

MEASUREMENT OF

COLOR AND JETS

$$R = \frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)} = 3 \cdot \sum_i Q_i^2$$

\uparrow 3 colors!
 \nwarrow available quarks (energetically)

CONVENTIONALLY: RED, GREEN, BLUE
 $r \quad g \quad b$

\Rightarrow three versions of every quark

$u \rightarrow \begin{matrix} u_r \\ u_g \\ u_b \end{matrix}$ three copies
 same mass, charge, spin
 different color
 $\bar{q} \rightarrow \begin{matrix} \bar{q}_r \\ \bar{q}_g \\ \bar{q}_b \end{matrix}$ anticolor for antiquarks

WHAT ABOUT THE PION?

$$|\pi^+\rangle = |u\bar{d}\rangle \xrightarrow{?} |u_r d_{\bar{r}}\rangle + |u_g d_{\bar{g}}\rangle + |u_b d_{\bar{b}}\rangle + \dots$$

YET WE SEE ONLY ONE PION

QCD POSTULATES THAT ONLY COLOR-NEUTRAL ("WHITE")
COMBINATIONS ARE PHYSICAL PARTICLES

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REMEMBER QCD ? THREE BUILDING BLOCKS

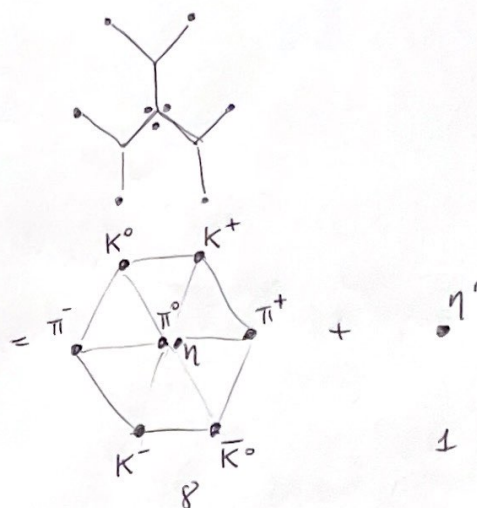
MESONS ARE $q\bar{q} \rightarrow \begin{pmatrix} u \\ d \\ s \end{pmatrix} \times \begin{pmatrix} \bar{u} \\ \bar{d} \\ \bar{s} \end{pmatrix}$

$SU(3) : 3 \otimes \bar{3} = 8 \oplus 1$
↑ ↑
 OCTET SINGLET

3:



$3 \times \bar{3} :$



SAME WITH COLOR

$\Rightarrow 8$ states form an octet \leftarrow these combinations have color $\neq 0$

+ 1 singlet

the singlet is colorless

this is the only color combination physical mesons can have \Rightarrow there is only one pair

DIFFERENCE BETWEEN

and

uds SU(3)

Color SU(3)

[4]

approximate symmetry

$$m_u \neq m_d \neq m_s$$

EXACT symmetry

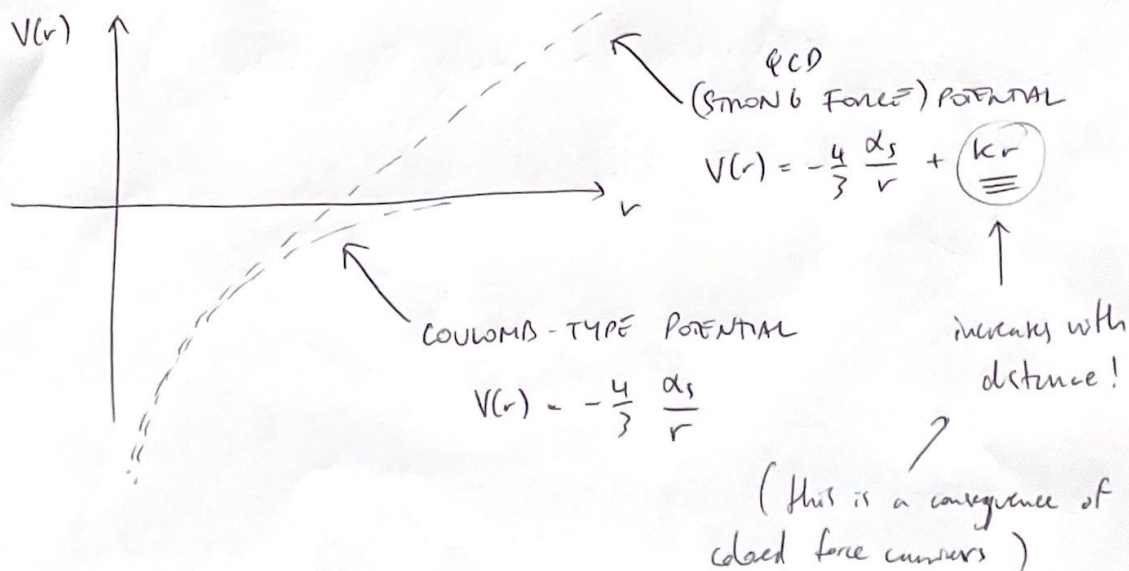
OTHER CONSEQUENCE: if you have a pair say $u\bar{d}$ you CANNOT extract a quark

