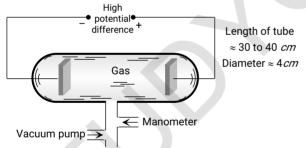


Electron, Photon, Photoelectric Effect and X-rays

Electric Discharge Through Gases

At normal atmospheric pressure, the gases are poor conductor of electricity. If we establish a potential difference (of the order of 30 kV) between two electrodes placed in air at a distance of few cm from each other, electric conduction starts in the form of sparks.

The discharge of electricity through gases can be systematically studied with the help of discharge tube shown below



As the pressure in sign discharge tube is gradually reduced, the following is the sequence of phenomenon that are observed.

- (1) At normal pressure no discharge takes place.
- (2) At the pressure 10 *mm* of *Hg*, a zig-zag thin red spark runs from one electrode to other and cracking sound is heard.



- (3) At the pressure 4 $\stackrel{\textbf{Fig. 25.2}}{\textit{mm.}}$ of $\stackrel{\textbf{Hg.}}{\textit{Hg}}$, an illumination is observed at the electrodes and the rest of the tube appears dark. This type of discharge is called dark discharge.
 - (4) When the pressure falls below 4 mm of Hg then the

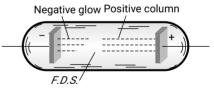
whole tube is filled with bright light called positive column and colour of light depends upon the nature of gas in the tube as shown in the following table.

Table 25.1: Colour for different gases

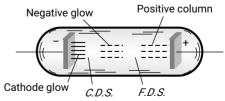
Gas	Air	Hı	№ 2	Ch	CO₂	Neon
Colour	Purple red	Blue	Red	Green	Bluish white	Dark red

(5) At a pressure of 1.65 mm of Hg:

Sky colour light is produced at the cathode it is called as negative glow. Positive column shrinks towards the anode and the dark space between positive column and negative glow is called Faradays dark space (FDS).



(6) At a pressure of 0.8Fign#5Ag: At this pressure, negative glow is detached from the cathode and moves towards the anode. The dark space created between cathode and negative glow is called as Crook's dark space. Length of positive column further reduced. A glow appear at cathode called cathode glow.



- (7) At a pressure of 0.0 Fig. 27.0 f Hg: The positive column splits into dark and bright disc of light called striations.
 - (8) At the pressure of 0.01 or 10^{-2} mm of Hg some invisible

2Electron, Photon, Photoelectric Effect and X-Rays

particles move from cathode whic tube on the opposite side of cath-These invisible rays emerging from rays.

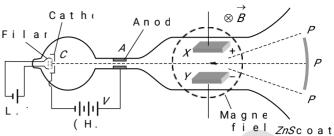
(9) Finally when pressu⁴r*mem* of *fHag*,p there is no discharge in tube.

Cathode Rays

- (1) Cathoderays, discovered by
- (2) They are streams of fast mov
- (3) They can be produced by us containing gas at a low pre² smsno f/leg. o
- (4) The cathoderays in the disc produced due to ionisation of gas due to collision of positive ions.
 - (5) Cathoderaystravel instrai
- (6) Cathode rays are emitted no surface. Their direction is inde anode.
- (7) Cathoderays exert mechanic strike.
- (8) Cathode rays produce heat v surface.
 - (9) Cathoderays produce fluore
- (10) When cathode rays strike : metal of high atomic weight X-arnady hsi emitted from the objects.
- (11) Cathoderays are deflected by a magnetic field.
- passed.
 - (13) Cathoderays can penetrate
- (14) Cathoderays are found to the found to t $t o \frac{1}{10} th$ of velocity of light.

J. J. Thomson's Experiment

- (1) It's working is based on th electron is subjected to \vec{t} he cro magnetic field, it experiences case the forces on the electrons these fields are equal and oppo undeflected.
- (2) When no field is applied, t illuminati Ponsat point
- (3) In the presence of any fiel electron beam deflected up por dpo wr
- (4) If both the fields are app adjusted such that electron bea produces il lumiP.nation at point



In this case; Electric **=** fe**g** ∈ Fs**e** v **@** ∈ eMag n **Fig. 2**:

$$\Rightarrow v = \frac{E}{B}$$
; $v = v e l o c i t y o f e l e c t r o n$

(5) As electron beamaccelerated f loss in potential energy appears as g If supproset he potential difference anode then, loss in ep/otential energy

And gain in kinetic energy= $\frac{1}{2}amtv^2$ anod

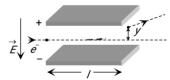
i.e.
$$eV = \frac{1}{2}mv^2 \implies \frac{e}{m} = \frac{v^2}{2V} \implies \frac{e}{m} = \frac{E^2}{2VB^2}$$

Thomson $f \stackrel{e}{\rightarrow} = 117d \times 10^{11} C / kg$.

