

Equation of state & pair-correlations in mass-imbalanced one-dimensional Fermi systems



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Schleching, March 7, 2017.



- ▶ Motivation: strongly interacting quantum gases
- ▶ Method: Hybrid Monte Carlo
- ▶ Results: discussion of EOS & pair-correlations
 - ▶ connection to RG method
- ▶ Summary, conclusions & future work

Strong interactions in ultracold gases

Phases of strongly interacting theories

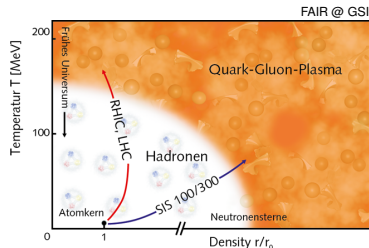


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- ▶ explore & understand phases of strongly-interacting theories
- ▶ universal behavior:
unitary fermi gas (UFG)
- ▶ ultracold quantum gases are a versatile tool to probe strongly-interacting systems (in experiment and theory)

A.N. Wenz *et al.*, *Science* **25**, 2013.

- ▶ many available methods, exchange with lattice QCD
- ▶ 1D often a benchmark for new methods



Strong interactions in ultracold gases

Phases of strongly interacting theories

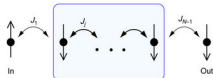


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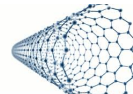
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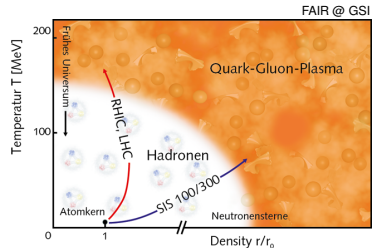
- ▶ many available methods, exchange with lattice QCD
- ▶ 1D often a benchmark for new methods, *but also interesting physics!*



O.V. Marchukov *et al.*, *Nature Comm.* **7**, 2016.



S. Iijima, *Nature* **354**, 1991.



Model:

Fermions on the lattice



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- Gaudin-Yang model for fermions in 1D

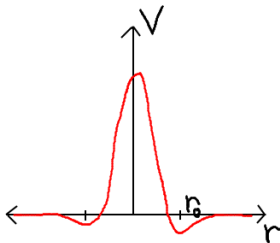
$$\hat{H} = - \sum_{s=\uparrow,\downarrow} \sum_x \hat{\psi}_s^\dagger(x) \frac{\hbar^2 \nabla^2}{2m_s} \hat{\psi}_s(x) + g \sum_x \hat{\psi}_\uparrow^\dagger(x) \hat{\psi}_\uparrow(x) \hat{\psi}_\downarrow^\dagger(x) \hat{\psi}_\downarrow(x)$$

Model:

Fermions on the lattice

- Gaudin-Yang model for fermions in 1D

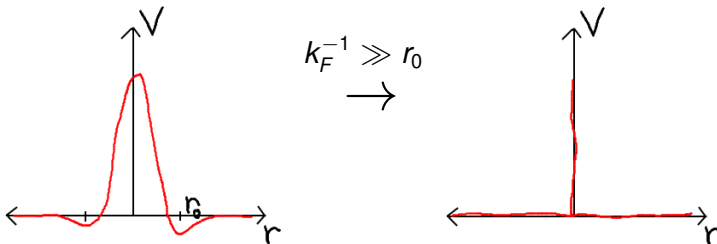
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Model: Fermions on the lattice

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Model:

Fermions on the lattice



- ▶ Gaudin-Yang model for fermions in 1D

$$\hat{H} = - \sum_{s=\uparrow,\downarrow} \sum_x \hat{\psi}_s^\dagger(x) \frac{\hbar^2 \nabla^2}{2m_s} \hat{\psi}_s(x) + g \sum_x \hat{\psi}_\uparrow^\dagger(x) \hat{\psi}_\uparrow(x) \hat{\psi}_\downarrow^\dagger(x) \hat{\psi}_\downarrow(x)$$

- ▶ problem: $\hat{H} = \hat{T} + \hat{V}$ not generally diagonalizable simultaneously
- ▶ in some cases exactly solvable with Bethe ansatz

M. Gaudin, *Phys. Lett. A* **24**, 1967.

C.N. Yang, *Phys. Rev. Lett.* **19**, 1967.