because it would create two colored states

\Rightarrow quark CONFINEMENT: they exist only

bound in colorless states

by the strong force potential

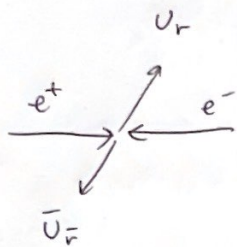


So let's look at ~~the~~

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$$e^+e^- \rightarrow (u\bar{u})$$

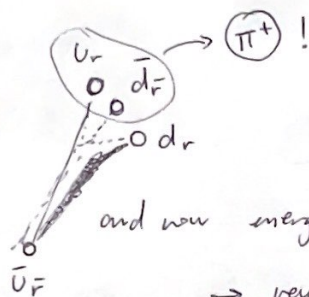
$\hookrightarrow u_r \bar{u}_r$ need a colorless combo



as they drift apart
potential energy between
them increases

\Rightarrow to the point it's energetically favored
to create a new $q\bar{q}$ pair from vacuum

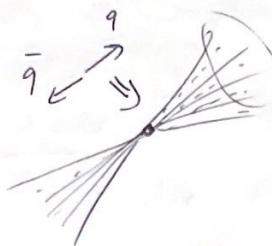
eg.:



and now energy buildup between \bar{u}_r and d_r
 \rightarrow new $q\bar{q}$ pair

and so on .. until all quarks are
"confined" in colorless configurations
 \rightarrow stable hadrons (HADRONIZATION)
(can be lots of them!)

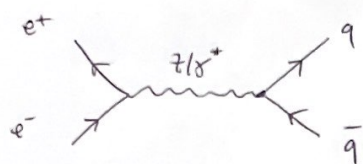
\Rightarrow hadronic JETS



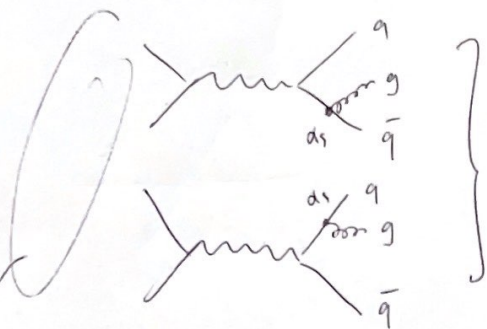
HADRONIZATION

OTHER CONSEQUENCE: three-jet events

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2 jets $e^+e^- \rightarrow q\bar{q}$



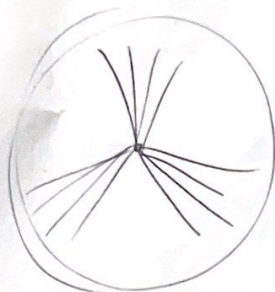
emission of hard gluon FSR

$e^+e^- \rightarrow q\bar{q}g$ 3 jets

FIRST OBSERVATION AT PETRA COLLIDER @ DESY

1981

e^+e^- with $\sqrt{s} > 30$ GeV



3 jet events

$$\frac{\sigma(e^+e^- \rightarrow 3 \text{ jets})}{\sigma(e^+e^- \rightarrow 2 \text{ jets})}$$

$\sim \alpha_s$ (JADE)

$\alpha_s(\sim 35 \text{ GeV}) \sim 0.14$

$\alpha_s(M_Z) \sim 0.11$ (ALEPH)

DIRECT EVIDENCE OF GLUONS EXISTENCE

ALSO: if this model is correct \rightarrow energy accumulated in the plane between q and \bar{q}

\Rightarrow 3rd jet (gluon) should be emitted in same plane as $q\bar{q} \Rightarrow$ 3 jets all on same plane \Rightarrow observed! (TASSO, JADE, etc)

