



Signals & Systems (ECN-203)

Lecture 6

(Fourier series representation of periodic signals)

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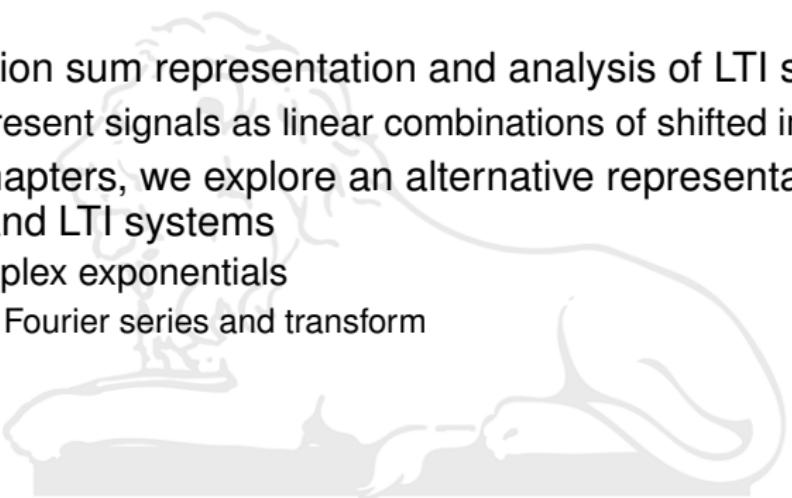


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Introduction



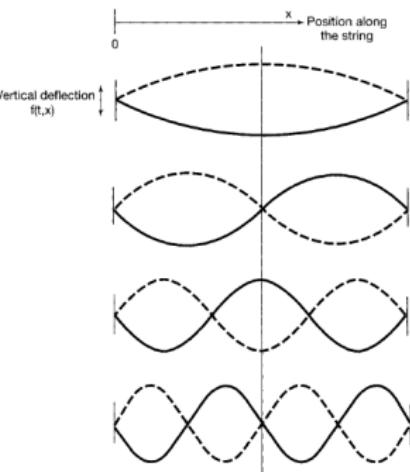
- ❑ Convolution sum representation and analysis of LTI systems:
 - ❑ Represent signals as linear combinations of shifted impulses
- ❑ In this chapters, we explore an alternative representation for signals and LTI systems
 - ❑ Complex exponentials
 - ❑ Fourier series and transform



A historical perspective



- ❑ Development of Fourier analysis has a long history involving several individuals and study of various physical phenomena
- ❑ Babylonians use “trigonometric sums” to describe periodic phenomena in order to predict astronomical events
 - ❑ 2000-1600 BC
- ❑ Modern history: L. Euler examined the motion of a vibrating string in 1748
- ❑ The “normal modes” are harmonically related sinusoidal functions of x
- ❑ Euler predicted configuration of a vibrating string as a linear combination of these normal modes



A historical perspective



- ❑ This property isn't particularly useful unless it is generalized
 - ❑ i.e., unless a large class of interesting functions could be represented by linear combinations of complex exponentials
- ❑ Bernoulli supported this on physical grounds
 - ❑ But did not pursue this mathematically
- ❑ Lagrange strongly criticized this
 - ❑ Based on his belief that it was impossible to represent signals with corners (discontinuous slopes) using trigonometric series
- ❑ Fourier presented his ideas in 1768
 - ❑ Based on the properties of heat propagation and diffusion
- ❑ Fourier took this type of representation one very large step farther:
 - ❑ Represented aperiodic signals as weighted integrals of sinusoids that are *not* all harmonically related



Applications of Fourier's work



- ❑ Fourier's work has enormous impact on so many disciplines within the fields of mathematics, science, and engineering
 - ❑ Theory of integration
 - ❑ Point-set topology
 - ❑ Eigen function expansions
- ❑ Application areas:
 - ❑ Motion of the planets and the periodic behavior of the earth's climate
 - ❑ Alternating-current sources
 - ❑ Waves in the ocean
 - ❑ Signals transmitted by radio and television stations

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Why represent a signal as a sum of other signals



- ❑ Advantageous to represent signals as linear combinations of basic signals that possess the following two properties:
 - ❑ The set of basic signals can be used to construct a broad and useful class of signals
 - ❑ Response of an LTI system to each signal should be simple
 - ❑ Convenient representation for the response of the system to any signal constructed as a linear combination of the basic signals
- ❑ Both of these properties are fulfilled by the set of complex exponential signals in continuous and discrete time
 - ❑ e^{st} in continuous time and z^n in discrete time
 - ❑ s and z are complex numbers

The second property



- ❑ Response of an LTI system to a complex exponential input is the same complex exponential with only a change in amplitude
 - ❑ $e^{st} \rightarrow H(s)e^{st}$
 - ❑ $z^n \rightarrow H(z)z^n$
 - ❑ $H(s)$ and $H(z)$ are complex functions of s and z respectively
- ❑ This is the Eigenvalue-Eigenfunction problem
 - ❑ e^{st} and z^n are Eigenfunctions
 - ❑ $H(s)$ and $H(z)$ are Eigenvalues

Complex exponential as an input to continuous-time LTI system



- ❑ Impulse response: $h(t)$
- ❑ Input: $x(t) = e^{st}$
- ❑ Output:

$$\begin{aligned}y(t) &= \int_{-\infty}^{\infty} h(\tau)x(t-\tau)d\tau \\&= \int_{-\infty}^{\infty} h(\tau)e^{s(t-\tau)}d\tau \\&= e^{st} \int_{-\infty}^{\infty} h(\tau)e^{-s\tau}d\tau \\&= H(s)e^{st}\end{aligned}$$

- ❑ $H(s) = \int_{-\infty}^{\infty} h(\tau)e^{-s\tau}d\tau$
- ❑ Assumption: This integral converges

