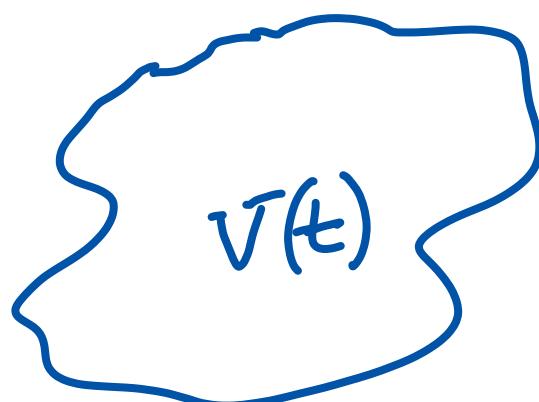
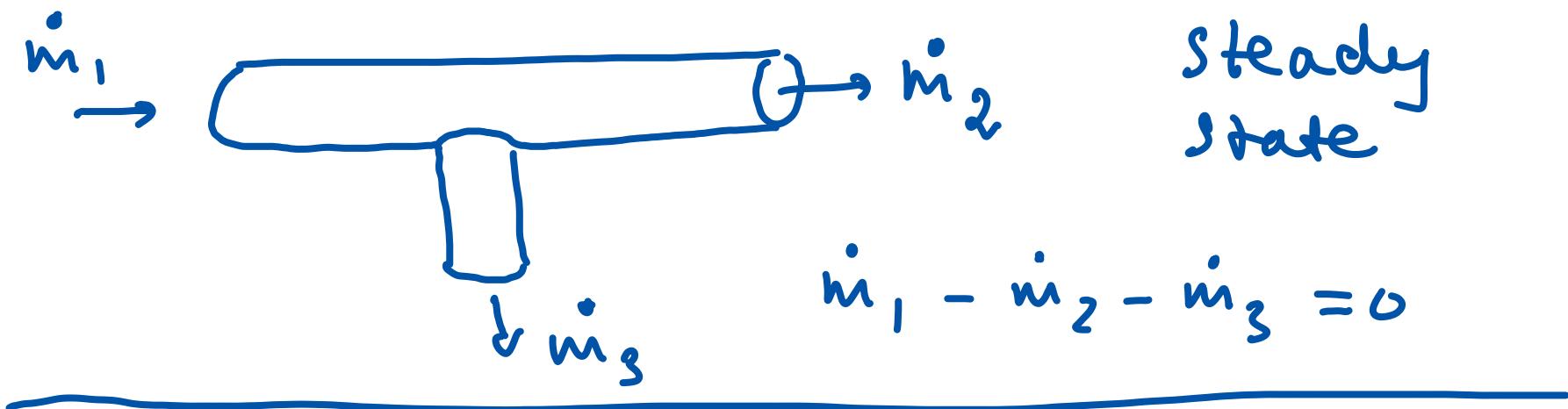


Fluid Mechanics 1

Lecture #2:

Mass Conservation.



blob of fluid

mass inside blob:

$$\text{Mass}(t) \equiv \int_{V(t)} \rho dV$$

$\rho(\vec{x}, t)$: mass density ($\frac{\text{kg}}{\text{m}^3}$)

Physics: $\text{Mass}(t) = \text{constant}$

$$\Rightarrow \frac{d}{dt} \text{Mass}(t) = 0 \Rightarrow \frac{d}{dt} \int_{V(t)} \rho dV = 0$$

$$\frac{d}{dt} \int_{V(t)} \rho dV = 0$$

Mathematical ToolBox

How to convert this into a useful equation?

