

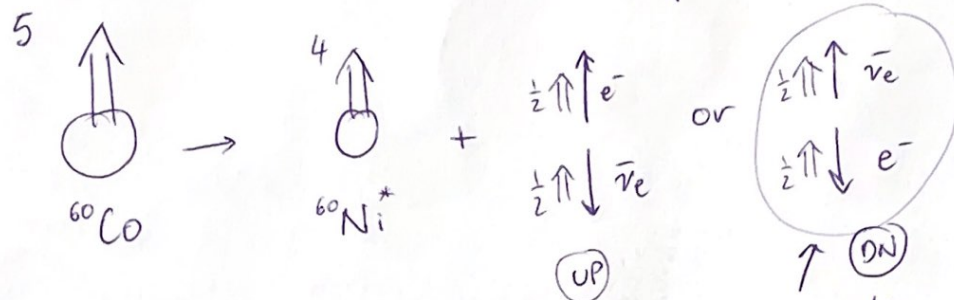
HELICITY OF THE NEUTRINO

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WU \rightarrow PARITY IS VIOLATED IN WEAK INT.

\rightarrow CROSS SECTIONS CAN DEPEND ON

$$\vec{\sigma} \cdot \vec{p}$$



WU says this one is favored

in (UP):

$h:$

$$e^-: \begin{matrix} \uparrow \uparrow \uparrow \vec{\sigma} \\ \rightarrow \vec{p} \end{matrix} \quad +1$$

$$\bar{\nu}_e: \begin{matrix} \leftarrow \uparrow \uparrow \uparrow \vec{\sigma} \\ \rightarrow \vec{p} \end{matrix} \quad -1$$

let's define helicity:

$$h = \frac{\vec{\sigma} \cdot \vec{p}}{|\vec{\sigma} \cdot \vec{p}|}$$

in (DN):

$h:$

$$e^-: \begin{matrix} \leftarrow \uparrow \uparrow \uparrow \vec{\sigma} \\ \rightarrow \vec{p} \end{matrix} \quad -1$$

$$\bar{\nu}_e: \begin{matrix} \rightarrow \uparrow \uparrow \uparrow \vec{\sigma} \\ \rightarrow \vec{p} \end{matrix} \quad +1$$

IN GENERAL h is NOT an invariant \rightarrow depends on frame

$\Rightarrow \exists$ boost β, γ that flips \vec{p} so
 $\begin{matrix} \vec{\sigma} \\ \Rightarrow \\ \vec{p} \end{matrix} \xrightarrow{\text{boost}} \begin{matrix} \vec{\sigma} \\ \Rightarrow \\ \vec{p} \end{matrix} \quad h=+1 \rightarrow h=-1$

HOWEVER in SM $m_\nu = 0 \Rightarrow V=C$

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\Rightarrow A boost that flips its \vec{p}

\Rightarrow h of neutrinos is invariant

i.e. if a neutrino is produced with $h=-1$
it will have $h=-1$ forever

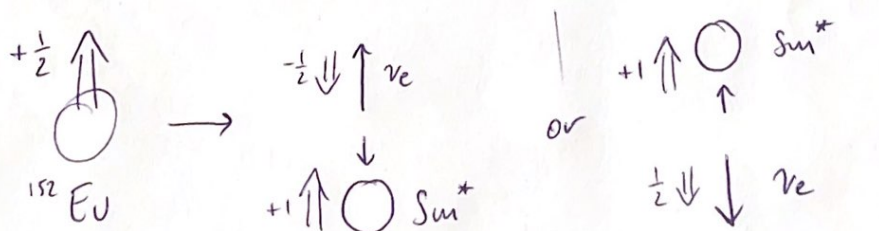
now, weak int. can distinguish between the two

if parity is maximally violated the weak
force will interact with only one $h=\pm 1$

\Rightarrow GOLDHABER EXPERIMENT (1957)

To measure the helicity of the neutrinos

Based on $^{152}\text{Eu} (+e^-) \xrightarrow{\text{electron capture}} \nu_e + ^{152}\text{Sm}^*$



here: $h(\nu) = -1$
 $h(\text{Sm}^*) = -1$

$h(\nu) = +1$
 $h(\text{Sm}^*) = +1$

$h(\nu) = h(\text{Sm}^*) !$

Then: EM decay of Su^*

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$$\tau(Su^*) \sim 10^{-14} s$$

$$^{152}Su^* \rightarrow ^{152}Su + \gamma$$

$$SPIN: \quad +1 \quad \quad 0 \quad +1$$

$$\Rightarrow \text{eg.} \quad \begin{array}{c} +1 \\ \uparrow \\ \bigcirc \end{array} \quad \quad 0 \quad \quad \begin{array}{c} \uparrow \\ \uparrow \\ \bigcirc \end{array} +1$$

