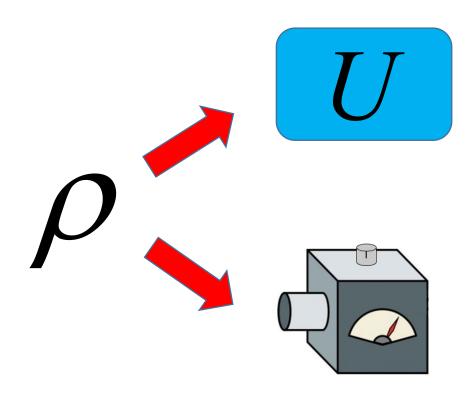
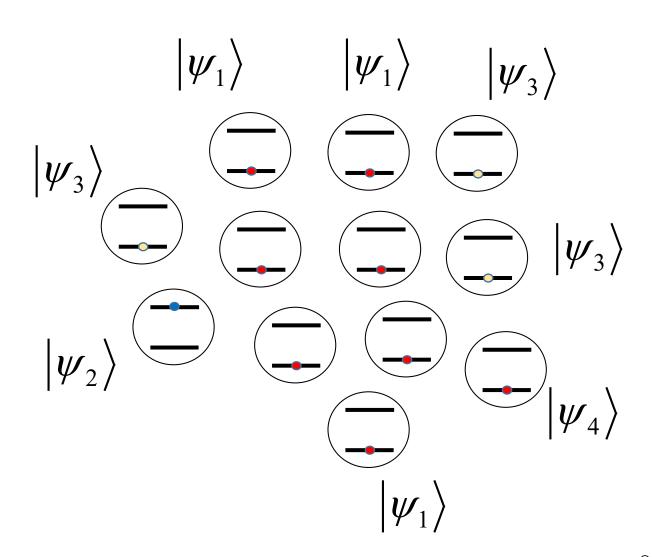
22. Applications of the Density Matrix



Recap: Density matrix definition

$$\rho = \sum_{n} p_{n} |\psi_{n}\rangle \langle \psi_{n}|$$

$$p_n=egin{array}{ll} ext{Probability of occurrence of} \ ext{state} \ \ket{\psi_n} \end{array}$$



What to do with the density matrix?

The density matrix ρ describes the quantum state of a noisy system.

Everything we did with the wavefunction $|\psi
angle$, we can also do with the density matrix!

- Measurements
- Observables
- Unitary operations (quantum gates)



We will show how these work with the density matrix.

Unitary operations

Recall that for a pure state, under a Hamiltonian H the state evolves like

$$|\psi\rangle \to U|\psi\rangle$$

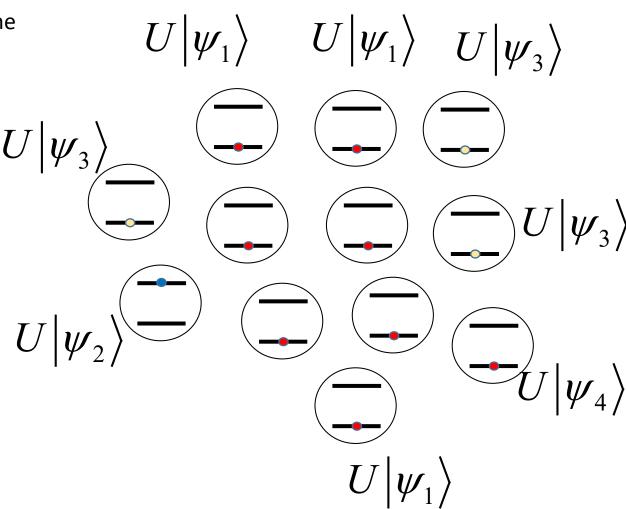
$$U = e^{-iHt/\hbar}$$

So the density matrix evolves as

$$\rho \to \rho' = \sum_{n} p_{n} U |\psi_{n}\rangle \langle \psi_{n}|U^{\dagger}$$

Or simply

$$\rho \rightarrow \rho' = U \rho U^{\dagger}$$



The state of a system is such that it is 50% in the state $|+\rangle$ and 50% in the state $|0\rangle$. An X gate is applied to the system. What is the new density matrix of the system?

