



# Rational Extension of Anisotropic Harmonic Oscillator Potentials in Higher Dimensions

Rajesh Kumar<sup>1,2</sup> Rajesh Kumar Yadav<sup>2</sup> Avinash Khare<sup>3</sup>

<sup>1</sup>Department of Physics, Model College, Dumka-814101, India

<sup>2</sup>Department of Physics, S. K. M. University, Dumka-814110, India

<sup>3</sup>Department of Physics, Savitribai Phule Pune University, Pune-411007, India



## Abstract

We present a first-order **supersymmetric rational extension (RE)** of the **anisotropic quantum harmonic oscillator (QAHO)** in **multiple dimensions**, covering full-line, half-line, and mixed configurations. The exact solutions are expressed in terms of exceptional orthogonal polynomials, and the extended potentials remain isospectral to the conventional QAHO.

## Introduction

The rational extension (RE) of the quantum anisotropic harmonic oscillator (QAHO) in higher dimensions remains unexplored in the literature. Additionally, few studies have extended supersymmetry (SUSY) to higher dimensions [1-5]. This work aims to bridge this gap by examining the RE of QAHO in two and higher dimensions.

## Supersymmetric (SUSY) Formalism

The formalism of SUSY involves two Hamiltonians that are related through **supercharge operators**. The main components are:

$$H - \epsilon = \begin{cases} A^\dagger A = -\frac{d^2}{dx^2} + V^-(x) - \epsilon, & \text{say } H^- \\ AA^\dagger = -\frac{d^2}{dx^2} + V^+(x) - \epsilon, & \text{say } H^+ \end{cases} \quad (1)$$

These Hamiltonians represent the **supersymmetric partners** of a quantum system, where:

- **Factorization:** The Hamiltonians are factored into operators  $A^\dagger$  and  $A$ .
- **Partner Potentials:** The partner potentials are related by the **superpotential**  $W(x)$ , with:

$$V^-(x) = V^+(x) + W'(x)$$

- **Energy Spectrum:** The two Hamiltonians share the same spectrum, except for the ground state energy.

The superpotential is defined in terms of seedless function having  $\epsilon_m$  eigen value in SE as

$$W(x) = \begin{cases} -\frac{d}{dx} \ln[\phi_m(x)], & \text{for } \epsilon_m = E_0 \\ +\frac{d}{dx} \ln[\phi_m(x)], & \text{for } \epsilon_m < E_0 \end{cases} \quad (2)$$

SUSY ensures that the system's **energy levels** are **isospectral**, except for the **lowest eigenvalue** (ground state). This structure allows for solvability and the generation of new quantum systems.

## SUSY in Multiple Dimension

In higher dimensions, the second derivative  $\frac{d^2}{dx^2}$  is replaced by the Laplacian  $\nabla^2$ , and the scalar superpotential  $W(x)$  becomes a vector superpotential  $\vec{W}$ . Linear operators are expressed in a frame-independent form [1,2]:

$$A = \hat{e}^+ \cdot (\vec{\nabla} + \vec{W}), \quad \hat{e}^+ = \sum_{i=1}^D c_i \hat{e}_i \quad (3)$$

$$A^\dagger = (\vec{\nabla}^\dagger + \vec{W}^\dagger) \cdot \hat{e}^{+\dagger}, \quad \hat{e}^{+\dagger} = \sum_{i=1}^D c_i^\dagger \hat{e}_i \quad (4)$$

The SUSY partner potentials in Higher dimension The  $m$ -dependent SUSY partner potential  $V^-(x, y, \dots, m_1, m_2, \dots)$  will then be given by

$$V^- = V^+(x_1, x_2, \dots) - 2 \sum_k^D \left| \frac{\partial^2 \phi_k(x_k)}{\partial x_k^2} \right| + Q_D \quad (5)$$

The superpotential is

$$W_k = - \left| \frac{\partial}{\partial x_k} \ln[\phi_{m_k}(x_k)] \right| \quad (6)$$

$Q_D$  for for  $D = 1$  is 0 and for  $D = 2$  is

$$Q_2 = 2i \left( W_x \frac{\partial}{\partial y} - W_y \frac{\partial}{\partial x} \right) \prod_k \phi_{m_k}(x_k) \quad (7)$$

