Sources: Alexei Kitaev's Characterization of Quantum Channels notes, John Preskill's chapter 3 notes, Robert Griffiths' Quantum Channels, Kraus Operators, POVMs notes.

GOAL: Prove that a superoperator T is completely positive iff \check{T} is Hermitian and positive (i.e. positive semi-definite).

1. Prove that T is completely positive if \check{T} is Hermitian and positive. Recall that the matrix \check{T} is related to the superoperator T in the following way:

$$\check{T}_{JJ'} = \check{T}_{ik:j'k'} = T_{ij':kk'},$$
(1)

where $T_{jj';kk'}$ constitutes the elements of the four-dimensional tensor T.

Because \check{T} is Hermitian, we can write a spectral decomposition for it, $\check{T} = MAM^{\dagger}$, where the columns (rows) of M (M^{\dagger}) correspond to the eigenvectors of \check{T} , and the real, non-negative values in the diagonal matrix A are its eigenvalues (non-negative since \check{T} is positive).

Let's rewrite \check{T} as $\Sigma_a \check{M}_a \check{M}_a^{\dagger}$, where \check{M}_a (\check{M}_a^{\dagger}) corresponds to a column (row) vector of M (M^{\dagger}) multiplied by the square root of its eigenvalue, which is also non-negative and real. $\Sigma_a \check{M}_a \check{M}_a^{\dagger}$ is also of course Hermitian and positive, which we can show in a different way:

$$Hermitian: (\Sigma_a \breve{M}_a \breve{M}_a^{\dagger})^{\dagger} = \Sigma_a \breve{M}_a \breve{M}_a^{\dagger}$$
 (2)

$$Positive: \Sigma_a \langle \psi | \check{M}_a \check{M}_a^{\dagger} | \psi \rangle = \Sigma_a ||\check{M}_a^{\dagger} | \psi \rangle ||^2 \ge 0 \to \Sigma_a \check{M}_a \check{M}_a^{\dagger} \ge 0 \quad (3)$$

In index notation,

$$\check{T}_{jk;j'k'} = \check{T}_{JJ'} = \Sigma_a(\check{M}_a)_J(\check{M}_a^{\bullet})_{J'} = \Sigma_a(\check{M}_a)_{jk}(\check{M}_a^{\bullet})_{j'k'}.$$
(4)

We can now reconstruct $T_{jj';kk'}$ by swapping the two middle indices of $\check{T}_{jk;j'k'}$:

$$T_{jj';kk'} = \Sigma_a(M_a)_{j;k}(M_a^*)_{j';k'}.$$
 (5)

Recall that for an arbitrary operator $X \in \mathbf{L}(\mathcal{L})$:

$$(TX)_{ii'} = \Sigma_{kk'} T_{ii':kk'} X_{kk'}. \tag{6}$$

Combining the two previous equations, we get:

$$(TX)_{jj'} = \sum_{akk'} (M_a)_{j:k} X_{kk'} (M_a^*)_{j':k'}. \tag{7}$$

Let's rewrite this without the indices:

$$TX = \Sigma_a M_a X M_a^{\dagger}. \tag{8}$$

This is very reminiscent of the equivalence between quantum channels in the Kraus representation and quantum channels as CPTP linear superoperators, except that we have not implemented the trace-preserving condition yet. Regardless, we can check whether or not T is a completely positive map by checking if TX and $(T \otimes I)X'$ preserve Hermiticity and positivity.

First, we let $X = X^{\dagger}$ and $X \geq 0$. We can expand X so that $X = \sum_{i} p_{i} |\phi_{i}\rangle \langle \phi_{i}|$, for $p_{i} \geq 0$. We verify that TX also holds these properties:

$$(TX)^{\dagger} = (\Sigma_a M_a X M_a^{\dagger})^{\dagger} = \Sigma_a M_a X M_a^{\dagger} = TX \tag{9}$$

$$\langle \psi | TX | \psi \rangle = \Sigma_{ai} p_i \langle \psi | M_a | \phi_i \rangle \langle \phi_i | M_a^{\dagger} | \psi \rangle$$
$$= \Sigma_a p_i || \langle \phi_i | M_a^{\dagger} | \psi \rangle ||^2 \ge 0 \to TX \ge 0 \quad (10)$$

Consequently, T is a positive map.

