

1.

Maxwell equations:

$$\begin{cases} \nabla \times \vec{E} = -j\omega \vec{B} & \text{-- ①} \\ \nabla \cdot \vec{D} = \rho & \text{-- ②} \\ \nabla \times \vec{H} = \vec{J} + j\omega \vec{D} & \text{-- ③} \\ \nabla \cdot \vec{B} = 0 & \text{-- ④} \end{cases}$$

由 ①, ④ 可延伸出

$$\begin{cases} \vec{B} = \nabla \times \vec{A} \\ \vec{E} + j\omega \vec{A} = -\nabla \phi \end{cases} \quad \text{代入 ③ 中有}$$

$$\frac{1}{\mu} \nabla \times (\nabla \times \vec{A}) = \vec{J} - j\omega \epsilon (\nabla \phi + j\omega \vec{A})$$

$$\nabla \times (\nabla \times \vec{A}) = \nabla (\nabla \cdot \vec{A}) - \nabla^2 \vec{A}$$

$$\Rightarrow \nabla^2 \vec{A} + \omega^2 \mu \epsilon \vec{A} = -\mu \vec{J} + \nabla (\nabla \cdot \vec{A} + j\omega \mu \epsilon \phi)$$

引入洛伦兹规范使得

$$\nabla \cdot \vec{A} = -j\omega \mu \epsilon \phi, \text{ 代入上式即有}$$

$$\nabla^2 \vec{A} + \omega^2 \mu \epsilon \vec{A} = -\mu \vec{J}$$

同理代入 ② 中有

$$-\nabla \cdot \epsilon (\nabla \phi + j\omega \vec{A}) = \rho$$

$$\Rightarrow \nabla^2 \phi + \omega^2 \mu \epsilon \phi = -\rho/\epsilon - j\omega (\nabla \cdot \vec{A} + j\omega \mu \epsilon \phi)$$

引入洛伦兹规范

$$\nabla^2 \phi + \omega^2 \mu \epsilon \phi = -\rho/\epsilon$$

2.

$$\text{通解: } \nabla^2 g(r) + \beta^2 g(r) = 0 \quad \text{-- ①} \quad (r \neq 0)$$

记 $g(r) = u(r)/r$, 在球坐标下有

$$\frac{1}{r^2} \frac{d}{dr} \left(r^2 \frac{du(r)}{dr} \right) + \beta^2 \frac{u(r)}{r} = 0$$

$$\Rightarrow \frac{d^2 u(r)}{dr^2} + \beta^2 u(r) = 0$$

$$\text{特征方程为 } x^2 + \beta^2 = 0, \quad x = e^{\pm j\beta r}$$

\therefore 独立解为

$$g_1(r) = C_1 e^{-j\beta r}/r, \quad g_2(r) = C_2 e^{j\beta r}/r$$

$r=0$ 处,

$$\nabla^2 g + \beta^2 g = -\delta(0)$$

$$\iint (\nabla^2 g + \beta^2 g) dv = -1$$

$$\nabla^2 g = \nabla \cdot \nabla g$$

$$\iint \nabla^2 g dv = \iint \nabla \cdot (\nabla g) dv = \iint \nabla g \cdot d\vec{S}$$

代入 $g_1(r)$ 有

$$\begin{aligned} \iint \nabla g \cdot d\vec{S} &= C_1 \iint \frac{-j\beta r e^{-j\beta r} - e^{-j\beta r}}{r^2} \cdot d\vec{S}, \quad r \rightarrow 0 \\ &= -C_1 \cdot 4\pi = -1 \quad \therefore C_1 = \frac{1}{4\pi} \end{aligned}$$

$$\therefore g_1(r) = \frac{e^{-j\beta r}}{4\pi r}$$

$$\text{同理 } g_2(r) = \frac{e^{j\beta r}}{4\pi r}$$

为方程两个独立解.

3.

$$\vec{D} = -\nabla \times \vec{F} = \epsilon \vec{E} \quad \therefore \vec{E} = -\frac{1}{\epsilon} \nabla \times \vec{F}$$

$$\nabla \times \vec{H} = j\omega \vec{D} = -j\omega \nabla \times \vec{F}$$

$$\therefore \nabla \times (\vec{H} + j\omega \vec{F}) = 0 \Rightarrow \vec{B} + j\omega \vec{F} = \nabla \phi'$$

$$\begin{cases} \vec{D} = -\nabla \times \vec{F} \\ \vec{H} + j\omega \vec{F} = \nabla \phi' \end{cases} \quad \begin{cases} \vec{E} = -\frac{1}{\epsilon} \nabla \times \vec{F} \\ \vec{B} + j\omega \mu \vec{F} = \nabla \phi' \end{cases}$$

场方程

$$\begin{aligned} \nabla \times \vec{E} &= -j\omega \vec{B} \\ \nabla \times \vec{H} &= j\omega \vec{D} \end{aligned} \Rightarrow [\nabla^2 + \mu \epsilon \omega^2] \begin{bmatrix} \vec{E} \\ \vec{H} \end{bmatrix} = 0$$

$$\begin{cases} \nabla^2 \vec{E} + \mu \epsilon \omega^2 \vec{E} = 0 \\ \nabla^2 \vec{H} + \mu \epsilon \omega^2 \vec{H} = 0 \end{cases}$$

仿照电流辐射场推导有

$$g(\vec{r}, \vec{r}') = e^{-jk|\vec{r}-\vec{r}'|}/|\vec{r}-\vec{r}'|, \quad k^2 = \mu\epsilon\omega^2$$

$$\phi'(\vec{r}) = \int \frac{\rho_m}{\mu} g(\vec{r}, \vec{r}') dV' \quad \vec{F}(\vec{r}) = \int -\epsilon \vec{M}(\vec{r}') g(\vec{r}, \vec{r}') dV'$$

$$\vec{D} = -\nabla \times \vec{F}$$

$$\vec{H} = \nabla \phi'(\vec{r}) - j\omega \vec{F}$$