## How Does SUSY Partner Potential Differ?

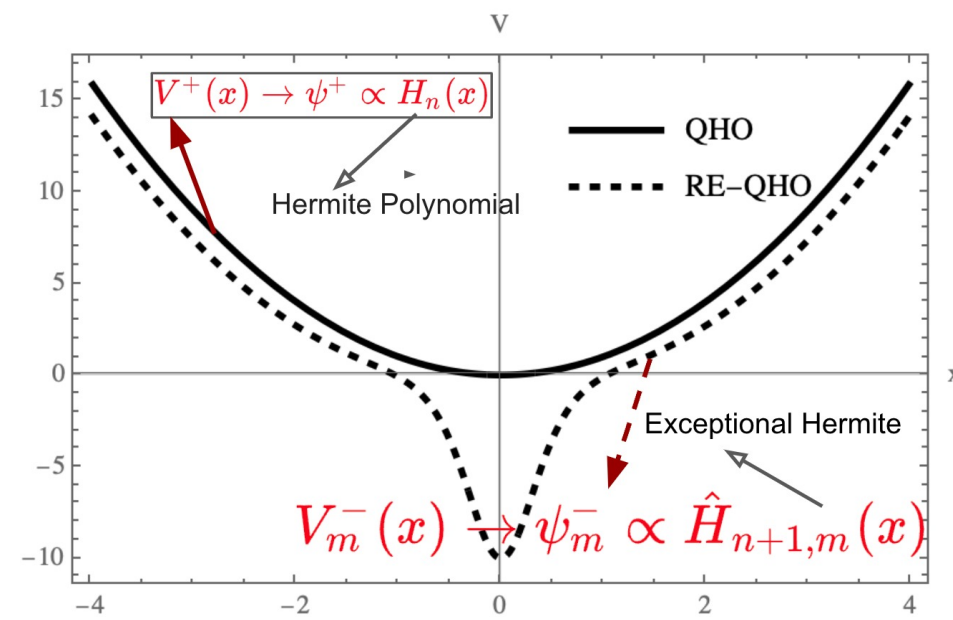


Figure 1. The solid line represents the potential of the quantum harmonic oscillator (QHO), while the dashed line represents its supersymmetric rational extension (RE-QHO).

Following combinations of potentials are used to construct SUSY partner in 2D and 3D

1. **Full-line:**

$$V^+(x) = \frac{1}{4} \omega_x^2 x^2, \quad -\infty < x < \infty, \quad (8)$$

Partner potential of  $V^+(x)$  is shown in figure above.

2. **Half-Line:**

$$V_h^+(x) = \begin{cases} \frac{1}{4} \omega_x^2 x^2 & , \quad x > 0 \\ \infty & , \quad x \leq 0. \end{cases} \quad (9)$$

There are two possible solutions, but only one is physically acceptable. However, the seedless function constructed using both solutions leads to two different SUSY potentials, given by

$$V_h^-(x, m, \alpha = \pm 1/2)$$

Table below list the expression for few  $m$  with energies  $E_{h,n,m}^-(\alpha) = E_{h,n,m}^+(\alpha) = 2(n + m + 1 + \alpha)\omega_x$

m	$V_h^-(x, m, -\frac{1}{2})$	$V_h^-(x, m, \frac{1}{2})$
0	$\frac{x^2 \omega_x^2}{4} - \omega_x$	$\frac{x^2 \omega_x^2}{4} + \frac{2}{x^2} - \omega_x$
1	$\frac{x^2 \omega_x^2}{4} + \frac{4\omega_x}{x^2 \omega_x + 1} - \frac{8\omega_x}{(x^2 \omega_x + 1)^2} - \omega_x$	$\frac{x^2 \omega_x^2}{4} + \frac{4\omega_x}{x^2 \omega_x + 3} - \frac{24\omega_x}{(x^2 \omega_x + 3)^2} + \frac{2}{x^2} - \omega_x$
2	$\frac{x^2 \omega_x^2}{4} + \frac{192x^2 \omega_x^2}{(x^4 \omega_x^2 + 6x^2 \omega_x + 3)^2} + \frac{8(x^2 \omega_x^2 - 3\omega_x)}{x^4 \omega_x^2 + 6x^2 \omega_x + 3} - \omega_x$	$\frac{x^2 \omega_x^2}{4} + \frac{320x^2 \omega_x^2}{(x^4 \omega_x^2 + 10x^2 \omega_x + 15)^2} + \frac{8(x^2 \omega_x^2 - 5\omega_x)}{x^4 \omega_x^2 + 10x^2 \omega_x + 15} + \frac{2}{x^2} - \omega_x$

Table 1. Potentials  $V_h^-(x, m, \alpha)$  (which are only valid for  $x > 0$ ) for different  $m$  when  $\alpha$  equals  $-\frac{1}{2}$  and  $\frac{1}{2}$  respectively.

