

PHYS 630 - Advanced Electricity and Magnetism  
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## Homework 7

**Problem 1**

Estimate the deflection angle for an electron of energy  $\epsilon = 1$  keV passing at a distance  $\ell = 1 \mu\text{m}$  from the proton.

*Solution.* The deflection angle  $\theta$  for a charged particle in a Coulomb field can be estimated using the formula:

$$\theta \approx \frac{2b}{\ell}.$$

where  $b$  is the impact parameter given by  $b = \frac{k|q_1 q_2|}{mv^2}$ .  
For the velocity  $v$ , we use the energy which is

$$\epsilon = \frac{1}{2}mv^2 \implies v = \sqrt{\frac{2\epsilon}{m}} = \sqrt{\frac{2 \times 1.60 \times 10^{-16}}{9.11 \times 10^{-31}}} \approx 1.87 \times 10^7 \text{ m/s}.$$

Calculating  $b$ , we have

$$b = \frac{ke^2}{mv^2} = \frac{(8.99 \times 10^9)(1.60 \times 10^{-19})^2}{(9.11 \times 10^{-31})(1.87 \times 10^7)^2} \approx 7.224 \times 10^{-13} \text{ m}.$$

Finally, calculating  $\theta$ , we have

$$\theta \approx \frac{2b}{\ell} = \frac{2(7.224 \times 10^{-13})}{10^{-6}} \approx 1.44 \times 10^{-6} \text{ radians}.$$

Therefore, the deflection angle is approximately  $1.44 \times 10^{-6}$  radians or about 0.297 arcseconds.

This small angle makes sense because the electron's energy is relatively high (1 keV) and the distance of closest approach is large ( $1 \mu\text{m}$ ). ■

**Problem 2**

It's a bright sunny day. The Sun is aligned with the zenith. You are holding a parasol of radius  $r = 50$  cm. Find the force from the Sun (assume complete absorption). Look up luminosity of the Sun and distance to the Sun.

*Solution.* The luminosity of the sun is  $L_{\odot} = 3.828 \times 10^{26} \text{ W}$  and the distance to the sun is  $1 \text{ AU} = 1.496 \times 10^8 \text{ km}$ .

The solar irradiance at Earth's distance is

$$I = \frac{L_{\odot}}{4\pi d^2} = \frac{3.828 \times 10^{26}}{4\pi(1.496 \times 10^{11})^2} \approx 1361 \text{ W/m}^2.$$

This value is known as the solar constant.

The area of the parasol is

$$A = \pi r^2 = \pi(0.5)^2 = 0.785 \text{ m}^2.$$

The power received by the parasol is

$$P = IA = (1361)(0.785) = 1068.4 \text{ W}.$$

Assuming complete absorption, the force from the Sun on the parasol is

$$F = \frac{P}{c} = \frac{1068.4}{2.998 \times 10^8} = 3.56 \times 10^{-6} \text{ N}.$$

Therefore, the Sun exerts a force of approximately  $3.56 \mu\text{N}$  on the parasol.

This is a very small force, roughly equivalent to the weight of a few micrograms on Earth, but it is responsible for important effects like radiation pressure in space and comet tail formation. ■

**Problem 3**

An electromagnetic wave is propagating along  $z$  and is polarized along  $x$ , so that its vector potential is

$$\mathbf{A} = A_0 (\cos(\omega t - k_z z), 0, 0),$$

where  $\omega = k_z c$ .

Calculate the averaged Poynting flux (over period of oscillation). Represent this linearly polarized wave as a sum of two circularly polarized waves and write the vector potential for  $\mathbf{A}_R$  and  $\mathbf{A}_L$ , the right and left-polarized components. There are media where right and left circular polarizations propagate with different speeds (called gyrotropic). After passing through such a media one of the waves (does not matter which, right or left) is delayed in phase by  $\pi$  (so that it becomes  $\propto \omega t - k_z z + \pi$ ). Find the new polarization after the wave leaves that gyrotropic medium.

