

Angular Momentum of a System of N particles

consider a system of N particles $\alpha=1, 2, \dots, N$ with $\vec{l}_\alpha = \vec{r}_\alpha \times \vec{p}_\alpha$



The total \vec{L} is
$$\vec{L} = \sum_{\alpha=1}^N \vec{l}_\alpha = \sum_{\alpha=1}^N \vec{r}_\alpha \times \vec{p}_\alpha$$

How does \vec{L} change with time?

$$\dot{\vec{L}} = \sum_{\alpha} \dot{\vec{l}}_{\alpha} = \sum_{\alpha} \vec{r}_{\alpha} \times \vec{F}_{\alpha}$$
 Remember, $\vec{F}_{\alpha} = \sum_{\beta \neq \alpha} \vec{F}_{\alpha\beta} + \vec{F}_{\alpha}^{ext}$

$$\dot{\vec{L}} = \sum_{\alpha} \vec{r}_{\alpha} \times \left[\sum_{\beta \neq \alpha} \vec{F}_{\alpha\beta} + \vec{F}_{\alpha}^{ext} \right] = \sum_{\beta \neq \alpha} \vec{r}_{\alpha} \times \vec{F}_{\alpha\beta} + \sum_{\alpha} \vec{r}_{\alpha} \times \vec{F}_{\alpha}^{ext}$$

We can use the same trick as before: (reorganize the sum)

$$\sum_{\alpha \neq \beta} \vec{r}_{\alpha} \times \vec{F}_{\alpha\beta} = \sum_{\alpha} \left(\sum_{\beta < \alpha} \vec{r}_{\alpha} \times \vec{F}_{\alpha\beta} + \sum_{\beta > \alpha} \vec{r}_{\alpha} \times \vec{F}_{\alpha\beta} \right)$$

$\vec{F}_{\alpha\beta} = -\vec{F}_{\beta\alpha}$

$$= \sum_{\alpha} \sum_{\beta < \alpha} \vec{r}_{\alpha} \times \vec{F}_{\alpha\beta} - \sum_{\alpha} \sum_{\beta > \alpha} \vec{r}_{\alpha} \times \vec{F}_{\beta\alpha}$$

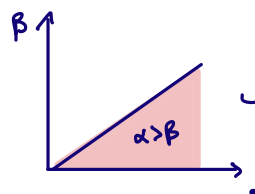
Next, reindex:

$$\begin{aligned} \alpha &\rightarrow \beta \\ \beta &\rightarrow \alpha \end{aligned}$$

$$= \sum_{\alpha} \sum_{\beta < \alpha} \vec{r}_{\alpha} \times \vec{F}_{\alpha\beta} - \sum_{\beta} \sum_{\alpha > \beta} \vec{r}_{\alpha} \times \vec{F}_{\alpha\beta}$$

Note:
$$\sum_{\beta} \sum_{\alpha > \beta} f(\alpha, \beta) = \sum_{\alpha} \sum_{\beta < \alpha} f(\alpha, \beta)$$

These sums are the same



This region ($\alpha > \beta$) is identical to the $\beta < \alpha$ region

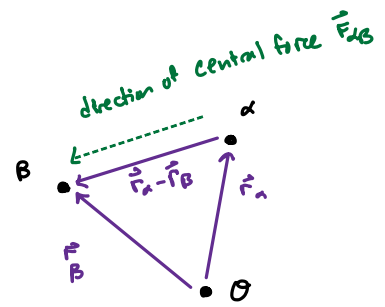
$$= \sum_{\alpha} \sum_{\beta < \alpha} \vec{r}_{\alpha} \times \vec{F}_{\alpha\beta} - \sum_{\alpha} \sum_{\beta < \alpha} \vec{r}_{\beta} \times \vec{F}_{\alpha\beta}$$

$$= \sum_{\alpha} \sum_{\beta < \alpha} \left(\vec{r}_{\alpha} \times \vec{F}_{\alpha\beta} - \vec{r}_{\beta} \times \vec{F}_{\alpha\beta} \right)$$

$$= \sum_{\alpha} \sum_{\beta < \alpha} (\vec{r}_{\alpha} - \vec{r}_{\beta}) \times \vec{F}_{\alpha\beta}$$

$$\Rightarrow \dot{\vec{L}} = \sum_{\alpha} (\vec{r}_{\alpha} - \vec{r}_{\beta}) \times \vec{F}_{\alpha\beta} + \sum_{\alpha} \vec{r}_{\alpha} \times \vec{F}_{\alpha}^{ext}$$

$$\dot{\vec{L}} = \sum_{\alpha \neq \beta} (\vec{r}_\alpha - \vec{r}_\beta) \times \vec{F}_{\alpha\beta} + \vec{r}_\alpha \times \vec{F}_\alpha^{\text{ext}}$$