If one includes the relativistic v $(m = m_0 / \sqrt{1 - v^2 / c^2})$, then specific charge decreases with the increase in its ve

(6) The deflection of an electron ;/=wLheenroeth of e pr=ch pl

(12) Cathode rays ionise the gadeflection of electromr=inspteheed foifet electron.

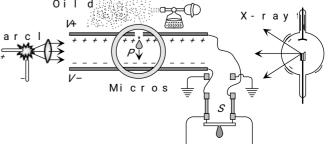


Millikans Oi^FigDr² op Experiment

- (1) Millikan performed the pione 6 for the precise measurement of the ch
- (2) By applying suitable electri plates, the charged oil droplets cou even held stationary in the field o time. He found that the charge on an o integral multiple of an exile¹ m[®]Centary c
 - (3) In this experiment charge on th

$$q = \frac{6 \frac{\pi}{V} \sqrt{(v_1 + v_2) d}}{V} \left[\frac{9 \frac{\pi}{V}}{2 g(\rho - \sigma)} \right]$$
Oild Atomi

X-ray



qas. This is done*qh*omyonfnesaisnugrliynigonise ion of the gas.

wherηe Coefficient of νη i= sTceorsmin. velocity of drop when no electric applied between the plates.

V = Potential difference d etSeparation be ptwelen spil taopt=eDse,onisli, ty (

Positive Rays

When potential difference is a of a dischar g^{*} nenntouf*htg*e ,(& Dectronsar (the perforated cathode. As they m energy. These energetic electron of the gas in the discharge tube, positive ions so for med at variou anode, travel towards the cathod the positive ions when reach the the holes in the cathode and a fai from each hole on the backside of positive rays, which are coming ou

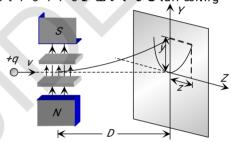


- (1) Positive_{F lig}y ₂ are positive experimental gas does not have is has isotopes then positive rays ar different masses.
- (2) They travel in straight line placed in their path. But the spee smaller than that of cathoderays.
- (3) They are deflected by elect the deflections are small as compa
- (4) They show a spectrum of vel ions move with different velociti much less than that of cathoderay
- (5q)/m ratio of these rays depend gas in the tube (while in cas/enios) constant and doesn't depend on th *q/m*forhydrogenis maximum.
- (6) They carry energy and moment positive rays is more than that of
- (7) The value of charge on posi multiple of electronic charge.
- (8) They cause ionisation (whi produced by cathoderays).

Thomson's Mass Spectrograp

It is used to measure atomic mas

- (1) The positive ions are produce platue s, Terminal velocity of drop hand side. These ions are accelerate of the positive ions pass through th This fine ray of positive ions East sub magneti & a fnidetlhden allowed to strike a $(\vec{F} \parallel \vec{B} \quad \mathbf{b} \ \mathbf{u}\vec{F}\mathbf{t}$ \vec{B} o $|\vec{r}|$
 - (2) If the initial moth ion of the io electric and magnetic fields are ap force due to electric fixe la atii ss ainn dt dhue magnetic fielz-dditestailcom.g



deflecti**p_in_{a.} g**ue The qELD

The deflection due to maze net ic f

From equation $z^2 (\exists k) \frac{q}{a} y n d (ii)$,

wherke= $\frac{B^2LD}{c}$; This is the equation of means all the charged particles movi but of sq/ammaealue will strike t/mzpelsacrree e on a parabolic track as shown in the a

(3) All the positi*q/m* enoivoin nsgowfistahmoeli. f velocity lie on the same parabola. Hi the valyandoThe ions of different spec on different parabola.

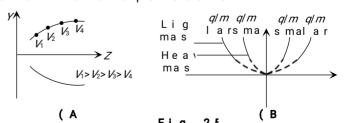
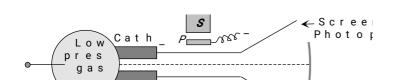


Fig. 2 t (4) The number of parabola tells t present in the given ionic beam.

Bainbridge Mass Spectrograph



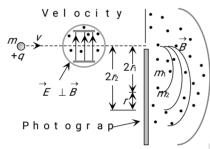
In Bainbridge mass spectrograp velocity are selected by using a v are subjected to a uniform magnet velocity of the particles. The par isotopes follow different circul

(1 W/e locity sTehleepcotsoirt: ive ions h velorogiettys isolated from all other chamber the electric and magneti the particle moves undeflected. F is $v = \frac{E}{R}$ and $\mathbf{d}\mathbf{f}$, \mathbf{B} and $\mathbf{d}\mathbf{f}$ should be mutually \mathbf{f} each other.

