

Isospin

$$p = |1/2, 1/2\rangle \quad |n\rangle = |1/2, -1/2\rangle$$

$$pn, \text{ deuteron} = |0, 0\rangle?$$

$$\begin{pmatrix} p \\ n \end{pmatrix} \text{ doublet } I=1/2$$

$$|\pi^{\pm, 0}\rangle = \text{triplet } I=1$$

$$p+p \rightarrow p+p \text{ elastic}$$

$$p \rightarrow \boxed{\text{triplet}} \text{ inelastic processes}$$

① $p+p \rightarrow d+\pi^+$ a)

$$Q \quad 1 \quad 1 \quad 1 \quad 1$$

$$B \quad 1 \quad 1 \quad 2 \quad 0$$

$$I_3 \quad +1/2 \quad +1/2 \quad 0 \quad +1$$

$$0 \oplus 1 = 1 \quad |I, I_3\rangle$$

initial
 $|I, I_3\rangle = |1, +1\rangle$

final state: $|1, 1\rangle$

strong interaction conserves Isospin.

$$\langle f | H_I | i \rangle$$

$$\langle 1, 1 | H_I | 1, 1 \rangle$$

② $p+n \rightarrow d+\pi^0$

$$Q \quad 1 \quad 0 \quad 1 \quad 0$$

$$B \quad 1 \quad 1 \quad 2 \quad 0$$

$$I_3 \quad +1/2 \quad -1/2 \quad 0 \quad 0$$

initial state

$$I_3 = 0$$

$p+n$

$$\frac{1}{2} \otimes \frac{1}{2} = 0 \oplus 1$$

initial state.

$$|p+n\rangle = \frac{1}{\sqrt{2}}(|1, 0\rangle + |0, 0\rangle)$$

final state: $I_3 = 0 \quad d+\pi^0 \quad |0, 0\rangle \times |1, 0\rangle$
 $|1, 0\rangle$

$$|i\rangle = \frac{1}{\sqrt{2}}(|1, 0\rangle + |0, 0\rangle) \xrightarrow{H_I} |f\rangle = |1, 0\rangle.$$

$$\sigma_a \propto |M_a|^2 \rho(p+p \rightarrow d+\pi^+)$$

$$\sigma_b \propto |M_b|^2 \rho(p+n \rightarrow d+\pi^0)$$

$$\# \text{acc} = \sigma \times \frac{dn_p}{dt} \times d \times n_B$$

$$\frac{\sigma_a}{\sigma_b} = \frac{\# p+p \rightarrow d+\pi^+}{\# p+n \rightarrow d+\pi^0} = \frac{|M_a|^2}{|M_b|^2} \frac{\rho(d\pi^+)}{\rho(d\pi^0)}$$

$$\frac{\Delta m_{pn}}{m_p} \approx 2 \text{ MeV} / 938 \text{ MeV}$$

$$\frac{\Delta m_{\pi}}{m_{\pi}} \approx 5 \text{ MeV} / 135 \text{ MeV}$$

$$\Rightarrow \frac{\rho(\pi^+d)}{\rho(\pi^0d)} \approx 1$$

$$M_a = \langle 1,1 | H | 1,1 \rangle \quad p+p \rightarrow d+\pi^+$$

$$M_b = \langle 1,0 | H | \frac{1}{\sqrt{2}}(1,1,0) + 1,0,0 \rangle$$

$$p+n \rightarrow d+\pi^0$$

$$H_{\text{strong}} \text{ conserves isospin. } \Rightarrow M_b = \frac{1}{\sqrt{2}} \langle 1,0 | H | 1,1,0 \rangle + \frac{1}{\sqrt{2}} \langle 1,0 | H | 1,0,0 \rangle$$

$$\Rightarrow \frac{\sigma_a}{\sigma_b} \approx \frac{1}{|\frac{1}{\sqrt{2}}|^2} \approx 2 \quad \text{confirmed by } \frac{\# \pi^+d}{\# \pi^0d}$$

