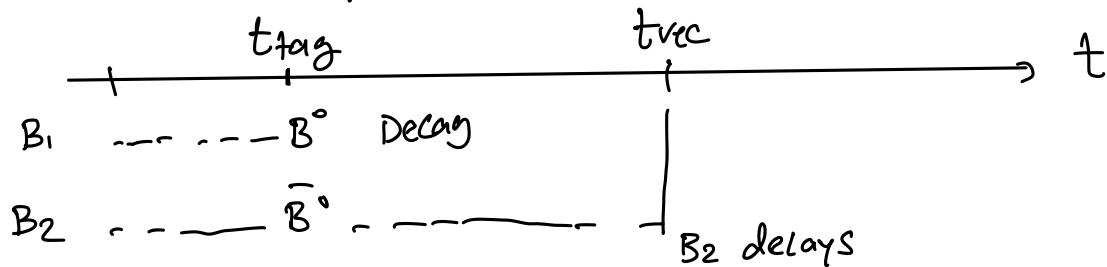


CP violation at B factory

$e^+ e^- \rightarrow I(GS) \rightarrow |B^0 \bar{B}^0\rangle \rightarrow B_1 \rightarrow \text{CP eigen state}$

$B_2 \rightarrow \text{tag/flavor eigen state}$

$$|B^0 \bar{B}^0\rangle = \frac{1}{\sqrt{2}} (|B^0\rangle |\bar{B}^0\rangle - |\bar{B}^0\rangle |B^0\rangle)$$



t_{tag} : time of delay $B \rightarrow$ flavor eigenstate

$$w^{\pm} \times \begin{cases} u & u \\ d & s \\ \pi^+ & K^+ \\ \nu & \nu_e \nu_\mu \nu_\tau \end{cases} \ell^\mp e^\pm \mu^\pm \tau^\pm$$

$$\bar{b} \quad \bar{c} \quad D^-$$

$$Q = m_B - m_D - m_X$$

$$= 5.279 - 1.86 - m_X.$$

$$B^0 \rightarrow D^- \pi^+$$

$$D^- K^+$$

$$D^- \ell^+ \nu$$

$$B^0 \rightarrow \ell^+ \quad \text{tagger for } \bar{b} \text{ quark.}$$

$$\bar{B}^0 \rightarrow \ell^- \quad \text{flavor tag for } b \text{ quark.}$$

$$\text{For } t = t_{tag} \quad B_2 = \bar{B}^0$$

$t > t_{tag}$ \bar{B}^0 can oscillate or decay.

$t_{rec} > t_{tag}$: B_2 decays.

2 scenarios:

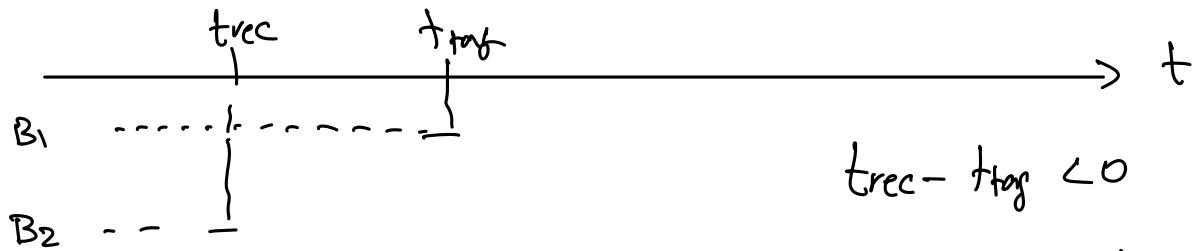
$$\begin{cases} B_1 = B^0 \\ B_2 = \bar{B}^0 \end{cases}$$

no oscillation

$$\begin{cases} B_1 = B^0 \\ B_2 = B^0 \end{cases}$$

B_2 oscillated after t_{tag} .

$$t_{\text{rec}} > t_{\text{tag}} \Rightarrow \Delta t = t_{\text{rec}} - t_{\text{tag}} > 0$$



$$P(B^0 \rightarrow B^0) = e^{-\beta t} (1 + \cos \Delta m \Delta t) \quad \Delta t = t_{\text{rec}} - t_{\text{tag}}$$

$$P(B^- \rightarrow \bar{B}^0) = P(\bar{B}^0 \rightarrow B^0) = e^{-\beta t} (1 - \cos \Delta m \Delta t)$$

$$\Delta t = \frac{\Delta z}{\beta v c}$$

Δz : distance between
decay vertices of B_1 and B_2 .

$$\sqrt{s} = 10.58 \text{ GeV} = m(Y(4S))$$

$$Q = m(Y(4S)) - 2m_B \Rightarrow P_B^* \leq 300 \text{ MeV.}$$

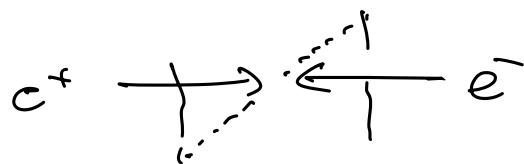
$$(\beta v)_B^* = \frac{300 \text{ MeV}}{5 \text{ GeV}} = \frac{0.3}{5} \approx 0.06$$

$$\tau_B = 1.5 \text{ fs} = 1.5 \times 10^{-15} \text{ s.}$$

$$\langle d \rangle = \beta v c \tau = 0.06 \times 3 \times 10^{10} \frac{\text{cm}}{\text{s}} \times 1.5 \times 10^{-15} \text{ s.}$$

$$\approx 10^{-6} \text{ cm.}$$

If symm e^+e^- beam:



$\approx 10^{-2} \text{ mm.}$ not useful/easy
to measure experimentally

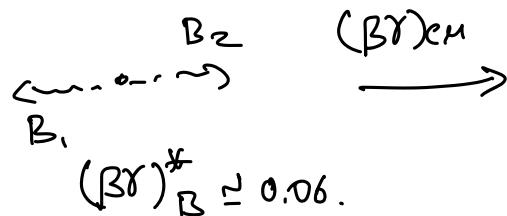
use Lorentz boost to increase Δz .

$$e^- \xrightarrow{9 \text{ GeV}} \xleftarrow{3.1 \text{ GeV}} e^+ \rightarrow z \quad P_{\text{tot}} = 9 - 3.1 = 5.9 \text{ GeV.} \rightarrow z$$

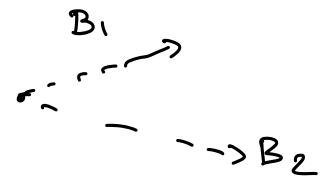
$$(\beta\gamma)_{CM} = \frac{S \cdot g}{\sqrt{s}} = \frac{S \cdot g}{10.58 \text{ GeV}} = 0.56$$

center of mass.