(2A) nalysing that mbies rc: hamber to Biasg applied perpendicular to the dire As a result the particles move alo

$$r = \frac{mE}{qBB}$$
 $\Rightarrow \frac{q}{m} = \frac{E}{BB \cdot r}$ als $\frac{r_1}{r_2} = \frac{m_1}{m_2}$

In this way the particles of diff $\lambda_{Thermal neutron} = \frac{6.62 \times 10^{-34}}{\sqrt{2} \times 1.67 \times 10^{-27} \times 1.38 \times 10^{-23} T} = \frac{30.83}{\sqrt{T}} \mathring{A}$ recles of different radii and rea at e. (4 Ratio of wavelength of phTohteon circles of different radii and rea plate.



Separation breqtqwezentwo traces $= d = 2r_2 - 2r_1 = \frac{2v(m_2 - m_1)}{aB'}$

Matter Waves (de-Broglie Wa

According to de-Broglie a mo someti mes acts as a wave and somet

The wave associated with moving wave or de-Broglie wave and it pro packets with group velocity.

(1d) e - Broglie w-acvceolrednigntght:o de - [thè wavel ength of de-Broglie wave

$$\lambda = \frac{h}{p} = \frac{h}{mv} = \frac{h}{\sqrt{2 \, mE}} \Rightarrow \lambda \propto \frac{1}{p} \propto \frac{1}{v} \propto \frac{1}{\sqrt{E}}$$

Wherher Plank's concentsolats as noto, fthevera Speed of the EnpEanretrigcyloef, the particl

The smallest wavelength whose is thay tof - rays.

The wavelength of matter wave microscopic particles like $e\alpha$ + ϵ part*eta*clies of the 10 T¹⁰ daner of

(2 de-Broglie wavelength assoc partiTcH ees nergy of a charged parti potential kolis£ f e≀ntvê ntopve

Hence de - Brogli λ e $\frac{h}{v}$ a $\frac{h}{\sqrt{2mE}}$ e $\frac{h}{\sqrt{2mqV}}$

$$\lambda_{Electron} = \frac{12.27}{\sqrt{V}} \mathring{A}$$
, $\lambda_{Proton} = \frac{0.286}{\sqrt{V}} \mathring{A}$,

$$\lambda_{Deutron} = \frac{0.202}{\sqrt{V}} \mathring{A}, \qquad \lambda_{\alpha-particle} = \frac{0.101}{\sqrt{V}} \mathring{A}$$

(3)de-Broglie wavelength associa partiFcdi Meustron de - Broglie wavelengthis

$$\hat{\mathcal{A}}_{Neutron} = \frac{0.286 \times 10^{-10}}{\sqrt{E \text{ (in } eV)}} m = \frac{0.286}{\sqrt{E \text{ (in } eV)}} \hat{\mathcal{A}}$$

Energy of thermal neutrons at ordi

$$E = kT \Rightarrow \lambda = \frac{h}{\sqrt{2mkT}} ; \text{ wher } T = Absolut$$

t emper fat Blook t z man' s1.080×m0524 stowlet/kelvin,

$$\lambda_{Thermal neutron} = \frac{6.62 \times 10^{-34}}{\sqrt{2 \times 1.67 \times 10^{-27} \times 1.38 \times 10^{-23} T}} = \frac{30.83}{\sqrt{T}} \dot{A}$$

wavelength of a phÆoitsogniorfe2_ppe⇒tb-grgy

While the wavelength of an elKiesctro . Therefore, for the s

t h e r
$$\frac{\lambda_{nk}}{\lambda_{e}} = \frac{c}{E} \sqrt{2mK} = \sqrt{\frac{2mc^{2}K}{E^{2}}}$$

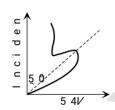
Characteristics of Matter Way

- (1) Matter wave represents the p particle in space.
 - (2) Matter waves are not electroma
- (3) de-Brogile or matter wave is i on the material particle. It means, wave is associated with every moving or uncharged).
- (4) Practical observation of matt when the de-Broglie wavelength is of particles.
- (5) Electron microscope works on waves.
- (6) The phase velocity of the matt than the speed of the light.
- (7) Matter waves can propagate in not mechanical waves.
- (8) The number of de-Broglier[†] wave orbital entectronis
- (9) Only those circular orbits a stable whose circumference is int Broglie wavelength associated witl

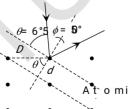
Davision and Germer Experimen

(1) It is used to study the scatter or to verify the wave nature of elec emitted by electron gun is made t al ong cubical axis a*t*V/carpyasrttailcbuelhaa containing inter**⊘antoi**nanin⁄og bleist ance three dimensional diffraction gr beamobtained from electron gun.

- (2) The diffracted beam of ele detector which can be positioned about the point of incidence. The electrons can also be varied by ch the electron gun.
- (3) According to classical phys beam of electrons at all scatter Davisson and Germer, found that beam of electrons was not the same angles of scattering. It is malatm 5 4voltpotential difference.



