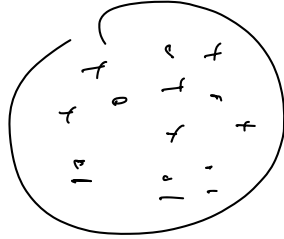


Elastic Scattering and Hadron Structure

Atomic structure \rightarrow Nucleus.

Atomic scale $1^{\circ} \text{A}^{\circ} \quad 10^{-10} \text{m}$.

Thompson theory



Rutherford proved: + charge confined in space: nucleus.

Size of nucleus $1 \text{ fm} \quad 10^{-15} \text{m}$.

Scattering of $\alpha + \text{Au} \rightarrow \alpha + \text{Au}$.

1910

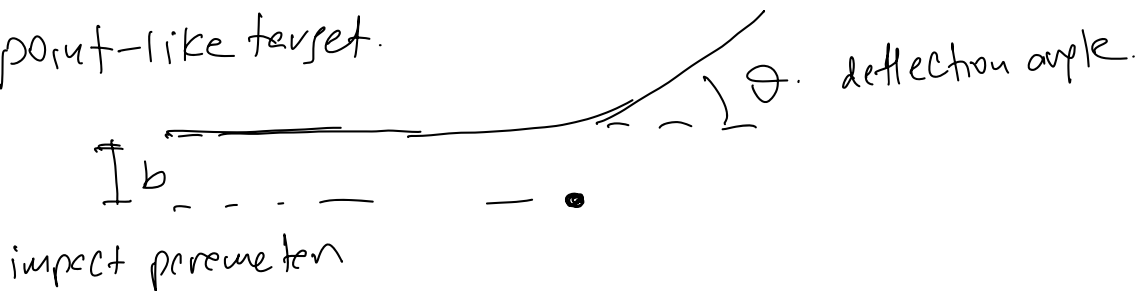
1) probe: α particles.

2) heavy target: $\text{Au} \Rightarrow$ neglect target recoil.

pure elastic scattering.

α particles have $K = E - m = 5-7 \text{ MeV}$. from α decay

3) point-like target.



a) measure $\frac{d\sigma}{d\Omega}(\theta)$ differential cross section.

reactions vs. θ .

Rutherford: ignore target recoil.

α : mass 3.7 GeV . $K = 7 \text{ MeV}$ highly non relativistic

e^- as probe: 1) ultra relativistic already @ 10 MeV .
2) spin to be accounted (helicity)

Experimental method:

- 1) predict $d\sigma/d\Omega$.
- 2) measure $d\sigma/d\Omega$ vs θ .
- 3) look for deviations.

probe + T \rightarrow probe + T elastic scattering.

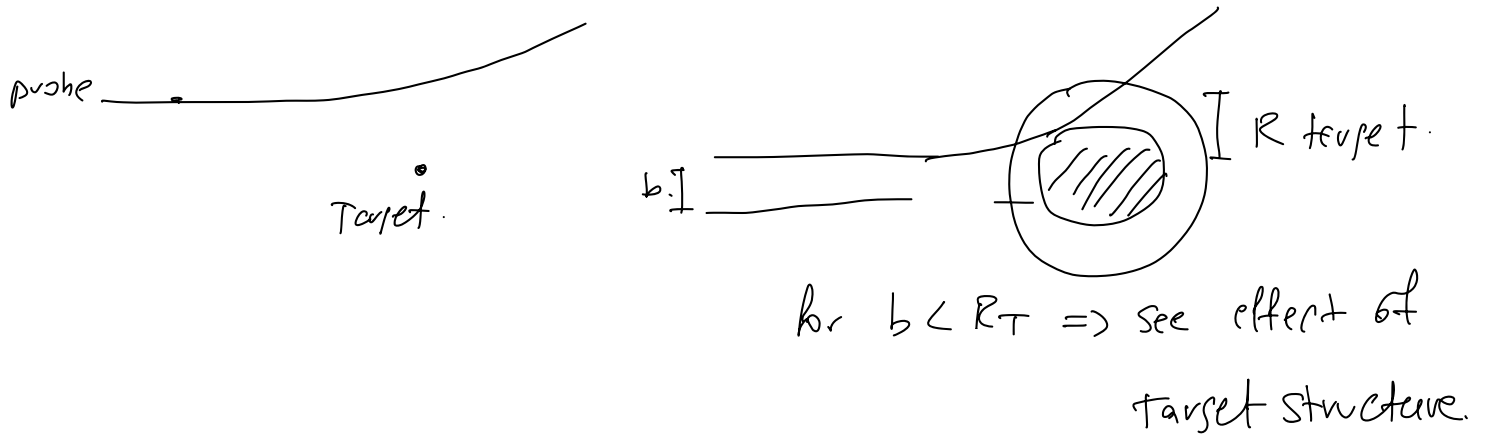
\rightarrow probe + T^{*} excitation.

