uncertainty related to TPC track reconstruction efficiency.

Figure 8.2: π^\pm TPC track reconstruction efficiency as a function of p_T ($|\eta| < 0.7$, $|V_z| < 80 \text{ cm}$) for embedded MC samples with $\langle \text{BBC_AND} \rangle = 700 \text{ kHz}$ and $\langle \text{BBC_AND} \rangle = 1400 \text{ kHz}$. The efficiencies from corresponding no-pile-up MC samples were also shown. Additionally, the differences from Eq. 8.2 were drawn in the bottom of each plot.

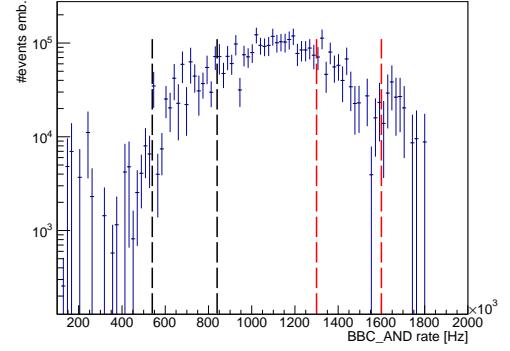
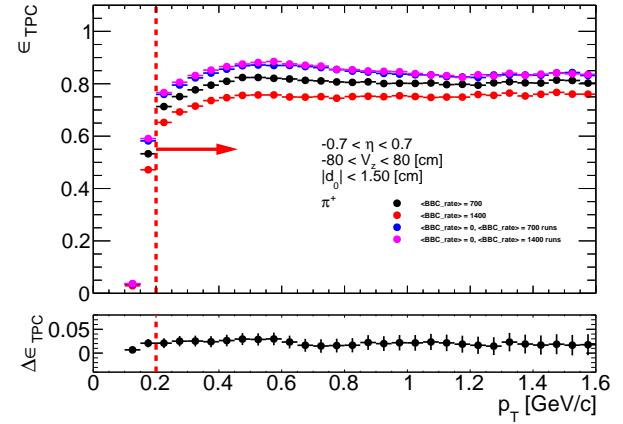
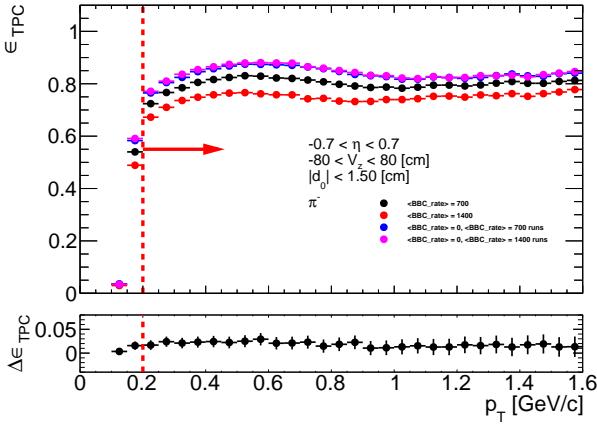
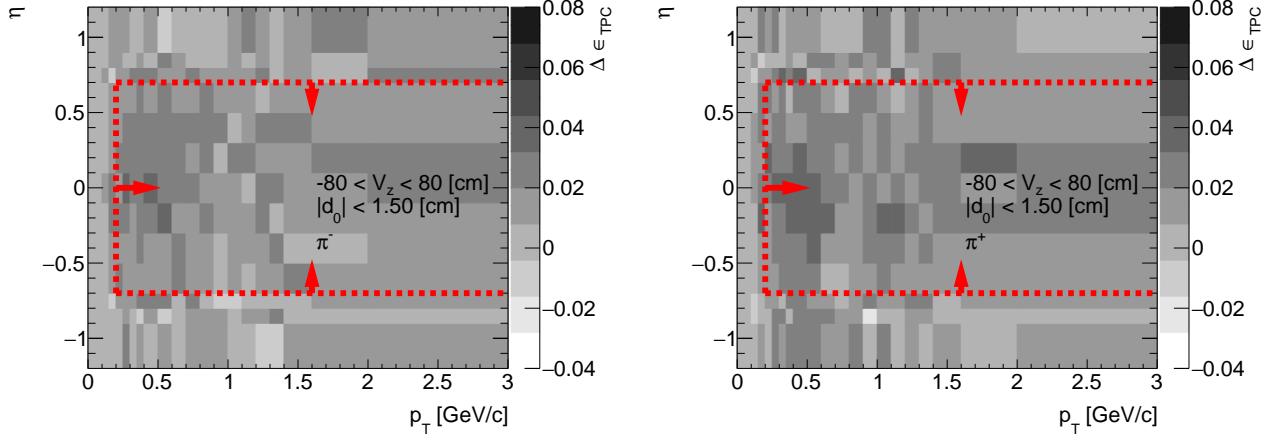


Figure 8.1: Number of events in embedded MC as a function of BBC_AND rate. The black and red lines represent the events with $\langle \text{BBC_AND} \rangle = 700 \text{ kHz}$ and $\langle \text{BBC_AND} \rangle = 1400 \text{ kHz}$, respectively.

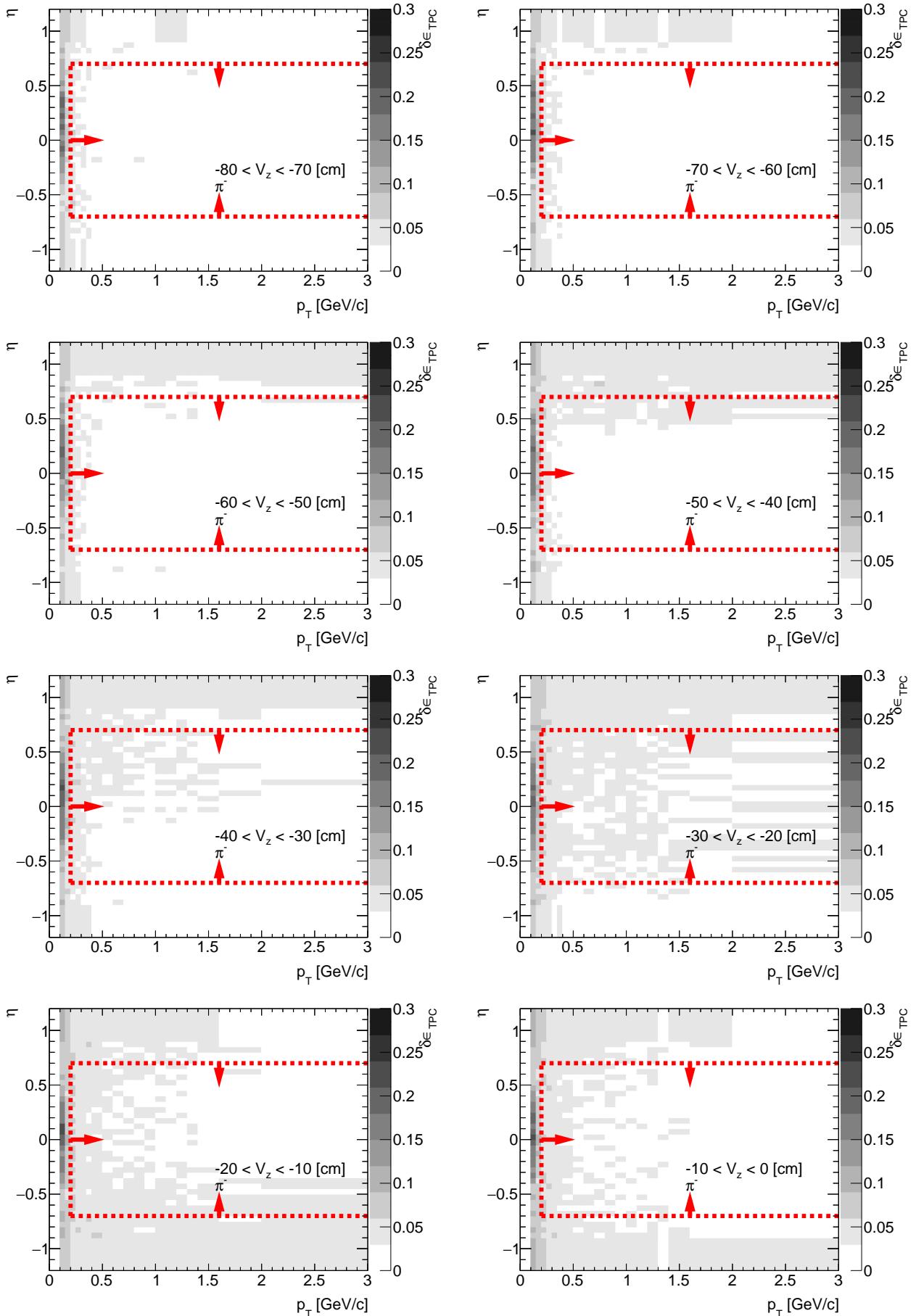
Figure 8.3: The difference $\Delta\epsilon_{TPC} = \Delta\epsilon_{TPC}^{1400 \text{ kHz}} - 2 \cdot \Delta\epsilon_{TPC}^{700 \text{ kHz}}$ for π^\pm as a function of p_T and η ($|V_z| < 80 \text{ cm}$).



8.1.2 Dead material effect on TPC track reconstruction efficiency

The amount of dead material in front of TPC differs up to 20% between data and simulation (see Sec. 7). First, the amount of lost particles, $\delta\epsilon_{TPC}$, due to the interaction with dead material in front of TPC was estimated using no-pile-up MC samples. The results for π^- in CD are shown in Fig. 8.4. Then the symmetric systematic uncertainty to the TPC track reconstruction efficiency due to dead material was introduced as $\pm 0.2 \cdot \delta\epsilon_{TPC}$. In Fig. 8.5 the systematic uncertainty is shown for each particle species in CD as a function of p_T ($|\eta| < 0.7, |V_z| < 80 \text{ cm}$). The results for other particles and SD are shown in Figs. in Appendix D.

Figure 8.4: The amount of lost π^- due to the interaction with dead material in front of TPC. Each plot represents the fraction of lost π^- , $\delta\epsilon_{TPC}$ (z-axis), as a function of true particle pseudorapidity η (y-axis) and transverse momentum p_T (x-axis) in single z -vertex bin.



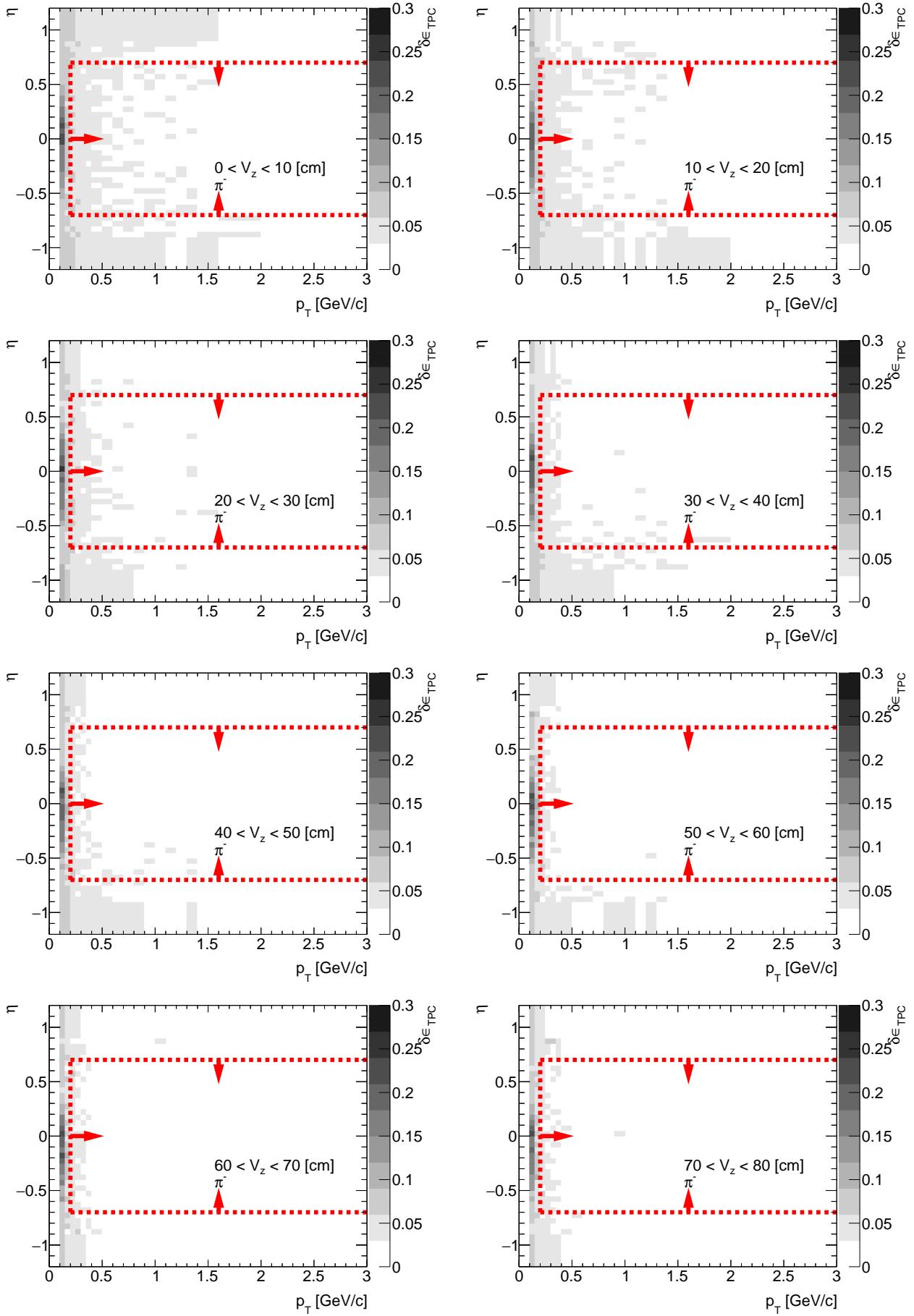
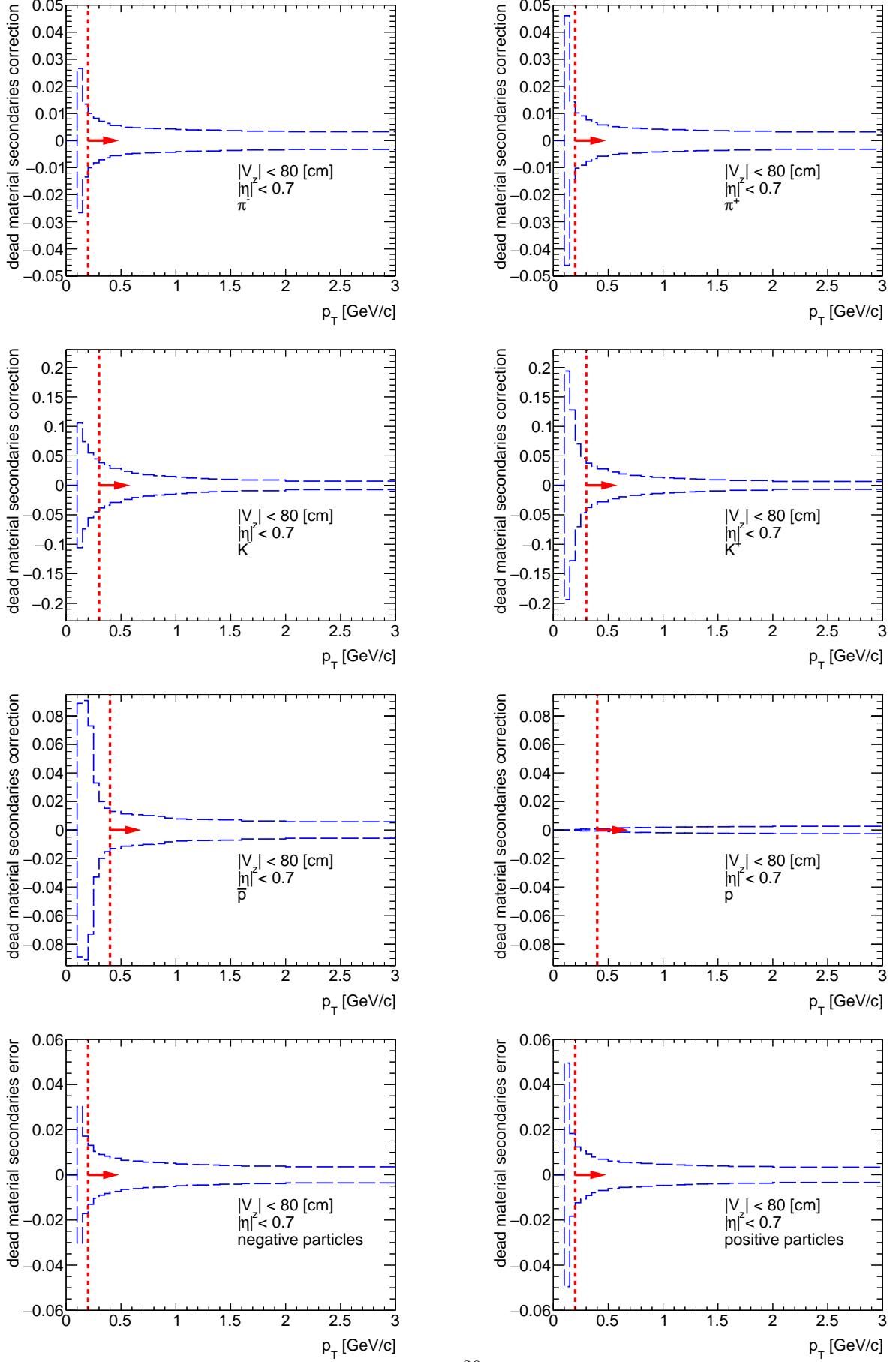


Figure 8.5: The systematic uncertainty to the TPC track reconstruction efficiency due to amount of dead material in front of TPC using MC samples for CD. Each plot represents the systematic uncertainty as a function of true particle p_T ($|\eta| < 0.7, |V_z| < 80 \text{ cm}$) for given particle species: $\pi^-, \pi^+, K^-, K^+, \bar{p}$ and p . It was also calculated for negative and positive particles without identification.



8.2 TOF matching efficiency

8.2.1 Embedding (pile-up) effect

The approach to calculate the systematic uncertainty on TOF matching efficiency related to pile-up was quite similar to the one used for TPC track reconstruction efficiency (Sec. 8.1.1). However, the TOF matching efficiency is conditional and depends on TPC track reconstruction efficiency. Since that, the difference between high and low pile-up runs is given by:

$$\Delta\epsilon_{TOF}^{1400/700 \text{ kHz}} = \frac{N_{TPC-TOF}^{no-pile-up}}{N_{TPC}^{no-pile-up}} - \frac{N_{TPC-TOF}^{pile-up}}{N_{TPC}^{pile-up}} \quad (8.3)$$

where:

$N_{TPC-TOF}^{pile-up}$ - number of reconstructed tracks, matched with MC tracks and TOF hit in pile-up MC,

$N_{TPC-TOF}^{no-pile-up}$ - number of reconstructed tracks, matched with MC tracks and TOF hit in no-pile-up MC,

$N_{TPC}^{pile-up}$ - number of reconstructed tracks, matched with MC tracks in pile-up MC,

$N_{TPC}^{no-pile-up}$ - number of reconstructed tracks, matched with MC tracks in no-pile-up MC.

Next the difference between high and low pile-up events was calculated with the formula similar to the one given by Eq. 8.2 and is shown in Figs. 8.6 and 8.7. The origin of $N_{TPC-TOF}$ increase is not known (it may be due to lack of pile-up in TPC or TOF). Since that, it is impossible to correctly calculate the statistical error for $\Delta\epsilon_{TOF}$. Nevertheless, $\Delta\epsilon_{TOF}$ is smaller than 0.5% and can be neglected in comparison with other systematic uncertainties.

Figure 8.6: π^\pm TOF matching efficiency as a function of p_T ($|\eta| < 0.7$, $|V_z| < 80$ cm) for embedded MC samples with $\langle BBC_AND \rangle \geq 700$ kHz and $\langle BBC_AND \rangle \geq 1400$ kHz. The efficiencies from corresponding no-pile-up MC samples were also shown. Additionally, the differences from Eq. 8.2 were drawn in the bottom of each plot.

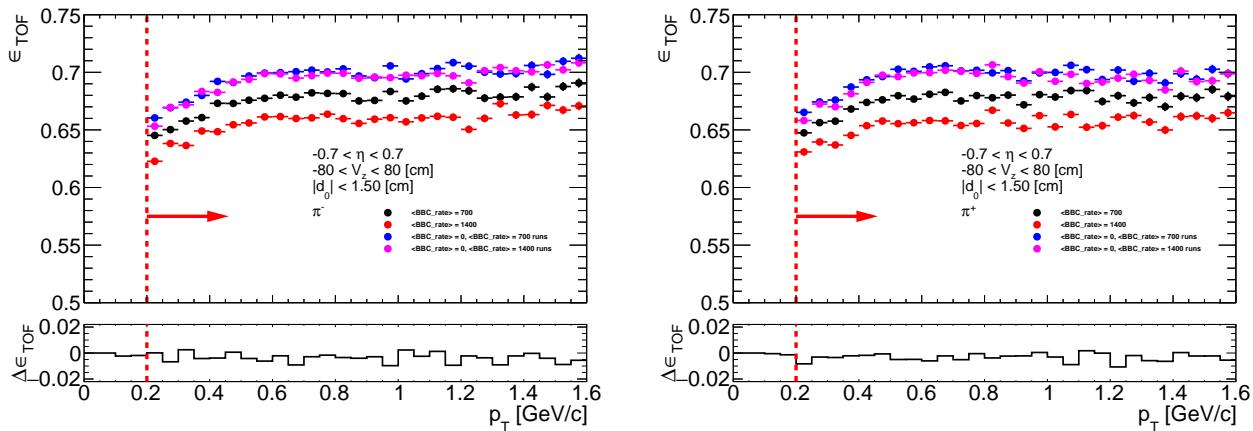
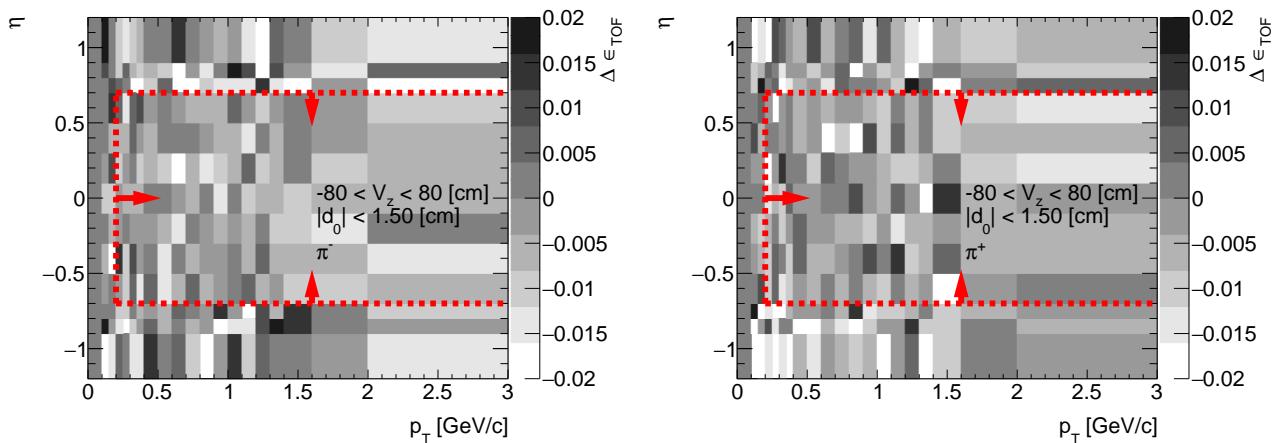


Figure 8.7: The difference $\Delta\epsilon_{TOF} = \Delta\epsilon_{TOF}^{1400 \text{ kHz}} - 2 \cdot \Delta\epsilon_{TOF}^{700 \text{ kHz}}$ for π^\pm as a function of p_T and η ($|V_z| < 80$ cm).



8.2.2 Simulation accuracy (absolute error on TOF efficiency)

The efficiency of TOF hit reconstruction and matching with the TPC tracks that was used in our analyses was taken directly from the STAR simulation. This makes the inaccuracies in the description of real detector geometry and its response propagating to physics results and introducing a bias. We decided to estimate the systematic uncertainty connected with the accuracy of description of the TOF system in (data-embedded) STAR simulation by extracting the TOF efficiency in the very same way from the data and embedded MC and comparing the results.

Unfortunately, in 2015 there were no low-luminosity (heavy-ion) runs that would imply lack of off-time pile-up tracks in TPC and thus would allow calculating the TOF efficiency in straightforward way, namely by dividing number of selected TPC tracks that were matched with TOF hits by number of all selected TPC tracks (matched or unmatched). For this reason a variation of the “tag and probe” method was developed and used.

The mentioned tag&probe method made use of the CEP of $\pi^+\pi^-$ events. In short, events with forward proton track on each side of STAR and with a TOF-matched primary TPC track (tag) were selected. Among the remaining primary TPC tracks in the same vertex the opposite-sign TPC track (probe) was chosen as the one which provides the minimum total transverse momentum of all four tracks. The probe was checked whether it has been matched with the TOF hit or not. The ratio of the matched TPC tracks to all (matched or unmatched) TPC tracks defined the TOF efficiency. The method is illustrated in the sketch in Fig. 8.8. The detailed procedure of the TOF efficiency extraction is listed below:

1. Data from RP_CPT2 triggers were used. Embedded $\pi^+\pi^-$ GenEx MC sample was subjected to the same trigger conditions. Since there is no simulation of the TOF trigger we used in MC analysis L0 TOF multiplicity reconstructed from the offline hits according to TOF trigger system description (see Ref. [8]). In Fig. 8.9 we show performance of implemented method of L0 TOF multiplicity reconstruction. One can find satisfactory correlation of the true multiplicity and the reconstructed one.
2. Events were selected with the following set of cuts from nominal CEP analysis (Ref. [1]): **C1**, **C2** (TPC vertex), **C4** (RP tracks), **C5** (TPC-RP vertex matching), **C6** (veto signal in large BBC), providing significant reduction of background events.
- 3.

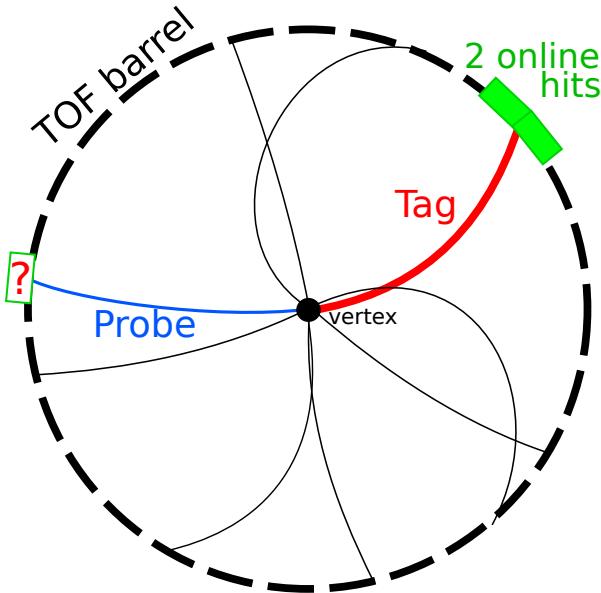


Figure 8.8: Sketch of the cross section of the central detector and CEP event with off-time pile-up tracks with drafted tag&probe method used to determine the TOF hit reconstruction and matching efficiency.

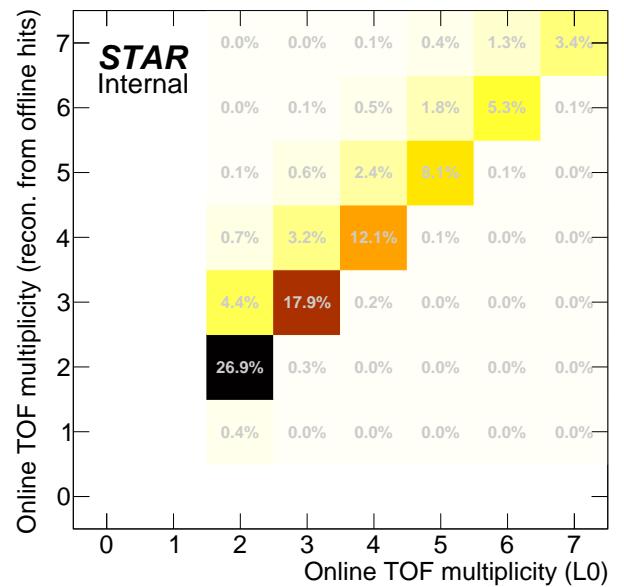


Figure 8.9: Comparison of original online TOF multiplicity at L0 (horizontal axis) and TOF multiplicity reconstructed from the offline hits (vertical axis) for preselected RP_CPT2 triggers (at least 1 primary vertex with TOF-matched TPC tracks).

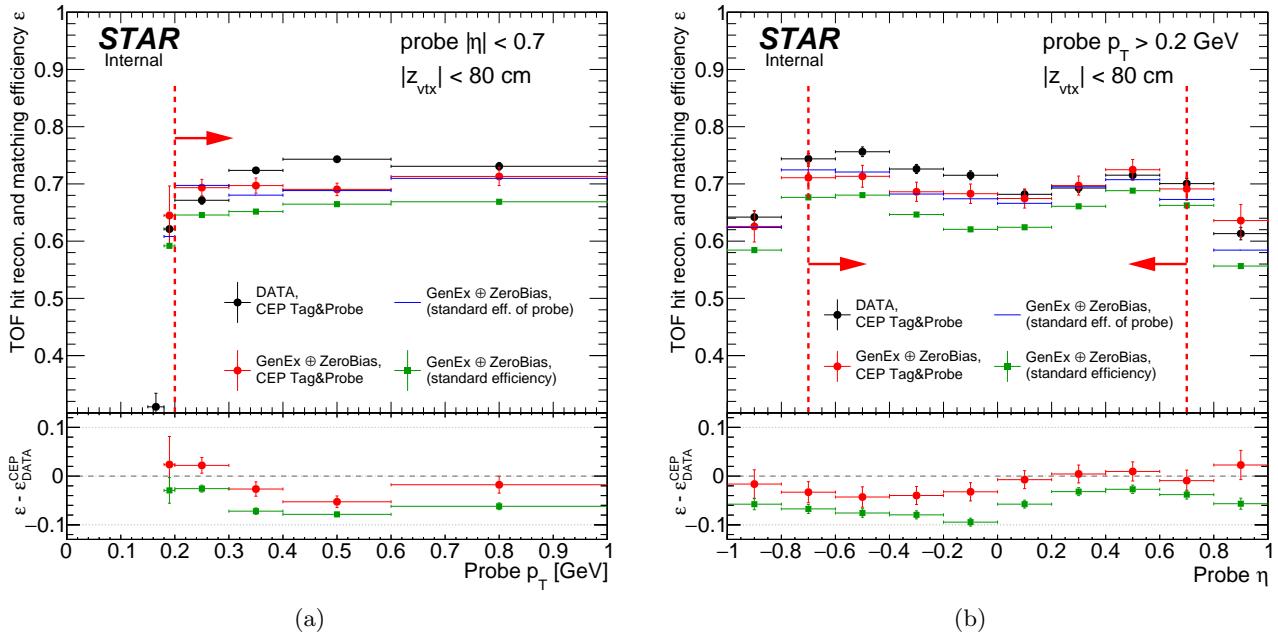


Figure 8.10: Tag&Probe TOF efficiency from CEP data (black points) compared with the result from embedded CEP MC (red points) as a function of TPC track p_T (8.10a) and η (8.10b). Green points represent the TOF efficiency calculated in the standard way from embedded CEP MC sample. Blue lines denote the TOF efficiency calculated in the standard way solely from the selected probe tracks that were matched to primary pions at the true level. Difference between red points and blue lines show the potential bias of tag and probe method.

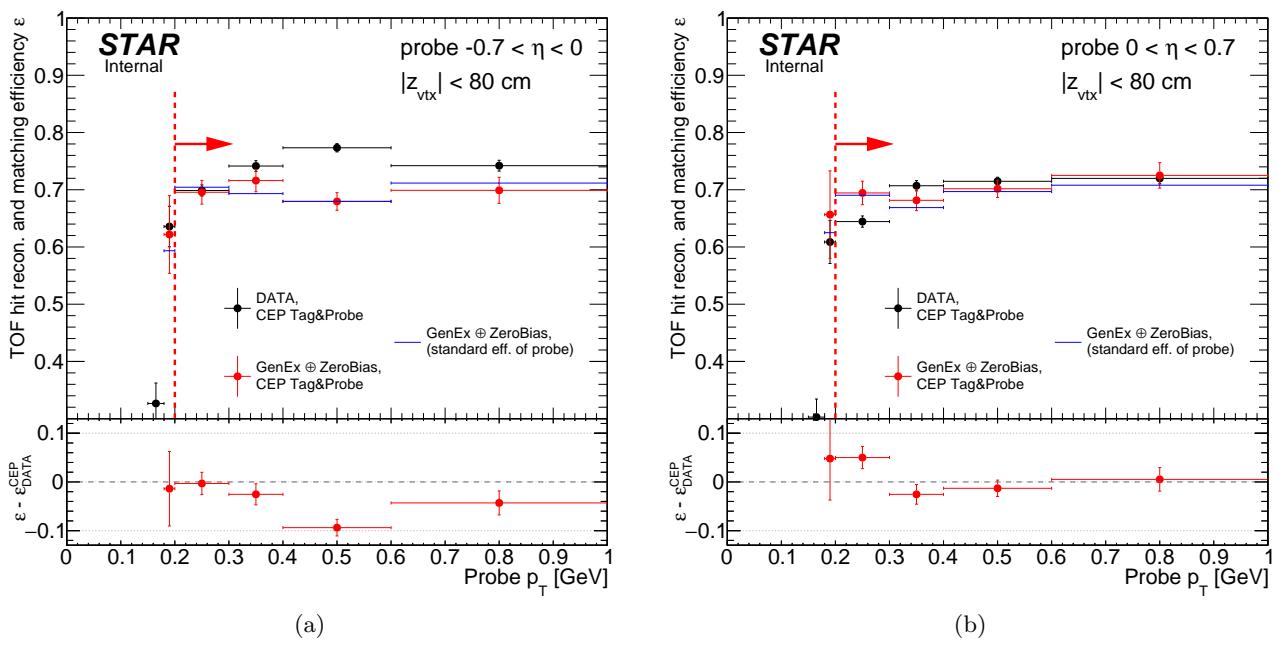
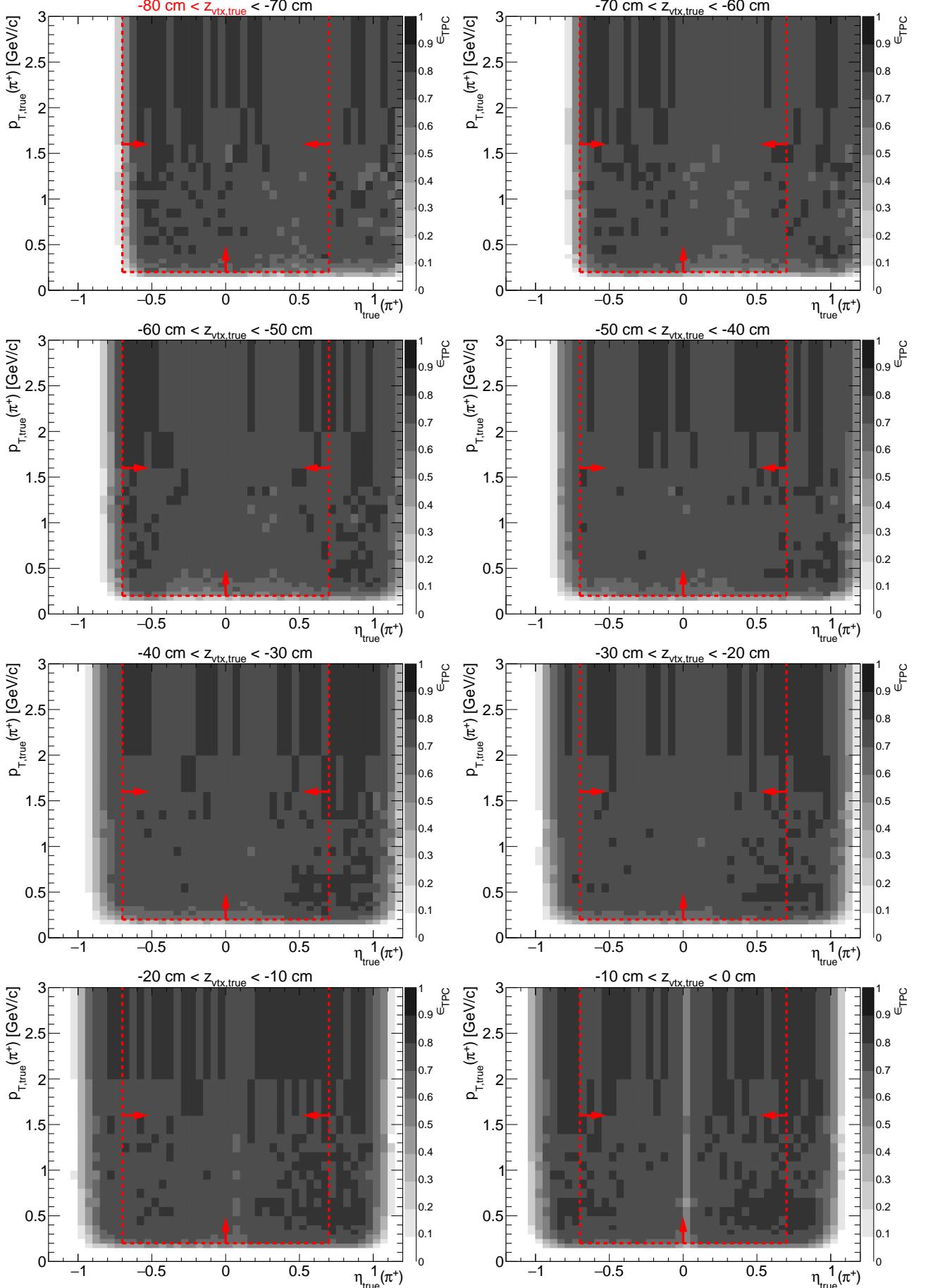


Figure 8.11: .

Appendix A

TPC track reconstruction efficiency

Figure A.1: TPC acceptance and reconstruction efficiency of π^+ . Each plot represents the TPC efficiency ϵ_{TPC} (z -axis) as a function of true particle pseudorapidity η (x -axis) and transverse momentum p_T (y -axis) in single z -vertex bin whose range is given at the top. Red lines and arrows indicate region accepted in analyses.