So if γ emitted in direction of flight

$$h = +1 \quad \xRightarrow{\quad} \quad \begin{array}{c} \bullet \\ Su \end{array} \quad \xRightarrow{\quad} \quad h(\gamma) = h(Su^*) (=h(\nu))$$

$Su^* \quad \quad \gamma$

$$h = -1 \quad \xleftarrow{\quad} \quad \begin{array}{c} \bullet \\ Su \end{array} \quad \xleftarrow{\quad} \quad h(\gamma) = h(Su^*)$$

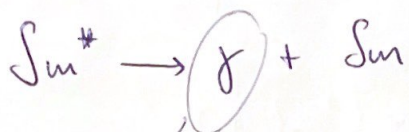
$Su^* \quad \quad \gamma$

TRANSFERRING THE HELICITY

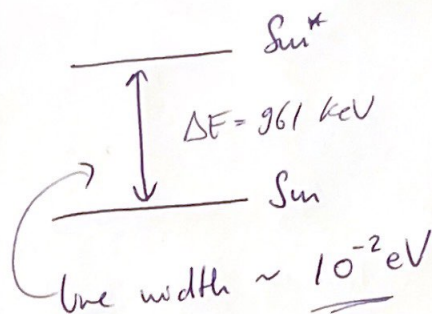
$$\text{from } \nu \text{ to } Su^* \text{ to } \gamma$$

Now

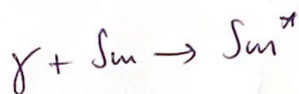
[4]



$$E(\gamma) = 961 \text{ keV}$$



if this ~~photon~~ encounters another Si it will not be reabsorbed by



why? Because part of the energy ($\sim 3.2 \text{ eV}$) is lost to give it to the recoiling Si

LAB
FRAME

(s.i.)
 Si^{*}
•
At rest

(s.f.)
 $Si \leftarrow \gamma$

$$\Rightarrow \text{in LAB frame } E(\gamma) < 961 \text{ keV} = \Delta E(Si^{*}, Si)$$

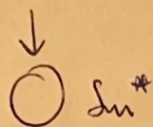
HOWEVER

E_0 ○ →

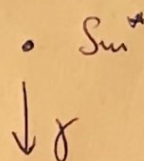
↑ z
↓
○ Si^{*}

← Si^{*} produced with boost in -z direction

LAB



CM



so if γ emitted in $-z$ (in CM)

\Rightarrow it will get a boost in $-z$

\Rightarrow E increases thanks to boost

\Rightarrow goes back to resonant 961 keV

THIS ONLY HAPPENS WHEN γ emitted in direction of flight of Sun^*

RECAP: Selecting γ that decay in direction of Sun^*

\Rightarrow (1) They have same h as neutrons

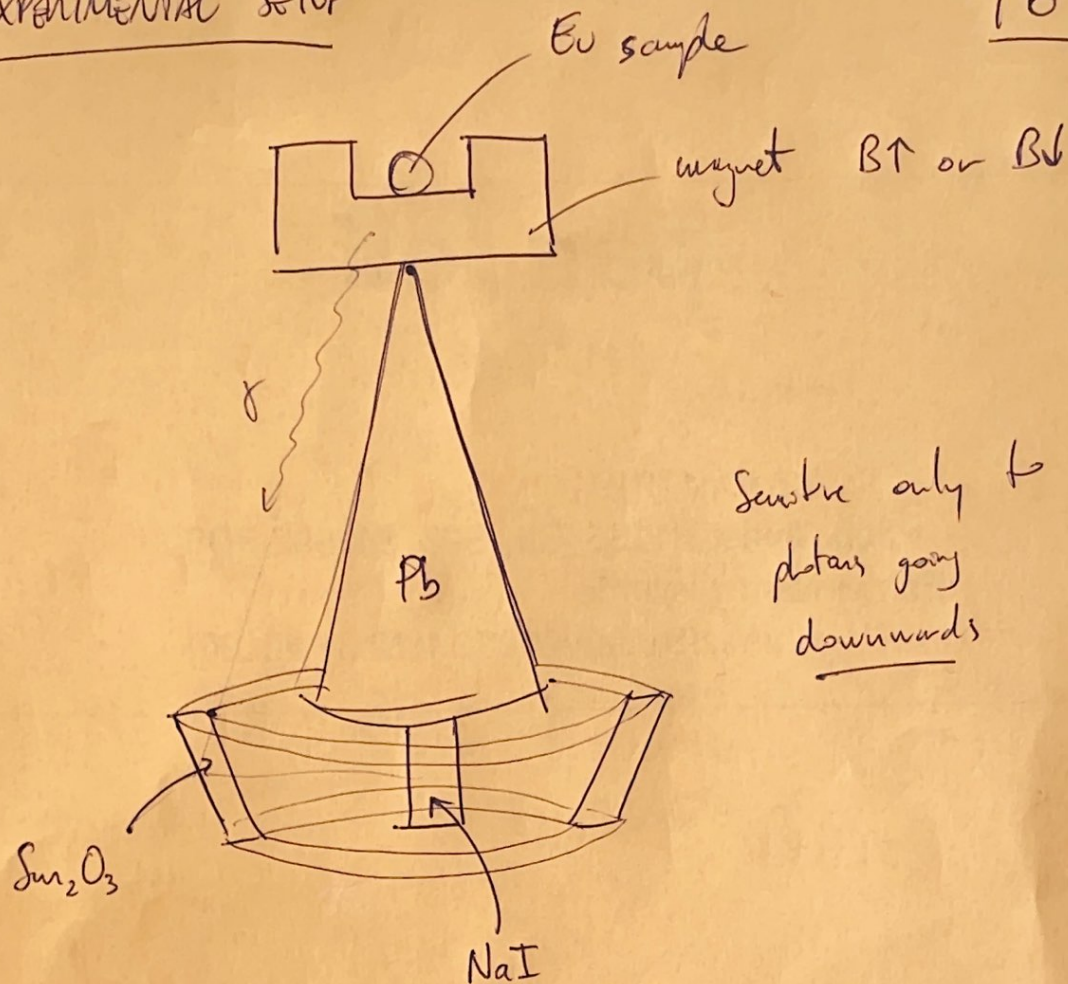
(2) they have $E = \Delta E(\text{Sun}^*, \text{Sun})$