\rightarrow probe + new particles. inelastic scattering.
 $\neq T$

if $\left. \frac{d\sigma}{d\Omega} \right|_{\text{pred.}} \neq \left. \frac{d\sigma}{d\Omega} \right|_{\text{meas.}}$

- \rightarrow calculation incorrect
- \rightarrow structure in the target
- \rightarrow new interactions.

classical Reth.

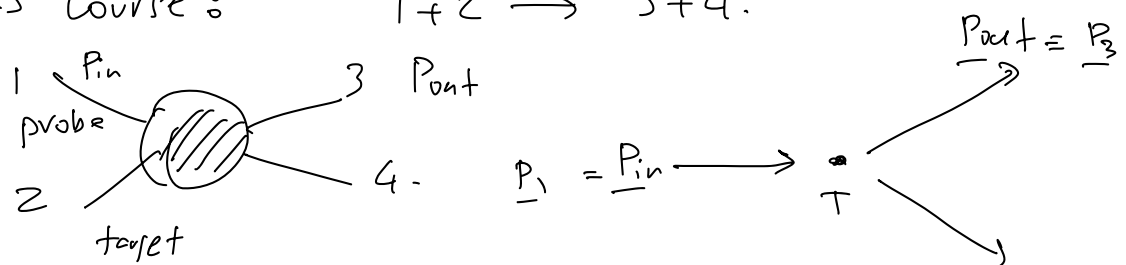


Experimentally: measure

- 1) # events $\frac{d\sigma}{d\Omega}$.
- 2) angle θ of deviation.
- 3) E1 energy of the probe after scattering

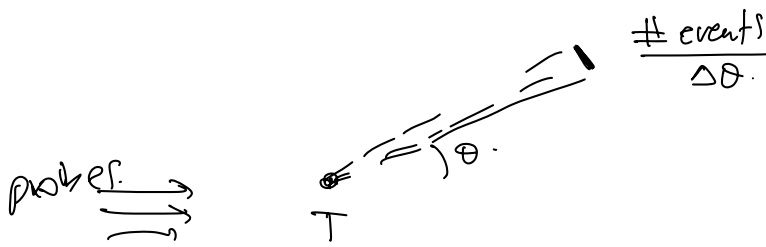
Lecture 2 of this course:

1 + 2 \rightarrow 3 + 4.

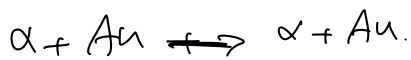


$$\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2} \frac{1}{(M_T)^2} \frac{|\vec{P}_{out}|}{|\vec{P}_{in}|} |M|^2$$

General Mott Formula
for 2 body scattering
in LAB frame.



