

HIGGS AT HADRON COLLIDER

Higgs Properties and Precision Test

Lecture 12

DIPARTIMENTO DI FISICA



SAPIENZA
UNIVERSITÀ DI ROMA

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Fisica delle Particelle Elementari, Anno Accademico 2015-16

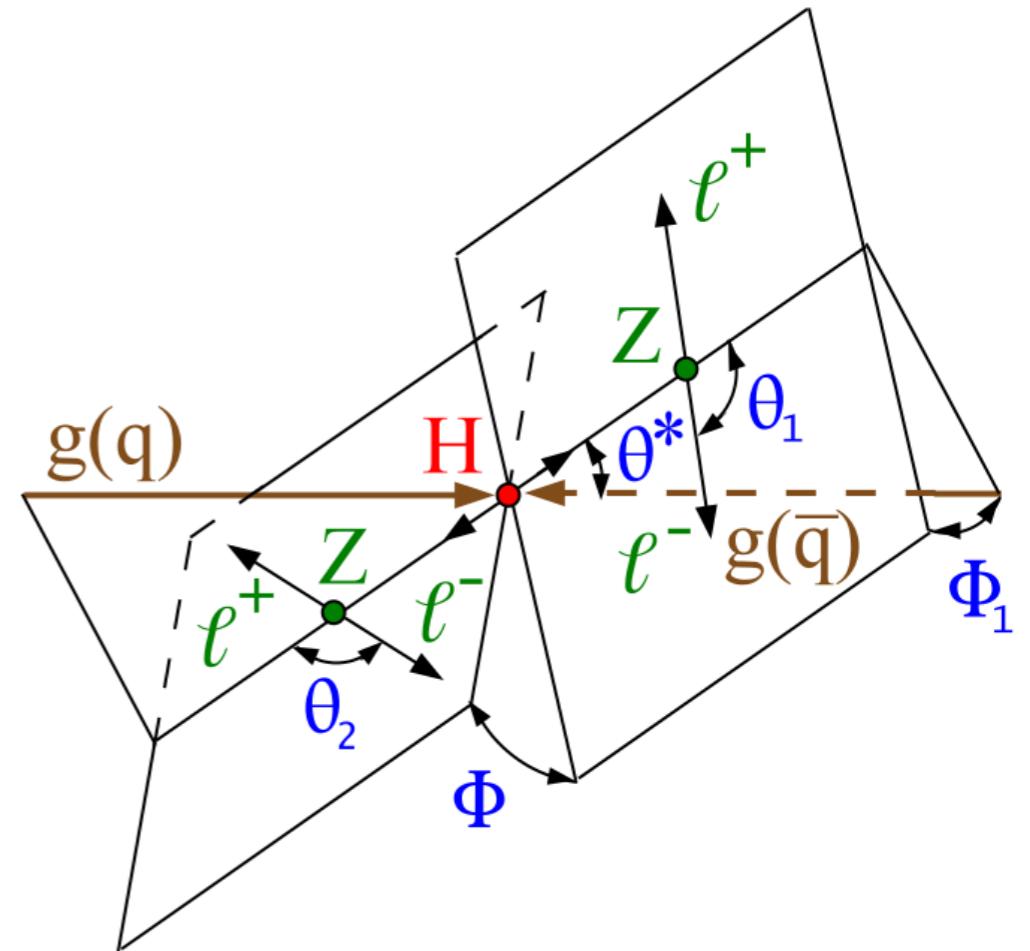
<http://www.roma1.infn.it/people/rahatlou/particelle/>

WHY AND WHICH Boson?

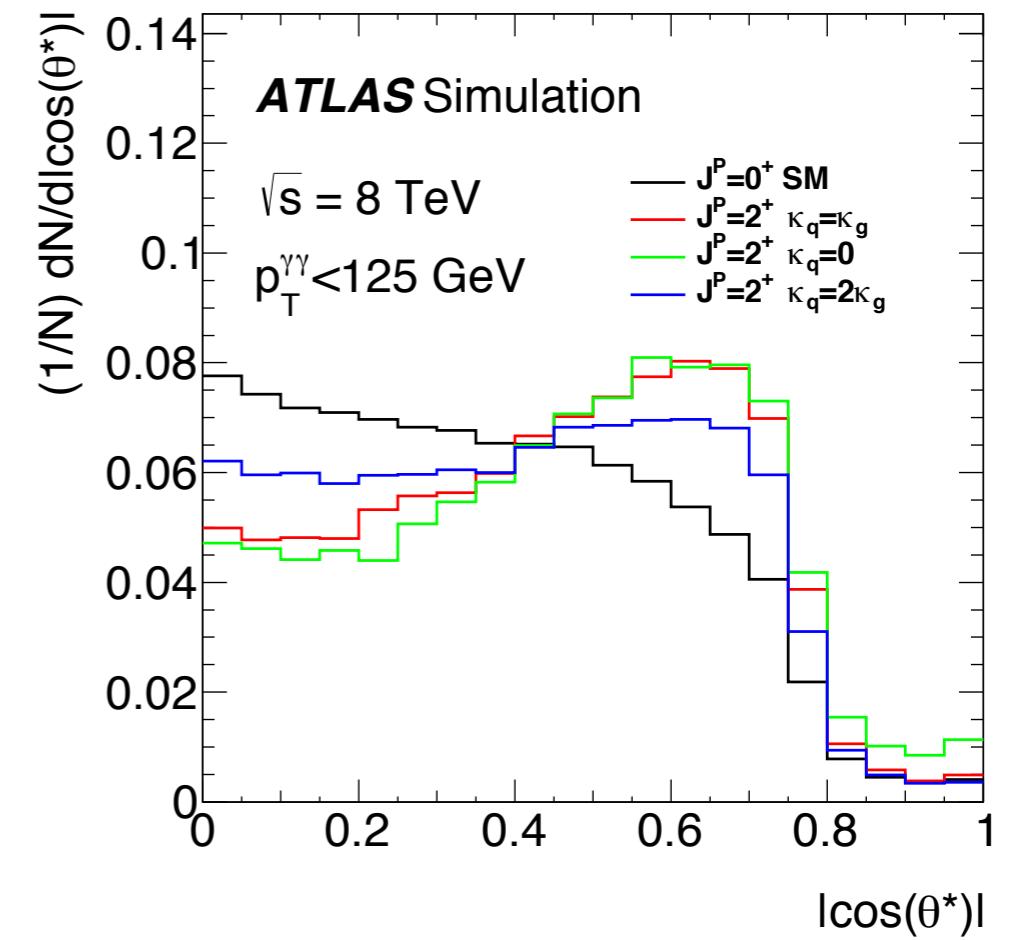
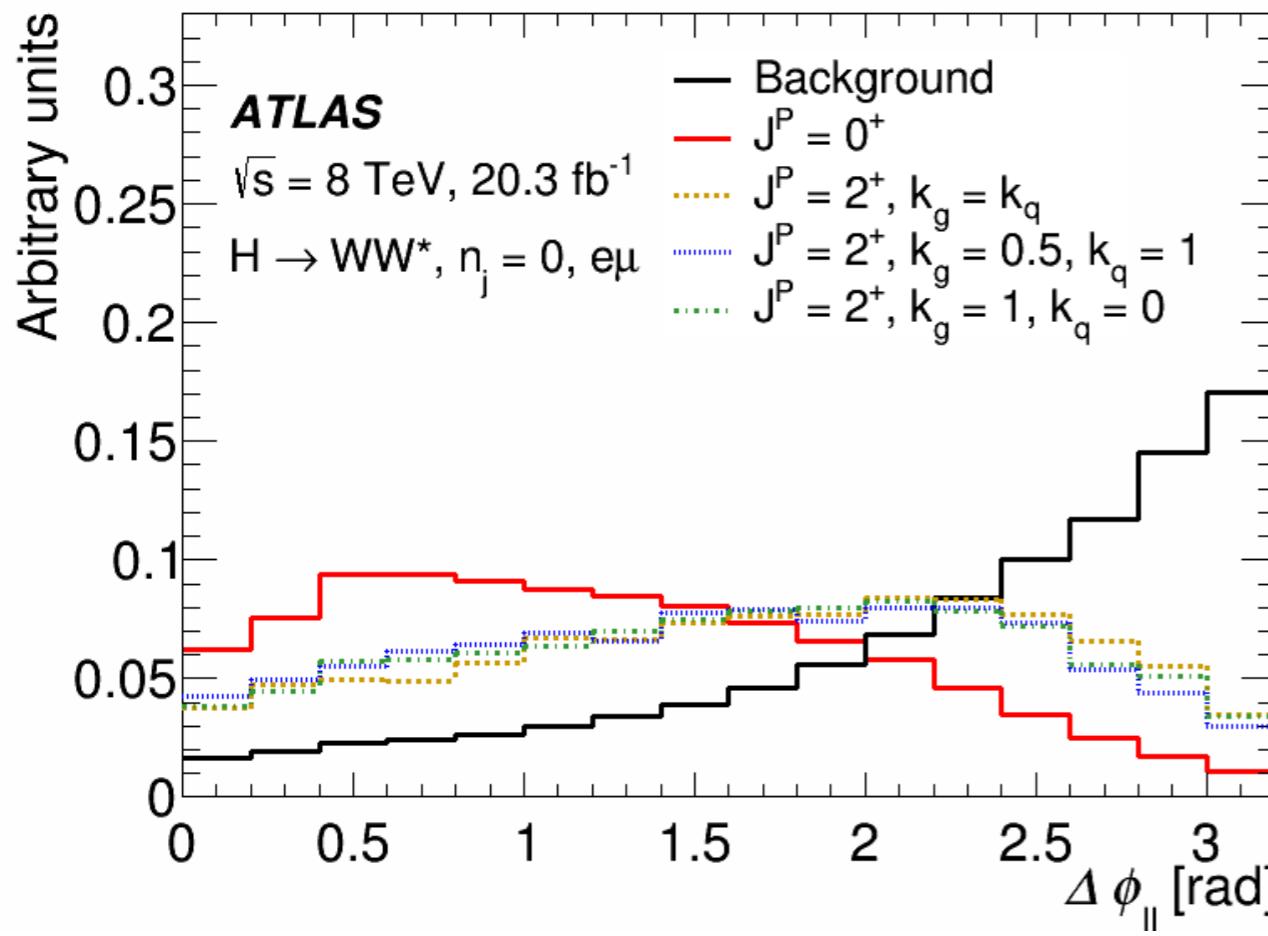
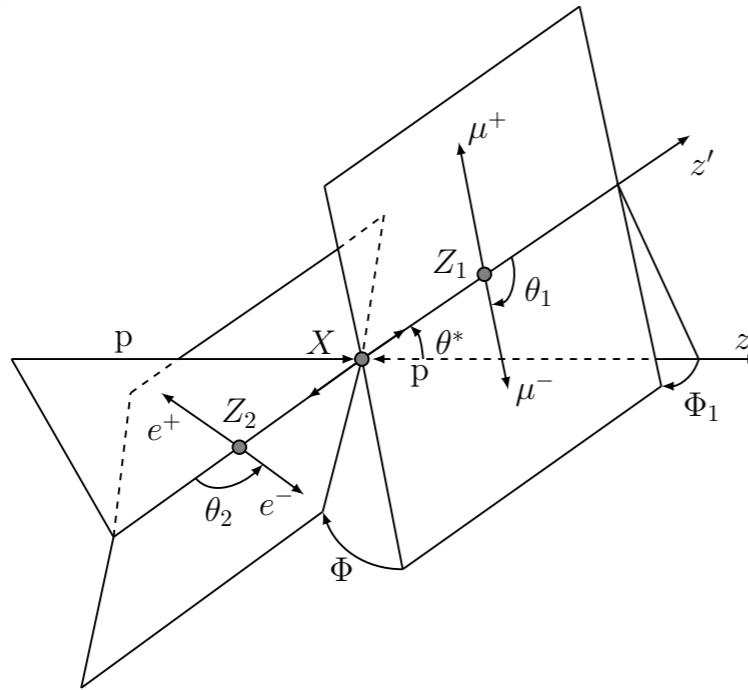
- Is it statistically significant?
 - use of p-value. Put threshold at 5σ
- Is it a boson?
 - determined by decay products
- Which is the mass?
 - to be studied with modes with good resolution, i.e. ZZ and $\gamma\gamma$
- Is it “the” SM Higgs boson?
 - study compatibility of $BR^*\sigma$ in different decay modes
- Is it “a” Higgs boson?
 - study the Higgs couplings to different particles
 - search for new bosons

SPIN AND PARITY

- SM Higgs is a scalar boson with $J^P = 0^+$
- Observation of $H \rightarrow \gamma\gamma$ decays implies $J \neq 1$
 - Landau-Yang theorem
- Observation of WW and ZZ decays disfavors CP-odd hypothesis
- Use angular kinematic variables from final states to distinguish different spin and parity hypotheses
 - $m_{||}, \Delta\Phi_{||}, m_T, \cos\theta^*$
- No variable alone powerful enough
- Combine all information with multivariate method e.g. neural network or boosted decision tree

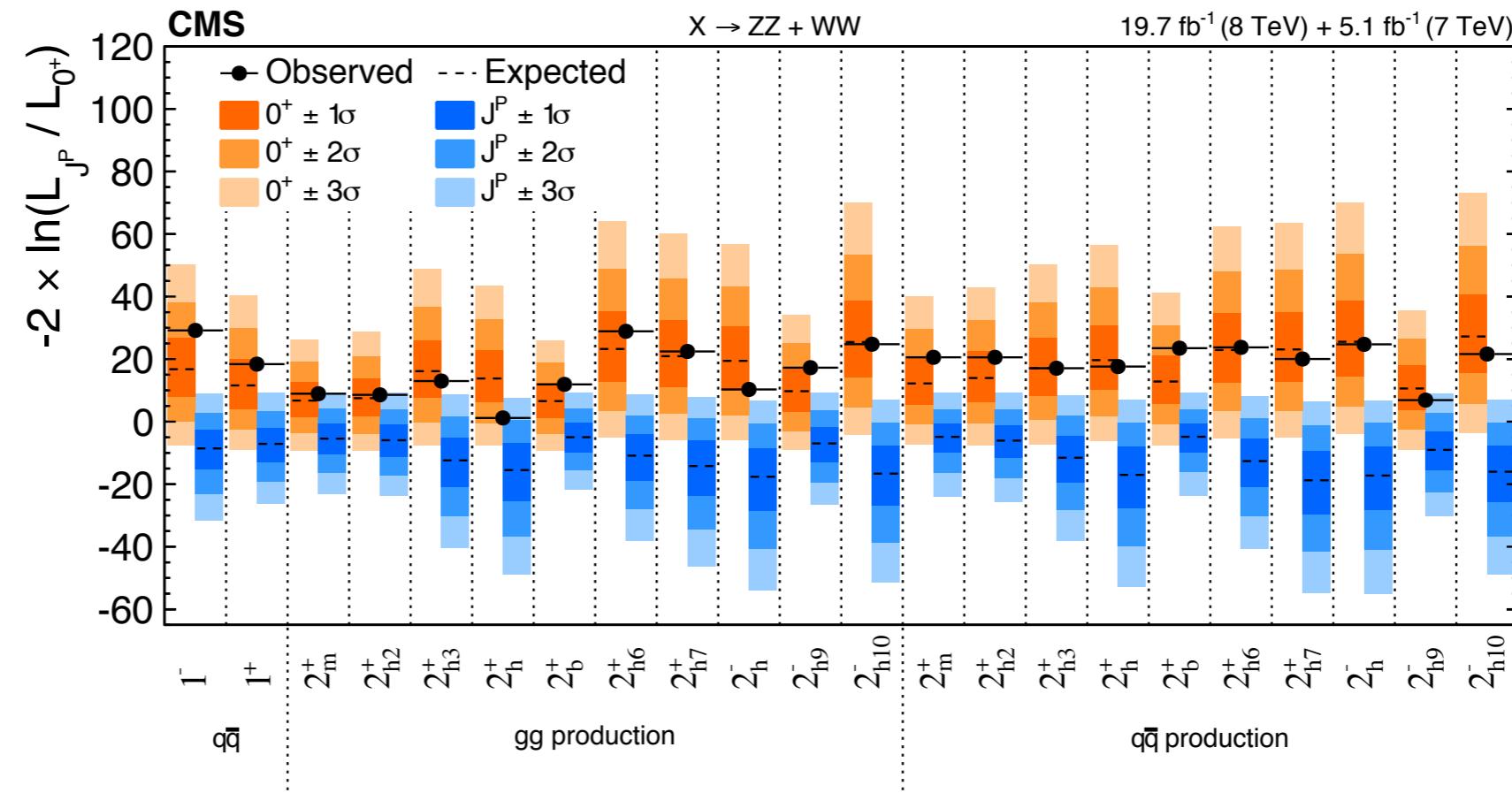
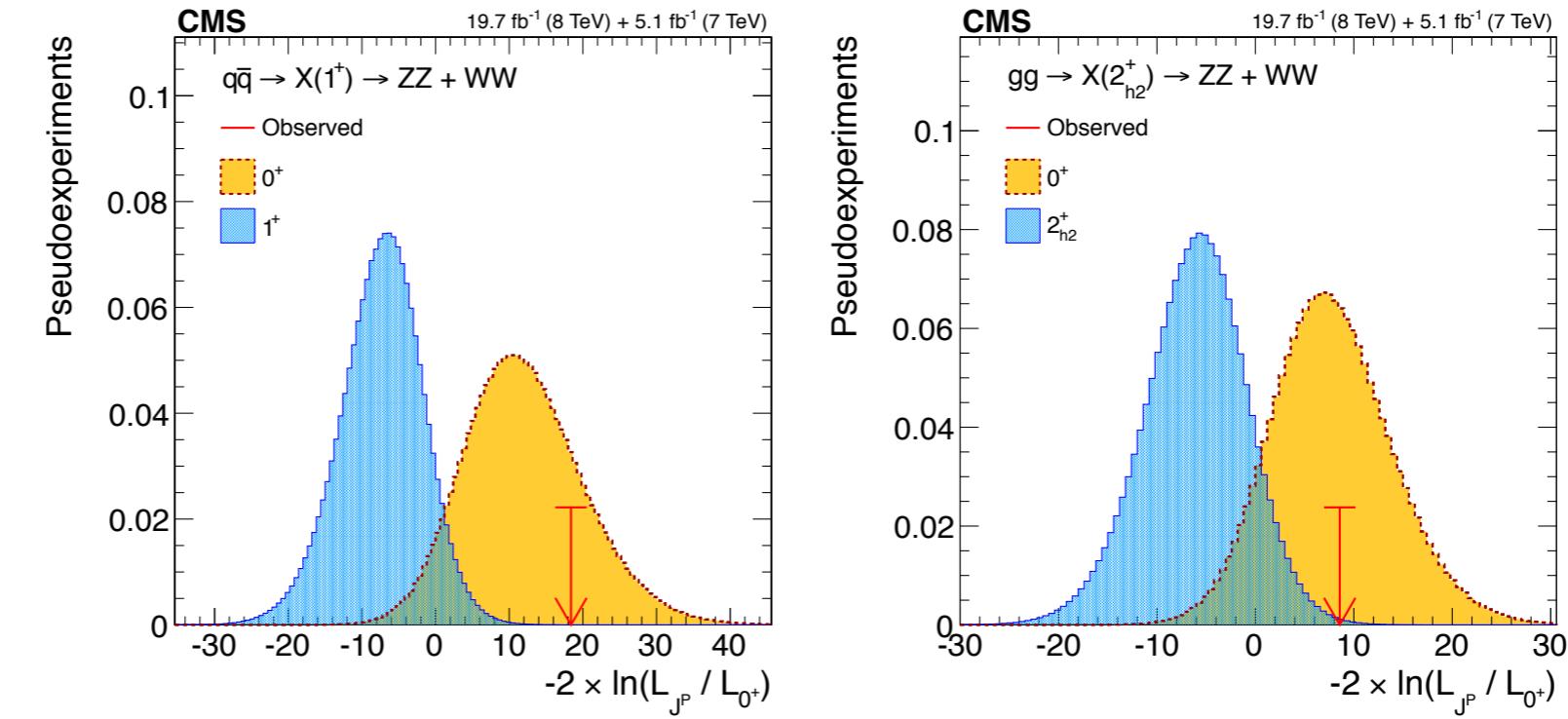