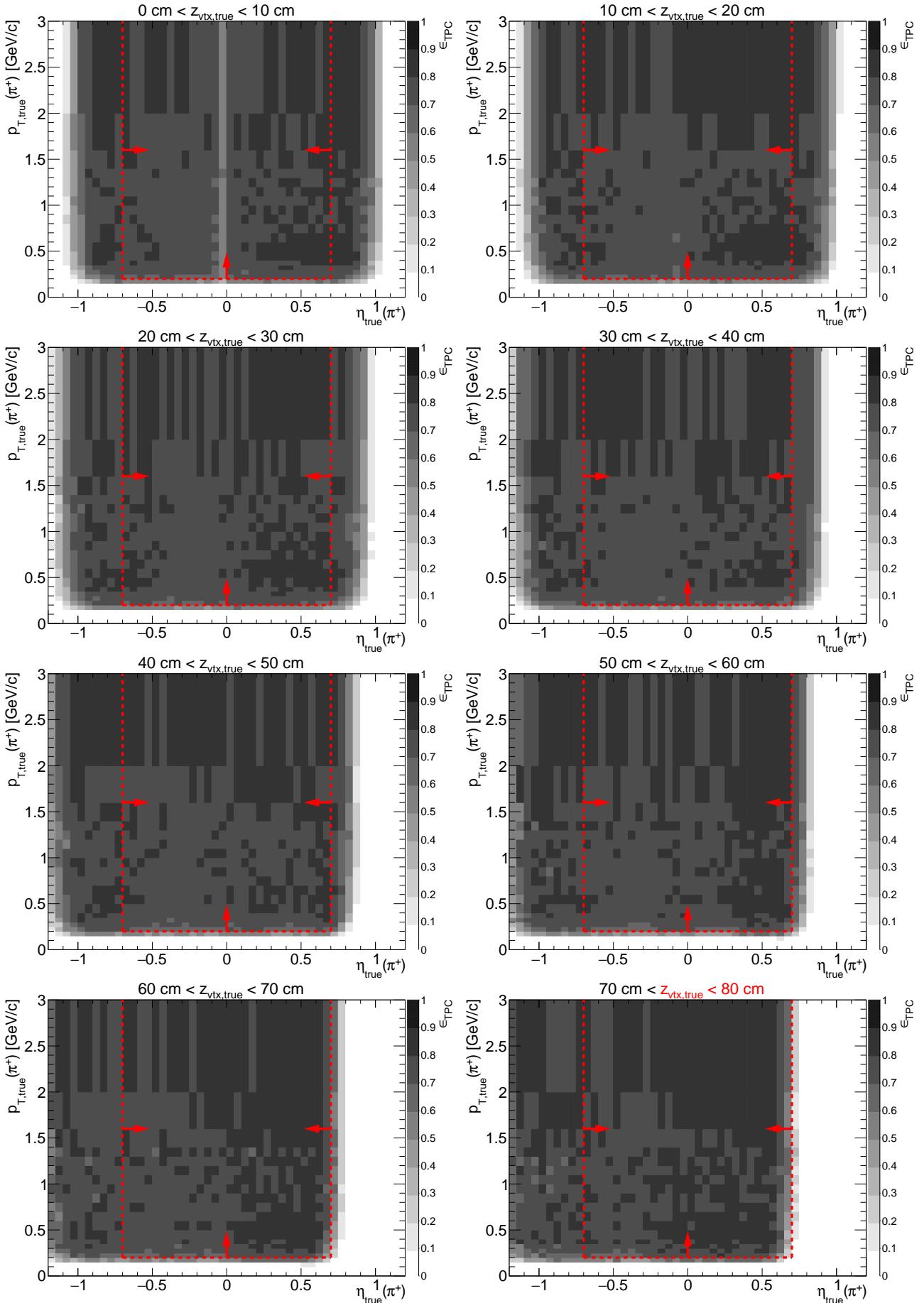
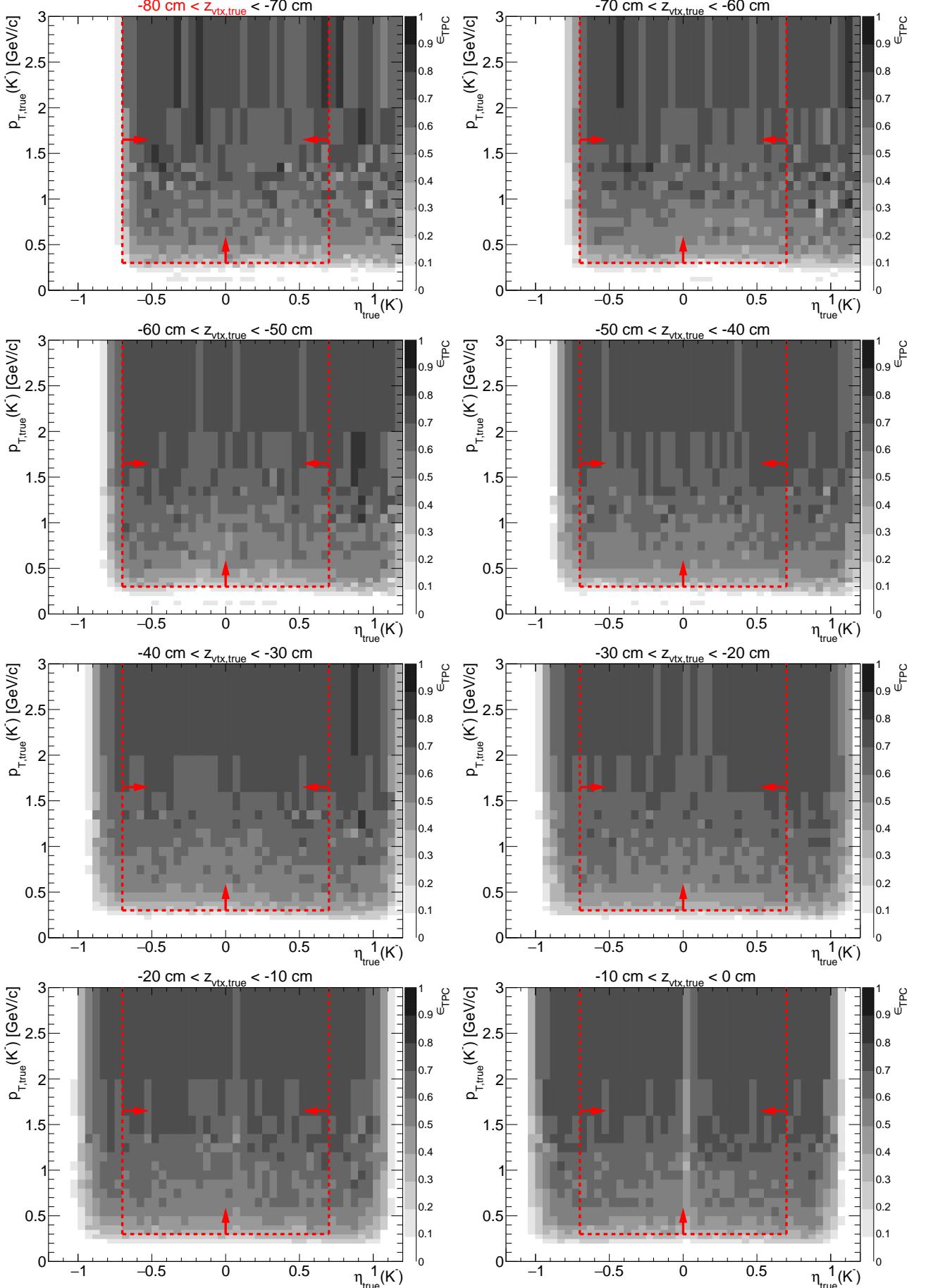


Figure A.2: TPC acceptance and reconstruction efficiency of K^- . Each plot represents the TPC efficiency ϵ_{TPC} (z -axis) as a function of true particle pseudorapidity η (x -axis) and transverse momentum p_T (y -axis) in single z -vertex bin whose range is given at the top. Red lines and arrows indicate region accepted in analyses.



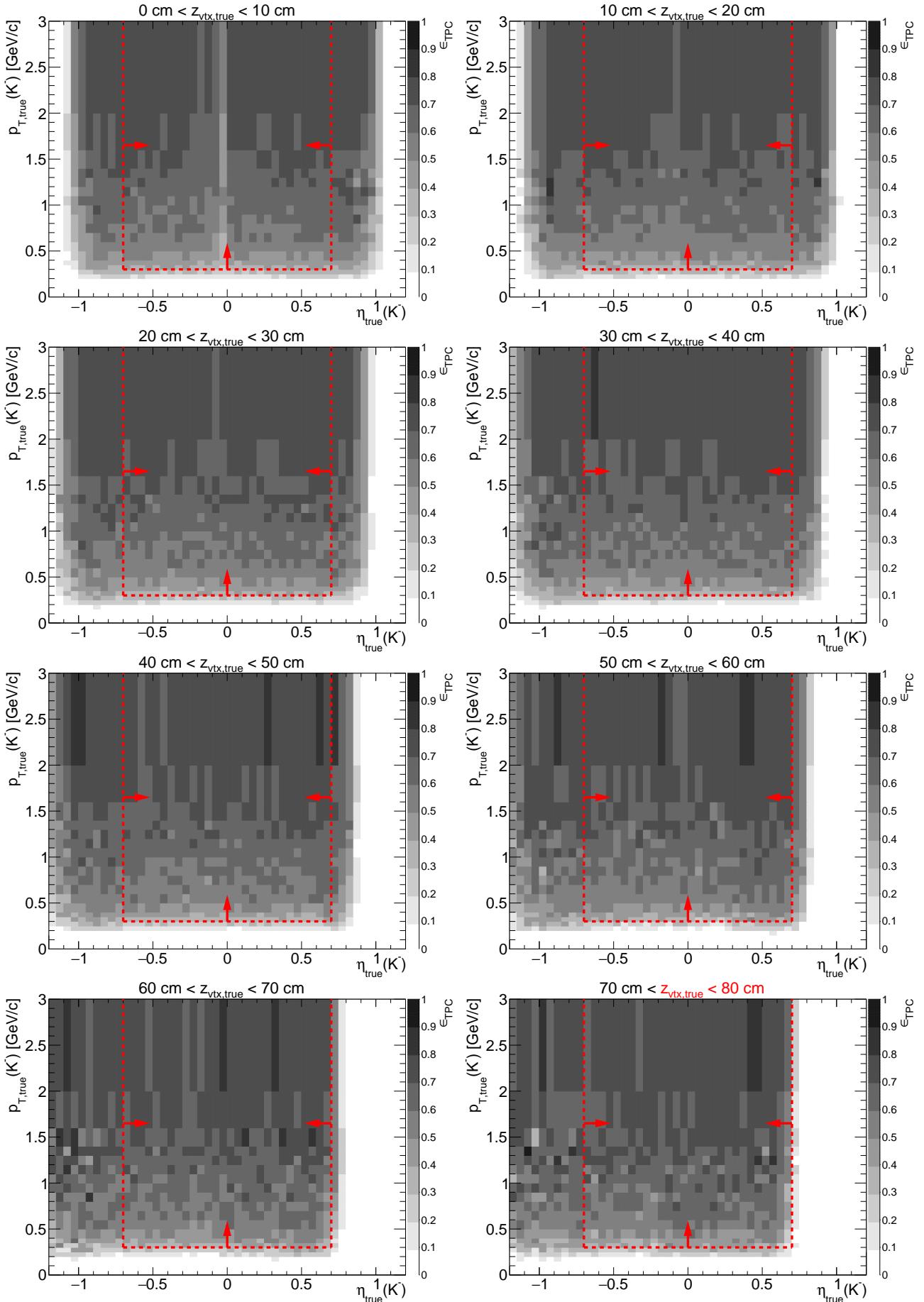
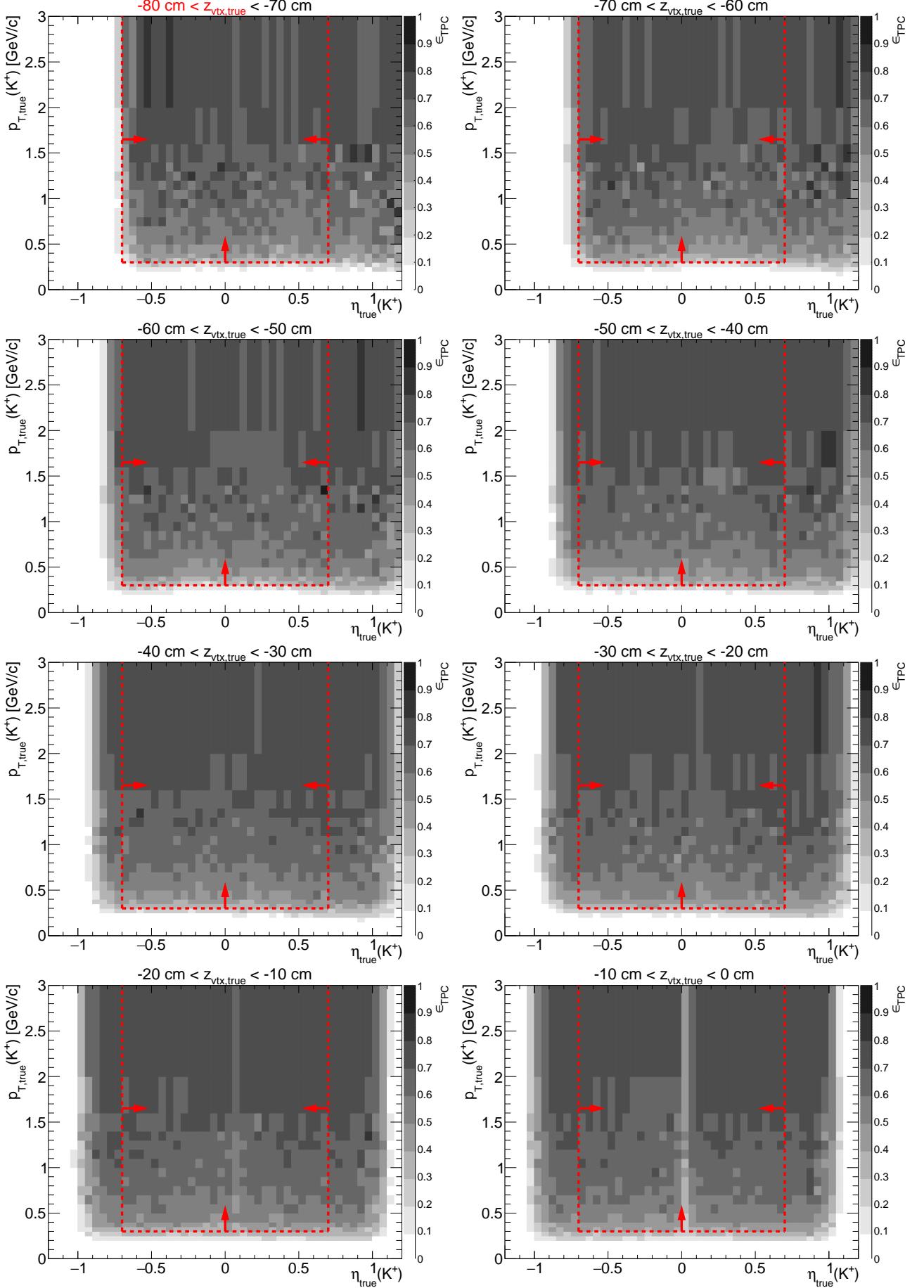


Figure A.3: TPC acceptance and reconstruction efficiency of K^+ . Each plot represents the TPC efficiency ϵ_{TPC} (z -axis) as a function of true particle pseudorapidity η (x -axis) and transverse momentum p_T (y -axis) in single z -vertex bin whose range is given at the top. Red lines and arrows indicate region accepted in analyses.



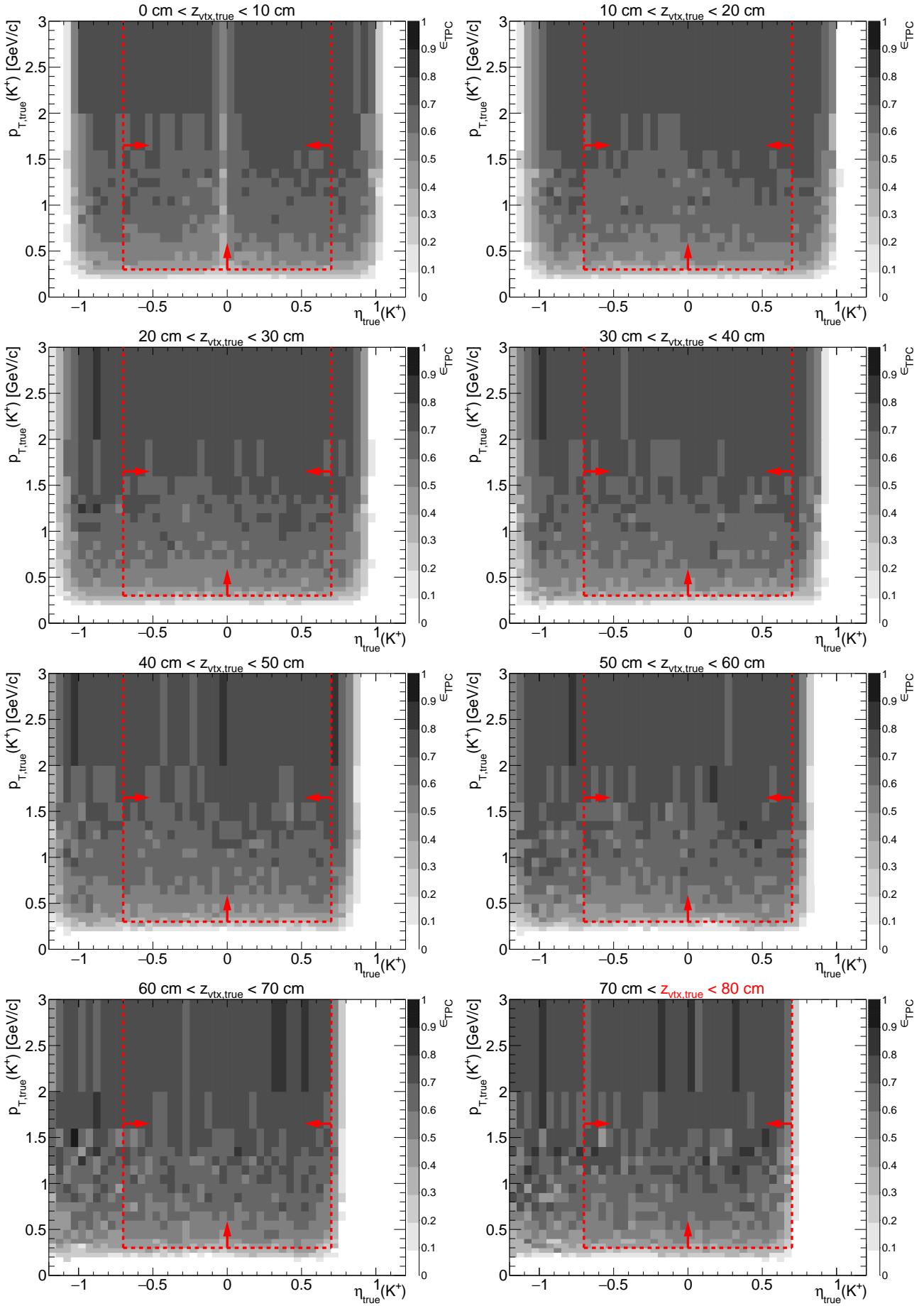
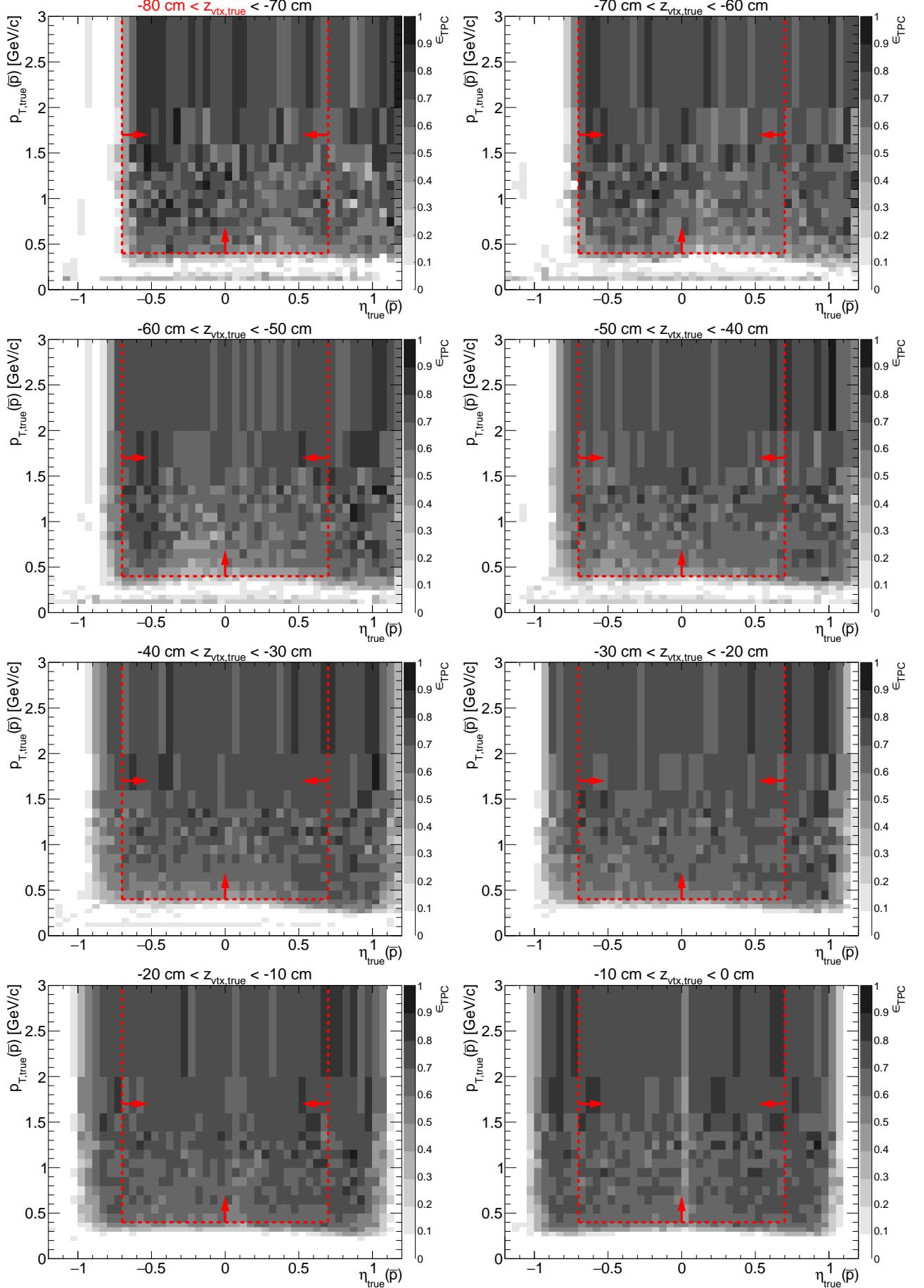


Figure A.4: TPC acceptance and reconstruction efficiency of \bar{p} . Each plot represents the TPC efficiency ϵ_{TPC} (z -axis) as a function of true particle pseudorapidity η (x -axis) and transverse momentum p_T (y -axis) in single z -vertex bin whose range is given at the top. Red lines and arrows indicate region accepted in analyses.



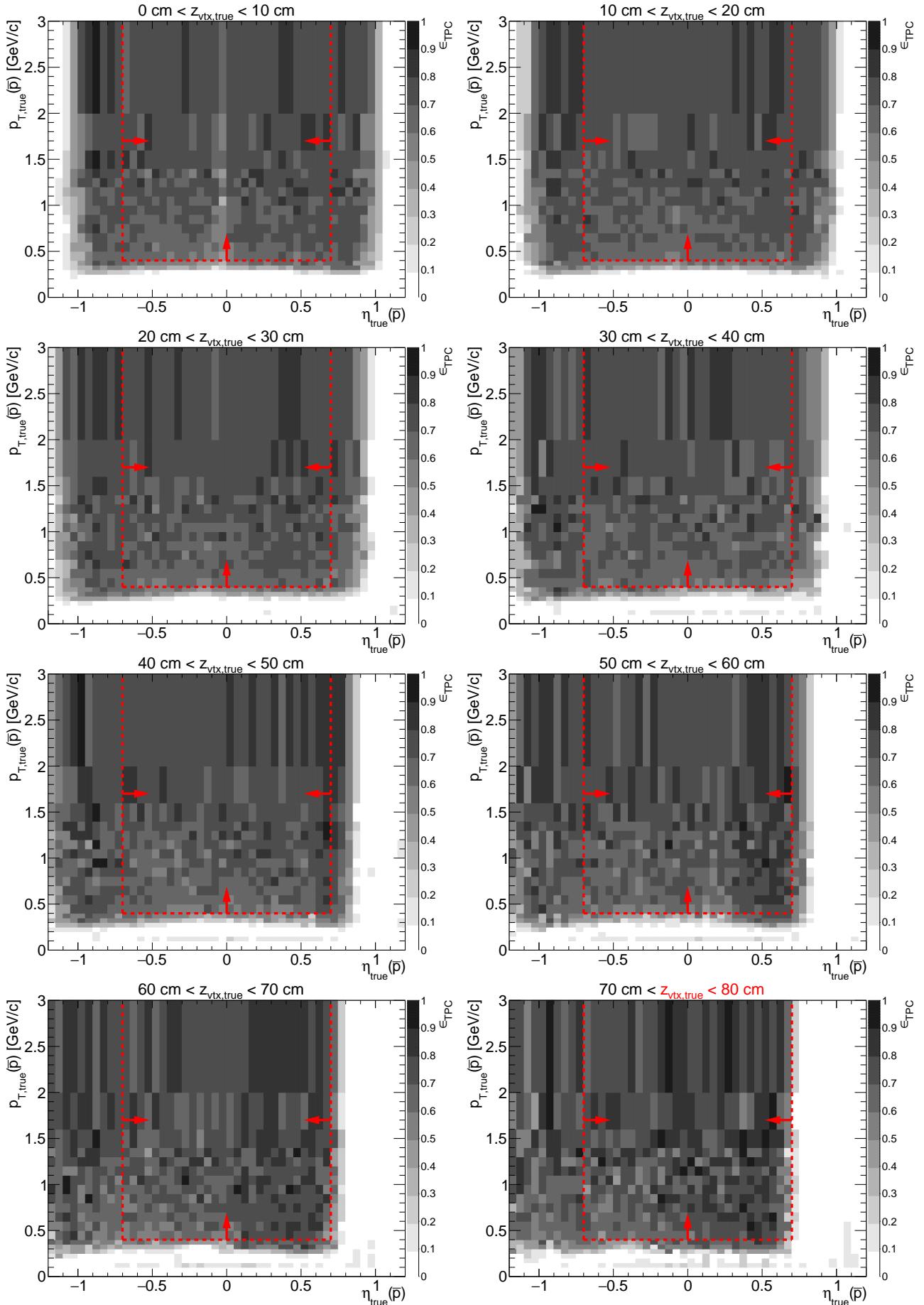
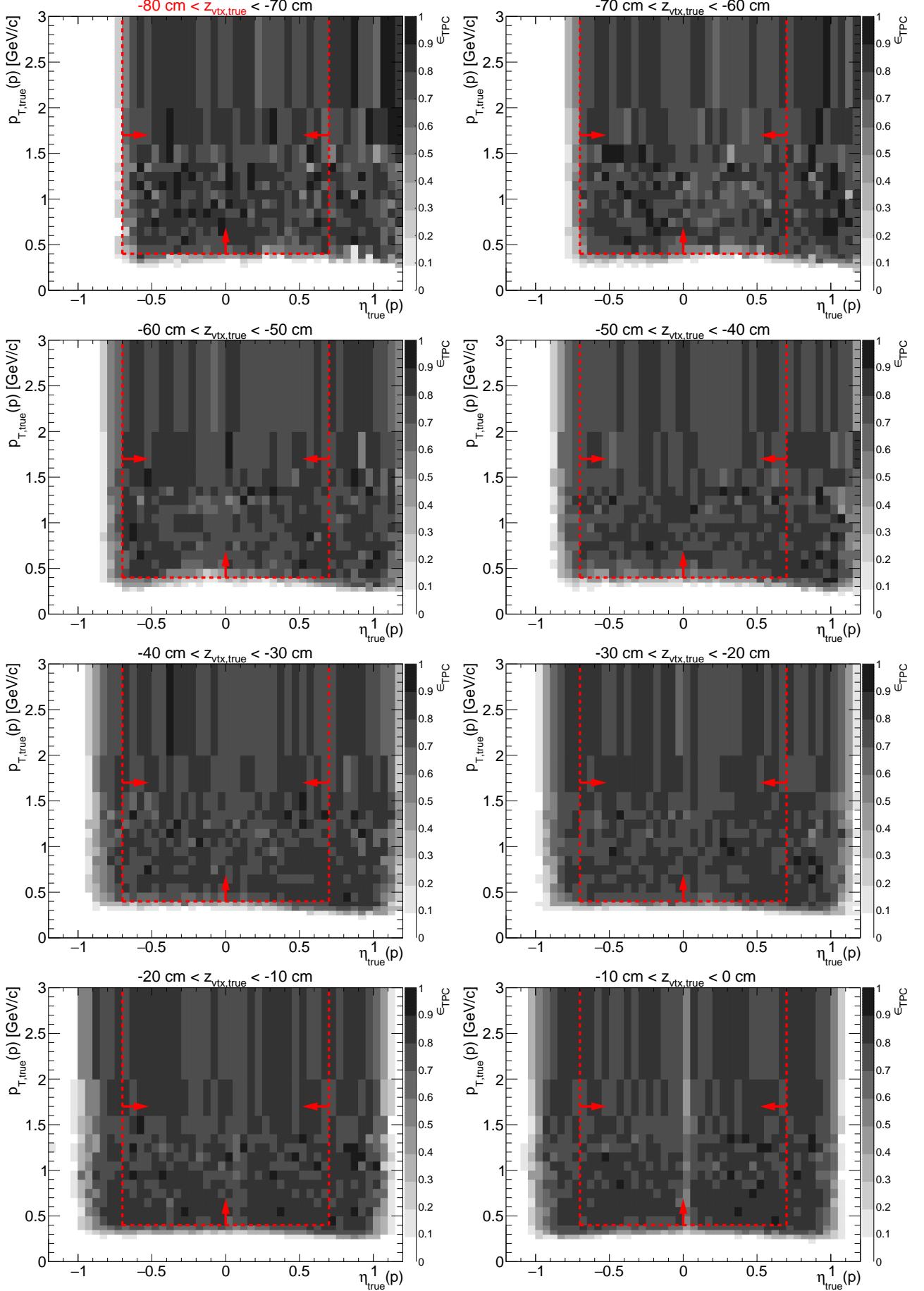
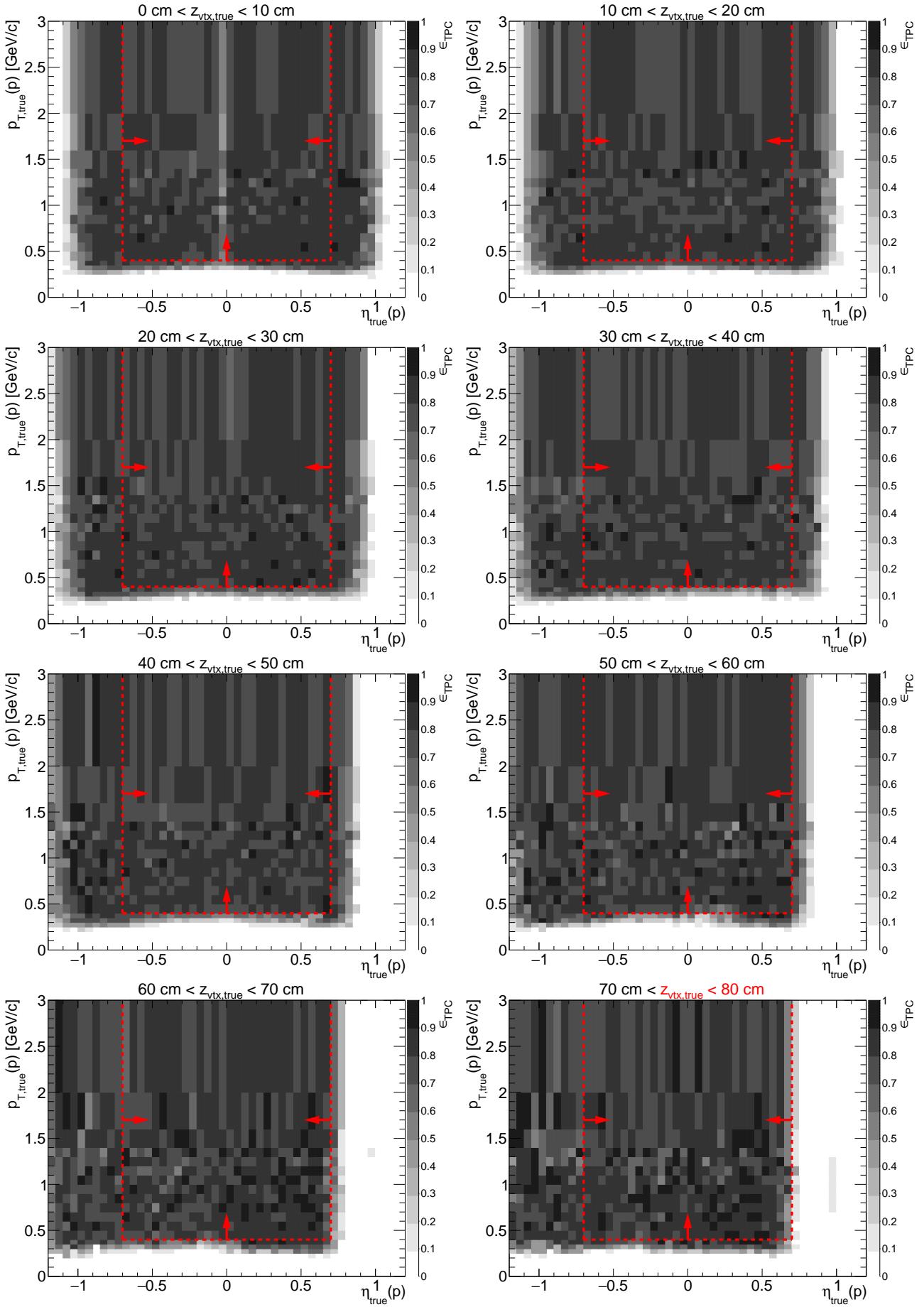


Figure A.5: TPC acceptance and reconstruction efficiency of p . Each plot represents the TPC efficiency ϵ_{TPC} (z -axis) as a function of true particle pseudorapidity η (x -axis) and transverse momentum p_T (y -axis) in single z -vertex bin whose range is given at the top. Red lines and arrows indicate region accepted in analyses.

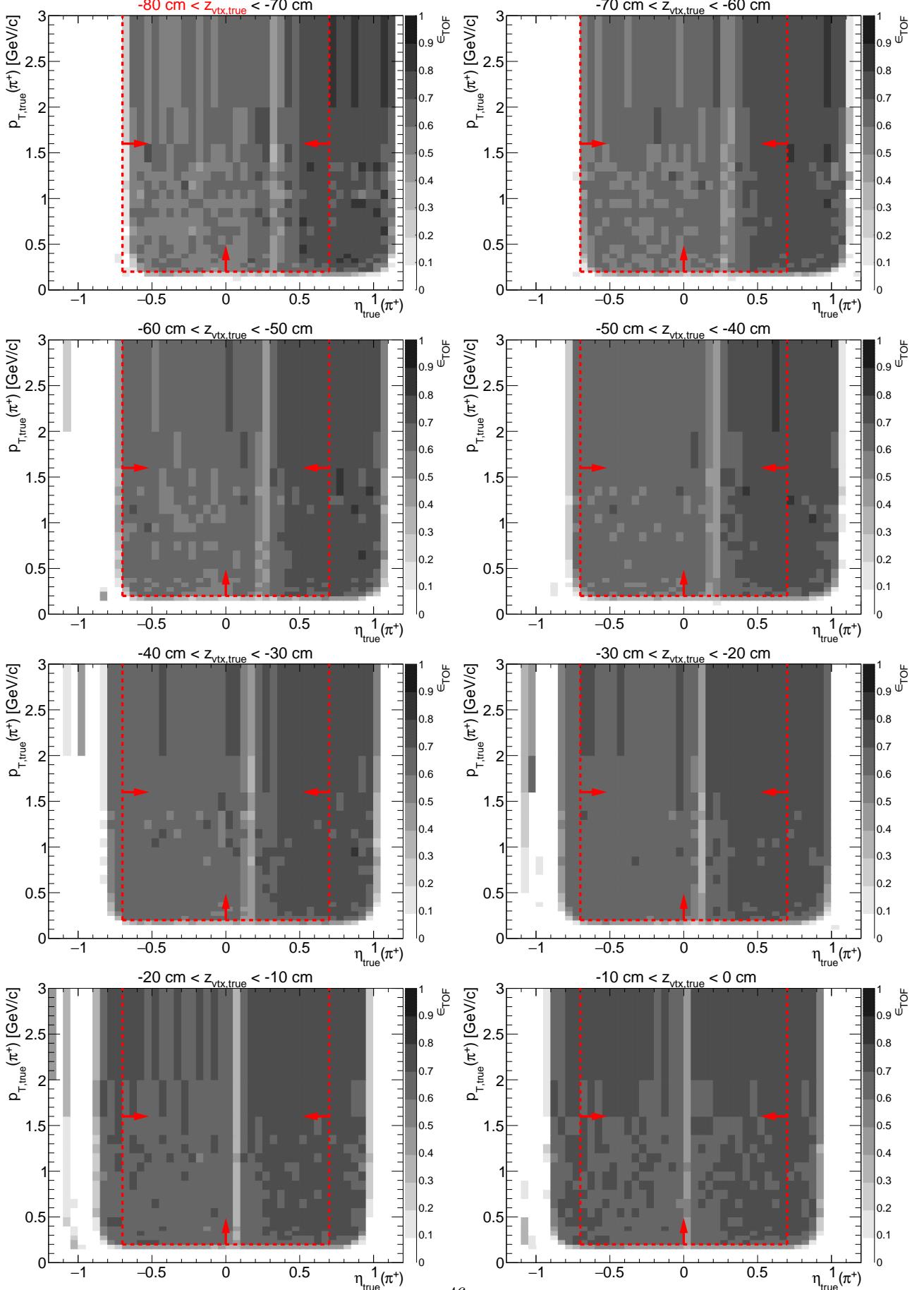




Appendix B

TOF hit reconstruction and matching efficiency

Figure B.1: TOF acceptance, reconstruction and matching efficiency of π^+ . Each plot represents the TOF efficiency ϵ_{TOF} (z -axis) as a function of true particle pseudorapidity η (x -axis) and transverse momentum p_T (y -axis) in single z -vertex bin whose range is given at the top. Red lines and arrows indicate region accepted in analyses.