JETS

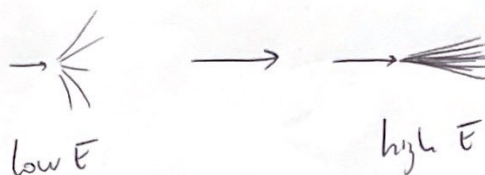
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higher initial energy

⇒ (1) more particles in jet

(2) particles are more collimated



IDEALLY if all particles of a jet are grouped together

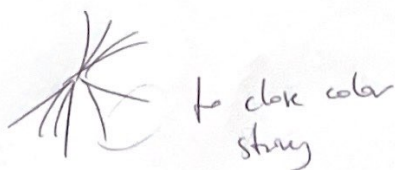
$$\Rightarrow P_{\text{jet}} = \sum_i P_i$$

However jets are never completely independent!

$$e^+e^- \rightarrow (q\bar{q}) \rightarrow 2 \text{ jets}$$

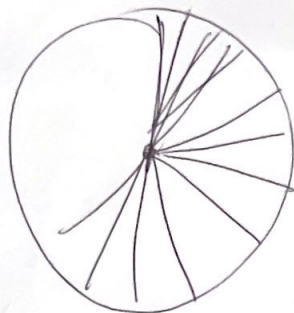
↑
 there is color connection between $q\bar{q}$

⇒ the particle content of 2 jets is correlated



BUT YOU SEE THIS

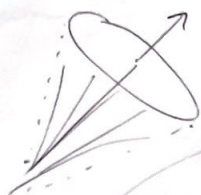
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How MANY JETS?

⇒ 1st PROBLEM: NEED JET CLUSTERING ALGORITHM

SIMPLE EXAMPLE: CMS ITERATIVE CONE



1. OPEN CONE AROUND HIGH ENERGY "SEED"
↳ FIXED RADIUS
2. SUM ALL MOMENTA OF PARTICLES IN CONE
3. COMPUTE NEW AXIS
4. OPEN NEW CONE
5. ITERATE TILL CONE AXIS DOESN'T CHANGE MUCH

How LARGE?

× TOO SMALL → YOU MISS PARTICLES FROM JET
× TOO LARGE → YOU HAVE PROBLEMS SEPARATING
CLOSE-BY JETS
(ALSO YOU COLLECT NOISE IN CONE)

WILL NEVER BE 100% EFFICIENT IN CLUSTERING PARTICLES

⇒ EVEN IF YOU RECONSTRUCT ALL HADRONS
PERFECTLY, SMALL INEFFICIENCY
DUE TO CLUSTERING (LEFT-OUTS)

JET RESPONSE

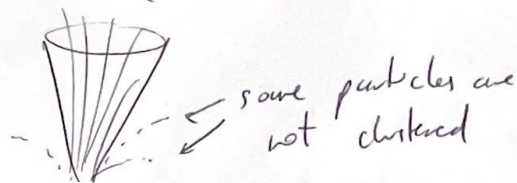
$$R_{jet} = \frac{|\vec{p}(jet, reco)|}{|\vec{p}(parton, true)|} = \frac{\left| \sum_i^{(in\ jet)} \vec{p}_i(hadron_i, reco) \right|}{|\vec{p}(parton, true)|} \quad \boxed{9}$$

↑
quark or gluon that started hadronization

$$= \frac{\left| \sum_{i \in jet} \vec{p}_i(hadron_i, reco) \right|}{\left| \sum_{i \in jet} \vec{p}_i(hadron_i, true) \right|} \cdot \frac{\left| \sum_{i \in jet} \vec{p}_i(hadron_i, true) \right|}{|\vec{p}(parton, true)|}$$

↑
detector response
to hadrons

↑
jet clustering efficiency
($\epsilon \sim 95-99\%$)



AKA JET ENERGY SCALE

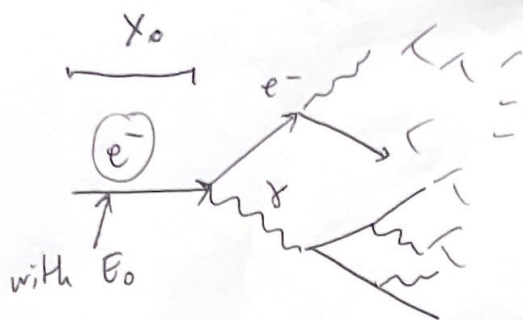
↑
in general: • difficult to get very close to 1
• depends on energy