ANGULAR VARIABLES



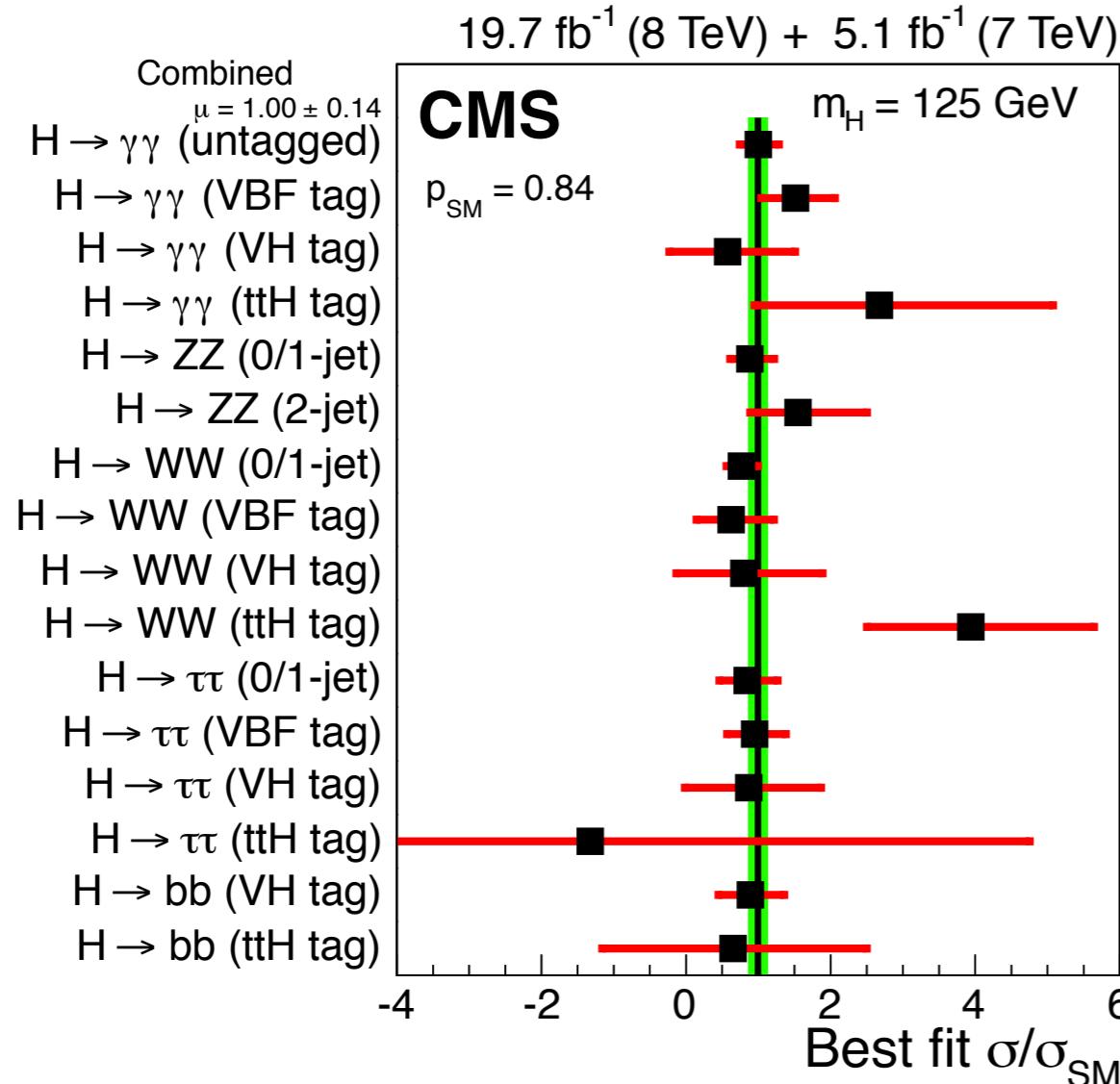
ATLAS: arXiv:1506.05669

CMS: arXiv:1411.3441



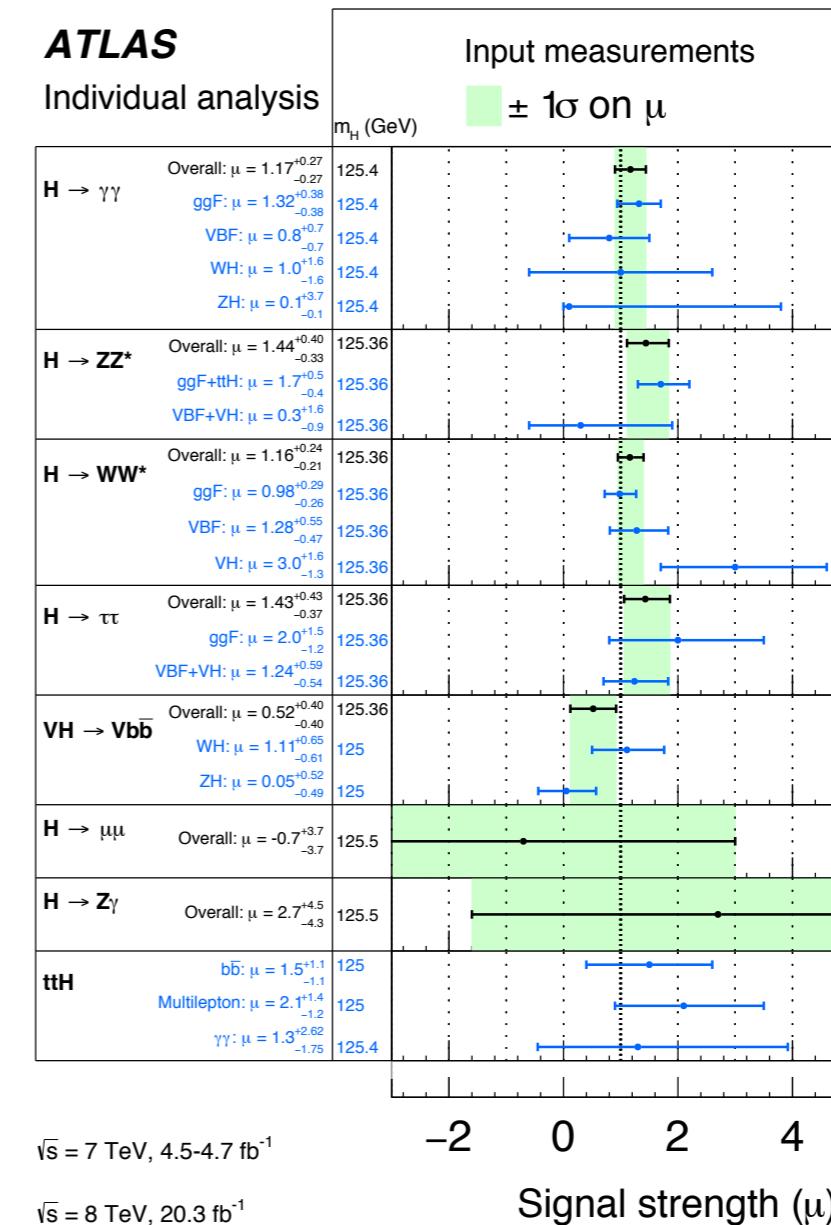
COMPATIBILITY WITH STANDARD MODEL

CMS: Eur. Phys. J. C 75 (2015) 212



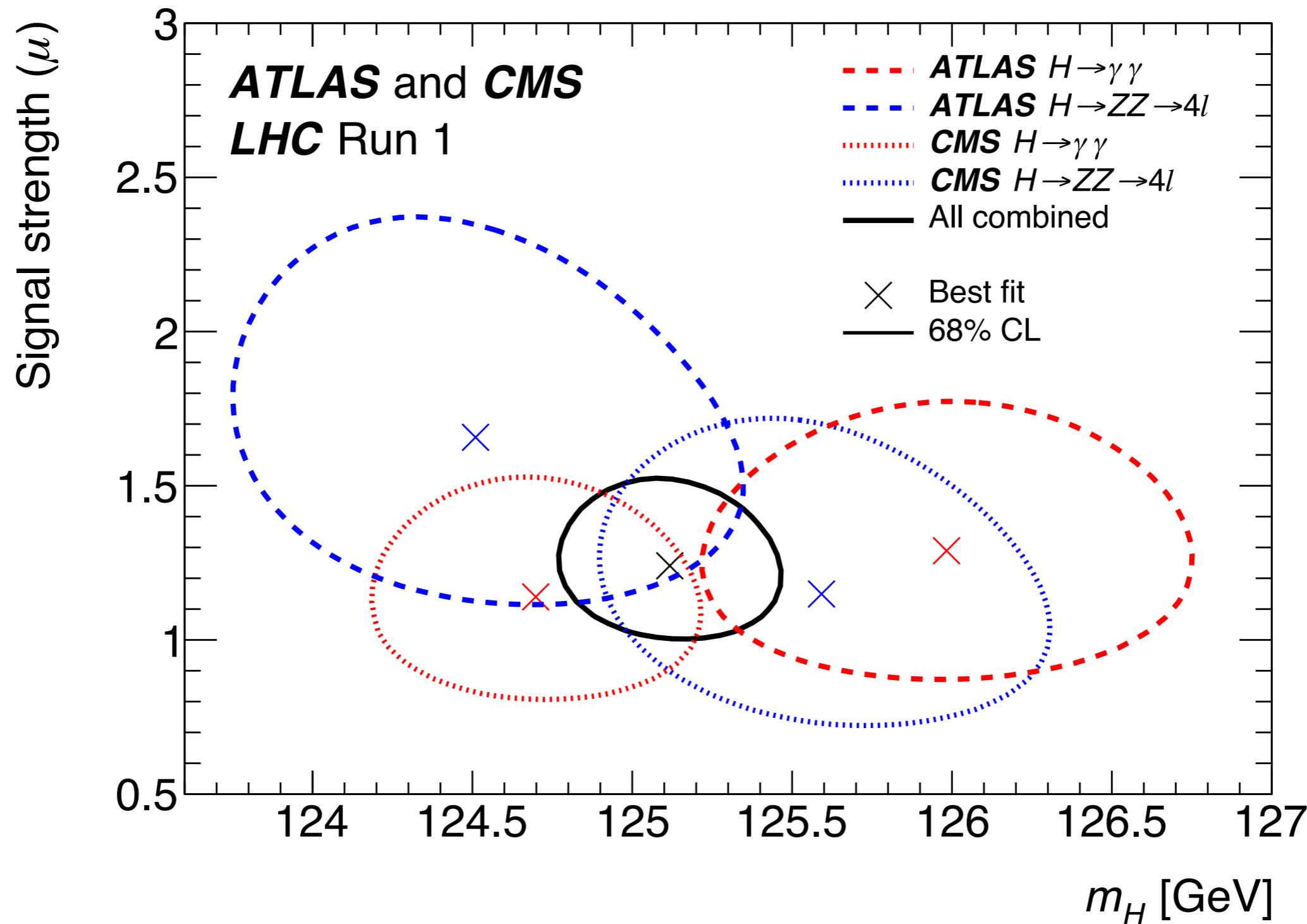
- Best fit SM strength
 - CMS $\mu = \sigma/\sigma_{\text{SM}} = 1.00 \pm 0.14$
 - ATLAS $\mu = \sigma/\sigma_{\text{SM}} = 1.17 \pm 0.27$
- Good agreement among modes

