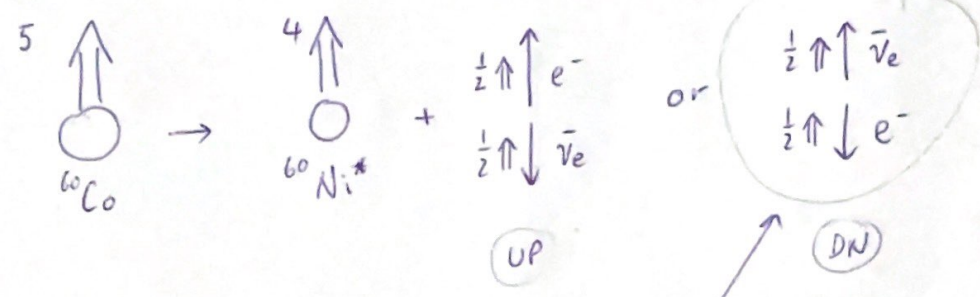


WU → PARITY VIOLATED IN WEAK INT.
(1956)

→ CROSS SECTIONS CAN DEPEND ON $\vec{\sigma} \cdot \vec{p}$



in (UP):
 e^- : $\begin{matrix} \uparrow \uparrow \uparrow \\ \downarrow \downarrow \downarrow \end{matrix} \quad h = +1$
 $\bar{\nu}_e$: $\begin{matrix} \uparrow \uparrow \uparrow \\ \downarrow \downarrow \downarrow \end{matrix} \quad h = -1$

in (DN):
 e^- : $\begin{matrix} \uparrow \uparrow \uparrow \\ \downarrow \downarrow \downarrow \end{matrix} \quad h = -1$
 $\bar{\nu}_e$: $\begin{matrix} \uparrow \uparrow \uparrow \\ \downarrow \downarrow \downarrow \end{matrix} \quad h = +1$

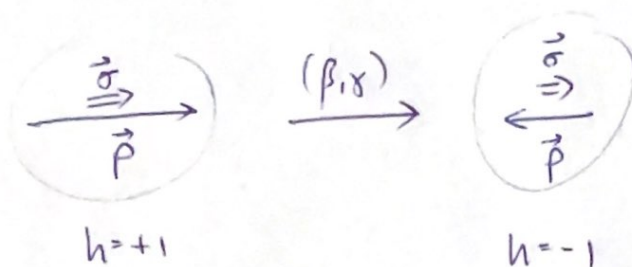
let's define helicity:
$$h = \frac{\vec{\sigma} \cdot \vec{p}}{|\vec{\sigma}| |\vec{p}|}$$

IN GENERAL h is NOT an invariant

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→ depends on frame

⇒ \exists boost (β, γ) that flips \vec{p} so



HOWEVER in SM $m_\nu = 0 \Rightarrow v=c$

⇒ \nexists boost that flips its \vec{p}

⇒ h of neutrinos is invariant

i.e. if a neutrino is produced with $h=-1$

⇒ it will have $h=-1$ forever

now weak int. can distinguish between the two

if parity is maximally violated

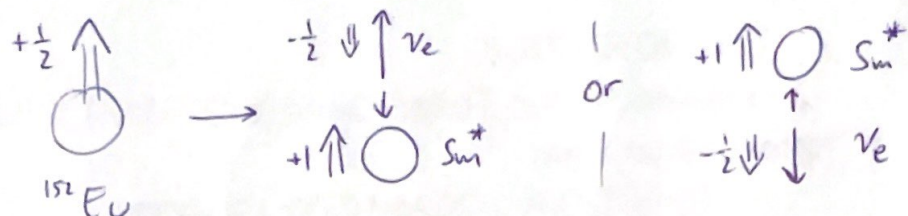
⇒ weak force will interact with only one
 $h = \pm 1$

GOLDHABER EXPERIMENT (1957)

