NUMERICAL METHODS FOR SOLVING WAVE EQUATION

Prepared By: Ramiz Mahmmud

ID: 18203035

MS Session: 2021-2022

INTRODUCTION TO WAVE EQUATION

A second order wave equation:

$$\frac{\partial^2 u}{\partial t^2} = c^2 \nabla^2 u$$

A second order 1-D wave equation:

$$\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$$

A first order 1-D wave equation also known as 1-D linear convection equation:

$$\frac{\partial u}{\partial t} + c \frac{\partial u}{\partial x} = 0$$

SOLUTION OF CONVECTION EQUATION

Consider the pure initial value problem,

$$\frac{\partial u}{\partial t} + c \frac{\partial u}{\partial x} = 0$$

Initial condition: $u(x,0) = f(x), -\infty < x < \infty$

Exact solution: u(x, t) = f(x - c t)

EULER EXPLICIT METHOD

Converting PDE into FDE (FTCS):

$$\frac{u_j^{n+1} - u_j^n}{\Delta t} + C\left(\frac{u_{j+1}^n - u_{j-1}^n}{2\Delta x}\right) = 0$$

or,
$$u_j^{n+1} = u_j^n - \frac{1}{2} \left(\frac{C\Delta t}{\Delta x} \right) \left(u_{j+1}^n - u_{j-1}^\eta \right)$$

Putting
$$\frac{C\Delta t}{\Delta x} = r$$
 we get,

Or,
$$u_j^{n+1} = u_j^n - \frac{1}{2}r(u_{j+1}^n - u_{j-1}^\eta)$$
 (1)

VON NEUMANN STABILITY ANALYSIS OF EULER EXPLICIT METHOD

Using Von Neumann stability analysis, we can put in (1)

$$u_i^n = G^n e^{ikj\Delta x}$$

Where G is growth factor, k is wave number and $i = \sqrt{-1}$

A method is stable if $|G| \leq 1$

Equation (1) reduces to,

$$G^{n+1}e^{ikj\Delta x} = G^n e^{ikj\Delta x} - \frac{r}{2} \left(G^n e^{ik(j+1)\Delta x} - G^n e^{ik(j-1)\Delta x} \right)$$

Or,
$$G = 1 - \frac{r}{2} \left(e^{ik\Delta x} - e^{-ik\Delta x} \right)$$

VON NEUMANN STABILITY ANALYSIS OF EULER EXPLICIT METHOD(CONT.)

Using the relation
$$e^{i\theta}-e^{-i\theta}=2i\sin\theta$$
 we get, $G=1-ir\sin(k\Delta x)$ So,
$$(|G|)^2=GG^*=\{1-ir\sin(k\Delta x)\}\{1+ir\sin(k\Delta x)\}$$

$$=1+r^2\sin^2(k\Delta x)$$
 So, $G\geq 1$

Thus Euler Explicit Method for convection equation is unconditionally unstable.

UPSTREAM (FIRST ORDER UPWIND OR WINDWARD) DIFFERENCING METHOD

Changing Forward space to Backward space in Euler explicit method:

Forward time backward space (FTBS):

$$\frac{u_j^{n+1} - u_j^n}{\Delta t} + c \frac{u_j^n - u_{j-1}^n}{\Delta x} = 0$$
Or,
$$u_j^{n+1} = (1 - r)u_j^n + ru_{j-1}^n$$
 (2)

This method is conditionally stable. Using Von Neumann stability analysis we get

$$0 \le r \le 1$$

LAX METHOD

Euler Explicit Method:
$$u_{j}^{n+1} = u_{j}^{n} - \frac{1}{2}r(u_{j+1}^{n} - u_{j-1}^{n})$$

Substituting $u_j^n = (u_{j+1}^n + u_{j-1}^n)/2$ in Euler Explicit equation we get,

$$u_{j+1}^{n} = \frac{1}{2} (u_{j+1}^{n} + u_{j-1}^{n}) - \frac{r}{2} (u_{j+1}^{n} - u_{j-1}^{n})$$

or,
$$u_j^n = \frac{1}{2} \{ (1-r)u_{j+1}^n + (1+r)u_{j-1}^n \}$$
 (3)

VON NEUMANN STABILITY ANALYSIS OF LAX METHOD

Using Von Neumann stability analysis we get,

$$G^{n+1}e^{ikj\Delta x} = \frac{1}{2} \left\{ (1-r)G^n e^{ik(j+1)\Delta x} + (1-r)G^n e^{ik(j-1)\Delta} \right\}$$

Or,
$$G = \frac{1}{2} \{ (1-r)e^{ik\Delta x} + (1+r)e^{-ik\Delta x} \}$$

Or,
$$G = \frac{1}{2} \{ (e^{ik\Delta x} + e^{-ik\Delta x}) - r(e^{ik\Delta x} - e^{-ik\Delta x}) \}$$

Or,
$$G = \cos(k\Delta x) - ir\sin(k\Delta x)$$

VON NEUMANN STABILITY ANALYSIS OF LAX METHOD (CONT.)