(4) If the de-Broglie waves exi should be dif #Frraacytse.d Uassing the B $2 d \sin \theta = n \lambda$, we can determine the wavel wher e/ = distance between dif $\theta = \frac{(180 - \phi)}{2}$ = glancing angle for incide



The distandFeigb.e2tsween diWfcfrryasotta this experd'i=6ne 9d' at nidsthe Bragg' °s Ta hi g i v e s7 = f 1b $\cancel{R} = 2 \times 0.91 \times 10^{-10} \sin 65^{\circ} = 1.65 \mathring{A}$

Now the de-Broglie wavelength using the $f \lambda = \frac{12.27}{mu} \pm \frac{12.27}{454} = 1.67 \text{ Å}$ Broglie hypothesis is verified.

(5) The Bragg's formula can be re

$$\therefore \theta = 90 - \frac{\phi}{2} \text{ and } d = D \cos \theta = D \sin \frac{\phi}{2}$$

Using
$$\theta = \cos \frac{\phi}{2}$$

$$2 d \sin \theta = \lambda \implies 2(D \sin \frac{\phi}{2}) \cdot \cos \frac{\phi}{2} = \lambda \implies D s i \phi r \lambda$$

Heisenberg Uncertainty Prince

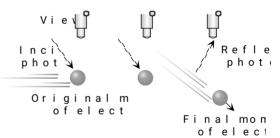
- (1) According to Heisenberg's un i mpossi ble to measure si multaneous l moment um of the particle.
- (2) $\triangle x$ eatn Δx be the uncertainty in th measurement of the position and mom ; where $h = 6 \text{ and 3d} 0^{3} \text{ }^4 \text{ J-s} \text{ is}$

the Planck' $\$\frac{h}{2\pi}$ \div 100\$ $1 10 $$^{-1}$ 1 $1 -sec)

A more rigorous t r \triangle e. \triangle pt $\ge \frac{\hbar}{2} \left($ on $\frac{1}{4\pi} g\right)$ i ves

(3) $\Delta x = f \cdot 0 + h \Delta p = n \infty a \cdot n \cdot d \Delta i p = 0 + h \Delta e = n \infty$

i.e., if we are able to measure the ex particle (say an electron) then t measurement of the linear momentum o Similarly, if we are able to measure of the p aiα, Δρi=c 0 ,e then we can not measu position of the particle at that time



An electron**fic**gan2ntot be observed w mo me n t u m

- (4) Uncertainty principle success
- (i) Non-existence of electrons in
- (ii) Finite size of spectral lines
- (5) The Heisenberg uncertainty pr to energy and time, angular mome displace me $\Delta t \ge H \frac{h}{2\pi}$ nce $\Delta t . \Delta \theta \ge \frac{h}{2\pi}$ and

(6) If the radius orthehnetnhuecoloreouls aibs finding the electron ∆ix ≠n Staind de utnhoe enrutcali in moment Δψ m ins

Photon

6Electron, Photon, Photoelectric Effect and X-Rays

According to Ei enstein's quant the bundles (packets or quanta) o called a photon and possessing ene

(1E) n e r g y of p̄ h e t o yn o f p h o t o n i s g

 $E = h \nu = \frac{hc}{a}$; wher ce Speed of $h \models i Pglhat n$, constant× ± 0 % 4 . \mathcal{L} 6ec, ν = Freque \mathcal{L} 2; \mathcal{L} = i Wavelength of light.

In elect
$$\mathbb{E}(\text{eV}) = V + \frac{hc}{e\lambda} + \frac{12375}{\lambda(\mathring{A})} \approx \frac{12400}{\lambda(\mathring{A})}$$

(2N) lass of phAocttouna: I y rest mass (zero. But it's effective massis q

$$E = mc^2 = h \nu$$
 $\Rightarrow m = \frac{E}{c^2} = \frac{h \nu}{c^2} = \frac{h}{c\lambda}$. This mass

known as kinetic mass of the photo (3M) omentum of the photon

Moment
$$\mu$$
 and $m \times c = \frac{E}{c} = \frac{h \nu}{c} = \frac{h}{\lambda}$

(4N) umber of emitted ephotmbas of emitted per second from a source c of wave 1λ eam gdt pho Rvie sr give $(n) = \frac{P}{E} = \frac{P}{h \nu} = \frac{P\lambda}{hc}$ whe E = energy of each photon

(51) ntensity/) of finleir gghyt c(rossing)

$$i.e.$$
 $I = \frac{E}{At} = \frac{P}{A}$ $\left(\frac{E}{t} = P = \text{radiation power}\right)$

At a dis/tfar no cmea point sou/Picnet ce fn s given $\frac{1}{\sqrt{2r}} \Rightarrow I \propto \frac{1}{r^2}$

(6N) umber of photons falm) lifegspel power of rad Eiast ti hoen ea med rgy of $a = p^{P}$ hot

Photo-Electric Effect

The photo-electric effect is (called photo-electrons when lig from the surface, the electron mu the incident radiation to overcon in the material of the surface.