Complex exponential as an input to discrete-time LTI system



- Impulse response: $h[n]$
- Input: $x[n] = z^n$
- Output:

$$\begin{aligned}y[n] &= \sum_{k=-\infty}^{\infty} h[k]x[n-k] \\&= \sum_{k=-\infty}^{\infty} h[k]z^{n-k} \\&= z^n \sum_{k=-\infty}^{\infty} h[k]z^{-k} \\&= H(z)z^n\end{aligned}$$

- $H(z) = \sum_{k=-\infty}^{\infty} h[k]z^{-k}$
- Assumption: This sum converges

Sum of complex exponentials as an input to a LTI system



- ❑ Let the input to a continuous-time LTI system is represented as a linear combination of complex exponentials
 - ❑ $x(t) = \sum_k a_k e^{s_k t}$
- ❑ Output: $y(t) = \sum_k a_k H(s_k) e^{s_k t}$
- ❑ If the input to a discrete-time LTI system is represented as a linear combination of complex exponentials
 - ❑ $x[n] = \sum_k a_k z_k^n$
- ❑ Output: $y[n] = \sum_k a_k H(z_k) z_k^n$

Example



- ❑ Consider a time-shift LTI system
 - ❑ $y(t) = x(t - 3)$
- ❑ $x(t) = e^{j2t} \rightarrow y(t) = e^{j2(t-3)} = e^{-6j} e^{j2t}$
 - ❑ Eigen function = e^{j2t}
 - ❑ Eigen value, $H(2j) = e^{-6j}$
- ❑ Impulse response, $h(t) = \delta(t - 3) \rightarrow$
 $H(s) = \int_{-\infty}^{\infty} h(\tau) e^{-s\tau} d\tau = \int_{-\infty}^{\infty} \delta(\tau - 3) e^{-s\tau} d\tau = e^{-3s}$
- ❑ $x(t) = \cos(4t) + \cos(7t) = \frac{1}{2}e^{j4t} + \frac{1}{2}e^{-j4t} + \frac{1}{2}e^{j7t} + \frac{1}{2}e^{-j7t}$
- ❑

$$\begin{aligned}y(t) &= \frac{1}{2}e^{-3 \times j4} e^{j4t} + \frac{1}{2}e^{-3 \times -j4} e^{-j4t} + \frac{1}{2}e^{-3 \times j7} e^{j7t} + \frac{1}{2}e^{-3 \times -j7} e^{-j7t} \\&= \frac{1}{2}e^{j4(t-3)} + \frac{1}{2}e^{-j4(t-3)} + \frac{1}{2}e^{j4(t-7)} + \frac{1}{2}e^{-j4(t-7)} \\&= \cos(4(t - 3)) + \cos(7(t - 3))\end{aligned}$$

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Linear combinations of harmonically related complex exponentials



- ❑ A signal $x(t)$ is periodic with period T if:
 - ❑ $x(t + T) = x(t), \forall t$
 - ❑ $\omega_0 = \frac{2\pi}{T}$ is the fundamental frequency
- ❑ Complex exponential $e^{j\omega_0 t}$ is periodic with fundamental frequency ω_0 and fundamental period $T = \frac{2\pi}{\omega_0}$
- ❑ Harmonically related complex exponentials of $e^{j\omega_0 t}$ are
$$\phi_k(t) = e^{jk\omega_0 t} = e^{jk\frac{2\pi}{T}t}, k = 0, \pm 1, \pm 2, \dots$$
 - ❑ Fundamental frequency of these exponentials is a multiple of ω_0
 - ❑ Each is periodic with period T (although not the fundamental period)
 - ❑ $k = \pm N \rightarrow N^{\text{th}}$ harmonic components
- ❑ Linear combination of harmonically related complex exponentials:
$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} = \sum_{k=-\infty}^{\infty} a_k e^{jk\frac{2\pi}{T}t}$$
 - ❑ Periodic with period T
 - ❑ This is the Fourier series representation

Example

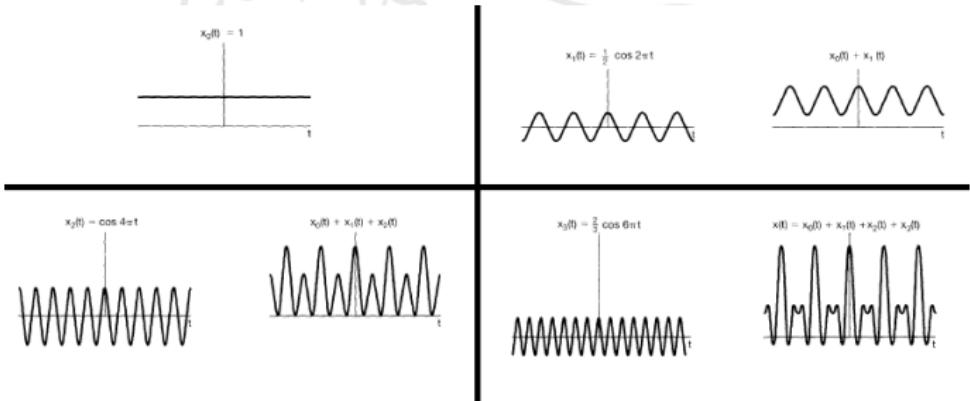


□ $x(t) = \sum_{k=-3}^3 a_k e^{jk2\pi t}$

□ $a_0 = 1, a_1 = a_{-1} = \frac{1}{4}, a_2 = a_{-2} = \frac{1}{2}, a_3 = a_{-3} = \frac{1}{3}$