Rutherford Experiment

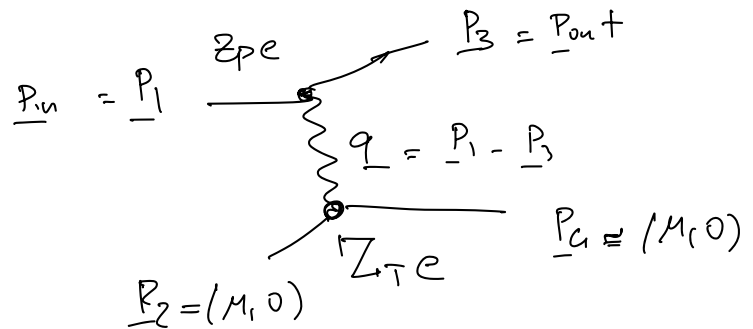


$$\mu \sim Z_P Z_T \frac{1}{q^2}$$

$$= \frac{Z_P Z_T \alpha}{q^2}$$

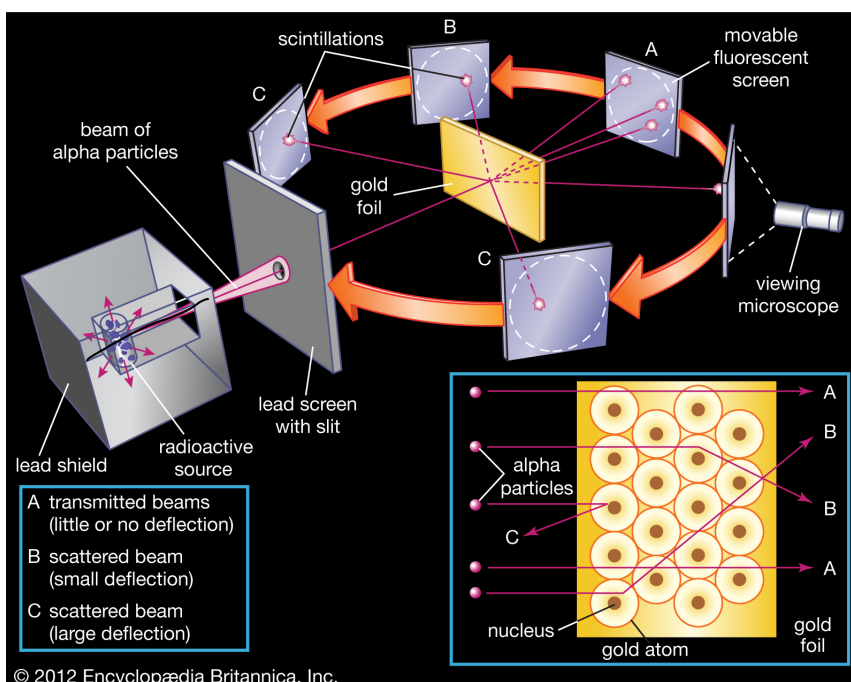
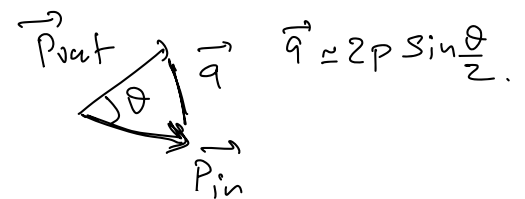
$$\Rightarrow \frac{d\sigma}{d\Omega} \sim (Z_P Z_T \alpha)^2 \frac{1}{q^4}$$

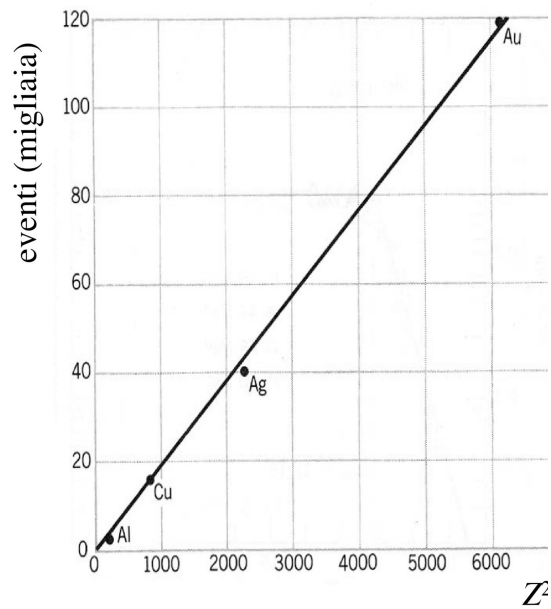
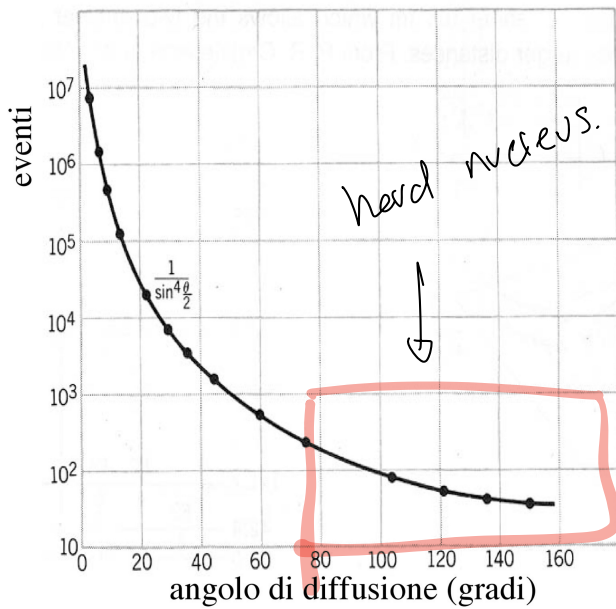
$$\frac{d\sigma}{d\Omega} \sim \alpha^2 (Z_P Z_T)^2 \frac{1}{p^4 \sin^4 \frac{\theta}{2}}$$