3

To measure helicity of neutrinos

Based on $^{152}\text{Eu} (+e^-) \xrightarrow{\text{electron capture}} \nu_e + ^{152}\text{Sm}^*$
 $(p + e^- \rightarrow \nu_e + n)$



$$h(\nu) = -1$$

$$h(\text{Sm}^*) = -1$$

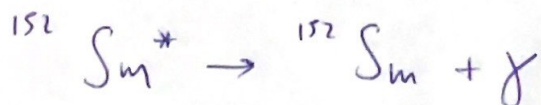
$$h(\nu) = +1$$

$$h(\text{Sm}^*) = +1$$

$$\Rightarrow h(\nu) = h(\text{Sm}^*) !$$

then: EM decay of Sm^*

$$\tau(\text{Sm}^*) \sim 10^{-14} \text{ s}$$



$$\text{Spin: } +1 \rightarrow 0 + 1$$



So if γ emitted in direction of flight

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$$h = +1 \quad \xRightarrow{S_{in}^*} \quad \begin{array}{c} \bullet \\ S_{in} \end{array} \xRightarrow{\gamma} \quad h(\gamma) = h(S_{in}^*) (=h(v))$$

$$h = -1 \quad \xRightarrow{S_{in}^*} \quad \begin{array}{c} \bullet \\ S_{in} \end{array} \xRightarrow{\gamma} \quad h(\gamma) = h(S_{in}^*) (=h(v))$$

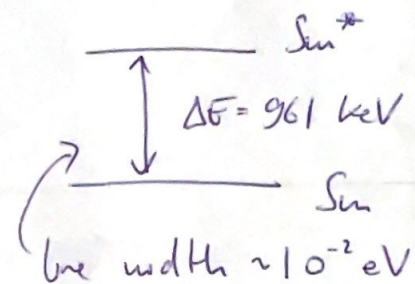
TRANSFERRING THE HELICITY

from v to S_{in}^* to γ

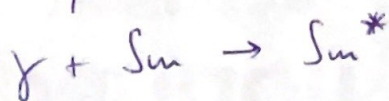
now



$$E(\gamma) = 961 \text{ keV}$$



if this γ encounters another S_{in} it will NOT be reabsorbed by



why? Because part of the energy ($\sim 3.2 \text{ eV}$) is lost to give it to recoiling S_{in}

LAB
FRAME

(S.I.)

Sm^*
•
AT REST

(S.F.)

Sm γ
← →

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⇒ in LAB frame $E(\gamma) < 961 \text{ keV} = \Delta E(Sm, Sm^*)$

HOWEVER

E_0
○ —

↑ v
↓

○ Sm^* ← Sm^* produced with
boost in the
-z direction

LAB

↓
○ Sm^*

CDM

• Sm^*

↖ ↗

so if γ emitted in -z

⇒ it will get a boost in -z

⇒ E increases

⇒ goes back to resonant 961 keV

IT JUST HAPPENS to regain exactly the energy
it needs

THIS ONLY HAPPENS WHEN γ emitted in direction of Sm^*

RECAP: Selecting γ that are emitted in
direction of Sun^* flight:

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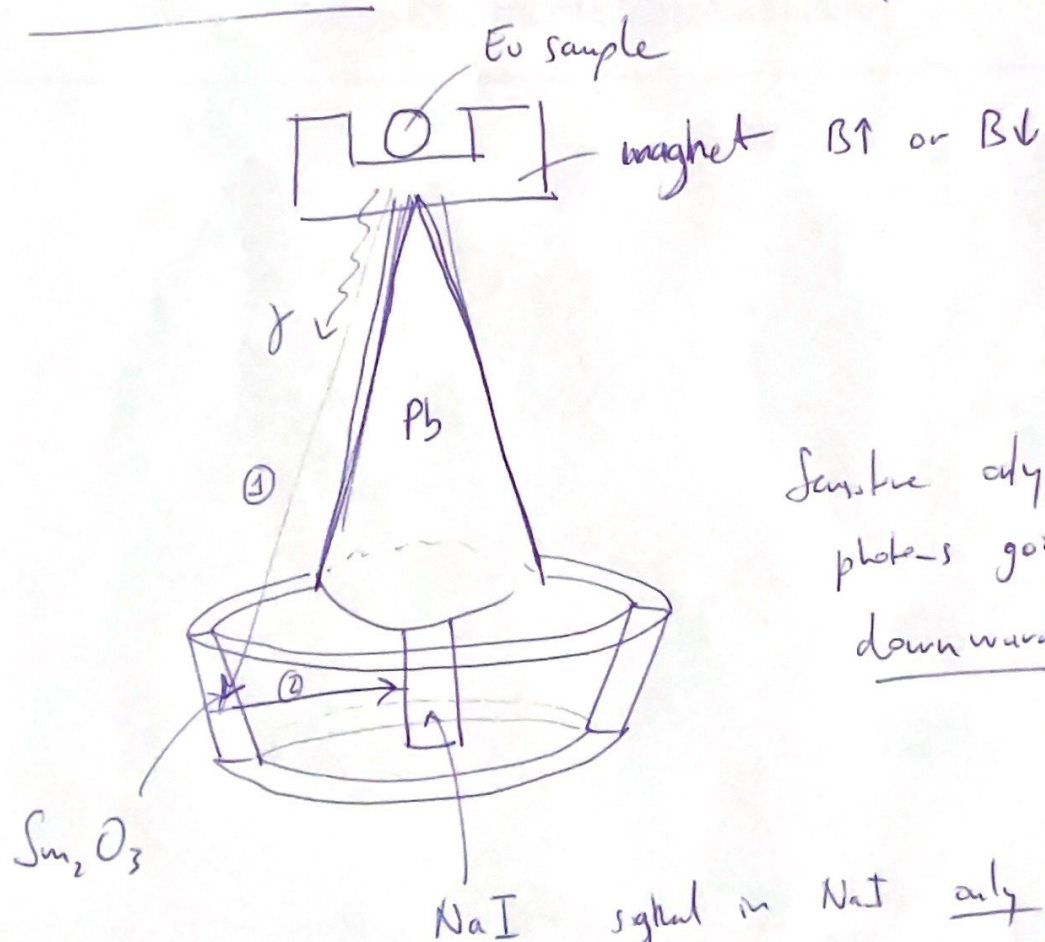
→ ① They have same h as neutrons

② They have $E = \Delta E(Sun, Sun^*)$

↑
resonance

⇒ can be reabsorbed
by Sun

EXPERIMENTAL SETUP

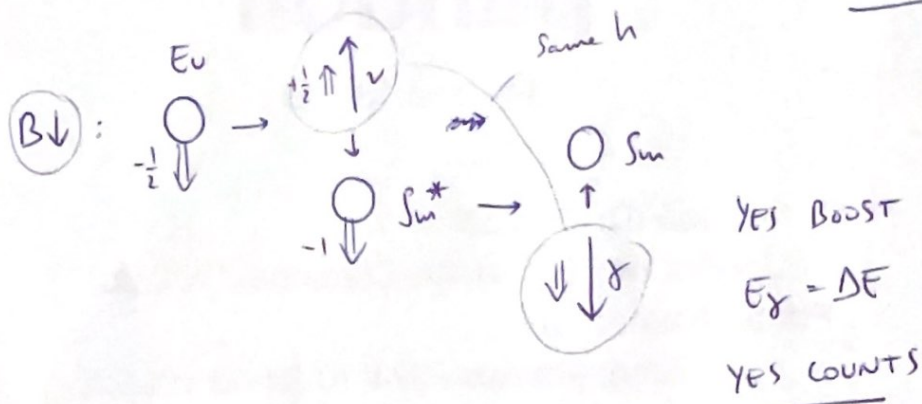
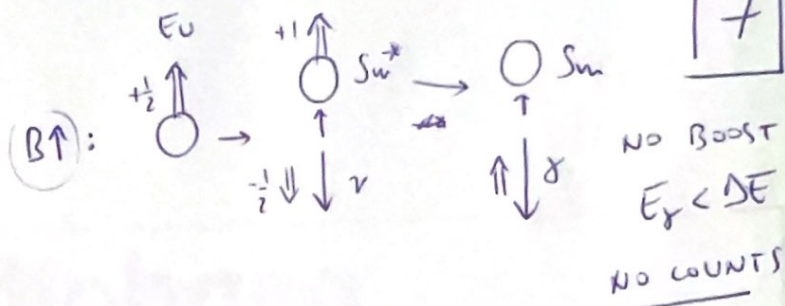