ATLAS: arxiv:1507.04548

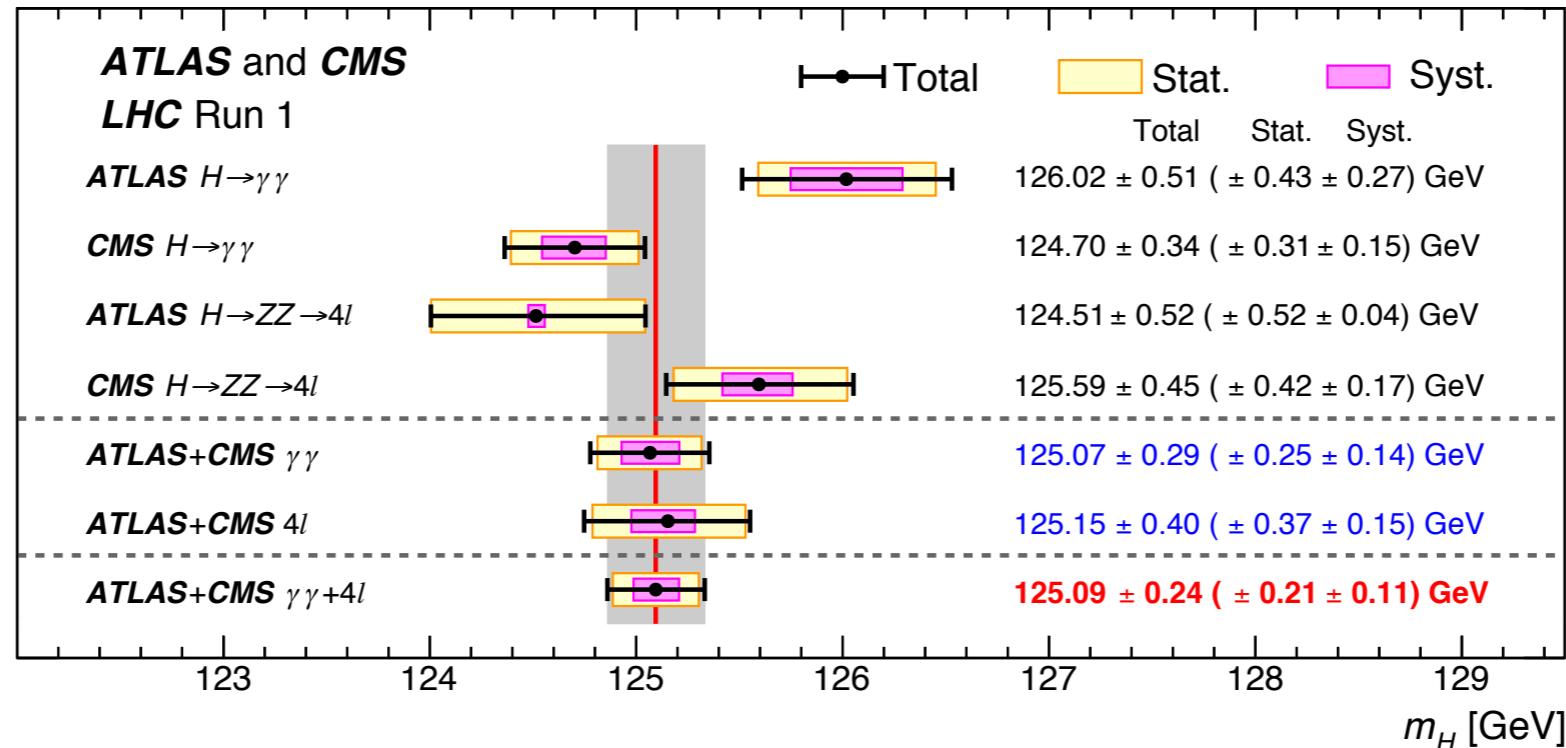
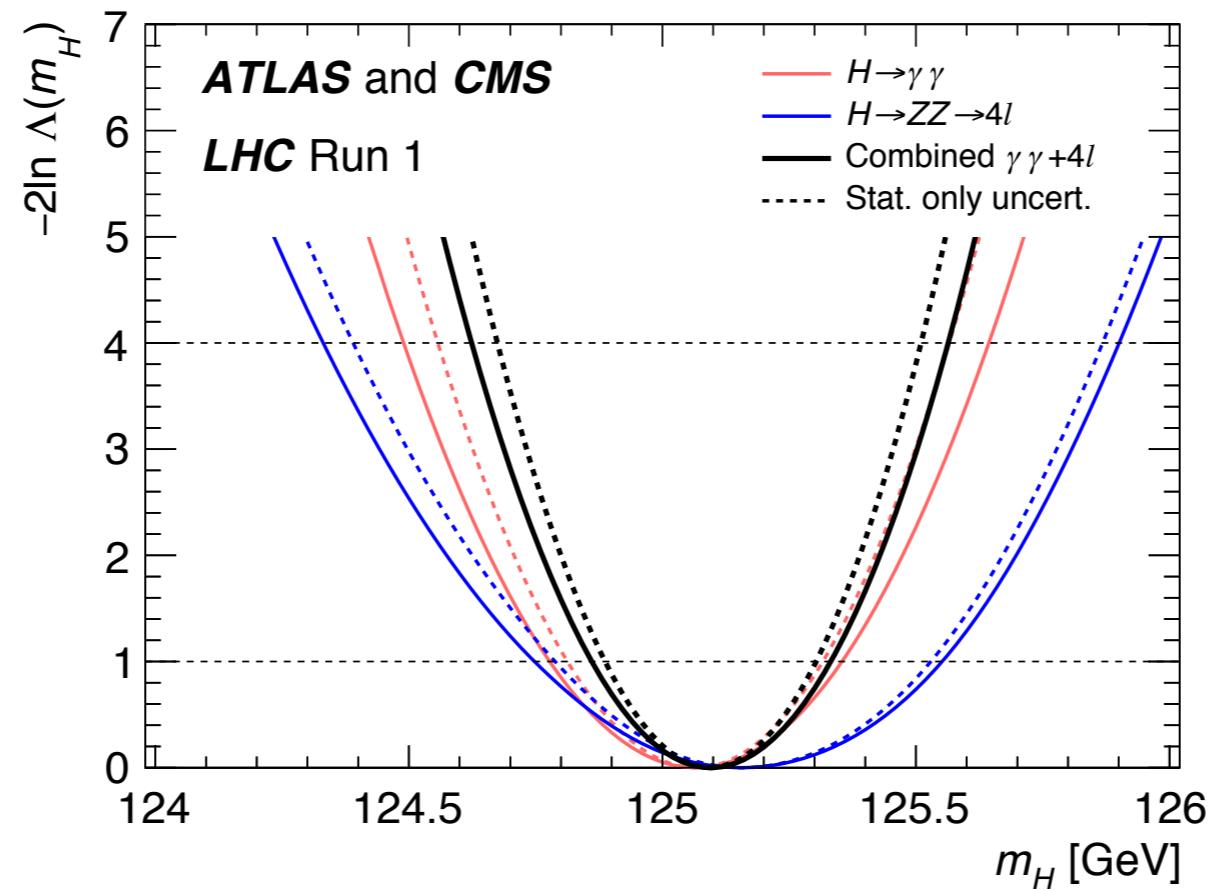


ATLAS and CMS Combination

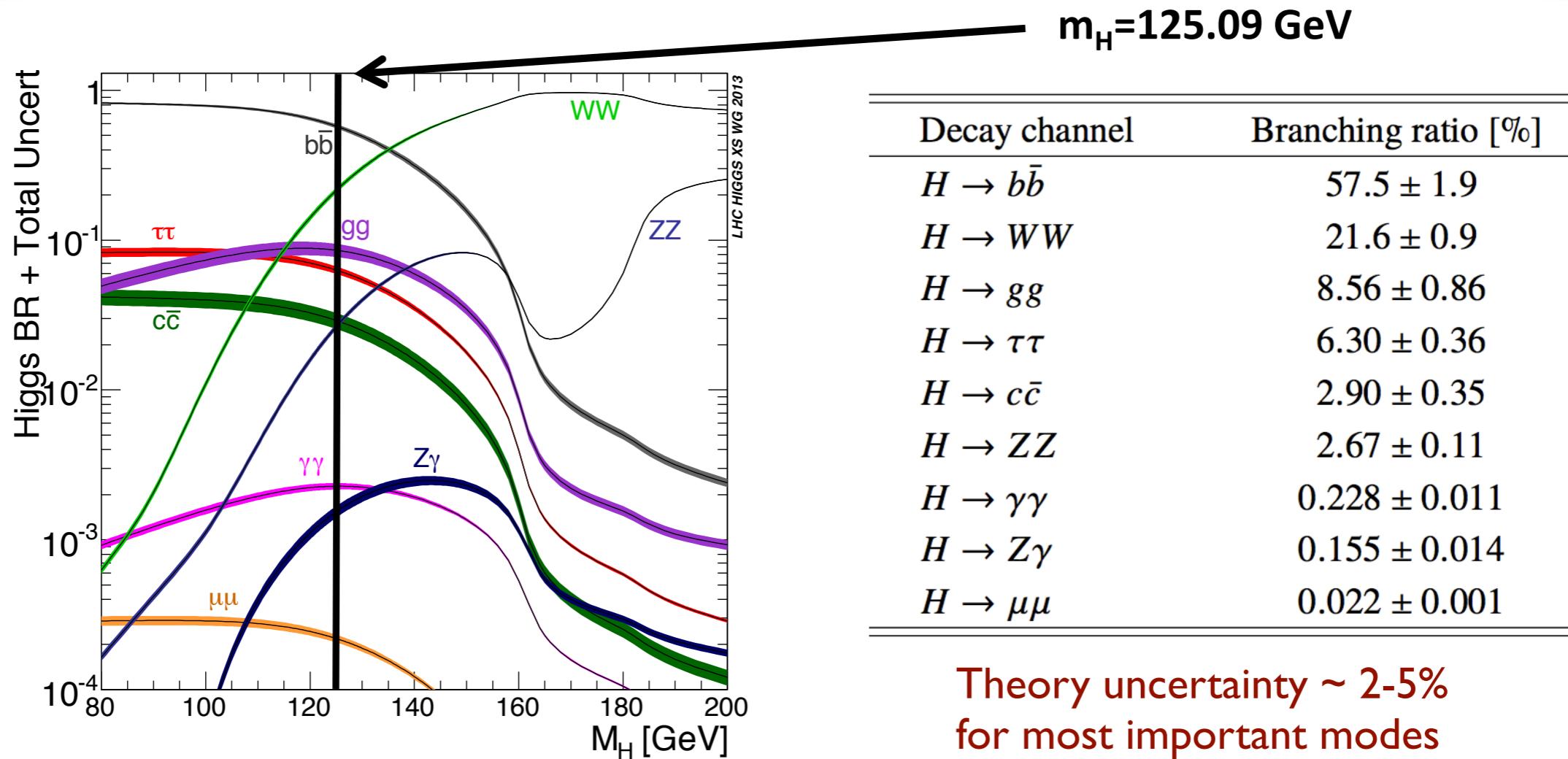
MASS VS SIGNAL STRENGTH



HIGGS MASS RUN 1



HIGGS WIDTH



- Natural width of Higgs boson expected to be very small $\sim 4.1 \text{ MeV}$
 - Mass resolution larger by $\times 10\text{-}100$!
- Measurement of Higgs width very challenging

COMPATIBILITY WITH STANDARD MODEL

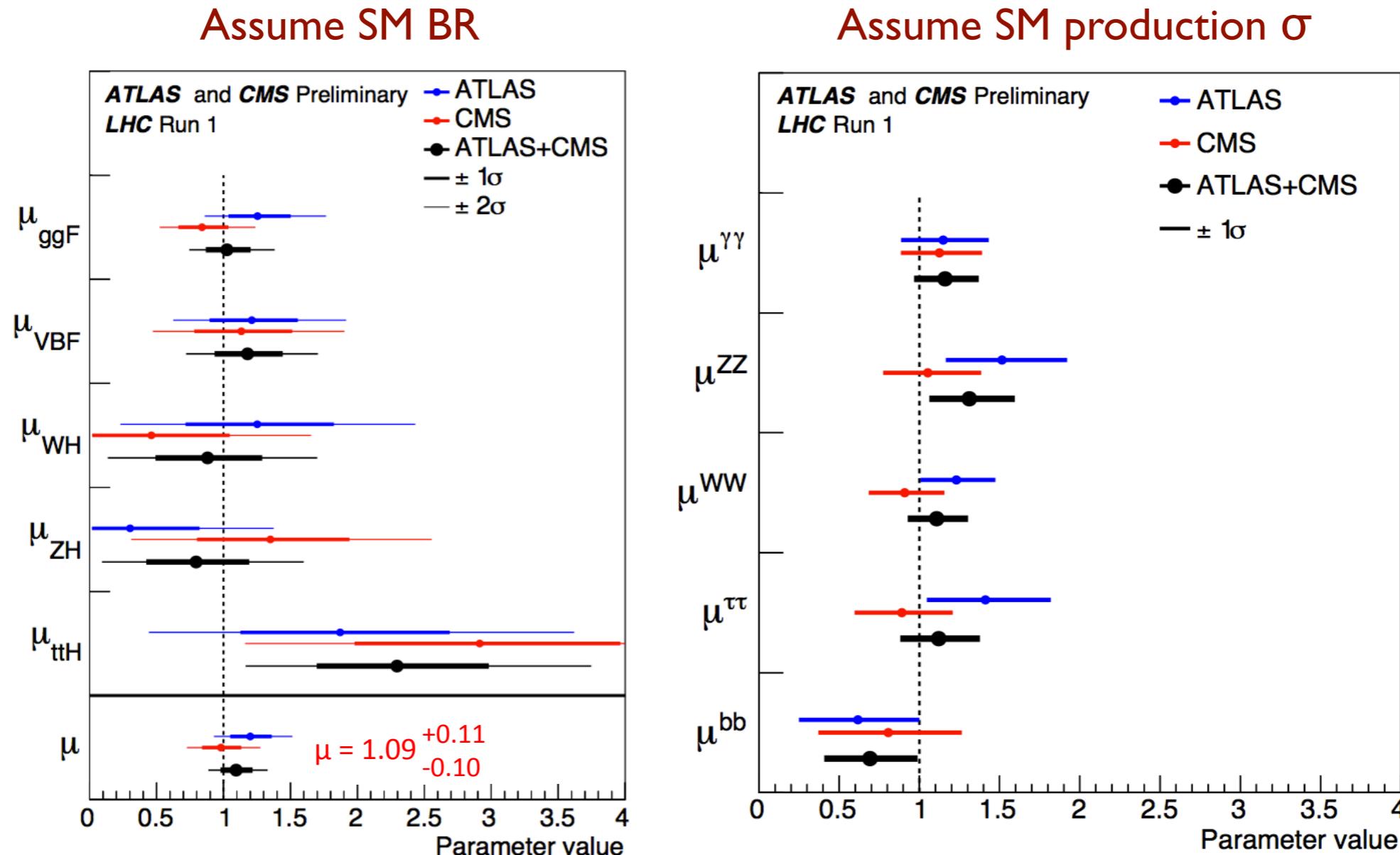
- Measure spin and parity of new boson and compare to Standard Model predicted scalar boson
 - angular distribution of decay products depends on these quantum numbers
- Assume a SM-like Higgs with $J^P = 0+$ with narrow width
 - decouple production and decay process
- Measure cross section and branching fractions and compare to SM
 - expect small deviations from SM predictions
 - at LHC only measure the product $\sigma \times BR$ not each separately
 - ▶ unless you make further assumptions
- Three general approaches to test compatibility
 - Measure of signal strengths μ for σ and BR relative to SM
 - Introduce modifier K for each coupling constant relative to SM (κ -framework)
 - Generic parameterization of σ and BR , measuring a mixture of σ and μ

$$\sigma_i \cdot BR^f = \frac{\sigma_i(\vec{\kappa}) \cdot \Gamma^f(\vec{\kappa})}{\Gamma_H}$$

SIGNAL STRENGTH TEST

$$\mu_i = \frac{\sigma_i}{\sigma_i^{\text{SM}}} \quad \text{and} \quad \mu^f = \frac{\text{BR}^f}{\text{BR}_{\text{SM}}^f} \quad \mu_i^f \equiv \frac{\sigma_i \cdot \text{BR}^f}{(\sigma_i \cdot \text{BR}^f)_{\text{SM}}} = \mu_i \times \mu^f$$

- Most constrained parameterization
 - one single signal strength μ
 - assuming also same value at 7 and 8 TeV

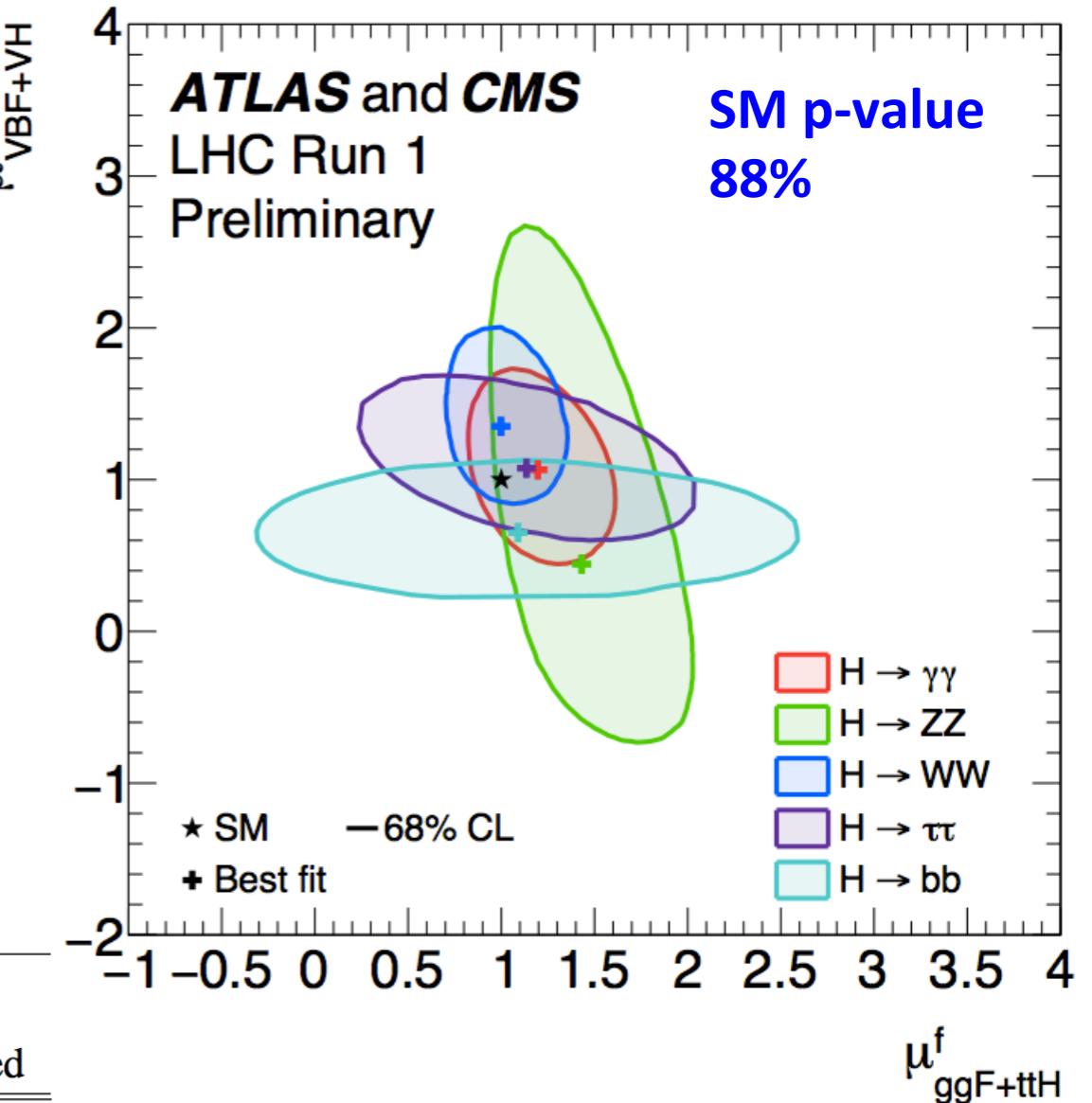


Bosons (μ_V)vs Fermions(μ_F)