□

$$\begin{aligned}x(t) &= 1 + \frac{1}{4}(e^{j2\pi t} + e^{-j2\pi t}) + \frac{1}{2}(e^{j4\pi t} + e^{-j4\pi t}) + \frac{1}{3}(e^{j6\pi t} + e^{-j6\pi t}) \\&= 1 + \frac{1}{2}\cos(2\pi t) + \cos(4\pi t) + \frac{2}{3}\cos(6\pi t)\end{aligned}$$



Fourier series of real periodic signals



- Let signal $x(t)$ is real and periodic with period T

- $$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t}$$

- $$x(t) = x^*(t) = \sum_{k=-\infty}^{\infty} a_k^* e^{-jk\omega_0 t}$$

- Substitute $k = -k$

- $$x^*(t) = \sum_{k=-\infty}^{\infty} a_{-k}^* e^{jk\omega_0 t} = x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t}$$

- $$a_{-k}^* = a_k \rightarrow a_{-k} = a_k^*$$

-

$$\begin{aligned} x(t) &= \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} = a_0 + \sum_{k=1}^{\infty} (a_k e^{jk\omega_0 t} + a_{-k} e^{-jk\omega_0 t}) \\ &= a_0 + \sum_{k=1}^{\infty} (a_k e^{jk\omega_0 t} + a_k^* e^{-jk\omega_0 t}) = a_0 + \sum_{k=1}^{\infty} 2\Re(a_k e^{jk\omega_0 t}) \end{aligned}$$

- $a_k = A_k e^{j\theta_k}$

- $$x(t) = a_0 + \sum_{k=1}^{\infty} 2\Re(A_k e^{j\theta_k} e^{jk\omega_0 t}) = a_0 + \sum_{k=1}^{\infty} 2A_k \cos(k\omega_0 t + \theta_k)$$

Determining Fourier series coefficients



- $x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t}$
- $x(t)e^{-jn\omega_0 t} = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} e^{-jn\omega_0 t}$
-

$$\begin{aligned}\int_0^T x(t)e^{-jn\omega_0 t} dt &= \int_0^T \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} e^{-jn\omega_0 t} dt \\&= \sum_{k=-\infty}^{\infty} a_k \int_0^T e^{j(k-n)\omega_0 t} dt \\&= \sum_{k=-\infty}^{\infty} a_k \left[\int_0^T (\cos((k-n)\omega_0 t) dt \right. \\&\quad \left. + j \int_0^T \sin((k-n)\omega_0 t)) dt \right]\end{aligned}$$

Determining Fourier series coefficients



- If $k \neq n$
 - $\cos((k-n)\omega_0 t)$ and $\sin((k-n)\omega_0 t)$ are periodic sinusoids with fundamental period $\frac{T}{|k-n|}$
 - Integral from 0 to T consists of integer number of periods
 - $\int_0^T (\cos((k-n)\omega_0 t) dt = \int_0^T \sin((k-n)\omega_0 t) dt = 0$
- If $k = n$
 - $\int_0^T (\cos((k-n)\omega_0 t) dt = \int_0^T 1 dt = T$
 - $\int_0^T (\sin((k-n)\omega_0 t) dt = \int_0^T 0 dt = 0$
- $$\int_0^T e^{j(k-n)\omega_0 t} dt = \begin{cases} T, & k = n \\ 0, & k \neq n \end{cases}$$
- $a_n = \frac{1}{T} \int_0^T x(t) e^{-jn\omega_0 t} dt$

Fourier series summary



- ❑ Synthesis equation

- ❑ $x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} = \sum_{k=-\infty}^{\infty} a_k e^{jk \frac{2\pi}{T} t}$

- ❑ Analysis equation:

- ❑ $a_k = \frac{1}{T} \int_T x(t) e^{-jk\omega_0 t} dt = \frac{1}{T} \int_T x(t) e^{-jk \frac{2\pi}{T} t} dt$

Example 1



□ $x(t) = \sin(\omega_0 t)$

□ $x(t) = \frac{1}{2j} e^{j\omega_0 t} - \frac{1}{2j} e^{-j\omega_0 t}$

□

$$a_k = \begin{cases} \frac{1}{2j}, & k = 1 \\ -\frac{1}{2j}, & k = -1 \\ 0, & \text{Otherwise} \end{cases}$$

Example 2



□ $x(t) = 1 + \sin(\omega_0 t) + 2\cos(\omega_0 t) + \cos(2\omega_0 t + \frac{\pi}{4})$

□

$$x(t) = 1 + \frac{1}{2j}(e^{j\omega_0 t} - e^{-j\omega_0 t}) + (e^{j\omega_0 t} + e^{-j\omega_0 t}) \\ + \frac{1}{2j}(e^{j(2\omega_0 t + \frac{\pi}{4})} + e^{-j(2\omega_0 t + \frac{\pi}{4})})$$

□ $x(t) = 1 + (1 + \frac{1}{2j})e^{j\omega_0 t} + (1 - \frac{1}{2j})e^{-j\omega_0 t} + \frac{1}{2j}e^{j\frac{\pi}{4}}e^{j(2\omega_0 t)} + \frac{1}{2j}e^{-j\frac{\pi}{4}}e^{-j(2\omega_0 t)}$

□

$$a_k = \begin{cases} 1, & k = 0 \\ 1 + \frac{1}{2j} = 1 - \frac{1}{2}j, & k = 1 \\ 1 - \frac{1}{2j} = 1 + \frac{1}{2}j, & k = -1 \\ \frac{1}{2j}e^{j\frac{\pi}{4}} = \frac{1}{2\sqrt{2}}(1 + j), & k = 2 \\ \frac{1}{2j}e^{-j\frac{\pi}{4}} = \frac{1}{2\sqrt{2}}(1 - j), & k = -2 \\ 0, & \text{Otherwise} \end{cases}$$