*Solution.* We first calculate the fields and Poynting flux for the initial wave. From  $\mathbf{A}$  we can find  $\mathbf{B}$  and  $\mathbf{E}$ , as follows

$$\begin{aligned}\mathbf{B} &= \nabla \times \mathbf{A} = A_0 (0, -k_z \cos(\omega t - k_z z), 0), \\ \mathbf{E} &= -\frac{\partial \mathbf{A}}{\partial t} = A_0 \omega (\sin(\omega t - k_z z), 0, 0).\end{aligned}$$

The Poynting vector is given by

$$\mathbf{S} = \frac{1}{\mu_0} \mathbf{E} \times \mathbf{B}$$

Substituting, we get

$$\mathbf{S} = \frac{A_0^2 \omega k_z}{\mu_0} (0, 0, \sin(\omega t - k_z z) \cos(\omega t - k_z z)).$$

Time averaging, using  $\langle \sin(\alpha) \cos(\alpha) \rangle = \frac{1}{2}$ , we have

$$\langle \mathbf{S} \rangle = \frac{A_0^2 \omega k_z}{2\mu_0} (0, 0, 1)$$

Now for circular polarization decomposition: A linearly polarized wave can be represented as a sum of right and left circular polarizations:

$$\begin{aligned}\mathbf{A}_R &= \frac{A_0}{2} (\cos(\omega t - k_z z), -i \cos(\omega t - k_z z), 0), \\ \mathbf{A}_L &= \frac{A_0}{2} (\cos(\omega t - k_z z), i \cos(\omega t - k_z z), 0).\end{aligned}$$

We can check that  $\mathbf{A}_R + \mathbf{A}_L = \mathbf{A}$ , since the imaginary parts cancel.

After passing through the gyrotropic medium, one component gets phase shifted by  $\pi$ : If right component is shifted:

$$\begin{aligned}\mathbf{A}'_R &= \frac{A_0}{2} (\cos(\omega t - k_z z + \pi), -i \cos(\omega t - k_z z + \pi), 0), \\ \mathbf{A}'_L &= \frac{A_0}{2} (\cos(\omega t - k_z z), i \cos(\omega t - k_z z), 0).\end{aligned}$$

The resulting wave is

$$\mathbf{A}' = \mathbf{A}'_R + \mathbf{A}'_L = A_0 (0, \cos(\omega t - k_z z), 0).$$

The resulting wave is linearly polarized along  $y$  axis. This is an interesting result. The gyrotropic medium has rotated the polarization by 90 degrees. This is the basic principle behind optical isolators and other devices that manipulate polarization states. ■

**Problem 4**

An electromagnetic wave  $\mathbf{A}_1$  is propagating along  $z$  and is polarized along  $x$ , so that its vector potential is

$$\mathbf{A}_1 = A_0 (\cos(\omega t - k_z z), 0, 0),$$

where  $\omega = k_z c$ .

Another identical electromagnetic wave  $\mathbf{A}_2$  propagating along  $z$  and is polarized along  $x$ , is added. Calculate averaged Poynting flux (over period of oscillation) of the two waves together.

*Solution.* Let's calculate the total fields and Poynting flux for two identical waves. Starting with the vector potentials, we have

$$\mathbf{A}_1 = A_0 (\cos(\omega t - k_z z), 0, 0),$$

$$\mathbf{A}_2 = A_0 (\cos(\omega t - k_z z), 0, 0).$$

The total vector potential is

$$\mathbf{A}_{tot} = \mathbf{A}_1 + \mathbf{A}_2 = 2A_0 (\cos(\omega t - k_z z), 0, 0).$$

From this, we can find  $\mathbf{B}$  and  $\mathbf{E}$ , given by

$$\mathbf{B}_{tot} = \nabla \times \mathbf{A}_{tot} = 2A_0 (0, -k_z \cos(\omega t - k_z z), 0),$$

$$\mathbf{E}_{tot} = -\frac{\partial \mathbf{A}_{tot}}{\partial t} = 2A_0 \omega (\sin(\omega t - k_z z), 0, 0).$$

The Poynting vector is given by

$$\mathbf{S}_{tot} = \frac{1}{\mu_0} \mathbf{E}_{tot} \times \mathbf{B}_{tot}.$$

Substituting, we have

$$\mathbf{S}_{tot} = \frac{4A_0^2 \omega k_z}{\mu_0} (0, 0, \sin(\omega t - k_z z) \cos(\omega t - k_z z)).$$

Time averaging, using  $\langle \sin(\alpha) \cos(\alpha) \rangle = \frac{1}{2}$ , we have

$$\langle \mathbf{S}_{tot} \rangle = \frac{2A_0^2 \omega k_z}{\mu_0} (0, 0, 1).$$

Note that

$$\langle \mathbf{S}_{tot} \rangle = 4 \langle \mathbf{S}_1 \rangle = 4 \langle \mathbf{S}_2 \rangle.$$

The total Poynting flux is four times larger than the flux of each individual wave. This is because the amplitudes of  $\mathbf{E}$  and  $\mathbf{B}$  fields doubled and because the Poynting flux is proportional to the product of  $\mathbf{E}$  and  $\mathbf{B}$ , hence the factor of 4. This demonstrates constructive interference of two identical waves leading to intensity enhancement. ■