The photoelectric effect was f Hertz and it was investigated in d Philipp Lenard.

The photoelectric effect is b conservation of energy.

(1) Work function (or th///es:Toble minimum energy of incident radia electrons from metallic surface that surface.

$$W_0 = h \, V_0 = \frac{hc}{\lambda_0}$$
 Joules; $\nu_0 = T h r e s h o l d f r e q u$

 λ_0 = Threshold wavelength

Work function i $M(eV) = e^{\frac{hc}{h}} = \frac{12375}{e^{\frac{h}{h}}} = \frac{12375}{h}$ olt

Table 25.2: Work function of se

Element	Work fun (<i>eV</i>)	Element	Workfun (<i>eV</i>)		
Platinu	6.4	Aluminu	4 . 3		
Gold	5 . 1	Silver	4 . 3		
Nickel	5 . 1	Sodium	2.7		
Carbon	5.0	Lithium	2.5		
Silicon	4 . 8	Potassi	2.2		
Copper	4.7	Cesium	1 . 9		

(21) hreshold f nb) e Tip bu e moc yn i(mum freque incident radiations required to ej surface is defined as threshold freq

If incident/<fντ=eNqcuphot/oelectroner

For most metal sthethreshold freq (corresponding to waveleng to/th)s, bbeuttw for potassium and cesium oxides λitis between 400*nm*)nd 700

(3T) hreshold wa&) eTheen mgatxhi (mum wa ve l of incident radiations required to normally per second is called int { metallic surface is defined as thres

If incident λw+al₀v⇔eNloepnhgothoelect rone

(4E) instein's photoekke ot dio ge quaEti photoelectric effect is the result o between photon and electron in which absorbed

$$E = h\nu_0$$

$$K.E. = 0$$

$$E = h\nu$$

$$K.E. = max$$

$$e^-$$

$$Met$$

Einstein's pFhg.t 2 leE= 4Mc+tKmeicequatio where $_{\text{max}} = \frac{1}{2} m v_{\text{max}}^2 =$ maximum kinetic emitted el ectrons.

Experimental Setup for Photo

(1) Two conducting eleQotarnoddeest, htoh (P) are enclosed in an evacuated glass

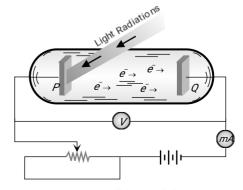


Fig. 25

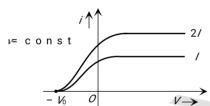
(6
$$y_0 = \frac{h}{e}(\nu - \nu_0) = \frac{hc}{e} \left(\frac{1}{\lambda} - \frac{1}{\lambda_0} \right) = 12375 \left(\frac{1}{\lambda} - \frac{1}{\lambda_0} \right)$$

- (2) The battery or other sourc creates an electric field in the d
- (3) Light of certain wavelengt surface of cathode causes a currer photoelectric current.
- (4) As potential difference incalsoincreases till saturation is
- (5) When polarity of bi.e.pteQfigsias negative pwor.tpelnatAf) eaellectrons star towards the cathode.
- (6) At a particular negaQthio veel pec will reach Qtahmed ptlhaet ecurrent will negative potensttioaplpinagc pd betneont te kd.a.b. Maximum kinetic energy of photoel potential wilk mat = h(le vt, l)eefvore be

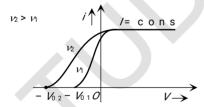
Effect of Intensity and Fre

(1E) f fect of il mft telmes intyrensity of (while it's frequency is kept the higher value, showing that more elunit time. But the sl⁄ctdoopeps indyr pózechta en mg t

Interossinby of incidexnot ophortoenm photoelectrocophorto biumeent



(2**Effect of frlefate a qy**e:ncy of increases, (keeping intensity increases but their is no change i



Important FF p. r2mulae for Phot

(1)
$$\nu = h \nu_0 + K_{\text{max}}$$
 $a \kappa \eta_{\text{max}} d = e V_0$

(2)
$$m_{\text{max}} = eV_0 = h(\nu - \nu_0)$$
 $\Rightarrow \frac{1}{2} m v_{\text{max}}^2 = h(\nu - \nu_0)$

$$(3)_{\text{max}} = \sqrt{\frac{2h(\nu - \nu_0)}{m}}$$

$$(4)_{\text{max}} = \frac{1}{2} m v_{\text{max}}^2 = e V_0 = hc \left(\frac{1}{\lambda} - \frac{1}{\lambda_0} \right) = hc \left(\frac{\lambda_0 - \lambda}{\lambda \lambda_0} \right)$$

$$(5)_{\text{max}} = \sqrt{\frac{2 hc}{m} \frac{(\lambda_0 - \lambda)}{\lambda \lambda_0}}$$

