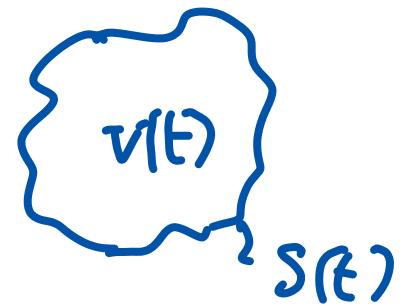


Fluid Mechanics 1

Lecture #3 :

Continuity Equation



Integral formulation of mass conservation *see video*

$$\frac{d}{dt} \int_{V(t)} \rho dV = \int_{\bar{V}(t)} \frac{\partial \rho}{\partial t} d\bar{V} + \int_{\mathcal{S}(t)} \rho u_j \cdot n_j d\mathcal{S} = 0$$

Global equation.

We also need a detailed equation. *flow field*

We will derive detailed equation from the global equation

video mathematics: gradient and divergence, and Gauss

1 Gradient $\phi(\vec{x}, t)$ e.g. $\rho(\vec{x}, t)$

$\vec{\nabla} \phi \equiv \begin{pmatrix} \frac{\partial \phi}{\partial x} \\ \frac{\partial \phi}{\partial y} \\ \frac{\partial \phi}{\partial z} \end{pmatrix} = \left(\frac{\partial \phi}{\partial x} \frac{\partial \phi}{\partial y} \frac{\partial \phi}{\partial z} \right)^T$

Theorem.

Note: input is scalar (ϕ)
output is vector ($\vec{\nabla} \phi$)

2 Divergence.

$$\vec{u}(\vec{x}, t)$$

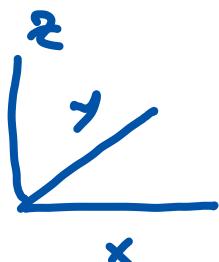
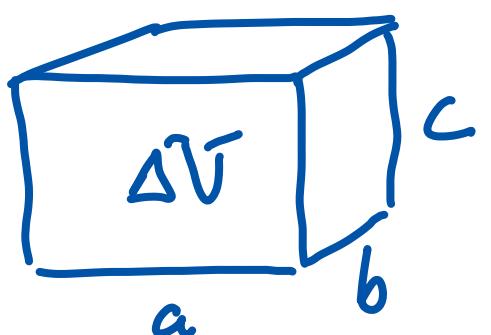
$$\begin{aligned}
 \vec{\nabla} \cdot \vec{u} &\equiv \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \\
 &\equiv \frac{\partial u_1}{\partial x_1} + \frac{\partial u_2}{\partial x_2} + \frac{\partial u_3}{\partial x_3} \\
 &\equiv \sum_{j=1}^3 \frac{\partial u_j}{\partial x_j} \equiv \frac{\partial u_j}{\partial x_j}
 \end{aligned}$$

ESC

Note: input is vector \vec{u}
output is scalar $\vec{\nabla} \cdot \vec{u}$

What is the physical meaning of the divergence?

Consider a sufficiently small blob:



$$\frac{d}{dt} \Delta V = ?$$

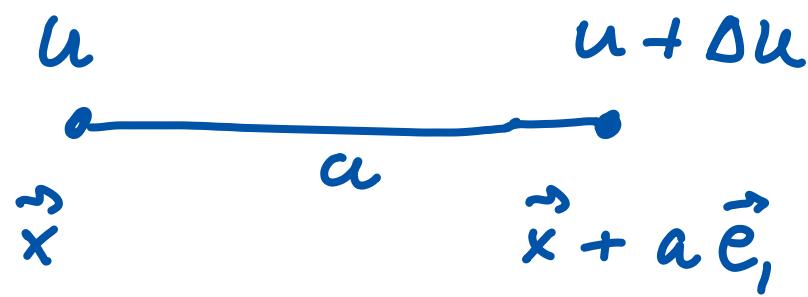
$$\frac{d}{dt} \Delta V = \frac{d}{dt} (abc)$$

chain rule

$$\begin{aligned}
 \frac{d}{dt} \Delta V &= \frac{da}{dt} bc + a \frac{db}{dt} c + ab \frac{dc}{dt}
 \end{aligned}$$

Side "a":

$$\vec{e}_1 = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$$



$$\Rightarrow \frac{da}{dt} = (u + \Delta u) - u = \Delta u \quad \Delta u?$$

$$\Delta u \equiv u(x + a \vec{e}_1, y, z, t) - u(x, y, z, t)$$

$$= \frac{\partial u}{\partial x}(x, y, z, t) a + \mathcal{O}(a^2)$$

↑ Taylor series approximation.