Example 2

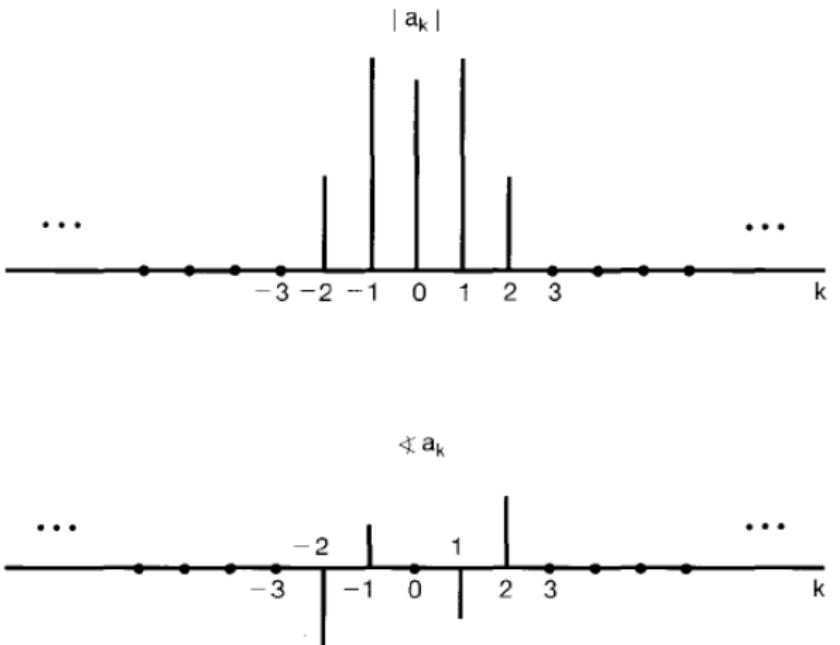


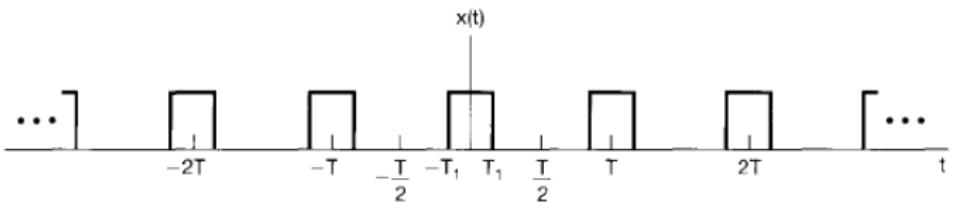
Figure 3.5 Plots of the magnitude and phase of the Fourier coefficients of the signal considered in Example 3.4.

Example 3



- ❑ $x(t)$ is a periodic square wave defined over one period as

$$x(t) = \begin{cases} 1, & |t| < T_1 \\ 0, & T_1 < |t| < \frac{T}{2} \end{cases}$$



- ❑ $a_0 = \frac{1}{T} \int_{-T_1}^{T_1} 1 dt = \frac{2T_1}{T}$

- ❑ $a_k = \frac{1}{T} \int_{-T_1}^{T_1} e^{-jk\omega_0 t} dt = \frac{-1}{jk\omega_0 T} \left(e^{jk\omega_0 T_1} - e^{-jk\omega_0 T_1} \right) = \frac{\sin(k\omega_0 T_1)}{k\pi}$

Example 3



- $T = 4T_1 \rightarrow x(t)$ is a square wave that is unity for half the period and zero for half the period

-

$$a_k = \begin{cases} \frac{1}{2}, & k = 0 \\ \frac{1}{\pi}, & k = 1, -1 \\ \frac{-1}{3\pi}, & k = 3, -3 \\ \frac{1}{5\pi}, & k = 5, -5 \\ \dots \\ 0, & \text{Otherwise} \end{cases}$$

Example 3

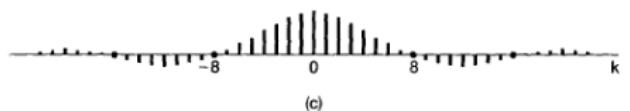
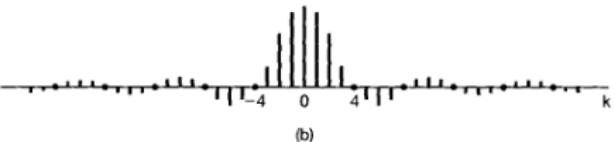
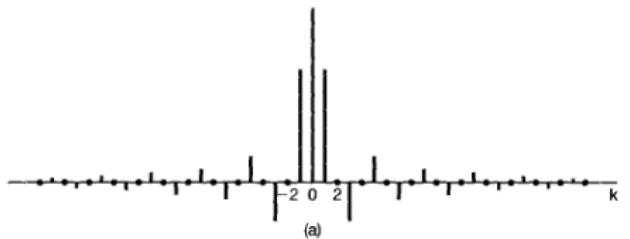


Figure 3.7 Plots of the scaled Fourier series coefficients $T a_k$ for the periodic square wave with T_1 fixed and for several values of T : (a) $T = 4T_1$; (b) $T = 8T_1$; (c) $T = 16T_1$. The coefficients are regularly spaced samples of the envelope $(2 \sin \omega T_1)/\omega$, where the spacing between samples, $2\pi/T$, decreases as T increases.