- Can also fit μ_V^f vs μ_F^f per decay:
 - $\mu_V^f = \mu_{VBF+VH}^f$
 - $\mu_F^f = \mu_{ggF+ttH}^f$
- μ_V/μ_f can be measured in the different decay channels and combined:

$$\mu_V/\mu_f = 1.06^{+0.35}_{-0.27}$$

Parameter	ATLAS+CMS observed	ATLAS+CMS expected unc.	ATLAS observed	CMS observed
μ_V/μ_F	$1.06^{+0.35}_{-0.27}$	$+0.34$ -0.26	$0.91^{+0.41}_{-0.30}$	$1.29^{+0.67}_{-0.46}$
$\mu_F^{\gamma\gamma}$	$1.13^{+0.24}_{-0.21}$	$+0.21$ -0.19	$1.18^{+0.33}_{-0.29}$	$1.03^{+0.30}_{-0.26}$
μ_F^{ZZ}	$1.29^{+0.29}_{-0.25}$	$+0.24$ -0.20	$1.54^{+0.44}_{-0.36}$	$1.00^{+0.33}_{-0.27}$
μ_F^{WW}	$1.08^{+0.22}_{-0.19}$	$+0.19$ -0.17	$1.26^{+0.29}_{-0.25}$	$0.85^{+0.25}_{-0.22}$
$\mu_F^{\tau\tau}$	$1.07^{+0.35}_{-0.28}$	$+0.32$ -0.27	$1.50^{+0.66}_{-0.49}$	$0.75^{+0.39}_{-0.29}$
μ_F^{bb}	$0.65^{+0.37}_{-0.28}$	$+0.45$ -0.34	$0.67^{+0.58}_{-0.42}$	$0.64^{+0.54}_{-0.36}$



SM p-value
62%

K-Framework Test

- Scale Higgs boson couplings by coupling modifiers κ

$$\kappa_j^2 = \sigma_j / \sigma_j^{\text{SM}} \quad \text{or} \quad \kappa_j^2 = \Gamma_j / \Gamma_{\text{SM}}^j$$

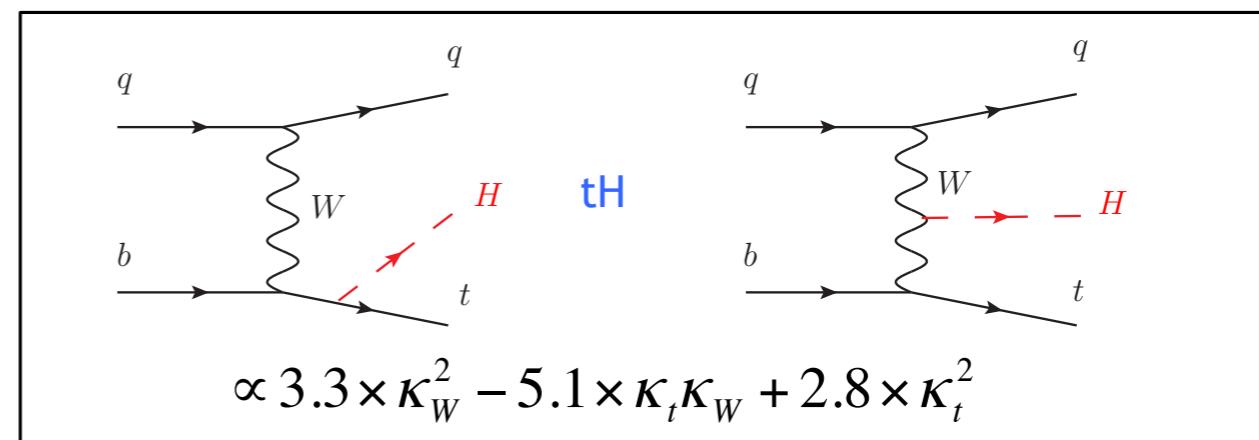
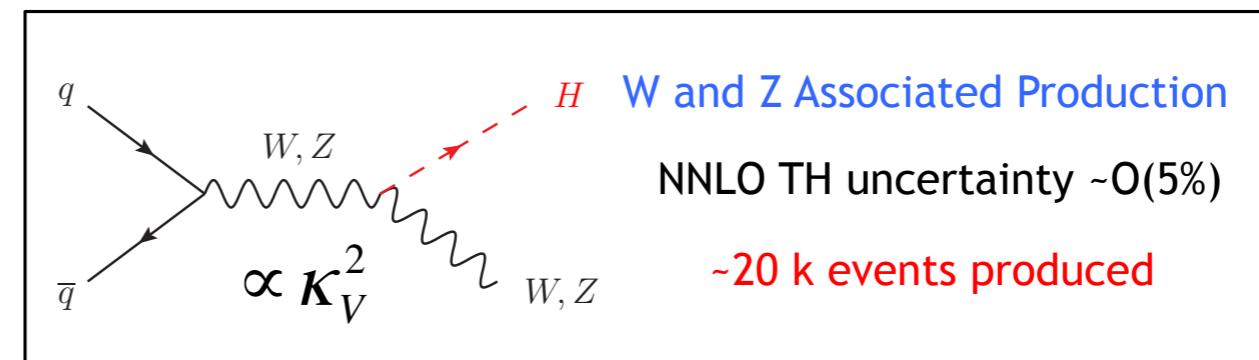
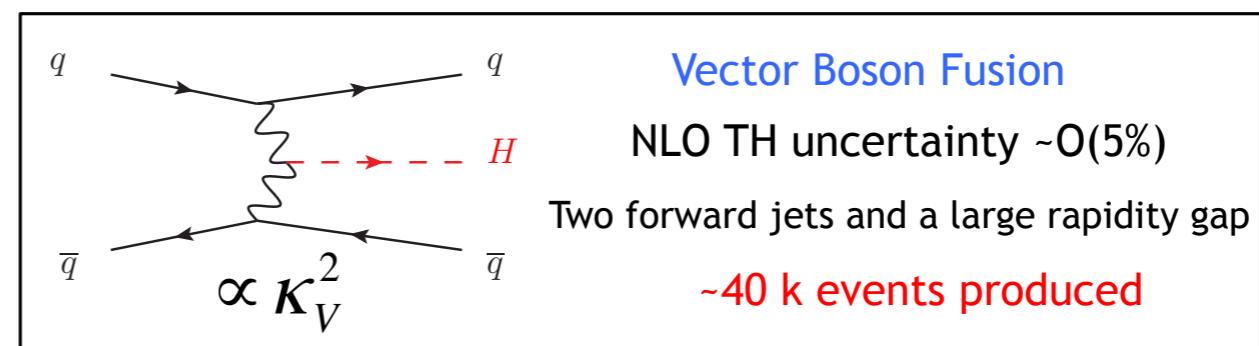
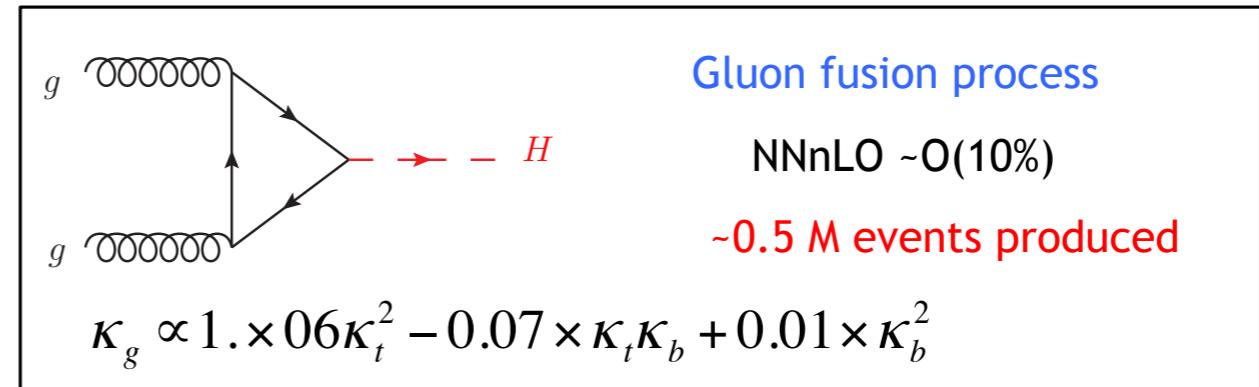
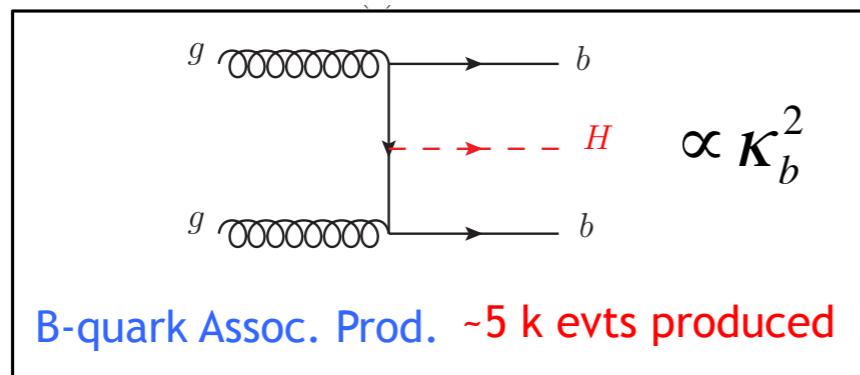
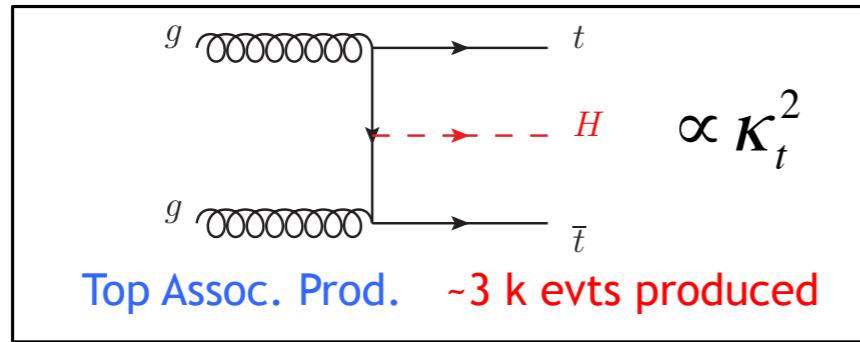
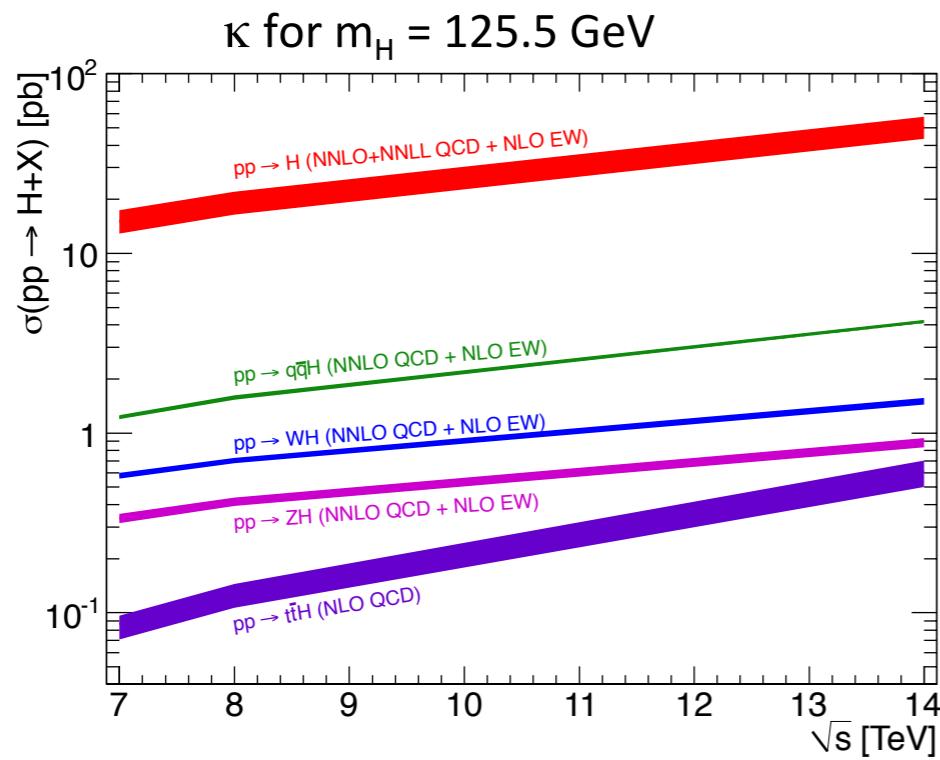
$$\kappa_H^2 = \sum_j \text{BR}_{\text{SM}}^j \kappa_j^2 \quad \Gamma_H = \frac{\kappa_H^2 \cdot \Gamma_H^{\text{SM}}}{1 - \text{BR}_{\text{BSM}}}$$

- Allow for possible new invisible or undetected Higgs decays
 - Invisible final states possibly due to new non-SM particles
 - undetected decays could be not yet measured and small decays like $cc\bar{c}\bar{c}$
- New decays and particles have two possible effects
 - $m_X < m_H/2$: BR_{BSM} will be affected
 - $m_X > m_H/2$: κ_i will be affected