In the LAB



$$(\beta\gamma)_B^* \approx 0.06.$$



$$\Delta t \approx \frac{\Delta z}{\beta c}$$

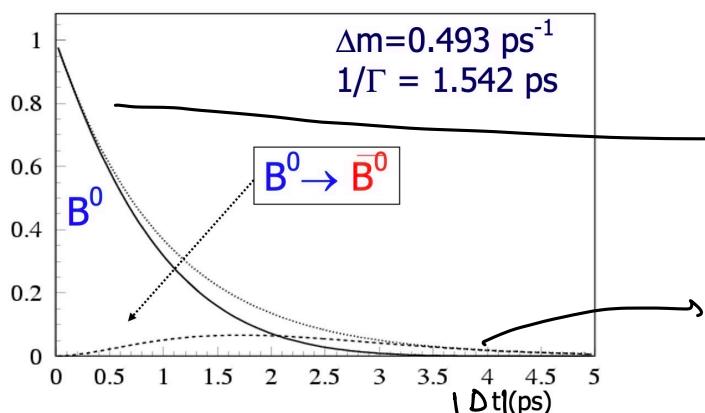
$$\Delta z \approx 300 \mu m.$$

this can be measured.

$$A = \frac{\#(B^0 \bar{B}^0) - \#(B^0 B^0 \text{ or } \bar{B}^0 \bar{B}^0)}{\#(B^0 \bar{B}^0) + \#(B^0 B^0 \text{ or } \bar{B}^0 \bar{B}^0)}$$

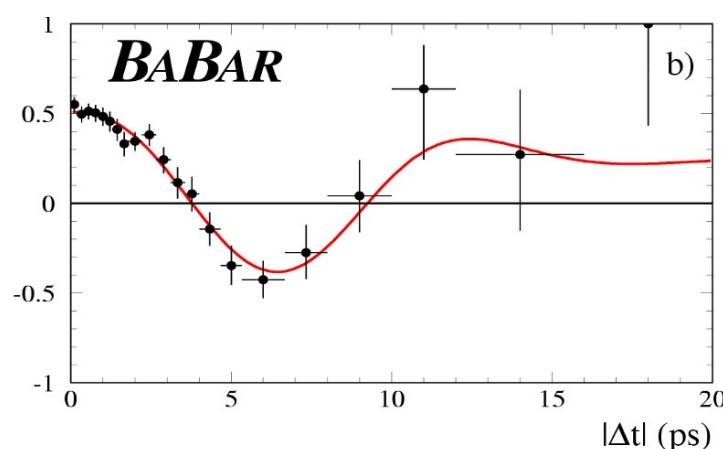
oscillation amplitude.

$$= \frac{(1 + \cos \Delta m \Delta t) - (1 - \cos \Delta m \Delta t)}{(1 + \cos \Delta m \Delta t) + (1 - \cos \Delta m \Delta t)} = \cos \Delta m \Delta t$$



$$(1 + \cos \Delta m \Delta t) \frac{1}{2} e^{-\Gamma t}$$

$$(1 - \cos \Delta m \Delta t) \frac{1}{2} e^{-\Gamma t}$$



$$0.5 = 1 - \tilde{\omega} w$$

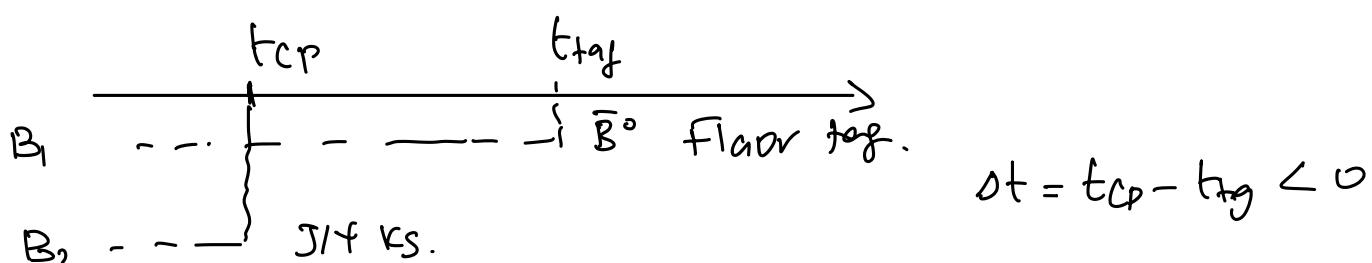
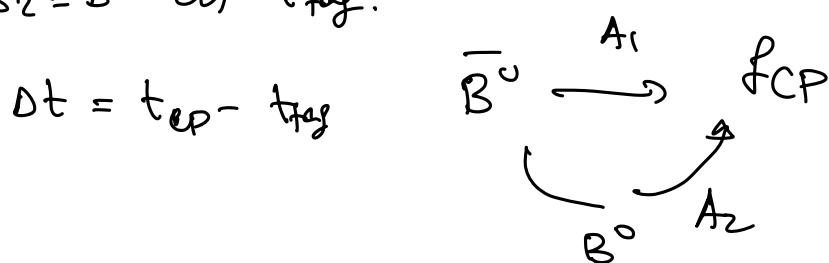
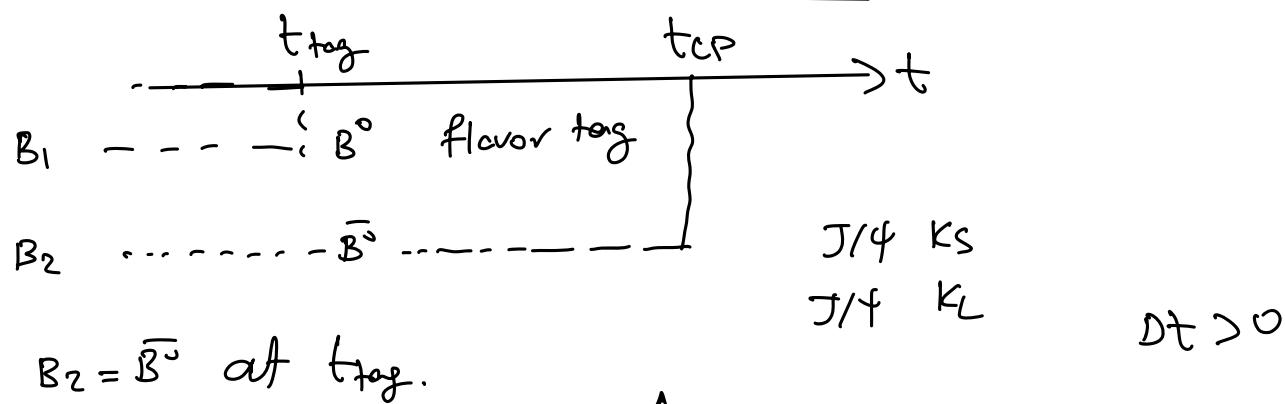
w: prob of wrong flavor tag.

