

Resonance search analysis of 2016 HPS spring run data.

Bump hunt folks

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Introduction

The Heavy Photon Search (HPS) experiment has capability to search for a so called heavy photon (A') with two complementary methods.

1 Data set

Describe the data, beam energy, beam current, target runs, etc.

2 Event Selections

This section describes all the cuts that are applied to get the final vertex candidate distribution. The main goal of event selection cuts is to maximize signal sensitivity.

In this analysis only events with “Pair1” trigger (see [1] for the description of HPS triggers) are used.

2.1 Cluster timing cuts

The readout window of ECal FADC data is 200 ns. Clusters coming from the physics events, that generated the trigger, are located in a narrow time range (few ns width because of the trigger jitter) in the readout window around $t = 56$ ns. In Fig.1 shown “time vs Energy” distributions of ECal clusters in the Top (Left)

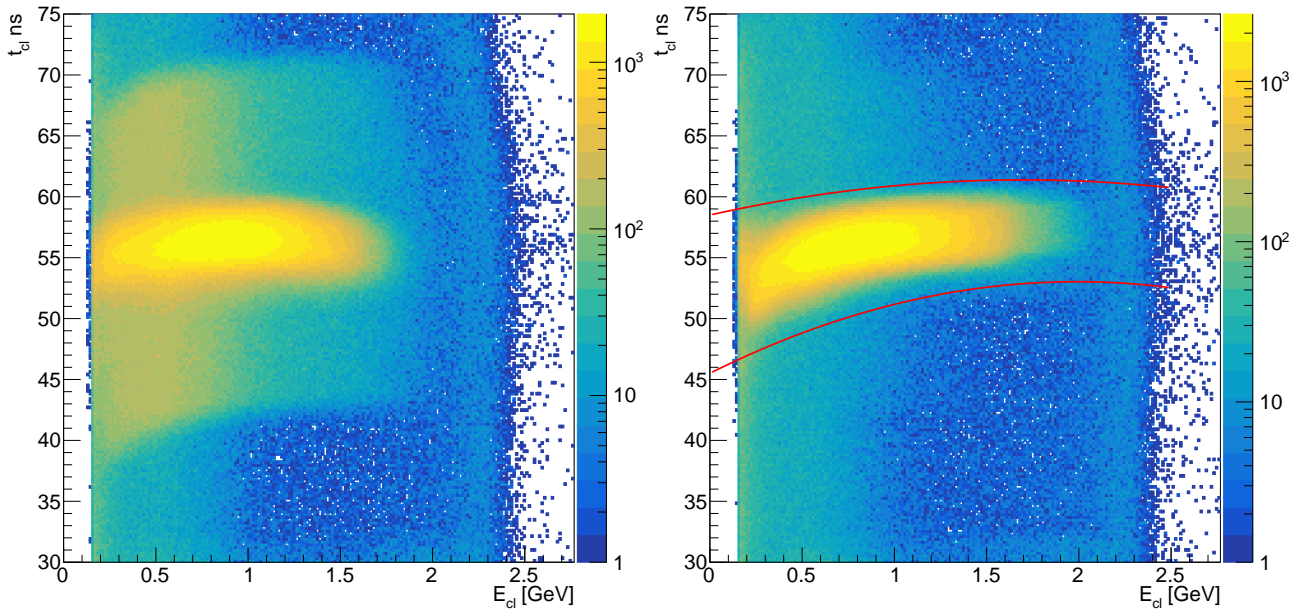


Figure 1: “time vs Energy” distributions of ECal clusters in the Top (Left) and Bottom (Right) half. Red curves in the in the right plot indicate cuts that are applied to clusters in the bottom half for the initial cluster selection. See the text for the description of the difference between left and right plots.

and Bottom (Right) halves. The bulge of events in the right plot are clusters that generated the trigger, and also trigger time is defined by these clusters. The noticeable energy dependence is due to the so called “time walk Corrections” [2]. During initial event selection only clusters that are inside the outlined red curves are used, since the rest are accidentals that didn’t come from the beam bunch generating the

trigger. One can notice that for the clusters in the top half, in addition to the central bulge, there is an extra occupancy of events in region ($40 \text{ ns} < t_{cl} < 70 \text{ ns}$). This is because the coincidence time between clusters in the “Pair” trigger was 12 ns [1], and the trigger time is determined by the bottom cluster. Unlike to clusters in the bottom half, in the initial event selection, we have not cut on time of the top cluster, but rather we have applied cut on the cluster time difference between top and bottom clusters.