Special case: if ∇ is not moving

$$\Rightarrow \frac{d}{dt} \int_{V(t)} \rho d\nu = \int_V \frac{\partial \rho}{\partial t} d\nu$$

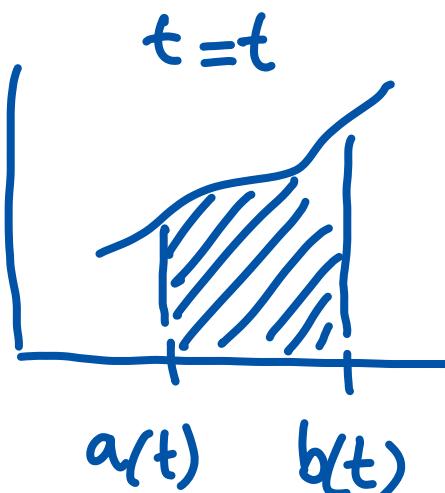
|
ordinary derivative

|
partial derivative.

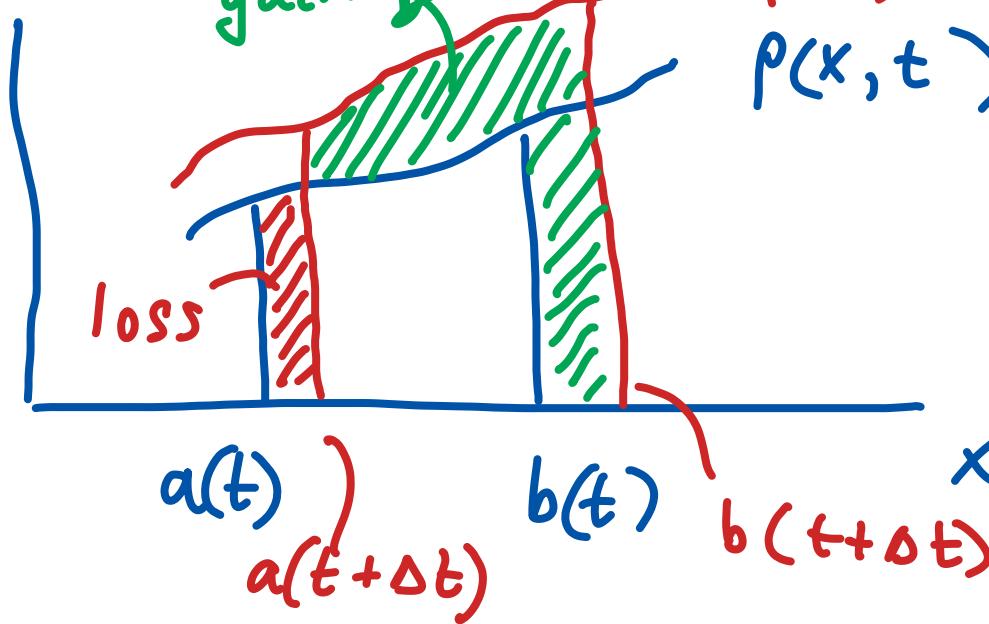
(*) Demonstration → minivideo Leibniz Reynolds
How to take the motion of ∇ into account?

1D example: $\rho(x,t)$

$$M(t) \equiv \int_{a(t)}^{b(t)} \rho(x,t) dx$$



$$\frac{dM}{dt} = ?$$



$$\begin{aligned} a(t + \Delta t) \\ = a(t) + \frac{da}{dt} \Delta t + \dots \end{aligned}$$

$$\frac{dM}{dt} = \lim_{\Delta t \rightarrow 0} \frac{M(t + \Delta t) - M(t)}{\Delta t}$$

width

$= \lim_{\Delta t \rightarrow 0} \frac{1}{\Delta t} \left\{ -\rho(a(t), t) \cdot \frac{da}{dt} \cdot \Delta t + \rho(b(t), t) \cdot \frac{db}{dt} \cdot \Delta t + \int_a(t) \frac{\partial \rho}{\partial t} dx + \underbrace{\text{error}}_{\Delta t^2} \right\}$

behaves like Δt^2

$$\Rightarrow \frac{dM}{dt} = -\rho(a(t), t) \frac{da}{dt} + \rho(b(t), t) \frac{db}{dt} + \int_a(t) \frac{\partial \rho}{\partial t} dx + 0$$

Leibniz' rule

Check dimensions.