$$\Rightarrow \frac{da}{dt} = \frac{\partial u}{\partial x} a + \dots$$

Similar expressions for $\frac{db}{dt}, \frac{dc}{dt}$

$$\Rightarrow \frac{d}{dt}(\Delta \vec{V}) = \left(\frac{\partial u}{\partial x} a \right) bc + a \left(\frac{\partial v}{\partial y} b \right) c + ab \left(\frac{\partial w}{\partial z} c \right) + \dots$$

$$= \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \right) abc + \dots$$

$$\Rightarrow \frac{d}{dt} \Delta \vec{V} = (\vec{\nabla} \cdot \vec{u}) \Delta \vec{V} + \dots$$

$\Delta \vec{V} \rightarrow 0 :$

$$\frac{1}{\Delta \vec{V}} \frac{d}{dt}(\Delta \vec{V}) = \vec{\nabla} \cdot \vec{u}$$

$\Rightarrow \vec{\nabla} \cdot \vec{u}$ is the relative change with time of a sufficiently small material blob.

Question: $\frac{1}{\Delta V} \frac{d}{dt} \Delta \bar{V} \stackrel{?}{=} \frac{d}{dt}$

absolutely NOT!

note $\frac{d}{dt} \Delta \bar{V} \equiv \lim_{\Delta t \rightarrow 0} \frac{\Delta \bar{V}(t + \Delta t) - \Delta \bar{V}(t)}{\Delta t}$

Similarly: $\frac{\sin x}{x} = \sin x = 6$ (joke)

But: $\frac{1}{\Delta V} \frac{d}{dt} \Delta \bar{V} \equiv \frac{d}{dt} (\ln \Delta \bar{V})$

check with chain rule:

$$\frac{d}{dt} \ln \Delta \bar{V} = \frac{d}{d \Delta \bar{V}} \ln \Delta \bar{V} \cdot \frac{d \Delta \bar{V}}{dt} = \frac{1}{\Delta \bar{V}} \frac{d \Delta \bar{V}}{dt}.$$

What does all of this mean?

Suppose $\rho = \text{const}$. \Rightarrow incompressible.

$$\Rightarrow \frac{d}{dt} \Delta V = 0 \Rightarrow \boxed{\vec{F} \cdot \vec{u} = 0}$$

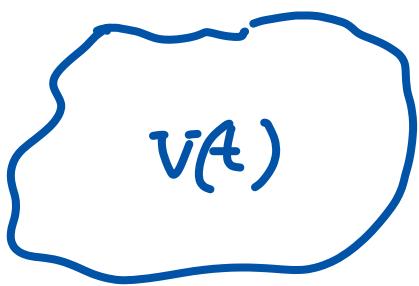
Why do we use partial derivatives?

$$\frac{du}{dx} + \frac{dv}{dy} + \frac{dw}{dz} \quad \text{wrong!}$$

$$\frac{\partial u}{\partial x} = \lim_{\Delta x \rightarrow 0} \frac{u(x+\Delta x, y, z, t) - u(x, y, z, t)}{\Delta x}$$

Sofar: sufficiently small blob δV

What about an arbitrary blob V ?



$$\frac{dV}{dt} = ?$$

$$\frac{dV}{dt} = \frac{d}{dt}(V) = \frac{d}{dt}\left(\sum_j \delta V_j\right)$$

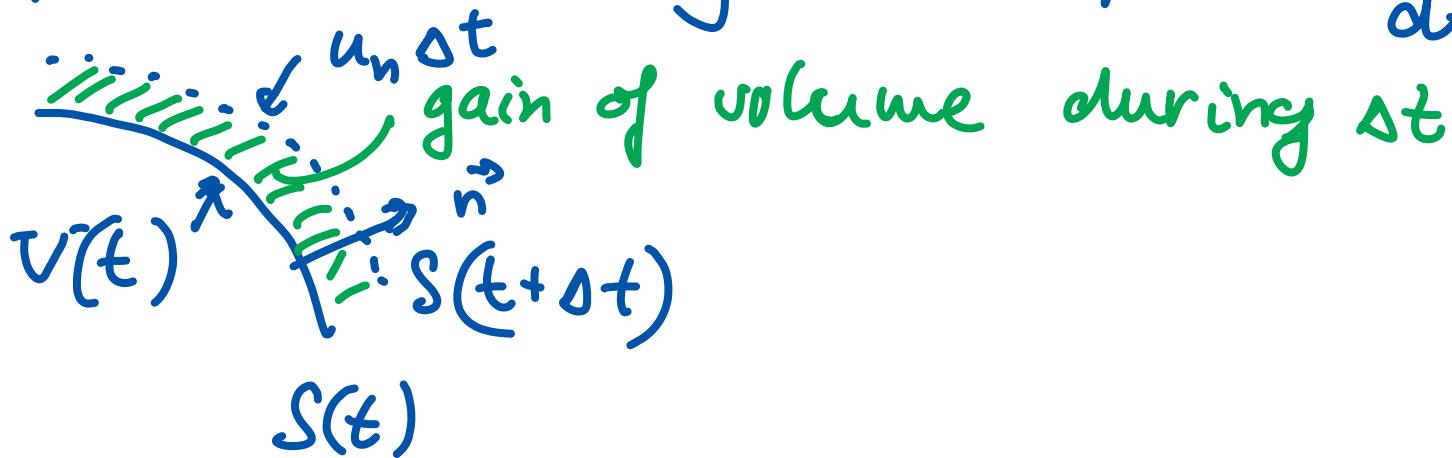
$$= \sum_j \frac{d}{dt}(\delta V_j) = \sum_j (\vec{\nabla} \cdot \vec{u})_j \delta V_j$$