2.1.1 Ad hoc ECal time corrections

The next step is to cut pairs of top-bottom clusters that are far from each other in terms of time. During the analysis it was found that ECal cluster times can be improved, in particular the dashed

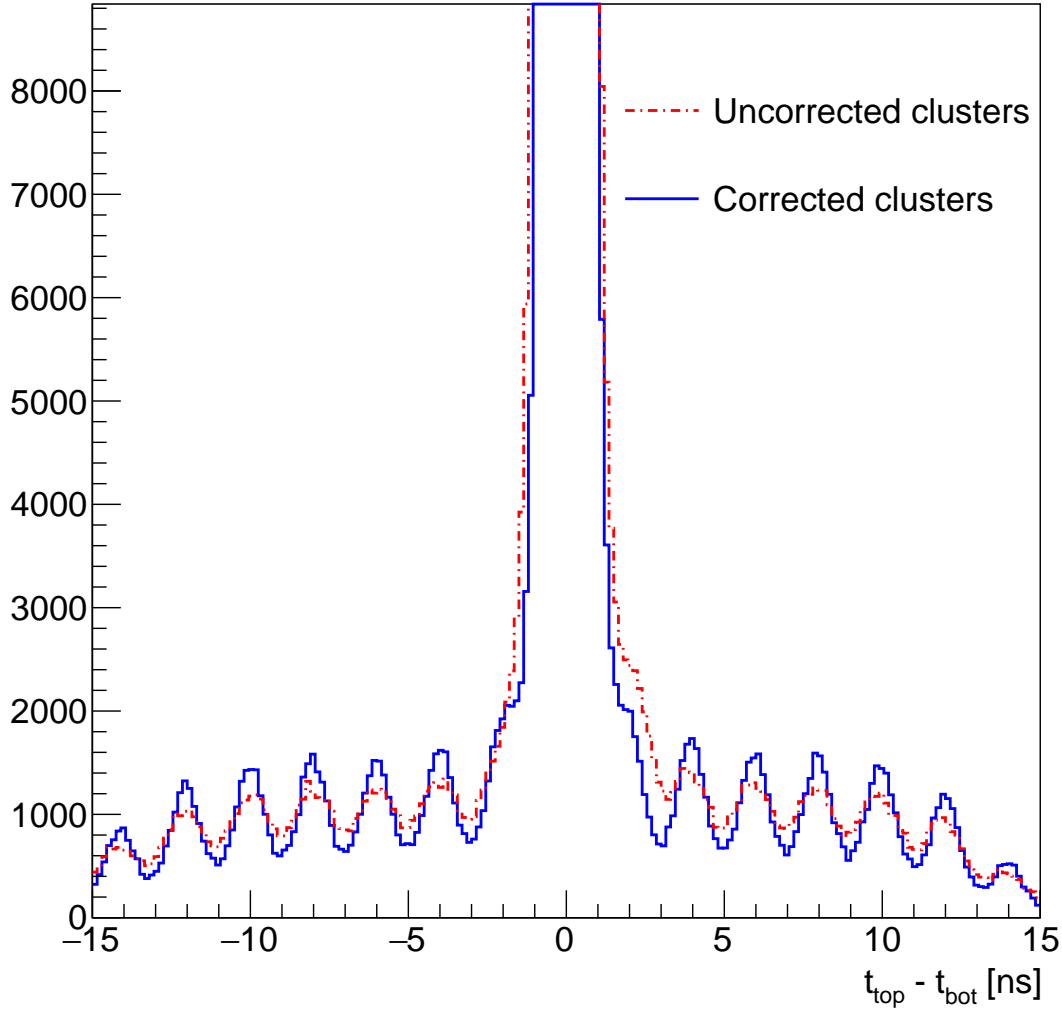


Figure 2: Time difference between top and bottom clusters at high E_{sum} region. Dashed red lines show uncorrected clusters, and the blue curve show corrected clusters.

red histogram in Fig.2 shows the time difference between top and bottom clusters¹ at high E_{sum} region ($1.9 \text{ GeV} < E_{\text{sum}} < 2.4 \text{ GeV}$). As one can see there is a bump at around 2 ns, while at -2 ns there is no clear bump. This suggests that time offsets of some crystals might be wrong. To check this, for each crystal the time difference between that crystal and its pair (in opposite half) crystal is constructed.