$$M: \frac{kg}{m^3} m = \frac{kg}{m^2}$$

$$\Rightarrow \frac{dM}{dt}: \frac{kg}{m^2 s}$$

$$\rho \frac{da}{dt}: \frac{kg}{m^3} \frac{m}{s} \cancel{s}$$

$$\int_a^b \frac{\partial \rho}{\partial t} dx: \frac{kg}{m^3 s} \cancel{m} \cancel{s}$$

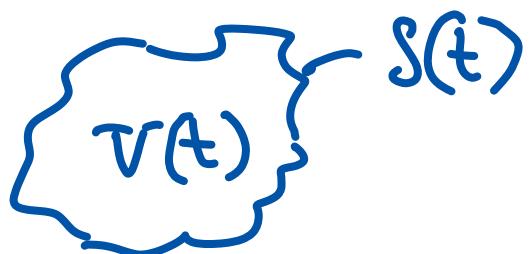
Remark: $a(t + \Delta t) = a(t) + \frac{da}{dt}(t) \Delta t + O(\Delta t^2)$

Taylor series.

order symbol
of Landau.

How does this translate to

$$\frac{d}{dt} \int_{V(t)} \rho dV ?$$



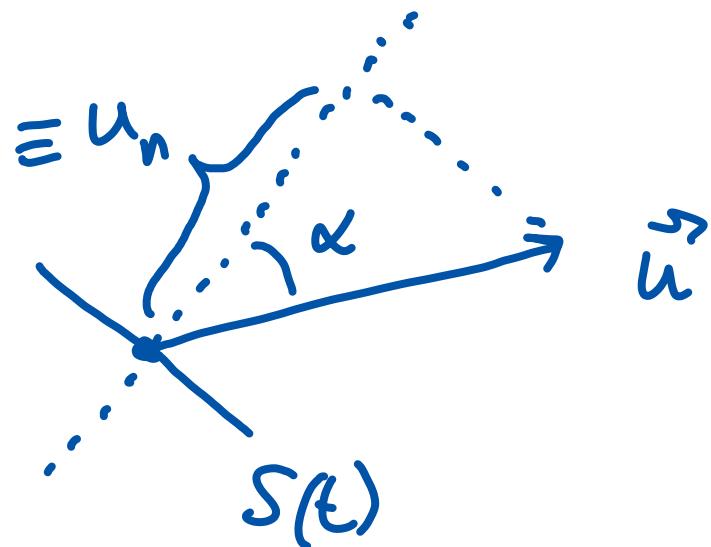
Answer:

$$\frac{d}{dt} \int_{V(t)} \rho(\vec{x}, t) dV = \int_{V(t)} \frac{\partial \rho}{\partial t} dV + \int_{S(t)} \rho u_n dS$$

Reynold's - Leibniz
transport theorem.

outward
 u_n = normal velocity
 $\vec{v}(t)$ $S(t)$ min, max

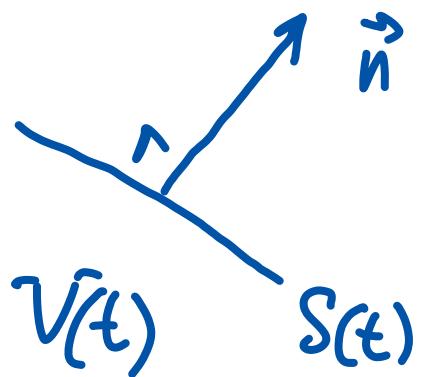
Mathematical Toolbox u_n
and inner product
How to compute u_n ?



u_n is scalar,
not a vector.

$$u_n = |\vec{u}| \cos \alpha$$

To compute $\cos \alpha$, we need the
outward unit normal vector.



\vec{n} :

- outward

$$|\vec{n}| = 1$$

- perpendicular to S (normal).