The state of a system is such that it is 50% in the state $|+\rangle$ and 50% in the state $|0\rangle$. An X gate is applied to the system. What is the new density matrix of the system?

Initial state:

$$\rho = \frac{1}{2} |+\rangle \langle +|+\frac{1}{2}|0\rangle \langle 0| = \frac{1}{4} \begin{pmatrix} 3 & 1 \\ 1 & 1 \end{pmatrix}$$

Final state:

$$\rho' = X \rho X^{\dagger} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \frac{1}{4} \begin{pmatrix} 3 & 1 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \frac{1}{4} \begin{pmatrix} 1 & 1 \\ 1 & 3 \end{pmatrix}$$

Measurements

Recall that under a measurement the state evolves like

$$\left|\psi\right\rangle \to \frac{M_{j}\left|\psi\right\rangle}{\sqrt{\left\langle\psi\left|M_{j}^{\dagger}M_{j}\right|\psi\right\rangle}}$$

So the density matrix evolves as

$$\rho \to \rho' = \frac{\sum_{n} p_{n} M_{j} |\psi_{n}\rangle \langle \psi_{n}| M_{j}^{\dagger}}{N}$$

Where N is the normalization factor

$$N = \operatorname{Tr}\left[\sum_{n} p_{n} M_{j} | \psi_{n} \rangle \langle \psi_{n} | M_{j}^{\dagger}\right]$$
$$= \sum_{n} p_{n} \langle \psi_{n} | M_{j}^{\dagger} M_{j} | \psi_{n} \rangle$$

The state of a system is such that it is 50% in the state $|+\rangle$ and 50% in the state $|0\rangle$, is now measured in the $\{|0\rangle,|1\rangle\}$ basis, and the $|0\rangle$ outcome is obtained. What is the new density matrix? $\rho \to \rho' = \frac{M_j \rho M_j^\dagger}{\mathrm{Tr}(M_j \rho M_j^\dagger)}$

The state of a system is such that it is 50% in the state $\ket{+}$ and 50% in the state $\ket{0}$, is now measured in the $\ket{\ket{0},\ket{1}}$

basis, and the $|0\rangle$ outcome is obtained. What is the new density matrix?

Initial state:

$$\rho = \frac{1}{2} |+\rangle \langle +| + \frac{1}{2} |0\rangle \langle 0| = \frac{1}{4} \begin{pmatrix} 3 & 1 \\ 1 & 1 \end{pmatrix}$$

$$M_0 = |0\rangle\langle 0|$$

Applying the measurement operators

$$M_0 \rho M_0^{\dagger} = \frac{3}{4} |0\rangle \langle 0|$$

$$Tr(M_0 \rho M_0^{\dagger}) = \frac{3}{4}$$

$$\rho' = |0\rangle\langle 0|$$

$$\rho \rightarrow \rho' = \frac{M_j \rho M_j^{\dagger}}{\text{Tr}(M_j \rho M_j^{\dagger})}$$

Observables

Recall that observables are calculated for a pure state by

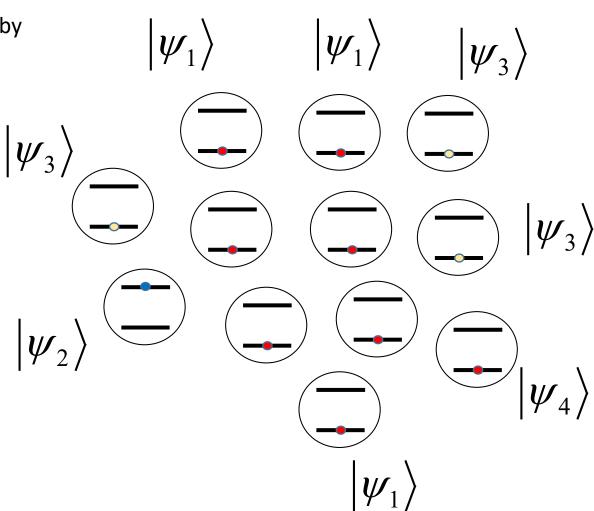
$$\langle A \rangle = \langle \psi \, | \, A \, | \, \psi \rangle$$