¹For the sake of better visualization, the plot shows doesn't fully show the entire central peak.

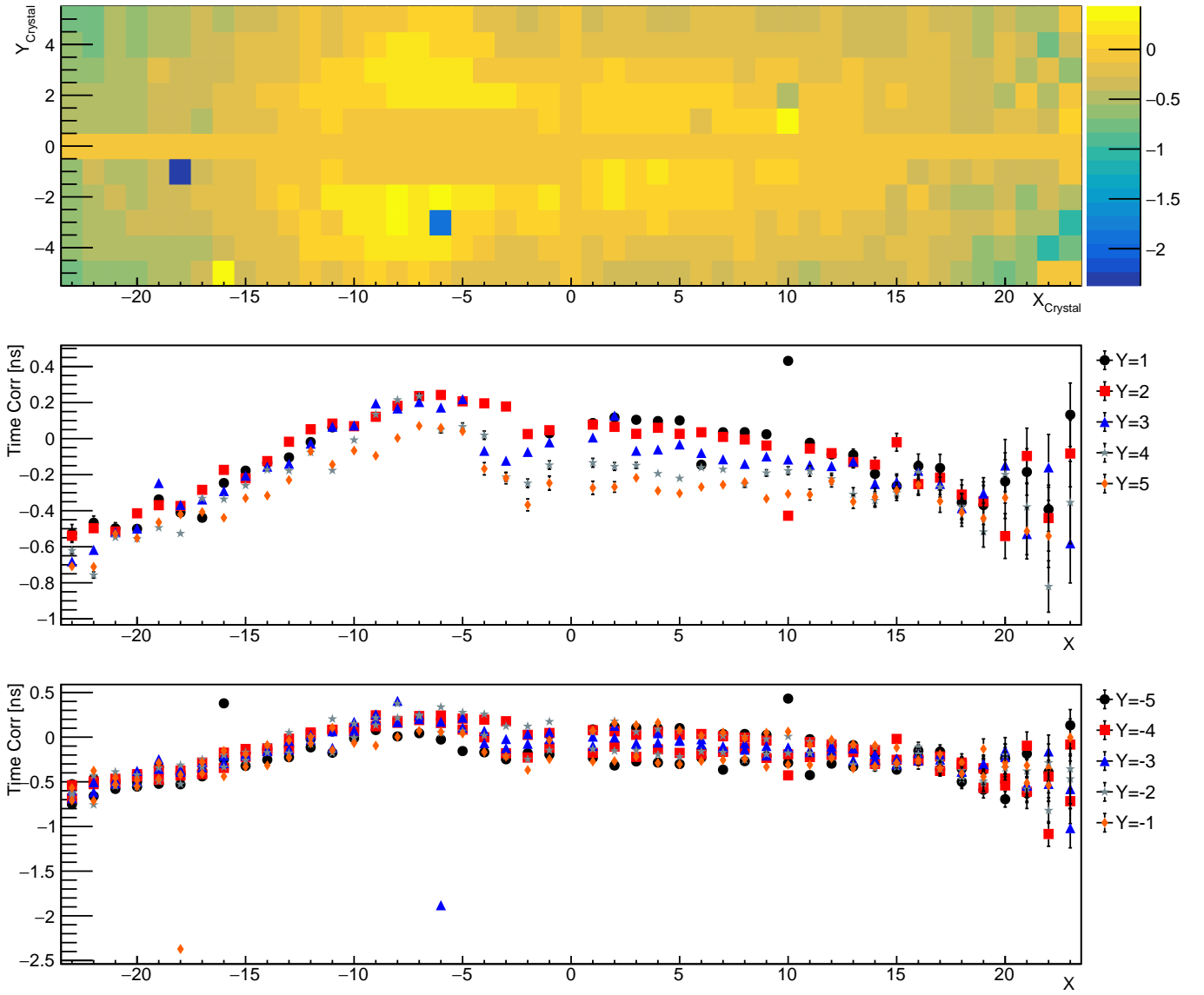


Figure 3: Time Corrections for each crystal.