negligible recoil.

$$\Rightarrow |\vec{P}_1| = |\vec{P}_3| = p.$$

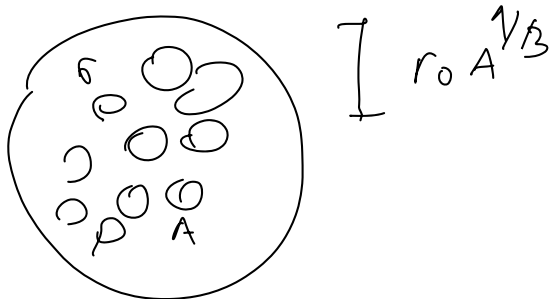
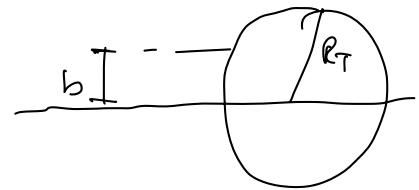




Rutherford: $R_N < 40 \text{ fm}$.

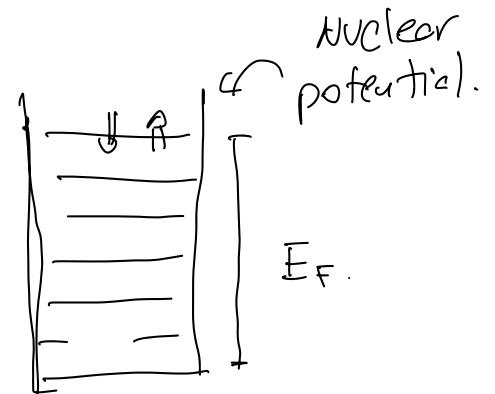
today we know
 $R_N \approx r_0 A^{1/3}$
 $r_0 \approx 1.2 \text{ fm}$.

Could not prob structure of nucleus:



Assume Fermi gas of nucleons:

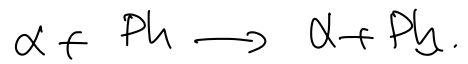
Fermi energy: max kin. energy of nucleons in the nucleus.



$$E_F \approx \frac{P_F^2}{2m_N} \approx 30 \text{ MeV}. \quad P_F \approx 250 \text{ MeV}.$$

To see the structure \Rightarrow accelerate α up to 20-30 MeV.
 around 1950's.

Rev. Mod. Phys 33, 190 (1961)