$$P(\text{surv}) = e^{-t/\tau} \quad t = 7\tau \Rightarrow P(\text{surv}) = e^{-7}$$

$$\#(B^0 \bar{B}^0) = (1-w) N(B^0 \bar{B}^0) + w N(B^0 \bar{B}^0)$$

w: prob. of making a mistake in flavor tag.

CP violation in interference



$J/\psi = c\bar{c}$ bound state

$$K_S = \frac{1}{\sqrt{2}} (K^0 + \bar{K}^0) \leq K_1 \quad K^0: \bar{s}d \quad \bar{K}^0 = s\bar{d}$$

$$K_L = \frac{1}{\sqrt{2}} (K^0 - \bar{K}^0)$$

$$\begin{aligned} CP(J/\psi K_S) &= P(J/\psi K_S) C(J/\psi K_S) \\ &= (-)^L P(J/\psi) P(K_S) C(J/\psi) C(K_S) \end{aligned}$$

$$S \begin{matrix} \downarrow & \downarrow \end{matrix} \Gamma_{US} \rightarrow \begin{matrix} B^0 & \bar{B}^0 \\ 0 & 0 \end{matrix} \Rightarrow L=1.$$

$$CP(J/\psi K_L) = (-1)^L CP(K_S) CP(J/\psi)$$

$\sim \pm 1.$

$$CP(K_L) = -1.$$

$$P(J/\psi) = -1. \quad CP(J/\psi) = +1.$$

$$C(J/\psi) = (-1)^{L+S} = (-1)^{0+1} = -1.$$

$$CP|J/\psi K_S\rangle = (-1)|J/\psi K_S\rangle$$

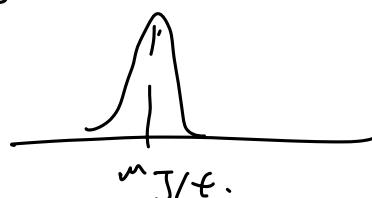
$$CP|J/\psi K_L\rangle = (+1)|J/\psi K_L\rangle$$

$$J/\psi \rightarrow e^+ e^- \quad \text{invariant mass } M_{ll}$$

$$K_S \rightarrow \pi^+ \pi^- \quad M_{\pi\pi}.$$

$$e^+ e^- \rightarrow B \text{tag} (\ell, K_C, \dots) \quad B_{CP}(\bar{\ell}^+, \pi^+ \pi^-)$$

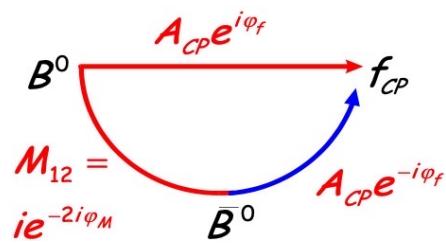
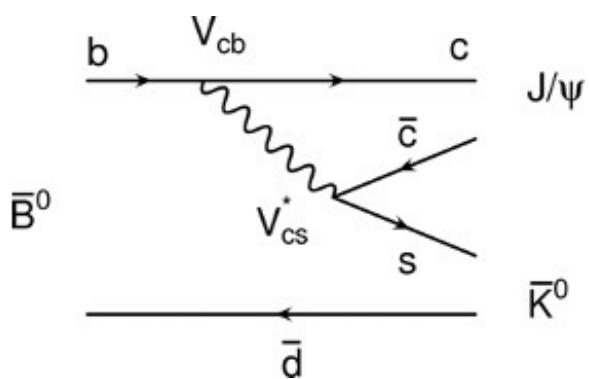
$$e^+ e^-$$



$m_{\pi\pi}$

$$B \rightarrow \ell \ell \pi \pi \quad \text{Clear experimental signature.}$$

m_B

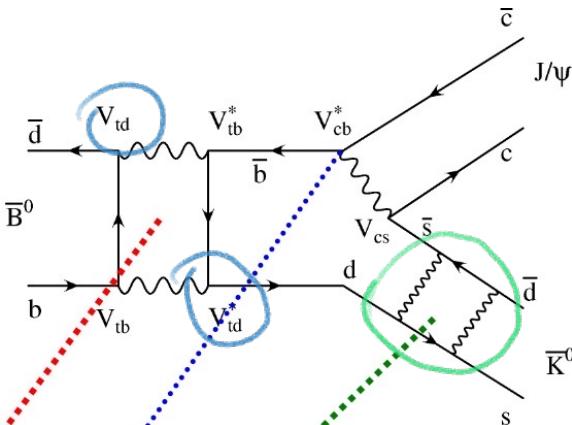
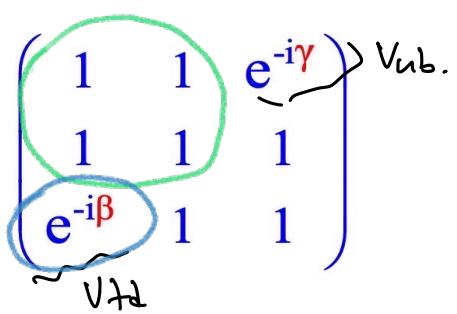


NO complex phase in V_{cb} , Yes in Wolfenstein parameter.