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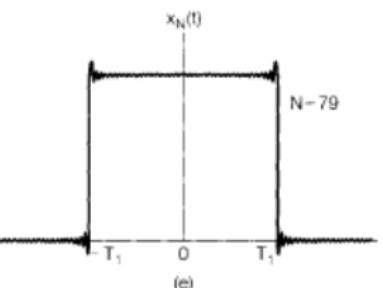
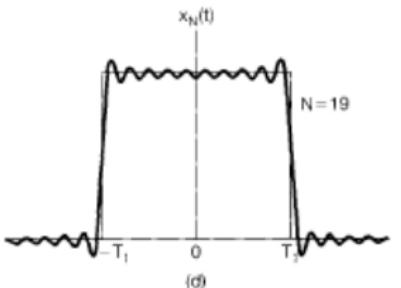
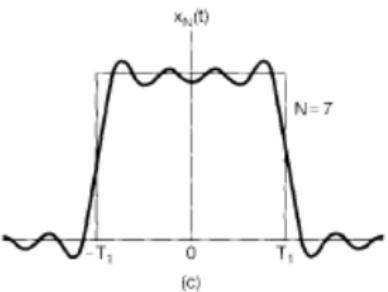
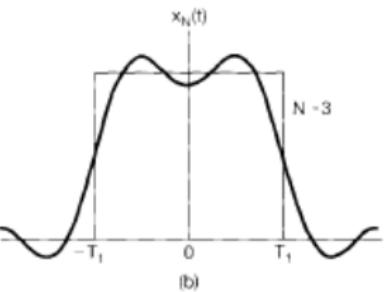
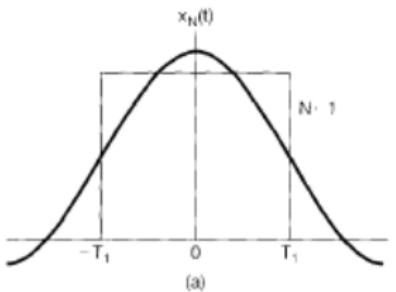
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Square wave Fourier series example



- ❑ $x(t)$ is discontinuous while each of its harmonic components is continuous
 - ❑ Is it possible to represent a discontinuous function as a sum of continuous functions?
- ❑ Consider the problem of approximating a given periodic signal by a linear combination of *finite* number of harmonically related complex exponentials
 - ❑ $x_N(t) = \sum_{k=-N}^N a_k e^{ik\omega_0 t}$
 - ❑ Approximation error, $e_N(t) = x(t) - x_N(t) = x(t) - \sum_{k=-N}^N a_k e^{ik\omega_0 t}$
 - ❑ Energy in the error over one period, $E_N = \int_T |e_N(t)|^2 dt$
- ❑ It can be shown that use of Fourier coefficients as a_k s minimizes E_N

Square wave Fourier series example



Convergence of Fourier series



- ❑ A periodic signal $x(t)$ have a Fourier series representation if all $a_k = \frac{1}{T} \int_T x(t) e^{-jk\omega_0 t} dt$ are finite
 - ❑ i.e., the integral does not diverge
- ❑ In some cases, even if all a_k s are finite, the resulting infinite series may not converge to the original signal $x(t)$
- ❑ If $\lim_{N \rightarrow \infty} E_N = 0$, $x(t)$ has a Fourier series representation
- ❑ Periodic signals having finite energy over a single period, $\int_T |x(t)|^2 dt < \infty$ are representable through Fourier series

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Linearity



- ❑ Fourier series coefficients of $x(t)$ are denoted by a_k
 - ❑ Notation: $x(t) \xleftrightarrow{\text{FS}} a_k$
- ❑ Let $x(t)$ and $y(t)$ are two periodic signals, both with period T
 - ❑ $x(t) \xleftrightarrow{\text{FS}} a_k$, and $y(t) \xleftrightarrow{\text{FS}} b_k$
- ❑ $z(t) = Ax(t) + By(t) \xleftrightarrow{\text{FS}} c_k = Aa_k + Bb_k$

Time shift



- If $x(t) \xrightarrow{\text{FS}} a_k$, then $y(t) = x(t - t_0) \xrightarrow{\text{FS}} b_k = ?$
- $b_k = \frac{1}{T} \int_T y(t) e^{-jk\omega_o t} dt = \frac{1}{T} \int_T x(t - t_0) e^{-jk\omega_o t} dt$
- Substitute $\tau = t - t_0$
 - $b_k = \frac{1}{T} \int_T x(\tau) e^{-jk\omega_o(\tau+t_0)} d\tau = e^{-jk\omega_o t_0} \frac{1}{T} \int_T x(\tau) e^{-jk\omega_o \tau} d\tau = e^{-jk\omega_o t_0} a_k$
- $x(t) \xrightarrow{\text{FS}} a_k \rightarrow x(t - t_0) \xrightarrow{\text{FS}} e^{-jk\omega_o t_0} a_k$
 - When a periodic signal is shifted in time, the magnitudes of its Fourier series coefficients remain unaltered

Time reversal



- If $x(t) \xleftrightarrow{\text{FS}} a_k$, then $y(t) = x(-t) \xleftrightarrow{\text{FS}} b_k = ?$
- $x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} \rightarrow x(-t) = \sum_{k=-\infty}^{\infty} a_k e^{-jk\omega_0 t}$
- Substitute $k = -m$
 - $x(-t) = \sum_{m=-\infty}^{\infty} a_{-m} e^{jm\omega_0 t}$
 - $b_k = a_{-k}$
- $x(t) \xleftrightarrow{\text{FS}} a_k \rightarrow x(-t) \xleftrightarrow{\text{FS}} a_{-k}$
- If $x(t)$ is even, i.e., $x(-t) = x(t) \rightarrow a_{-k} = a_k$
- If $x(t)$ is odd, i.e., $x(-t) = -x(t) \rightarrow a_{-k} = -a_k$

Time scaling



- ❑ Time scaling changes the period of the underlying signal
- ❑ $x(t) \rightarrow$ period T and fundamental frequency $\omega_0 = \frac{2\pi}{T}$
 - ❑ $x(\alpha t) \rightarrow$ period $\frac{T}{\alpha}$ and fundamental frequency $\alpha\omega_0$
- ❑ Time-scaling applies directly to each of the harmonic components of $x(t)$
 - ❑ Fourier coefficients for each of those components remain the same
- ❑ $x(\alpha t) = \sum_{k=-\infty}^{\infty} a_k e^{jk(\alpha\omega_0)t}$
 - ❑ Fourier series representation has changed because of the change in the fundamental frequency