The Top 2D plot of Fig. 3 shows mean values of each of crystals. The middle and the bottom plots show same mean values as a function of crystal X index for Top and Bottom crystal respectively. Different markers show different rows. As one can see crystals $(-18, -1)$ and $(-6, -3)$ are shifted from their immediate neighbors by about 2 ns. There are some other crystals which are shifted significantly too (but less than 2ns). These include for example crystals $(-16, -5)$, $(10, -5)$, $(10, 1)$ and $(10, 2)$. In addition to this, we see that there is a general crystal X index (and slightly Y index) dependence too. In reality crystal X index is correlated to the charged particle energy too, and the original dependence might be not on X but on energy. Studying it is out of the scope of this note, and here for each crystal we have corrected the time, by subtracting these calculated mean values from the reconstructed cluster time. After the correction the cluster time difference is depicted by blue solid line in Fig.2. One can see that the excess of events at 2 ns disappeared. Dips and peaks between bumps indicating difference beam bunches also got sharper, which is an indication of an improvement of the cluster time resolution.

2.1.2 Fitting Cluster time difference

After correction of individual cluster times, the Top-Bottom cluster time difference was fitted with a following function:

$$F = \sum_{i=0}^{N_{\text{peak}}} a_i \cdot \left(\text{Gaus}(x - \mu_i^1, \sigma_i^1) + b \cdot \text{Gaus}(x - \mu_i^2, \sigma_i^2) \right) \quad (1)$$

where N_{peak} is the number of peaks. Each peak is described by the sum of two Gaussian functions $\text{Gaus}(x - \mu_i^1, \sigma_i^1)$ and $\text{Gaus}(x - \mu_i^2, \sigma_i^2)$ with their amplitude ratio "b". The parameter "b" is the same for all peaks. In the fit, free parameters are a_i , μ_i^1 , σ_i^1 , μ_i^2 , σ_i^2 , b.

The fit result is shown in Fig.4. Different peak components of the function are depicted by different

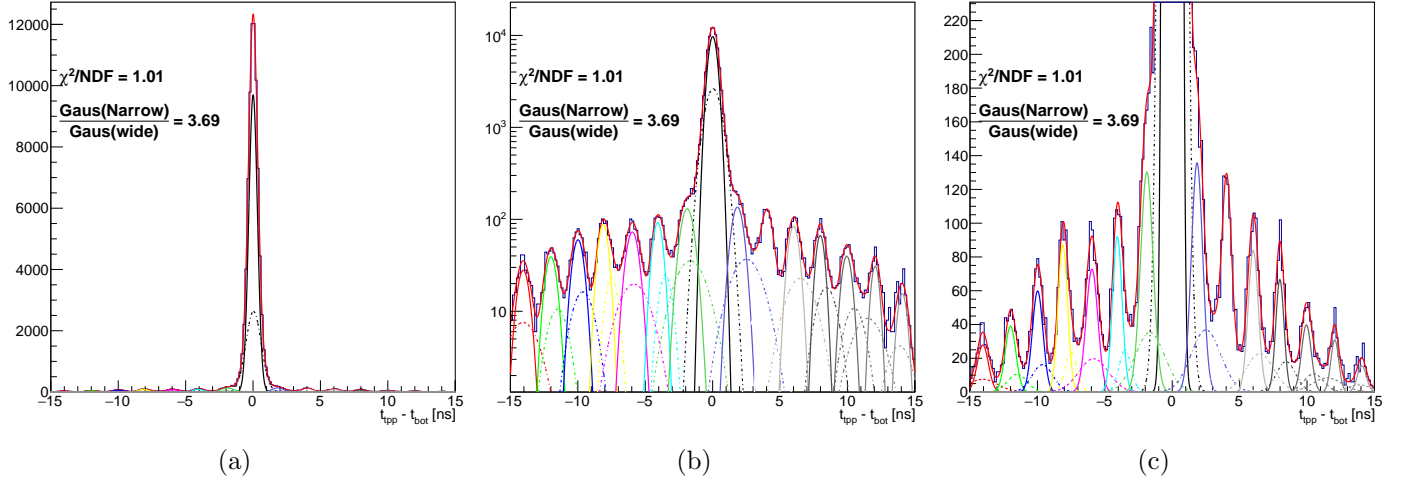


Figure 4: Fit of the Top and Bottom cluster time difference. Left: linear scale, Middle: Log scale and right: Linear scale but includes only low magnitude peaks. Main Gaussian functions are represented by solid lines, and the secondary Gaussian (with wider width and lower magnitude) are represented by dashed lines.

color. The main Gaussian function of each peak is represented by a solid line, while the secondary Gaussian is represented with a dashed line. One can see that this function fits the distribution reasonably well.