$$\text{Another experiment } \pi^+ + (p/n) \rightarrow$$

$$I \quad I_3$$

$$I \quad I_3$$

$$\pi^+ + p \rightarrow \pi^+ + p$$

$$(a)$$

$$3/2 \quad 3/2$$

$$|3/2, 3/2\rangle \quad (1/2, 1/2)$$

$$\pi^+ + p \rightarrow \pi^+ + n$$

$$\begin{array}{cc|cc} B & 0 & 1 & 0 & 1 \\ Q & +1 & +1 & +2 & 0 \\ I_3 & & & & \end{array}$$

$$\times$$

$$\text{No } \pi^+n$$

$$\pi^+ + p \rightarrow n + \pi^+ \pi^+$$

$$I_3 \quad +1 \quad 1/2 \quad -1/2 \quad +1 \quad +1$$

$$M = \langle f | H | i \rangle$$

$$= \langle n\pi\pi | H_S | p\pi^+ \rangle$$

$$\pi^+ + n \longrightarrow \pi^+ + n \quad (d)$$

$$\pi^+ + n \longrightarrow \pi^0 + p \quad (g)$$

$$\pi^- + p \longrightarrow \pi^- + p \quad (c)$$

$$\pi^- + p \longrightarrow \pi^0 + n \quad (j)$$

$$\pi^- + n \longrightarrow \pi^- + n \quad (f)$$

$$I_3 = -1 \quad -1/2$$

$$-3/2 \Rightarrow |3/2, -3/2\rangle$$

$$(c) \quad \pi^- + p \longrightarrow \pi^- + p.$$

$$I_3 = -1 + 1/2 = -1/2.$$

$$\text{tot } I_3 = -1/2.$$

$$|\pi^- + p\rangle = \alpha |3/2, -1/2\rangle + \beta |1/2, -1/2\rangle$$

Table 3.3. Clebsch-Gordan coefficients in pion-nucleon scattering

Pion	Nucleon	$I = 3/2$				$I = 1/2$	
		$I_3 = 3/2$	$1/2$	$-1/2$	$-3/2$	$1/2$	$-1/2$
π^+	p	1					
π^+	n		$\sqrt{\frac{1}{3}}$			$\sqrt{\frac{2}{3}}$	
π^0	p		$\sqrt{\frac{2}{3}}$			$-\sqrt{\frac{1}{3}}$	
π^0	n			$\sqrt{\frac{2}{3}}$			$\sqrt{\frac{1}{3}}$
π^-	p			$\sqrt{\frac{1}{3}}$			$-\sqrt{\frac{2}{3}}$
π^-	n				1		

$$\mathcal{M} = (\langle 3/2, -1/2 | \alpha + \langle 1/2, -1/2 | \beta) (H_S) (\alpha | 3/2, -1/2 \rangle + \beta | 1/2, -1/2 \rangle)$$

H_S conserves isospin

$$\Rightarrow \mathcal{M}_C = A \langle 3/2 | H_S | 3/2 \rangle + B \langle 1/2 | H_S | 1/2 \rangle$$

$$= \frac{1}{3} \mathcal{M}_{3/2} + \frac{2}{3} \mathcal{M}_{1/2}$$

$$\Rightarrow |\mathcal{M}_C|^2 = \frac{1}{9} |\mathcal{M}_3 + 2\mathcal{M}_1|^2$$

$$\pi^- + p \longrightarrow \pi^- + p.$$

$$(a) \quad \pi^+ + p \rightarrow \pi^+ + p.$$

$$|\mathcal{M}_a|^2 = |\mathcal{M}_3|^2$$

$$\mathcal{M}_a = \mathcal{M}_{3/2}$$

$$(j) \quad \pi^- + p \rightarrow \pi^0 + n.$$

$$\mathcal{M}_j \sim \frac{\sqrt{2}}{3} \mathcal{M}_{3/2} - \frac{\sqrt{2}}{3} \mathcal{M}_{1/2}$$

$$\sigma_j \propto |\mathcal{M}_j|^2 \sim \frac{2}{9} |\mathcal{M}_3 - \mathcal{M}_1|^2$$

Case 1)
if $|\mathcal{M}_1| > |\mathcal{M}_3| \Rightarrow \sigma_a(\pi^+ + p \rightarrow \pi^+ + p) \ll \sigma(\pi^-)$

$$\#(\pi^+ + p) \approx 0. \Rightarrow \text{not confirmed by data.}$$

Case 2)