LET'S SEE WHY

REMEMBER EM SHOWERS

SPALTING CONTINUES
TIL ENOUGH ENERGY

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ALL ENERGY IS
POTENTIALLY DETECTABLE

ALL ENERGY TRANSFORMED INTO LOW-ENERGY e^\pm OR γ
 \Rightarrow IN HOMOGENEOUS ECAL ALL "STORED" IN EXCITATION
 OF ELECTRONS \rightarrow DE-EXCITATION PRODUCES
 SCINTILLATION LIGHT $\rightarrow N_\gamma \propto E_0$

$V E_0$!

\Rightarrow with a properly designed ECAL

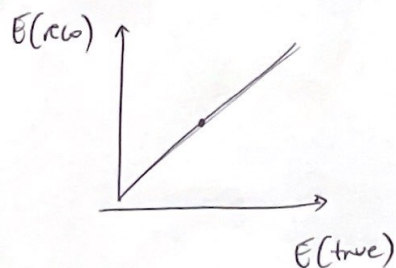
- NO HOLES
- DEEP ENOUGH TO CONTAIN FULL SHOWER



ECAL RESPONSE IS LINEAR

\Rightarrow you need to calibrate it in principle only

ONCE $E(\text{reco}) = \underline{k \cdot N_\gamma}$



WHAT ABOUT HADRONS

WIMBLES COSMOMETRY
SECTION 2.3

III

π^+ entering the detector

charged \rightarrow will ionize (like a μ)

then eventually encounter a nucleus \rightarrow strong int.

dominate effect

many possible outcomes

eg: turn into 15 new hadrons

\uparrow These will continue
and might have other
strong int.

A STRONG. INT.
ON AVERAGE EVENT

λ_{int} (like X_0 for EM)

\rightarrow multiplication

\rightarrow hadronic shower

IN PRINCIPLE CONCEPTUALLY SIMILAR TO EM SHOWERS

BUT: ① INVISIBLE ENERGY

part of the energy of the hadronic shower
is fundamentally undetectable

② particle composition of shower
changes with energy

\Rightarrow hadron calorimeters are intrinsically
non-linear

③ size: $\lambda_{int} > X_0$ eg. For (p) $X_0 = 0.89 \text{ cm}$
 $\lambda_{int} > 20 \text{ cm}$

① INVISIBLE ENERGY

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When a hadron hits a nucleus it can release nucleons from the nucleus

REMEMBER NUCLEAR STARS FROM EMULSION EXPERIMENTS?

Spallation: incoming hadron collides with nucleus, nucleus collides with other nucleus (cascade effect), some pions might be created in the process (if enough E), some nucleons might escape the nucleus

to be able to escape the nucleus (bound state) enough E must be transferred to overcome nuclear binding energy

TYPICALLY
30-40%
OF HADRONIC
INTERACTIONS
ARE INVISIBLE

THIS ENERGY IS LOST FOR CALORIMETRIC PURPOSES \rightarrow INVISIBLE ENERGY

DP CHROMIUM: ULTIMATE ABSORBER

\rightarrow NUCLEAR INTERACTION INDUCED RADIATION

\rightarrow TRY TO RECOVER INVISIBLE ENERGY

OTHER SOURCE OF INVISIBLE ENERGY IN JETS:

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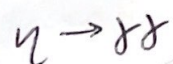
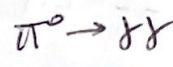