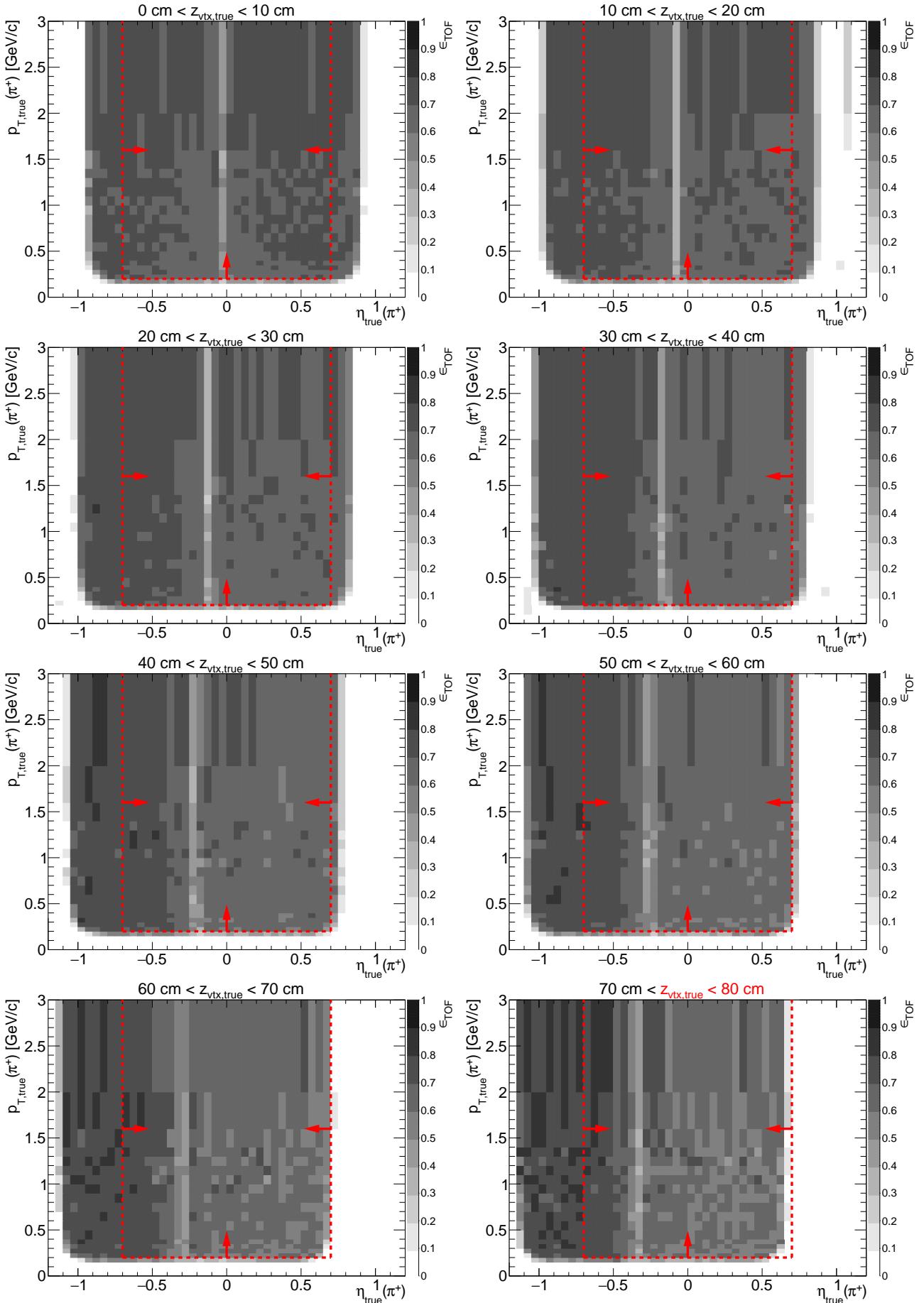
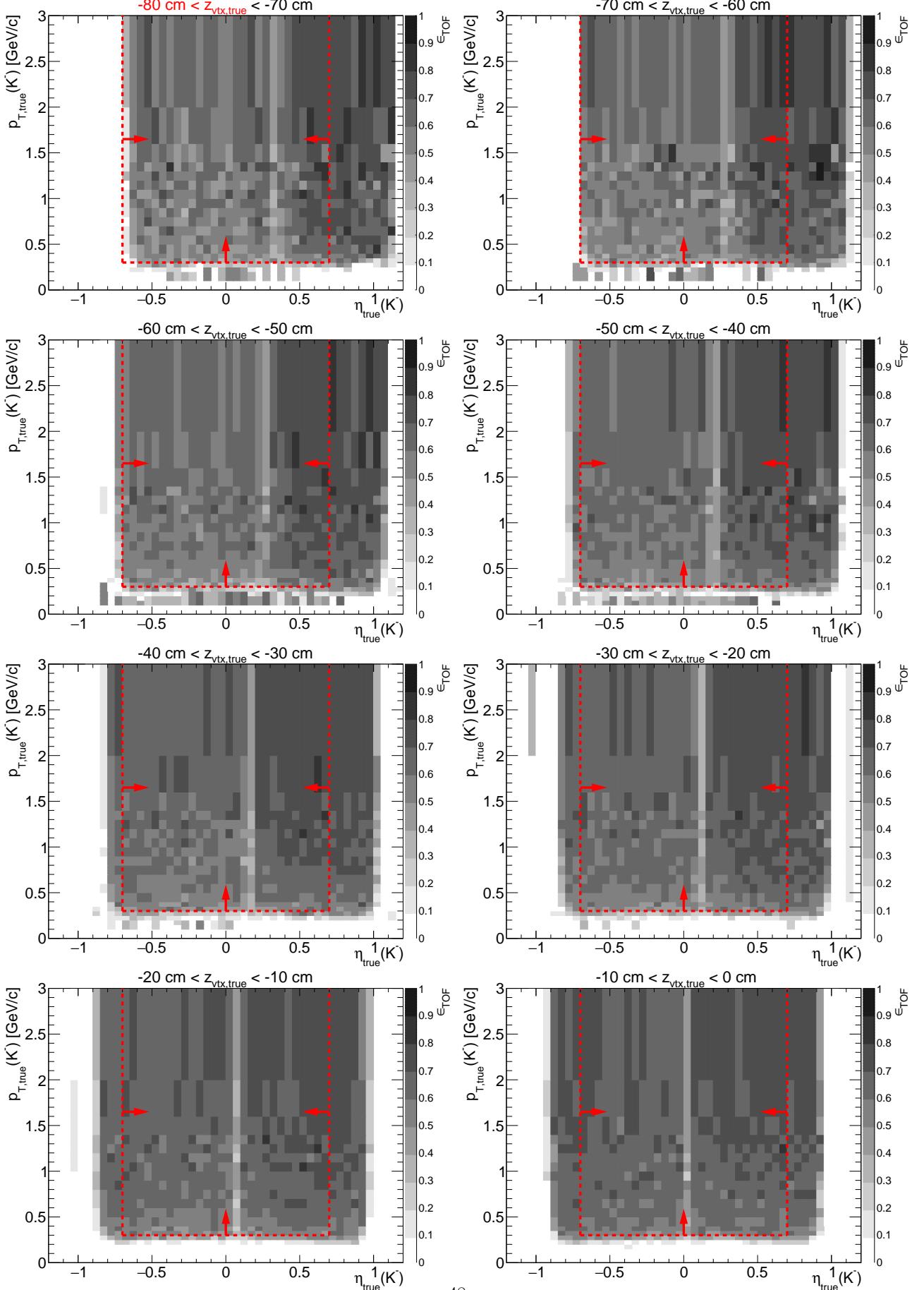


Figure B.2: TOF acceptance, reconstruction and matching efficiency of K^- . Each plot represents the TOF efficiency ϵ_{TOF} (z -axis) as a function of true particle pseudorapidity η (x -axis) and transverse momentum p_T (y -axis) in single z -vertex bin whose range is given at the top. Red lines and arrows indicate region accepted in analyses.



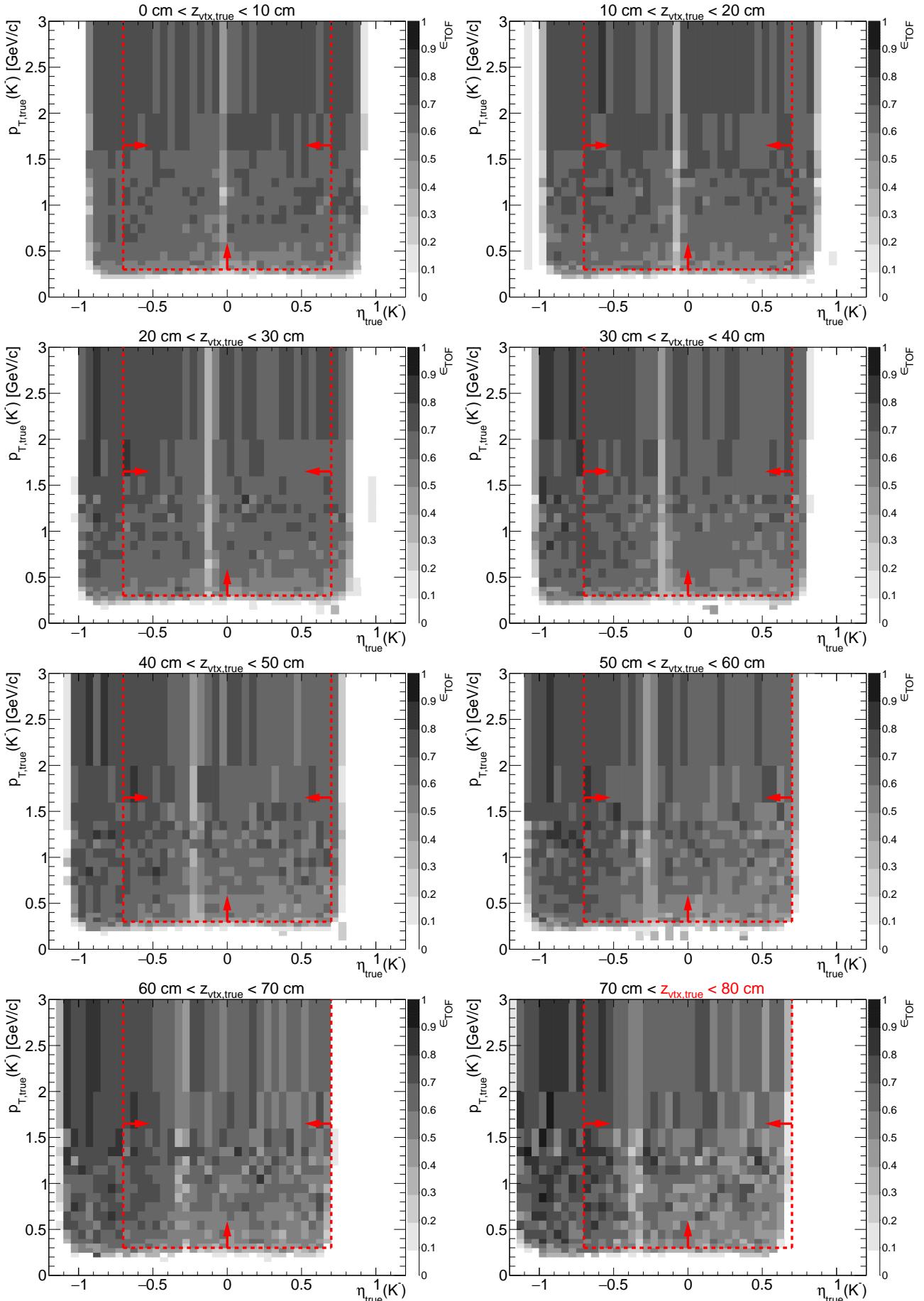
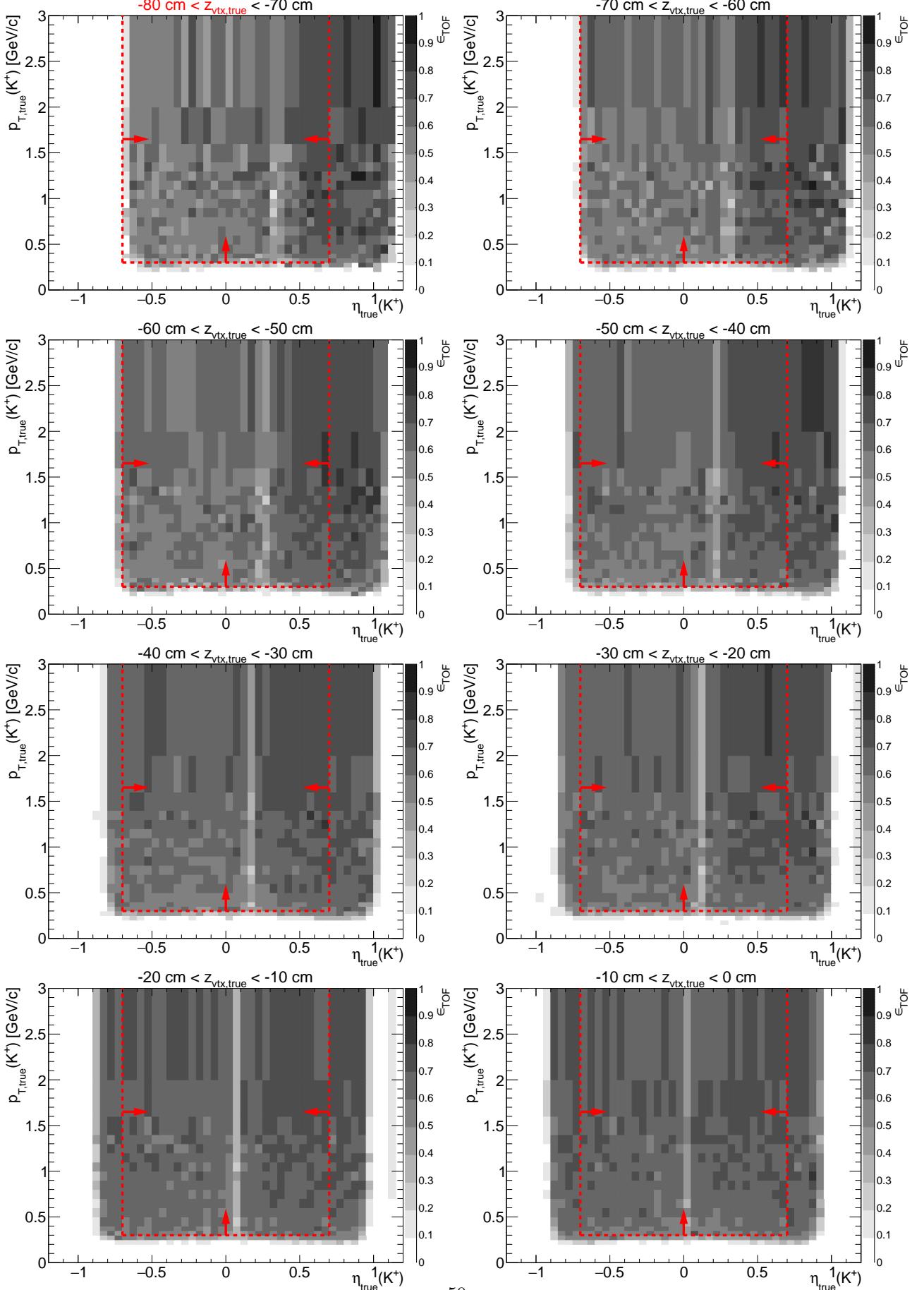


Figure B.3: TOF acceptance, reconstruction and matching efficiency of K^+ . Each plot represents the TOF efficiency ϵ_{TOF} (z -axis) as a function of true particle pseudorapidity η (x -axis) and transverse momentum p_T (y -axis) in single z -vertex bin whose range is given at the top. Red lines and arrows indicate region accepted in analyses.



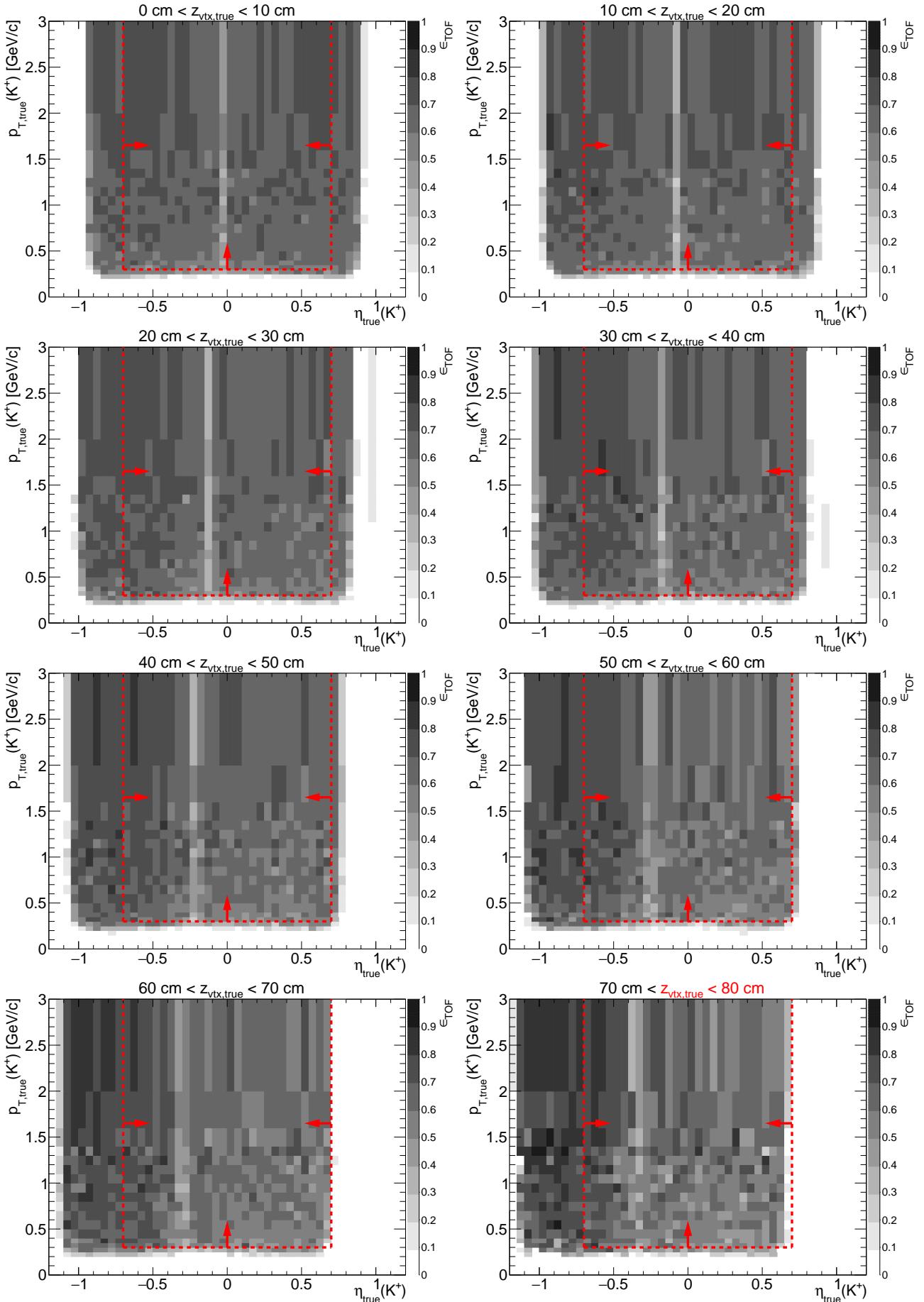
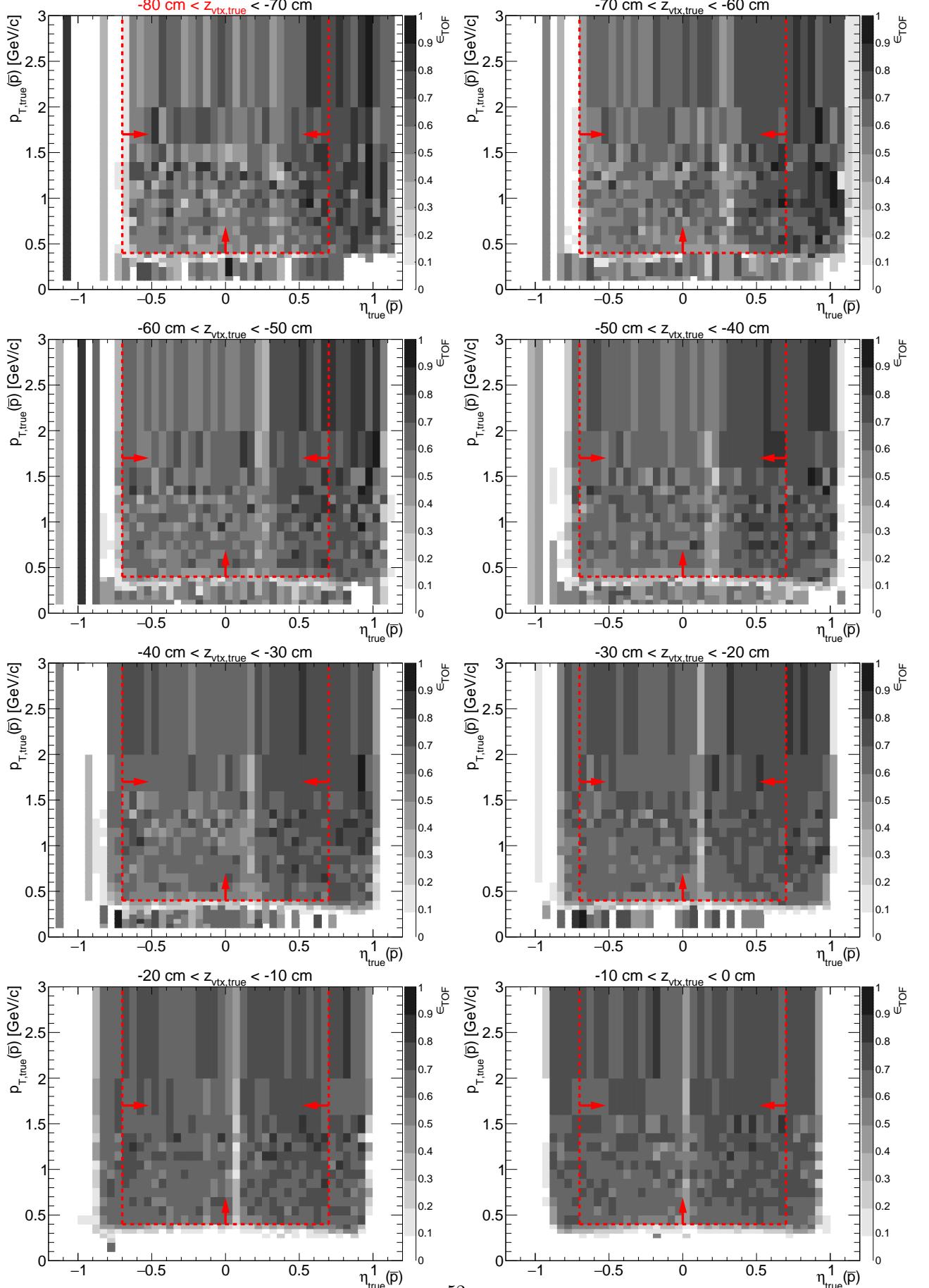


Figure B.4: TOF acceptance, reconstruction and matching efficiency of \bar{p} . Each plot represents the TOF efficiency ϵ_{TOF} (z -axis) as a function of true particle pseudorapidity η (x -axis) and transverse momentum p_T (y -axis) in single z -vertex bin whose range is given at the top. Red lines and arrows indicate region accepted in analyses.



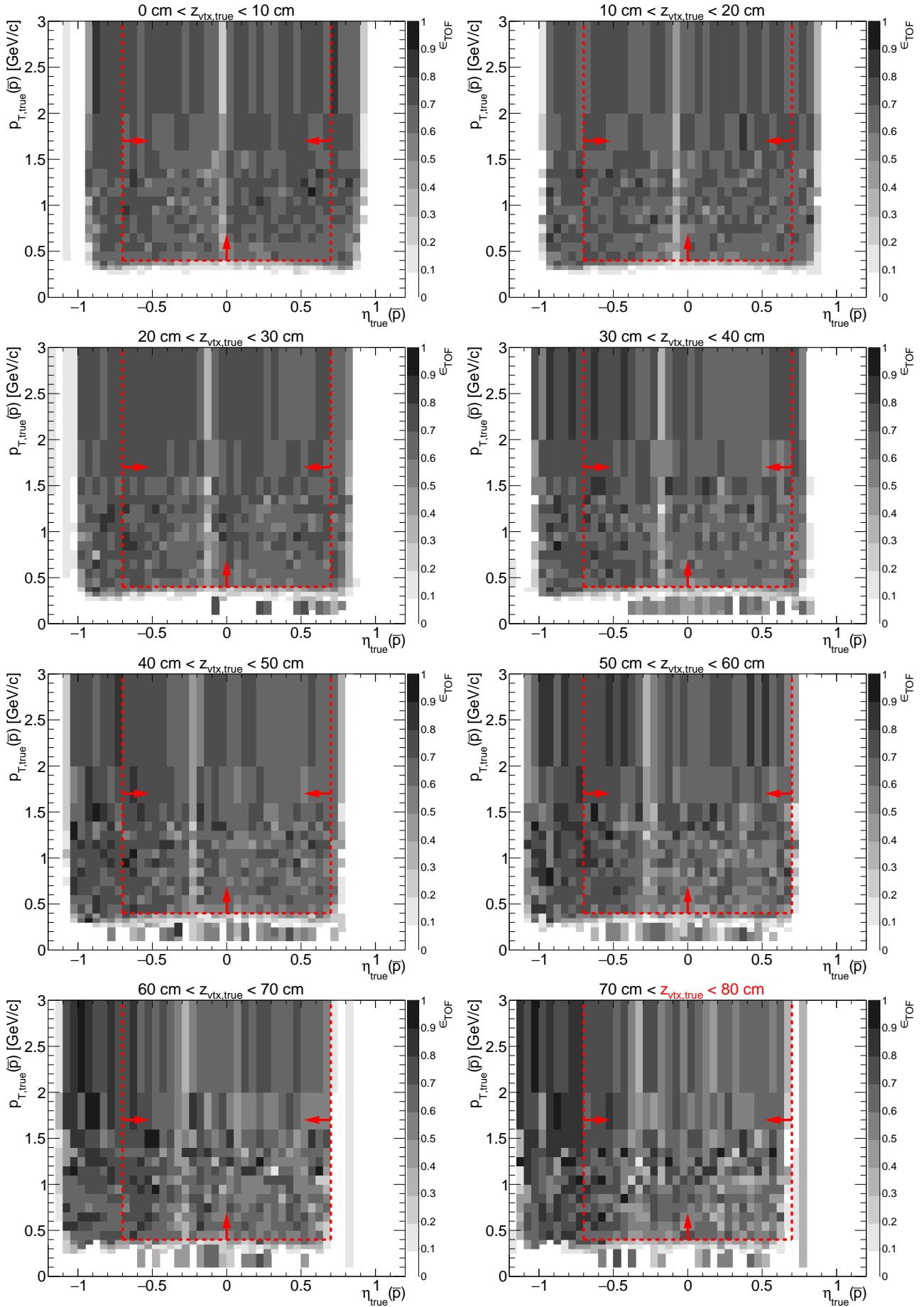
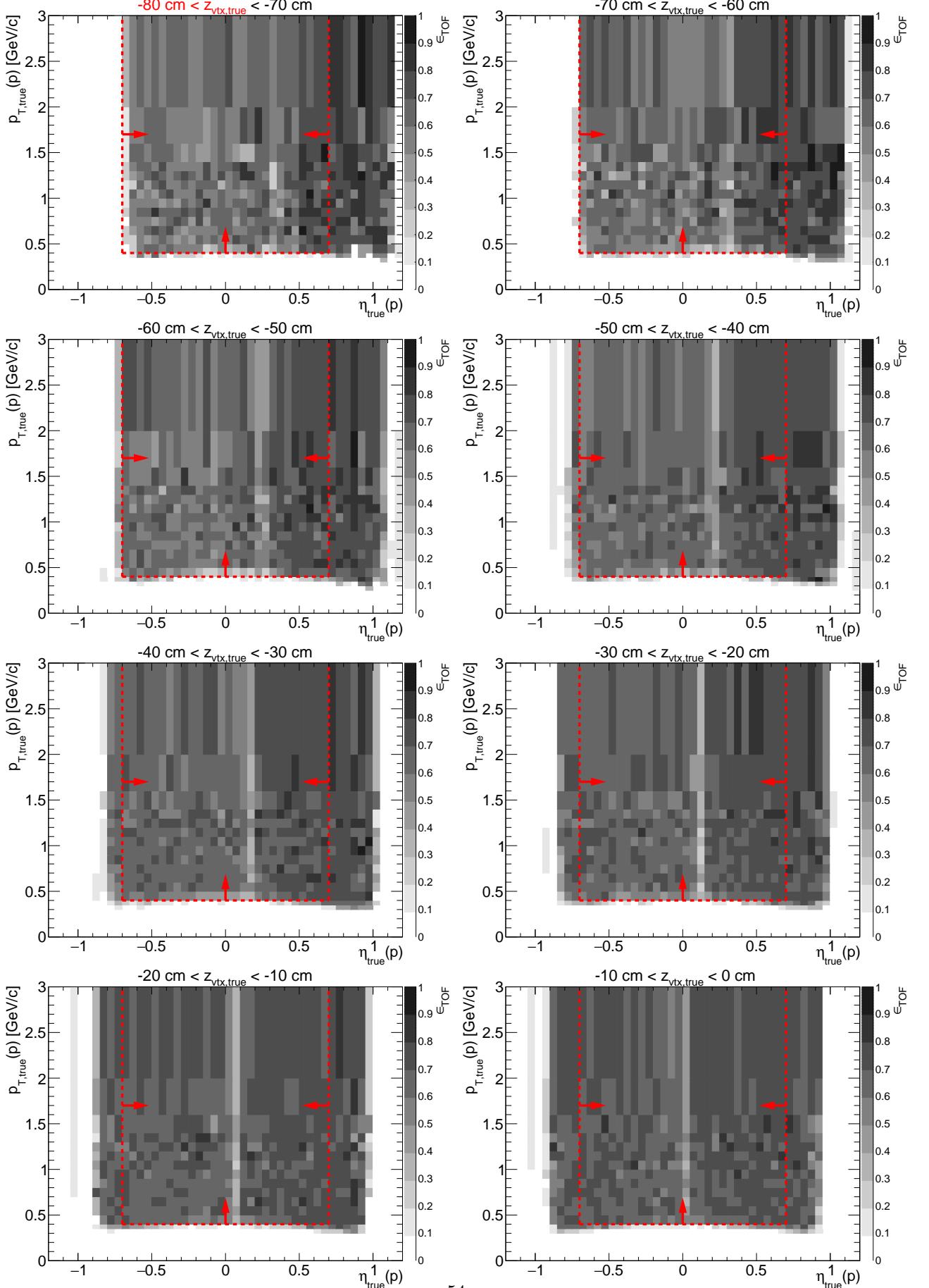
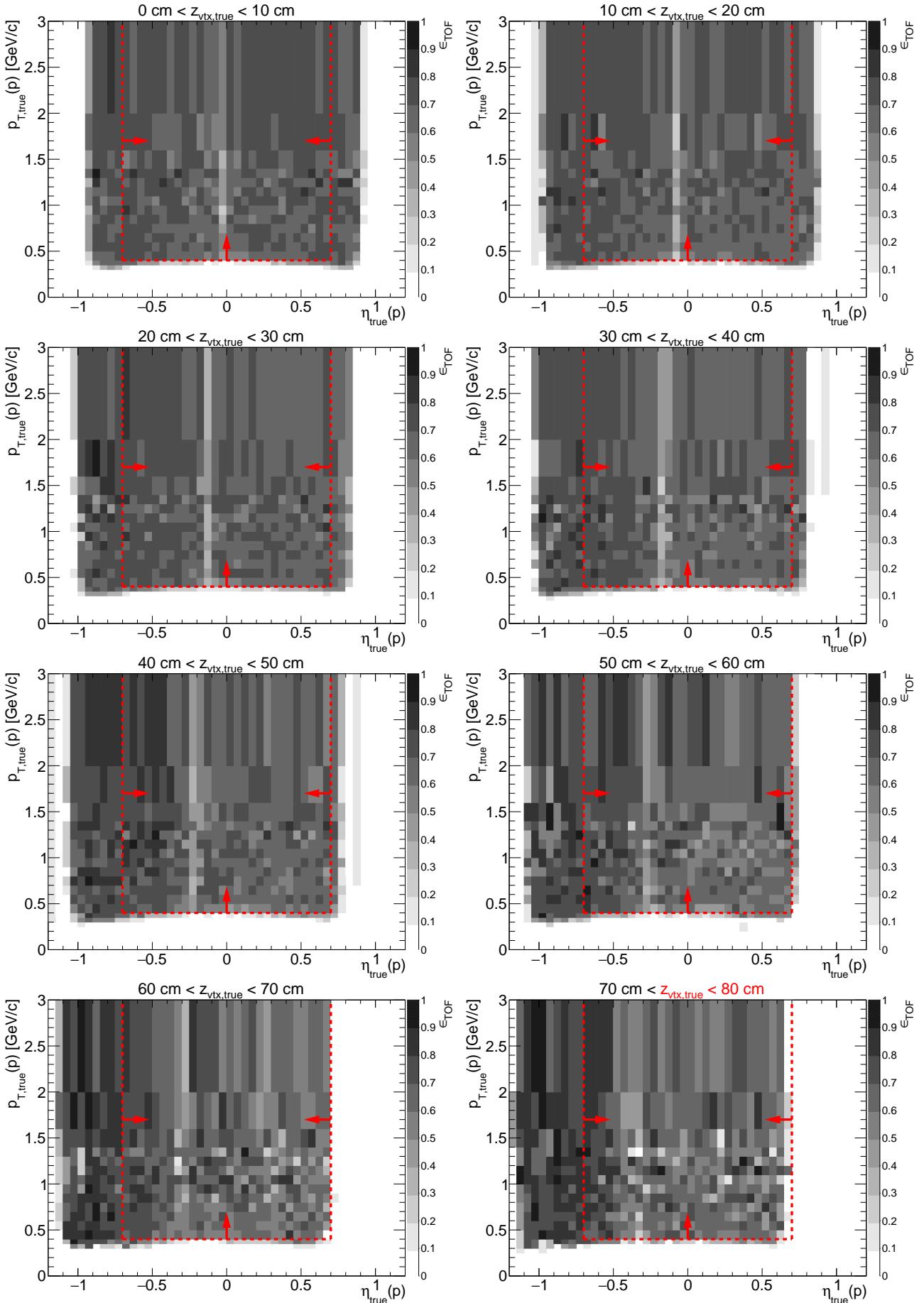


Figure B.5: TOF acceptance, reconstruction and matching efficiency of p . Each plot represents the TOF efficiency ϵ_{TOF} (z -axis) as a function of true particle pseudorapidity η (x -axis) and transverse momentum p_T (y -axis) in single z -vertex bin whose range is given at the top. Red lines and arrows indicate region accepted in analyses.

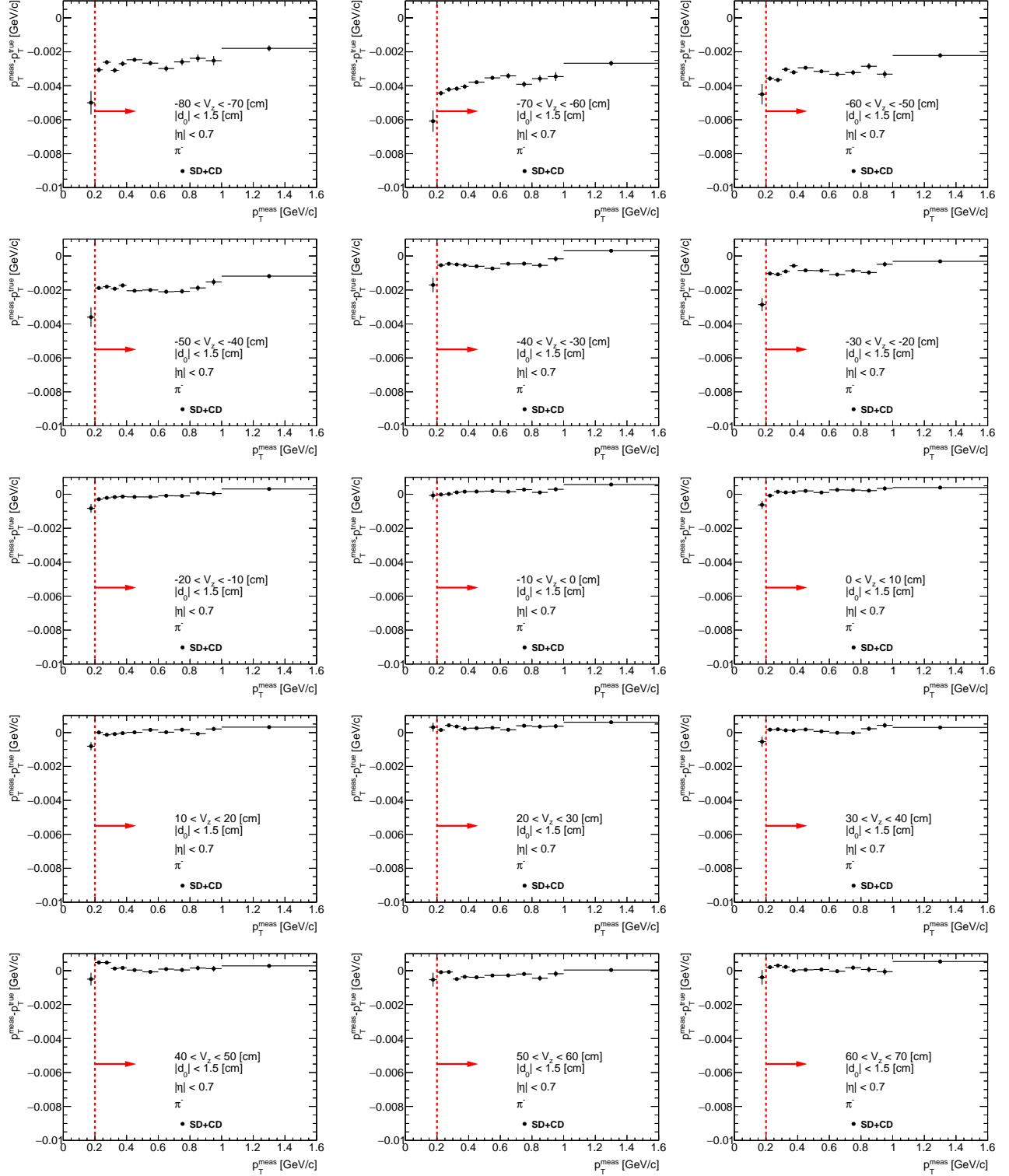




Appendix C

Energy Loss Correction

Figure C.1: Energy loss correction $p_T^{meas} - p_T^{true}$ for π^- as a function of reconstructed transverse momentum p_T^{meas} ($|\eta| < 0.7$) in single z -vertex bin whose range is given on each plot. Red lines and arrows indicate region accepted in analyses. One can notice an offset of about $3 - 4$ MeV for negative z -vertex. It is a known issue with STAR simulation where HFT support material is badly described for negative $z < -30$ cm.



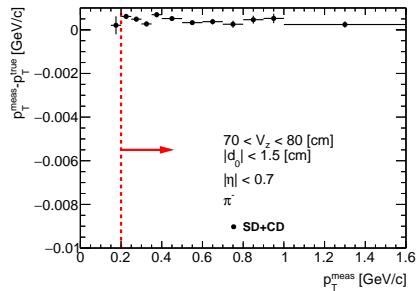


Figure C.2: Energy loss correction $p_T^{meas} - p_T^{true}$ for π^+ as a function of reconstructed transverse momentum p_T^{meas} ($|\eta| < 0.7$) in single z -vertex bin whose range is given on each plot. Red lines and arrows indicate region accepted in analyses. One can notice an offset of about 3 – 4 MeV for negative z -vertex. It is a known issue with STAR simulation where HFT support material is badly described for negative $z < -30$ cm.