Compt on Effect

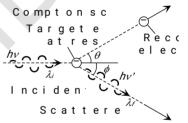
- (1) The scattering of a photon by Compton effect.
 - (2) The energy and momentum is cons
- (3) Scattered photon will have wavelength) as compare to incident p
- (4) The energy lost by the photon i kinetic energy.
- (5) The change in wavelength due called Compton shift. Com

$$\lambda_f - \lambda_i = \Delta \lambda = \frac{h}{m_0 c} (1 - \cos \phi)$$

I
$$\mathcal{B} = 0$$
, $\Delta \lambda = 0$

$$\phi = 9^{\circ} \Omega \Delta \lambda = \frac{h}{m_0 c} = 0.24 \text{ nm}$$

 $\phi = 180^{\circ}$, $\Delta \lambda = \frac{2h}{m_0 c} = 0.48 \text{ nm}$ (called Compton wa



X-Rays

Fig. 25

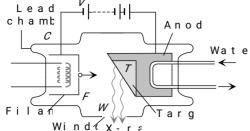
- (1) X-rays were discovered by sciethey are also called Ront genrays.
- (2) Rontgen discovered that whe discharge tu b³ en/n os f/k/g a p td 1 p 60 tential dif kept k/12/5 hen some unknown radiations (by anode.
- (3) There are three essential production of X-rays.
 - (i) A source of electron
 - (ii) An arrangement to accelerate
- (iii) A target of suitable materia high melting point on which these hig

Coolidge X - Ray Tube

- (1) It consists of a highly evacua cathode and target (also known as fill the cathode consist of a tungsten for coated with oxides of bariumor stroof electrons even at low temperat surrounded by a molybdenum cylinder w.r.t. the target.
- (2) The target (It is a material of melting point and high thermal conduor molybdenumis embedded in a copper

8Electron, Photon, Photoelectric Effect and X-Rays

(3) The face of the tatqetthe sis electron strę, am.



- (4) The filan field t 25 heated by pa A high potentiaa1l0/dl/ntfof8/eØr eiscaep (olied the target and cathode to acceler emitted by filament. The stream (are focussed on the target.
- (5) Most of the energy of the el heat (above 98%) and only a fract electrons (about 2%) is used to pr
- (6) During the operation of the is produced in this target, this copper anode to the cooling fins f radiation and convection.
- (7¢ ontrol of intensinty eonfs iXt-yrai number of X-ray photons produce Absorption of X-Rays intensity of X-rays emitted is electrons emitted per second from increased by increasing tinheresity iof IX-a rays ∝ Filament current
- (8¢ ontrol of quality or ple meaty) Quality of X-rays implies the pen can be controlled by varying the the cathode and the target.

For large potential differen electrons will be large and hence of X-rays.

Table 25.3: Types of X-

Hard X-rays	Soft X-rays
More penetration	Less penetration
More frequency of	Less frequency o
≈1 0° <i>Hz</i>	≈1 0 ⁶ Hz
Lesser wavelengÅt	More wavelengÅt h-
-4Å)	10ÅD)

Properties of X-Rays

- (1) X-rays are electromagneti rangeÅ 0-1.01Å0.
- (2) The wavelength of X-rays is (This is the only difference betwee get for med.
 - (3) X-raysare invisible.
 - (4) They travel in a straight li

- (5) X-rays are measured in Ront ger power).
- (6) X-rays carry no charge so the magnetic field and electric field.

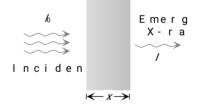
$$(7)\lambda_{Gama\ rays} < \lambda_{X-rays} < \lambda_{UV\ rays}$$

- (8) They used in the study of crysta
- (9) They ionise gases
- (10) X-rays do not pass through hea
- (11) They affect photographic pla
- (12) Long exposure to X-raysisinj
- (13) Leadisthebest absorber of X-
- (14) For X-ray photograph BaSO4 fi shuma the best absorber.
 - (15) They produce photoelectrice
 - (16) X-raysare not emitted by hydr
- (17) These cannot be used in Radar reflected by the target.
- (18) They showall the important pr reflection, refraction, interfere

X-rays are absorbed when they inci Intensity oXf reamye≔sr/₀gg¯é^xnt

So intensity of absorbed X-rays $I' = I_0 - I = I_0 (1 - e^{-\mu x})$

where thickness of absorp mpetdii coefficient



 $\mu \propto \lambda^3$; (λ = Wavelength of X-ray)

 $\mu \propto \overline{\nu}^3 \ (\nu = \text{Frequency of X-ray})$

 $\mu \propto Z^4 (Z = Atomic number of target)$

Classification of X-Rays

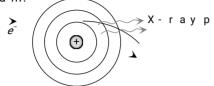
In X-ray tube, when high speed ele they penetrate the target. They los comestorest inside the metal. The e stopped makes several collisions wi the wavelength of light. Hence thAt each collision one of the followi

- (1) Continuous X-rays
- (2) Characteristic X-rays

Continuous X-Rays

As an electron passes close to t of the target, the electron is def figure. This results in decelera of X-rays.

spectrum.