\Rightarrow NO complex phase in $\bar{B}^0 \rightarrow J/\psi K_S$

$$V_{CKM} = \begin{pmatrix} V_{ud} & & \\ & V_{ub} & \\ & & V_{tb} \end{pmatrix} + \mathcal{O}(\lambda^4)$$

$$V_{ud} = \begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix}$$



$$\lambda_{\psi K_S} = \frac{q_B}{p_B} \frac{\bar{A}_{\psi K_S}}{A_{\psi K_S}} = \frac{V_{tb}^* V_{td}}{V_{tb} V_{td}^*} \frac{V_{cb} V_{cs}}{V_{cb}^* V_{cs}} \frac{V_{cs} V_{cd}}{V_{cs}^* V_{cd}} = - \frac{V_{tb}^* V_{td}}{V_{tb} V_{td}^*} \frac{V_{cb} V_{cd}}{V_{cb}^* V_{cd}}$$

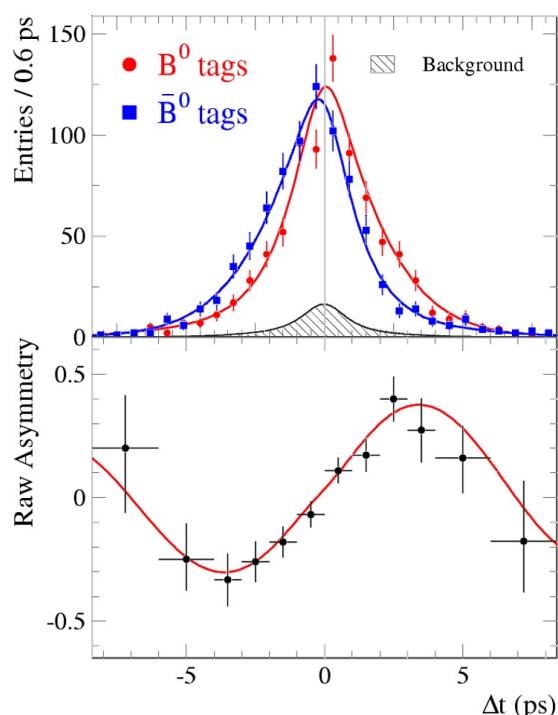
$$e^{-i2\beta}$$

$f(\bar{B}^0 \rightarrow J/\psi K_S)$ } function $\sin(\Delta m \Delta t) \sin 2\beta$.

$f(B^0 \rightarrow J/\psi K_S)$

$\#(\bar{B}^0 \text{ tagged} \rightarrow J/\psi K_S)$. function

$\#(B^0 \text{ tagged} \rightarrow J/\psi K_S)$



$\bar{B}^0 \text{ tag} \Rightarrow B_{CP} = \bar{B}^0 \text{ at tag.}$

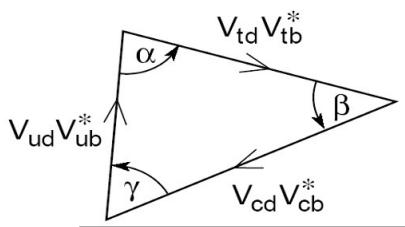
$$e^{-pt} (1 + \sin 2\beta \sin \Delta m \Delta t)$$

$\bar{B}^0 \text{ tag} \Rightarrow B_{CP} = B^0 \text{ at tag}$

$$e^{-pt} (1 - \sin 2\beta \sin \Delta m \Delta t)$$

$$A_{CP} = \frac{\#(B^0 \text{ tag}) - \#(\bar{B}^0 \text{ tag})}{\# + \#}$$

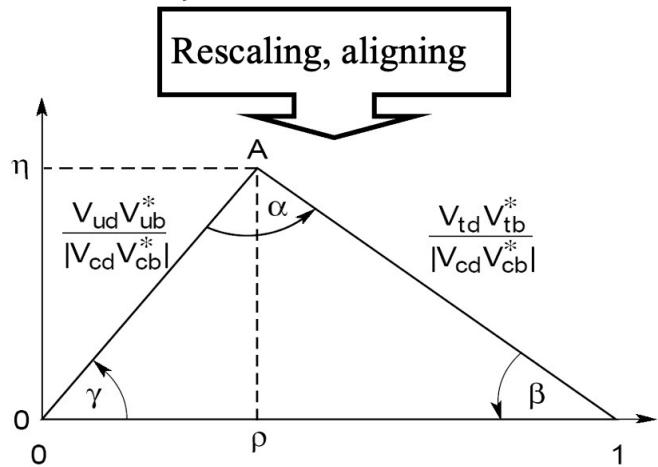
$$\sin 2\beta = 0.741 \pm 0.067_{(stat)} \pm 0.033_{(syst)}$$



amplitude $\neq 0$.

$\Rightarrow \sin 2\beta \neq 0$.

$\Rightarrow \beta \neq 0$.



$\Rightarrow V_{td}$ is complex.

\Rightarrow Kobayashi-Maskawa mechanism for CP correct.

