

# Fluid Mechanics 1

## Lecture #3:

### Continuity Equation



Integral formulation of mass conservation *see video*

$$\frac{d}{dt} \int_{V(t)} \rho dV = \int_{V(t)} \frac{\partial \rho}{\partial t} dV + \int_{S(t)} \rho u_j n_j dS = 0$$

Global equation.

We also need a detailed equation.

We will derive detailed equation from the global equation

*video mathematics. gradient and divergence and Gauss*

Gradient

$$\vec{\nabla} \phi \equiv \begin{pmatrix} \frac{\partial \phi}{\partial x} \\ \frac{\partial \phi}{\partial y} \\ \frac{\partial \phi}{\partial z} \end{pmatrix} \equiv \left( \frac{\partial \phi}{\partial x} \quad \frac{\partial \phi}{\partial y} \quad \frac{\partial \phi}{\partial z} \right)^T$$

*Theorem!*



Note: input is scalar ( $\phi$ )  
output is vector ( $\vec{\nabla} \phi$ )

## 2 Divergence.

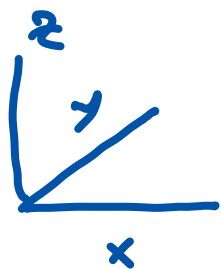
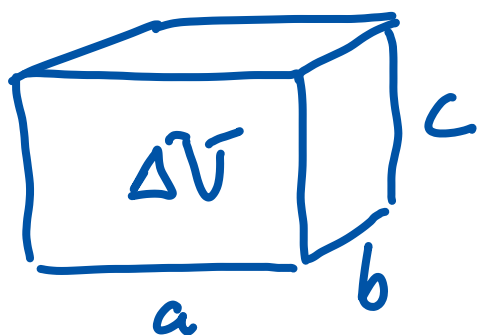
$$\vec{u}(\vec{x}, t)$$

$$\begin{aligned}\vec{\nabla} \cdot \vec{u} &\equiv \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \\ &\equiv \frac{\partial u_1}{\partial x_1} + \frac{\partial u_2}{\partial x_2} + \frac{\partial u_3}{\partial x_3} \\ &\equiv \sum_{j=1}^3 \frac{\partial u_j}{\partial x_j} \equiv \underbrace{\frac{\partial u_j}{\partial x_j}}_{ESC}\end{aligned}$$

Note: input is vector  $\vec{u}$   
output is scalar  $\vec{\nabla} \cdot \vec{u}$

What is the physical meaning of the divergence?

Consider a sufficiently small blob:



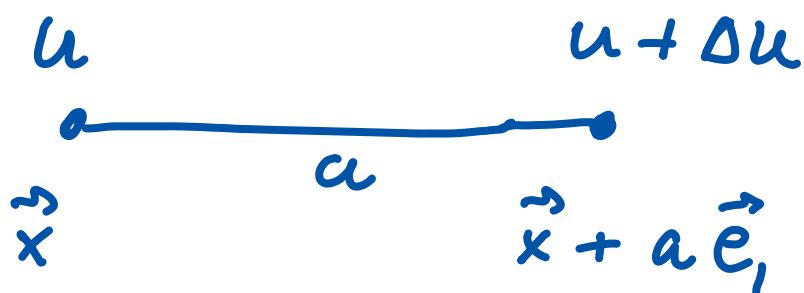
$$\frac{d}{dt} \Delta \bar{V} = ?$$

$$\frac{d}{dt} \Delta \bar{V} = \frac{d}{dt} (abc)$$

chain rule  $\uparrow = \frac{da}{dt} bc + a \frac{db}{dt} c + ab \frac{dc}{dt}$

Side "a":

$$\vec{e}_1 = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$$



$$\Rightarrow \frac{da}{dt} = (u + \Delta u) - u = \Delta u \quad \Delta u?$$

$$\Delta u \equiv u(x + a\vec{e}_1, y, z, t) - u(x, y, z, t)$$

$$= \frac{\partial u}{\partial x}(x, y, z, t) a + \mathcal{O}(a^2)$$

↑ Taylor series approximation.

$$\Rightarrow \frac{da}{dt} = \frac{\partial u}{\partial x} a + \dots$$

similar expressions for  $\frac{db}{dt}$ ,  $\frac{dc}{dt}$

$$\begin{aligned} \Rightarrow \frac{d}{dt}(\Delta \vec{V}) &= \left( \frac{\partial u}{\partial x} a \right) bc + a \left( \frac{\partial v}{\partial y} b \right) c + ab \left( \frac{\partial w}{\partial z} c \right) + \dots \\ &= \left( \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \right) abc + \dots \end{aligned}$$

$$\Rightarrow \frac{d}{dt} \Delta \vec{V} = (\vec{\nabla} \cdot \vec{u}) \Delta \vec{V} + \dots$$

$\Delta \vec{V} \rightarrow 0$  :

$$\boxed{\frac{1}{\Delta \vec{V}} \frac{d}{dt}(\Delta \vec{V}) = \vec{\nabla} \cdot \vec{u}}$$

$\Rightarrow \vec{\nabla} \cdot \vec{u}$  is the relative change with time of a sufficiently small material blob.