(1M) i n i mum wa v e Whe en ig the electror whole of it's energy in a single co photon of maxim/nuahaniesneemii.geyted

$$\frac{1}{2} m v^2 = eV = h v_{\text{max}} = \frac{hc}{\lambda_{\text{min}}}$$

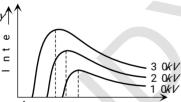
wherve velocity of electron befatom/, = potential difference thr accelear=ast peede, d of $k = 1 \frac{\partial y}{\partial x} = 3$

Maximum frequency of $r_{Max}d = \frac{eV}{a}t$ i on

Minimumwavelength = cut off way

$$\lambda_{\min} = \frac{hc}{eV} = \frac{12375}{V} \, \text{Å}$$

(21) ntensity wavelTehnegtchongtriampuho spectra consist of all the waveler wavelength are of different inte the intensity variation of diff accelerating voltages applied to



For each voltage, theintensit minimum wav & hin)e.n & t is e(s rapidly to then drops gradually.

The wavelength at which the depends on the accelerating volt voltage and vice-versa.

Characteristic X-Rays

Few of the fast moving electro penetrate the surface atoms of th out the tightly bound electrons e of the atom. Now when the electron created at that place.

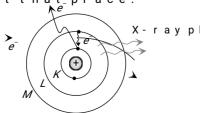
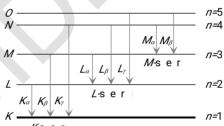


Fig. 25

To fill this vacancy electrons fro the created vacancies, we know that energy of the electron during dec≀ from a higher Ænteorlgoywoerrbeint Æn, roig nt ora bdii ta enerÆgyrÆg). Thus this energy differen The X-ray photons so emitted form of X-rays of very small but de depends upon the target material. Th of sharplines and is called characte

> (1 K, L, M, series: If the electron strik eject an elect Krosnh felb mofhtehe atom, a created Kisn htehlel. Immediately an elec outer shé-Islh, estal yjum/pisshteoltlh, ee mittin photon of energy equal to the energy two shells. Si milarly Misfhæni eljeuomtprsot K-shell, X-ray photon of higher ene photons emitted due to the juL, mL/p Nofel shell s to the KK₂, K₂hKellilnsegsioK/fostsehreie s of spectrum.



Kser If the electron sometimes of the targe L-shell of the target ato MM, Na.n. es heecltlr jumps t 🕭 **she**ll so that X-ray sphotons emitted.

The sephotons L-fsoer rmite be eof the spec s i milar way the M7 soermiaNetssie,ornie/teasf may be explained.

(21) n t e n s i t y - w a v e lt e c ne gr tt ha ġ n as ph ha r p l wavelengths, the intensity of K_{α_i} ray $\mathcal{K}_{\!\scriptscriptstyleeta}$ as shown in figure. These X-r characteristic X-rays. At other wav gradually and these X-rays are calle

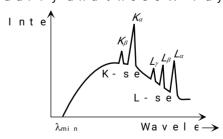


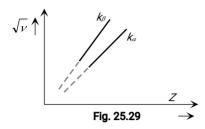
Fig. 25 Mosley's Law

Mosley studied the X-orhaayr **sapeetis** to number of a heavy elements and concl different elements are very similar

1 OE lectron, Photon, Photoelectric Effect and X-Rays

number, the spectral lines me r frequencies.

He also gave the fo $\sqrt{1}$ =0a $\sqrt{2}$ i- $\sqrt{2}$ i g relat



wherv = Frequency of Z = miAt bendd in ore targaetP, roportionab i tSycceesitag tc,(Shi el ding constant.