For a mixed state, we simply average over all the expectation values

$$\langle A \rangle = \sum_{n} p_{n} \langle \psi_{n} | A | \psi_{n} \rangle$$

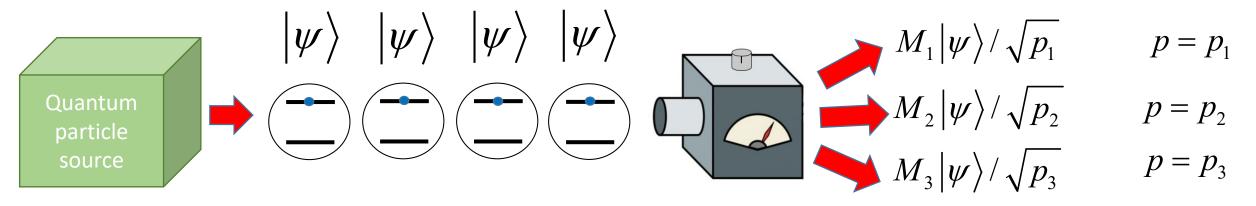
Or another way to write this is

$$\langle A \rangle = \operatorname{Tr}(A\rho)$$



Mixed state after measurement

How do we get a mixed state in the first place? One way is through measurement:



The density matrix of the ensemble after the measurement is

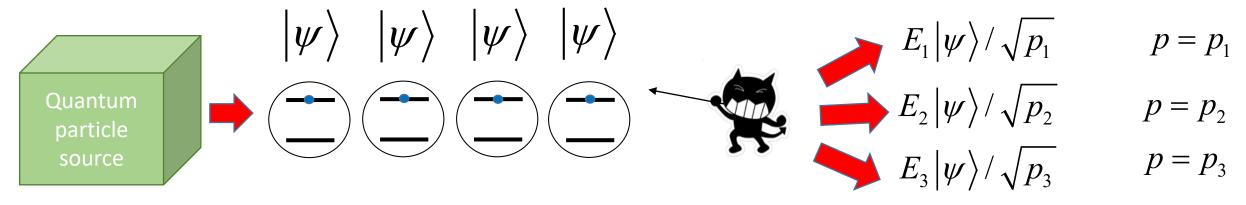
$$\rho' = \sum_{n} p_{n} |\psi_{n}\rangle \langle \psi_{n}| \qquad \text{In this case} \qquad |\psi_{n}\rangle = \frac{M_{n} |\psi\rangle}{\sqrt{p_{n}}}$$

$$\rho' = \sum_{n} M_{n} |\psi\rangle \langle \psi| M_{n}^{\dagger} \qquad \text{Probability cancels!}$$

$$\rho' = \sum_{n} M_{n} \rho M_{n}^{\dagger}$$

Error channels: decoherence

Actually having a physical error is very similar to a measurement



Using the same formula as for the measurement

$$\rho' = \sum_{n} E_{n} \rho E_{n}^{\dagger}$$

We can view decoherence as being a type of uncontrolled measurement

 $p_n = \operatorname{Tr}(E_n \rho E_n^{\dagger})$

Trace preserving operation

Even when an error occurs on the quantum state, we must still preserve normalization of the state, since it is still a quantum state



So we must have

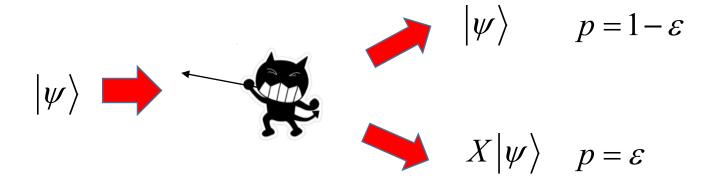
$$\operatorname{Tr}(\rho') = \operatorname{Tr}(\sum_{n} E_{n} \rho E_{n}^{\dagger}) = \operatorname{Tr}(\sum_{n} E_{n}^{\dagger} E_{n} \rho) = 1$$