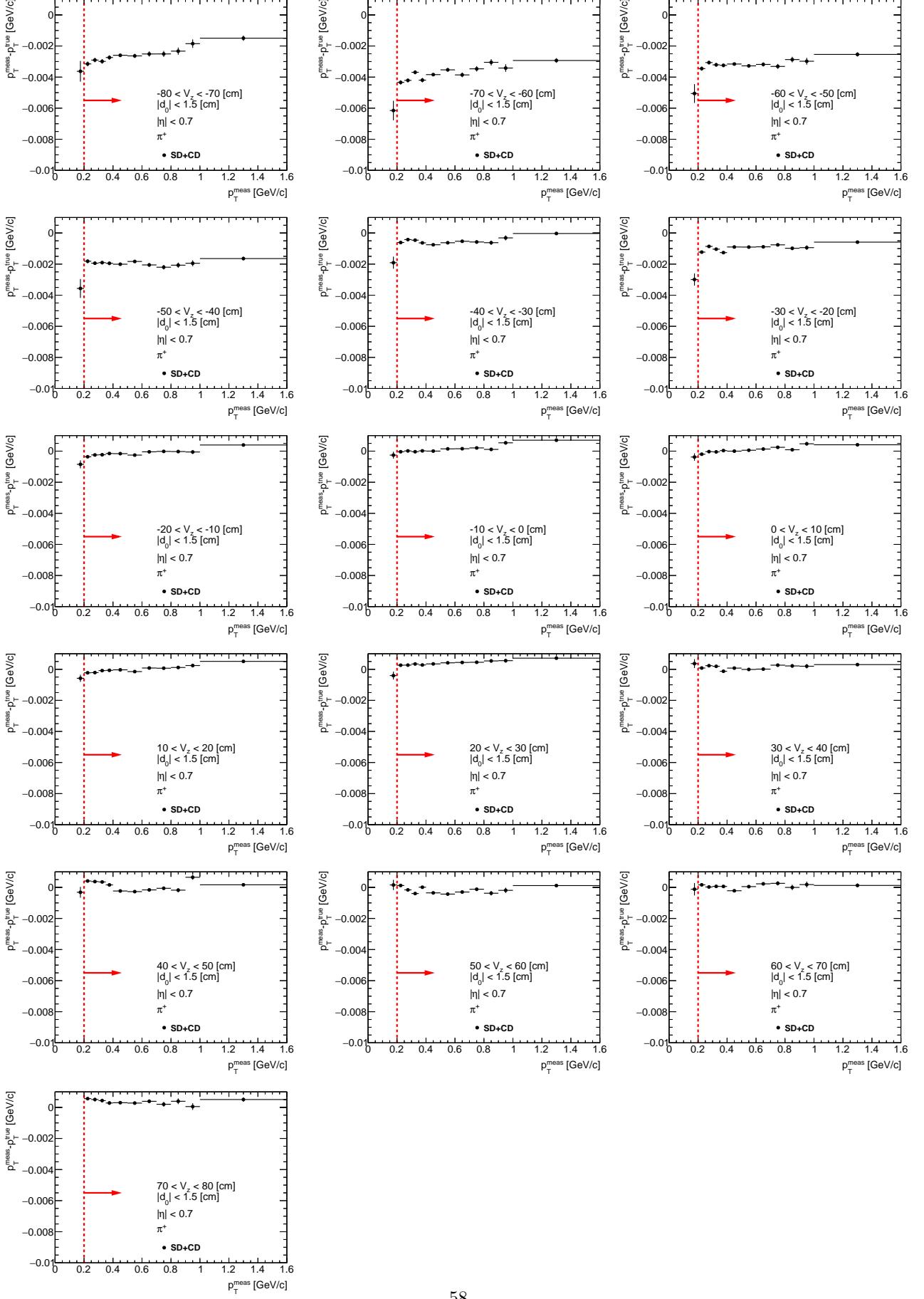


Figure C.3: Energy loss correction $p_T^{meas} - p_T^{true}$ for K^+ as a function of reconstructed transverse momentum p_T^{meas} ($|\eta| < 0.7$) in single z -vertex bin whose range is given on each plot. Red lines and arrows indicate region accepted in analyses.

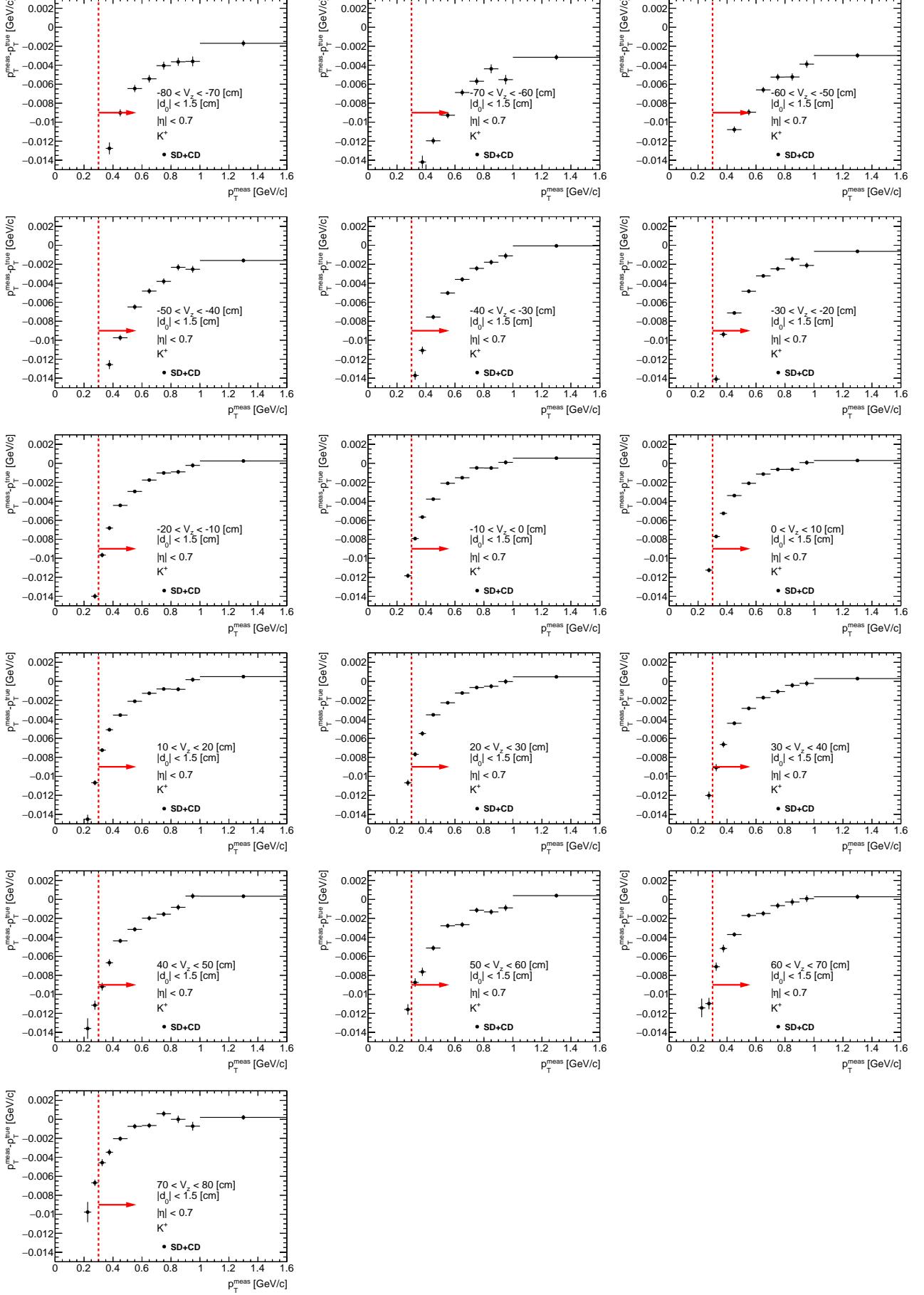


Figure C.4: Energy loss correction $p_T^{meas} - p_T^{true}$ for \bar{p} as a function of reconstructed transverse momentum p_T^{meas} ($|\eta| < 0.7$) in single z -vertex bin whose range is given on each plot. Red lines and arrows indicate region accepted in analyses.

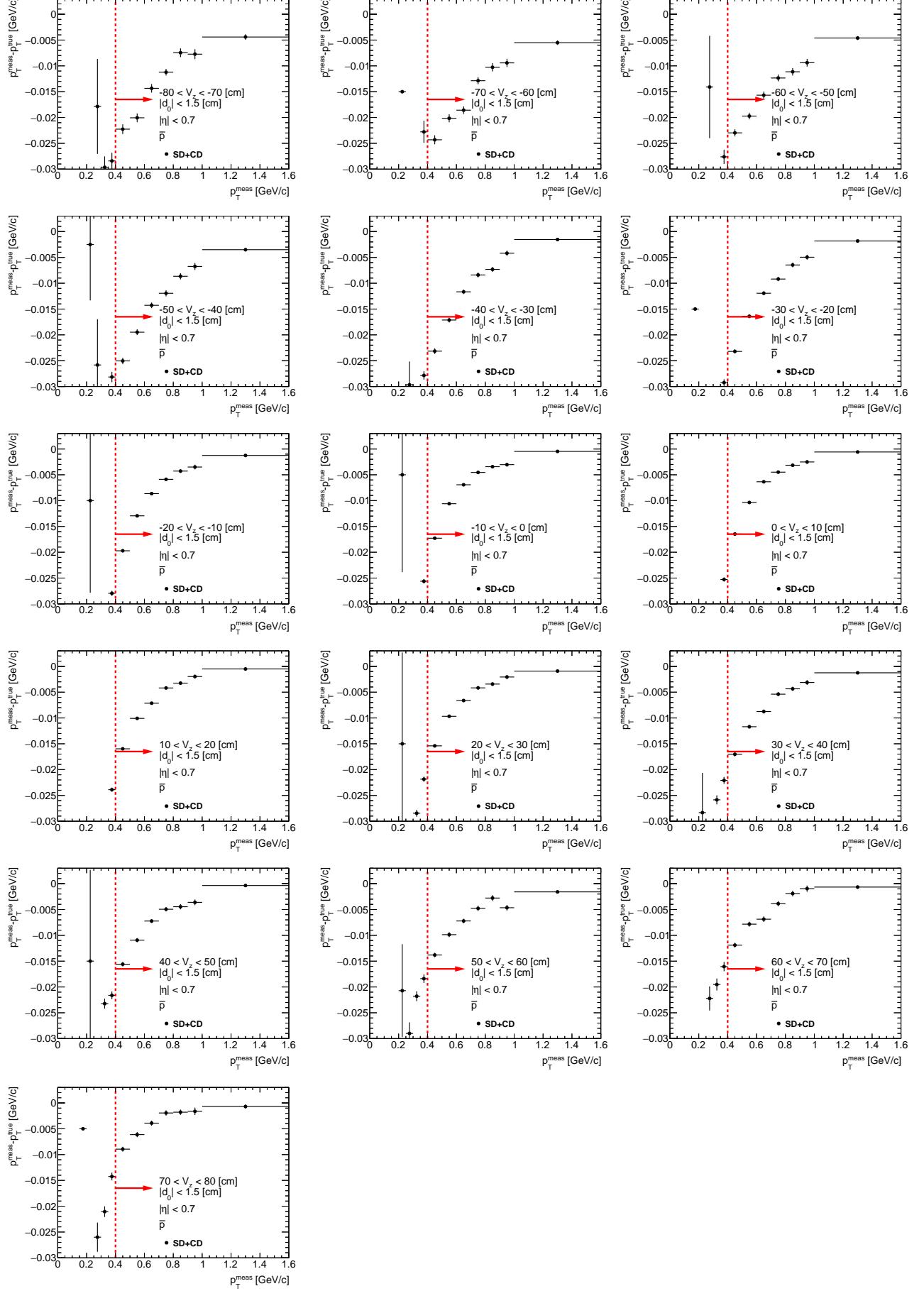


Figure C.5: Energy loss correction $p_T^{meas} - p_T^{true}$ for p as a function of reconstructed transverse momentum p_T^{meas} ($|\eta| < 0.7$) in single z -vertex bin whose range is given on each plot. Red lines and arrows indicate region accepted in analyses.

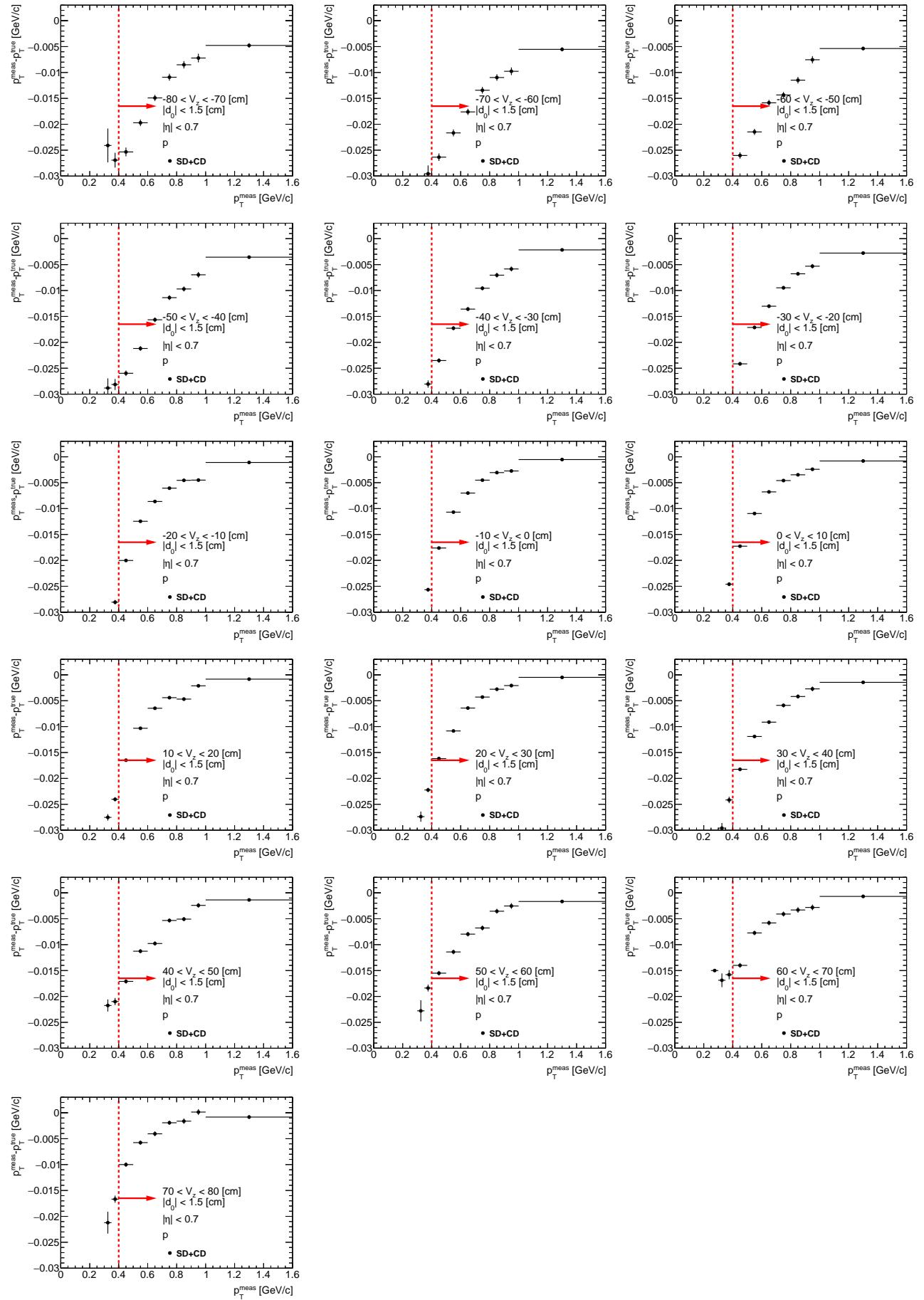


Figure C.6: Energy loss correction $p_T^{meas} - p_T^{true}$ for negative particles as a function of reconstructed transverse momentum p_T^{meas} ($|\eta| < 0.7$) in single z -vertex bin whose range is given on each plot. Red lines and arrows indicate region accepted in analyses.

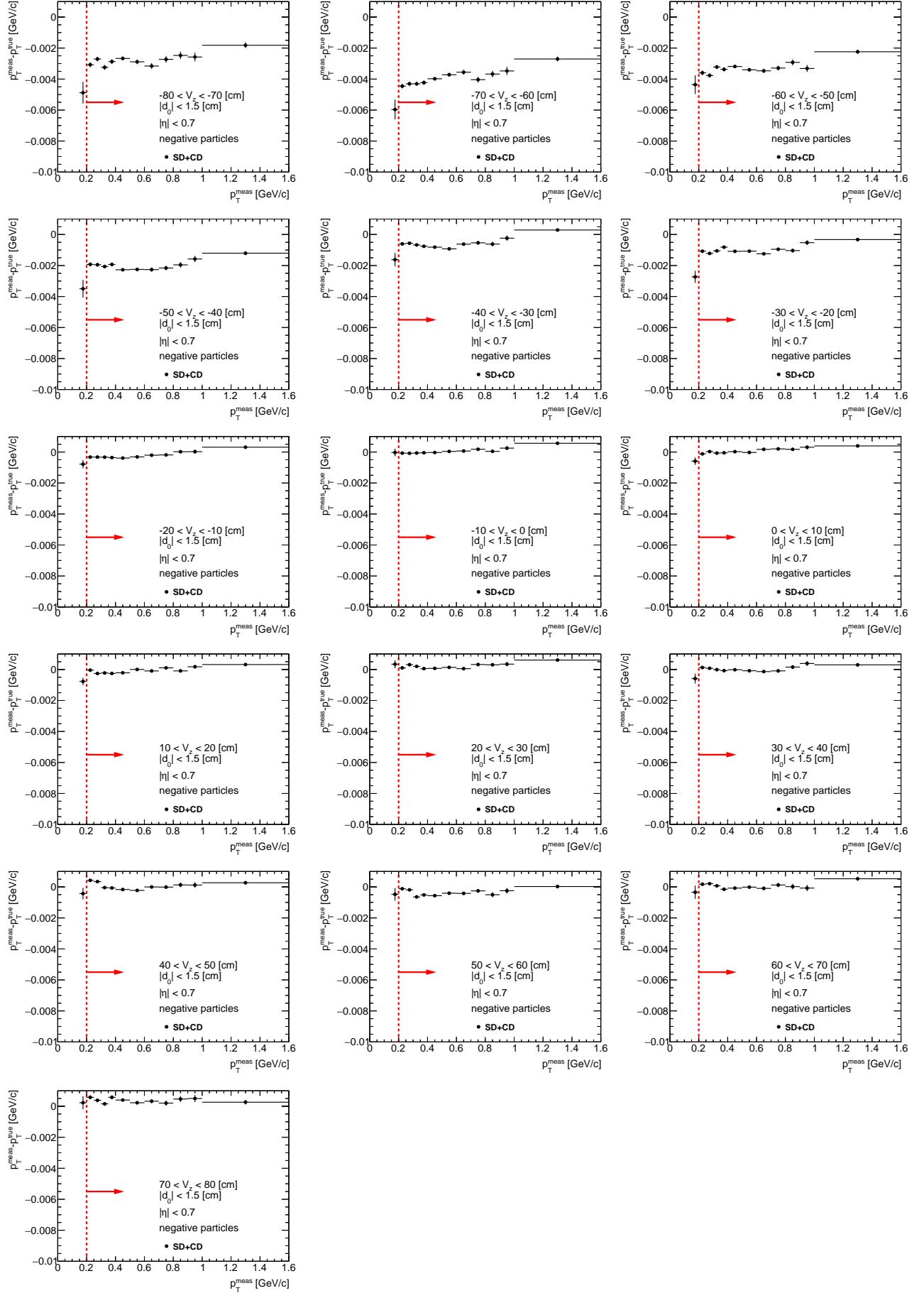


Figure C.7: Energy loss correction $p_T^{meas} - p_T^{true}$ for positive particles as a function of reconstructed transverse momentum p_T^{meas} ($|\eta| < 0.7$) in single z -vertex bin whose range is given on each plot. Red lines and arrows indicate region accepted in analyses.

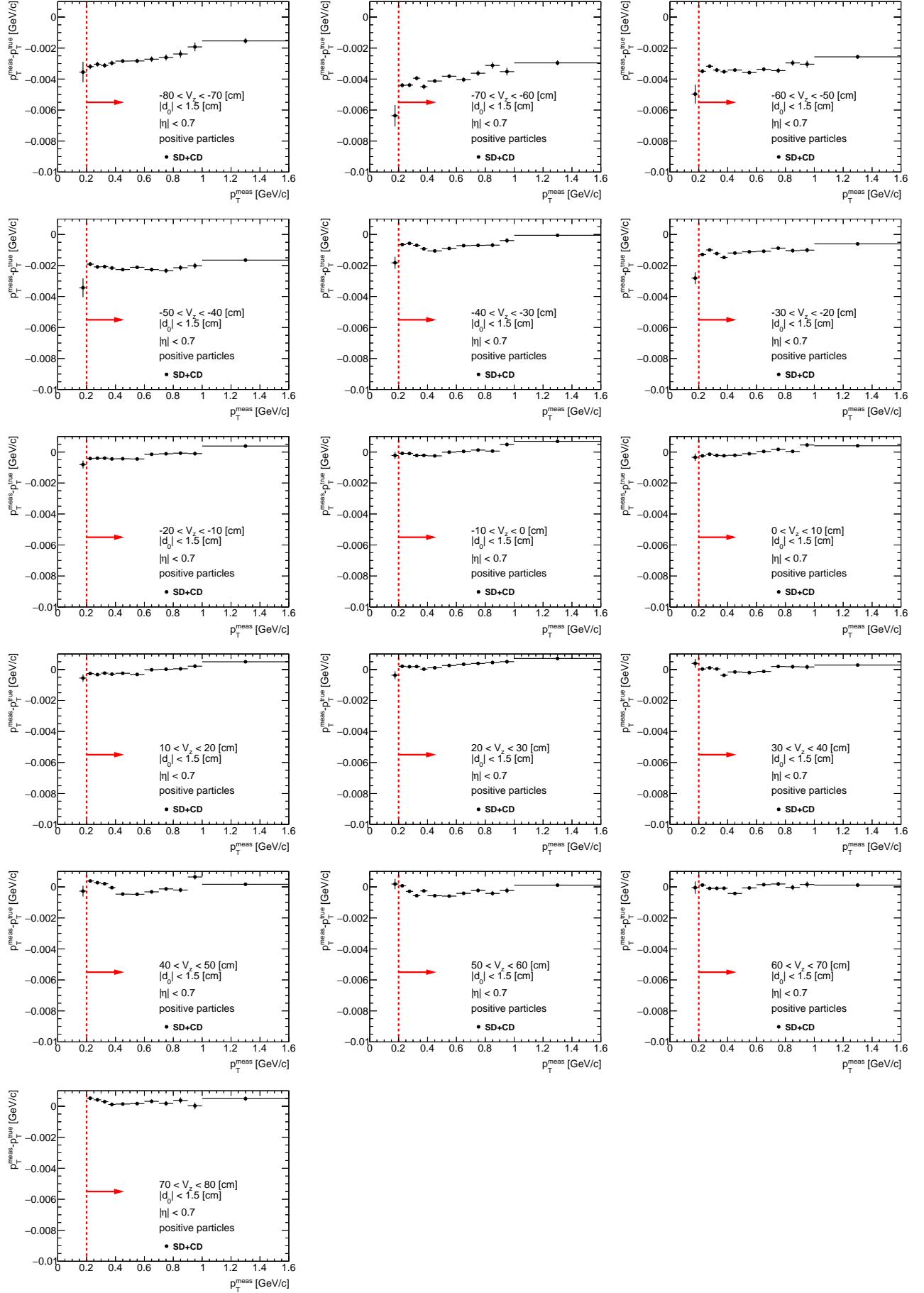
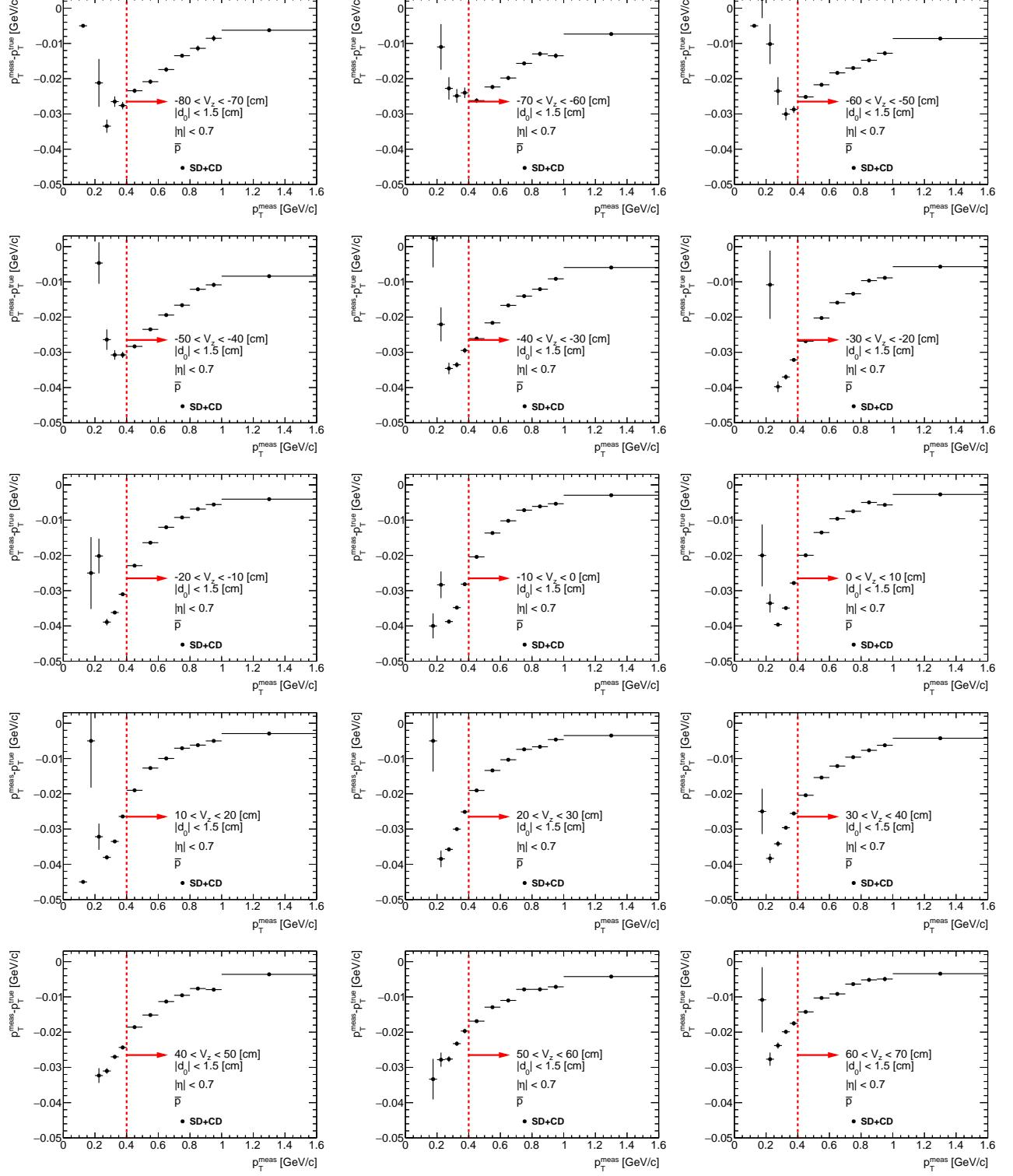


Figure C.8: Energy loss correction $p_T^{meas} - p_T^{true}$ for \bar{p} as a function of reconstructed global track transverse momentum p_T^{meas} ($|\eta| < 0.7$) in single z -vertex bin whose range is given on each plot. Red lines and arrows indicate region accepted in analyses. One may need the energy loss correction for reconstructed global proton and antiproton tracks to estimate the knock-out proton background. During the reconstruction, global tracks are corrected only for energy losses in TPC, whereas primary tracks have the information of dead material inside and in front of TPC. Since that, there is an offset of about 4 MeV for global tracks.



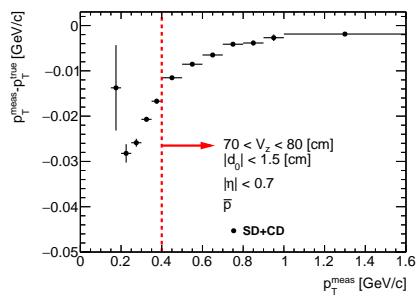
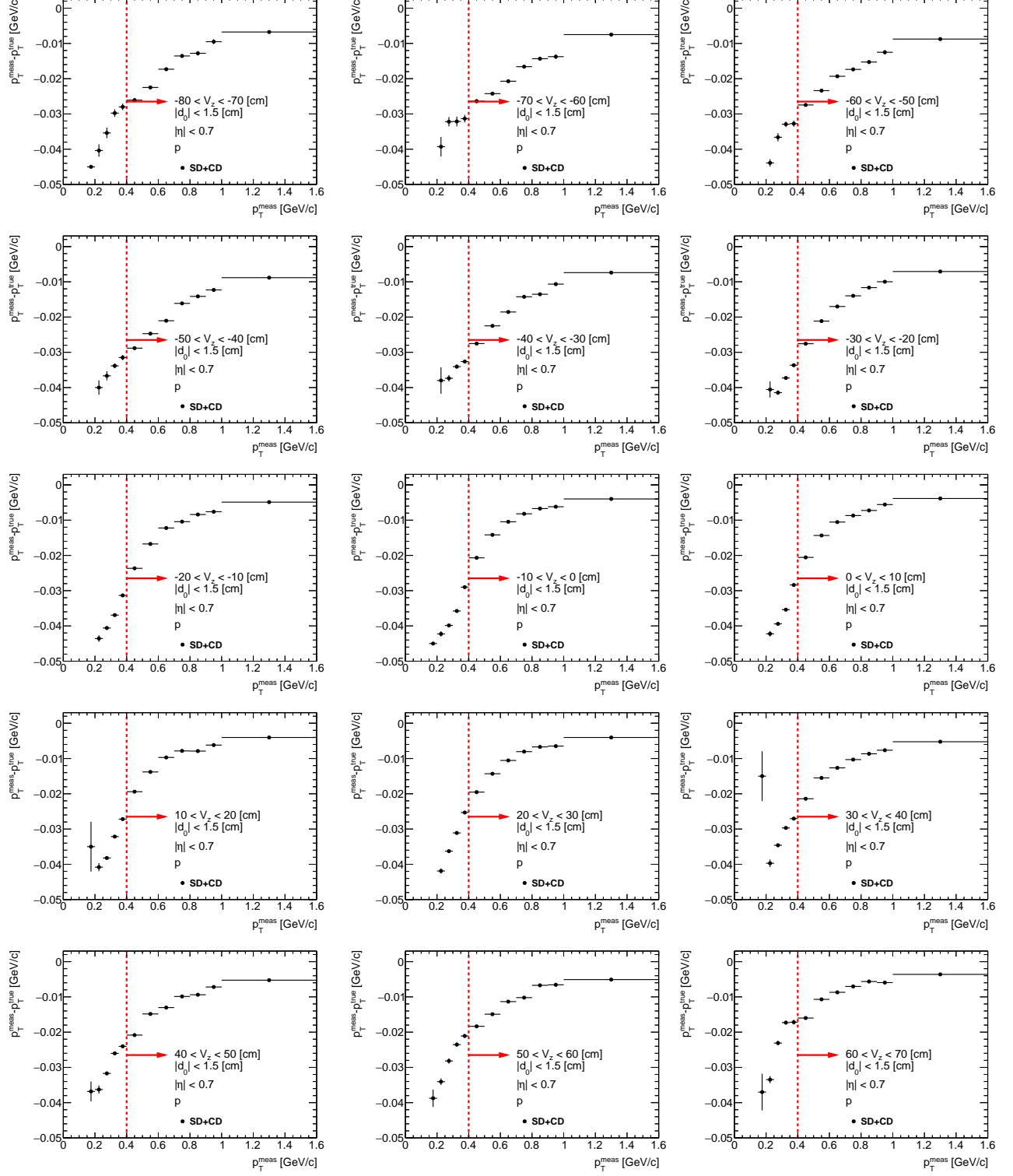
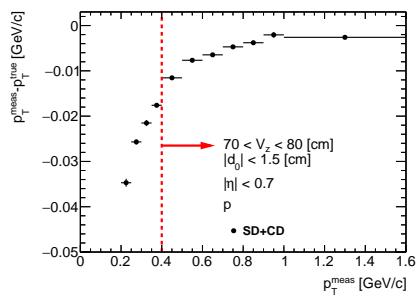


Figure C.9: Energy loss correction $p_T^{meas} - p_T^{true}$ for p as a function of reconstructed global track transverse momentum p_T^{meas} ($|\eta| < 0.7$) in single z -vertex bin whose range is given on each plot. Red lines and arrows indicate region accepted in analyses. One may need the energy loss correction for reconstructed global proton and antiproton tracks to estimate the knock-out proton background. During the reconstruction, global tracks are corrected only for energy losses in TPC, whereas primary tracks have the information of dead material inside and in front of TPC. Since that, there is an offset of about 4 MeV for global tracks.

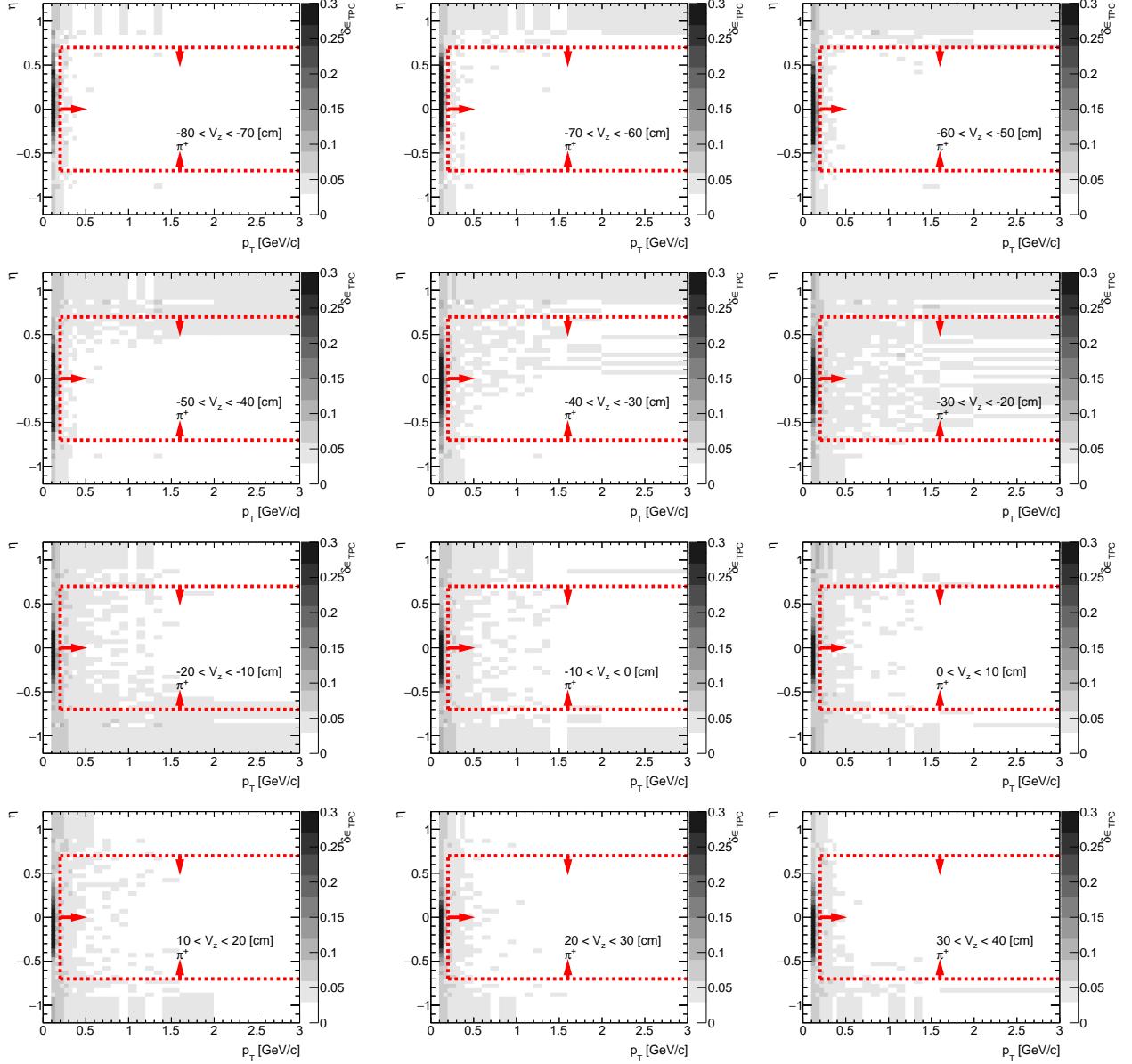




Appendix D

Dead material effect on TPC track reconstruction efficiency

Figure D.1: The amount of lost π^- due to the interaction with dead material in front of TPC in CD MC sample. Each plot represents the fraction of lost π^+ , $\delta\epsilon_{TPC}$ (z -axis), as a function of true particle pseudorapidity η (y -axis) and transverse momentum p_T (x -axis) in single z -vertex bin.



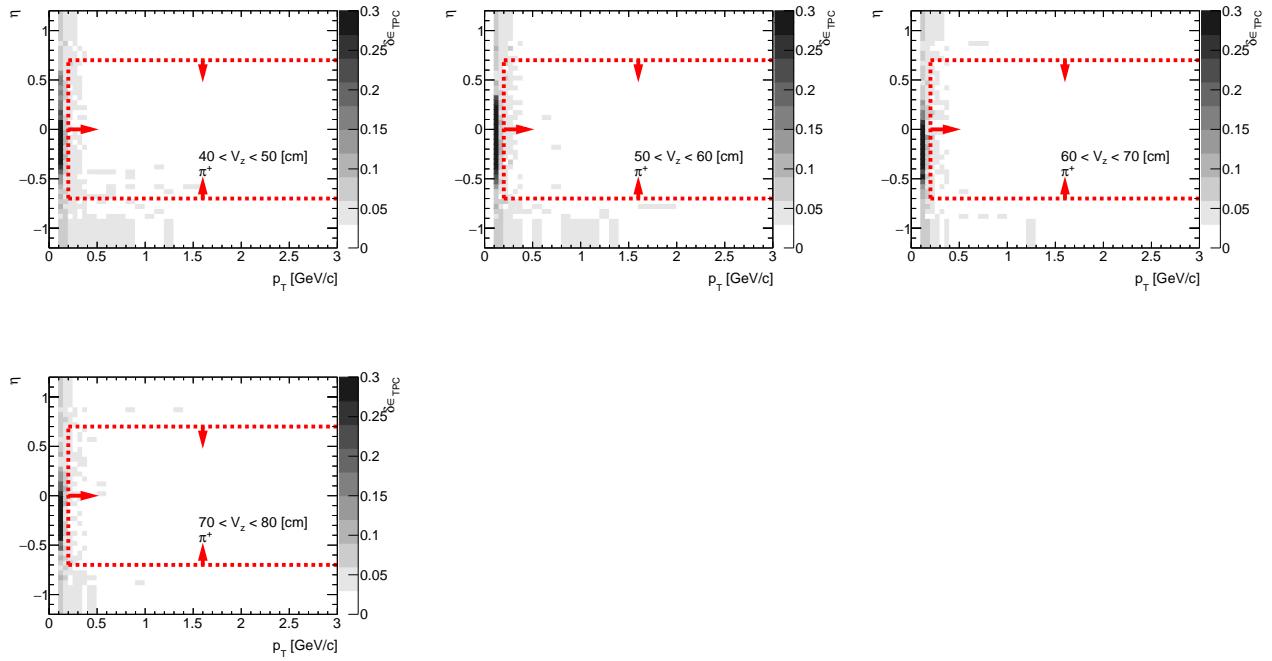


Figure D.2: The amount of lost K^- due to the interaction with dead material in front of TPC in CD MC sample. Each plot represents the fraction of lost K^- , $\delta\epsilon_{TPC}$ (z -axis), as a function of true particle pseudorapidity η (y -axis) and transverse momentum p_T (x -axis) in single z -vertex bin.