↑
resonance

\Rightarrow can be reabsorbed by Sun

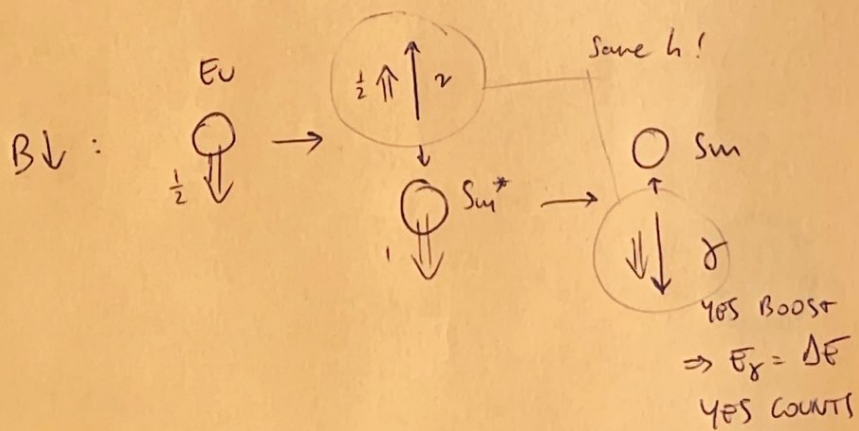
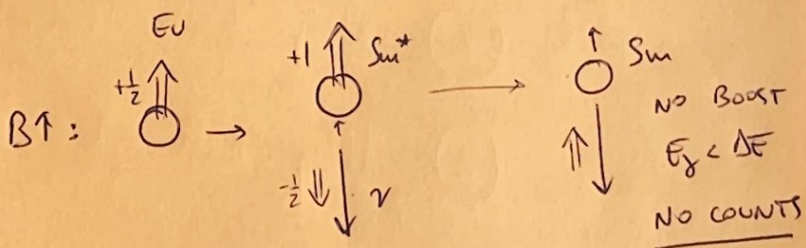
EXPERIMENTAL SETUP

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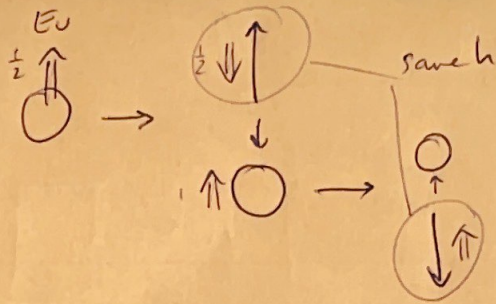


CASES

if $h(\gamma) = +1$

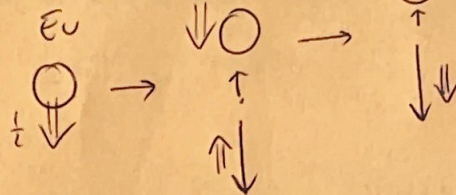


if $h(v) = -1$

$$B \uparrow =$$


YES BOOT
 $\Rightarrow E = \Delta E$
YES COUNTS

BL:



NO BOX
NO COUNTS

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	$h_v = +1$	$h_v = -1$
$B \uparrow$	NO COUNTS	YES COUNTS
$B \downarrow$	YES COUNTS	NO COUNTS

DATA compatible with ONLY $h_2 = -1$

\Rightarrow NEUTRONS only exist with $k=1$

And $\lim_{k \rightarrow \infty} \lambda_k$ exists with $k = +1$

→ WEAK INTERACTIONS only see

normals with $h = -1$

(and antineutrinos with $h = +1$)

What about neutrinos with $h=+1$?