\Rightarrow complex phase in matrix.
CKM

KM mechanism \Rightarrow 3 generations of quarks/matter

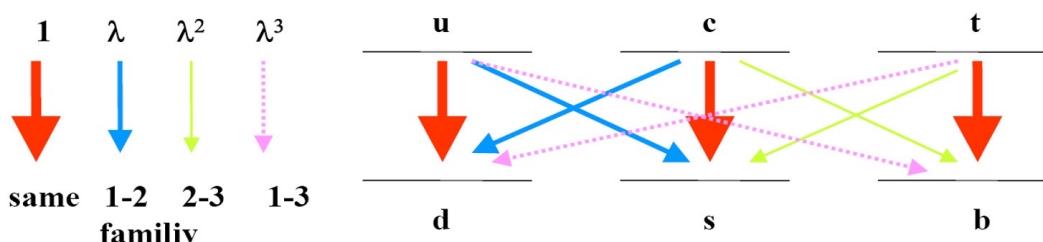
$$\begin{pmatrix} e^- \\ \nu_e \end{pmatrix}_L \quad \begin{pmatrix} \mu^- \\ \nu_\mu \end{pmatrix}_L \quad \begin{pmatrix} \tau^- \\ \nu_\tau \end{pmatrix}_L \quad \text{doublet of weak isospin.}$$

$$\begin{pmatrix} u \\ d \end{pmatrix}_L \quad \begin{pmatrix} c \\ s \end{pmatrix}_L \quad \begin{pmatrix} t \\ b \end{pmatrix}_L$$

$e_R, \bar{\nu}_R, u_R, d_R, \dots$ singlets of weak isospin.

$\xrightarrow{\text{family mass}}$

experimental observation.

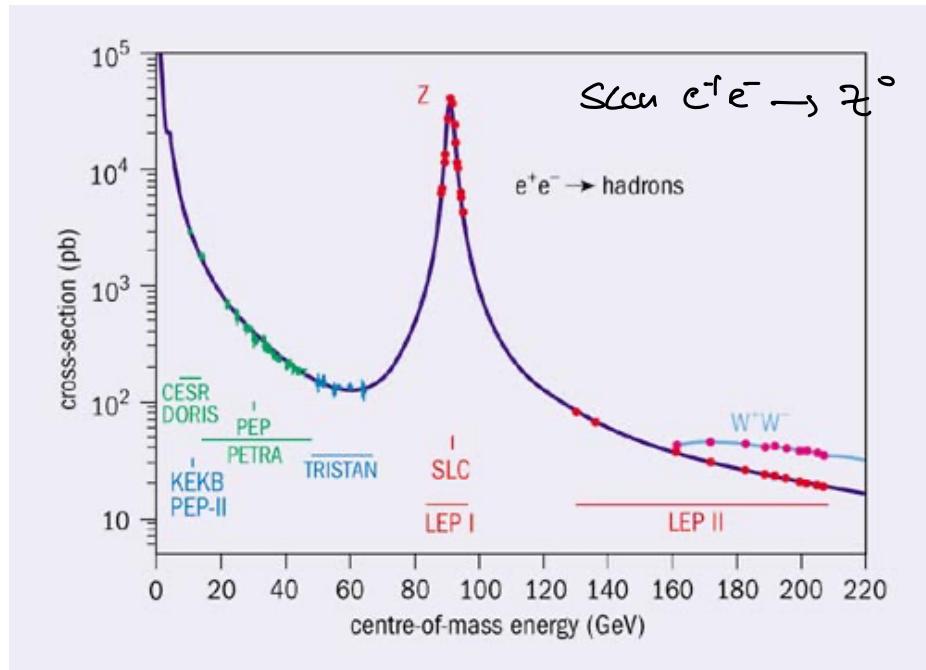
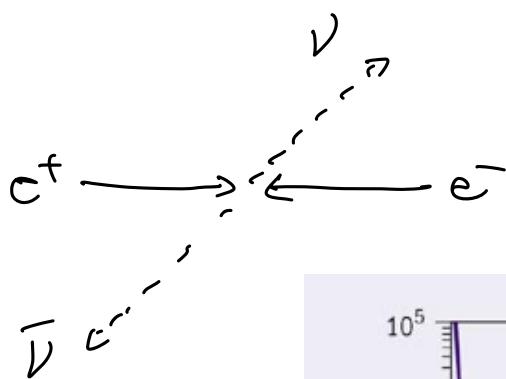


Weak interaction strength suppressed across families.

$\tilde{\chi}^0 \rightarrow$
 $e^+ e^-$
 $\mu^+ \mu^-$
 $\nu \bar{\nu}$
 $u \bar{u}$
 $d \bar{d}$
 $b \bar{b}$
 $x \bar{x}$

$e^+ e^- \rightarrow \tilde{\chi}^0 \rightarrow$ measure $\tilde{\chi}^0$
 $\sqrt{s} = m_{\tilde{\chi}^0}$ decay width.

In particular : $\tilde{\chi}^0 \rightarrow \nu \bar{\nu}$ Experimentally.
 $\nu_e \bar{\nu}_e$
 $\nu_\tau \bar{\nu}_\tau$
 $\nu_x \bar{\nu}_x$
 $\tilde{\chi}^0 \rightarrow$ invisible
measure at LEP.