(Z-b) = Effective atomic number aan od doesn't depend on the natu valuebsaroefas follows

$$b=1$$
 for K-series
 $b=7.4$ fo L-series
 $b=19.2$ foM-series

- (1) Mosley's law supported Bohr
- (2) It experimentally deterZ)miof elements.
- (3) This law established the i elements in periodic table by ato weight.
- (4) Gaps in Mose Al e y 4 '3 \$ 66 at, a 7 f2 o r 7 5 existence of newel ements which we
- (5) The atomic Onu, uA gndo ne dPt we exorfeestal tobe 29, 47 and 78 respectively.
- (6) When a vacancy Koschoelrks, itn htehre electron remakishe byli.n Atmhel Leosthreold feel an effect Z-v1e)ed to te rt Zee+ ro 6 m the n and efrom the reAfmaheingelect t-osmh,ε 🗷 A photon is not a material pa orbit is well K-osuhteslildoertbhiet.
- Wave length $\frac{1}{\lambda} = R(Z - b)^2 \left(\frac{1}{n_s^2} - \frac{1}{n_s^2} \right)$ and energy of X-re

$$\Delta E = h \nu = \frac{hc}{\lambda} = Rhc \left(Z - b\right)^2 \left(\frac{1}{n_1^2} - \frac{1}{n_2^2}\right)$$

(8) If transition $p_2 \neq 22$ kt/nex0=s $1p/k1/(\alpha - a c e n f e e$

(i a)=
$$\sqrt{\frac{3 RC}{4}}$$
 = 2.47 × 10 ¹⁵ Hz

(i
$$\dot{V}_K$$
) = RC $(Z-1)^2 \left(1 - \frac{1}{2^2}\right) = \frac{3 RC}{4} (Z-1)^2$

$$= 2.47 \times 10^{15} (Z-1)^2 Hz$$

(iii) In general the Kwalvien lees nagrteh

$$\frac{1}{\lambda_{\nu}} = R(Z-1)^2 \left(1 - \frac{1}{n^2}\right)$$
 where 2, 3..., 4,

While
$$K_{\alpha}f$$
 bin $n\lambda_{R_{\alpha}}e_{\alpha}=\frac{1216}{\left(Z-1\right)^{2}}$ \mathring{A}

(i VE)
$$_{\alpha} = 10.2(Z-1)^2 eV$$

Uses of X-Ravs

- (i) Instudy of crystal structure: determined using X-ray diffraction.
 - (ii) In medical science
 - (iii) Inradiography
 - (iv) Inradiotherapy
 - (v) In engineering
 - (vi) In laboratories
 - (vii) In detective depart ment
- (viii) In art the change occurring exàmined by X-rays.

Tips & Tricks

- Discovery of positive rays isotopes.
- ★ The de-Broglie wavelength of of an atomis equal to circumfere
- 🗷 A particle having zero rest m and momentum must travels with a light.
- As de-Broglie wavelength associsgiv $\lambda = \frac{h}{mv_{ms}} = \frac{h}{\sqrt{3 \ mkT}}$ (Energy
- at tempe T i as $E = \frac{3}{2}kT$)
- energy.
- ★ When a particle exhibits wave with a wave packet, rather then a
- By coating the metal surface wi stronti umoxi de i t's work functi
- ★ We must remember that intent radiation is inversely proport between source of light an & i,pe.h,
- $I \propto \frac{1}{d^2}$ S O $I \propto I \propto \frac{1}{d^2}$)
- ★ The photoelectric current car some inert gas like Argon into t emitted by cathode i on ise the ga current is increased.
- ∠ Production of X-ray is the r
 ∠ photoelectric effect.

At The thickness of medium at whi X-rays beco $oldsymbol{a}$ m/= $\frac{I_0}{2}$ is scalled half v

 $(x_1)_2$ and it is $x_{1/2} = \frac{0.693}{\mu}$ is

∠Continuos X-rays are produce called "Bremsstrahlung". It me radiation.

Æ The wavelength of characterion accelerating voltage. It de⊅pof the target material.

Med Nearly all metals emits photo UV light. But alkali metals lik rubidium and cesium emit photexposed to visible light.

Ø Oxide coated filament in vacuel electrons at relatively lowert

Æ Conduction of electricity in because colliding electrons ac toincrease in mean free path.

★ Kinetic energy of cathode ray and work function of cathode.

∠ Photoelectric effect is due t

Hydrogen at omdoes not emit X-levels are too close to each othe

★ The essential difference bay
thaptr, ays emits from nucleus whi
of atom.

Act There is no time delay between incidence beft photoetloenctrons are en as the light falls on metal surfa

If light were wave (not photon to eject a photoelectron out of t

∠ Doze of X-ray are measured in t free energy via ionisait on.

est Safe doze for human body per w∈ Rontgon is the amount of X×1rÓa/¶ se energy through igmoaniirz at iNd nP of 1

As The photoelectrons emitted f have different kinetic energi photons have same energy. This electrons do not exist in the sur

Those coming from below the suingetting themselves free.

Æ Einstein was awarded Nobel p photoelectric effect. **£** Uncertainty in the measuren photon within $t \Delta p = \frac{h}{2\pi d}$ cleus is

wher de= diameter of the Δ are αd = euncertainty in the measure ment