HIGGS COUPLINGS

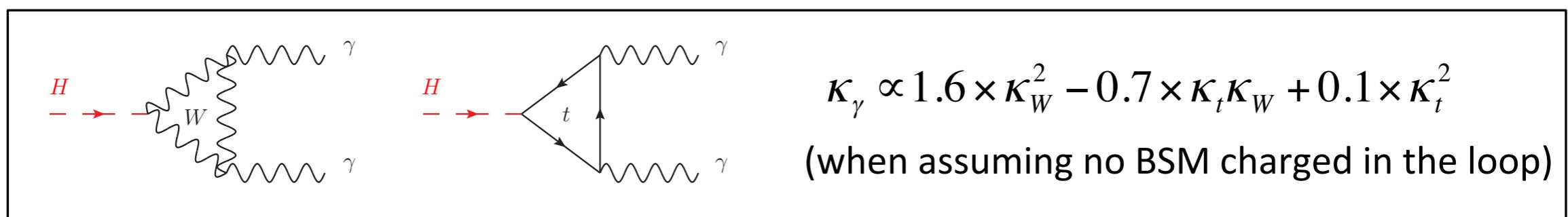
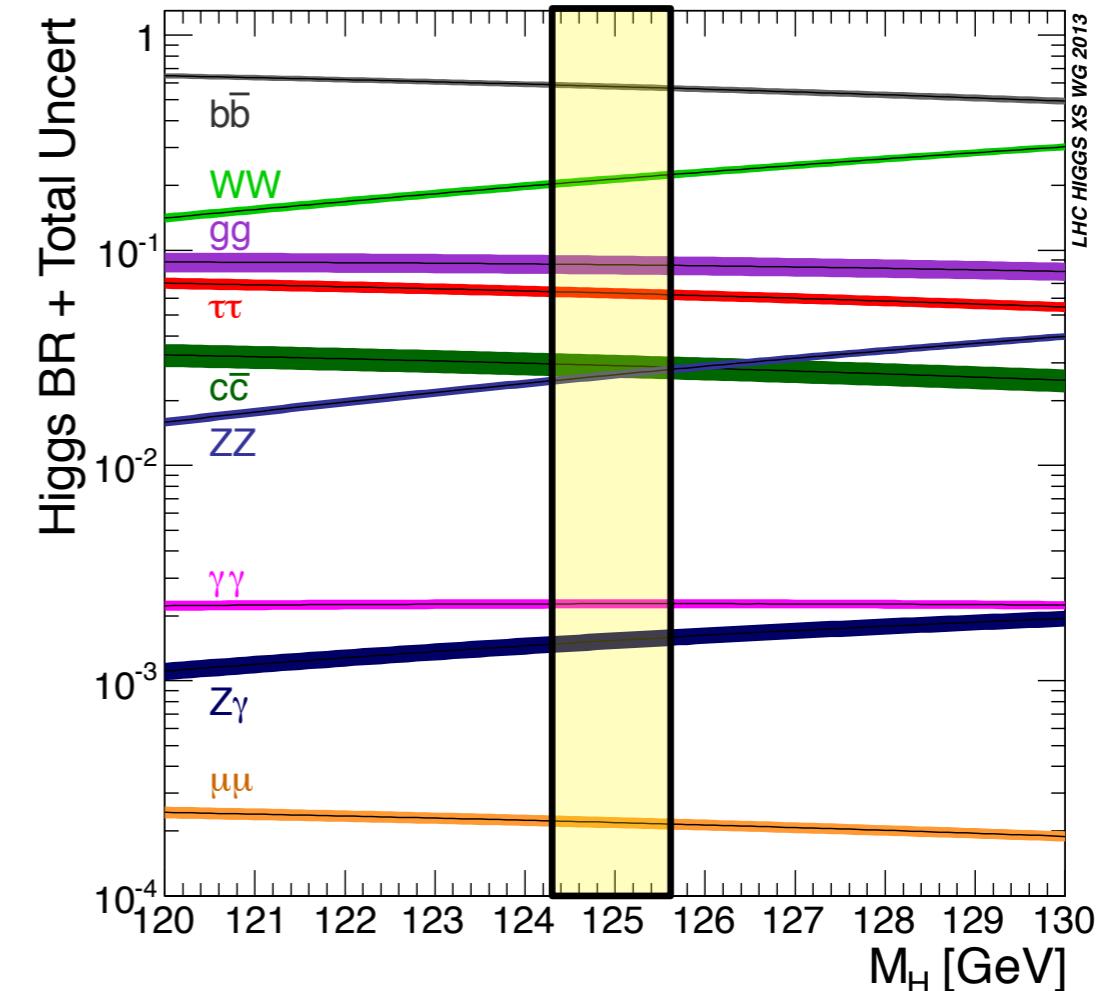
Production modes		Detectable decay modes		Currently undetectable decay modes	
$\frac{\sigma_{ggH}}{\sigma_{ggH}^{SM}}$	$= \begin{cases} \kappa_g^2(\kappa_b, \kappa_t, m_H) \\ \kappa_g^2 \end{cases}$	$\frac{\Gamma_{WW^{(*)}}}{\Gamma_{WW^{(*)}}^{SM}} = \kappa_W^2$		$\frac{\Gamma_{t\bar{t}}}{\Gamma_{t\bar{t}}^{SM}} = \kappa_t^2$	
$\frac{\sigma_{VBF}}{\sigma_{VBF}^{SM}}$	$= \kappa_{VBF}^2(\kappa_W, \kappa_Z, m_H)$	$\frac{\Gamma_{ZZ^{(*)}}}{\Gamma_{ZZ^{(*)}}^{SM}} = \kappa_Z^2$		$\frac{\Gamma_{gg}}{\Gamma_{gg}^{SM}} : \text{see Section 3.1.2}$	
$\frac{\sigma_{WH}}{\sigma_{WH}^{SM}}$	$= \kappa_W^2$	$\frac{\Gamma_{b\bar{b}}}{\Gamma_{b\bar{b}}^{SM}} = \kappa_b^2$		$\frac{\Gamma_{c\bar{c}}}{\Gamma_{c\bar{c}}^{SM}} = \kappa_t^2$	
$\frac{\sigma_{ZH}}{\sigma_{ZH}^{SM}}$	$= \kappa_Z^2$	$\frac{\Gamma_{\tau^-\tau^+}}{\Gamma_{\tau^-\tau^+}^{SM}} = \kappa_\tau^2$		$\frac{\Gamma_{s\bar{s}}}{\Gamma_{s\bar{s}}^{SM}} = \kappa_b^2$	
$\frac{\sigma_{t\bar{t}H}}{\sigma_{t\bar{t}H}^{SM}}$	$= \kappa_t^2$	$\frac{\Gamma_{\gamma\gamma}}{\Gamma_{\gamma\gamma}^{SM}} = \begin{cases} \kappa_\gamma^2(\kappa_b, \kappa_t, \kappa_\tau, \kappa_W, m_H) \\ \kappa_\gamma^2 \end{cases}$	$\frac{\Gamma_{\mu^-\mu^+}}{\Gamma_{\mu^-\mu^+}^{SM}} = \kappa_\tau^2$	$\frac{\Gamma_H}{\Gamma_H^{SM}} = \begin{cases} \kappa_H^2(\kappa_i, m_H) \\ \kappa_H^2 \end{cases}$	
		$\frac{\Gamma_{Z\gamma}}{\Gamma_{Z\gamma}^{SM}} = \begin{cases} \kappa_{(Z\gamma)}^2(\kappa_b, \kappa_t, \kappa_\tau, \kappa_W, m_H) \\ \kappa_{(Z\gamma)}^2 \end{cases}$			
Total width		Yukawa sector			
		$\Gamma_{t\bar{t}}$	τ	Z	Gauge sector
		Γ_{WW}	γ	W	
		Γ_{ZZ}	b	c	Mixed sector
		$\Gamma_{\tau\tau}$	t	\bar{t}	
		$\Gamma_{\mu\mu}$	g	γ	
		$\Gamma_{\gamma\gamma}$	b	t	
		$\Gamma_{Z\gamma}$	g	\bar{t}	
					Quark loop
					Loops (γ, g) are sensitive to BSM contributions.
Production Loops Interference Expression in fundamental coupling-strength scale factors					
$\sigma(ggF)$	✓	$b-t$	$\kappa_g^2 \sim 1.06 \cdot \kappa_t^2 + 0.01 \cdot \kappa_b^2 - 0.07 \cdot \kappa_t \kappa_b$		
$\sigma(VBF)$	-	-	$\sim 0.74 \cdot \kappa_W^2 + 0.26 \cdot \kappa_Z^2$		
$\sigma(WH)$	-	-	$\sim \kappa_W^2$		
$\sigma(q\bar{q} \rightarrow ZH)$	-	-	$\sim \kappa_Z^2$		
$\sigma(gg \rightarrow ZH)$	✓	$Z-t$	$\kappa_{ggZH}^2 \sim 2.27 \cdot \kappa_Z^2 + 0.37 \cdot \kappa_t^2 - 1.64 \cdot \kappa_Z \kappa_t$		
$\sigma(bbH)$	-	-	$\sim \kappa_b^2$		
$\sigma(ttH)$	-	-	$\sim \kappa_t^2$		
$\sigma(gb \rightarrow Wh)$	-	$W-t$	$\sim 1.84 \cdot \kappa_t^2 + 1.57 \cdot \kappa_W^2 - 2.41 \cdot \kappa_t \kappa_W$		
$\sigma(qb \rightarrow tHq')$	-	$W-t$	$\sim 3.4 \cdot \kappa_t^2 + 3.56 \cdot \kappa_W^2 - 5.96 \cdot \kappa_t \kappa_W$		
Partial decay width					
$\Gamma_{b\bar{b}}$	-	-	$\sim \kappa_b^2$		
Γ_{WW}	-	-	$\sim \kappa_W^2$		
Γ_{ZZ}	-	-	$\sim \kappa_Z^2$		
$\Gamma_{\tau\tau}$	-	-	$\sim \kappa_\tau^2$		
$\Gamma_{\mu\mu}$	-	-	$\sim \kappa_\mu^2$		
$\Gamma_{\gamma\gamma}$	✓	$W-t$	$\kappa_\gamma^2 \sim 1.59 \cdot \kappa_W^2 + 0.07 \cdot \kappa_t^2 - 0.66 \cdot \kappa_W \kappa_t$		
$\Gamma_{Z\gamma}$	✓	$W-t$	$\kappa_{Z\gamma}^2 \sim 1.12 \cdot \kappa_W^2 + 0.00035 \cdot \kappa_t^2 - 0.12 \cdot \kappa_W \kappa_t$		
Total decay width					
Γ_H	✓	$W-t$	$\kappa_H^2 \sim 0.57 \cdot \kappa_b^2 + 0.22 \cdot \kappa_W^2 + 0.09 \cdot \kappa_g^2 + 0.06 \cdot \kappa_t^2 + 0.03 \cdot \kappa_Z^2 + 0.03 \cdot \kappa_c^2 + 0.0023 \cdot \kappa_\gamma^2 + 0.0016 \cdot \kappa_{Z\gamma}^2 + 0.00022 \cdot \kappa_\mu^2$		
$b-t$					

HIGGS PRODUCTION MODES



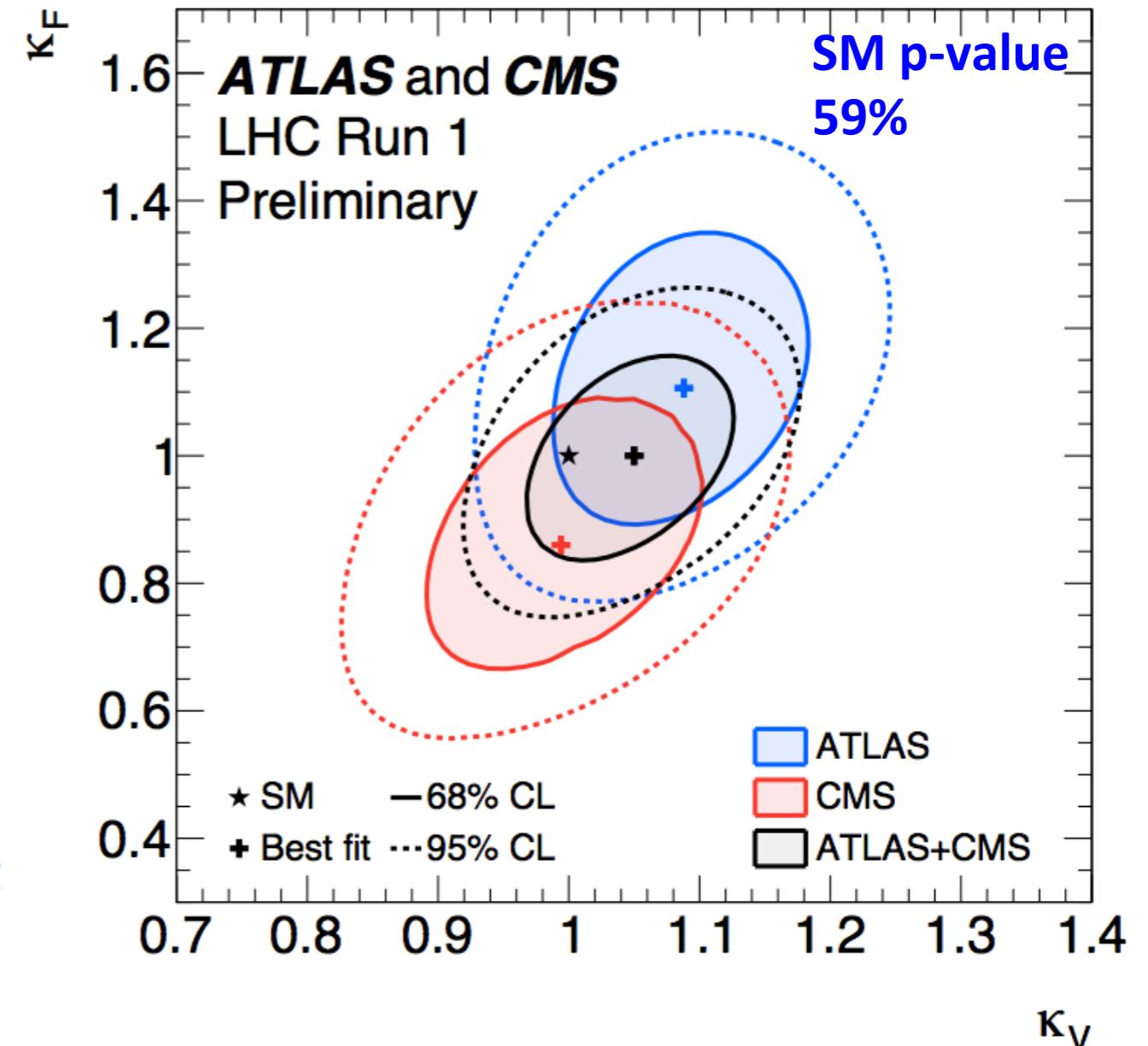
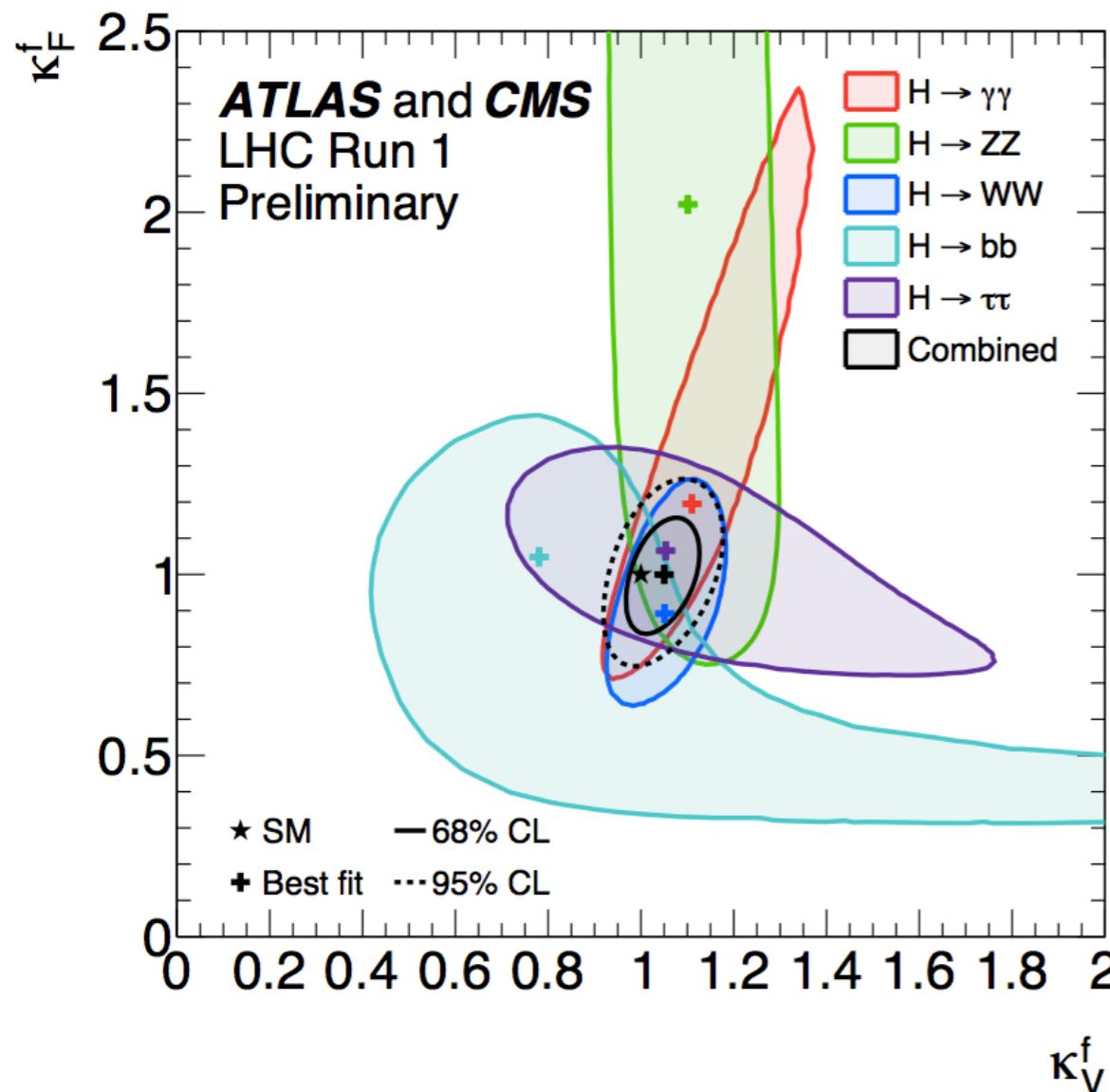
HIGGS DECAY CHANNELS

- Dominant: bb (57%) $\propto \kappa_b^2 / \kappa_H^2$
- WW channel (22%) $\propto \kappa_W^2 / \kappa_H^2$
- $\tau\tau$ channel (6.3%) $\propto \kappa_\tau^2 / \kappa_H^2$
- ZZ channel (3%) $\propto \kappa_Z^2 / \kappa_H^2$
- cc channel (3%) $\propto \kappa_c^2 / \kappa_H^2$
Extremely difficult
- The $\gamma\gamma$ channel (0.2%) $\propto \kappa_\gamma^2 / \kappa_H^2$