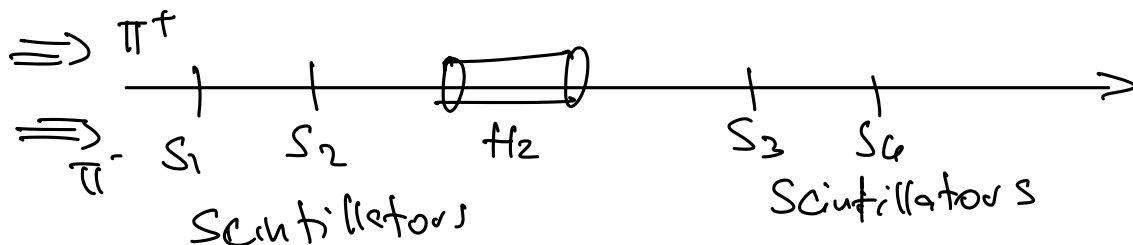
$$|\mathcal{M}_3| > |\mathcal{M}_1|$$

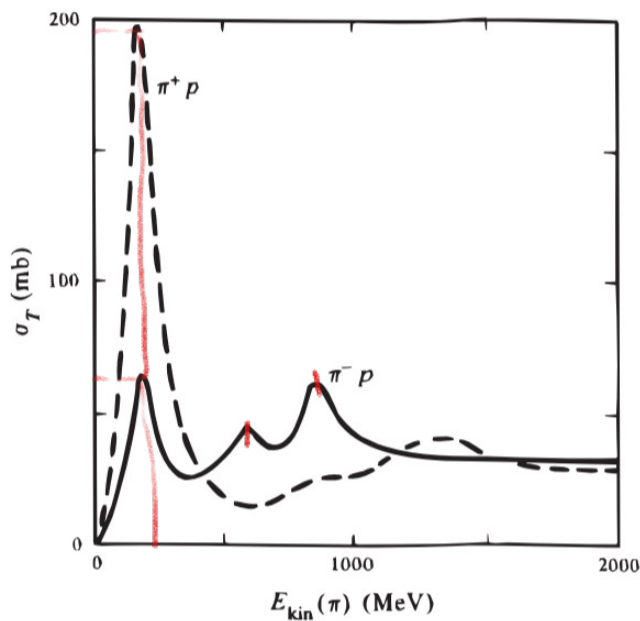
$$\mathcal{M}_a \approx \mathcal{M}_3 \Rightarrow \sigma(\pi^+ + p) \approx |\mathcal{M}_3|^2$$

$$\mathcal{M}_c \approx \frac{1}{\sqrt{3}} \mathcal{M}_3 \Rightarrow \sigma_c \approx \frac{1}{9} |\mathcal{M}_3|^2$$

$$\mathcal{M}_j \approx \frac{\sqrt{2}}{\sqrt{3}} \mathcal{M}_3 \Rightarrow \sigma_j \approx \frac{2}{9} |\mathcal{M}_3|^2$$

$$\frac{\#(\pi^+ + p \rightarrow \pi^+ + p)}{\#(\pi^- + p)} = \frac{|\mathcal{M}_3|^2}{\frac{1}{9} |\mathcal{M}_3|^2 + \frac{2}{9} |\mathcal{M}_3|^2} \approx 3$$





$$E_{kin} = E_{\pi} - m_{\pi}$$

Figure 5.35: Total cross section as a function of pion kinetic energy for the scattering of positive and negative pions from protons. (1 mb = 1 millibarn = 10^{-27} cm².)

	EM	Strong	Weak
P	✓	✓	✗
Q	✓	✓	✗
I		✓	
G		✓	
G parity			

$$Q |\pi^+\rangle = - |\pi^-\rangle \quad C_{q\bar{q}} = (-1)^{L+S}$$

$$Q |\pi^-\rangle = - |\pi^+\rangle$$

$$Q |\pi^0\rangle = + |\pi^0\rangle \quad C_{\pi^0} = (C_{\pi^+})^2 = +1$$

NB: $\pi^+ = u\bar{d} \neq q\bar{q}$

π^{\pm} are not eigenstates of Q because charged.