in the limit of $\delta V_j \rightarrow 0$:

$$\boxed{\frac{dV}{dt} = \int_V \vec{\nabla} \cdot \vec{u} dV}$$

dimensions
 $\frac{m^3}{s} = \frac{1}{m} \frac{m}{s} \cdot m^3$

Alternative way to compute $\frac{dV}{dt}$:



$$\Rightarrow \frac{d\bar{V}}{dt} = \lim_{\Delta t \rightarrow 0} \frac{\bar{V}(t + \Delta t) - \bar{V}(t)}{\Delta t}$$

$$= \lim_{\Delta t \rightarrow 0} \frac{1}{\Delta t} \left\{ \int_{S(t)} \vec{u} \cdot \vec{n} \Delta \vec{S} \right\}$$

$$\Rightarrow \boxed{\frac{d\bar{V}}{dt} = \int_{S(t)} \vec{u} \cdot \vec{n} dS'}$$

Equate the two expressions for $\frac{d\bar{V}}{dt}$:

$$\int_{\bar{V}(t)} \vec{\nabla} \cdot \vec{u} d\bar{V} = \int_{S(t)} \vec{u} \cdot \vec{n} dS'$$

Divergence Theorem of Gauss

Little guy in class room:

$$1 + 2 + 3 + \dots + 99 + 100 =$$

$$= (1 + 100) + (2 + 99) + \dots + (50 + 51)$$

$$= 50 \times 101 = \underline{\underline{5050}}$$

Little guy was Gauss.

Index notation: + ESC

$$\int_{V(t)} \frac{\partial u_j}{\partial x_j} dV = \int_{S(t)} u_j n_j dS$$

for arbitrary
vector fields

Back to mass conservation:

integral form: $\int_{V(t)} \frac{\partial \rho}{\partial t} d\bar{V} + \int_{S(t)} \rho u_j n_j dS = 0$

$$\int_{S(t)} (\rho u_j) n_j dS = \int_{\substack{V(t) \\ \text{Gauss}}} \frac{\partial}{\partial x_j} (\rho u_j) d\bar{V}$$

(note: vector field $\rho \vec{u} = \begin{pmatrix} \rho u \\ \rho v \\ \rho w \end{pmatrix}$).

$$\Rightarrow \int_{V(t)} \frac{\partial \rho}{\partial t} d\bar{V} + \int_{V(t)} \frac{\partial}{\partial x_j} (\rho u_j) d\bar{V} = 0$$

$$\Rightarrow \int_{V(t)} \left\{ \frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x_j} (\rho u_j) \right\} d\bar{V} = 0 \quad \begin{matrix} + \\ \equiv \end{matrix} \bar{V}, t$$

$$\Leftrightarrow \{ \} = 0 \quad \forall \vec{x}, t$$

L equivalence!

$$\Rightarrow \boxed{\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x_j} (\rho u_j) = 0}$$

Differential
Formulation
of Mass
Conservation.

Note, via ESC:

$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x_1} (\rho u_1) + \frac{\partial}{\partial x_2} (\rho u_2) + \frac{\partial}{\partial x_3} (\rho u_3) = 0$$

$$\Rightarrow \frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x} (\rho u) + \frac{\partial}{\partial y} (\rho v) + \frac{\partial}{\partial z} (\rho w) = 0$$

Vector notation: $\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{u}) = 0$

Special case: $\rho = \text{const}$ (incompressible fluid)

$$\frac{\partial \rho}{\partial t} = 0 \Rightarrow \boxed{\frac{\partial}{\partial x_j} (\rho u_j) = 0} \quad \frac{\partial \rho}{\partial x_1} \Rightarrow \text{etc}$$

$$\rho \frac{\partial u_j}{\partial x_j} = 0 \quad \rho \neq 0 \Rightarrow \boxed{\frac{\partial u_j}{\partial x_j} = 0}$$

Other special case:

Steady flow (but not $\rho = \text{const}$).

$$\frac{\partial \rho}{\partial t} = 0 \Rightarrow \boxed{\frac{\partial}{\partial x_j} (\rho u_j) = 0}$$