If all $\vec{F}_{\alpha\beta}$ are central forces, then $\vec{F}_{\alpha\beta} \parallel \vec{r}_\alpha - \vec{r}_\beta \Rightarrow$

$$\Rightarrow (\vec{r}_\alpha \times \vec{r}_\beta) \times \vec{F}_{\alpha\beta} = 0$$

$$\Rightarrow \dot{\vec{L}} = \vec{r}_\alpha \times \vec{F}_\alpha^{\text{ext}} = \vec{r}^{\text{ext}}$$

Conservation of angular momentum:

If a net external torque on an N-particle system = 0, then the total $\dot{\vec{L}} = \vec{r}_\alpha \times \vec{F}_\alpha^{\text{ext}} = \text{constant}$
(Assuming all internal forces are central forces, which is almost always true)

Moment of Inertia:

When a rigid body rotates on an axis (let's call it the z-axis), we can calculate L_z using the moment of inertia:

$$L_z = I\omega$$

$$\uparrow \\ \omega = \dot{\phi}$$

$$I = M\alpha^2$$

This is derived from $\vec{L} = \vec{r} \times \vec{p}$ in Ch.10

$$\hookrightarrow \int r^2 dm = \int r^2 \rho dV \quad \text{for solid objects}$$

Angular momentum about CM:

In an inertial frame, $\dot{\vec{L}} = \vec{r}^{\text{ext}}$

In any frame, $\dot{\vec{L}}(\text{about CM}) = \vec{r}^{\text{ext}}(\text{about CM})$, even when CM is accelerating!
(Proof in Ch.10)

Chapter 4: Kinetic Energy and Work

Kinetic energy: $T \equiv \frac{1}{2}mv^2 \rightarrow \text{useful! But why?}$

Let's take its derivative: $\dot{T} = \frac{1}{2}m \frac{d}{dt}(\vec{v} \cdot \vec{v})$

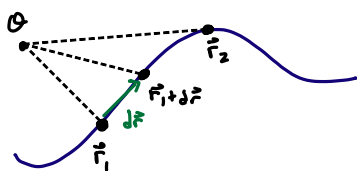
$$= \frac{1}{2}m (\dot{\vec{v}} \cdot \vec{v} + \vec{v} \cdot \dot{\vec{v}}) = m \vec{v} \cdot \dot{\vec{v}} = m \vec{a} \cdot \vec{v}$$

$$\dot{T} = \vec{F} \cdot \vec{v} = \vec{F} \cdot \dot{\vec{r}}$$

$$\frac{dT}{dt} = \vec{F} \cdot \frac{d\vec{r}}{dt}$$

$$dT = \vec{F} \cdot d\vec{r}$$

Change in the system: $dT = \vec{F} \cdot d\vec{r} = dW$: Work done by \vec{F} over a distance $d\vec{r}$

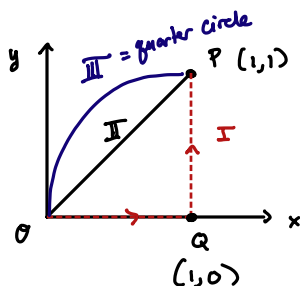


$$\Delta T \equiv T_2 - T_1 = \sum \vec{F} \cdot d\vec{r} \rightarrow \int_1^2 \vec{F} \cdot d\vec{r} \quad (\text{line integral})$$

$$\Delta T = T_2 - T_1 = \int_1^2 \vec{F} \cdot d\vec{r} \equiv W(1 \rightarrow 2) \quad \text{The Work-KE theorem}$$