Sensitive only to
photo- γ s going
downwards

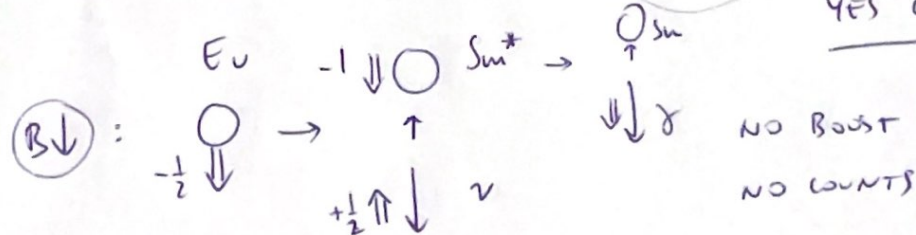
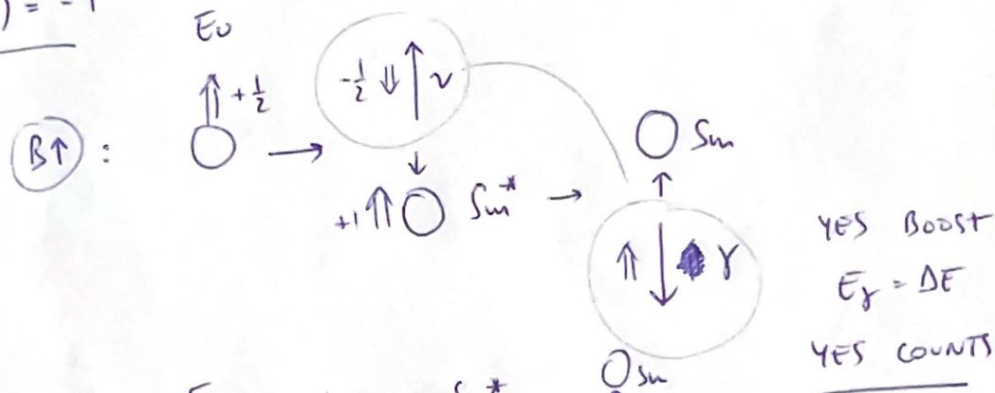
signal in NaI only
if γ emitted downwards
and with right energy

CASES

if $h(v) = +1$



if $h(v) = -1$



so

	$h_v = +1$	$h_v = -1$
B↑	NO COUNTS	YES COUNTS
B↓	YES COUNTS	NO COUNTS

DATA compatible
with only
 $h_v = -1$

⇒ WEAK INTERACTIONS only see
 NEUTRINOS with $h = -1$
 (and antineutrinos with $h = +1$)

WHAT ABOUT NEUTRINOS with $h = +1$?

$q = 0 \Rightarrow$ no EM interaction

color = 0 \Rightarrow ~~no~~ no strong

$m = 0 \Rightarrow$ no gravity (and no Higgs)

$h = +1 \Rightarrow$ no weak

⇒ neutrinos with $h = +1$ DO NOT EXIST IN SM

(may have been produced at big bang
 but then unable to interact with anything
 else \rightarrow so they have remained undetected)

BUT AS SAID HELICITY IS NOT INVARIANT

So if weak int. really "saw" only $h = -1$

\Rightarrow weak int are not invariant

\rightarrow ~~not~~ a good physical theory!