Multiplication of two signals



- Let $x(t)$ and $y(t)$ are both periodic with period T

- $x(t) \xrightarrow{\text{FS}} a_k$, and $y(t) \xrightarrow{\text{FS}} b_k$
- $x(t) \times y(t) \xrightarrow{\text{FS}} c_k = ?$

-

$$\begin{aligned} x(t) \times y(t) &= \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} \times \sum_{m=-\infty}^{\infty} b_m e^{jm\omega_0 t} \\ &= \sum_{k=-\infty}^{\infty} \left(\sum_{l=-\infty}^{\infty} a_l b_{k-l} \right) e^{jk\omega_0 t} \end{aligned}$$

- $c_k = \sum_{l=-\infty}^{\infty} a_l b_{k-l}$
 - Discrete-time convolution of the sequence representing the Fourier coefficients of $x(t)$ and $y(t)$

Conjugation and Conjugate Symmetry



- ❑ If $x(t) \xrightarrow{\text{FS}} a_k$, then $y(t) = x^*(t) \xrightarrow{\text{FS}} b_k = ?$
- ❑ $x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} \rightarrow x^*(t) = \sum_{k=-\infty}^{\infty} a_k^* e^{-jk\omega_0 t}$
- ❑ Substitute $k = -m$
 - ❑ $x^*(t) = \sum_{m=-\infty}^{\infty} a_{-m}^* e^{jm\omega_0 t}$
 - ❑ $b_k = a_{-k}^*$
- ❑ If $x(t)$ is real
 - ❑ $x^*(t) = x(t) \rightarrow a_k = a_{-k}^* \rightarrow a_{-k} = a_k^*$
 - ❑ Fourier series coefficients will be conjugate symmetric
- ❑ if $x(t)$ is real and even
 - ❑ $a_{-k} = a_k^*$ and $a_{-k} = a_k \rightarrow a_k = a_k^*$ (real and even coefficients)
- ❑ if $x(t)$ is real and odd
 - ❑ $a_{-k} = a_k^*$ and $a_{-k} = -a_k \rightarrow a_k = -a_k^*$ (purely imaginary and odd coefficients)

Parseval's relation



- Average power (i.e., energy per unit time) of the periodic signal $x(t)$ is given by:



$$\begin{aligned} P &= \frac{1}{T} \int_T |x(t)|^2 dt = \frac{1}{T} \int_T x(t)x^*(t)dt \\ &= \frac{1}{T} \int_T \left(\sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} \right) \left(\sum_{m=-\infty}^{\infty} a_m^* e^{jm\omega_0 t} \right) dt \\ &= \frac{1}{T} \int_T \sum_{k=-\infty}^{\infty} \left(\sum_{l=-\infty}^{\infty} a_l a_{k+l}^* \right) e^{jk\omega_0 t} dt \end{aligned}$$

- $e^{jk\omega_0 t}$ is periodic with period T for all values of k , hence $\int_T e^{jk\omega_0 t} dt = 0$
- Except for $k = 0$, for which it is constant
- $P = \frac{1}{T} \int_T \sum_{l=-\infty}^{\infty} a_l a_l^* dt = \sum_{l=-\infty}^{\infty} |a_l|^2$
- $\frac{1}{T} \int_T |x(t)|^2 dt = \sum_{k=-\infty}^{\infty} |a_k|^2$

Summary of properties



Property	Periodic signal	Fourier series coefficients
Linearity	$Ax(t) + By(t)$	$Aa_k + Bb_k$
Time shifting	$x(t - t_0)$	$a_k e^{-jk\omega_0 t_0}$
Frequency shifting	$e^{jM\omega_0 t} x(t)$	a_{k-M}
Conjugation	$x^*(t)$	a_{-k}^*
Time reversal	$x(-t)$	a_{-k}
Time scaling	$x(\alpha t)$, $\alpha > 0$ (periodic with period $\frac{T}{\alpha}$)	a_k
Periodic convolution	$\int_T x(\tau) y(t - \tau) d\tau$	$T a_k b_k$
Multiplication	$x(t)y(t)$	$\sum_{l=-\infty}^{\infty} a_l b_{k-l}$

Summary of properties

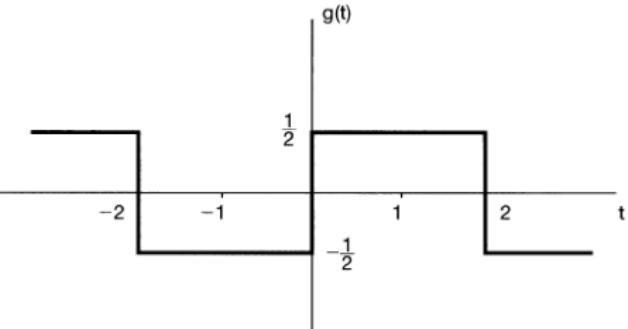


Property	Periodic signal	Fourier series coefficients
Differentiation	$\frac{d}{dt}x(t)$	$jk\omega_0 a_k$
Integration	$\int_{-\infty}^t x(\tau)d\tau$ (finite valued and periodic only if $a_0 = 0$)	$\frac{1}{jk\omega_0} a_k$
Conjugate symmetry for real signals	$x(t)$ is real	$a_k = a_{-k}^*$
	$x(t)$ is real and even	a_k real and even
	$x(t)$ is real and odd	a_k is purely imaginary and odd
Even-odd decomposition	$\begin{cases} x_e(t) = \frac{x(t)+x(-t)}{2} \\ x_o(t) = \frac{x(t)-x(-t)}{2} \end{cases}$	$\begin{cases} \text{Re}\{a_k\} \\ j\text{Im}\{a_k\} \end{cases}$
Parseval's relation	$\frac{1}{T} \int_T x(t) ^2 dt = \sum_{k=-\infty}^{\infty} a_k ^2$	