Then in order to determine the optimal cut on the cluster time difference, we will use the value, which maximizes the $\frac{S}{\sqrt{S + \text{Bgr}}}$ ratio, where "S" is the signal (in our case the central peak), and "S + Bgr" is the signal plus Background (the total fit function). The $\frac{S}{\sqrt{S + \text{Bgr}}}$ ratio as a function of cluster time difference cut is shown in Fig.5, where the maximum value at $\Delta t < 1.43$ ns is indicated by a vertical dashed line.

2.2 Two dimensional cuts

Some of event selections cuts described below are two dimensional cuts, i.e. the cut value depends on the value of another variable. In most of cases two dimensional cuts are implemented as a function of particle's momentum.

In general, to study the distribution of a given variable for a "signal like" particle, the rest of event selection cuts are applied, to make as clean as possible signal. The only exception is the two cluster time difference cut, which is described in section 2.1). Applying the rest of cuts except the one under the

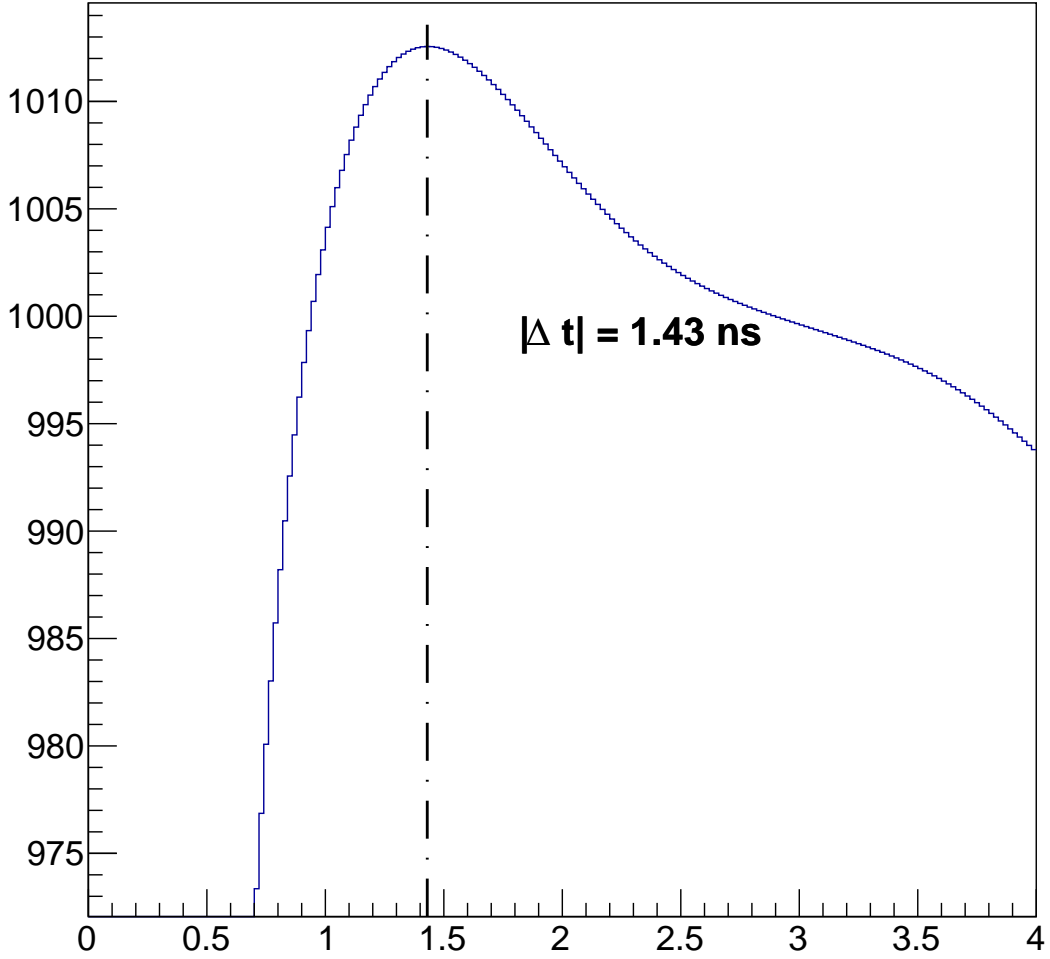


Figure 5: The $\frac{S}{\sqrt{S + Bgr}}$ ratio as a function of cluster time difference cut. The Dashed line indicates the maximum of the function.