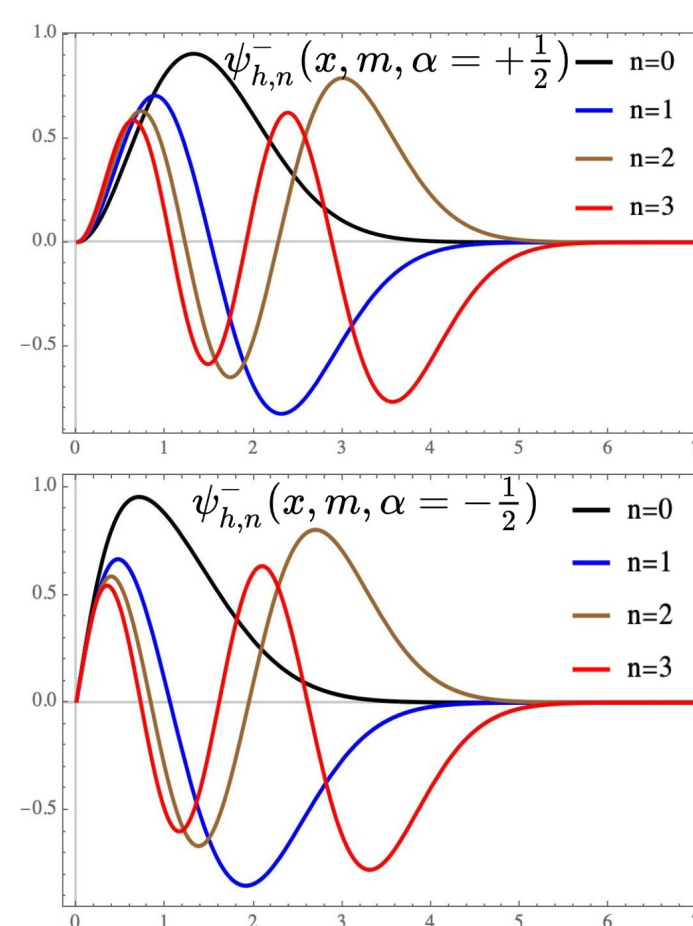


Figure 2. First few eigenstates for both RE-potentials.

## Rational Extension of Anisotropic QHO

There are three possible rational extensions of the 2D anisotropic harmonic oscillator, namely

1. 2D-Full-line oscillator

$$V^+(x, y) = \frac{1}{4} \omega_x^2 x^2 + \frac{1}{4} \omega_y^2 y^2; \quad -\infty < x < \infty, -\infty < y < \infty$$

2. 2D-Truncated oscillator and

$$V_h^+(x, y) = \begin{cases} \frac{1}{4} \omega_x^2 x^2 + \frac{1}{4} \omega_y^2 y^2 & \text{for } (x, y) \in (0, \infty) \times (0, \infty) \\ \infty & \text{otherwise} \end{cases}$$

3. combination of 1D-Full-line and 1D-Half-line oscillator.

$$V_{fh}^+(x, y) = \begin{cases} \frac{1}{4} \omega_x^2 x^2 + \frac{1}{4} \omega_y^2 y^2 & \text{for } y > 0; -\infty < x < \infty, \\ \infty & \text{for } y \leq 0 \end{cases}$$

## Results

For the above type of potential  $Q_D = 0$  for any dimension  $D$  so general form of SUSY partner potential, its wavefunction and energy eigenvalues are

$$V^-(x_1, x_2, \dots, m_1, m_2, \dots) = \sum_{n=1}^D V^-(x_n, m_n) \quad (10)$$

$$\psi_{n_1, n_2, \dots}^-(x_1, x_2, \dots, m_1, m_2, \dots) = \prod_k \psi_{n_k}^-(x_k, m_k) \quad (11)$$

$$E^-(m_1, m_2, \dots) = \sum_k^D E^-(m_k) \quad (12)$$

## Conclusions

The rationally extended potential is the sum of individual extensions, the energy is additive, and the wave function is their product.

## Acknowledgement

The author sincerely acknowledges SKMU University for financial support and thanks Dr. Rajesh Kumar Yadav and Dr. Avinash Khare for valuable discussions and insights.

## References

1. Fernández C, David J and Fernández-García, Nicolás. "Higher-order supersymmetric quantum mechanics." *AIP Conference Proceedings*, vol. 744, pp. 236-273, 2004. American Institute of Physics.
2. Das, Ashok, Okubo, S, and Pernice, SA. "Higher-dimensional SUSY quantum mechanics." *Modern Physics Letters A*, vol. 12, no. 08, pp. 581-588, 1997. World Scientific.
3. Andrianov, A. A., Borisov, N. V., and Ioffe, M. V. "The factorization method and quantum systems with equivalent energy spectra." *Physics Letters A*, vol. 105, no. 1-2, pp. 19-22, 1984.
4. Khare, Avinash, and Maharana, Jnanadeva. "Supersymmetric quantum mechanics in one, two and three dimensions." *Nuclear Physics B*, vol. 244, no. 2, pp. 409-420, 1984.
5. Leblanc, Martin, Lozano, G., and Min, H. "Extended superconformal Galilean symmetry in Chern-Simons matter systems." *Annals of Physics*, vol. 219, no. 2, pp. 328-348, 1992.
6. Rajesh Kumar, Rajesh Kumar Yadav, and Avinash Khare. "Rational Extension of Anisotropic Harmonic Oscillator Potentials in Higher Dimensions." *arXiv preprint*, arXiv:2411.02955, November 2024. Available at: arXiv:2411.02955.



Figure 3. Full Paper QR Code. Scan to access the full paper.