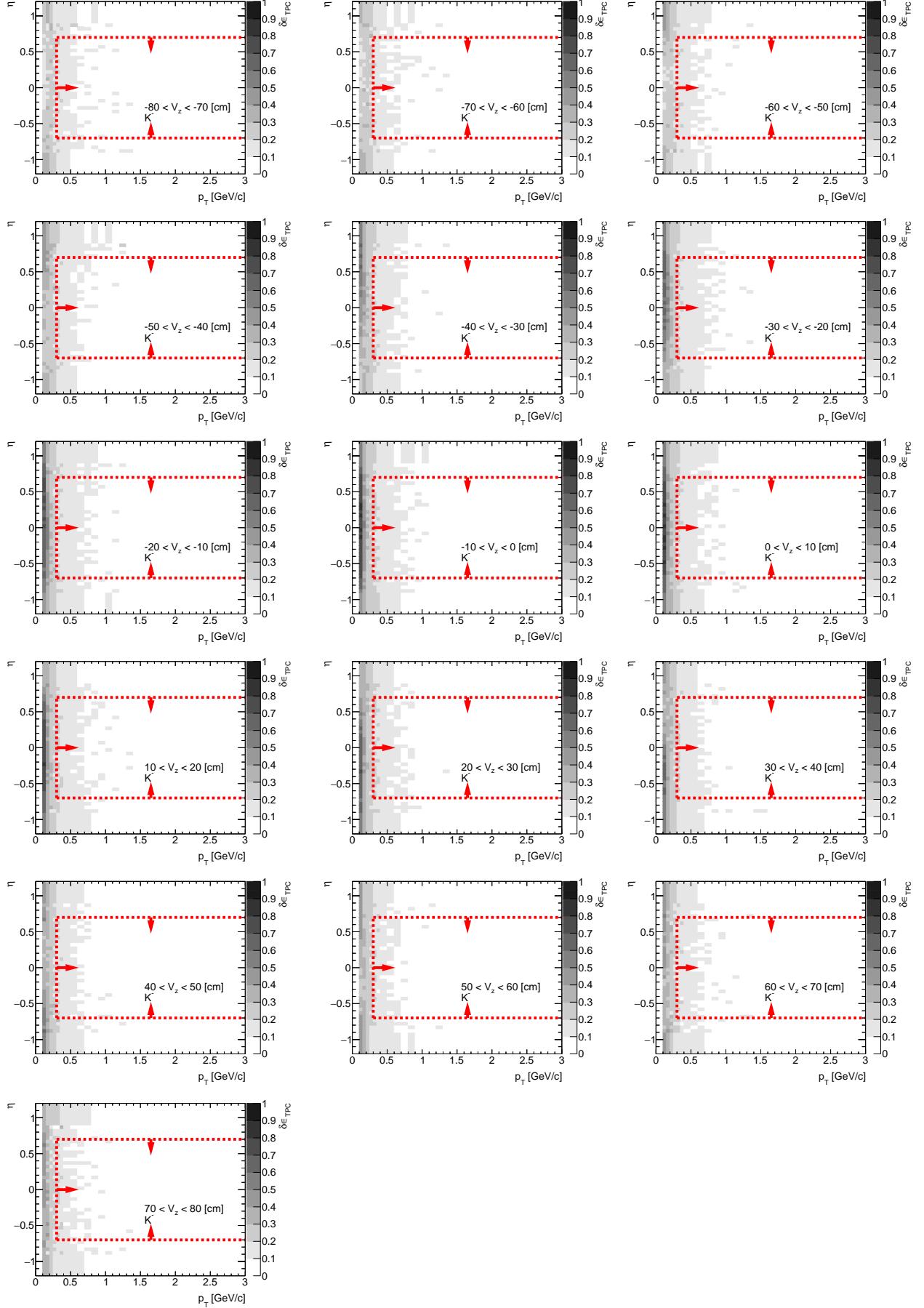


Figure D.3: The amount of lost K^+ due to the interaction with dead material in front of TPC in CD MC sample. Each plot represents the fraction of lost K^+ , $\delta\epsilon_{TPC}$ (z -axis), as a function of true particle pseudorapidity η (y -axis) and transverse momentum p_T (x -axis) in single z -vertex bin.

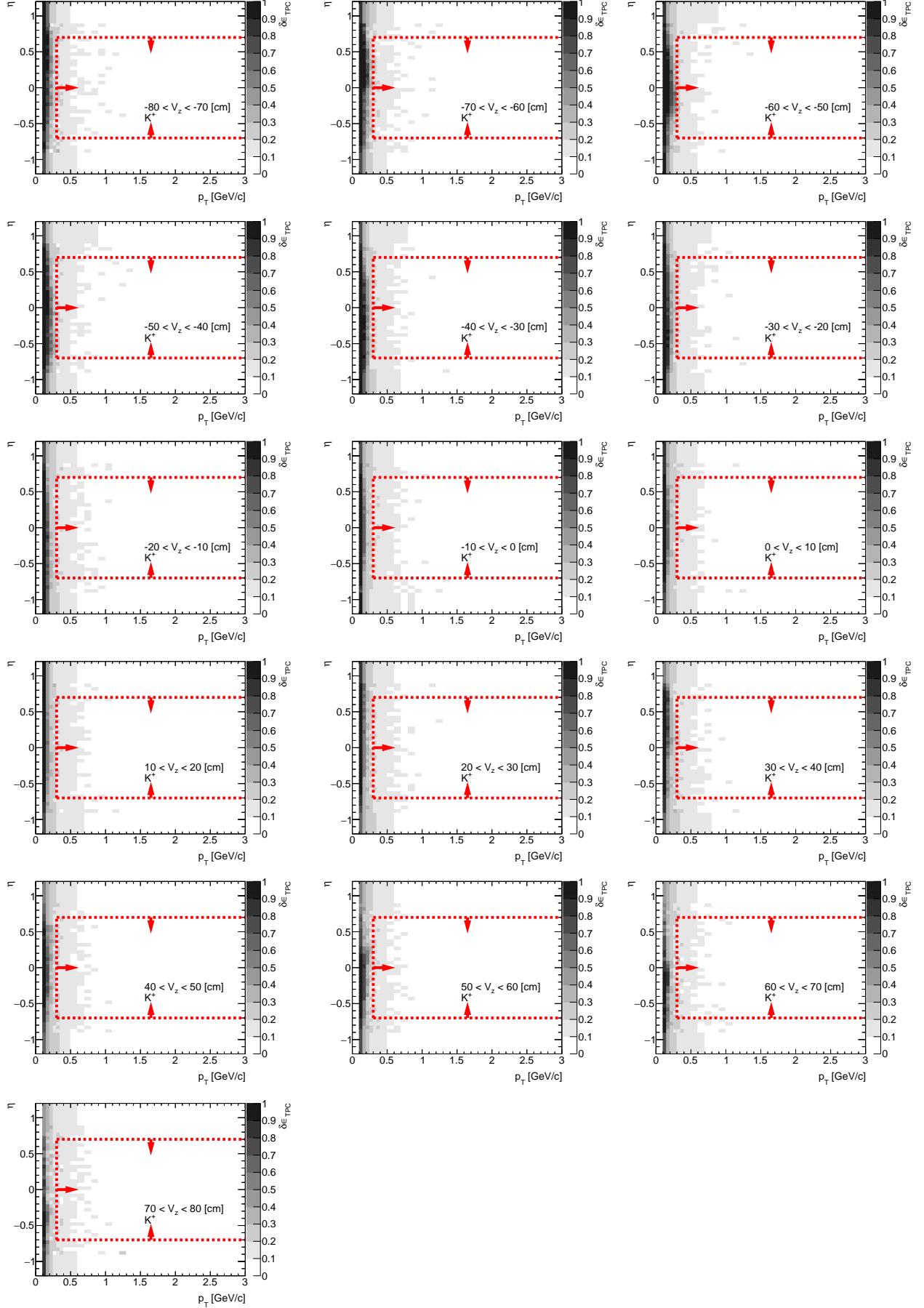


Figure D.4: The amount of lost \bar{p} due to the interaction with dead material in front of TPC in CD MC sample. Each plot represents the fraction of lost \bar{p} , $\delta\epsilon_{TPC}$ (z -axis), as a function of true particle pseudorapidity η (y -axis) and transverse momentum p_T (x -axis) in single z -vertex bin.

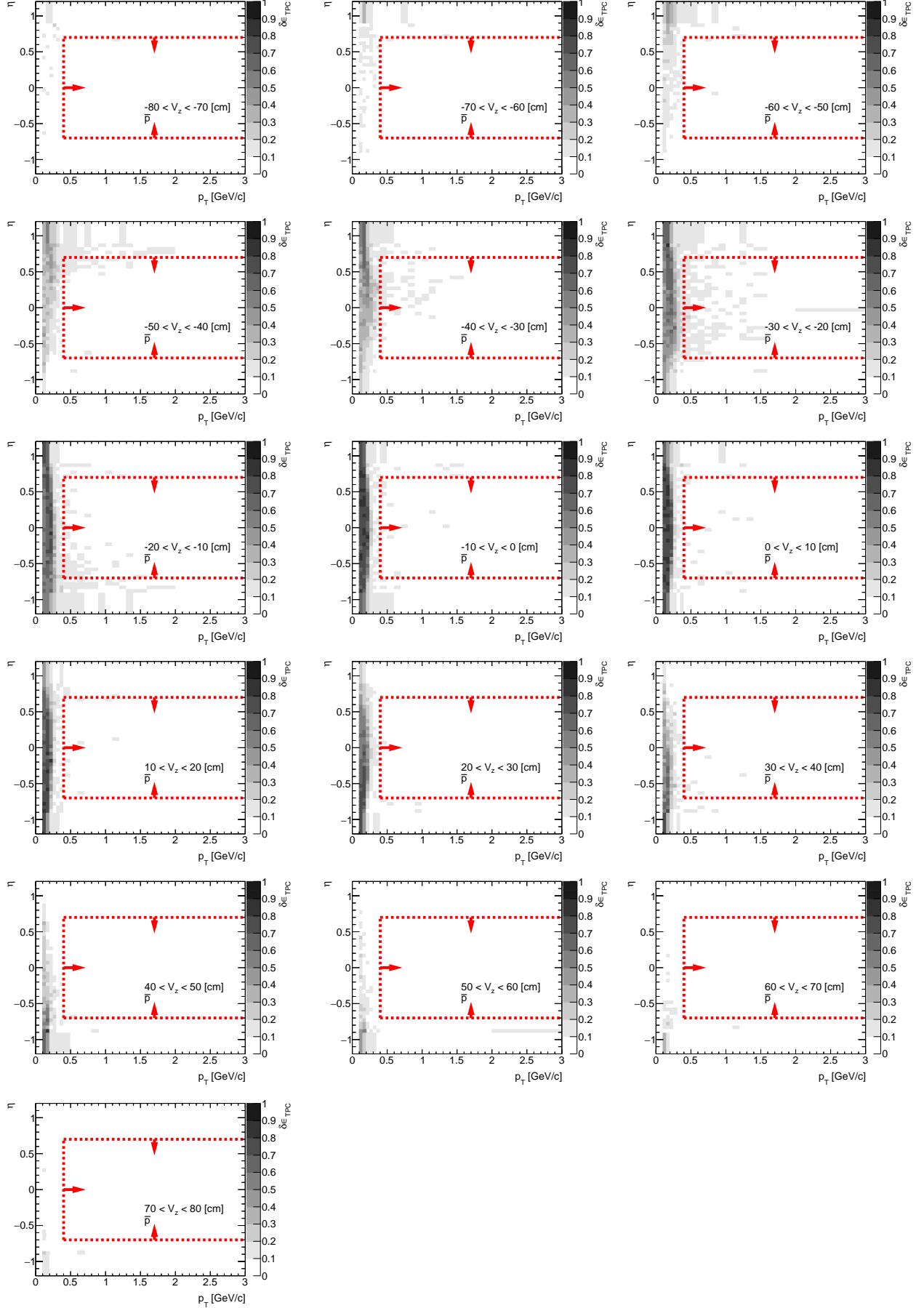


Figure D.5: The amount of lost p due to the interaction with dead material in front of TPC in CD MC sample. Each plot represents the fraction of lost p , $\delta\epsilon_{TPC}$ (z -axis), as a function of true particle pseudorapidity η (y -axis) and transverse momentum p_T (x -axis) in single z -vertex bin.

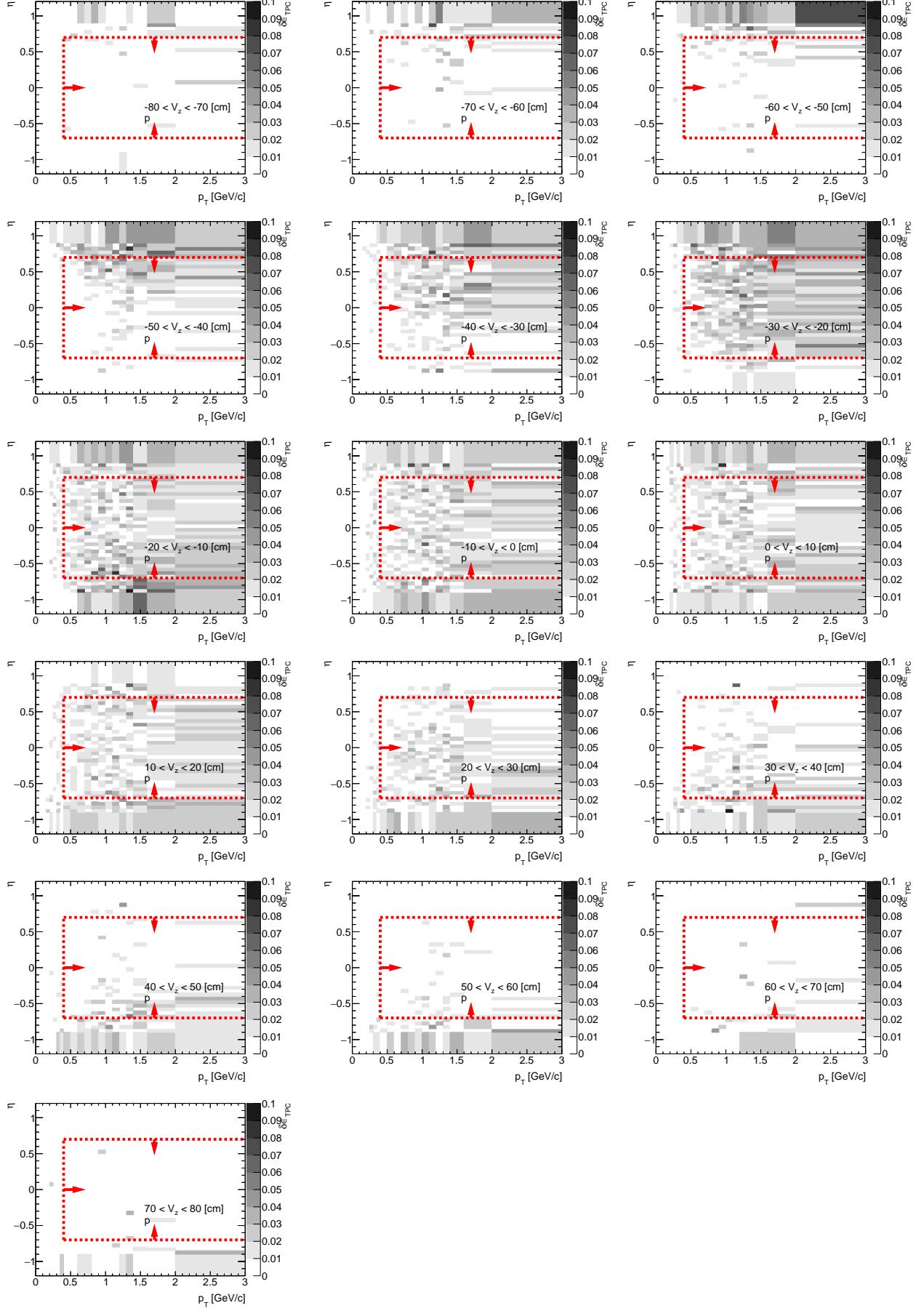


Figure D.6: The amount of lost negative particles due to the interaction with dead material in front of TPC in CD MC sample. Each plot represents the fraction of lost negative particles, $\delta\epsilon_{TPC}$ (z-axis), as a function of true particle pseudorapidity η (y-axis) and transverse momentum p_T (x-axis) in single z -vertex bin.

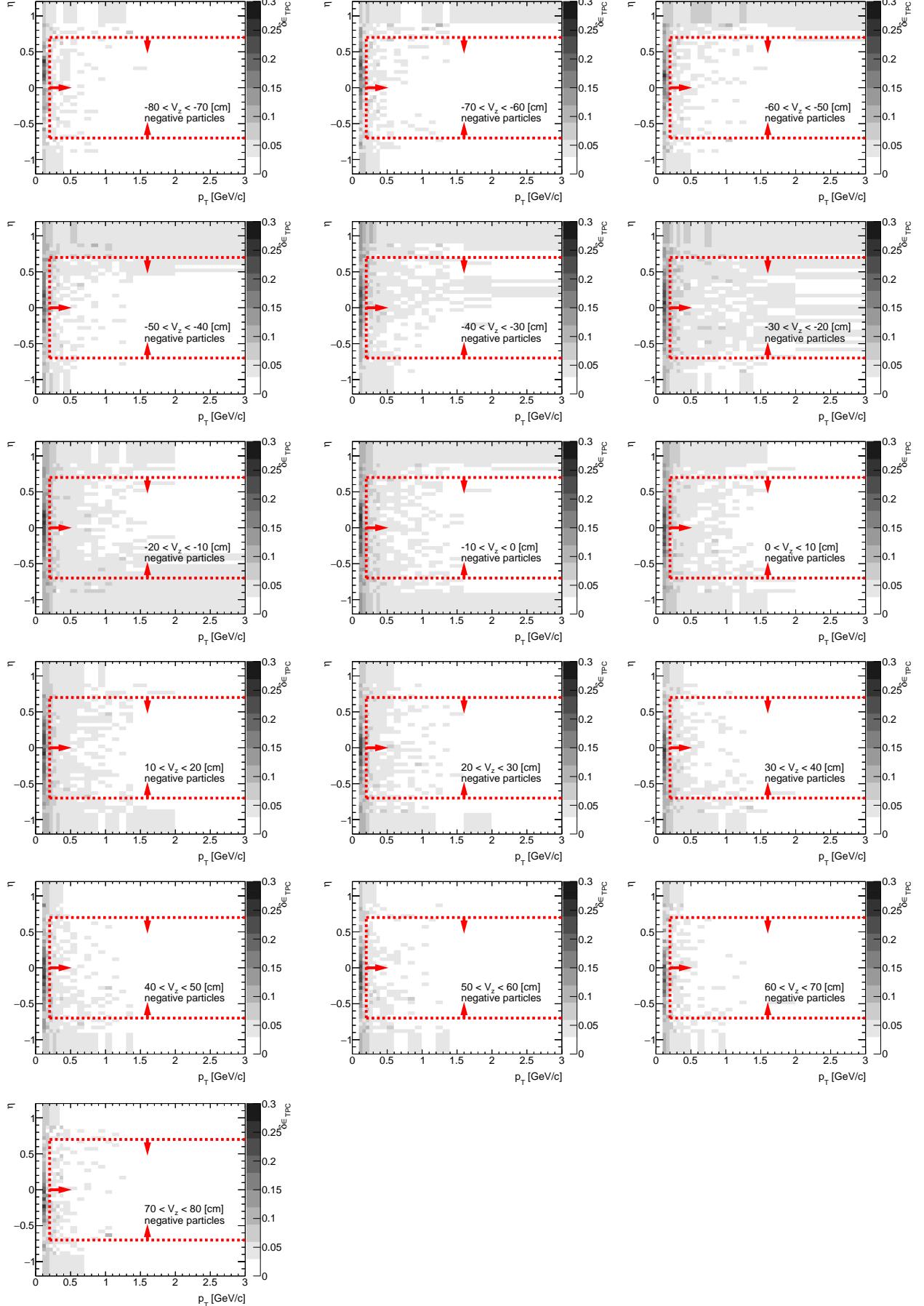


Figure D.7: The amount of lost positive particles due to the interaction with dead material in front of TPC in CD MC sample. Each plot represents the fraction of lost positive particles, $\delta\epsilon_{TPC}$ (z-axis), as a function of true particle pseudorapidity η (y-axis) and transverse momentum p_T (x-axis) in single z -vertex bin.

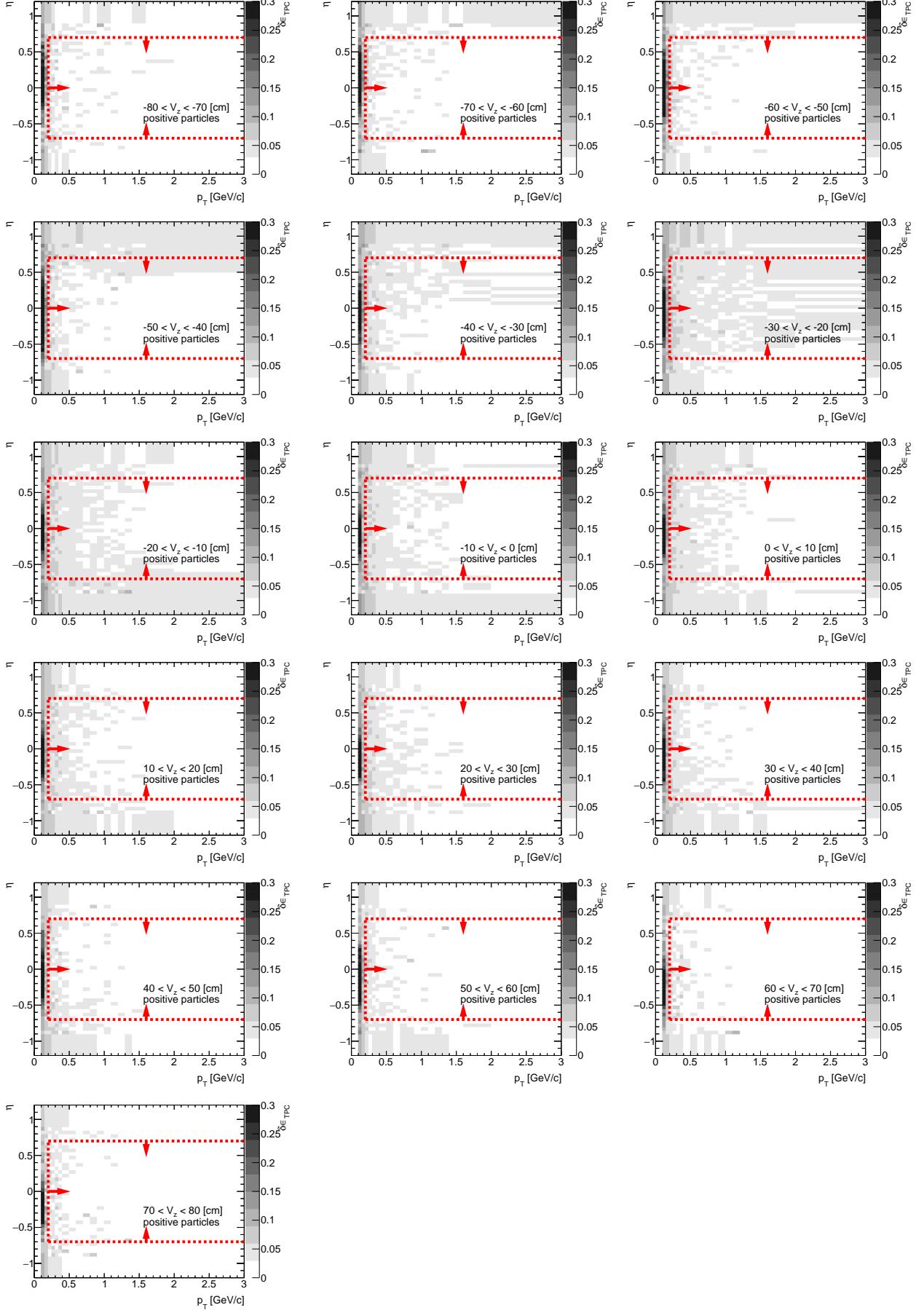


Figure D.8: The amount of lost π^- due to the interaction with dead material in front of TPC in SD MC sample. Each plot represents the fraction of lost π^- , $\delta\epsilon_{TPC}$ (z-axis), as a function of true particle pseudorapidity η (y-axis) and transverse momentum p_T (x-axis) in single z -vertex bin.

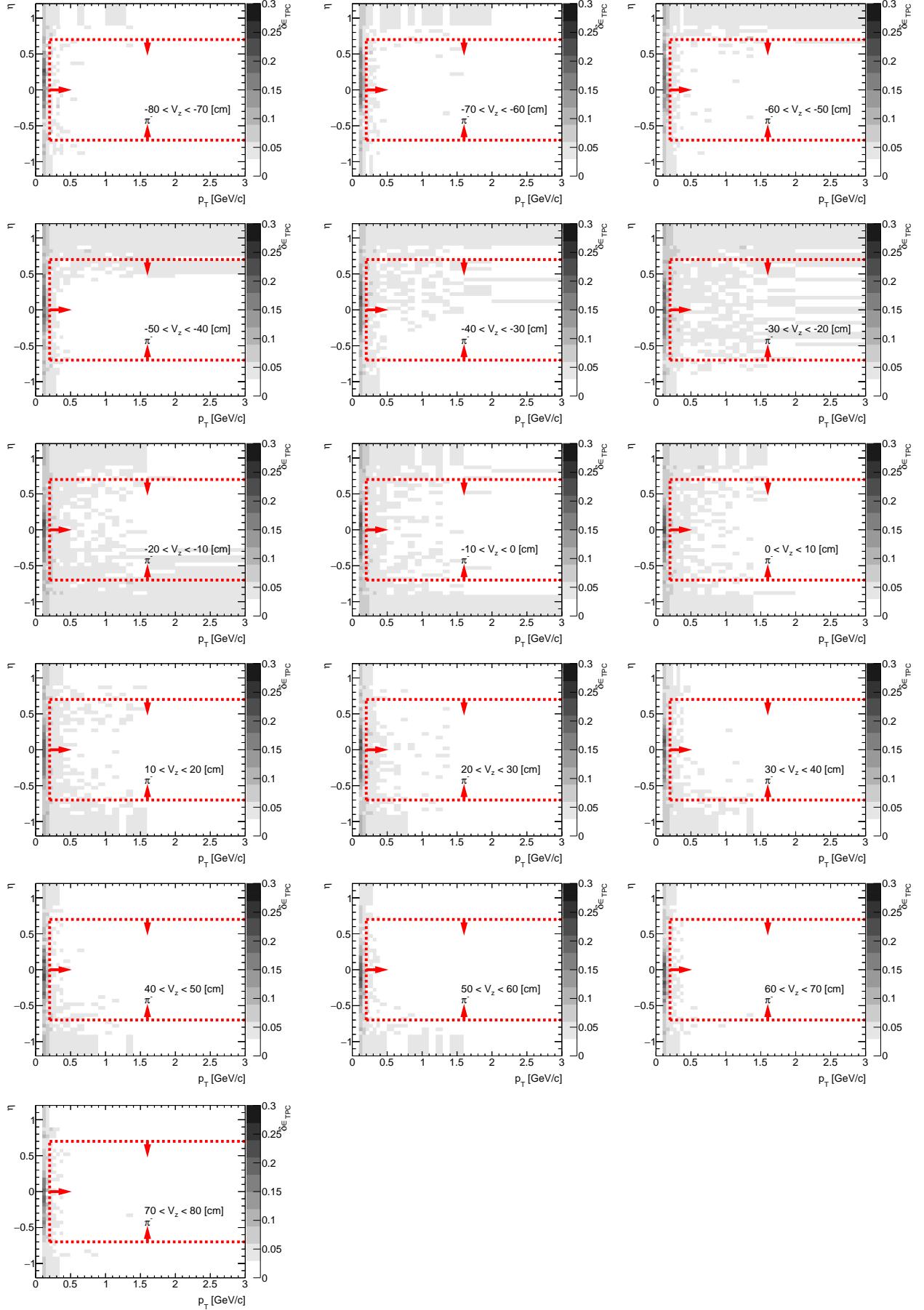


Figure D.9: The amount of lost π^+ due to the interaction with dead material in front of TPC in SD MC sample. Each plot represents the fraction of lost π^+ , $\delta\epsilon_{TPC}$ (z-axis), as a function of true particle pseudorapidity η (y-axis) and transverse momentum p_T (x-axis) in single z -vertex bin.

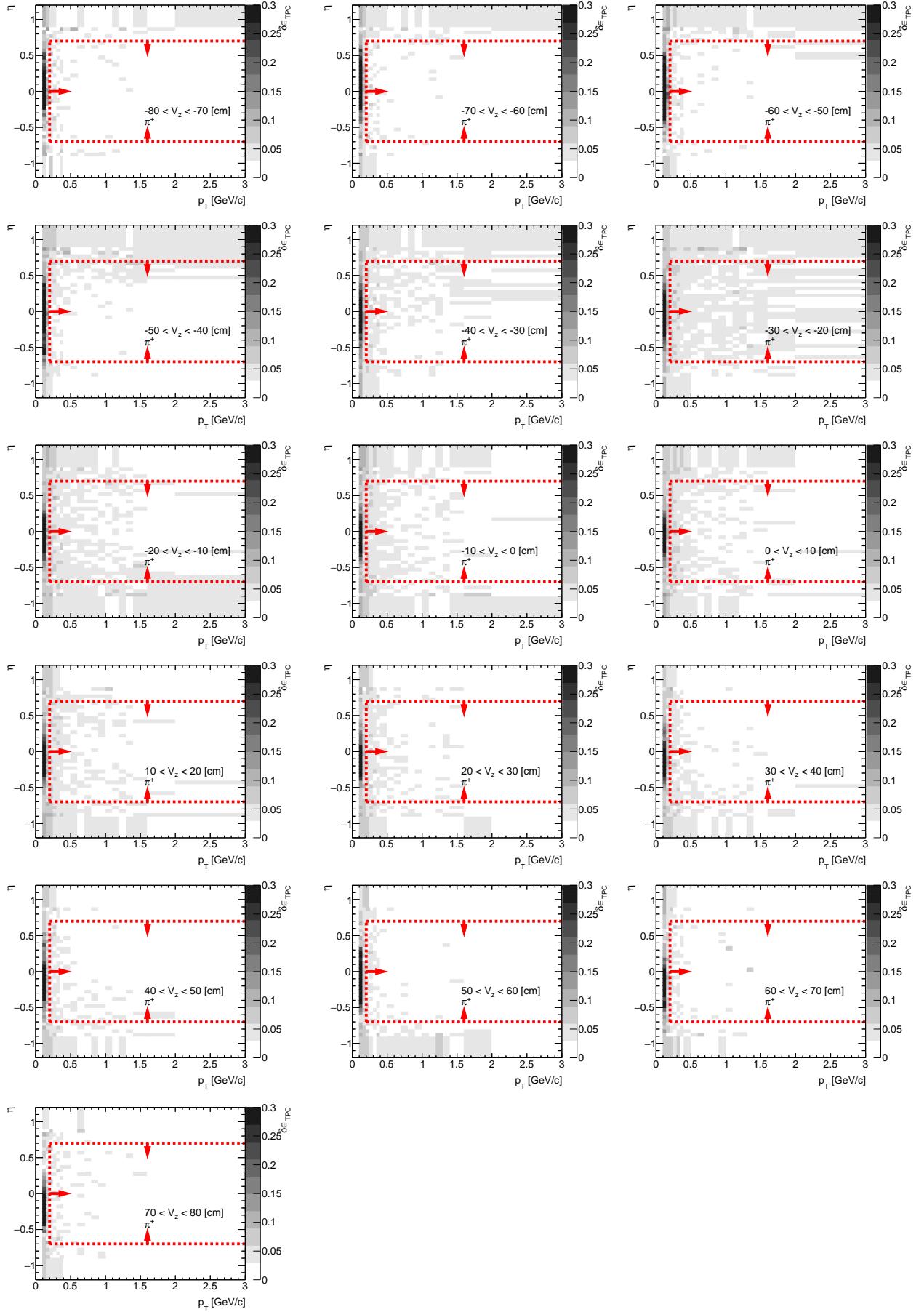


Figure D.10: The amount of lost K^- due to the interaction with dead material in front of TPC in SD MC sample. Each plot represents the fraction of lost K^- , $\delta\epsilon_{TPC}$ (z -axis), as a function of true particle pseudorapidity η (y -axis) and transverse momentum p_T (x -axis) in single z -vertex bin.

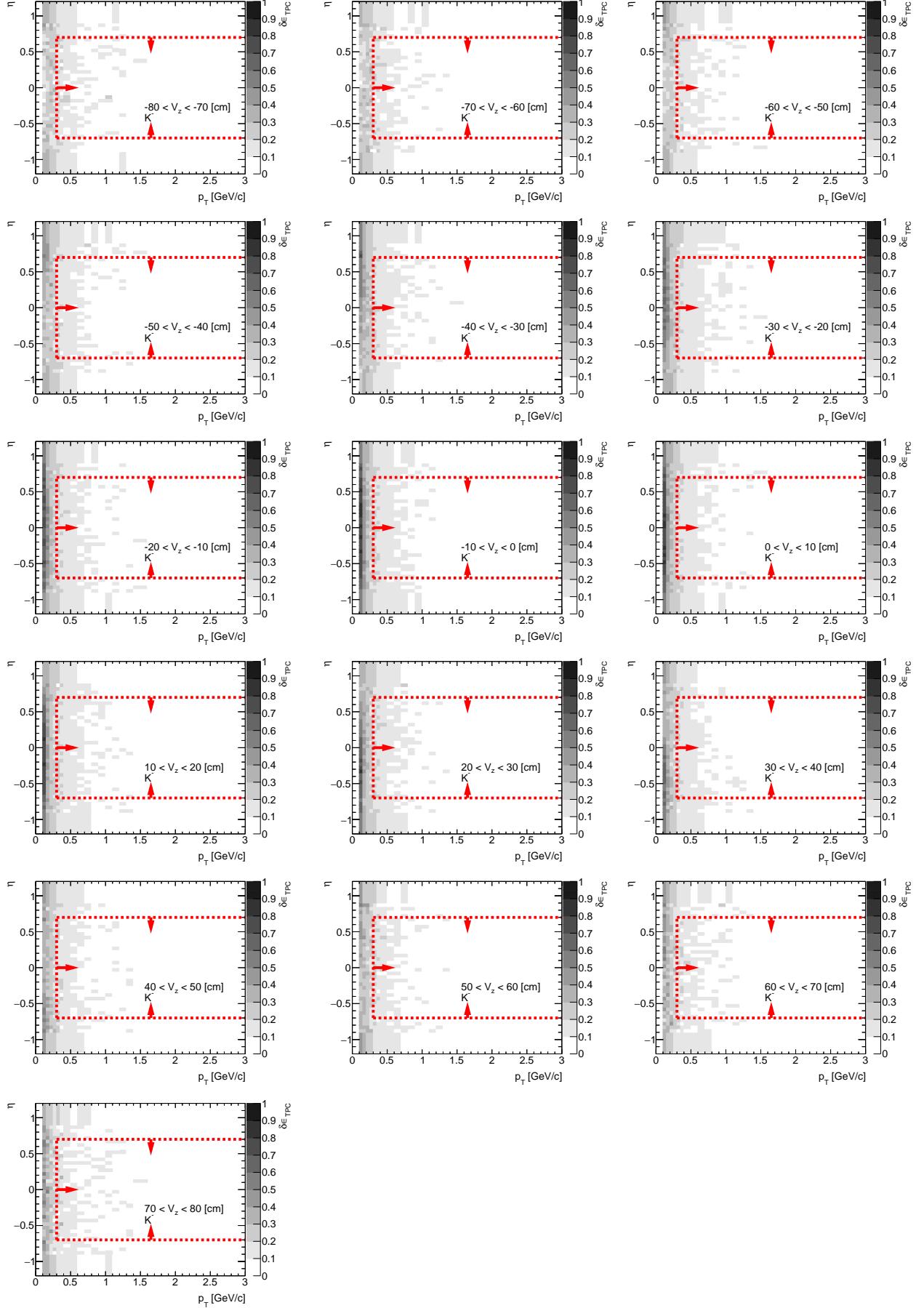


Figure D.11: The amount of lost K^+ due to the interaction with dead material in front of TPC in SD MC sample. Each plot represents the fraction of lost K^+ , $\delta\epsilon_{TPC}$ (z -axis), as a function of true particle pseudorapidity η (y -axis) and transverse momentum p_T (x -axis) in single z -vertex bin.

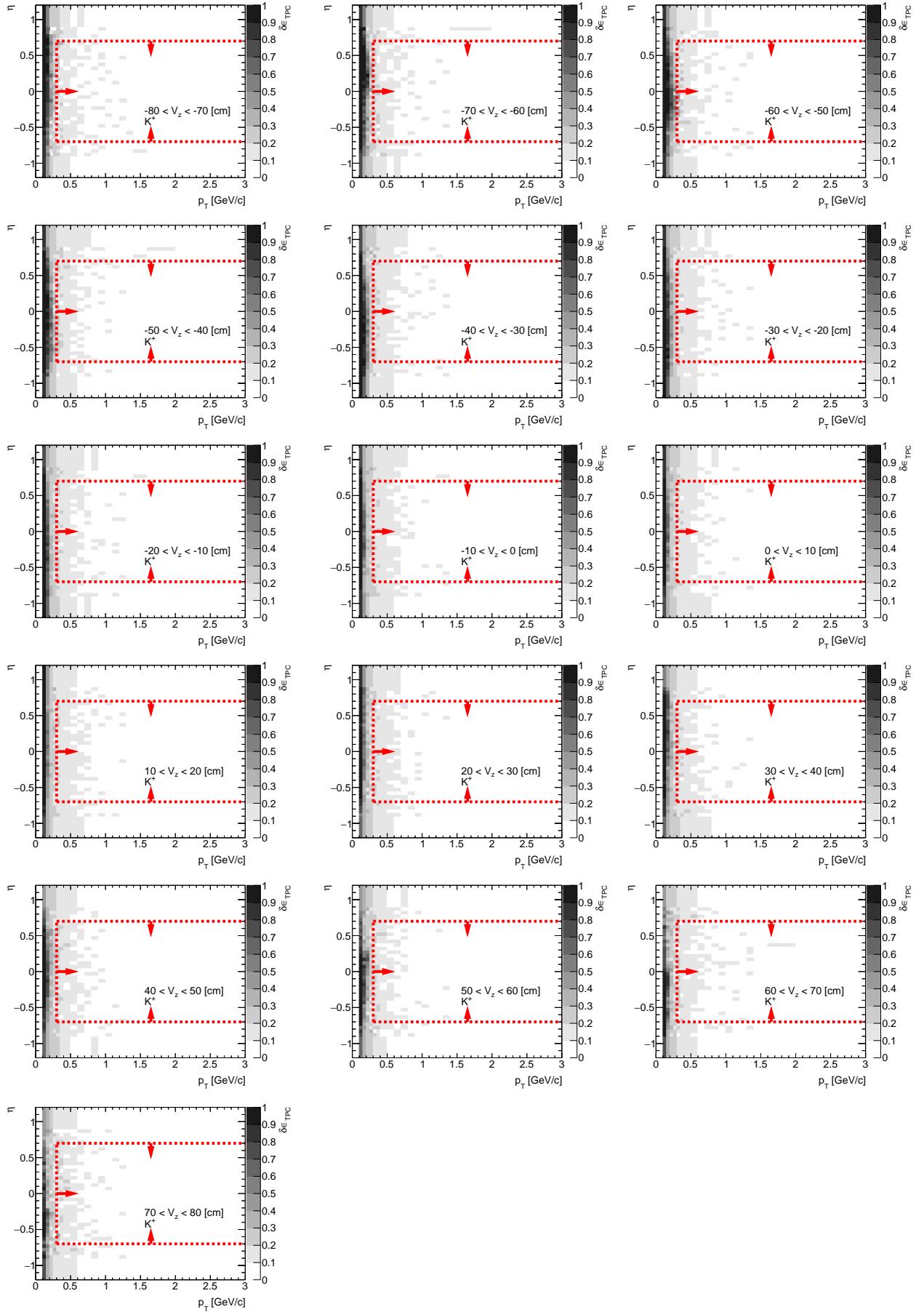


Figure D.12: The amount of lost \bar{p} due to the interaction with dead material in front of TPC in SD MC sample. Each plot represents the fraction of lost \bar{p} , $\delta\epsilon_{TPC}$ (z -axis), as a function of true particle pseudorapidity η (y -axis) and transverse momentum p_T (x -axis) in single z -vertex bin.