Question:  $\frac{1}{\cancel{\Delta V}} \frac{d}{dt} \cancel{\Delta V} \stackrel{?}{=} \frac{d}{dt}$

absolutely NOT!

note  $\frac{d}{dt} \Delta \bar{V} \equiv \lim_{\Delta t \rightarrow 0} \frac{\Delta \bar{V}(t + \Delta t) - \Delta \bar{V}(t)}{\Delta t}$

Similarly:  $\frac{\cancel{\sin} x}{\cancel{x}} = \sin x = 6$  (joke)

But:  $\frac{1}{\Delta \bar{V}} \frac{d}{dt} \Delta \bar{V} \equiv \frac{d}{dt} (\ln \Delta \bar{V})$

check with chain rule:

$$\frac{d}{dt} \ln \Delta \bar{V} = \frac{d}{d \Delta \bar{V}} \ln \Delta \bar{V} \cdot \frac{d \Delta \bar{V}}{dt} = \frac{1}{\Delta \bar{V}} \frac{d \Delta \bar{V}}{dt}.$$

What does all of this mean?

Suppose  $\rho = \text{const} \Rightarrow$  incompressible.

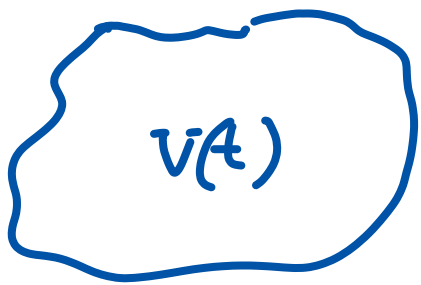
$$\Rightarrow \frac{d}{dt} \Delta \tilde{V} = 0 \quad \Rightarrow \boxed{\vec{\nabla} \cdot \vec{u} = 0}$$

Why do we use partial derivatives?

~~$\frac{du}{dx} + \frac{dv}{dy} + \frac{dw}{dz}$~~       wrong!

$$\frac{\partial u}{\partial x} \equiv \lim_{\Delta x \rightarrow 0} \frac{u(x + \Delta x, y, z, t) - u(x, y, z, t)}{\Delta x}$$

Sofar: sufficiently small blob  $\Delta V$   
 What about an arbitrary blob  $V$ ?



$$\frac{dV}{dt} = ?$$

$$\frac{dV}{dt} = \frac{d}{dt}(V) = \frac{d}{dt} \left( \sum_j \Delta V_j \right)$$

$$= \sum_j \frac{d}{dt}(\Delta V_j) = \sum_j (\vec{\nabla} \cdot \vec{u})_j \Delta V_j$$

in the limit of  $\Delta V_j \rightarrow 0$ :

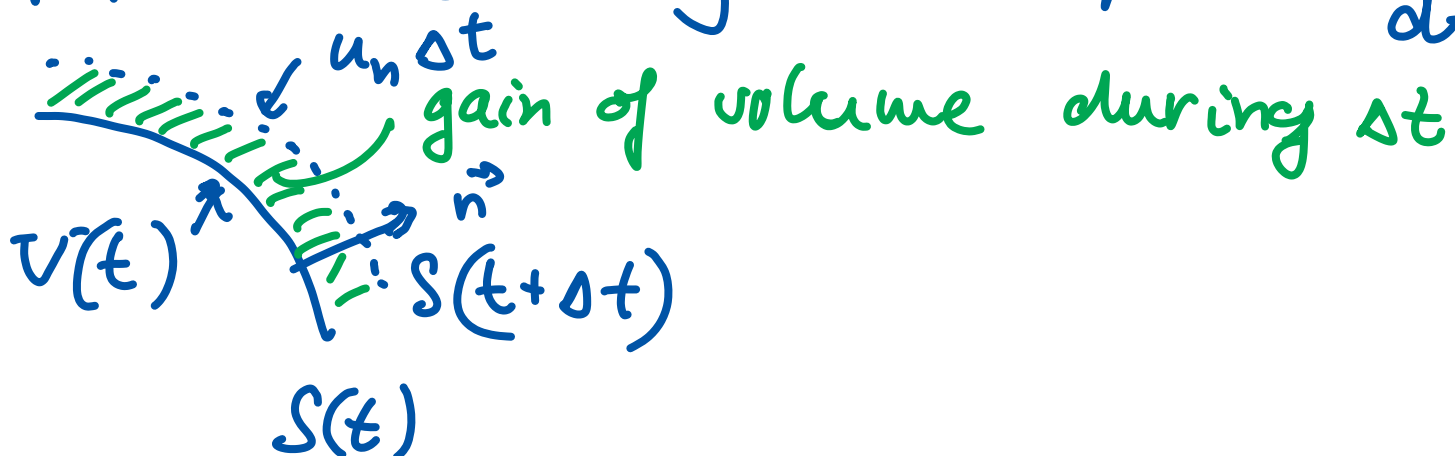
$$\boxed{\frac{dV}{dt} = \int_{V(t)} \vec{\nabla} \cdot \vec{u} \, dV}$$

