

$$h_k = \frac{1}{\sqrt{2 \sinh(\frac{\pi \omega}{a})}} (e^{\pi \omega / 2a} g_k + e^{-\pi \omega / 2a} g_{-k})$$

$$h_k^{(2)} = \frac{1}{\sqrt{2 \sinh(\frac{\pi \omega}{a})}} (e^{\pi \omega / 2a} \underline{g_k^{(2)}} + e^{-\pi \omega / 2a} \underline{g_{-k}^{(1)*}})$$

Unruh modes!! \rightarrow Some linear combination of +ve and -ve frequency

Imp

Rindler modes!!

We will see that Unruh modes can be written as some linear combination of +ve frequency Minkowskian modes (bez of our choice of branch cut!!)

~~$h_k^{(1,2)} \sim e^{i\omega(t \pm x)}$~~

$$\underline{h_k^{(1)}} \sim e^{i\omega(t+x)} e^{i\omega/a}$$

$$\underline{h_k^{(2)}} \sim (t+x)^{i\omega/a}$$

\leftarrow There will be some factors of ω !!

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Now, consider that t, x can be complex!!

After introducing the branch cut on my complex $(t-x)$ plane,

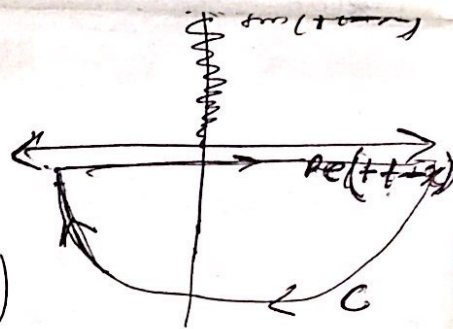
the func. is analytic and bounded everywhere else!!

We choose this to be in

the UHP \rightarrow we will see why!!

h_k mode is analytic
and bounded in LHP

Let's evaluate: $I = \oint h_k^{(1)} e^{-i\omega' z} dz$



$|\omega'| > 0$

$= 2\pi i \times \text{Res}(h_k^{(1)}) = 0$

So that the integrand vanishes of the big contour C

$\Rightarrow \int_{-\infty}^{\infty} h_k^{(1)} e^{-i\omega' z} dz + \int_C h_k^{(1)} e^{-i\omega' z} dz = 0$

bcz $e^{-i\omega' z}$ exponentially decrease for $\text{Im}(z) < 0$

$\Rightarrow \int_{-\infty}^{\infty} h_k^{(1)} e^{-i\omega'(z)} dz = 0$

$(t+ix)$ is real on this line !!

To enforce this;
we would require

$h_k^{(1)} = \int_0^{\infty} \chi(\omega) e^{-i\omega z} d\omega$ s.t. $\omega > 0$

Some linear combination of ve frequency (right-moving) Minkowskian modes !!

How is this enforced by I?

How is Evaluated this on the contour C again?

bcz then $\int_0^{\infty} \int_{-\infty}^{\infty} \chi(\omega) e^{-i(\omega+\omega')z} d\omega dz$

$\sum_{\omega} \int_{-\infty}^{\infty} \chi(\omega) \left\{ \cos(\omega+\omega')z - i \sin(\omega+\omega')z \right\}$