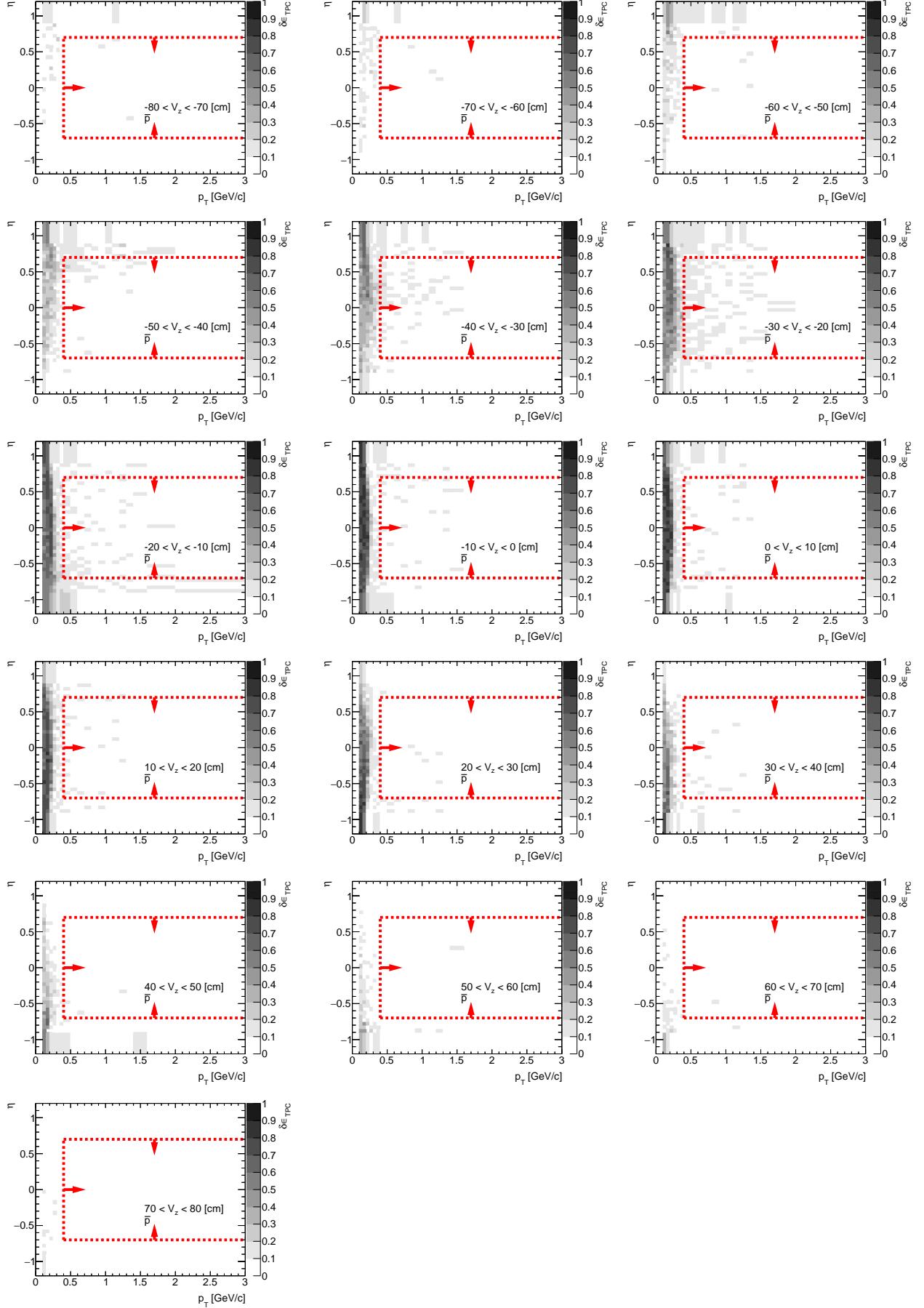


Figure D.13: The amount of lost p due to the interaction with dead material in front of TPC in SD MC sample. Each plot represents the fraction of lost p , $\delta\epsilon_{TPC}$ (z -axis), as a function of true particle pseudorapidity η (y -axis) and transverse momentum p_T (x -axis) in single z -vertex bin.

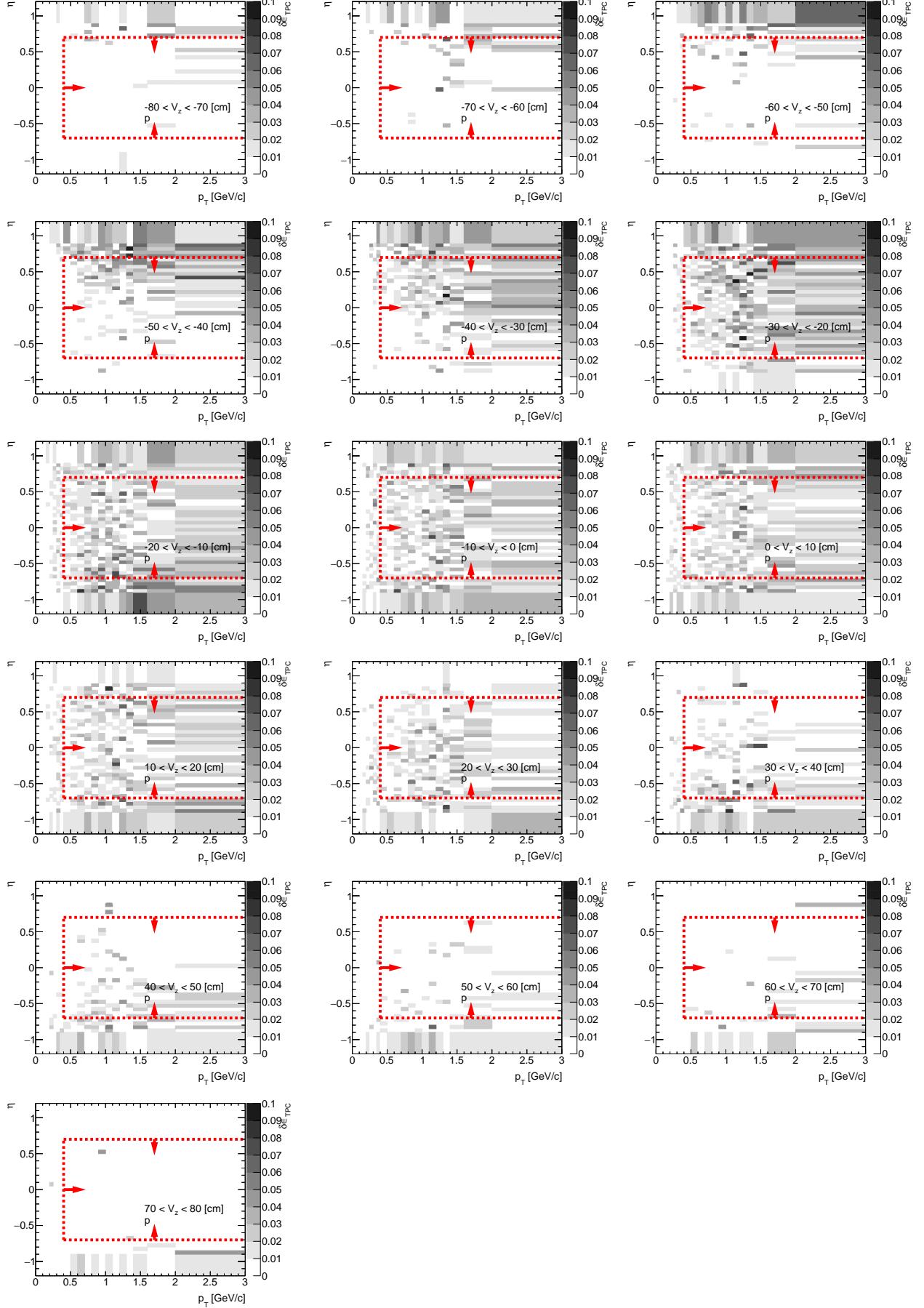


Figure D.14: The amount of lost negative particles due to the interaction with dead material in front of TPC in SD MC sample. Each plot represents the fraction of lost negative particles, $\delta\epsilon_{TPC}$ (z-axis), as a function of true particle pseudorapidity η (y-axis) and transverse momentum p_T (x-axis) in single z -vertex bin.

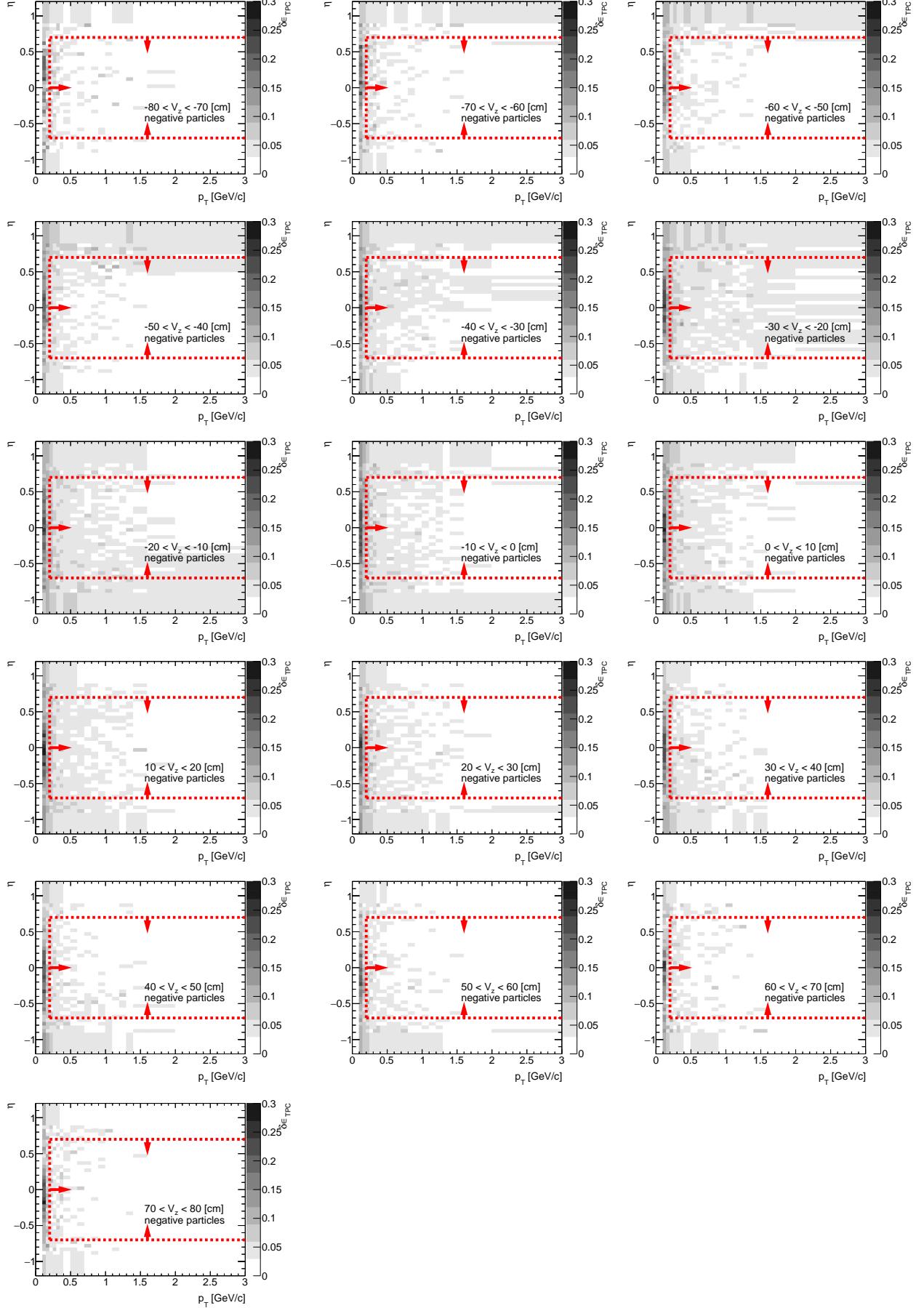


Figure D.15: The amount of lost positive particles due to the interaction with dead material in front of TPC in SD MC sample. Each plot represents the fraction of lost positive particles, $\delta\epsilon_{TPC}$ (z-axis), as a function of true particle pseudorapidity η (y-axis) and transverse momentum p_T (x-axis) in single z -vertex bin.

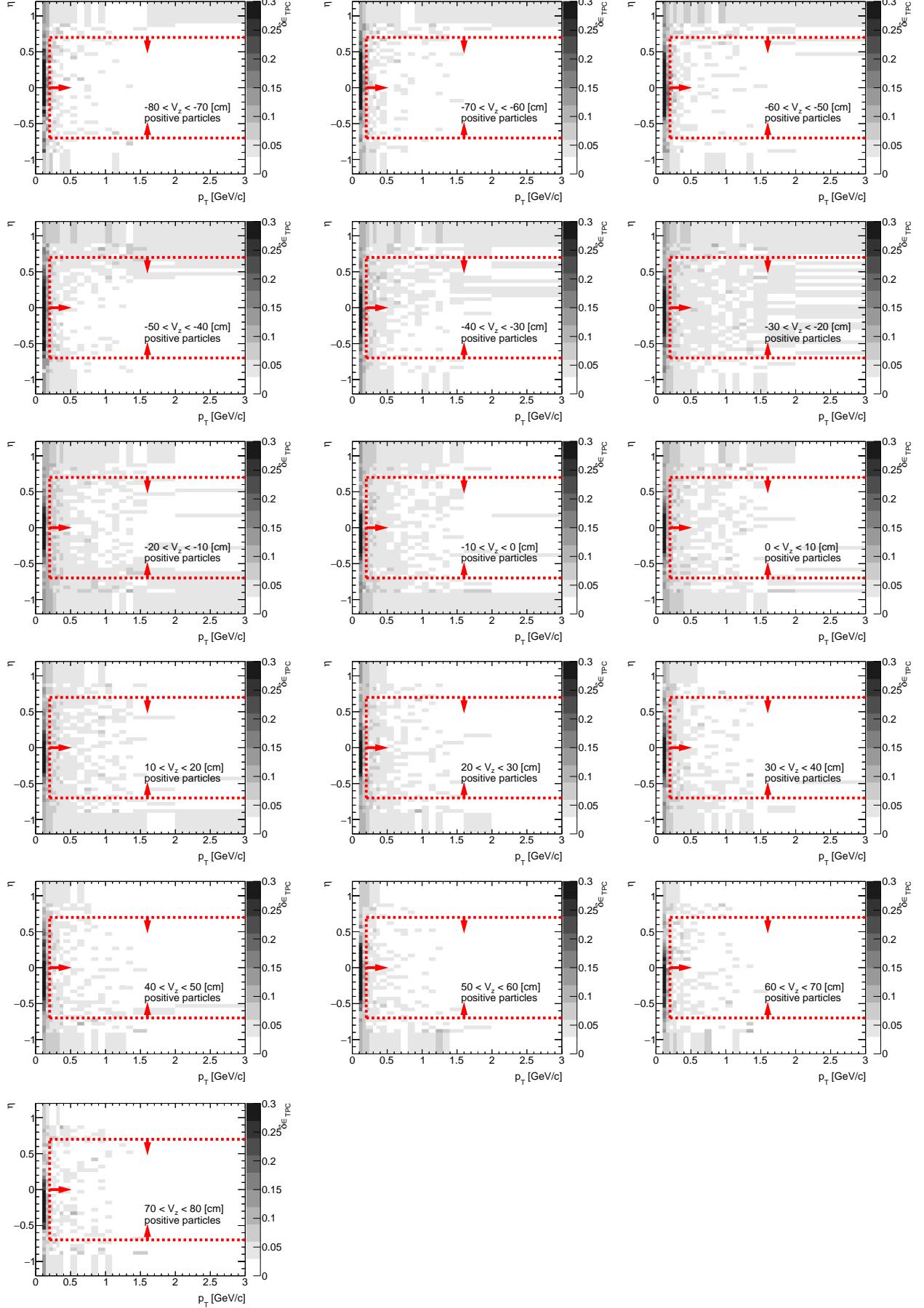
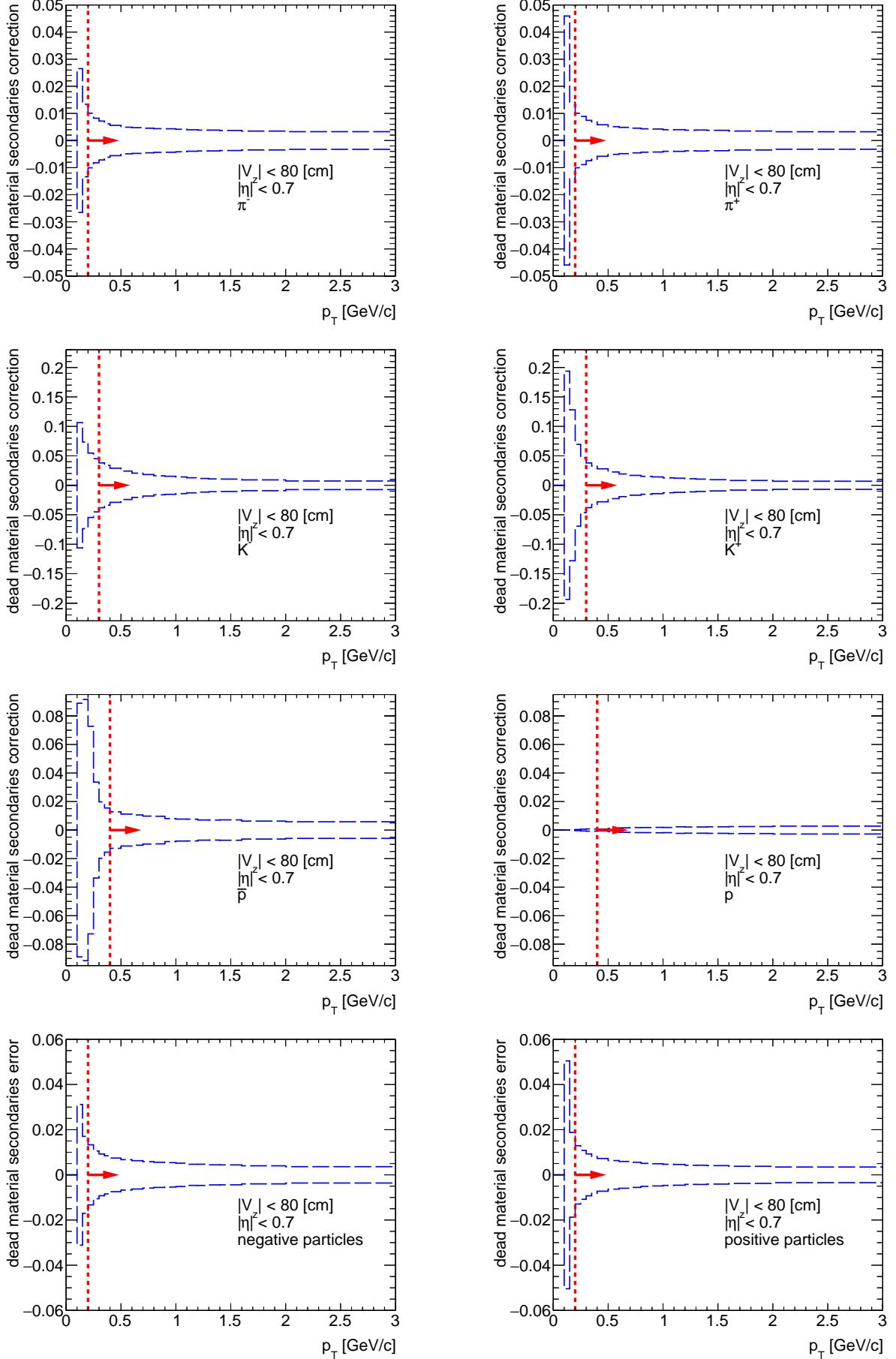


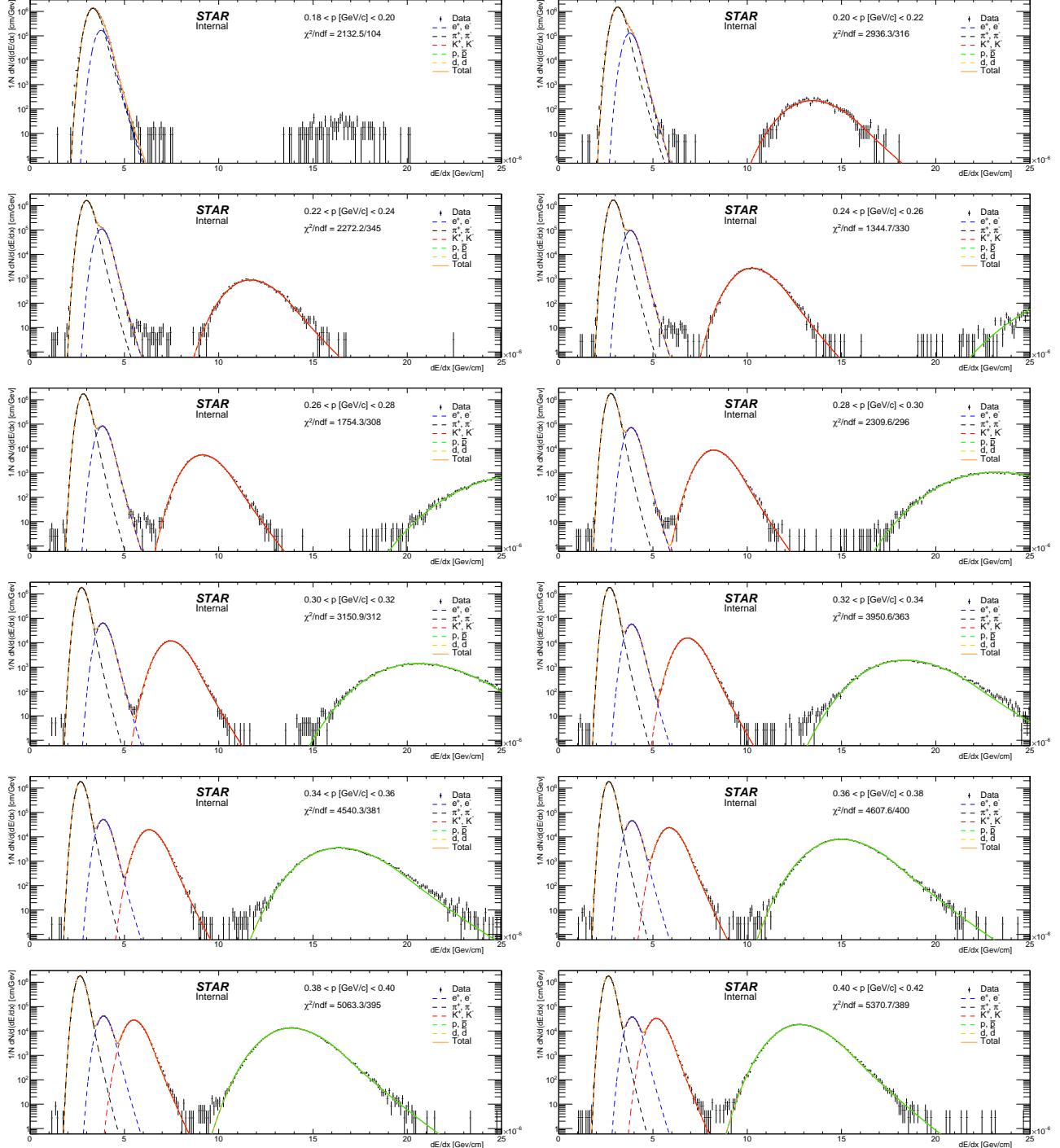
Figure D.16: The systematic uncertainty to the TPC track reconstruction efficiency due to amount of dead material in front of TPC using MC samples for SD. Each plot represents the systematic uncertainty as a function of true particle p_T ($|\eta| < 0.7, |V_z| < 80 \text{ cm}$) for given particle species: $\pi^-, \pi^+, K^-, K^+, \bar{p}$ and p . It was also calculated for negative and positive particles without identification.

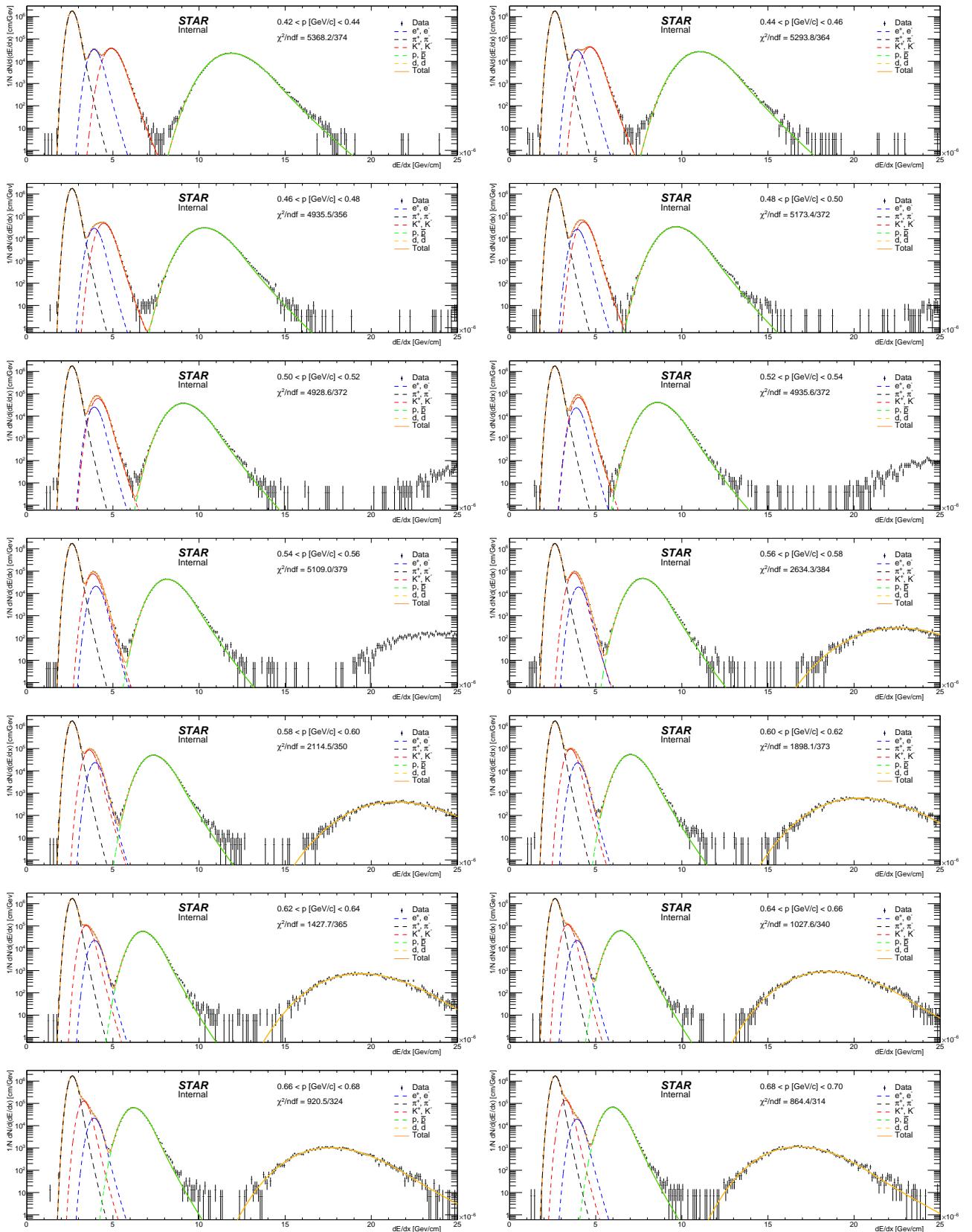


Appendix E

Fits to dE/dx spectra, comparison of dE/dx and n_X^σ between data and MC

Figure E.1: Fits of sum of functions from Eq. (5.6) corresponding to different particle species to dE/dx spectra in the data in momentum bins.





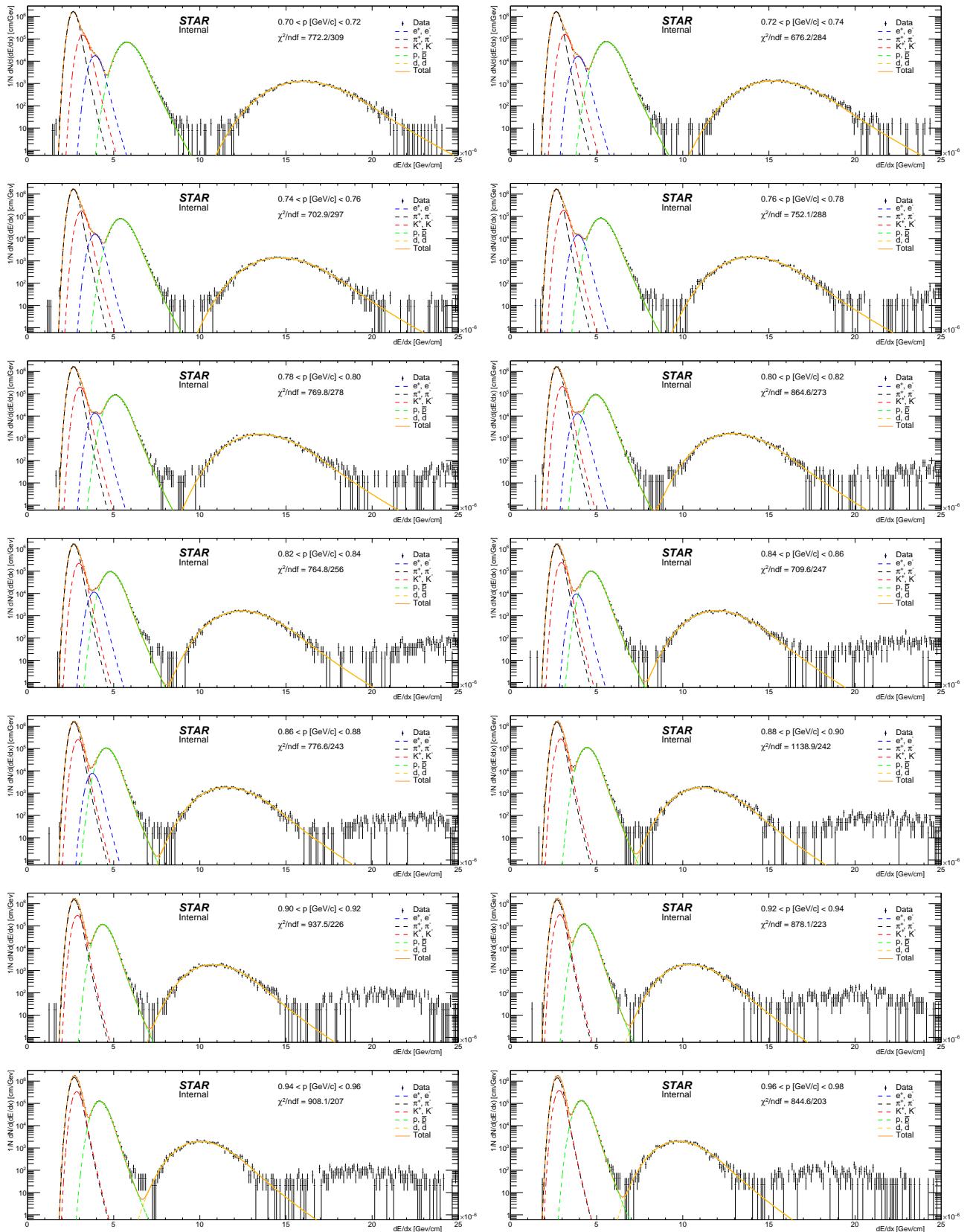
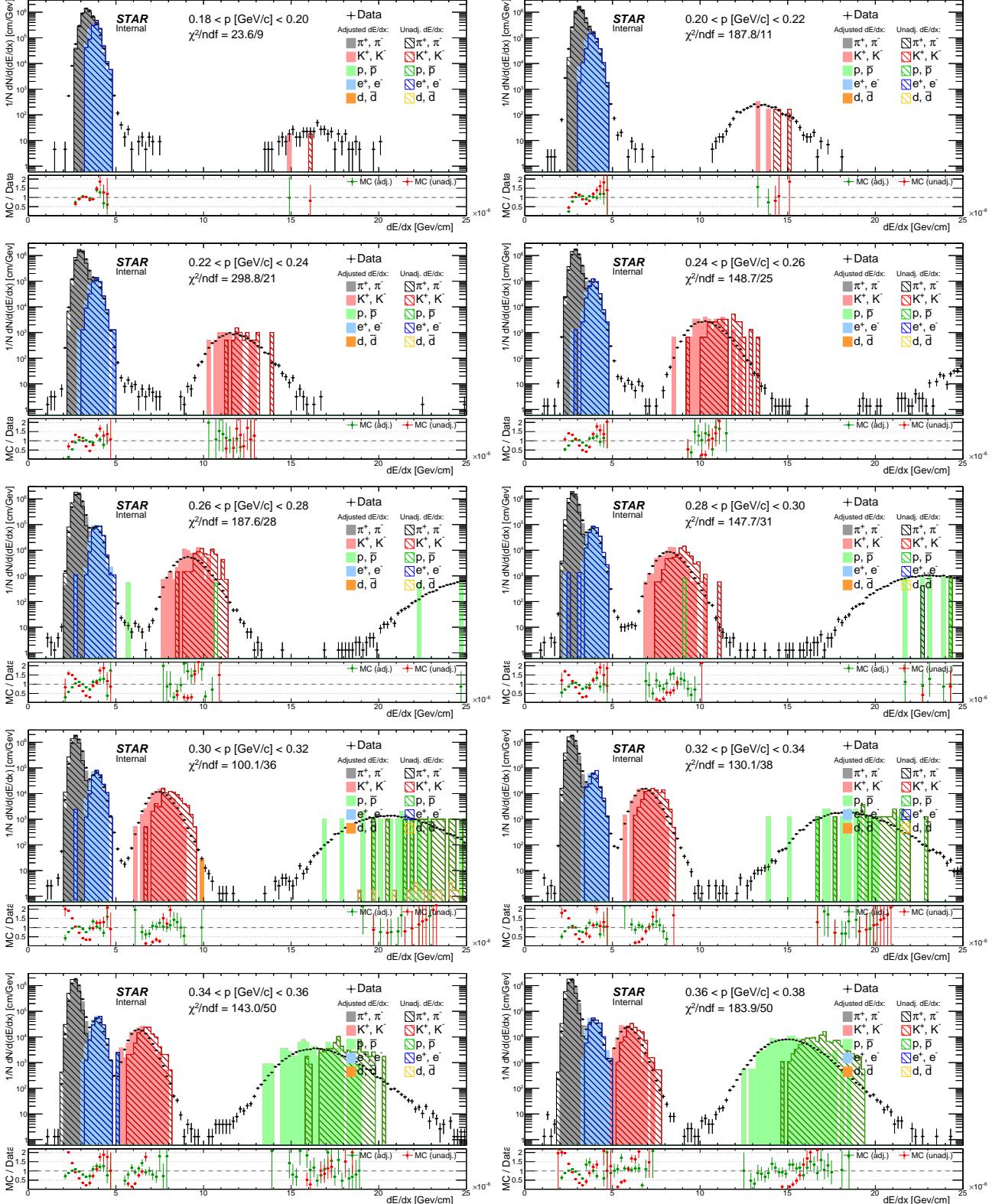
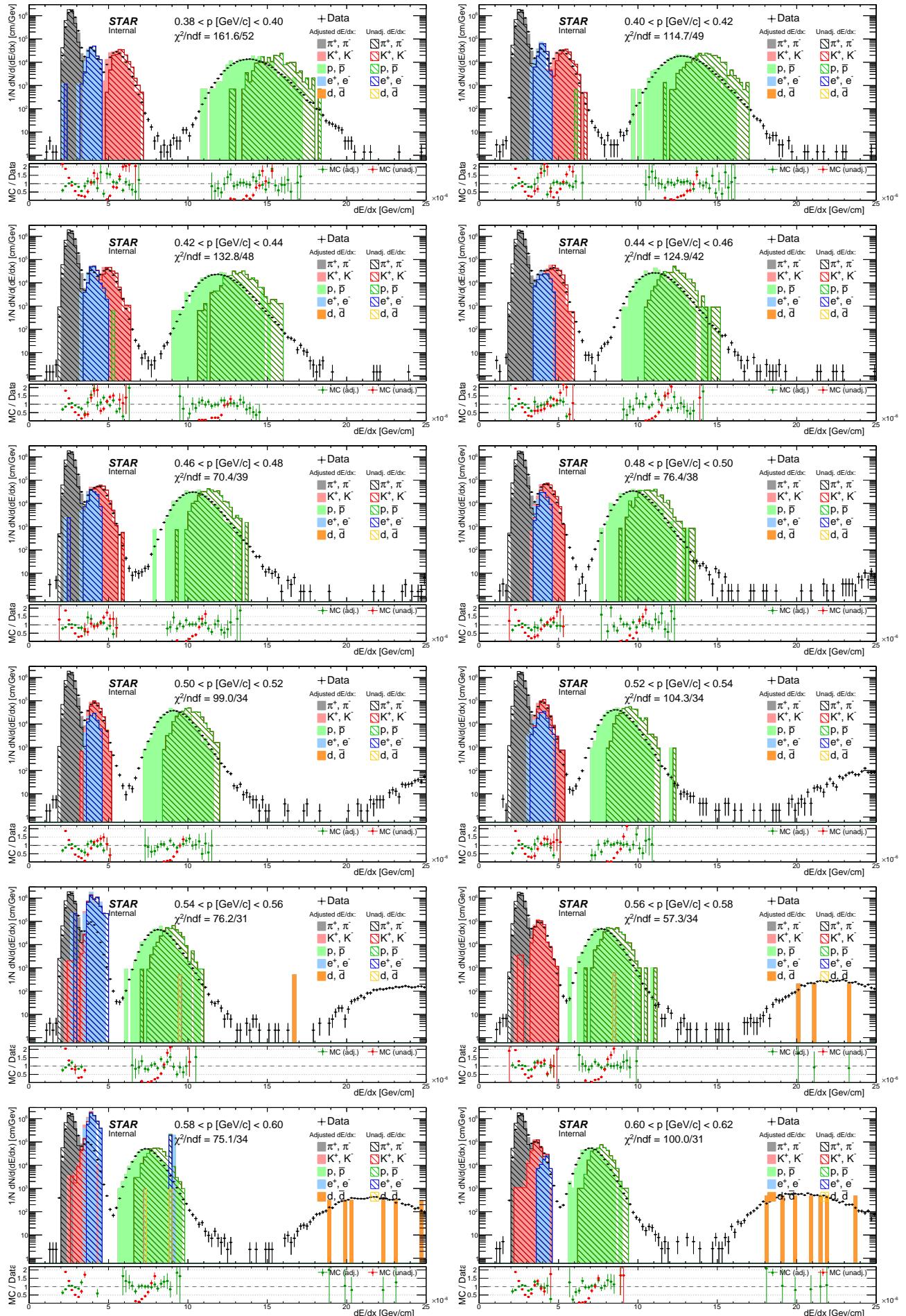
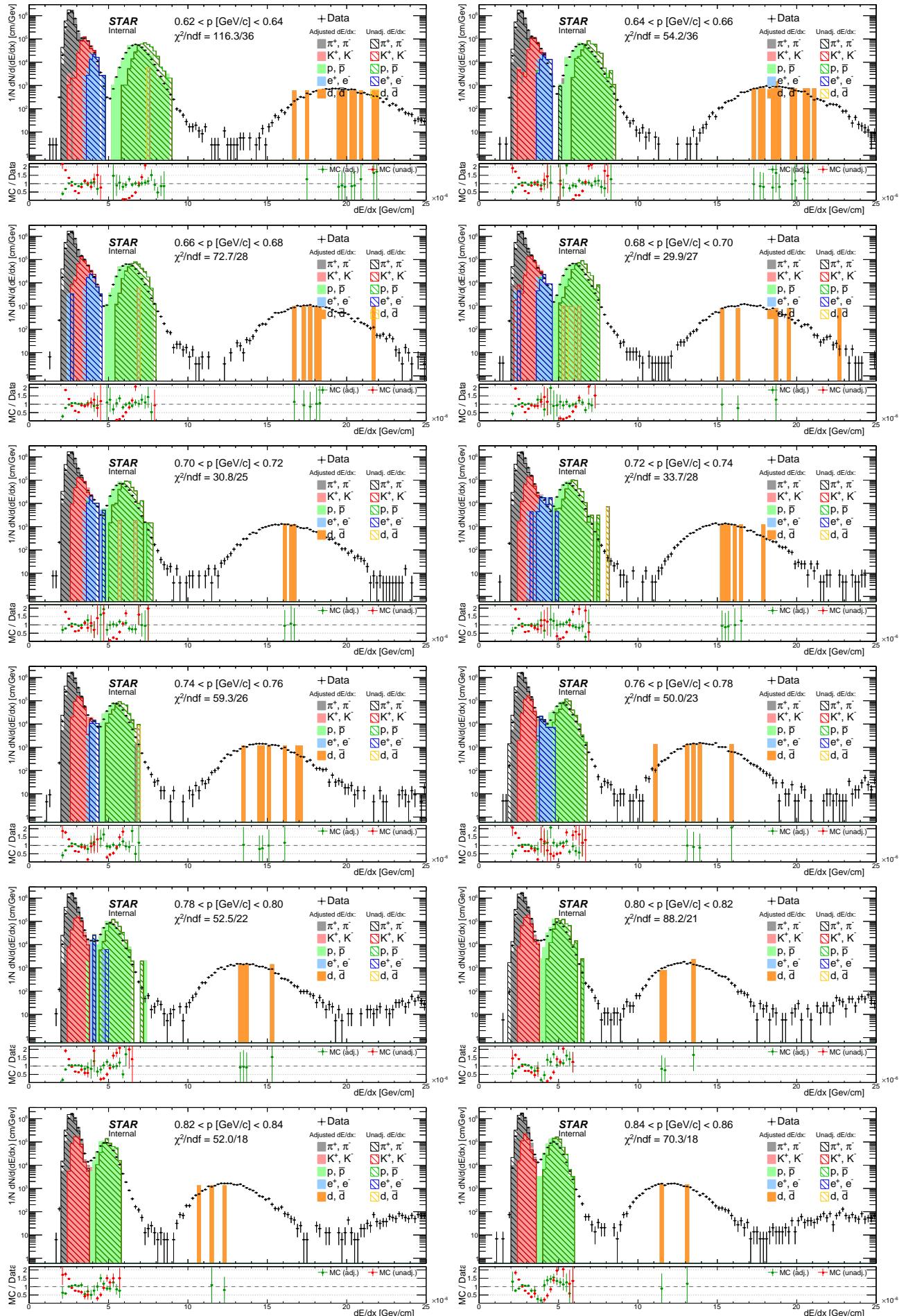


Figure E.2: Comparison of dE/dx spectrum between data and embedded MC (before and after dE/dx adjustment) in momentum bins.







