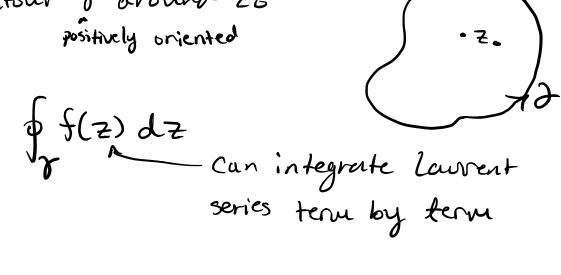
Residue Theory (Contour Integrals)
Suppose f(z) has an isolated singularity at Zo.

$$f(z) = \sum_{n=-\infty}^{+\infty} a_n (z-z_0)^n$$
 Laurent Series about 2.

Have a contour of around 20 positively oriented

Consider



The point: 
$$\iint_{\mathcal{T}} \{f(z) dz = \int_{n=-\infty}^{+\infty} a_n \oint_{\mathcal{T}} (z-z_0)^n dz$$

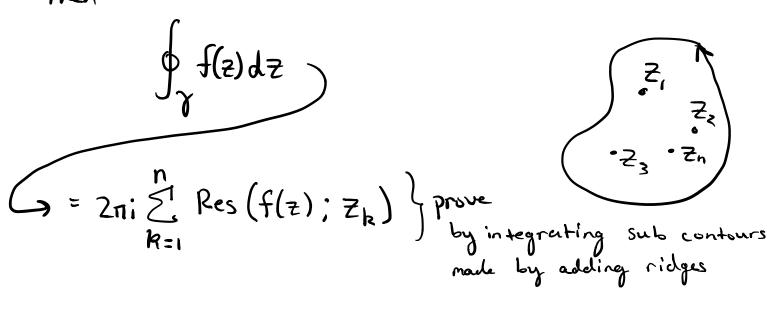
$$\int_{0}^{\infty} (z-z_{0})^{n} dz = \begin{cases} 0, & n \neq -1 \\ 2\pi i, & n = -1 \end{cases}$$

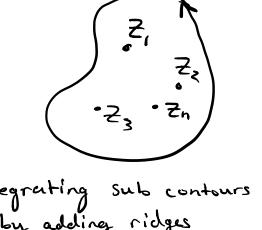
And the a., coefficient is called the residue of z at Zo.

$$a_{-1} = \text{Res}(f(z); z_{\circ})$$

So Calculating the contour boils down to finding the residue.

Suppose we had more than one singularity; Say, Z.,..., Zn (finite amount)





How to Calculate Residues

$$f(z) = z^3 e^{6/z} = z^3 \left[1 + \frac{6}{z} + \frac{1}{2!} \left(\frac{6}{z}\right)^2 + \frac{1}{3!} \left(\frac{6}{z}\right)^5 + \cdots\right]$$

2) If a simple pole at Zo,

$$f(z) = -(a_1)\frac{f(z)}{z-z_0} + q_0 + q_1(z-z_0) - \cdots$$

All that's left is the residue!

Example:

Res
$$\left(\frac{e^{z}}{\sin z}; \pi\right)$$

Sinz has a simple pole at Z=TI

Res = 
$$\lim_{z \to \pi} (z - \pi) e^{z}$$
 L'hospitals!  $\lim_{z \to \pi} (z - \pi) e^{z} + e^{z}$ 

3) If a pole of order m, then

Res 
$$(f(z); z_0) = \lim_{z \to z_0} \frac{1}{(m-1)!} \left(\frac{d}{dz}\right)^{m-1} (z-z_0)^m f(z)$$

Proof: Use laurent series - Page 310

Using Residues to find Trigonometric Integrals on [0,21]

xample 
$$I = Integral = \int_0^{2\pi} \frac{d\theta}{a + \cos \theta}$$
 where a >1 technically a function of a low finite

Let 
$$z=e^{i\theta}$$
  $\longrightarrow$   $dz=e^{i\vartheta}id\theta$ 

$$d\theta = -i\frac{dz}{e^{i\theta}} = -i\frac{dz}{z}$$

Then

$$\cos \theta = \frac{z + \frac{1}{2}}{2}$$

This converts I to an integral around unit circle

$$I = \int_{|z|=1}^{\frac{-i}{2}} \frac{dz}{a+z+\frac{1}{2}} = \int_{|z|=1}^{\frac{-2i}{2}+2az+1} dz$$

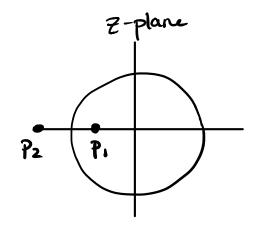
inside unit circle?

quadratic has two

$$-\frac{2a}{2} + \frac{\sqrt{4a^2 - 4}}{2} = -a + \sqrt{a^2 - 1}$$

Poles: 
$$P_1 = -at \sqrt{a^2-1}$$

$$Pz = -\alpha - \sqrt{\alpha^2 - 1}$$



So only residue from p. contributes to I.

I = 
$$2\pi i \leq \text{Res} \left(\text{singularities inside } \Upsilon\right)$$
  
=  $2\pi i \text{Res} \left[\frac{2}{i} \frac{1}{7^2 + 3c^2 + 1}; P_i\right]$ 

$$P_1 - P_2 = 2 - \overline{a^2 - 1}$$

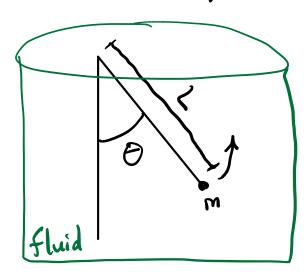
$$= \frac{4\pi}{P_1 - P_2} = \frac{2\pi}{\sqrt{\alpha^2 - 1}}$$

Wowza!

$$T(a) = \int_{0}^{2\pi} \frac{d\theta}{a + \cos \theta} = \frac{2\pi}{\sqrt{a^{2} - 1}}$$
 This is a neat

## Physics Application

Period of whirling, overdamped pendulum driven by a constant torque.



torque It (constant)

the torque will cause this to "whirl" in the sense that it will refate all around.

$$\frac{d}{dt}(ml+L\dot{\theta}) + b\theta' + mglsin(\theta) = \Gamma^{1}$$

damping so large, inertia term doesn't apply

$$T = \int dt = \int_0^{2\pi} \frac{d\theta}{\alpha - \sin\theta}, \quad \frac{d\theta}{dt} + \sin\theta = \alpha$$

$$dt = \begin{bmatrix} a - \sin\theta \\ d\theta \end{bmatrix}$$