investigation, will ensure the accidental background is minimal (negligible), and the resulting distribution will represent actual signal (the e^- , e^+ , (X) final state). In most of cases the distribution is not Gaussian, even when it represent a small momentum bin. In such cases, the conventional $\pm 3\sigma$ cuts will not be will not keep 99.7% but rather might cut more events. As an example in Fig.6 shown a toy distribution which is not a pure Gaussian, but rather has a tail on the left side. The Gaussian fit is shown on top if the histogram and $\pm 3\sigma$ limits are shown by vertical red lines. One can see that -3σ limit will cut several % of events rather than 0.3%. Instead it was decided to choose left and right cuts limits such that will keep 99% of the signal and will throw 0.5% of signal events from each side. In this particular case 99% cut limits are shown by blue vertical lines.

2.3 Track-Cluster Matching

The offline reconstruction code forms particles by matching tracks and clusters to each other, by utilizing spatial coordinate and time differences between tracks and clusters. In the offline reconstruction the matching is quite loose ("Better to keep junk, rather than throwing a good particle"). In this section

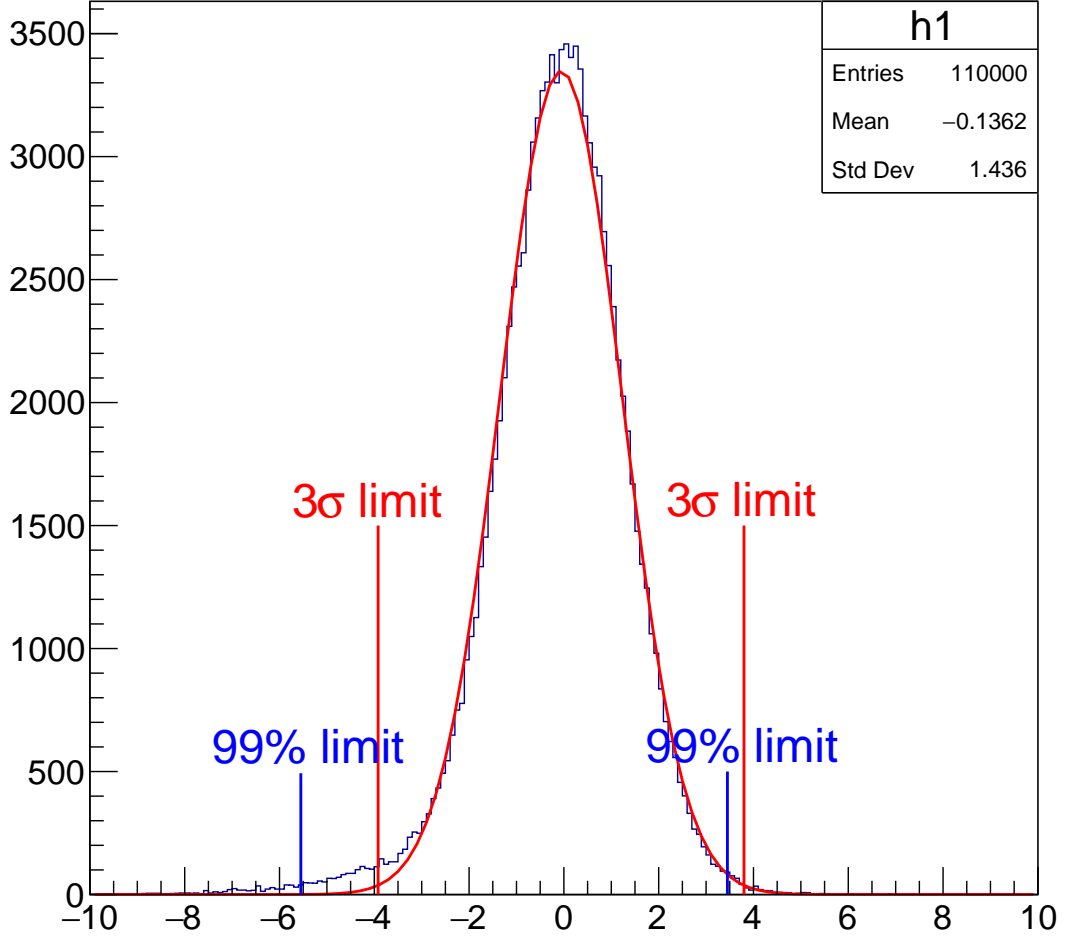


Figure 6: Illustration of 3σ cut limits vs 99% cut limits on a toy distribution.

spatial and time matching cuts are described. Both, time and position resolution of e^- and e^+ clusters depend on particle momentum. The precision of the track projected coordinate at the ECal face does depend on the track momentum too. Because of these reasons we studied track-cluster matching as a function of momentum.