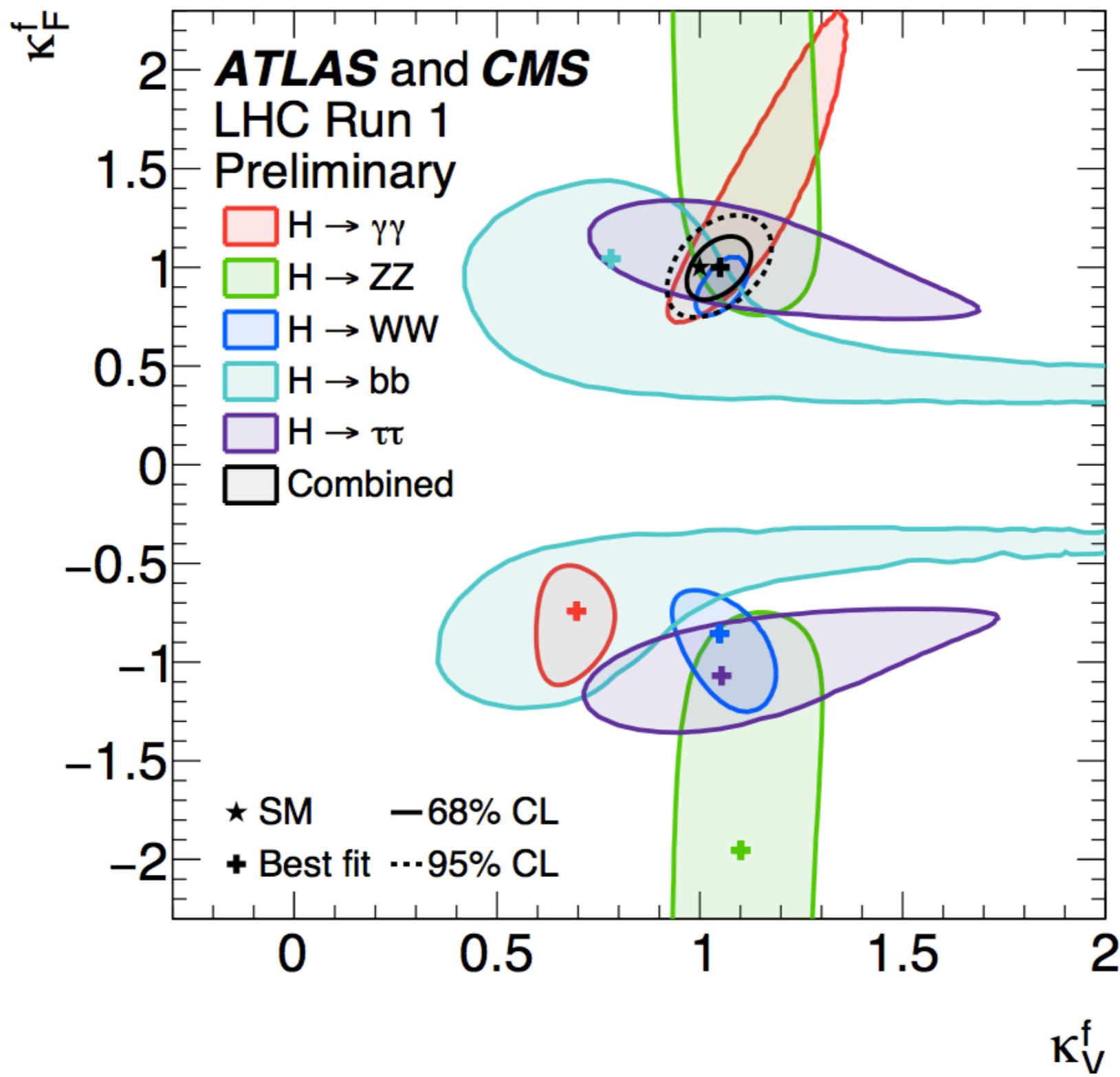


- The $Z\gamma$ (0.2%) $\kappa_{Z\gamma} \approx 1.12 \times \kappa_W^2 - 0.15 \times \kappa_t \kappa_W + 0.03 \times \kappa_t^2$
- The $\mu\mu$ channel (0.02%) $\propto \kappa_\mu^2 / \kappa_H^2$

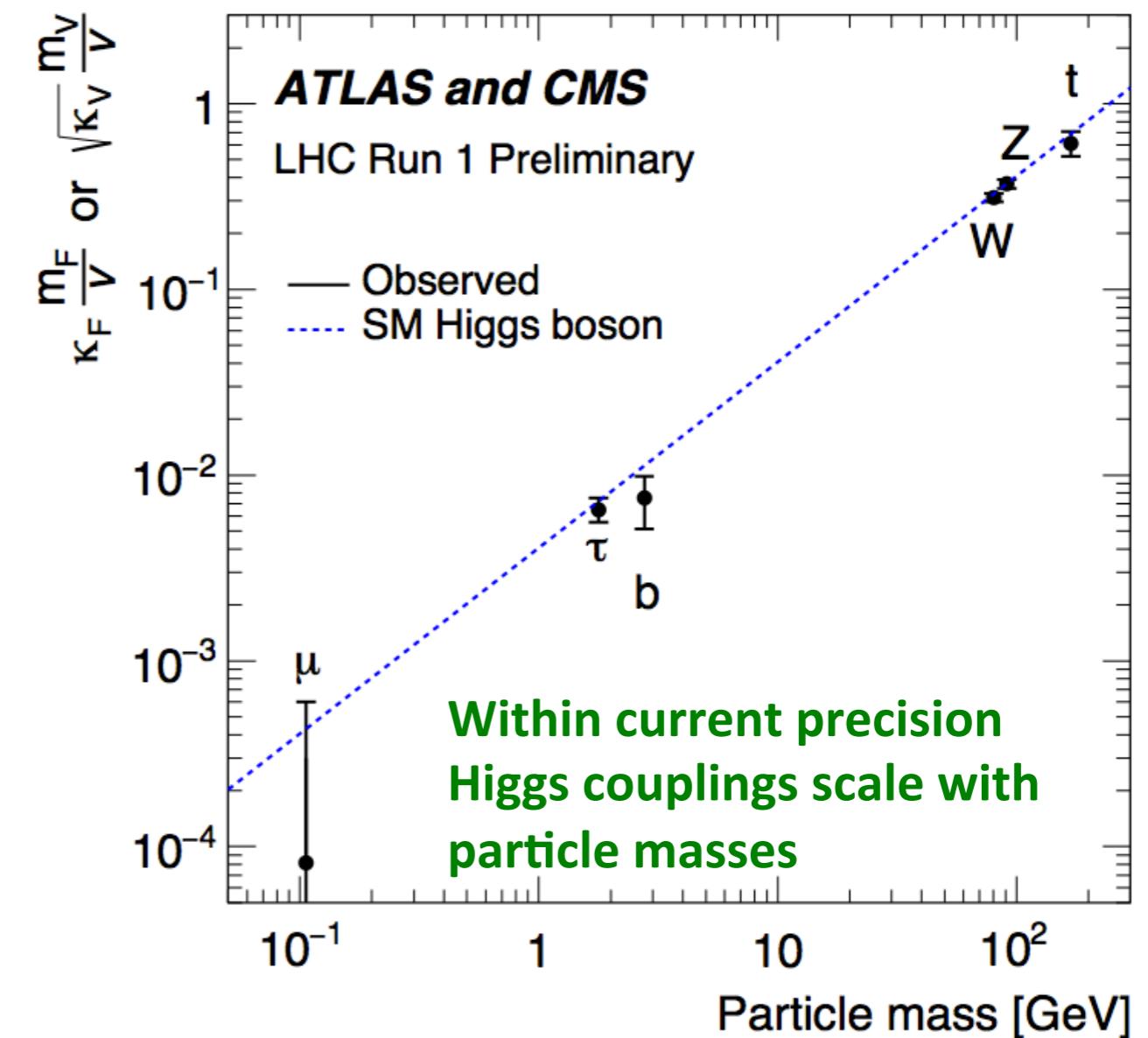
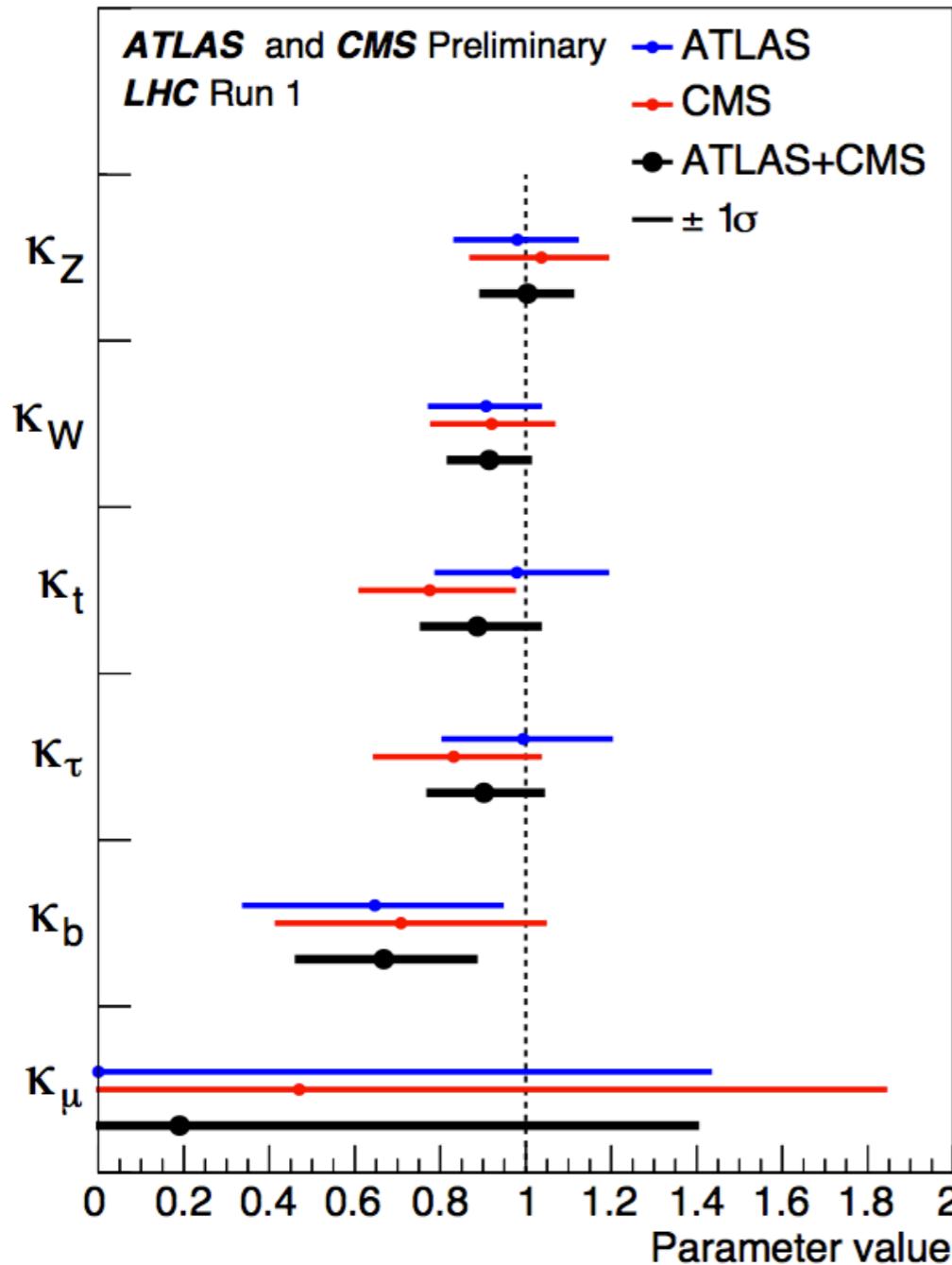
κ_F and κ_V Contours



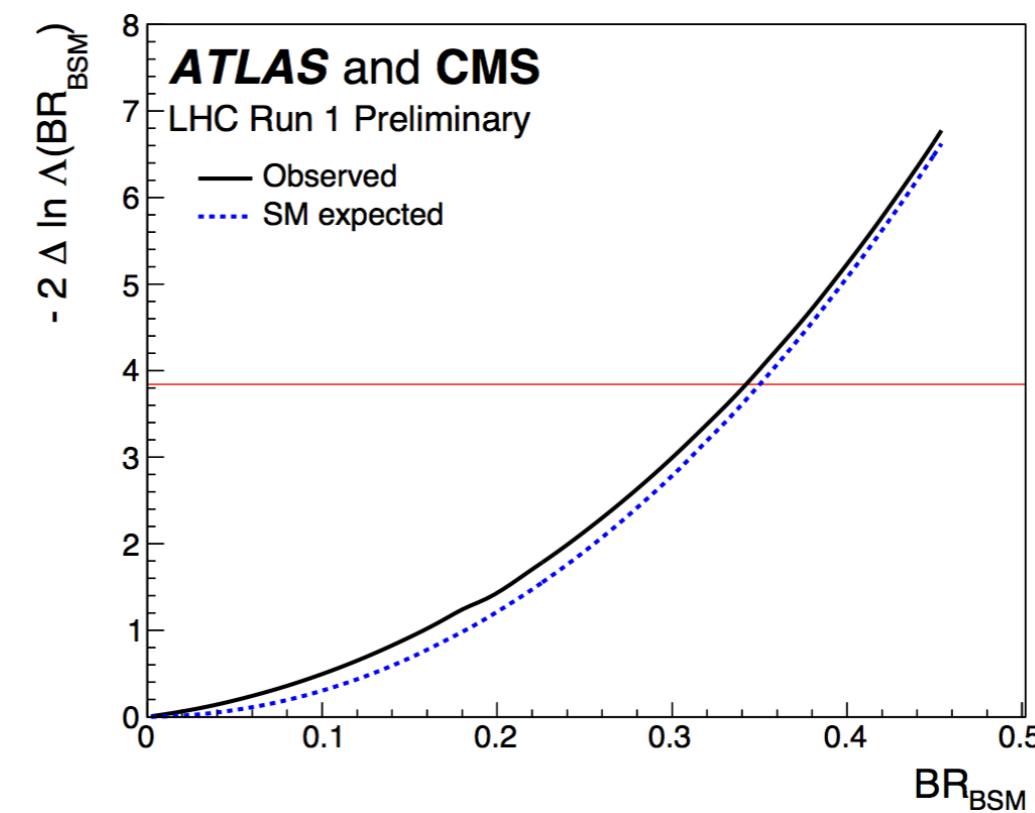
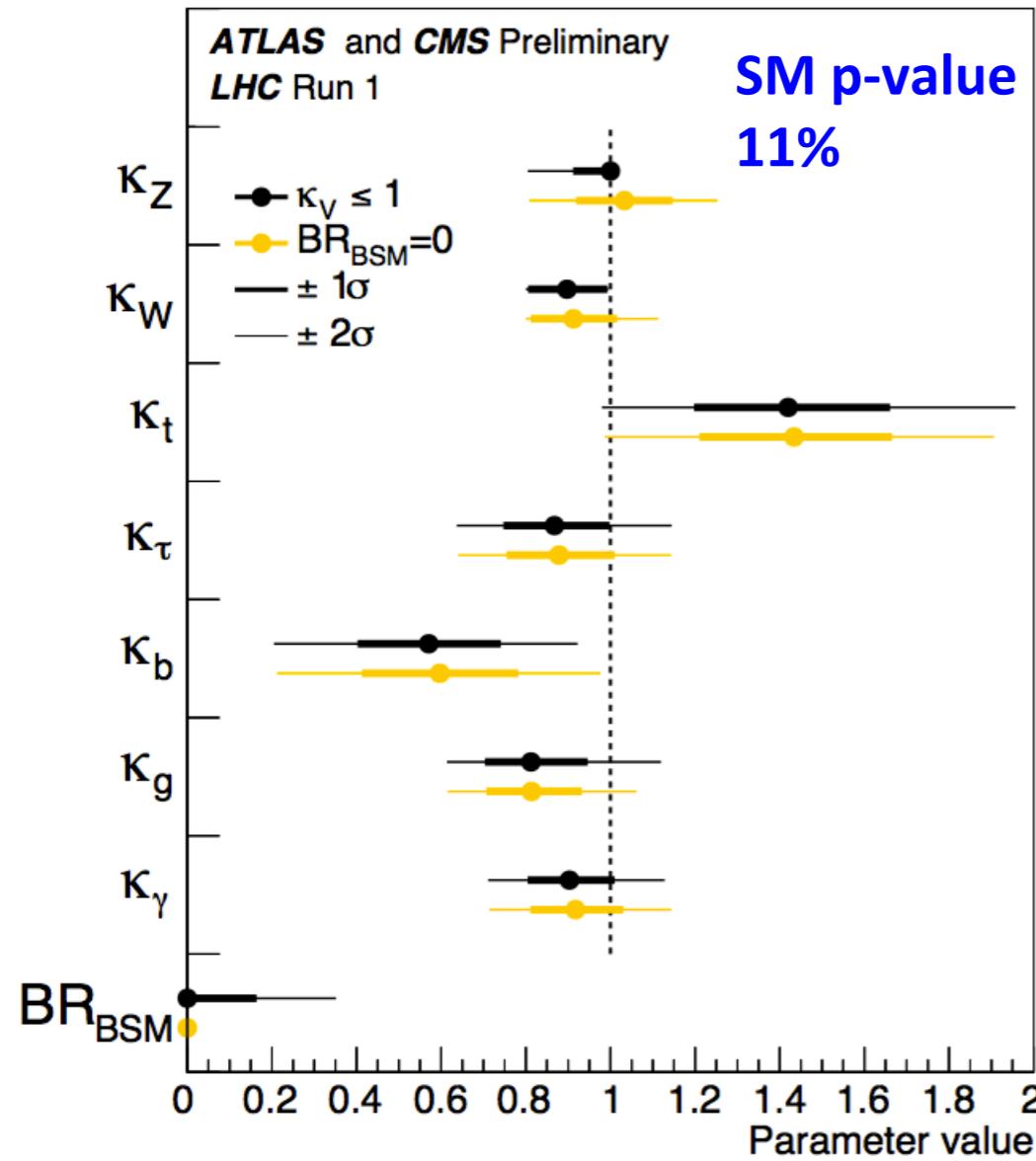
κ_F and κ_V Contours