$$\theta = 60^\circ$$

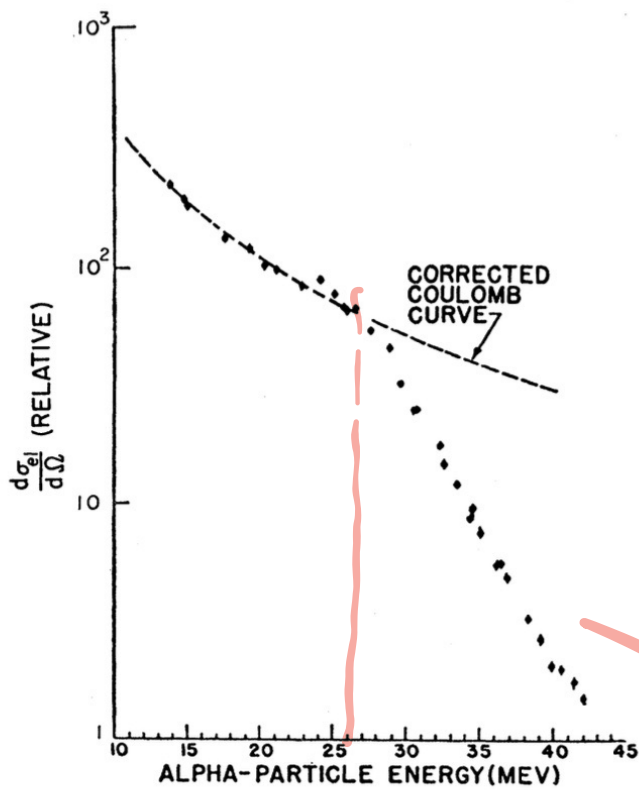
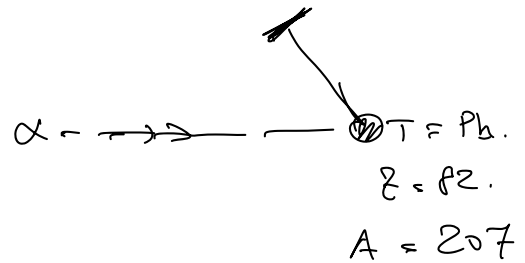


FIG. 5. Differential cross section for the elastic scattering of alpha particles by Pb at 60° as a function of the alpha-particle energy.



Exponential deviation from Coulomb scatt.

For $E > 25 \text{ MeV} \Rightarrow$ due e^- which are nonrelativistic.

Study $e^- + N \longrightarrow e^- + N$ elastic
 $e^- + N^*$ excitation.
 $e^- + H$ inelastic.

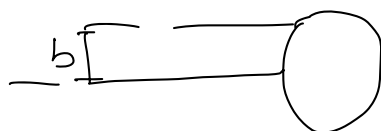
1910 α probes.

1960 e^- probes.

>1960 - 2000 e^-/ν probes.

QM: De Broglie wave length of particle of mass M ?

$$\Delta p \Delta x \approx 1 \Rightarrow \Delta p \approx \frac{1}{b}$$

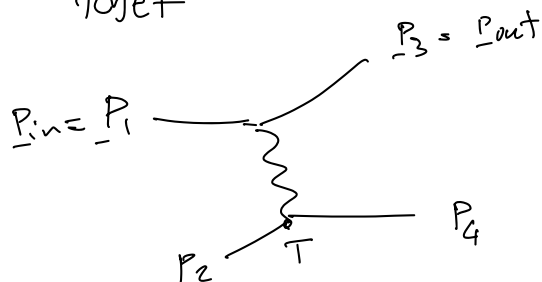
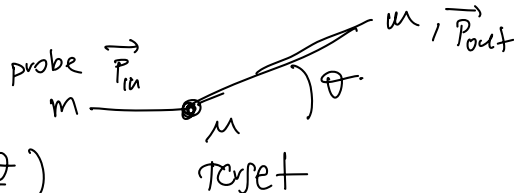


$$200 \text{ fm} \cdot \text{MeV} = \hbar c = 1.$$

e^- : spin 1/2 relativistic

$$\frac{d\sigma}{d\Omega} \sim \alpha^2 \frac{1}{q^4} (m^2 + P_{in}^2 \cos^2 \frac{\theta}{2})$$

Rutherford.
Mott Formula.



$$m^2 + P^2 \cos^2 \frac{\theta}{2} = m^2 + P^2 (1 - \sin^2 \frac{\theta}{2})$$

$$= \underbrace{m^2 + P^2} - P^2 \sin^2 \frac{\theta}{2} = E^2 - P^2 \sin^2 \frac{\theta}{2} = E^2 (1 - (\frac{P}{E})^2 \sin^2 \frac{\theta}{2})$$

$$\beta = \frac{P}{E}$$

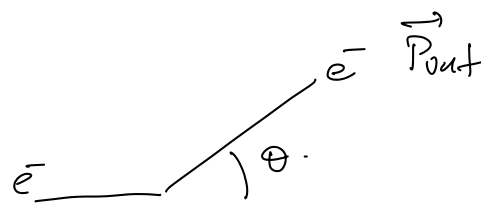
$$E^2 (1 - \beta^2 \sin^2 \frac{\theta}{2}).$$

$$\frac{d\sigma}{d\Omega} \Big|_{\text{Mott}} = \frac{d\sigma}{d\Omega} \Big|_{\text{Ruth.}} \underbrace{(1 - \beta^2 \sin^2 \frac{\theta}{2})}_{\text{because of } e^- \text{ spin.}}$$