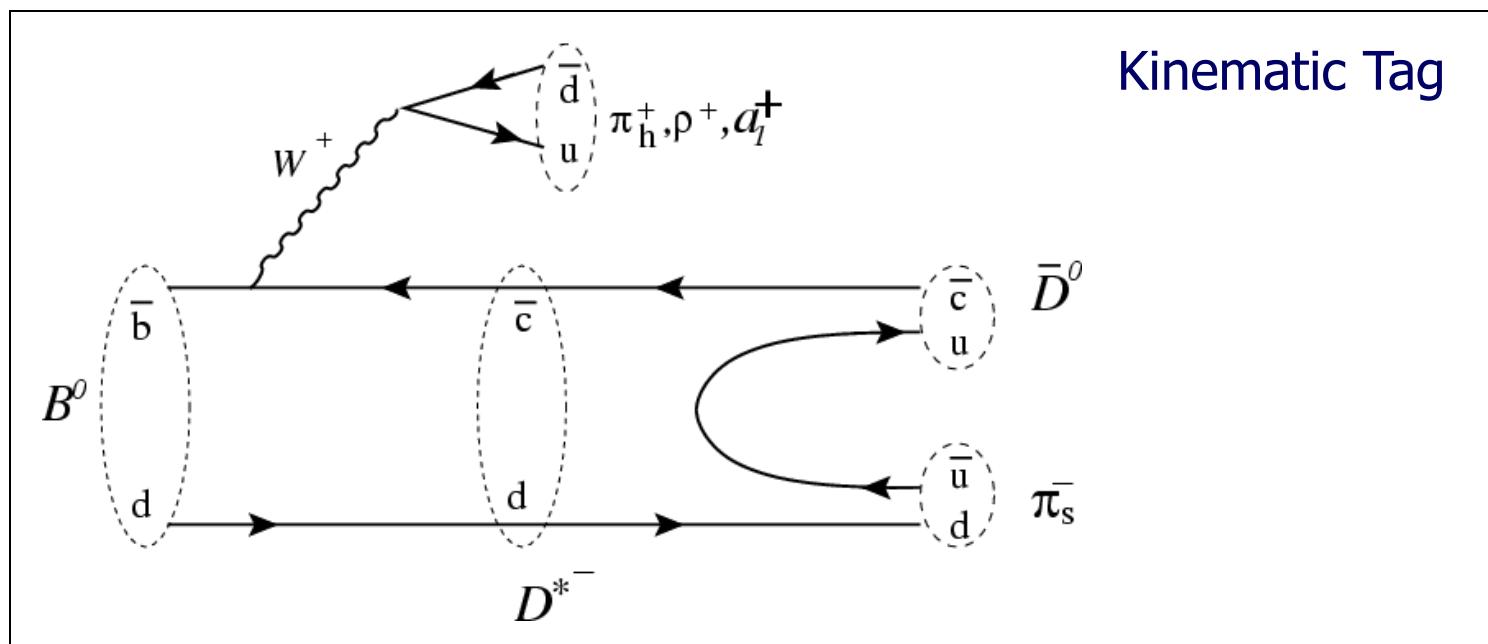
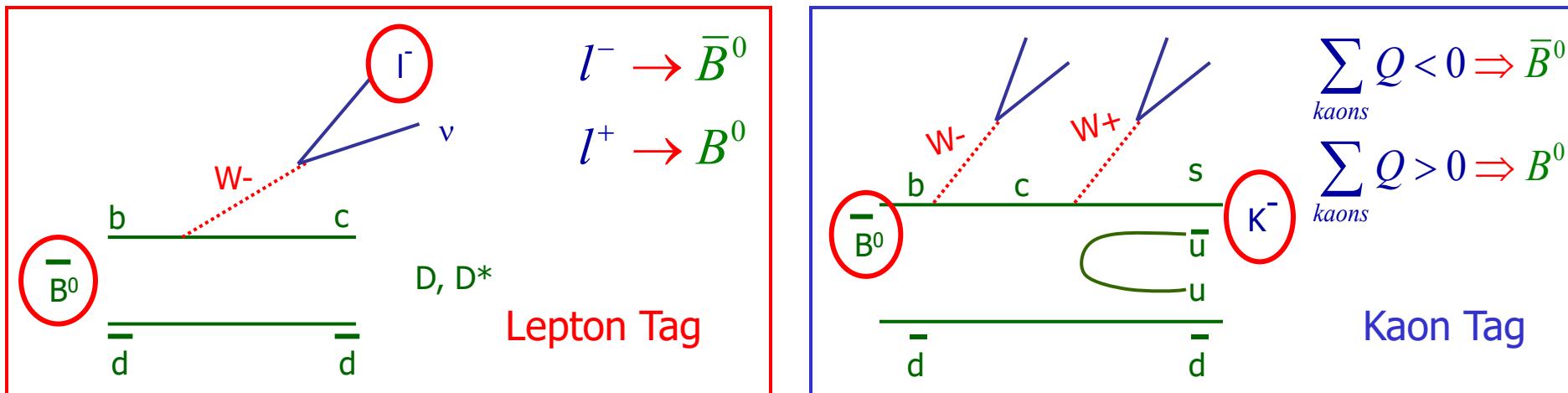


$\Gamma_{Z^0} \propto \# \text{ families of neutrinos.}$

LEP \Rightarrow 3 families of neutrinos.