COUPLING VS MASS

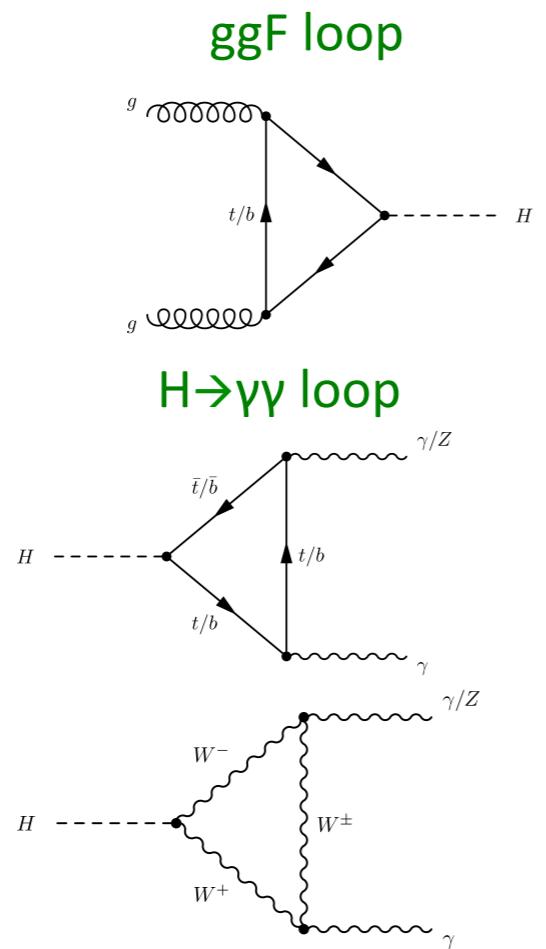


INVISIBLE HIGGS DECAYS

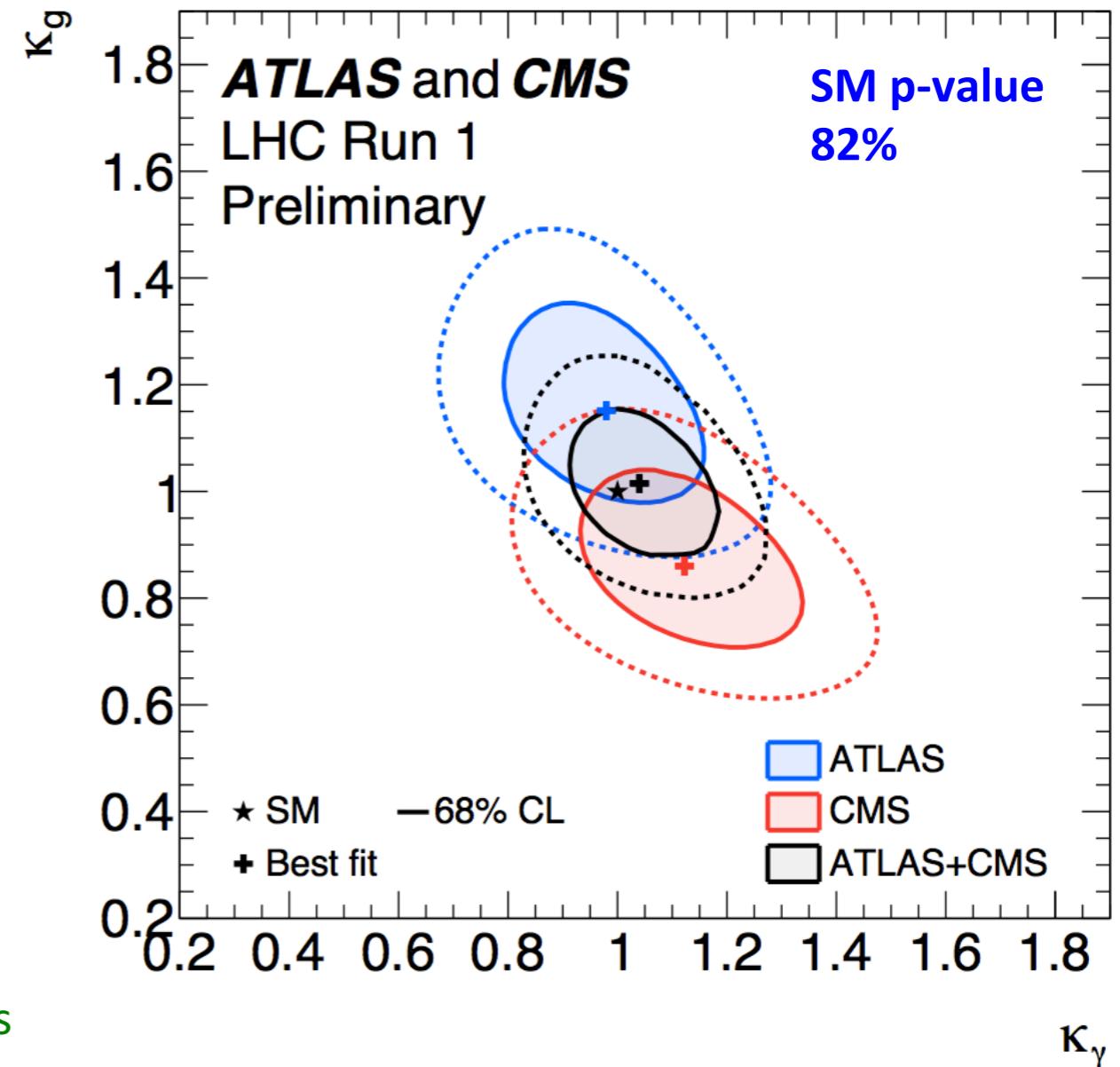


$BR_{BSM} < 0.34$ at 95% C.L. (assuming $\kappa_V \leq 1$)
 BR_{BSM} includes all possible non standard decays, visible or invisible

κ_g and κ_γ



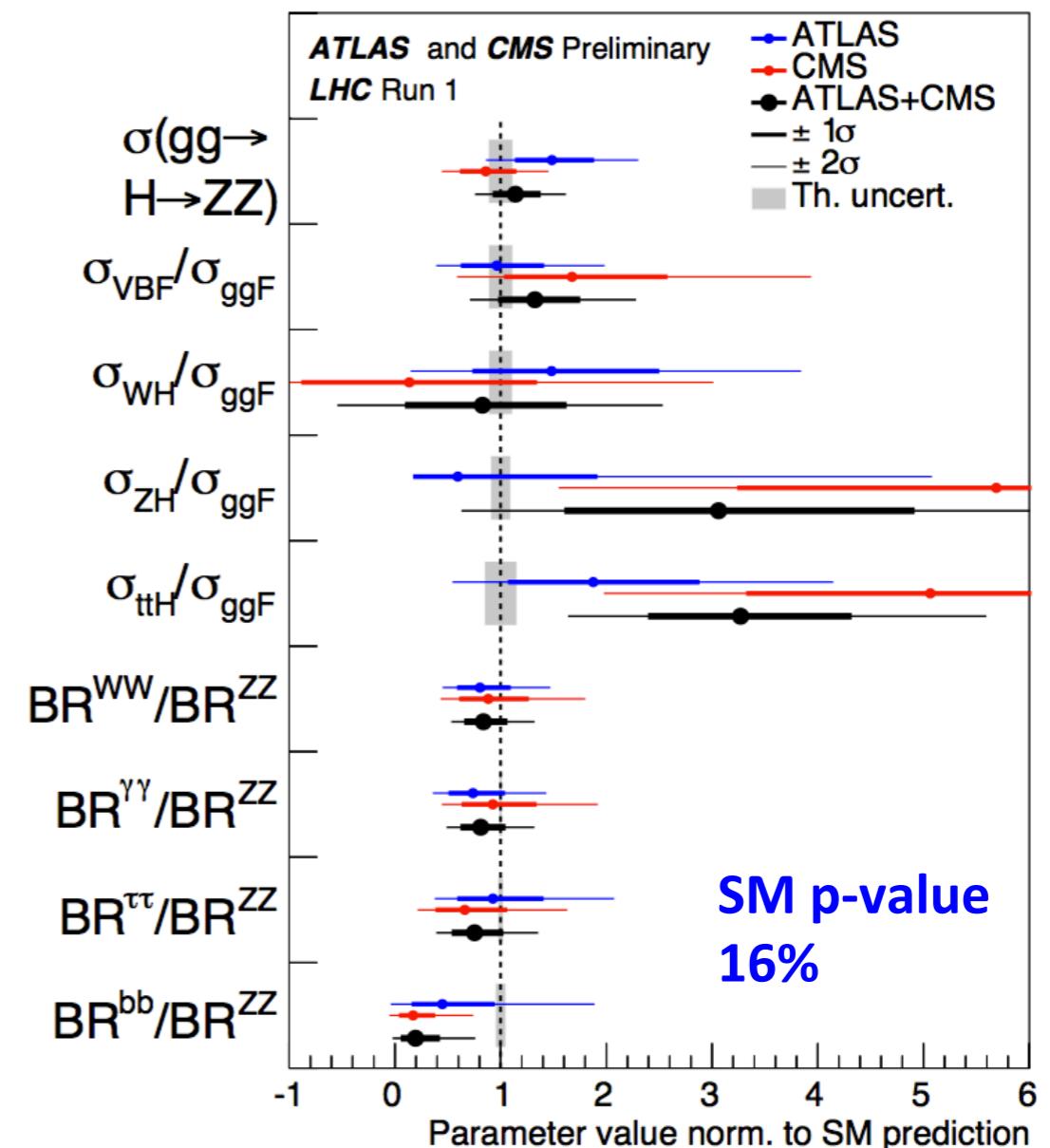
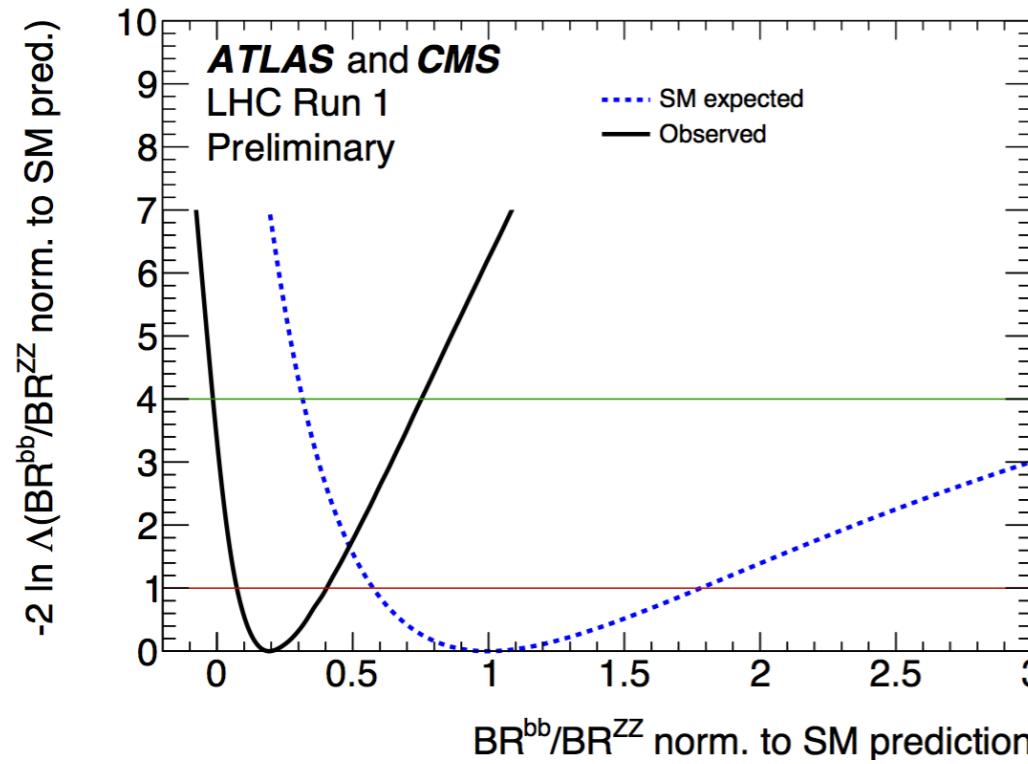
Additional heavy fermions or charged Higgs boson would modify the effective couplings



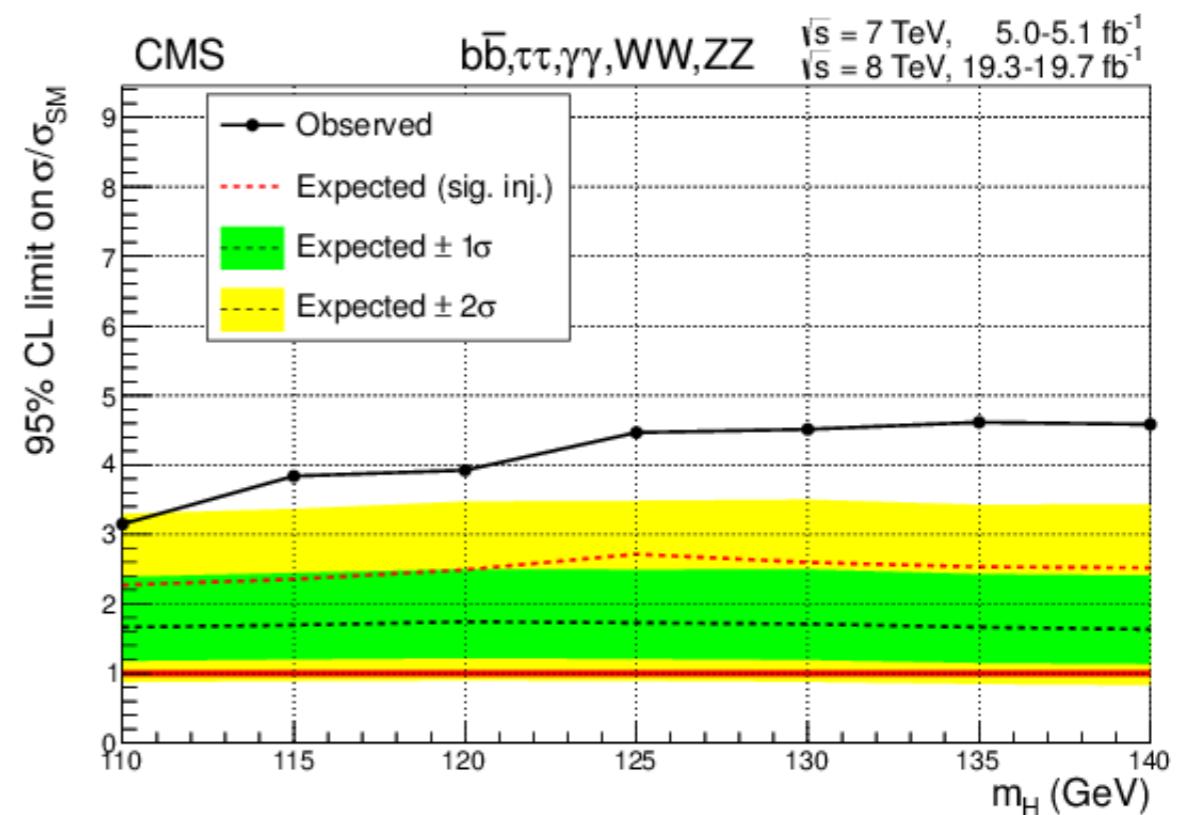
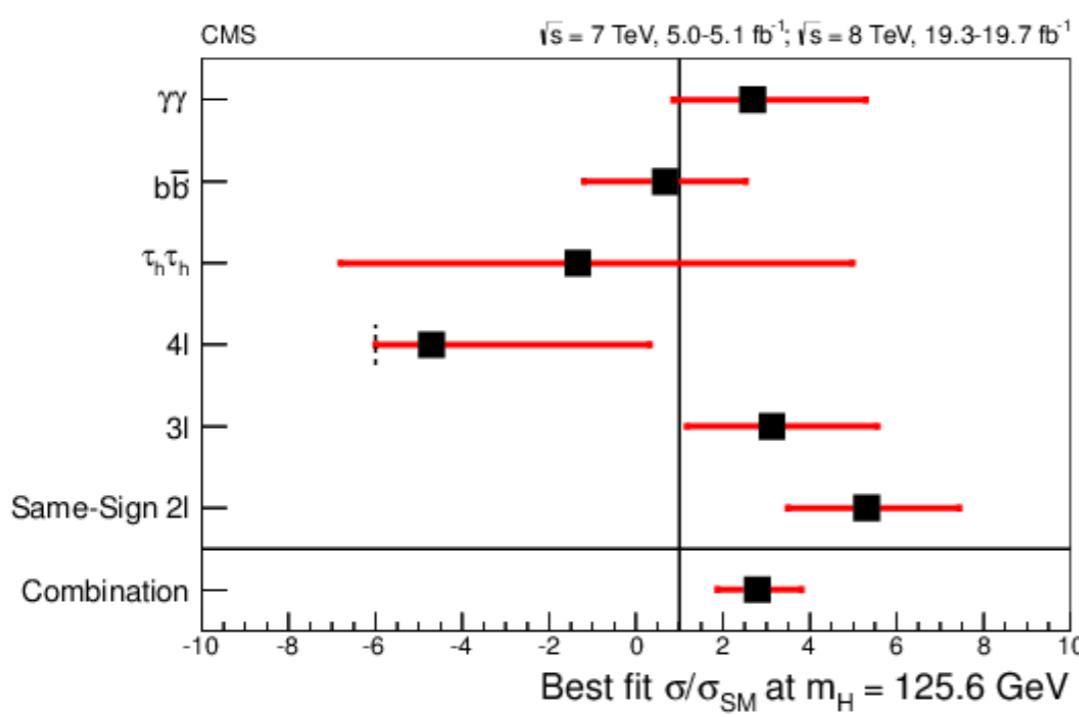
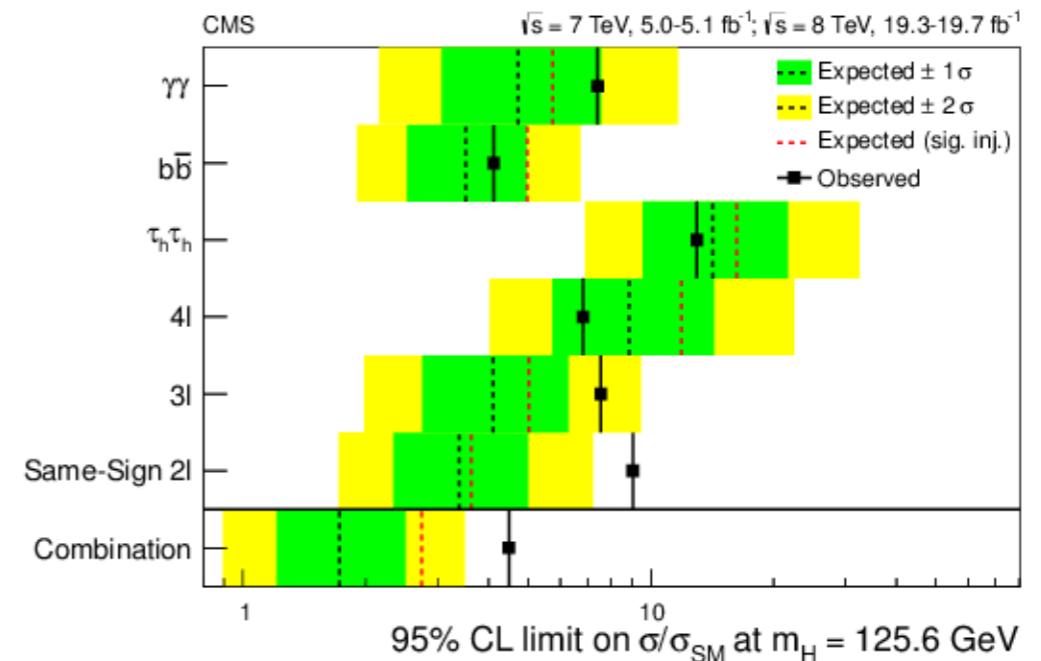
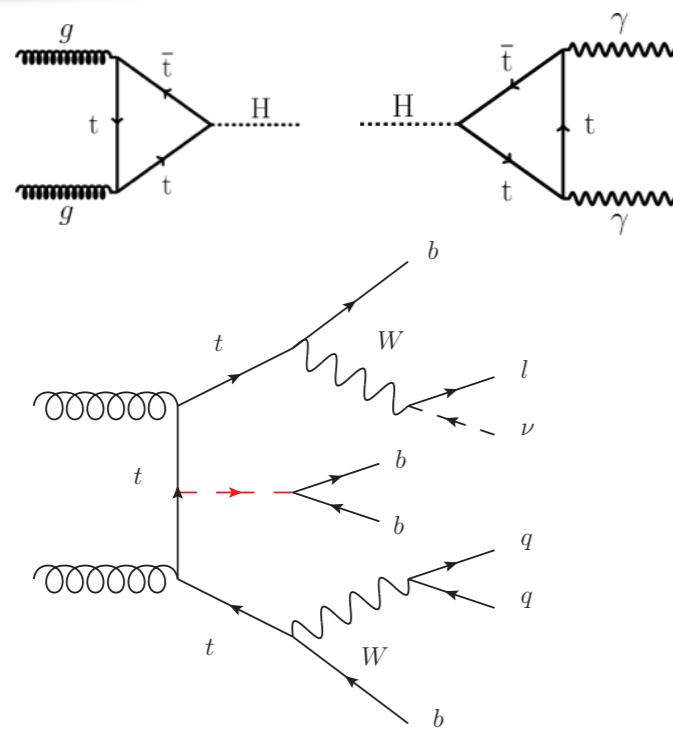
- Assume SM to be valid at tree level (no change to SM couplings)
- Consider modification only to main two loops
 - gluon-gluon fusion loop for production
 - top and W loop for Higgs decay to photons