π^0 is eigenstate of Q .

$$C_{\pi^0} = (-1)^{L+S} = +1$$

R_2 : Rotation around I_2 in isospin space.

$$R_2 = e^{i\pi I_2}$$

$$R_2 |\pi^+\rangle = (-1)^{I-I_3} |\pi^-\rangle.$$

$$R_2 |I, I_3\rangle = (-1)^{I-I_3} |I, -I_3\rangle$$

$$R_2 |\pi^+\rangle = (-1)^{1-1} |1, -1\rangle = +1 |\pi^-\rangle$$

$$R_z |\pi^+ \rangle = (-1)^{1-(-1)} |\pi^+ \rangle = +1 |\pi^+ \rangle$$

$$R_z |\pi^0 \rangle = (-1)^{1-0} |\pi^0 \rangle = (-1) |\pi^0 \rangle$$

$$(Q + R_z) |\pi^+ \rangle = (-1)(+1) |\pi^+ \rangle = (-1) |\pi^+ \rangle$$

$$(Q + R_z) |\pi^- \rangle = (-1)(+1) |\pi^- \rangle = (-1) |\pi^- \rangle$$

$$(Q + R_z) |\pi^0 \rangle = (+1)(-1) |\pi^0 \rangle = (-1) |\pi^0 \rangle$$

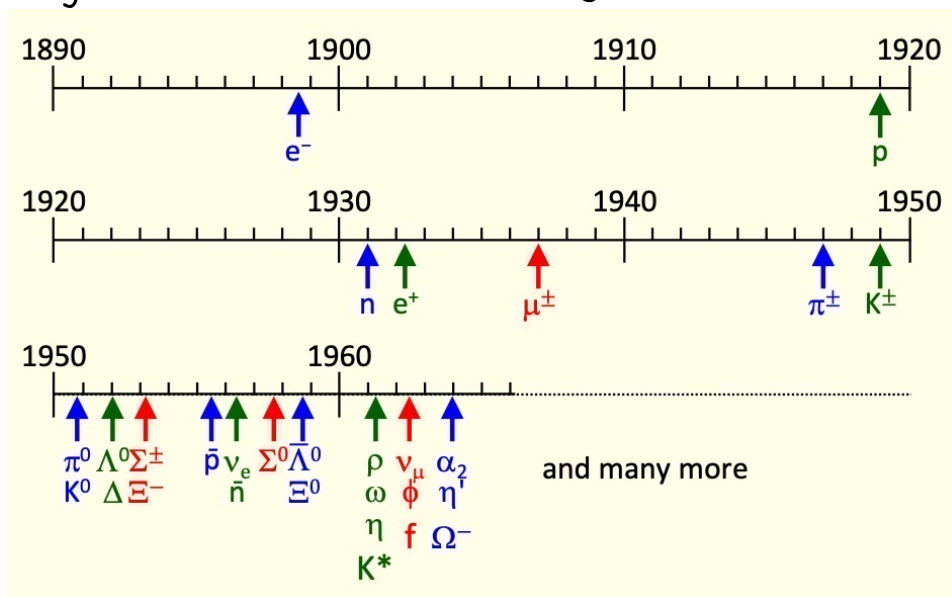
$$G = Q \times R_z$$

$$G = (-1)^{2+S+I}$$

for multiplet with L, S, I
the multiplet must be with $I = \text{integer}$

Same value on
Isospin multiplet

only conserved in strong interactions.



Strange particles :

produced strongly but decay
with long lifetime (weak interaction).

$$K^\pm, K^0, \Lambda$$

Gell Man - Nishijima Formula for Charge.

$$Q \rightarrow \frac{I_3}{2}$$

nucleus with Z proton + $A-Z$ neutron.