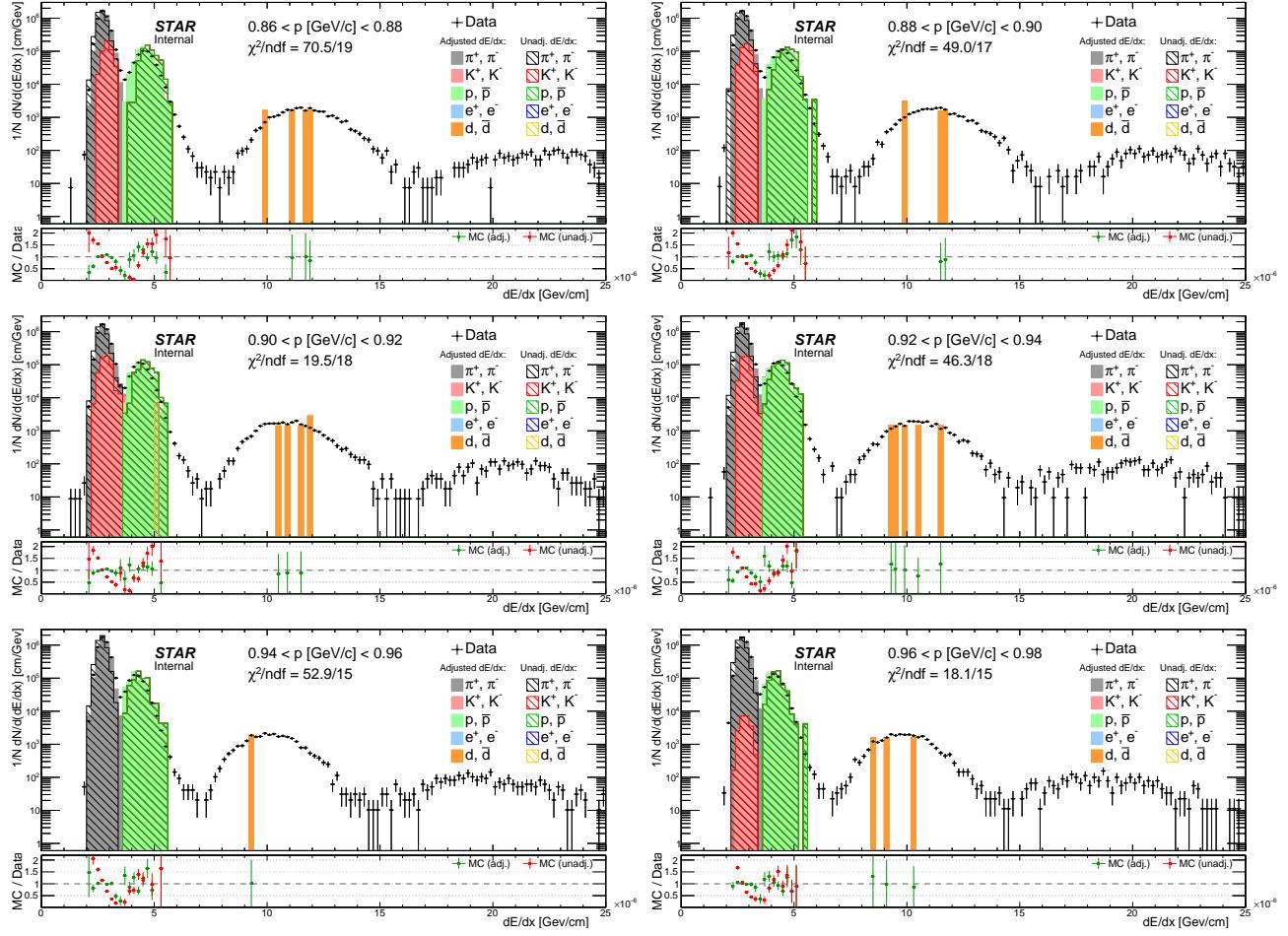
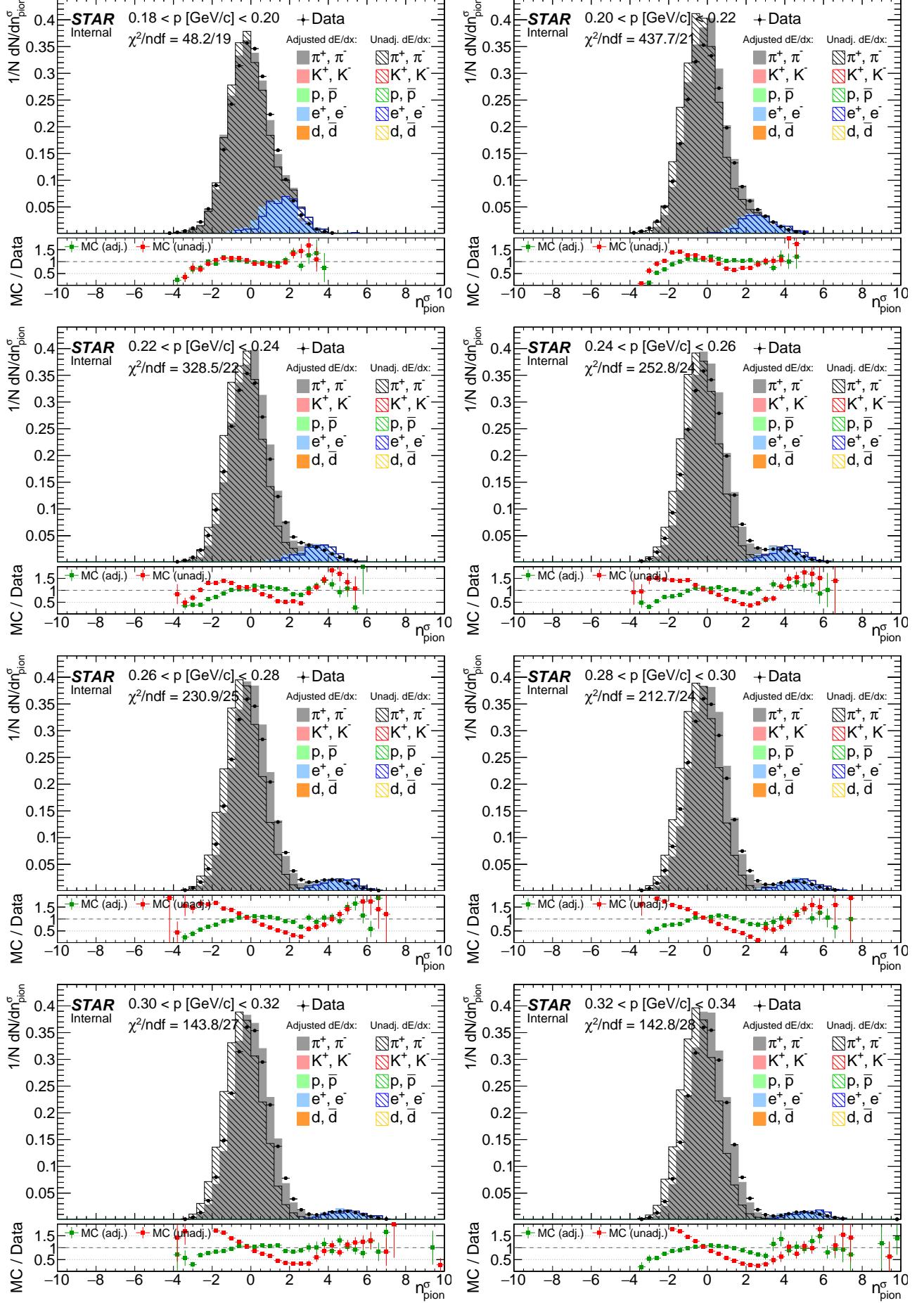
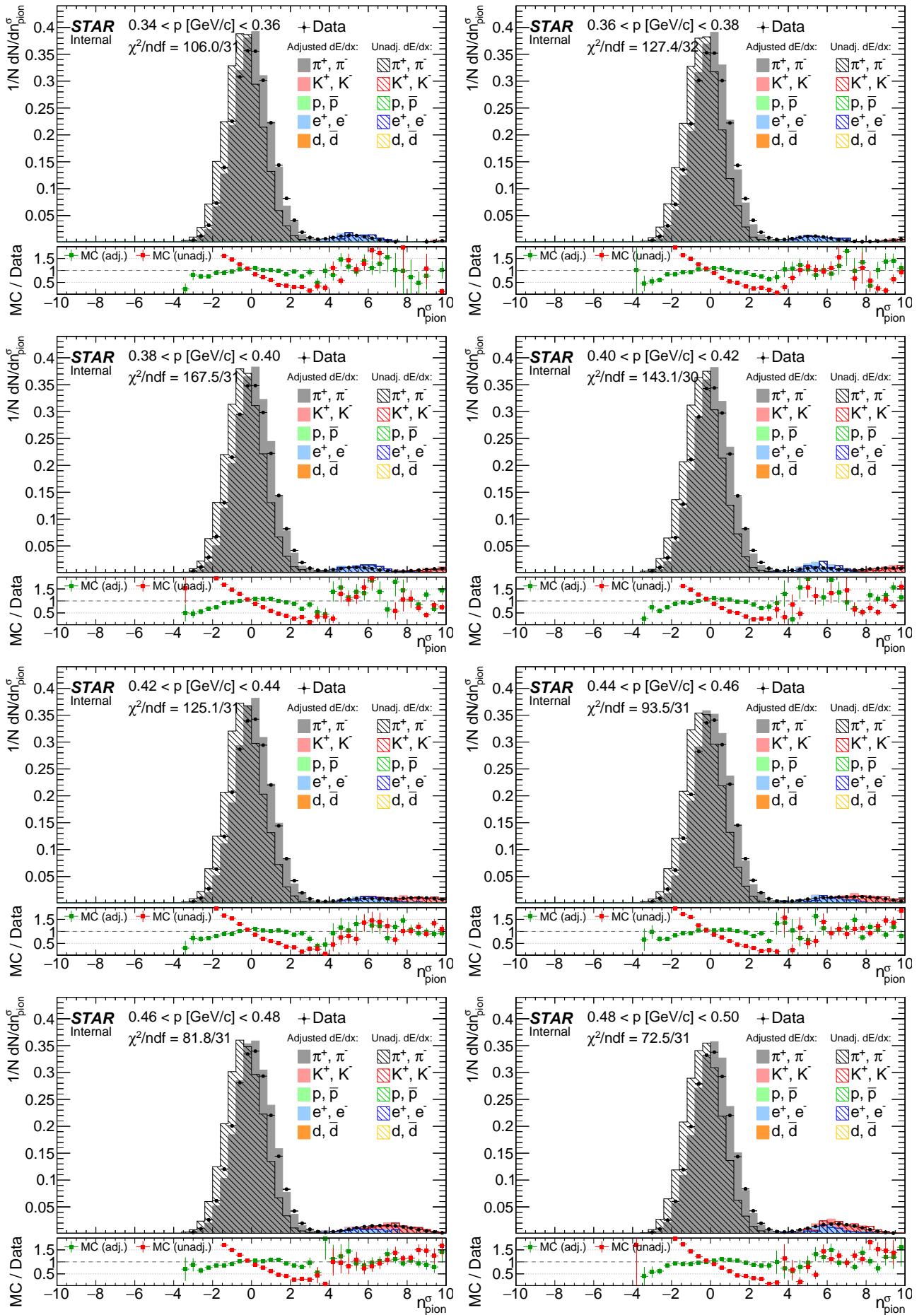
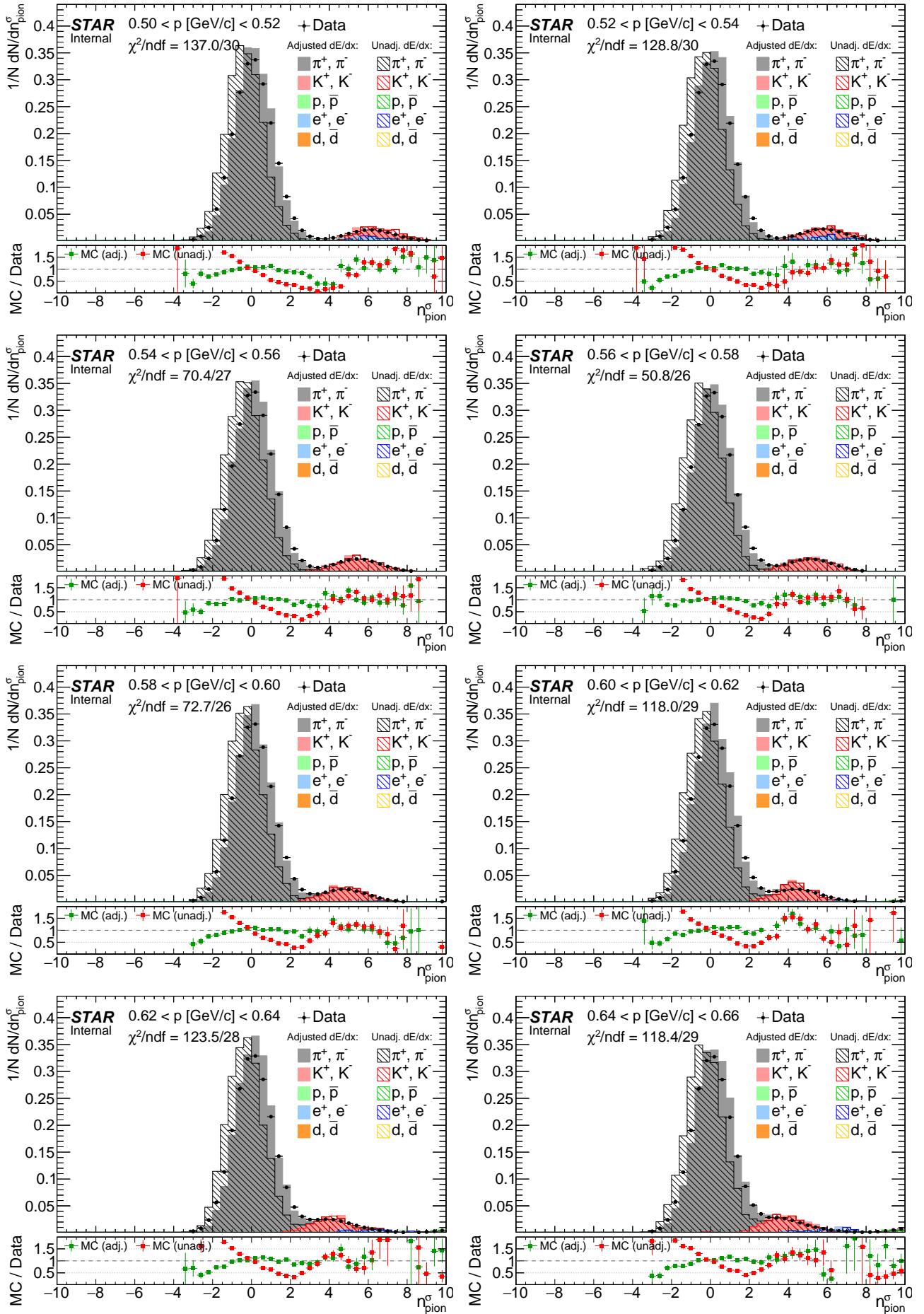
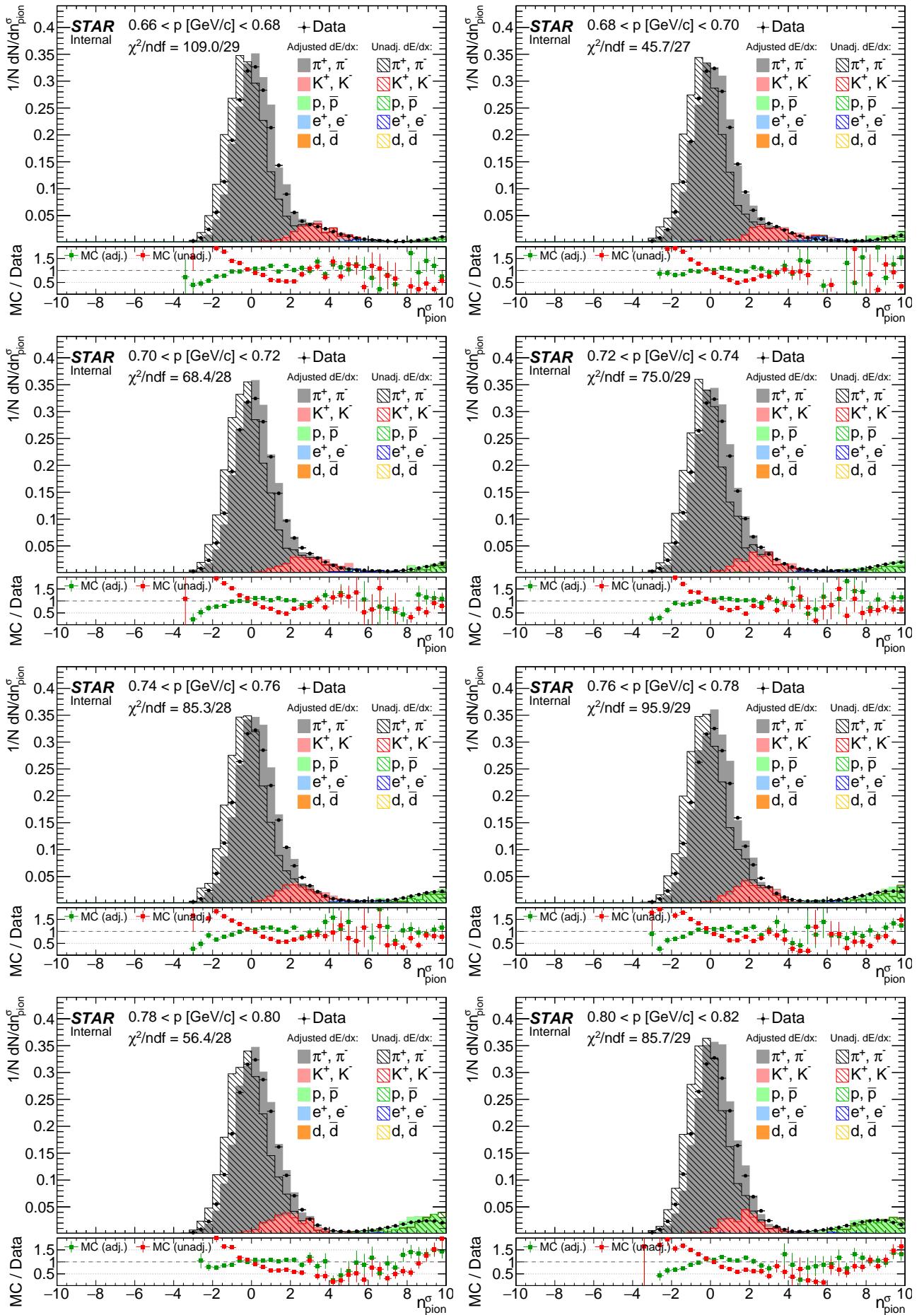


Figure E.3: Comparison of n_{pion}^{σ} distribution between data and embedded MC (before and after dE/dx adjustment) in momentum bins.









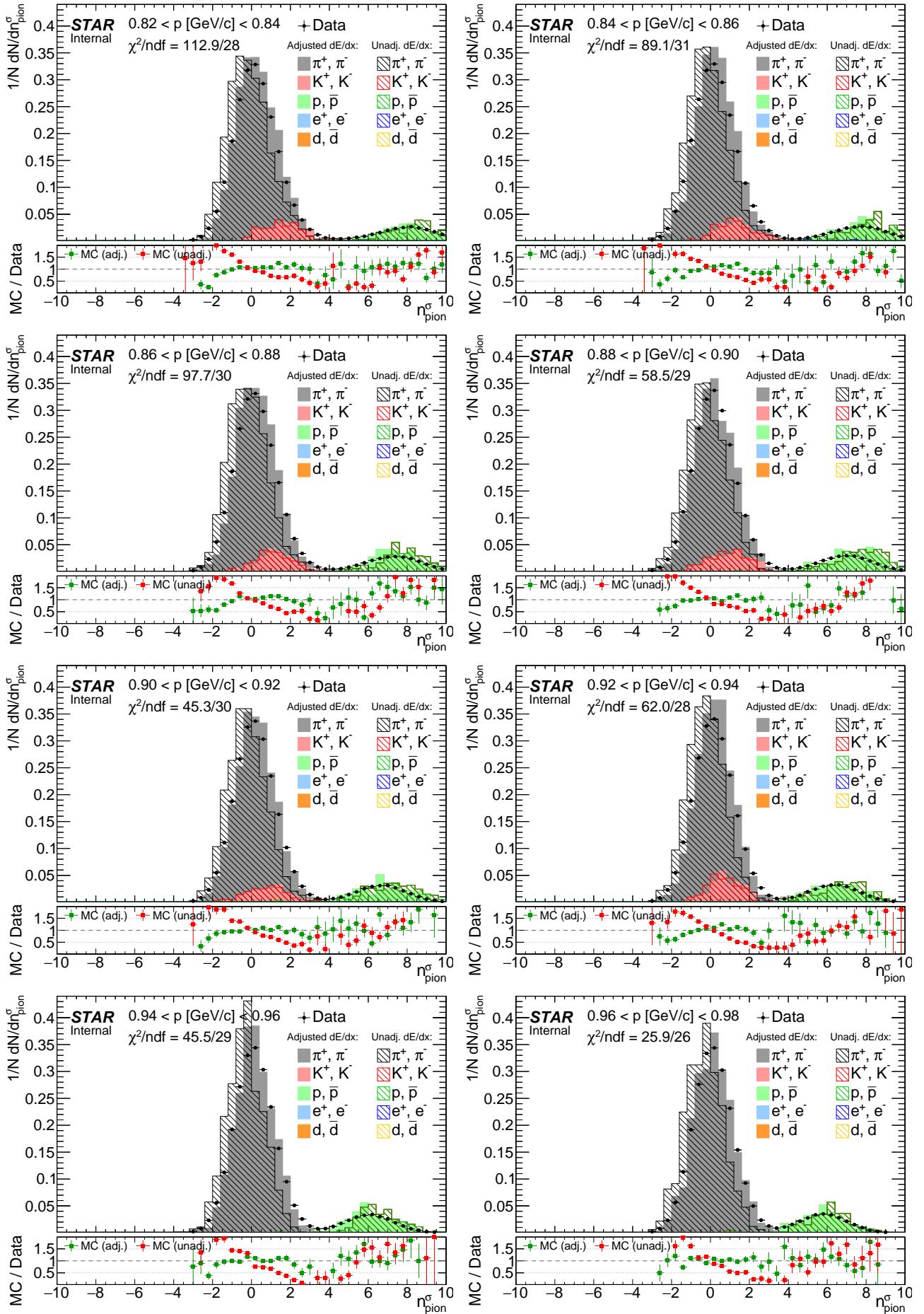
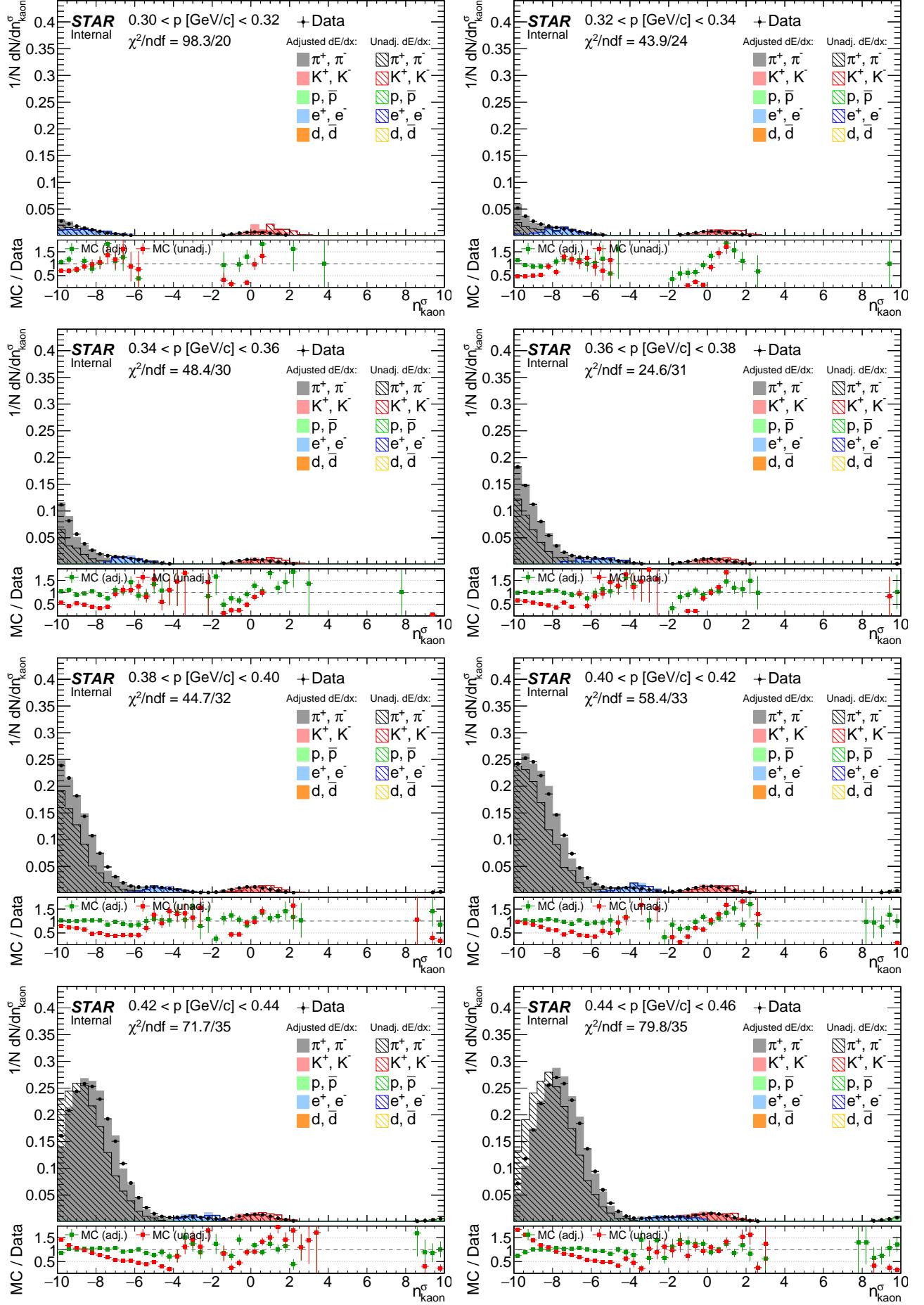
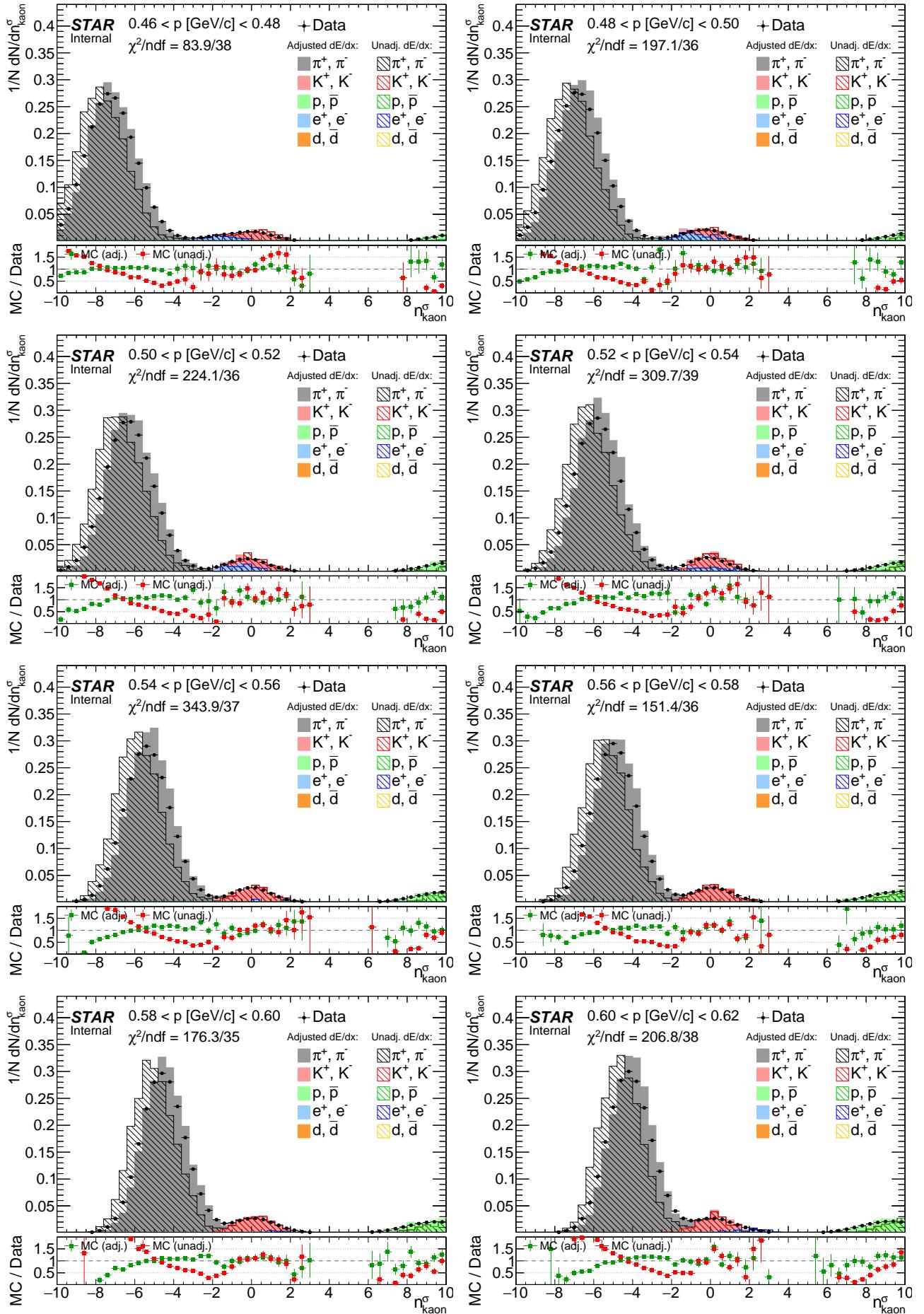
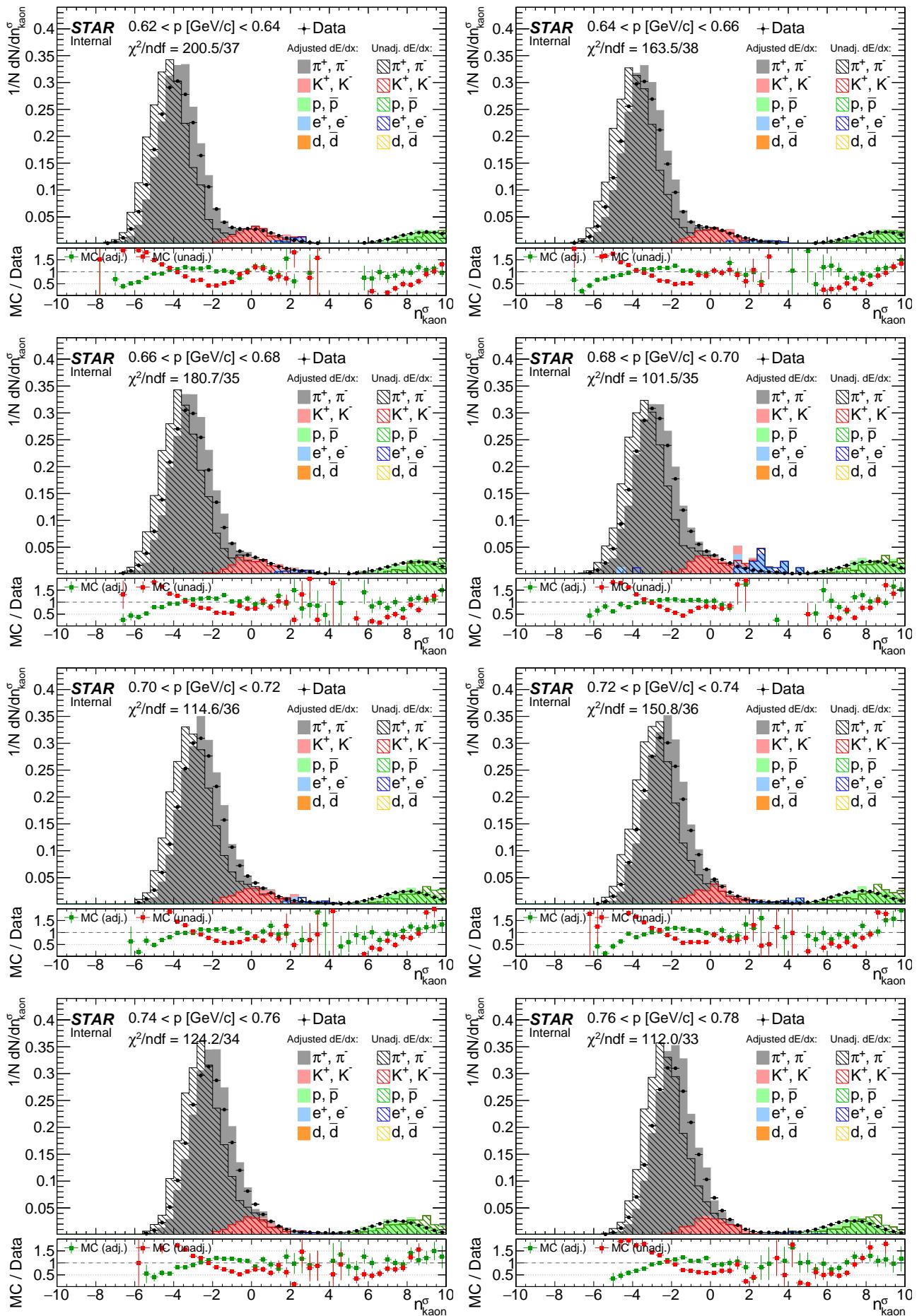
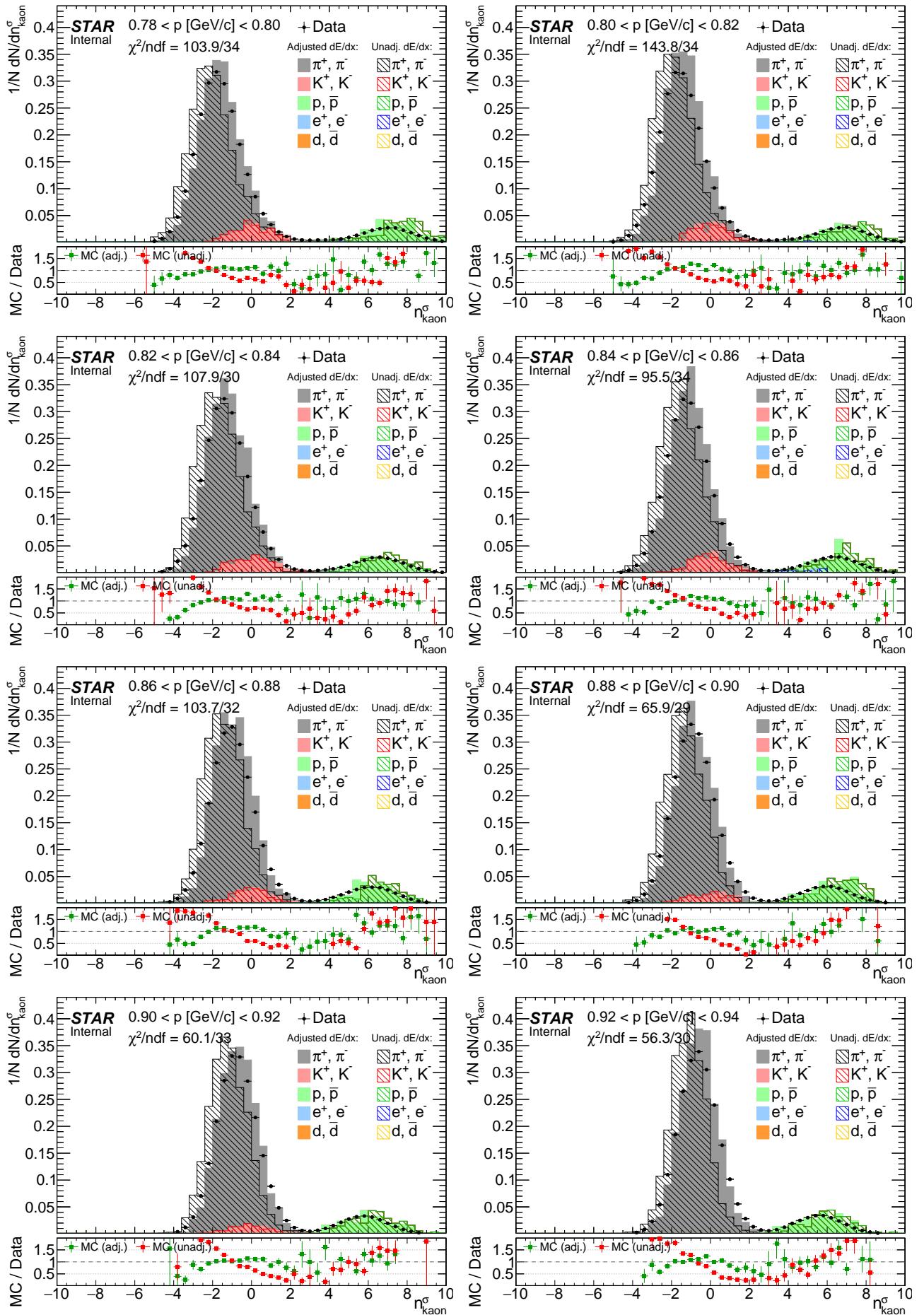


Figure E.4: Comparison of n_{kaon}^{σ} distribution between data and embedded MC (before and after dE/dx adjustment) in momentum bins.