\vec{n} : non-dimensional

We also need the inner-product:

$$\vec{a}, \vec{b} \in \mathbb{R}^3 \quad \vec{a} \cdot \vec{b} \equiv a_1 b_1 + a_2 b_2 + a_3 b_3$$

$$\vec{a} = \begin{pmatrix} a_1 \\ a_2 \\ a_3 \end{pmatrix} \quad \vec{b} = \begin{pmatrix} b_1 \\ b_2 \\ b_3 \end{pmatrix} \quad \equiv \sum_{j=1}^3 a_j b_j$$

$$\equiv \sum_j a_j b_j$$

video

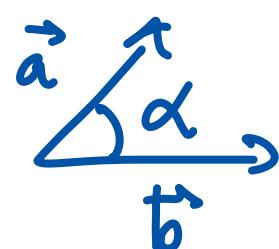
Separate, ESC

Einstein Summation convention
on repeated indices.

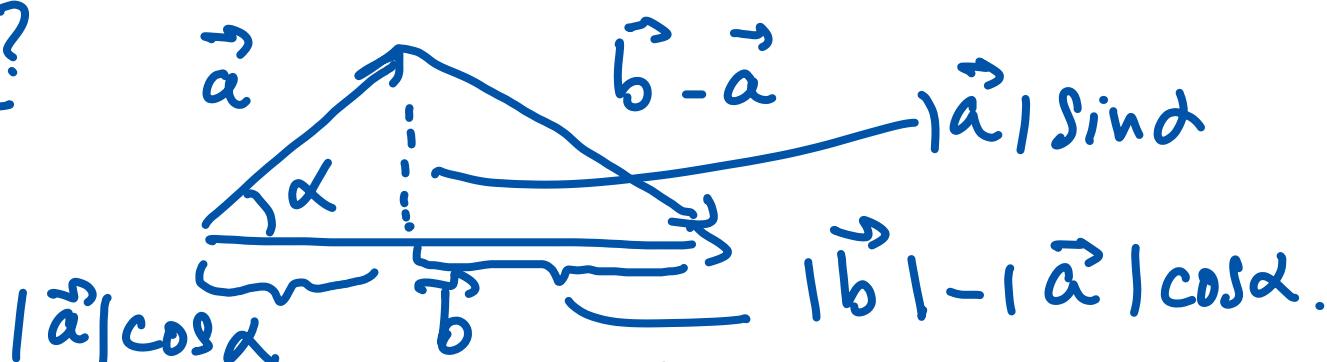
(ESC)

Property of the inner-product:

$$\vec{a} \cdot \vec{b} = |\vec{a}| |\vec{b}| \cos \alpha$$



Proof?



Idea: compute $|\vec{b} - \vec{a}|$ in two ways.

1 method:

$$\text{Pythagoras: } |\vec{b} - \vec{a}|^2 = |\vec{a}|^2 \sin^2 \alpha + (|\vec{b}| - |\vec{a}| \cos \alpha)^2$$

$$|\vec{a}| \equiv a \quad |\vec{b}| \equiv b$$

$$\Rightarrow |\vec{b} - \vec{a}|^2 = a^2 \sin^2 \alpha + b^2 - 2ab \cos \alpha + a^2 \cos^2 \alpha$$

$$= \underline{a^2 + b^2 - 2ab \cos \alpha}.$$

method 2: $|\vec{b} - \vec{a}|^2 = (\vec{b} - \vec{a}) \cdot (\vec{b} - \vec{a})$

$$= \vec{b} \cdot \vec{b} - \vec{a} \cdot \vec{b} - \vec{b} \cdot \vec{a} + \vec{a} \cdot \vec{a}$$

$$= |\vec{b}|^2 - 2\vec{a} \cdot \vec{b} + |\vec{a}|^2$$

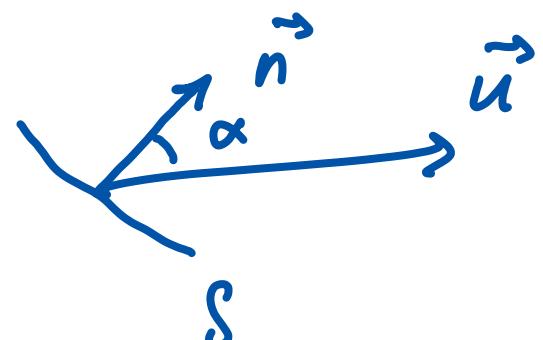
$$= \underline{a^2 + b^2 - 2\vec{a} \cdot \vec{b}}$$