Then the error channel must satisfy

$$\sum E_n^{\dagger} E_n = I$$

Example 1: bit flip channel

Consider an error occurs on a qubit with probability $\mathcal E$ where the qubit is flipped by applying X



The error channel for this is written

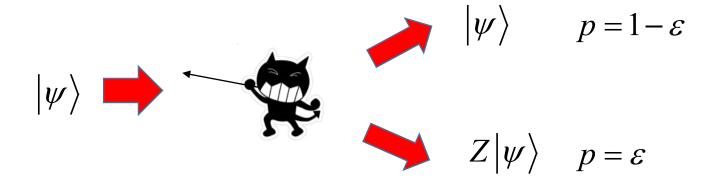
$$E_0 = \sqrt{1 - \varepsilon} I, E_1 = \sqrt{\varepsilon} X$$

This is trace-preserving

$$\sum_{n} E_{n}^{\dagger} E_{n} = (1 - \varepsilon)I + \varepsilon X^{2} = I$$

Example 2: phase flip channel

Consider an error occurs on a qubit with probability $\mathcal E$ where the qubit is flipped by applying Z



The error channel for this is written

$$E_0 = \sqrt{1 - \varepsilon} I, E_1 = \sqrt{\varepsilon} Z$$

This is trace-preserving

$$\sum_{n} E_{n}^{\dagger} E_{n} = (1 - \varepsilon)I + \varepsilon Z^{2} = I$$

Say the initial state is $\ket{+}$ and the phase flip channel acts on the state with probability ${\cal E}$. What is the final state?

$$E_0 = \sqrt{1 - \varepsilon} I, E_1 = \sqrt{\varepsilon} Z$$

$$\rho' = \sum_n E_n \rho E_n^{\dagger}$$

Say the initial state is $\ket{+}$ and the phase flip channel acts on the state with probability ${\cal E}$. What is the final state?

$$E_0 = \sqrt{1 - \varepsilon} I, E_1 = \sqrt{\varepsilon} Z \qquad \qquad \rho' = \sum_n E_n \rho E_n^{\dagger}$$

Initial state

$$\rho = |+\rangle\langle +| = \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix}$$

Apply

$$E_{0}\rho E_{0}^{\dagger} = (1-\varepsilon)|+\rangle\langle+| \qquad E_{1}\rho E_{1}^{\dagger} = \varepsilon Z|+\rangle\langle+|Z=\varepsilon|-\rangle\langle-|$$

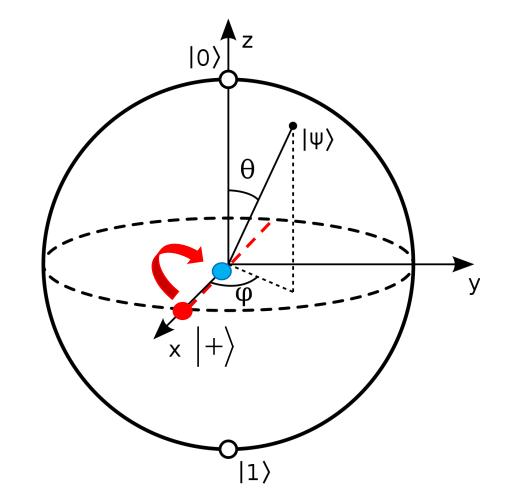
$$\rho' = (1-\varepsilon)|+\rangle\langle+|+\varepsilon|-\rangle\langle-| = \frac{1}{2}\begin{pmatrix}1 & 1-2\varepsilon\\1-2\varepsilon & 1\end{pmatrix}$$

Where is this mixed state on the Bloch sphere?