Method:

Hybrid Monte Carlo (HMC) in a nutshell

- ▶ ground-state partition function as a function of imaginary time

$$\mathcal{Z}(\beta) = \langle \psi_0(\beta) | \psi_0(\beta) \rangle = \langle \psi_0(0) | e^{-\beta \hat{H}} | \psi_0(0) \rangle$$

- ▶ rewrite with time-discretization & Hubbard-Stratonovich transform:
path-integral over *auxiliary field* σ

$$\mathcal{Z}(\beta) = \int \mathcal{D}\sigma \det U_{\sigma}^{(\uparrow)}(\beta) \det U_{\sigma}^{(\downarrow)}(\beta) = \int \mathcal{D}\sigma P[\sigma, \beta]$$

- ▶ instead of one very complicated many-body problem: many single-particle problems in an external field
- ▶ high-dimensional: integration with Monte Carlo
- ▶ "hybrid": smarter choice of new field-configurations (molecular dynamics)

S. Duane, A.D. Kennedy, B.J. Pendleton, D. Roweth, *Phys. Lett. B* **195**, 1987.

A. Bulgac, J.E. Drut, P. Magierski, *Phys. Rev. Lett.* **96**, 2006.



Method:

Hybrid Monte Carlo (HMC) in a nutshell



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When is it useful?

- ▶ transfer matrices $U_{\sigma}^{(s)}$ are equal for homogeneous case, otherwise: **sign-problem**

$$\mathcal{Z} = \int \mathcal{D}\sigma \det U_{\sigma}^{(\uparrow)} \det U_{\sigma}^{(\downarrow)} \equiv \int \mathcal{D}\sigma P[\sigma]$$

↓

$$\mathcal{Z} = \int \mathcal{D}\sigma \det^2 U_{\sigma}$$

- ▶ probability-measure positive-semidefinite

Results: Equation of state

ground-state energy vs. coupling strength

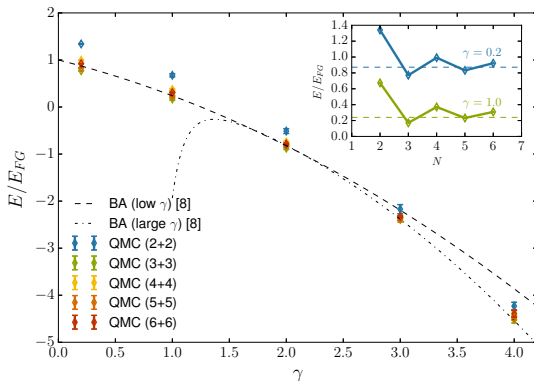
- ▶ ground-state energy experimentally accessible
- ▶ computed via

$$\langle \hat{H} \rangle = -\frac{\partial}{\partial \beta} \log \mathcal{Z}(\beta)$$

- ▶ excellent agreement with Bethe ansatz

M. Wadati, T. Iida,
Phys. Lett. A **360**, 2007.

- ▶ limit of large γ : bosonic behavior



LR, W.J. Porter, A.C. Loheac, and J.E. Drut, *Phys. Rev. A* **92**, 2015.

$$\gamma = g/n$$

Method II: Density functional theory + RG (DFT-RG)



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- ▶ Hohenberg-Kohn theorem: energy can be written as functional of the density
- ▶ ground-state energy via minimization

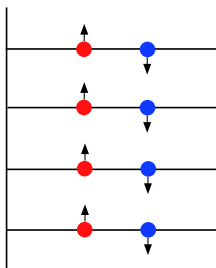
$$E_{\text{gs}} = \inf_{\rho} E[\rho]$$

- ▶ problem: no recipe for the energy functional
- ▶ typically: minimization of global ansatz for energy density functional
- ▶ DFT-RG: method to incorporate microscopic interactions

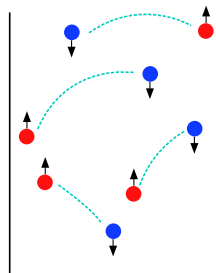
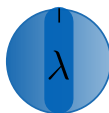
S. Kemler, J. Braun, *Journ. Phys. G* **40**, 2013.