$$\Rightarrow \vec{a} \cdot \vec{b} = ab \cos \alpha. \quad \square$$

Now return to computing $u_n = |\vec{u}| \cos \alpha$

Consider $\vec{u} \cdot \vec{n}$

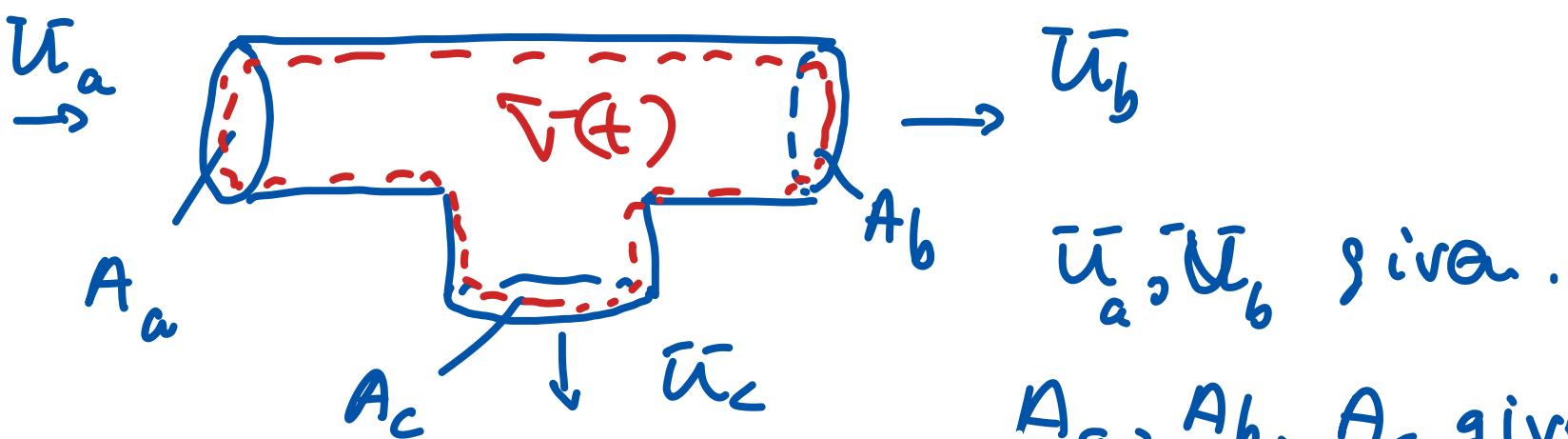
$$\vec{u} \cdot \vec{n} = |\vec{u}| |\vec{n}| \cos \alpha \quad |\vec{n}| \equiv 1$$



$$\Rightarrow \vec{u} \cdot \vec{n} = |\vec{u}| \cos \alpha. = u_n$$

$$\Rightarrow \boxed{u_n = \vec{u} \cdot \vec{n}} \quad \equiv u_j \cdot n_j. \quad \vec{n} = \begin{pmatrix} n_1 \\ n_2 \\ n_3 \end{pmatrix}$$

Application example .