$$\rho' = \frac{1}{2} \begin{pmatrix} 1 & 1 - 2\varepsilon \\ 1 - 2\varepsilon & 1 \end{pmatrix}$$

$$\langle X \rangle = 1 - 2\varepsilon$$

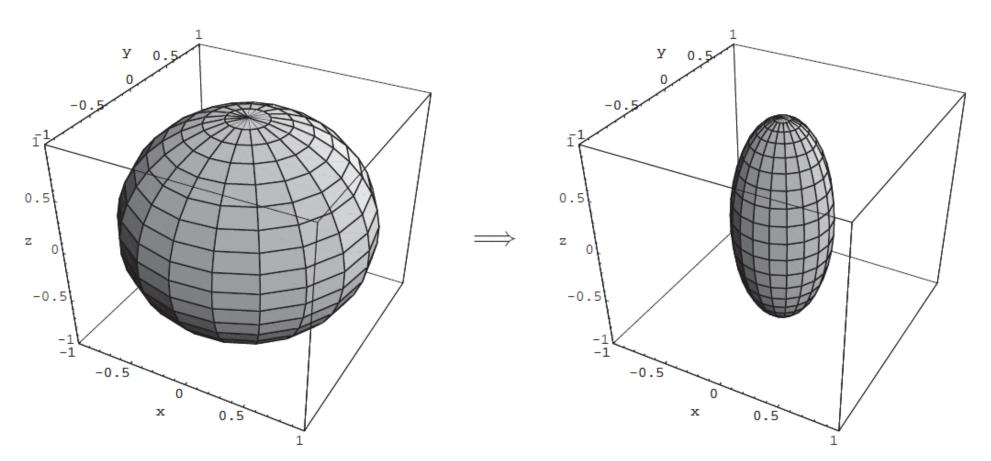
 $\langle Y \rangle = 0$
 $\langle Z \rangle = 0$



States are in the middle of the Bloch sphere. Pure states on the surface.

Dephasing channel

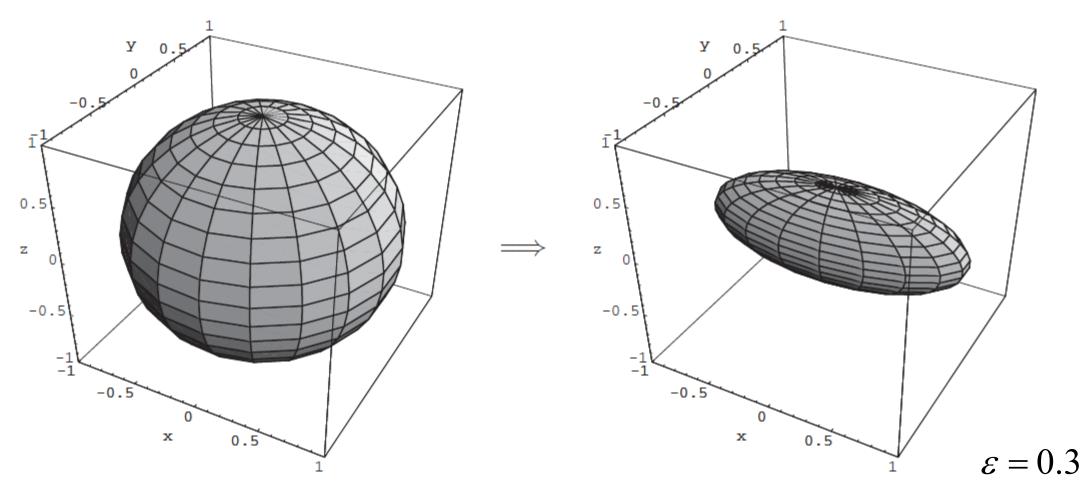
More generally we can visualize the evolution of the dephasing channel on a qubit as



 $\varepsilon = 0.3$

Bit flip channel

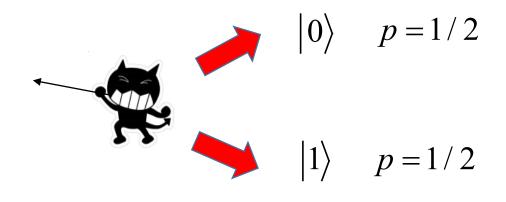
$$E_0 = \sqrt{1 - \varepsilon} I, E_1 = \sqrt{\varepsilon} X$$