Mott
with no recoil
in target.

suppose ultrarel. $e^- \Rightarrow \beta = 1$.

$$\frac{d\sigma}{d\Omega} \Big|_{\text{Mott}} = \frac{d\sigma}{d\Omega} \Big|_{\text{Ruth.}} \cos^2 \frac{\theta}{2}$$



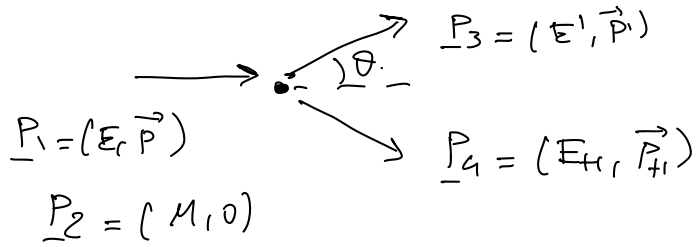
$$\theta = \pi \Rightarrow e^- \Rightarrow \bullet \quad \text{IN.}$$

$$e^- \Leftarrow \bullet \quad \text{FIN.}$$

$\beta = 1 \Rightarrow$ helicity conserved. $\Rightarrow \Delta J = 1$. not possible.

NB: Mott with zero recoil of target.

Study $e^- + N \rightarrow e^- + N/H$



Elastic: $e^- + N \rightarrow e^- + N$.

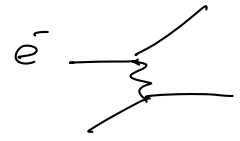
$$\underline{P}_4 = (M, \vec{P}_4)$$

Inelastic β .

hadronic system.

$$\left\{ \begin{array}{l} E_H = \sum_i^N E_i \\ \vec{P}_H = \sum_i^N \vec{P}_i \end{array} \right.$$

$$W = M_H = \left(\sum_i \underline{P}_i \right)^2$$



Example: $e^- + p \rightarrow e^- + p$.



Elastic case

$$\underline{P}_1 + \underline{P}_2 = \underline{P}_3 + \underline{P}_4$$

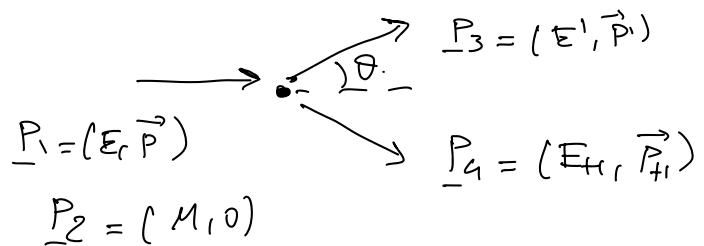
$$\underline{q} = \underline{P}_1 - \underline{P}_3$$

4-vector of mom.

transferred. to target.

$$\underline{P}_4 = \underline{P}_1 - \underline{P}_3 + \underline{P}_2$$

$$P_4^2 = M^2$$



Observable quantities:

$$\theta, E'$$

$$(\underline{P}_1 - \underline{P}_3 + \underline{P}_2)^2 = M_e^2 + M_e^2 + M^2 + \underbrace{2 \underline{P}_1 \cdot \underline{P}_2}_{2EM} - 2 \underline{P}_1 \cdot \underline{P}_3 - \underbrace{2 \underline{P}_3 \cdot \underline{P}_2}_{2E'M}$$

$$M^2 = 2m_e^2 + M^2 + 2EM - 2E'M - 2(EE' - \vec{P} \cdot \vec{P}')$$

$E \gg m_e$ relativistic regime.

$$E' \gg m_e, \quad |\vec{P}| \approx E, \quad |\vec{P}'| \approx E'$$

$$2m_e^2 \approx 0.$$

$$0 = EM - E'M - (EE' - EE' \cos \theta).$$

$$E' = \frac{E}{1 + \frac{E}{M}(1 - \cos \theta)}$$

$$1 - \cos \theta \approx 2 \sin^2 \frac{\theta}{2}.$$

Elastic scattering: E', θ are correlated.
not independent.