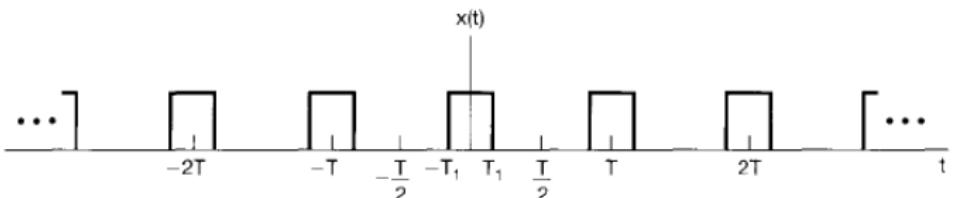
Example 1



- ❑ Consider the signal $g(t)$ with a fundamental period of 4
- ❑ Determine the Fourier series representation of $g(t)$



- ❑ Remember $x(t)$ from Slide 25? $x(t) = \begin{cases} 1, & |t| < T_1 \\ 0, & T_1 < |t| < \frac{T}{2} \end{cases}$



- ❑ $a_0 = \frac{2T_1}{T}$, $a_k = \frac{\sin(k\omega_0 T_1)}{k\pi}$ for $k \neq 0$

Example 1

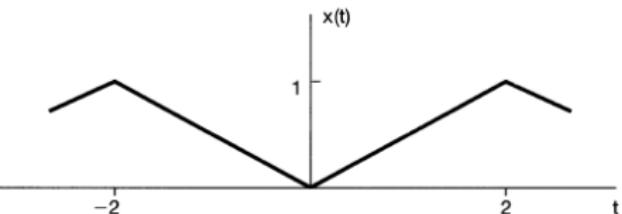


- ❑ Can we relate $g(t)$ with $x(t)$?
 - ❑ With $T = 4(\omega_0 = \frac{2\pi}{4})$ and $T_1 = 1$, $g(t) = x(t - 1) - \frac{1}{2}$
- ❑ Fourier Series coefficients of $x(t - 1)$ are $b_k = a_k e^{-jk\frac{\pi}{2}}$
 - ❑ Time shift property
- ❑ Fourier series coefficients of the dx offset term $\frac{1}{2}$ are:
$$c_k = \begin{cases} 0, & k \neq 0 \\ \frac{-1}{2}, & k = 0 \end{cases}$$
- ❑ Using linearity property, coefficients of $g(t) = x(t - 1) - \frac{1}{2}$ are:
 - ❑ $d_k = \begin{cases} a_k e^{-jk\frac{\pi}{2}}, & k \neq 0 \\ a_0 - \frac{1}{2}, & k = 0 \end{cases} = \begin{cases} \frac{\sin(k\frac{2\pi}{4})}{k\pi} e^{-jk\frac{\pi}{2}}, & k \neq 0 \\ 0, & k = 0 \end{cases}$

Example 2



- ❑ Consider the triangular wave signal $x(t)$ with period $T = 4$
- ❑ Determine its Fourier series representation
- ❑ $x(t)$ is the integral of $g(t)$ in the previous example
 - ❑ $x(t) = \int_{-\infty}^t g(\tau) d\tau$
- ❑ Fourier series coefficients of $x(t)$, $e_k = \frac{1}{jk\frac{\pi}{2}} d_k$
- ❑ $e_k = \frac{2\sin(k\frac{\pi}{2})}{jk^2\pi^2} e^{-jk\frac{\pi}{2}}, k \neq 0$
- ❑ $e_0 = \frac{1}{T} \int_T x(t) dt = \frac{1}{T} \times \frac{1}{2} \times T \times 1 = \frac{1}{2}$

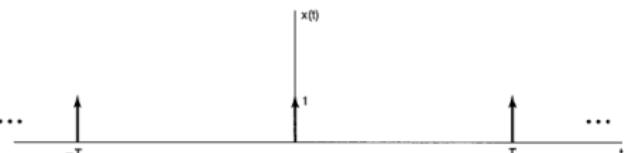


Example 3



- Consider a periodic train of impulses (impulse train)

- $x(t) = \sum_{k=-\infty}^{\infty} \delta(t - kT)$

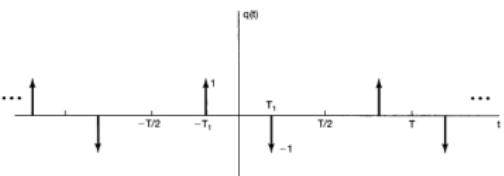
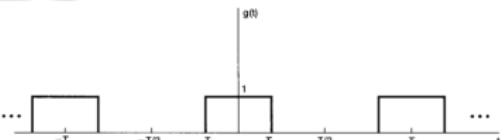


- $a_k = \frac{1}{T} \int_T x(t) e^{-jk\omega_0 t} dt = \frac{1}{T} \int_{-\frac{T}{2}}^{\frac{T}{2}} \delta(t) e^{-jk\omega_0 t} dt = \frac{1}{T}$

- All Fourier series coefficients of the impulse train are identical

- Consider the square-wave signal $g(t)$ and its derivative $q(t)$

- $q(t) = x(t + T_1) - x(t - T_1)$



Example 3



- ❑ Fourier series coefficients b_k of $q(t)$ are:

$$b_k = e^{jk\omega_0 T_1} a_k - e^{-jk\omega_0 T_1} a_k = (e^{jk\omega_0 T_1} - e^{-jk\omega_0 T_1}) a_k = \frac{2j\sin(k\omega_0 T_1)}{T}$$