Next, we consider X', which is an operator that exists on an augmented Hilbert space. Let $X' = X'^{\dagger}$ and $X' \geq 0$. We can expand X' so that $X' = \sum_i p_i' |\phi_i'\rangle \langle \phi_i'|$, for $p_i' \geq 0$. We verify that $(T \otimes I)X'$ also holds these properties:

$$((T \otimes I)X')^{\dagger} = (\Sigma_a(M_a \otimes I)X'(M_a^{\dagger} \otimes I))^{\dagger}$$
$$= \Sigma_a(M_a \otimes I)X'(M_a^{\dagger} \otimes I) = (T \otimes I)X' \quad (11)$$

$$\langle \psi | (T \otimes I)X' | \psi \rangle = \sum_{ai} p_i' \langle \psi | (M_a \otimes I) | \phi_i' \rangle \langle \phi_i' | (M_a^{\dagger} \otimes I) | \psi \rangle$$
$$= \sum_a p_i' || \langle \phi_i' | (M_a^{\dagger} \otimes I) | \psi \rangle ||^2 \ge 0 \to (T \otimes I)X' \ge 0 \quad (12)$$

Consequently, T is a completely positive map. T is completely positive if \check{T} is Hermitian and positive.

2. Brief aside regarding the equivalence between different definitions of quantum channels.

Let's now implement the trace-preserving property on our superoperator T:

$$\delta_{kk'} = \Sigma_j T_{jj,kk'} = \Sigma_{aj} (M_a)_{j;k} (M_a^*)_{j;k'} = (\Sigma_a M_a^{\dagger} M_a)_{kk'}. \tag{13}$$

In other words, $\Sigma_a M_a^{\dagger} M_a = I$. This is the exact condition which is necessary to ensure that the channel $\mathcal{E}(X) = \Sigma_a M_a X M_a^{\dagger}$ is trace-preserving.

We have therefore shown that the quantum channel that is represented by a CPTP linear superoperator is equivalent to the quantum channel $\mathcal{E}(X)$ in the Kraus representation.

3. Prove that if T is completely positive, \check{T} is Hermitian and positive.

If T is completely positive, the result of the map $(T \otimes I)X'$ will be a positive semi-definite operator for any positive semi-definite operator X', where X' lies on the augmented Hilbert space $\mathcal{L} \otimes \mathcal{N}$.

Since we are allowed to choose X', let $X' = \sum_{jkj'k'} |jk\rangle \langle j'k'| \, \delta_{jk}\delta_{j'k'}$, or in index notation, $(X')_{jkj'k'} = \delta_{jk}\delta_{j'k'}$. X' is the unnormalized maximally entangled state between both of our Hilbert spaces. We have constrained our Hilbert spaces $\mathcal L$ and $\mathcal N$ to be of the same dimension, which is a choice that we are allowed to make, since it does not impose any restrictions on T itself.

Let us compute $(T \otimes I)X'$:

$$((T \otimes I)X')_{jkj'k'} = \sum_{ll'} T_{jj';ll'} X'_{lkl'k'}$$

$$= \sum_{ll'} T_{jj';ll'} \delta_{lk} \delta_{l'k'} = T_{jj';kk'} = \breve{T}_{jk;j'k'} = \breve{T}_{JJ'} = \breve{T} \quad (14)$$

We can immediately conclude that \check{T} is positive semi-definite (Hermitian and positive), since T is completely positive. If T is completely positive, \check{T} is Hermitian and positive.

To conclude, a superoperator T is completely positive iff \check{T} is Hermitian and positive.

Additional note: this proof can also be done using the channel state duality.