for neutrinos with $m_\nu \leq \frac{M_Z}{2}$

Separating B^0 and \bar{B}^0 mesons

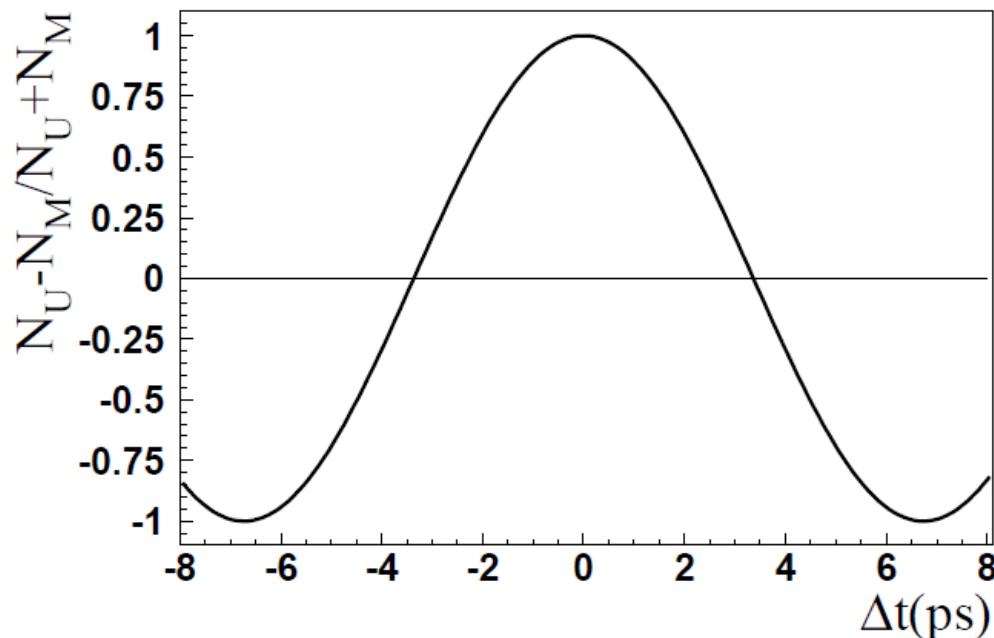
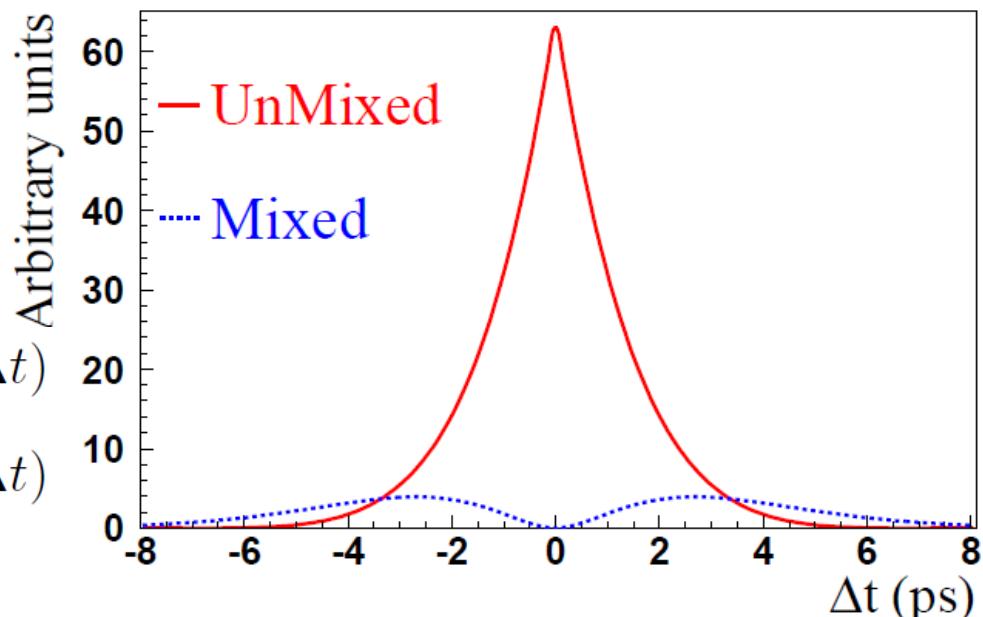


Time Dependent B Oscillation (Or Mixing) at $\Upsilon(4S)$

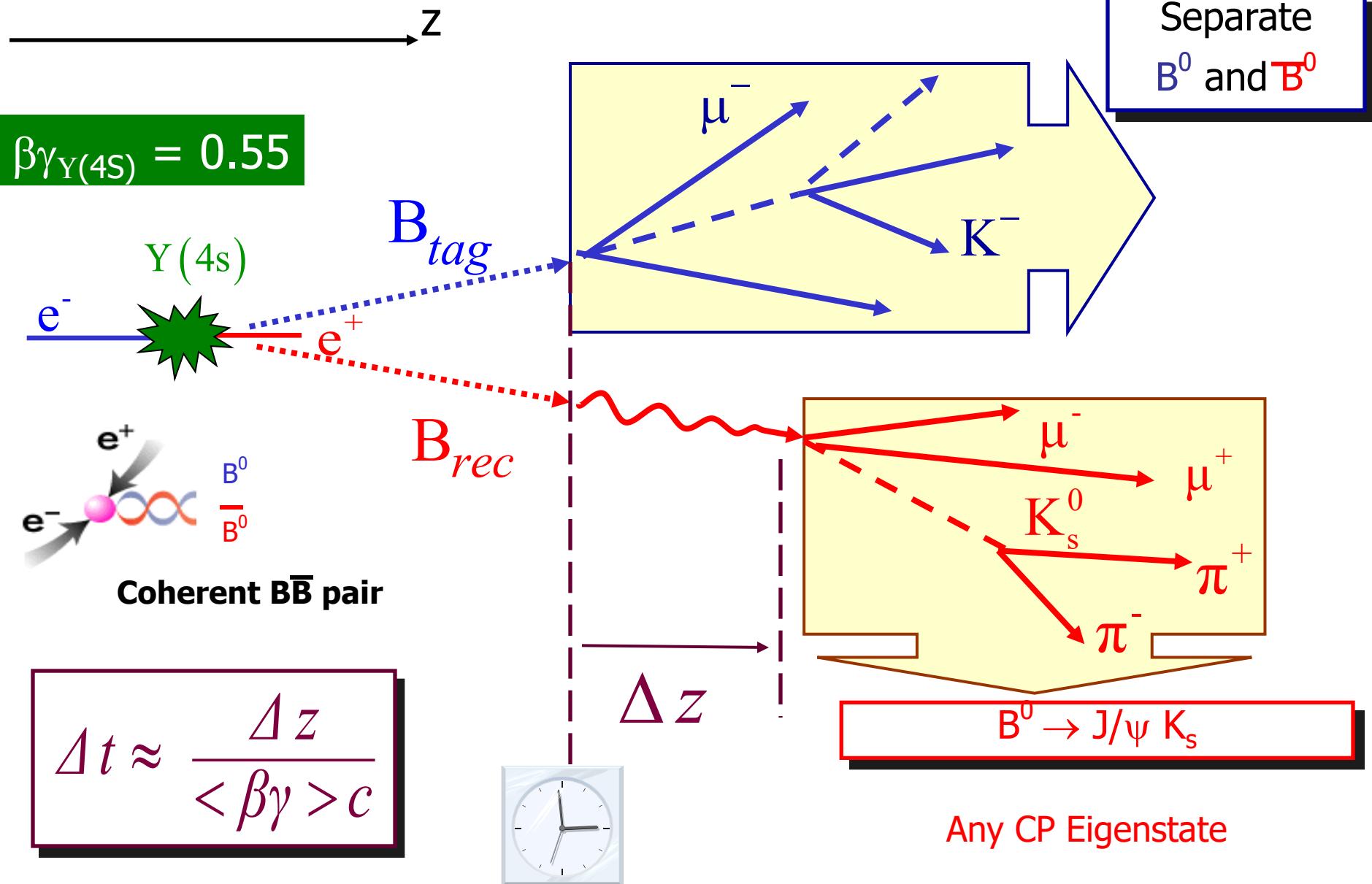
$$f_{\text{unmix}}(\Delta t) \propto e^{-\Gamma|\Delta t|} (1 + \cos \Delta m_d \Delta t)$$