WHAT MATTERS IS CHIRALITY

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$$\psi(x) = \underbrace{\frac{1}{2}(1+\gamma^5)}_{P_R} \psi + \underbrace{\frac{1}{2}(1-\gamma^5)}_{P_L} \psi$$

$$\equiv \underbrace{\psi_R(x)}_{\text{"RIGHT"}} + \underbrace{\psi_L(x)}_{\text{"LEFT"}} \quad \text{INVARIANT}$$

→ WEAK INTERACTIONS ONLY COUPLE TO LEFT FIELDS

now, for massless particles: HELICITY = CHIRALITY

$$\Rightarrow h = -1 \Leftrightarrow \text{LEFT}$$

$$h = +1 \Leftrightarrow \text{RIGHT}$$

for MASSIVE PARTICLES HELICITY \neq CHIRALITY

↑
depends on frame,
or β of particle in given frame

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let's take electron with β

$$\text{if } h(e) = -1 \quad \text{helicity} = -1$$

\Rightarrow this state has both e_L and e_R components

$$\Rightarrow h_e = -1 \neq \psi = e_L$$

there is $e_R \neq 0$ but suppressed by $(1-\beta)$

that is why $\pi^\pm \rightarrow \mu^\pm \nu_\mu$

$$\pi^- \rightarrow \mu^- \bar{\nu}_\mu$$

$$\pi^- \rightarrow e^- \bar{\nu}_e$$

$$m(\pi^-) \sim 140 \text{ MeV}$$

$$m(\mu^-) \sim 106 \text{ MeV}$$

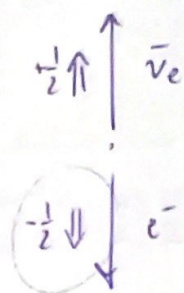
$$m(e^-) \sim 0.5 \text{ MeV}$$

$\Rightarrow \pi^- \rightarrow e^- \bar{\nu}_e$ is favored by phase space

BUT in π^- rest frame

π^-
•
Spin = 0

\Rightarrow



massless \Rightarrow can only
 \exists with $h = +1$

\Rightarrow to conserve spin $h(e^-) = +1$

only the LEFT part of this electron is
involved in WEAK int

for particles:

$h = -1$

$$\Rightarrow \text{hms } \psi_L \sim \frac{1}{2} (1 + \beta)$$

$$\psi_R \sim \frac{1}{2} (1 - \beta)$$

$h = +1$

$$\Rightarrow \text{hms } \psi_R \sim \frac{1}{2} (1 + \beta)$$

$$\psi_L \sim \frac{1}{2} (1 - \beta)$$

(opposite for ANTIPARTICLES)

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in $\pi^- \rightarrow e^- \bar{\nu}_e$

$$\beta_e = 0.99997 \quad (\text{ex}) \quad \text{and} \quad h = +1$$

$\Rightarrow \psi_L$ is heavily suppressed

for $\pi^- \rightarrow \mu^- \bar{\nu}_\mu$

$$\beta_\mu = 0.27$$

THIS IS WHY $R_\pi = \frac{\pi \rightarrow e \nu}{\pi \rightarrow \mu \nu} \sim 10^{-4}$

WEAK FORCE INTERACTS

ONLY WITH $\psi_L(x) \leftarrow$ PARTICLES

$\psi_R(x) \leftarrow$ ANTIPARTICLES

\Rightarrow ONLY LEFT (RIGHT) COMPONENT OF $\psi(x)$

PARTICIPATES ~~TO~~ WEAK INTERACTIONS

FOR PARTICLES (ANTIPARTICLES)