Now,
$$|G|^2 = \cos^2(k\Delta x) + r^2 \sin^2(k\Delta x)$$

Or,
$$|G|^2 = 1 - (1 - r^2) \sin^2(k\Delta x)$$

So, for stability,
$$1 - r^2 \ge 0$$
 or, $r \le 1$

So Lax method is stable for

$$r \le 1 \Rightarrow \frac{c\Delta t}{\Delta x} \le 1$$

Lax method is conditionally stable

EULER IMPLICIT METHOD

For implicit method Difference (FTCS) equation is:

$$\frac{u_j^{n+1} - u_j^n}{\Delta t} + C\left(\frac{u_{j+1}^{n+1} - u_{j-1}^{n+1}}{2\Delta x}\right) = 0$$

Or,

$$\frac{r}{2}u_{j+1}^{n+1} + u_j^{n+1} - \frac{r}{2}u_{j-1}^{n+1} = u_j^n$$
 (4)

Here only u_j^n is known. To find the others values we have to solve system of linear equation.

VON NEUMANN STABILITY ANALYSIS OF EULER IMPLICIT METHOD

Using Von Neumann stability analysis we get,

$$\frac{r}{2}G^{n+1}e^{ik(j+1)\Delta x} + G^{n+1}e^{ikj\Delta x} - \frac{r}{2}G^{n+1}e^{ik(j-1)\Delta x} = G^ne^{ijk\Delta x}$$

Or,
$$\frac{r}{2}Ge^{ik\Delta x} + G - \frac{r}{2}Ge^{ik\Delta x} = 1$$

Or,
$$G\left\{1+\left(e^{ik\Delta x}-e^{-ik\Delta x}\right)\frac{r}{2}\right\}=1$$

Or,
$$G\{1 + i r \sin(k\Delta x)\} = 1$$

Or,
$$G = \frac{1}{1 + ir \sin(k\Delta x)}$$

So,
$$|G|^2 = \frac{1}{1 + r^2 \sin^2(k\Delta x)} \le 1$$

Thus the Euler implicit method is unconditionally stable

LEAP FROG METHOD

Taking central time and central space difference formula we get equation for Leap Frog scheme.

$$\frac{u_j^{n+1} - u_j^{n-1}}{2\Delta t} + C \frac{u_{j+1}^n - u_{j-1}^n}{2\Delta x} = 0$$

Or,
$$u_j^{n+1} = u_j^{n-1} - r(u_{j+1}^n - u_{j-1}^n)$$
 (5)

Leap Frog method is stable for $|r| \leq 1$ and It is second order T.E scheme.

Leap Frog method is conditionally stable.

CRANK-NICOLSON METHOD

The difference equation for convection equation using Crank-Nicolson method is given

by:
$$\frac{u_j^{n+1} - u_j^n}{\Delta t} + c \frac{1}{2} \left[\frac{u_{j+1}^{n+1} - u_{j-1}^{n+1}}{2\Delta x} + \frac{u_{j+1}^n - u_{j-1}^n}{2\Delta x} \right] = 0$$

Or,
$$\frac{1}{4}r u_{j-1}^{n+1} - u_j^{n+1} - \frac{1}{4}r u_{j-1}^{n+1} = -u_j^n + \frac{1}{4}r \left(u_{j+1}^{\eta} - u_{j-1}^{n}\right)$$
 (6)

It is a second order accuracy method.

VON NEUMANN STABILITY ANALYSIS CRANK-NICOLSON METHOD

Using Von Neumann stability analysis, we get:

$$\frac{1}{4}r G^{n+1} e^{ik(j-1)\Delta x} - G^{n+1} e^{ikj\Delta x} - \frac{1}{4}r G^{n+1} e^{ik(j+1)\Delta x} = -G^n e^{ikj\Delta x} + \frac{1}{4}r(G^n e^{ik(j+1)\Delta x} - G^n e^{ik(j-1)\Delta x})$$

$$Or, -G\left(1 + \frac{r}{4}(e^{ik\Delta x} - e^{-ik\Delta x})\right) = -(1 - \frac{r}{4}(e^{ik\Delta x} - e^{-ik\Delta x}))$$

Or,
$$G = \frac{1 - \frac{ir}{2} \sin \phi}{1 + \frac{ir}{2} \sin \phi}$$
; where $\phi = k\Delta x$

VON NEUMANN STABILITY ANALYSIS FOR CRANK-NICOLSON METHOD(CONT.)

Or,
$$G = \frac{\left(1 - \frac{r^2 \sin^2 \phi}{4}\right) - ir \sin \phi}{1 + \frac{r^2 \sin^2 \phi}{4}}$$

So,
$$|G|^2 = \frac{1}{1 + \frac{r^2 \sin^2 \phi}{4}} \le 1$$

Thus Crank-Nicolson Method is unconditionally stable.

REFERENCE

- Computational Fluid Dynamics and Heat Transfer by John C. Tannehill, Dale A. Anderson, Richard H. Pletcher
- Computational Fluid Dynamics Volume I by Klaus A. Hoffmann, Steve T. Chiang

Thank You