$$f_{\text{mix}}(\Delta t) \propto e^{-\Gamma|\Delta t|} (1 - \cos \Delta m_d \Delta t)$$

$$\mathcal{A}_{\text{mix}}(\Delta t) = \frac{f_{\text{unmix}} - f_{\text{mix}}}{f_{\text{unmix}} + f_{\text{mix}}}$$



CP Violation in Interference between Mixing and Decay



Time-Evolution of B Decays to CP Eigenstates

- Probability of $|B^0\rangle|B^0\rangle \rightarrow |f_{CP}\rangle|f_{tag}\rangle$ depends on
 - Difference Δt between decay time of the two B mesons

- Decay amplitudes

$$A_{f_{CP}} = \langle f_{CP} | H | B^0, t \rangle$$

$$\bar{A}_{f_{CP}} = \langle f_{CP} | H | \bar{B}^0, t \rangle$$

- Oscillation parameter

$$\frac{q}{p} = -\frac{|M_{12}|}{M_{12}} = e^{-i2\beta}$$

- Flavor of tagging neutral B meson: B^0 or $B0$

- Convenient parameter to describe time evolution

- Takes into account combined effect of oscillation and decay

$$\lambda = \frac{q}{p} \frac{\bar{A}_{f_{CP}}}{A_{f_{CP}}} = \eta_{f_{CP}} \frac{q}{p} \frac{\bar{A}_{\bar{f}_{CP}}}{A_{f_{CP}}}$$

Time-Dependent Decay Rates to CP Eigenstates

$$f_{B_{\text{tag}}=B^0}(t_{\text{tag}}, t_{f_{CP}}) \propto e^{-\Gamma(t_{f_{CP}} - t_{\text{tag}})} \left\{ 1 + \frac{1 - |\lambda_{f_{CP}}|^2}{1 + |\lambda_{f_{CP}}|^2} \cos[\Delta m_d(t_{f_{CP}} - t_{\text{tag}})] \right. \\ \left. - \frac{2\mathcal{I}m\lambda_{f_{CP}}}{1 + |\lambda_{f_{CP}}|^2} \sin[\Delta m_d(t_{f_{CP}} - t_{\text{tag}})] \right\}$$

$$f_{B_{\text{tag}}=\bar{B}^0}(t_{\text{tag}}, t_{f_{CP}}) \propto e^{-\Gamma(t_{f_{CP}} - t_{\text{tag}})} \left\{ 1 - \frac{1 - |\lambda_{f_{CP}}|^2}{1 + |\lambda_{f_{CP}}|^2} \cos[\Delta m_d(t_{f_{CP}} - t_{\text{tag}})] \right. \\ \left. + \frac{2\mathcal{I}m\lambda_{f_{CP}}}{1 + |\lambda_{f_{CP}}|^2} \sin[\Delta m_d(t_{f_{CP}} - t_{\text{tag}})] \right\}$$

$$\lambda = \frac{q}{p} \frac{\bar{A}_{f_{CP}}}{A_{f_{CP}}} = \eta_{f_{CP}} \frac{q}{p} \frac{\bar{A}_{\bar{f}_{CP}}}{A_{f_{CP}}}$$

- Expression and complexity of λ depends on specific final states
 - Decay amplitudes A and \bar{A} can be more or less complicated depending on number of amplitudes contributing to total amplitude

Time-Dependent CP Asymmetry in Interference

$$a_{f_{CP}}(\Delta t) = \frac{f_{B_{\text{tag}}=B^0} - f_{B_{\text{tag}}=\bar{B}^0}}{f_{B_{\text{tag}}=B^0} + f_{B_{\text{tag}}=\bar{B}^0}} = \frac{1 - |\lambda_{f_{CP}}|^2}{1 + |\lambda_{f_{CP}}|^2} \cos \Delta m_d \Delta t$$

$$- \frac{2\mathcal{I}m\lambda_{f_{CP}}}{1 + |\lambda_{f_{CP}}|^2} \sin \Delta m_d \Delta t$$

- CP Violation occurs if

$$|\lambda| = \left| \frac{q}{p} \right| \left| \frac{\bar{A}}{A} \right| \neq 1$$

$$\left| \frac{q}{p} \right| = 1$$

No CP Violation
in Mixing

$$\left| \frac{\bar{A}}{A} \right| = 1$$

No Direct
CP Violation

- But even with $|\lambda|=1$ it is sufficient to have $\text{Im}\lambda \neq 0$

In Standard Model we expect $|\lambda| \approx 1$ in most of B decays

Simple Case with $|\lambda_{CP}|=1$

$$\Phi_M = \beta$$

$$A_f = A e^{i(\Phi_W + \delta)}$$

$$\bar{A}_f = \eta_{f_{CP}} A e^{i(-\Phi_W + \delta)}$$

$$\lambda_{f_{CP}} = \eta_{f_{CP}} e^{-2i(\Phi_W - \Phi_M)}$$

$$a_{f_{CP}}(\Delta t) = -\text{Im} \lambda_{f_{CP}} \sin \Delta m_d \Delta t = \eta_{f_{CP}} \sin 2\Phi \sin \Delta m_d \Delta t$$

- Very simple expression for CP violating asymmetry
- Amplitude of asymmetry defined by phase difference between mixing parameter q/p and ratio of decay amplitudes
- Complex phase Φ_M depends on specific final state
 - Can probe different angles of Unitarity triangle through different B decays