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$q=0 \Rightarrow$ no EM interaction

$\text{color}=0 \Rightarrow$ no strong

$m=0 \Rightarrow$ no gravity (and no Higgs)

$h=+1 \Rightarrow$ no weak

\Rightarrow neutrinos with $h=+1$ DO NOT EXIST

(may have been produced at big bang
but then unable to interact with anything else)

BUT AS SAID HELICITY IS NOT AN INVARIANT

So if WEAK really "saw" only $h=-1$ would
mean that WEAK int. are not invariant

\rightarrow not a good physical theory

WHAT MATTERS IS CHIRALITY

$$\begin{aligned} \psi(x) &= \underbrace{\frac{1}{2}(1+\gamma^5)}_{P_R} \psi + \underbrace{\frac{1}{2}(1-\gamma^5)}_{P_L} \psi \\ &\equiv \underbrace{\psi_R(x)}_{\text{"RIGHT"}} + \underbrace{\psi_L(x)}_{\text{"LEFT"}} \end{aligned}$$

INvariant

→ WEAK INTERACTIONS ONLY COUPLE TO LEFT FIELDS 9

FOR MASSLESS PARTICLES

$$\underline{\text{HELICITY} = \text{CHIRALITY}}$$

$$\Rightarrow h = -1 \Leftrightarrow \text{LEFT}$$

$$h = +1 \Leftrightarrow \text{RIGHT}$$

FOR MASSIVE PARTICLES HELICITY \neq CHIRALITY

depends on frame

on β of particle in given frame

~~Example~~ let's take an electron with β

$$\text{if } h(e) = -1 \quad \text{helicity} = 1$$

\Rightarrow this state has both e_L and e_R components

$$\Rightarrow h = -1 \neq e_L$$

there is $e_R \neq 0$ but suppressed by $(1-\beta)$

thus is why $\pi^\pm \rightarrow \mu^\pm \nu_\mu$

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$$\pi^- \rightarrow \mu^- \bar{\nu}_\mu$$

$$\pi^- \rightarrow e^- \bar{\nu}_e$$

$$m(\pi^-) \sim 140 \text{ MeV}$$

$$m(\mu^-) \sim 106 \text{ MeV}$$

$$m(e^-) \sim 0.5 \text{ MeV}$$

$\Rightarrow \pi^- \rightarrow e^- \bar{\nu}_e$ is favored by phase space

BUT IN π^- rest frame

π^-
•
spin=0

\Rightarrow

$+\frac{1}{2} \uparrow \uparrow$
 $\bar{\nu}_e$
 $-\frac{1}{2} \downarrow \downarrow$
 e^-

massless \Rightarrow can only \exists with
 $h=+1$

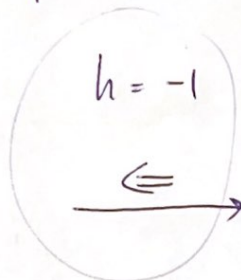
\Rightarrow to conserve spin $\Rightarrow h(e^-) = +1$

~~this is suppressed by (RFS)~~

only the left part of this is involved
in WEAK int

for particles:

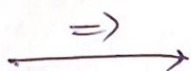
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$$\text{has } \psi_L \sim \frac{1}{2} (1 + \beta)$$

$$\psi_R \sim \frac{1}{2} (1 - \beta)$$

$$h = +1$$



$$\text{has } \psi_L \sim \frac{1}{2} (1 - \beta)$$

$$\psi_R \sim \frac{1}{2} (1 + \beta)$$

(OPPOSITE FOR ANTIPARTICLES)

$$\text{in } \pi^- \rightarrow e^- \bar{\nu}_e$$

$$\beta_e = 0.99997 \text{ (ex)}$$

\Rightarrow heavily suppressed

$$\text{for } \pi^- \rightarrow \mu^- \bar{\nu}_\mu$$

$$\beta_\mu = 0.27$$

$$\text{THIS IS WHY } R_\pi \equiv \frac{\pi \rightarrow e \nu}{\pi \rightarrow \mu \nu} \sim 10^{-4}$$