S. Kemler, M. Pospiech and J. Braun, *Journ. Phys G* **44**, 2017.

Method II: Density functional theory + RG (DFT-RG)



$\lambda = 0$



$\lambda = 1$

$$\partial_{\lambda} \Gamma_{\lambda} = \frac{1}{2} \rho_{\sigma} U_{\sigma\sigma'} \rho_{\sigma'} + \frac{1}{2} \text{Tr} \left\{ U_{\sigma\sigma'} \left[\left(\Gamma_{\lambda, \sigma' \sigma}^{(2)} \right)^{-1} + \rho_{\sigma'} \cdot \mathbb{1}_{(\sigma, x)} \right] \right\}$$

Results: Equation of state

ground-state energy for few-body systems



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- ▶ ground-state energy experimentally accessible
- ▶ computed via

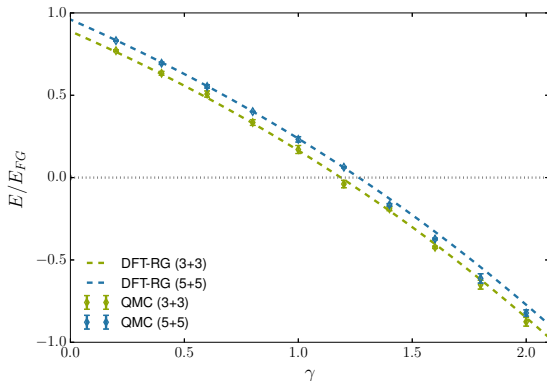
$$\langle \hat{H} \rangle = -\frac{\partial}{\partial \beta} \log \mathcal{Z}$$

- ▶ excellent agreement with Bethe ansatz

M. Wadati, T. Iida,
Phys. Lett. A **360**, 2007.

- ▶ limit of large γ : bosonic behavior
- ▶ perfect agreement with few-body DFT-RG

S. Kemler, M. Pospiech and J. Braun, *in preparation*



$$\gamma = g/n$$

Method revisited: Mass-imbalance and the sign problem



- ▶ typically subject to a **sign-problem**:
unequal masses \rightarrow unequal transfer matrices

$$\mathcal{Z} = \int \mathcal{D}\sigma \det U_{\sigma}^{(\uparrow)} \det U_{\sigma}^{(\downarrow)} = \int \mathcal{D}\sigma P[\sigma]$$



$$\mathcal{Z} = \int \mathcal{D}\sigma \det^2 U_{\sigma}$$

- ▶ exponentially increasing numerical effort with larger systems

Method revisited: Imaginary mass-imbalance to the rescue!



- ▶ sign-problem surmounted by rewriting

$$m_{\uparrow} = m_0(1 + i\delta m) \text{ and } m_{\downarrow} = m_0(1 - i\delta m)$$

J. Braun, J.E. Drut, D. Roscher., *Phys. Rev. Lett.* **114**, 2014.

$$\mathcal{Z} = \int \mathcal{D}\sigma \det U_{\sigma}^{(\uparrow)} \det U_{\sigma}^{(\downarrow)} \equiv \int \mathcal{D}\sigma P[\sigma]$$

↓

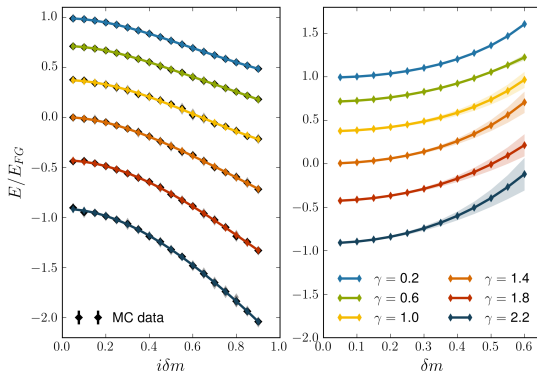
$$\mathcal{Z} = \int \mathcal{D}\sigma |\det U_{\sigma}|^2$$

- ▶ probability-measure again positive-semidefinite

Results: Equation of state

Mass-imbalanced systems

- ▶ experimental realization of ${}^6\text{Li}$ and ${}^{40}\text{K}$ mixture ($\delta m \sim 0.74$)
- ▶ real mass-imbalance via analytic continuation (Padé approximation)



LR, W.J. Porter, J. Braun, J.E. Drut, *in preparation*.