$$I_3 = Z \times \left(\frac{1}{2}\right) + (A-Z) \left(-\frac{1}{2}\right) = Z \frac{1}{2} + (A-Z) \left(-\frac{1}{2}\right) = \frac{A}{2} - \frac{Z}{2}$$

$$= Z - \frac{A}{2}$$

$$\Rightarrow Z = I_3 + \frac{A}{2}$$

$$\downarrow$$

$$Q = I_3 + \frac{B}{2}$$

Gell Mann - Nishijima

$$\pi^+ : +1 + \frac{0}{2} = +1$$

$$\text{strange particles} \Rightarrow Q = I_3 + \frac{B+S}{2}$$

Y
Hyper Charge

Name	π^\pm	π^0	K^\pm	K^0	η	p	n	Λ	$\Sigma^{\pm,0}$	Δ
Mass (MeV)	140	135	494	498	548	938	940	1116	1190	1232
Charge	± 1	0	± 1	0	0	1	0	0	$\pm 1, 0$	$2, \pm 1, 0$
Parity	-	-	-	-	-	+	+	+	+	+
Baryon n.	0	0	0	0	0	1	1	1	1	1
Spin	0	0	0	0	0	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{3}{2}$

Static Quark Model: classification of known hadrons based on symmetry.

Gell-Mann, Zweig 1961-1964

Eightfold way

Isospin, Spin $\frac{1}{2}$ particles $SU(2)$ symmetry.

$SU(N)$ operators:

$$U = e^{i a_j T_j} \quad j=1, \dots$$

$$\det = 1$$

$$\text{trace} = 0$$

T_j : generators.

Generators: $N^2 - 1$

$N-1$: are diagonal.

$SU(2)$: 3 generators S_x, S_y, S_z I_1, I_2, I_3

diagonal S_z

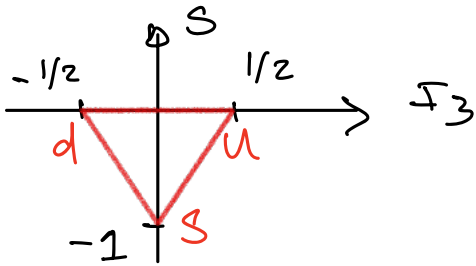
I_3

static quark model : \Rightarrow hadrons are multiplets of $SU(3)$

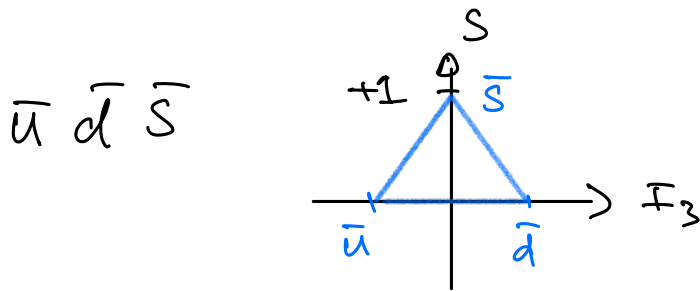
$SU(3)$: 8 generators

2 diagonal : I_3 , Strangeness.

Fundamental repp. of $SU(3)$ $u = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$ $d = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}$ $s = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$



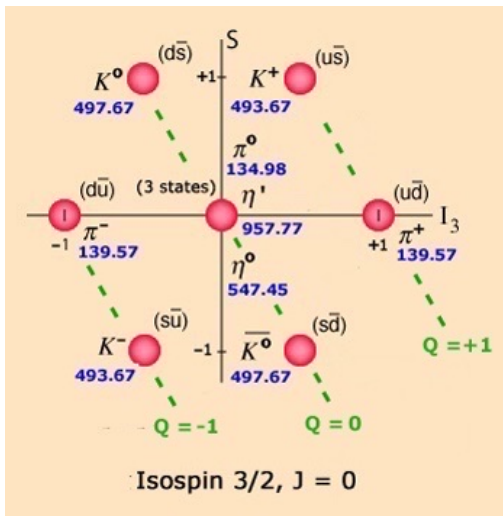
	I_3	S
u	$+1/2$	0
d	$-1/2$	0
s	0	-1



Baryon number

$B = \frac{1}{3}$ for quarks.

$= -1/3$ for antiquarks.



Fundamental repp:

3 or 3-bar

u \bar{u}
 d \bar{d}
 s \bar{s}

hadrons: $3 \otimes \bar{3} = 1 \oplus 8$

$u \times (\bar{u}, \bar{d}, \bar{s})$

In the middle.

$u\bar{u}, d\bar{d}, s\bar{s}$