- ❑ Using time-shifting and linearity properties

- ❑ Since $g(t)$ is integral of $q(t)$, the Fourier series coefficients c_k of $g(t)$ are:

- $c_k = \frac{1}{jk\omega_0} b_k = \frac{2j\sin(k\omega_0 T_1)}{jk\omega_0 T} = \frac{\sin(k\omega_0 T_1)}{\pi k}, K \neq 0$

- ❑ c_0 is the average value of $g(t)$ over one period

- $c_0 = \frac{2T_1}{T}$

Example 4



- ❑ Following facts about a signal $x(t)$ are known
 1. $x(t)$ is a real signal
 2. $x(t)$ is periodic with period $T = 4$, and it has Fourier series coefficients a_k
 3. $a_k = 0$ for $|k| > 1$
 4. The signal with Fourier coefficients $b_k = e^{-jk\frac{\pi}{2}} a_{-k}$ is odd
 5. $\frac{1}{4} \int_4 |x(t)|^2 dt = \frac{1}{2}$
- ❑ Determine the signal $x(t)$
- ❑ Fact 2 $\rightarrow \omega_0 = \frac{2\pi}{T} = \frac{\pi}{2}$
- ❑ Fact 3 $\rightarrow x(t) = a_0 + a_1 e^{j\frac{\pi}{2}t} + a_{-1} e^{-j\frac{\pi}{2}t}$
- ❑ Fact 1 $\rightarrow a_1 = a_{-1}^* \rightarrow x(t) = a_0 + 2\Re\{a_1 e^{j\frac{\pi}{2}t}\}$

Example 4



□ Fact 4 →

- a_{-k} corresponds to the signal $x(-t)$
- $b_k = e^{-jk\frac{\pi}{2}} a_{-k} = e^{-jk\omega_0} a_{-k}$ correspond to the signal
 $x(-(t-1)) = x(-t+1)$
- According to fact 4, $x(-t+1)$ is odd
- Since $x(t)$ is real, $x(-t+1)$ should also be real
- Real and odd $x(-t+1) \rightarrow$ purely imaginary and odd b_k s
 - $b_0 = 0$, and $b_{-1} = -b_1$

□ Fact 5 →

- Time-reversal and time-shift does not change average power
- $\frac{1}{4} \int_4 |x(t)|^2 dt = \frac{1}{4} \int_4 |x(-t+1)|^2 dt = \frac{1}{2}$
- Using Parseval's relation:
 $|b_1|^2 + |b_{-1}|^2 = |b_1|^2 + |-b_1|^2 = 2|b_1|^2 = \frac{1}{2} \rightarrow b_1 = \frac{j}{2}$ or $\frac{-j}{2}$
- $a_0 = 0$, $a_1 = e^{-j\frac{\pi}{2}} b_{-1} = -jb_{-1} = jb_1$
- $b_1 = \frac{j}{2} \rightarrow x(t) = -\cos(\frac{\pi t}{2})$
- $b_1 = \frac{-j}{2} \rightarrow x(t) = \cos(\frac{\pi t}{2})$

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complex exponentials as an input to a LTI system



- ❑ Response of an LTI system to a linear combination of complex exponentials takes a particularly simple form
 - ❑ $x(t) = e^{st} \rightarrow y(t) = H(s)e^{st}$
 - ❑ $H(s) = \int_{-\infty}^{\infty} h(\tau) e^{-s\tau} d\tau$
 - ❑ $h(t)$ is the impulse response
 - ❑ $H(s)$ is called system function
- ❑ Consider the specific case where $\text{Re}\{s\} = 0$, i.e., $s = j\omega$
- ❑ $H(j\omega) = \int_{-\infty}^{\infty} h(\tau) e^{-j\omega\tau} d\tau$ is called frequency response of the system

Response of a LTI system to a periodic signal



- ❑ Let $x(t)$ is a periodic signal: $x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t}$
- ❑ Output of a LTI system with $x(t)$ as input:
 $y(t) = \sum_{k=-\infty}^{\infty} a_k H(jk\omega_0) e^{jk\omega_0 t}$
 - ❑ Since each of the complex exponentials is an eigenfunction of the system
- ❑ $y(t)$ is also periodic with period T and Fourier series coefficients $a_k H(jk\omega_0)$

Example



- Input $x(t) = \sum_{k=-3}^3 a_k e^{jk2\pi t}$
 - $a_0 = 1, a_1 = a_{-1} = \frac{1}{4}, a_2 = a_{-2} = \frac{1}{2}, a_3 = a_{-3} = \frac{1}{3}$
- Impulse response: $h(t) = e^{-t} u(t)$
- Impulse response:

$$\begin{aligned} H(j\omega) &= \int_{-\infty}^{\infty} h(\tau) e^{-j\omega\tau} d\tau = \int_0^{\infty} e^{-\tau} e^{-j\omega\tau} d\tau \\ &= \int_0^{\infty} e^{-(1+j\omega)\tau} d\tau = \frac{1}{1+j\omega} \end{aligned}$$

- Output $y(t) = \sum_{k=-3}^3 b_k e^{jk2\pi t}$
 - $b_k = a_k H(jk2\pi) = \frac{a_k}{1+jk2\pi}$
 - $b_0 = 1, b_1 = \frac{1}{4(1+j2\pi)}, b_{-1} = \frac{1}{4(1-j2\pi)}, b_2 = \frac{1}{2(1+j4\pi)}, b_{-2} = \frac{1}{2(1-j4\pi)},$
 $b_3 = \frac{1}{3(1+j6\pi)}, b_{-3} = \frac{1}{3(1-j6\pi)}$

Thanks.