$$\gamma = g/n$$

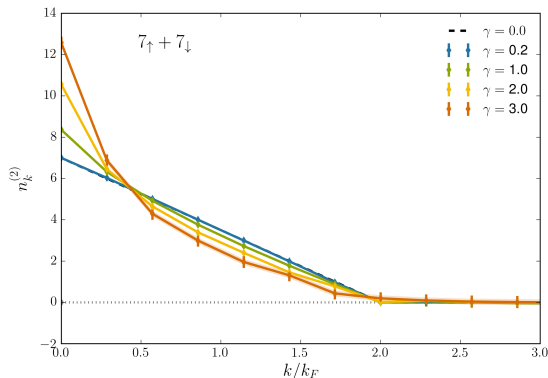
Results: Pair-correlation

Mass-balanced systems

- ▶ pairing properties through pair-correlation function

$$\rho_{(2)}(|x - x'|) = \langle \hat{\psi}_{\uparrow}^{\dagger}(x) \hat{\psi}_{\downarrow}^{\dagger}(x) \hat{\psi}_{\downarrow}(x') \hat{\psi}_{\uparrow}(x') \rangle$$

- ▶ smooth increase of peak at $k = 0$ with interaction strength
- ▶ formation of tightly-bound pairs
- ▶ crossover from BCS-pairing *Bose-Einstein condensate*



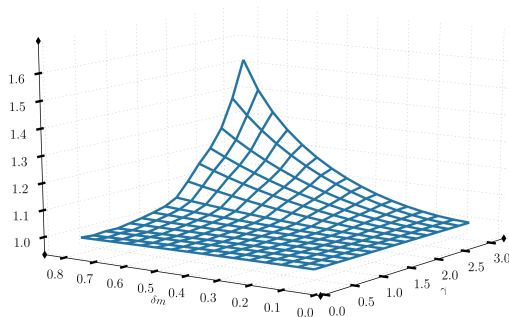
$$\gamma = g/n$$

LR, W.J. Porter, J. Braun, J.E. Drut, *in preparation*.

Results: Pair-correlation

Mass-imbalanced systems

- ▶ peak still at $k = 0$:
no FFLO-phase observed
- ▶ normalized peak height
to mass-balanced system:
influence of imbalance
larger at strong coupling



LR, W.J. Porter, J. Braun, J.E. Drut, *in preparation*.

$$\gamma = g/n$$

Results: Pair-correlation

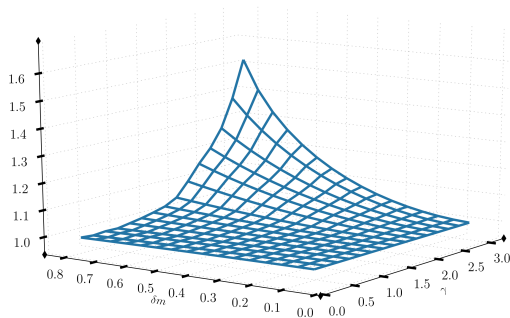
Mass-imbalanced systems



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PRELIMINARY



$$\gamma = g/n$$

LR, W.J. Porter, J. Braun, J.E. Drut, *in preparation*.

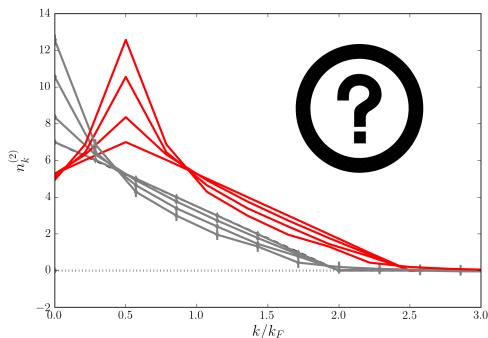
Outlook: Pair-correlation

Spin-imbalanced systems

- ▶ hallmark of Fulde-Ferrell-Larkin-Ovchinnikov phase:
peak at nonzero momentum $k = |k_F^{(\uparrow)} - k_F^{(\downarrow)}|$

P. Fulde, R.A. Ferrell,
Phys. Rev. **135**, 1964.

A.I. Larkin, Y.N. Ovchinnikov,
Sov. Phys. JETP **20**, 1965.





Summary

- ▶ precise calculation of energies and correlations
agreement for few-body systems, convergence to thermodynamic limit
- ▶ stable results for high mass-imbalances
- ▶ observed formation of bosonic pairs in the ground state
- ▶ no FFLO behavior in purely mass-imbalanced systems



Summary

- ▶ precise calculation of energies and correlations
agreement for few-body systems, convergence to thermodynamic limit
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Future work

- ▶ spin-imbalance (FFLO behavior expected)
- ▶ long-term goal: nonperturbative phase-diagrams of 3D unitary Fermi gas