→ SO CONCEPTUALLY IMPOSSIBLE TO DETECT 100% OF ENERGY OF HADRONIC SHOWERS

(2) PARTICLE COMPOSITION

AT EVERY ~~HADRONIC~~ ^{NUCLEAR} INTERACTION → NEW HADRONS PRODUCED

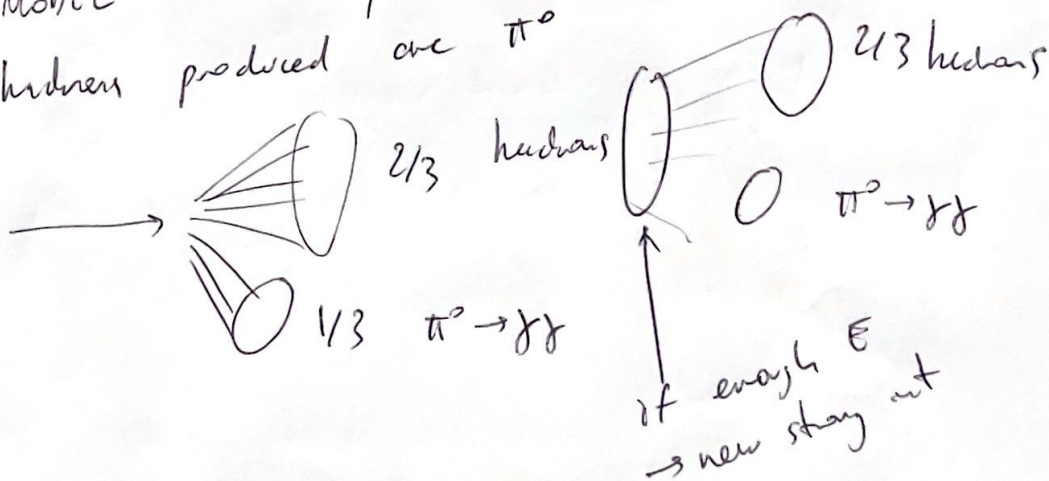
Some of these decay to photons:



→ hadronic showers have EM component

(EM component has ~~no~~ ^{NO} invisible E)

SIMPLE MODEL: at every interaction $\frac{1}{3}$ of the hadrons produced are π^0



So at every step $\frac{1}{3}$ of E

is STORED IN PHOTONS → IRREVERSIBLE PROCESS

"A ONE-WAY STREET"

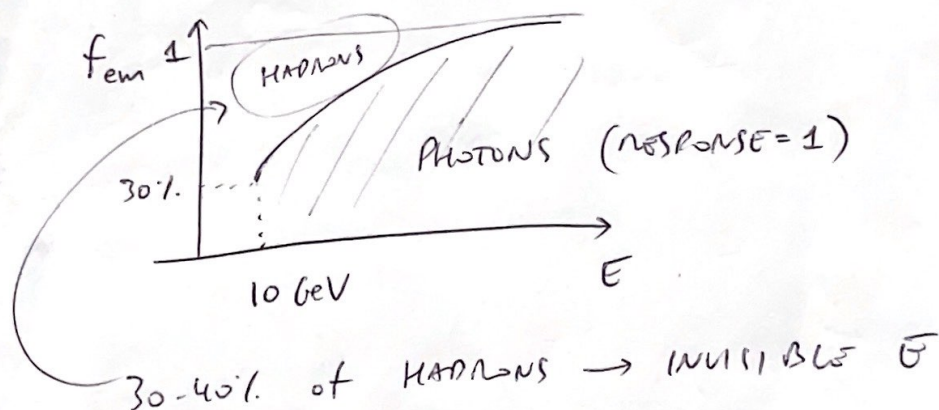
in this simple model at any step for 14
 which there's still enough E to produce
 new hadrons $\rightarrow \frac{1}{3}$ of energy goes to EM sector

$$\Rightarrow f_{em} = 1 - \left(1 - \frac{1}{3}\right)^n$$

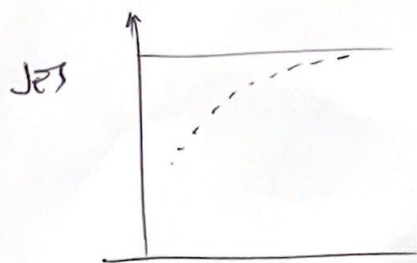
number of steps
 $n \propto E_0$

THIS IS A SIMPLE MODEL
 in reality $f_{90} < \frac{1}{3}$ (other particles can be produced)

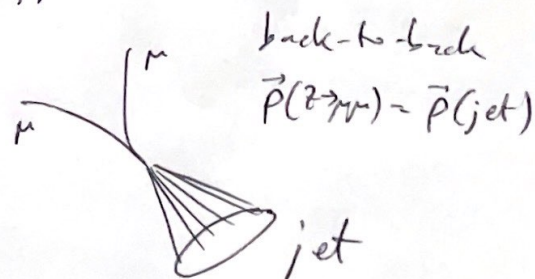
BUT NOT TOO FAR OFF



\Rightarrow JET ENERGY SCALE: NEED TO CALIBRATE AT ALL ENERGIES



$Z \rightarrow \mu\mu + \text{jet}$ events



FLUCTUATIONS! JESR \sim 5-10% AT BEST (e.g. $\text{CMS} < 1\%$ @ CM)