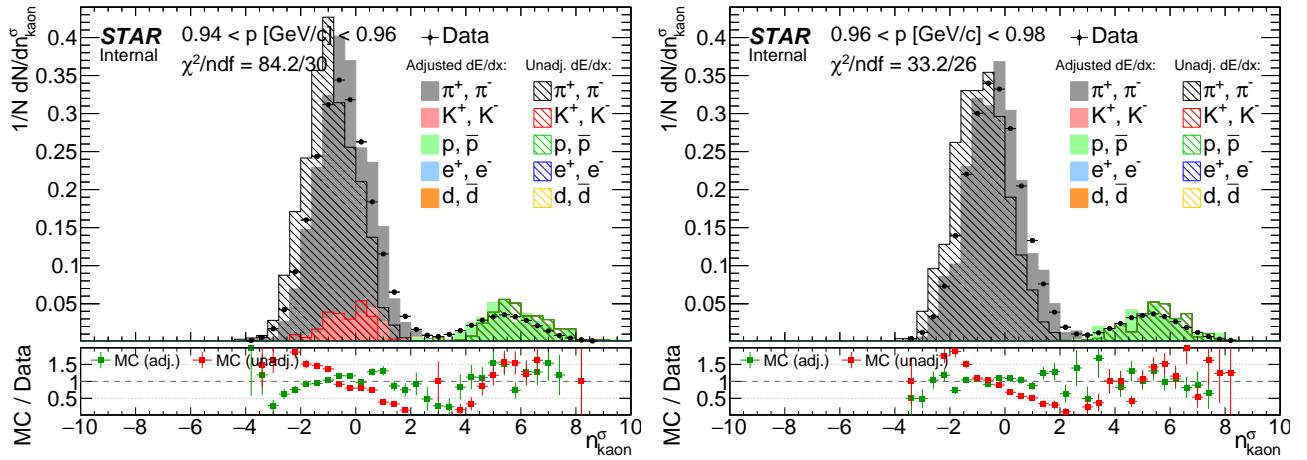
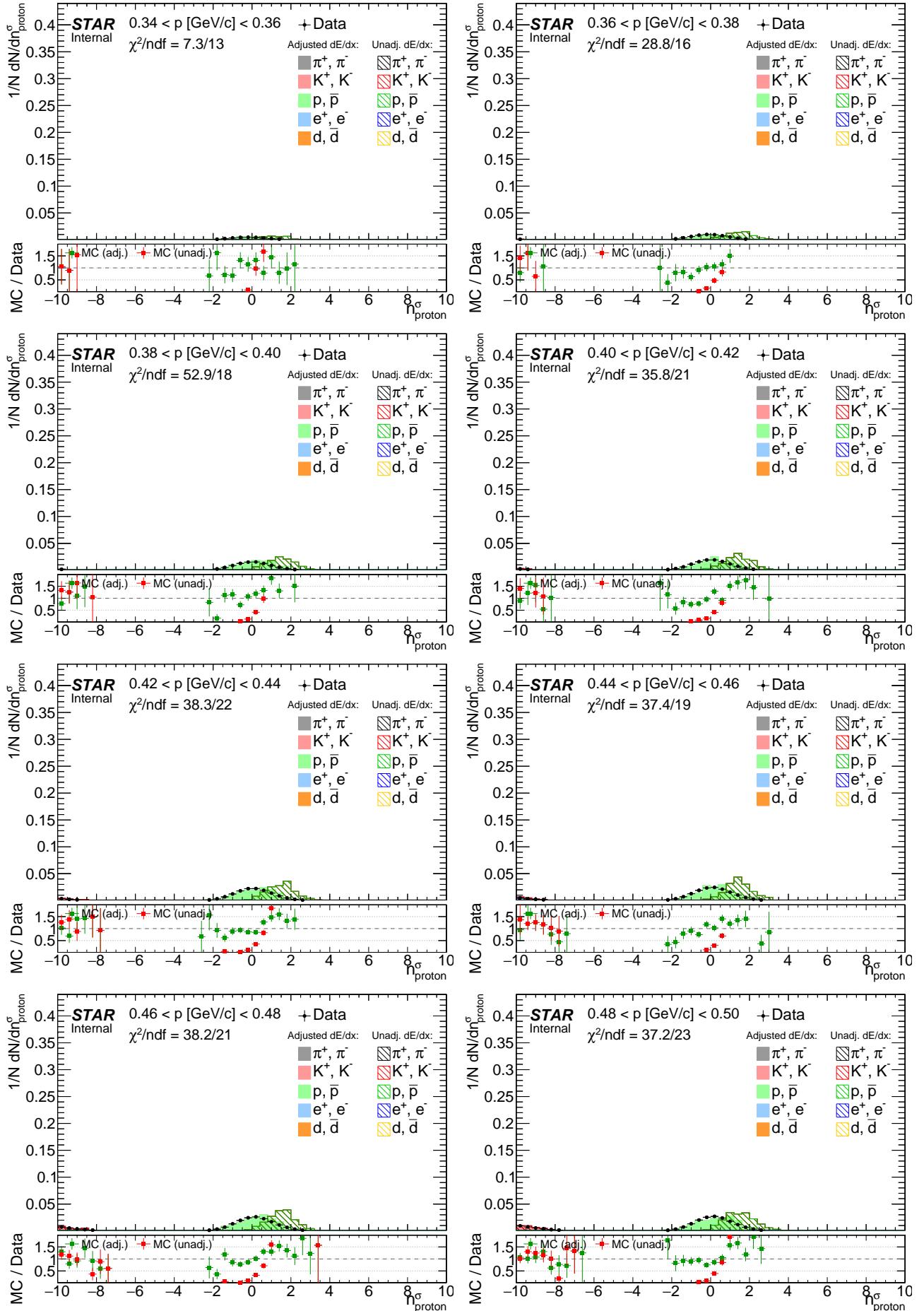
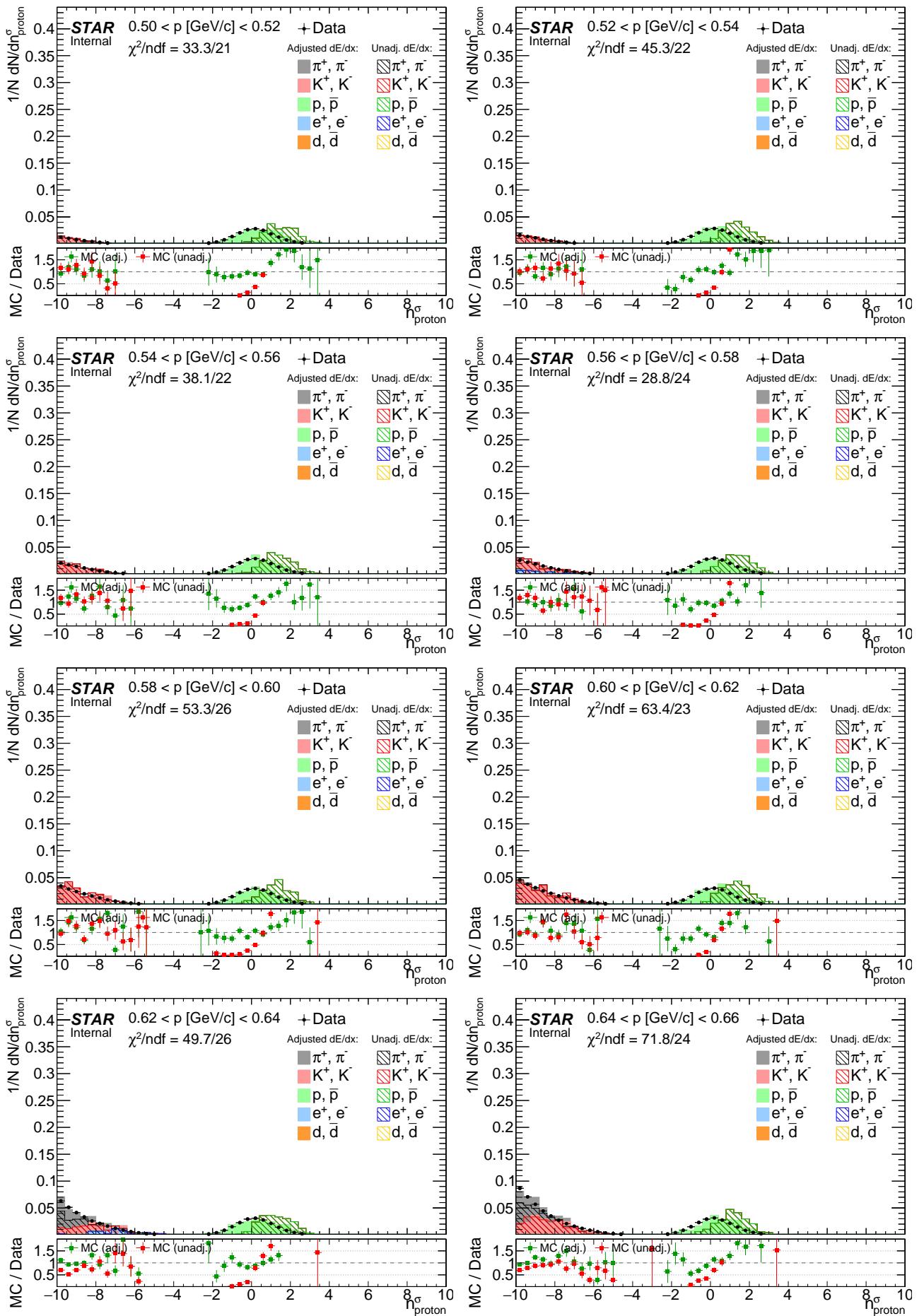
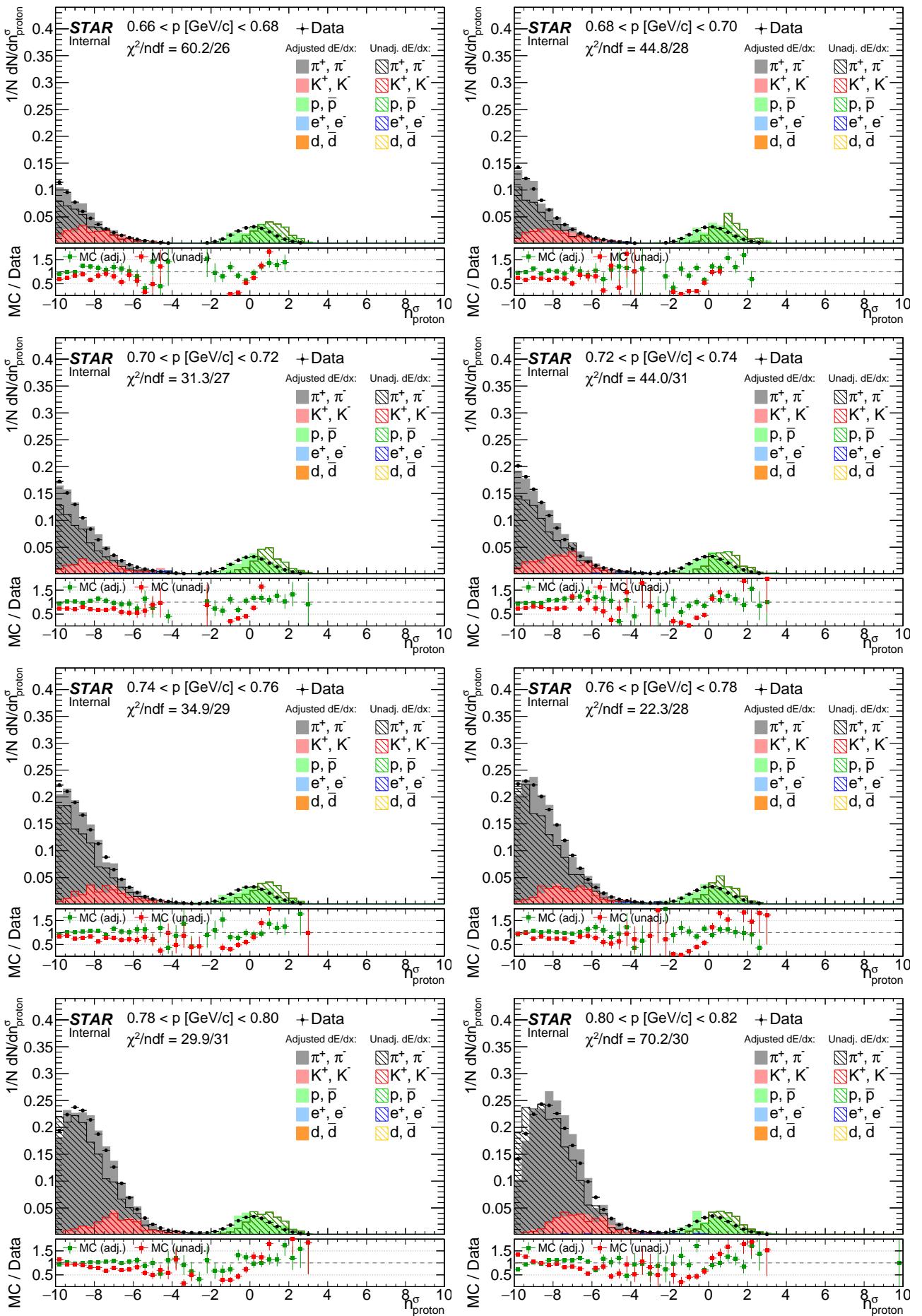
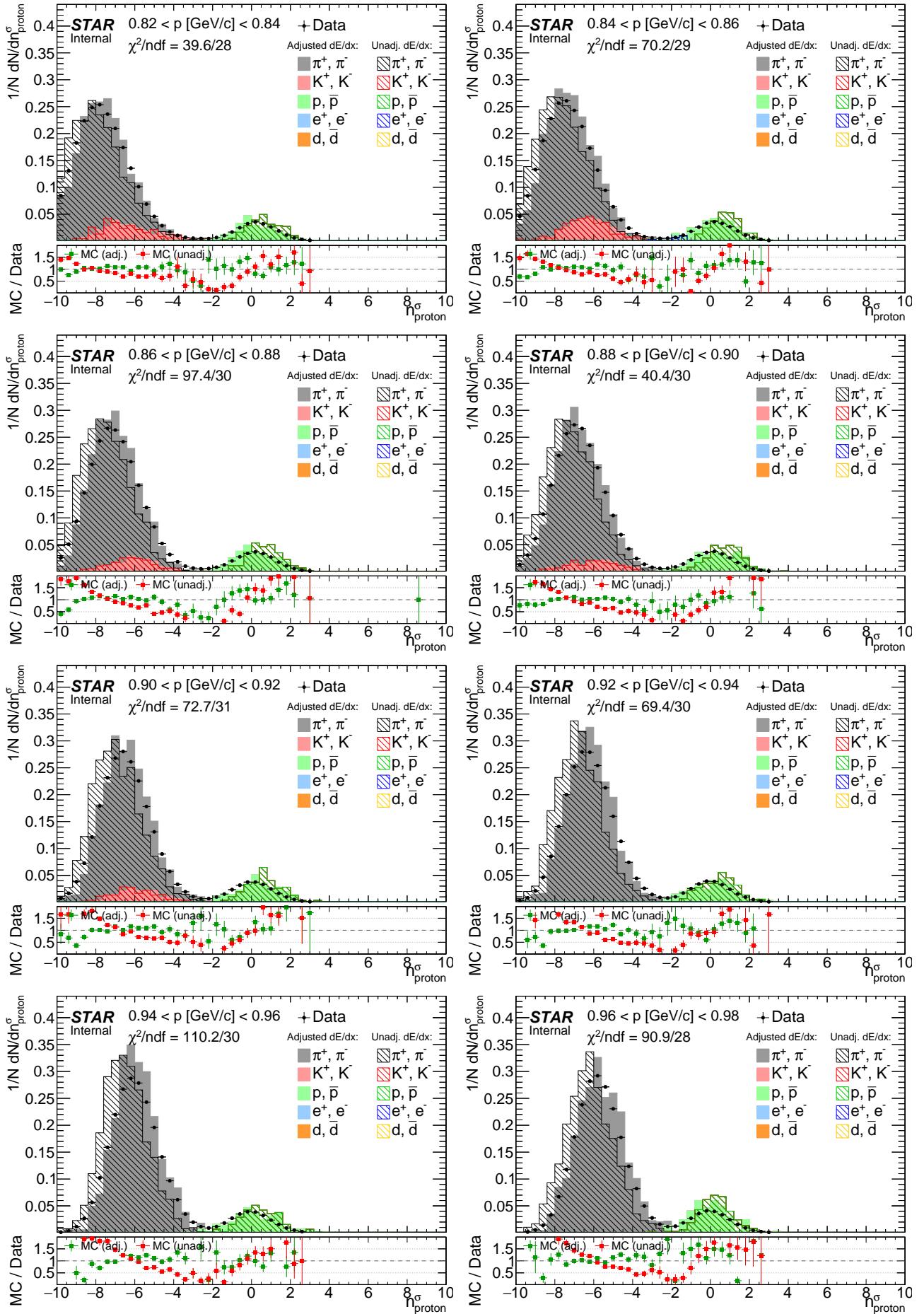


Figure E.5: Comparison of $n_{\text{proton}}^{\sigma}$ distribution between data and embedded MC (before and after dE/dx adjustment) in momentum bins.





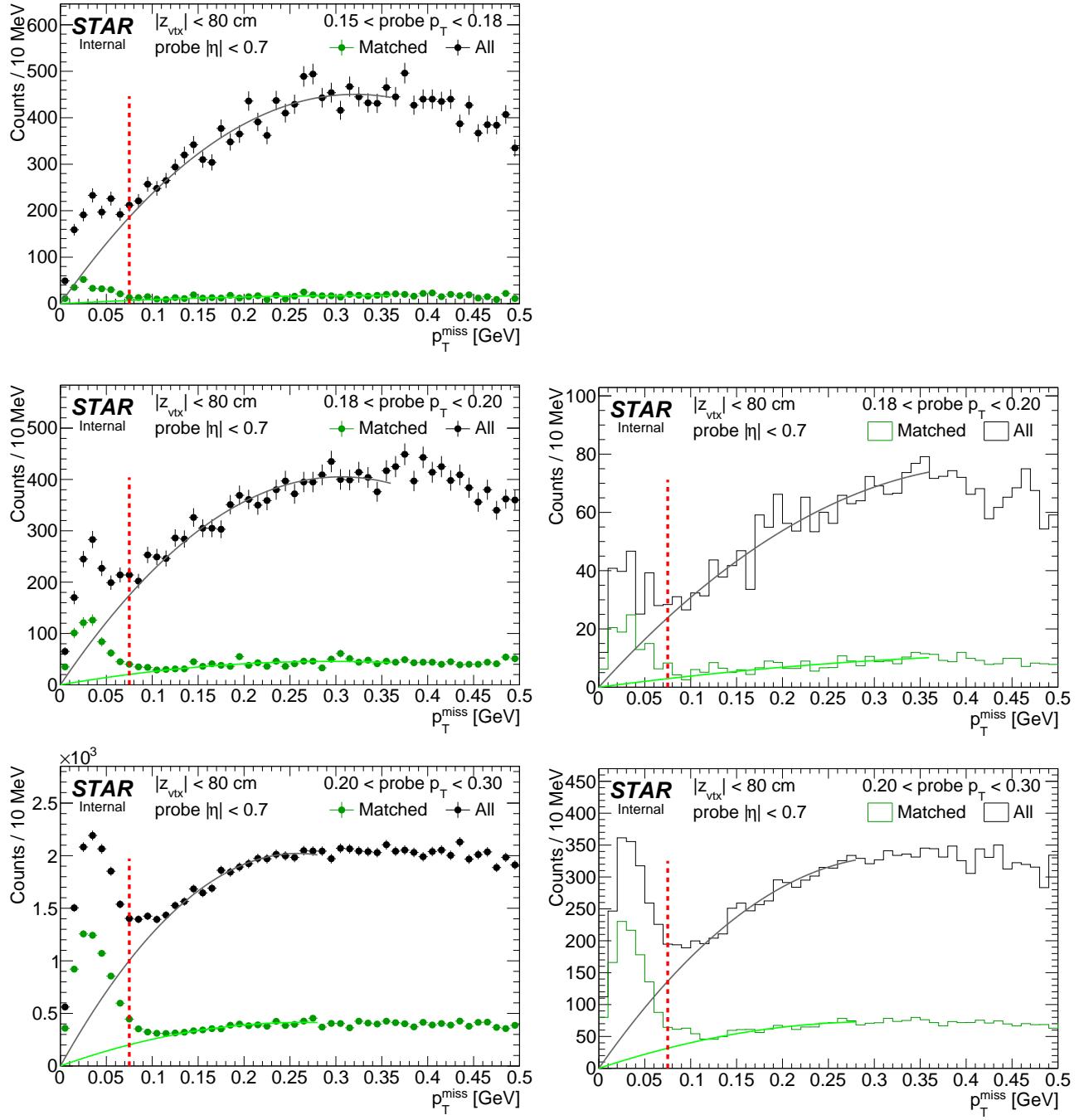




Appendix F

Tag&Probe fits for TOF hit reconstruction and matching efficiency

Figure F.1: lines and arrows indicate region accepted in analyses.



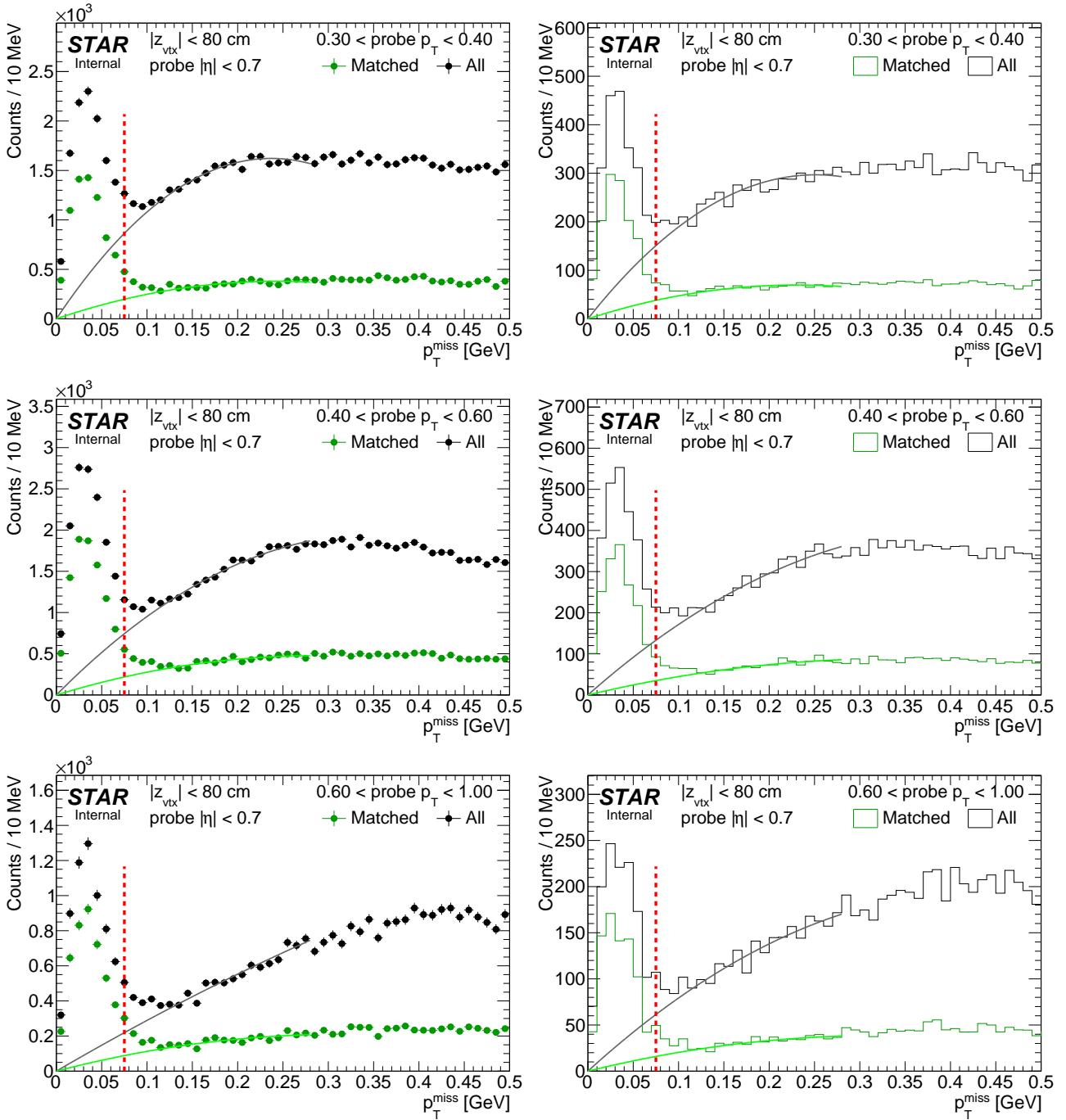
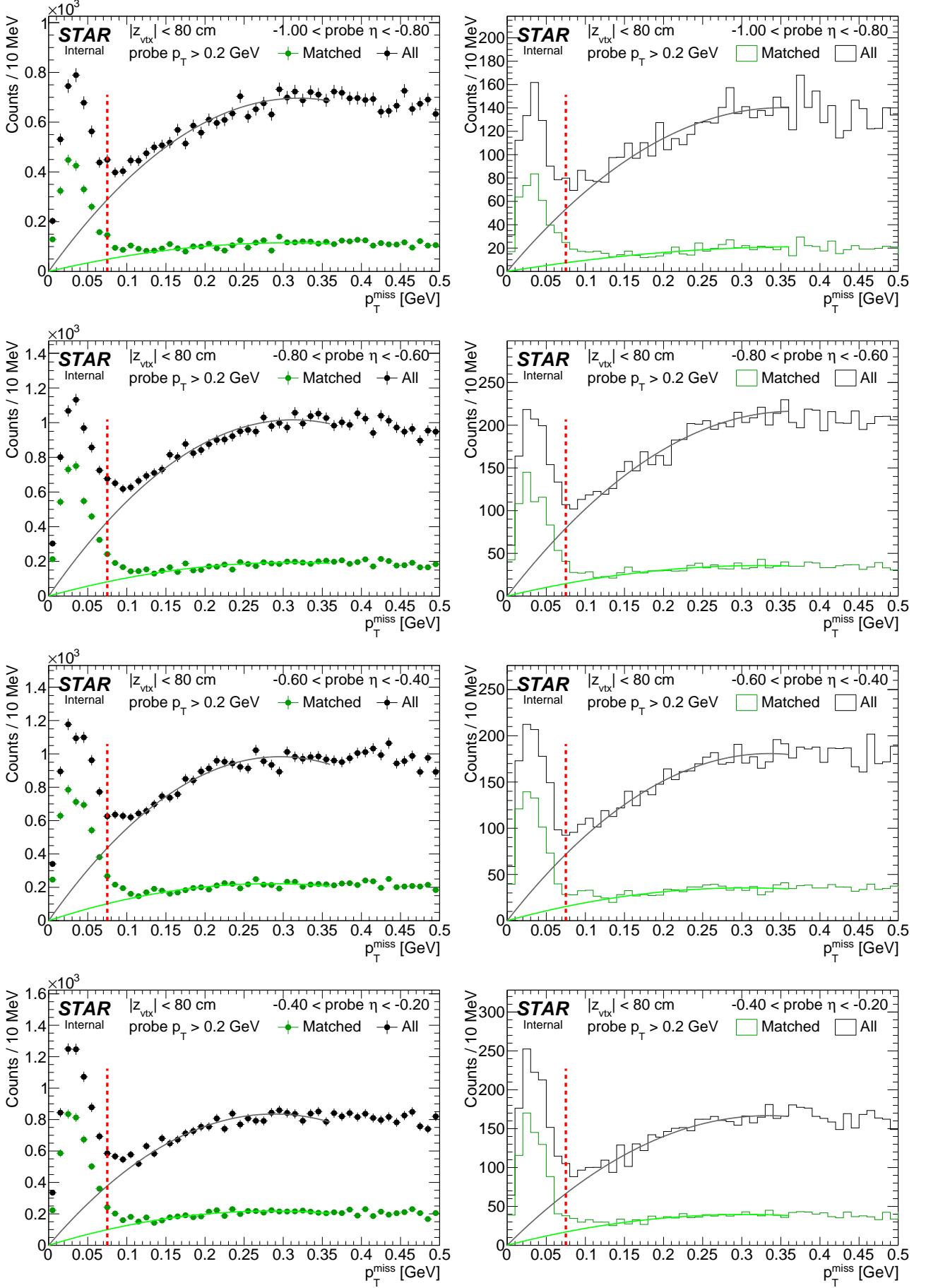
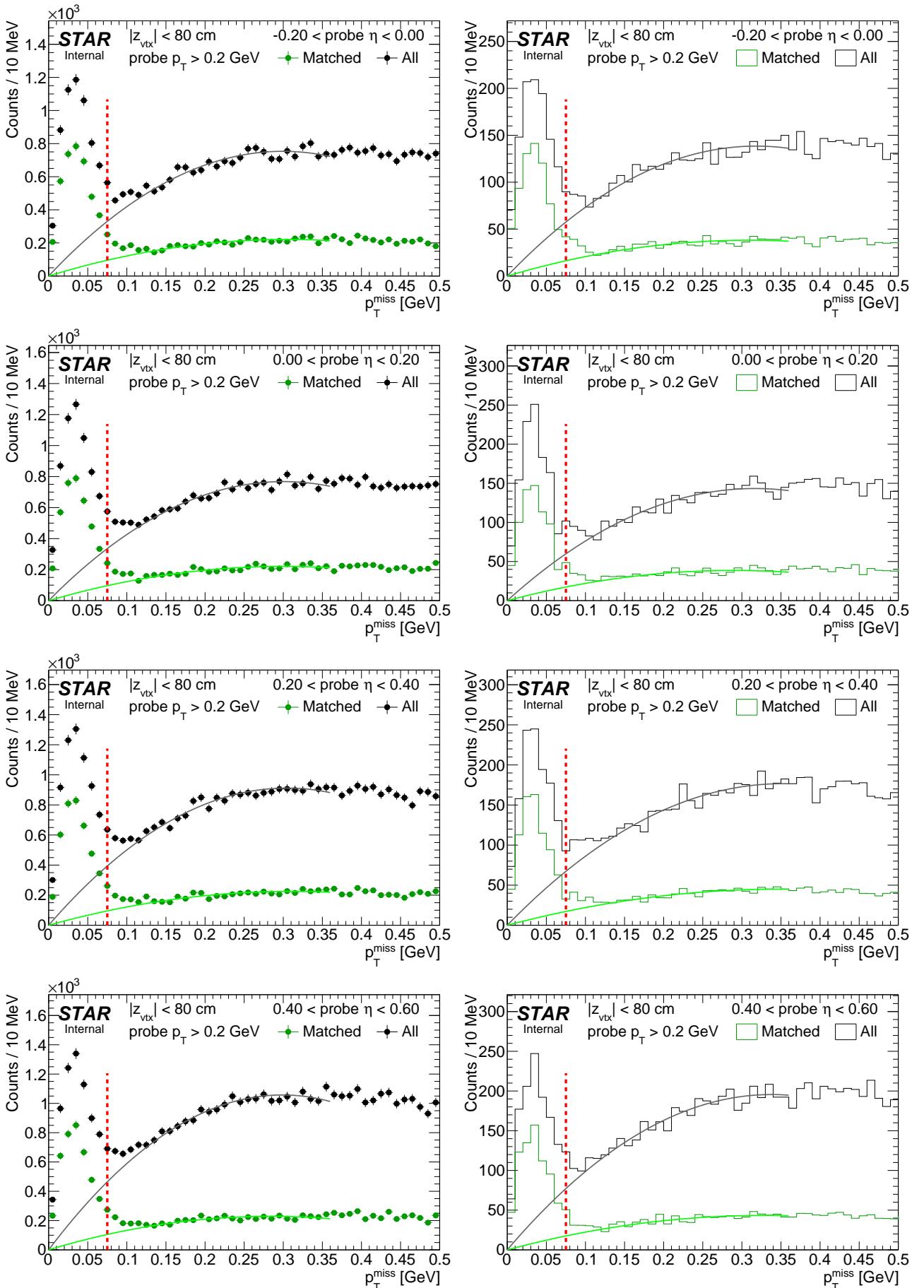
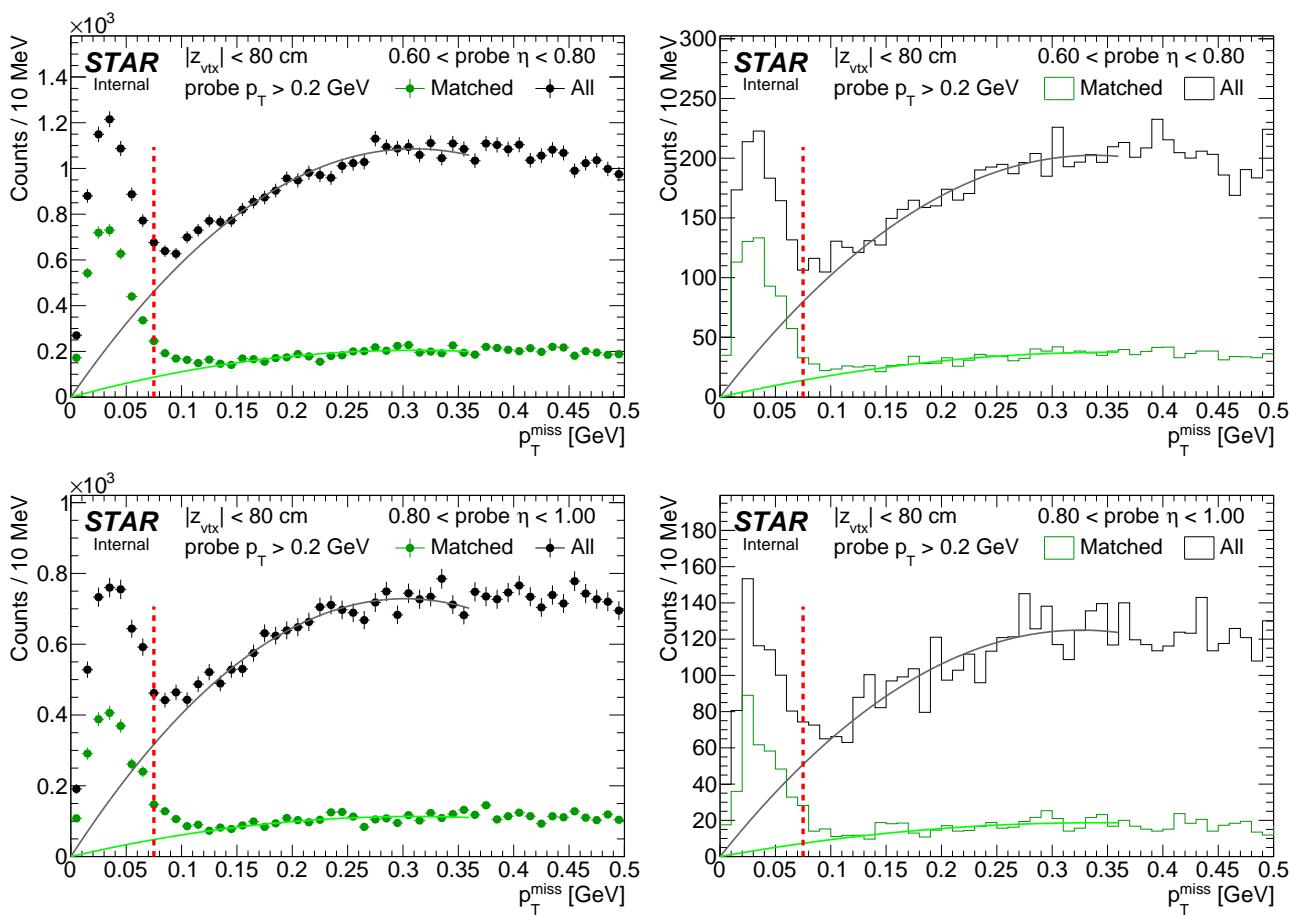


Figure F.2: TOF acceptance, reconstruction and matching efficiency of π^+ . Each plot represents the TOF efficiency ϵ_{TOF} (z -axis) as a func.







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