2.3.1 time matching

In addition to the momentum dependence we noticed also slight difference between top and bottom sectors, therefore two separate cuts are developed for each detector half. In Fig.7 shown Cluster-Track time difference as a function of particle momentum. Left plot represents particles in the bottom half, and the right plot represents particles in the top half of the detector. The area marked by Red dots represents the acceptance region.

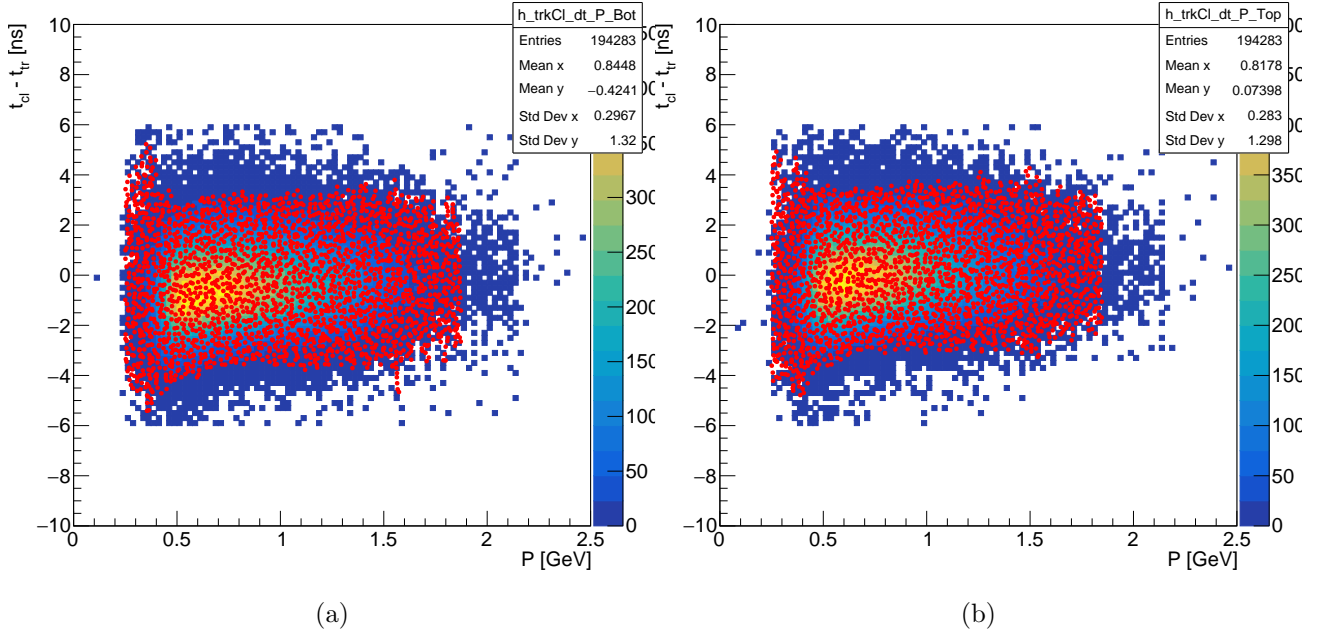


Figure 7: Cluster-Track time difference as a function of particle momentum. Left plot represents particles in the bottom half, and the right plot represents particles in the top half of the detector. The area marked by Red dots represents the acceptance region.

2.4 Track χ^2 cuts

2.5 WAB Suppression cuts

Describe L1 and d_0 cuts.

Mention why L1 is important, addition of suppressing WABs, it also significantly improves the mass resolution.

3 Parametrization of Mass resolution.

4 Bump hunt analysis

5 Study of systematics

Here goes studies on systematics

References

- [1] Kyle McCarty, Valery Kubarovsky and Benjamin Raydo, “Description and Tuning of the HPS Trigger”, HPS Note 2018-002
- [2] Holly Szumila-Vance, “HPS Ecal Timing Calibration for the Spring 2015 Engineering Run”, HPS Note 2015-011.