GENERIC PARAMETERIZATION

- No attempt to disentangle production cross section from decay BR
- Measure everything relative to
 - gluon-gluon fusion: production mode with smallest systematic uncertainty
 - $H \rightarrow ZZ$ decay: cleanest experimental decay mode



tth



How MANY HIGGS ARE OUT THERE?

- Higgs doublet was introduced to cure missing mass term for vector bosons
- Yukawa coupling to fermions, again added intentionally, cured origin of fermion mass
- However no constraint on how many Higgs doublet must/can/should exist!

Two Higgs Doublets

- Two doublets with opposite hypercharge

$$H_d = \begin{pmatrix} H_d^1 \\ H_d^2 \end{pmatrix} = \begin{pmatrix} \Phi_1^0 {}^* \\ -\Phi_1^- \end{pmatrix}, \quad H_u = \begin{pmatrix} H_u^1 \\ H_u^2 \end{pmatrix} = \begin{pmatrix} \Phi_2^+ \\ \Phi_2^0 \end{pmatrix}$$

$$Y = -1 \qquad \qquad \qquad Y = +1$$

- 8 total degrees of freedom
- Different coupling to quarks for each doublet
 - H_u to up quarks
 - H_d to down quarks

$$\mathcal{L}_{\text{Yukawa}} = -h_u^{ij}(\bar{u}_R^i u_L^j H_u^2 - \bar{u}_R^i d_L^j H_u^1) - h_d^{ij}(\bar{d}_R^i d_L^j H_d^1 - \bar{d}_R^i u_L^j H_d^2) + \text{h.c.}$$

- Slightly more complicated Higgs potential

$$V = \left(m_d^2 + |\mu|^2 \right) H_d^{i*} H_d^i + \left(m_u^2 + |\mu|^2 \right) H_u^{i*} H_u^i - m_{ud}^2 \left(\epsilon^{ij} H_d^i H_u^j + \text{h.c.} \right)$$

$$+ \frac{1}{8} \left(g^2 + g'^2 \right) \left[H_d^{i*} H_d^i - H_u^{j*} H_u^j \right]^2 + \frac{1}{2} g^2 |H_d^{i*} H_u^i|^2,$$

$$\epsilon^{12} = -\epsilon^{21} = 1 \text{ and } \epsilon^{11} = \epsilon^{22} = 0$$

FIVE HIGGS BOSONS!

- After EW symmetry breaking, 3 scalars are absorbed by W and Z
- 5 remaining scalar degrees of freedom and two vacuum expectation value

$$\langle H_d \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} v_d \\ 0 \end{pmatrix}, \quad \langle H_u \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_u \end{pmatrix} \quad v^2 \equiv v_d^2 + v_u^2 = 4m_W^2/g^2 = (246 \text{ GeV})^2$$
$$\tan \beta \equiv \frac{v_u}{v_d}$$

- Combination of remaining fields results in 5 physical scalars

$$H^\pm = H_d^\pm \sin \beta + H_u^\pm \cos \beta$$

CP-odd $A^0 = \sqrt{2} \left(\text{Im } H_d^0 \sin \beta + \text{Im } H_u^0 \cos \beta \right)$

CP-even $h^0 = -(\sqrt{2} \text{Re } H_d^0 - v_d) \sin \alpha + (\sqrt{2} \text{Re } H_u^0 - v_u) \cos \alpha ,$

$$H^0 = (\sqrt{2} \text{Re } H_d^0 - v_d) \cos \alpha + (\sqrt{2} \text{Re } H_u^0 - v_u) \sin \alpha ,$$

- Mixing angle α due to diagonalization of original H_u and H_d fields to obtain physical mass

HIGGS MASS HIERARCHY

- 5 mass terms + 2 angles + coupling terms for each new Higgs but only two independent variables
 - constraints provided by supersymmetry (will be discussed in next lecture)
- Typically use m_A and $\tan\beta$ to parameterize other parameters

$$m_{H^\pm}^2 = m_A^2 + m_W^2$$

$$m_{H,h}^2 = \frac{1}{2} \left(m_A^2 + m_Z^2 \pm \sqrt{(m_A^2 + m_Z^2)^2 - 4m_Z^2 m_A^2 \cos^2 2\beta} \right)$$

$$m_h \leq m_Z |\cos 2\beta| \leq m_Z$$

- For $m_A \gg m_Z$ h decouples from remaining for Higgs bosons!

$$m_h^2 \simeq m_Z^2 \cos^2 2\beta ,$$

$$m_A \simeq m_H \simeq m_{H^\pm}$$

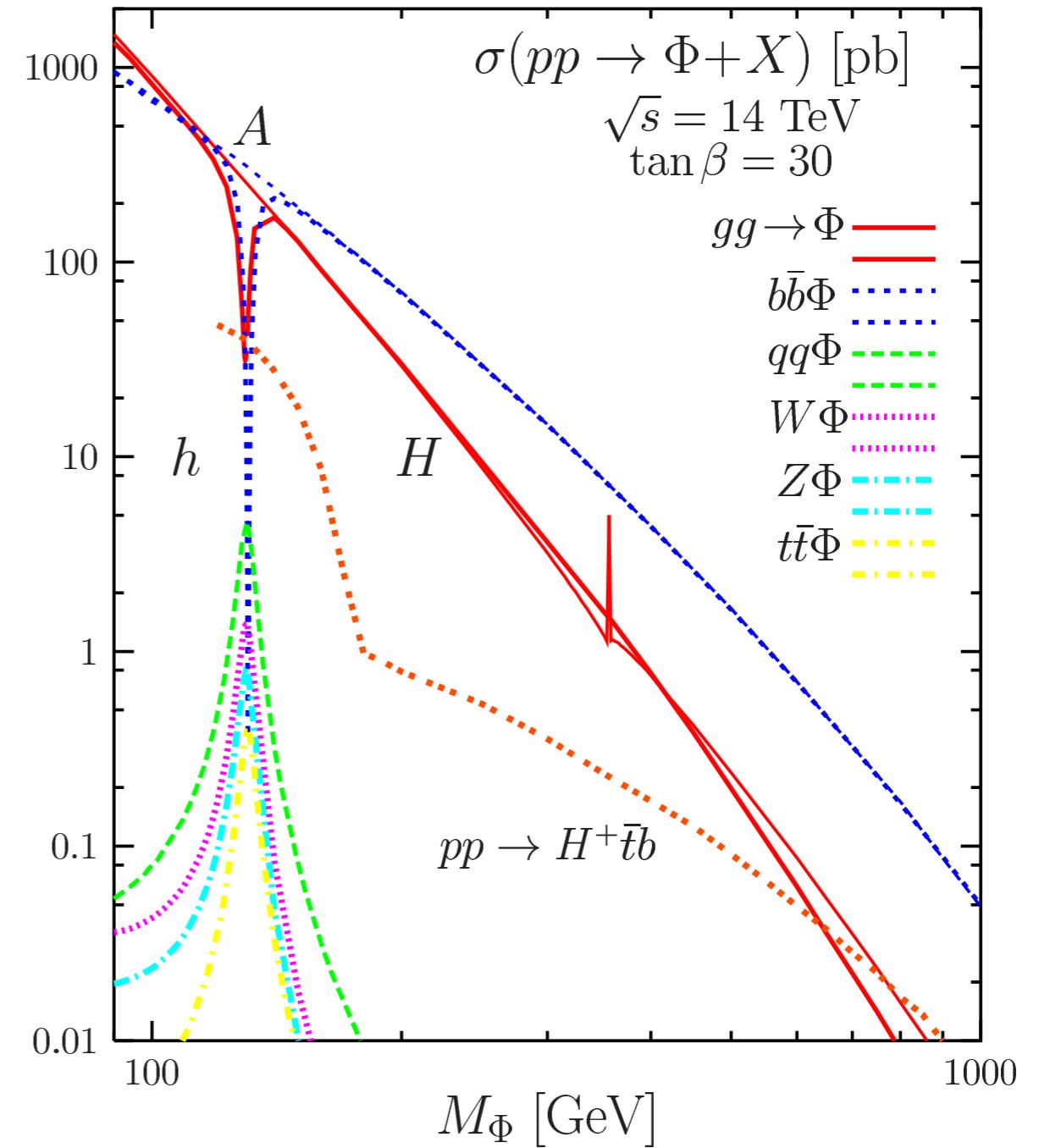
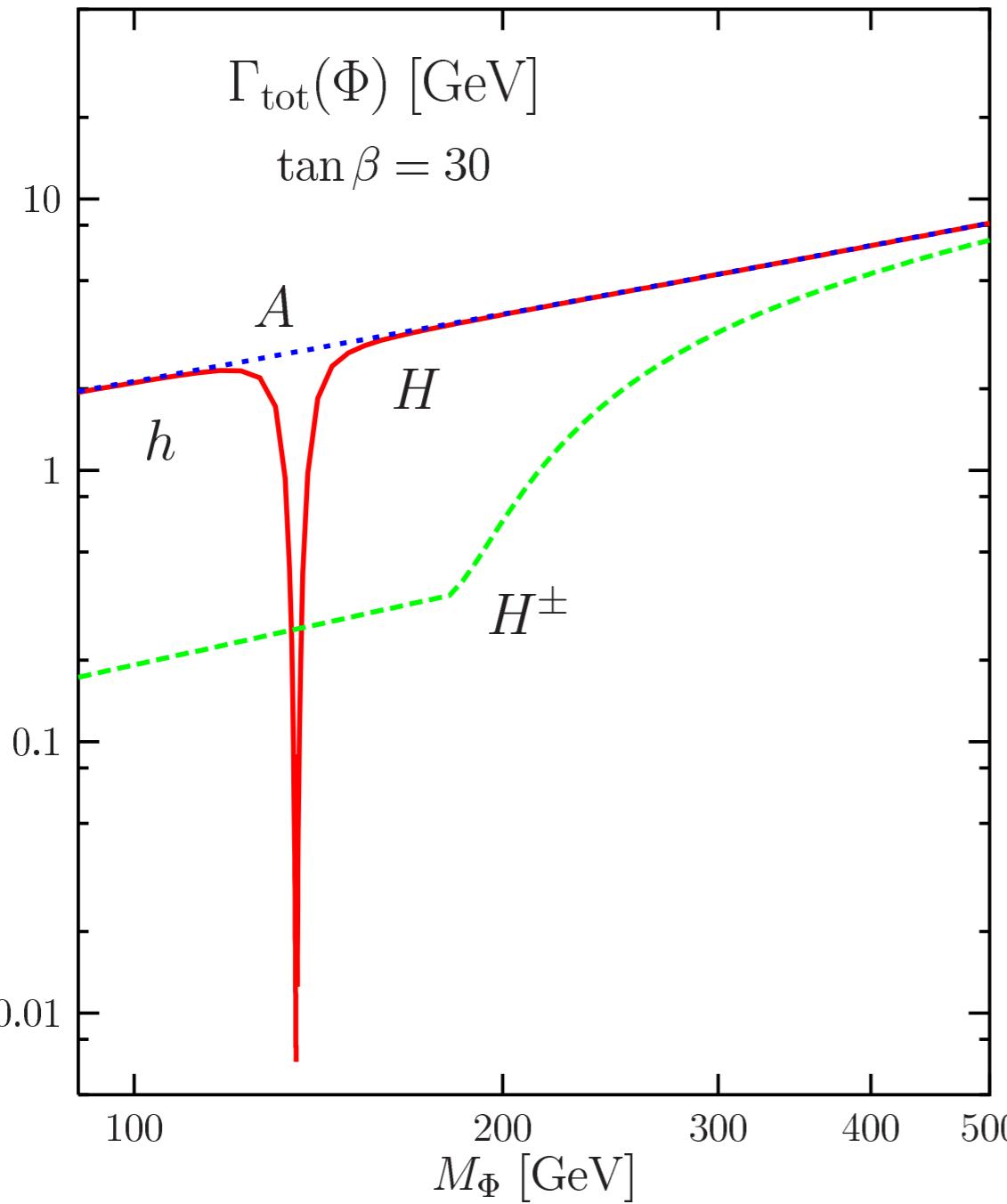
$$m_H^2 \simeq m_A^2 + m_Z^2 \sin^2 2\beta$$

$$m_{H^\pm}^2 = m_A^2 + m_W^2 ,$$

- Lowest Higgs h will behave like SM Higgs
 - similar coupling to fermions and vector bosons

$$\cos^2(\beta - \alpha) \simeq \frac{m_Z^4 \sin^2 4\beta}{4m_A^4} \sim 0$$

HIGGS PROPERTIES



- Enhanced cross section for large $\tan \beta$ compared to Standard Model