dimensional

$$\frac{m^3}{s} = \frac{1}{m} \frac{m}{s} \cdot m^3$$

ℓ.

Alternative way to compute  $\frac{dV}{dt}$ :



$$\Rightarrow \frac{d\bar{V}}{dt} = \lim_{\Delta t \rightarrow 0} \frac{\bar{V}(t + \Delta t) - \bar{V}(t)}{\Delta t}$$

$$= \lim_{\Delta t \rightarrow 0} \frac{1}{\cancel{\Delta t}} \left\{ \int_{S(t)} \vec{u} \cdot \vec{n} \cancel{\Delta t} dS \right\}$$

$$\Rightarrow \boxed{\frac{d\bar{V}}{dt} = \int_{S(t)} \vec{u} \cdot \vec{n} dS}$$

Equate the two expressions for  $\frac{d\bar{V}}{dt}$ :

$$\boxed{\int_{\bar{V}(t)} \vec{\nabla} \cdot \vec{u} d\bar{V} = \int_{S(t)} \vec{u} \cdot \vec{n} dS'}$$

## Divergence Theorem of Gauss

Little guy in class room:

$$1 + 2 + 3 + \dots + 99 + 100 =$$

$$= (1 + 100) + (2 + 99) + \dots + (50 + 51)$$

$$= 50 \times 101 = \underline{\underline{5050}}$$

Little guy was Gauss.

Index notation: + ESC

$$\int_{V(t)} \frac{\partial u_j}{\partial x_j} dV = \int_{S(t)} u_j n_j dS$$

for arbitrary  
vector fields

Back to mass conservation:

integral form:  $\int_{V(t)} \frac{\partial \rho}{\partial t} dV + \int_{S(t)} \rho u_j n_j dS = 0$

$$\int_{S(t)} (\rho u_j) n_j dS = \int_{V(t)} \frac{\partial}{\partial x_j} (\rho u_j) dV$$

↑ Gauss

(note: vector field is  $\rho \vec{u} = \begin{pmatrix} \rho u \\ \rho v \\ \rho w \end{pmatrix}$ ).

$$\Rightarrow \int_{V(t)} \frac{\partial \rho}{\partial t} dV + \int_{V(t)} \frac{\partial}{\partial x_j} (\rho u_j) dV = 0$$

$$\Rightarrow \int_{V(t)} \left\{ \frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x_j} (\rho u_j) \right\} dV = 0 \quad \forall V, t$$

$\equiv$

$$\Leftrightarrow \left\{ \right\} = 0 \quad \forall \vec{x}, t$$

↳ equivalence! ▼



$$\Rightarrow \boxed{\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x_j}(\rho u_j) = 0}$$

Differential Formulation of Mass Conservation.

Note, via ESC:

$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x_1}(\rho u_1) + \frac{\partial}{\partial x_2}(\rho u_2) + \frac{\partial}{\partial x_3}(\rho u_3) = 0$$

$$\Rightarrow \frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x}(\rho u) + \frac{\partial}{\partial y}(\rho v) + \frac{\partial}{\partial z}(\rho w) = 0$$

Vector notation:  $\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{u}) = 0$

Special case:  $\rho = \text{const}$  (incompressible fluid)

$$\frac{\partial \rho}{\partial t} = 0 \Rightarrow \frac{\partial}{\partial x_j}(\rho u_j) = 0 \quad \frac{\partial \rho}{\partial x_1} = 0 \text{ etc}$$

$$\rho \frac{\partial u_j}{\partial x_j} = 0 \quad \rho \neq 0 \Rightarrow \boxed{\frac{\partial u_j}{\partial x_j} = 0}$$

Other special case:

Steady flow (but not  $\rho = \text{const}$ ).

$$\frac{\partial \rho}{\partial t} = 0 \Rightarrow \boxed{\frac{\partial}{\partial x_j}(\rho u_j) = 0}$$