Given: $\rho = \text{constant}$

Given: steady . $\frac{\partial^2}{\partial t^2}() = 0$

use: $\int \frac{\partial p}{\partial t} dV + \int p u_j n_j dS = 0$

Steady, also $p = \text{const.}$

$\uparrow S(t)$ $\quad p = \text{const.}$

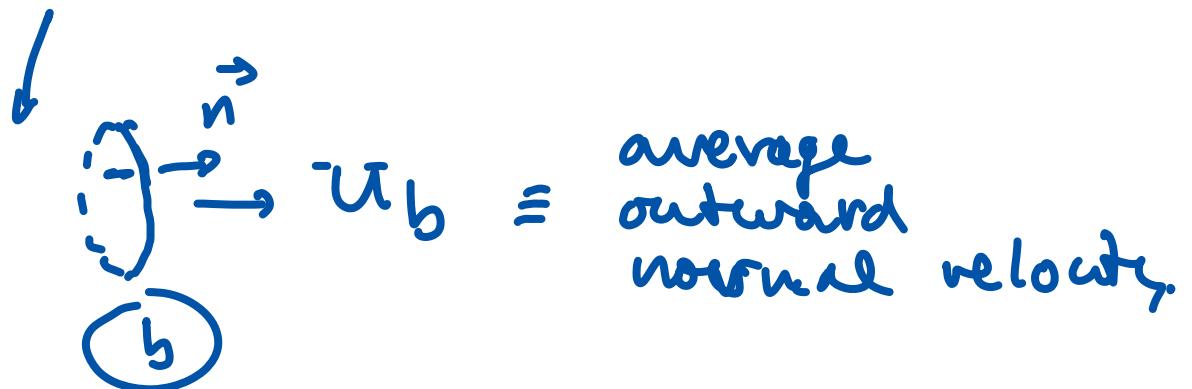
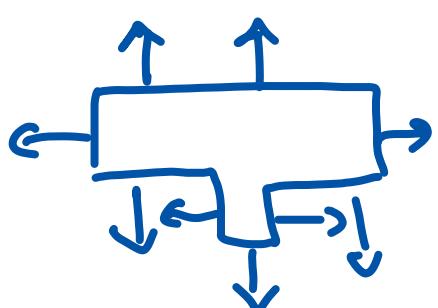
$$\Rightarrow \rho \int_{S(t)} u_j n_j dS = 0 \quad \rho \neq 0 \Rightarrow$$

||||| Solid Steady wall
 $\Rightarrow u_j \cdot n_j = 0$

$$S(t) = A_a + A_b + A_c + A_w$$

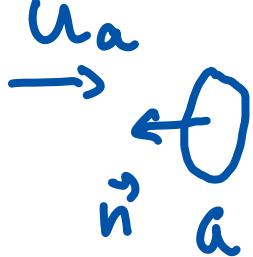
$$\Rightarrow \int_{A_a} \dots + \int_{A_b} \dots + \int_{A_c} \dots + \cancel{\int_{A_w} \dots} = 0 \quad (=0 \quad u_j n_j = 0)$$

$$\Rightarrow \int_{A_a} u_j n_j dS + \int_{A_b} u_j n_j d\ell + \int_{A_c} u_j n_j dJ = 0$$



$$\frac{1}{A_b} \int_{A_b} u_j n_j dS \equiv \bar{u}_b$$

$$\Rightarrow \text{2nd integral is } A_b \bar{u}_b$$



$\bar{u}_a \approx$ average inward normal velocity !

$$-\frac{1}{A_a} \int_{A_a} u_j n_j dS \equiv \bar{u}_a \Rightarrow \text{1st integral is } -A_a \bar{u}_a$$

3rd integral is $A_C \bar{u}_C$

$$\Rightarrow -A_a \bar{u}_a + A_b \bar{u}_b + A_c \bar{u}_c = 0$$

$$\Rightarrow \boxed{\bar{u}_c = \frac{A_a}{A_c} \bar{u}_a - \frac{A_b}{A_c} \bar{u}_b}$$

check dimensions: $\frac{A_a}{A_c} : 1 \quad \cancel{t}$

other checks: suppose $A_b=0$ \textcircled{b} is closed

$$\Rightarrow \bar{u}_c = \frac{A_a}{A_c} \bar{u}_a$$

$$\text{if } A_a = A_c \Rightarrow \bar{u}_c = \bar{u}_a \quad \cancel{t}.$$

\rightarrow Leibniz - Reynolds

- normal velocity and inner product
- Einstein summation convention