↑
This is the
total force

Practice with line integrals:



$$\vec{F} = y\hat{x} + 2x\hat{y}$$

$$W_I = \int \vec{F} \cdot d\vec{r} = \int_O^Q \vec{F} \cdot d\vec{r} + \int_Q^P \vec{F} \cdot d\vec{r}$$

$$d\vec{r} = dx\hat{x} + dy\hat{y}$$

From $O \rightarrow Q$, $dy = 0 \Rightarrow d\vec{r} = dx\hat{x} + \cancel{dy\hat{y}} = dx\hat{x}$

From $Q \rightarrow P$, $dx = 0 \Rightarrow d\vec{r} = \cancel{dx\hat{x}} + dy\hat{y} = dy\hat{y}$

$$W_I = \int_{O_x}^{Q_x} \underbrace{\vec{F} \cdot \hat{x}}_{=F_x} dx + \int_{Q_y}^{P_y} \underbrace{\vec{F} \cdot \hat{y}}_{=F_y} dy$$

$$W_I = \int_0^1 \underbrace{y}_{y=0 \text{ between } O \& Q} dx + \int_0^1 \underbrace{2x}_{x=1 \text{ between } Q \& P} dy = \int_0^1 \cancel{(0)} dx + \int_0^1 2(1) dy$$

$$W_I = \int_0^1 2 dy = 2y \Big|_0^1 = \boxed{2 = W_I}$$

Next, calculate W_{II} :

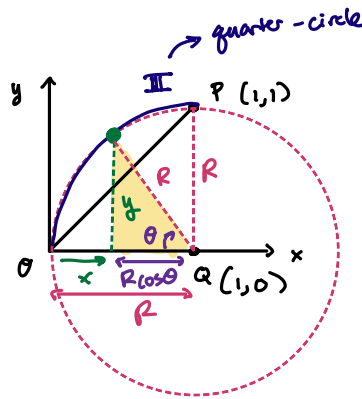
$$W_{II} = \int_{II} \vec{F} \cdot d\vec{r} \quad \text{on path II: } x=y \Rightarrow dx=dy$$

$$W_{II} = \int_I \vec{F} \cdot (dx\hat{x} + dy\hat{y}) = \int_{II} (y dx + 2x dy) = \int_{\theta_x}^{P_x} (x dx + 2x dx) = \int_0^1 3x dx = \left. \frac{3}{2} x^2 \right|_0^1 = \boxed{\frac{3}{2}} = W_{II}$$

use $x=y$
 $dx=dy$

W_{III} : On path III,

$$\begin{aligned} W_{III} &= \int_{III} \vec{F} \cdot d\vec{r} \\ &= \int_{III} F_x dx + F_y dy \\ &= \int_{III} y dx + 2x dy \end{aligned}$$



At point P, $\theta = \pi/2$

At θ , $\theta = 0$

$$\Rightarrow y = R \sin \theta \quad R = 1$$

$$\begin{aligned} y &= \sin \theta \\ dy &= \cos \theta d\theta \end{aligned}$$

$$\Rightarrow x + R \cos \theta = R \quad R = 1$$

$$x + \cos \theta = 1$$

$$dx - \sin \theta d\theta = 0$$

$$dx = \sin \theta d\theta$$

$$= \int_{\theta_0}^{P_\theta} \left[(\sin \theta) (\sin \theta d\theta) + 2(1 - \cos \theta) (\cos \theta d\theta) \right]$$

$$= \int_0^{\pi/2} (\sin^2 \theta - 2\cos \theta - 2\cos^2 \theta) d\theta = 2 \int_0^{\pi/2} \cos \theta d\theta + \int_0^{\pi/2} [(\sin^2 \theta - \cos^2 \theta) - \cos^2 \theta] d\theta$$

$$= 2 \sin \theta \Big|_0^{\pi/2} + \int_0^{\pi/2} \left[(-\cos 2\theta) - \frac{1}{2} (\cos 2\theta + 1) \right] d\theta$$

$$= 2 \sin \frac{\pi}{2} - 0 + \int_0^{\pi/2} \left(-\frac{3}{2} \cos 2\theta - \frac{1}{2} \right) d\theta$$

$$= 2(1) - \frac{3}{2} \frac{1}{2} \sin 2\theta \Big|_0^{\pi/2} - \frac{1}{2} \theta \Big|_0^{\pi/2}$$

$$= 2 - \frac{3}{4} (\sin \pi - 0) - \frac{\pi}{4} + 0$$

$$\boxed{W_{III} = 2 - \pi/4}$$