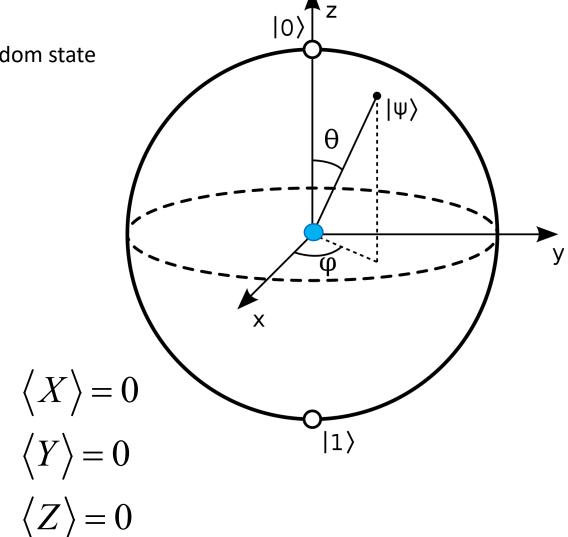


Completely mixed qubit state

The state at the middle of the Bloch sphere is a completely random state

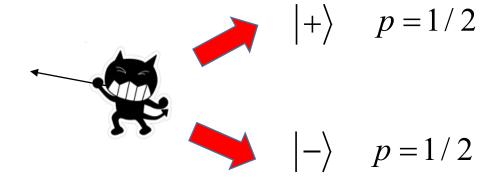
$$\rho = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$





Basis invariance of completely mixed state

Actually it doesn't have to be a mixture of $|0\rangle, |1\rangle$



$$\rho = \frac{1}{2} |+\rangle \langle +|+\frac{1}{2}|-\rangle \langle -|=\frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

Or ANY 50/50 mixture of orthogonal states will produce this same state

This is a special property of the completely mixed state

$$\rho = \frac{1}{2}I = \frac{1}{2}UIU^{\dagger} = \frac{1}{2}I$$

$$\rho = \frac{I}{d}$$
 For a d-dimensional system

Verify that a completely mixed state of the states

$$\frac{|\theta\rangle = \cos\theta |0\rangle + \sin\theta |1\rangle}{|\theta\rangle = \sin\theta |0\rangle - \cos\theta |1\rangle}$$
 gives
$$\rho = \frac{1}{2}I$$

Verify that a completely mixed state of the states

$$\frac{|\theta\rangle = \cos\theta |0\rangle + \sin\theta |1\rangle}{|\theta\rangle = \sin\theta |0\rangle - \cos\theta |1\rangle}$$
 gives
$$\rho = \frac{1}{2}I$$

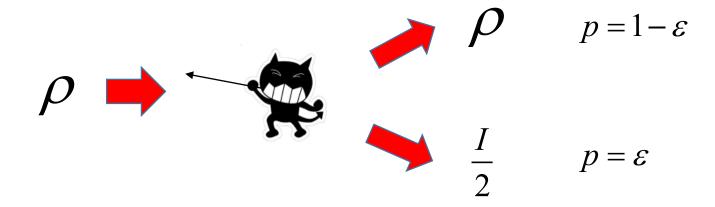
$$\rho = \frac{1}{2} |\theta\rangle \langle\theta| + \frac{1}{2} |\overline{\theta}\rangle \langle\overline{\theta}|$$

$$= \frac{1}{2} \begin{pmatrix} \cos^2\theta & \cos\theta\sin\theta \\ \cos\theta\sin\theta & \sin^2\theta \end{pmatrix} + \frac{1}{2} \begin{pmatrix} \sin^2\theta & -\cos\theta\sin\theta \\ -\cos\theta\sin\theta & \cos^2\theta \end{pmatrix}$$

$$= \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

Example 3: Depolarizing channel

A more symmetric type of error is the depolarizing channel



The final state is

$$\rho \to \varepsilon \frac{I}{2} + (1 - \varepsilon)\rho$$

$$\langle X \rangle = (1 - \varepsilon) \langle X \rangle_0$$
$$\langle Y \rangle = (1 - \varepsilon) \langle Y \rangle_0$$
$$\langle Z \rangle = (1 - \varepsilon) \langle Z \rangle_0$$

Depolarizing channel

