

# A Quaternion-Based Theory of Everything: Unifying Quantum Field Theory and Gravity with Dynamic Symmetry Breaking

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## Abstract

We present a novel Theory of Everything (TOE) that unifies the Standard Model (SM) of particle physics with gravity within a 4D spacetime using a quaternion-based framework. The measurable state is the quaternion wave function  $\psi = \psi_0 + \psi_1 I + \psi_2 J + \psi_3 K$ , a 4x4 real matrix, and the total quantum state  $Q = (M_1, M_2)$  is constructed from SU(4) matrices via the Cayley-Dickson process. The condition  $\psi_0^2 = |V|^2$ , where  $V = 2(\psi_1 I + \psi_2 J + \psi_3 K)$ , dynamically couples the scalar and vector components, reproducing SM masses (e.g., Higgs at 125 GeV) when normalized. SM gauge fields emerge from SU(4) projections, while gravity arises as  $T_{\mu\nu} = \text{Re}(Q^\dagger Q)$ , quantized as a spin-2 field. An operator  $q$  generates pilot waves via  $q^2$ , guiding fermion dynamics. The theory predicts a universe without dark matter particles, explores superluminal propagation, and offers testable cosmological and particle physics signatures.

## 1 Introduction

The quest to unify quantum mechanics and general relativity (GR) remains a central challenge in theoretical physics. The Standard Model (SM) of particle physics, based on the gauge group  $\text{SU}(3) \times \text{SU}(2) \times \text{U}(1)$ , successfully describes strong, weak, and electromagnetic interactions, while GR governs gravity through the Einstein field equations. However, these frameworks are incompatible at quantum scales, where quantum field theory (QFT) predicts singular behavior and GR resists quantization without modification. This discord has spurred diverse approaches, such as string theory, which introduces extra dimensions, and loop quantum gravity, which posits discrete spacetime. In contrast, this paper presents a novel theory of everything (TOE) that unifies SM forces and gravity within a conventional 4D spacetime using a quaternion-based wave function, eschewing additional particles or dimensions beyond those observed.

The theory begins with a quaternion  $\psi = \psi_0 + \psi_1 I + \psi_2 J + \psi_3 K$ , a 4x4 real matrix defining the measurable state, and a total quantum state  $Q = (M_1, M_2) = M_1 + M_2 k$ , where  $M_1$  and  $M_2$  are SU(4) matrices constructed via the Cayley-Dickson process

( $k^2 = -1$ ). The scalar component  $\psi_0$  is dynamically coupled to the fermion vector part  $V = 2(\psi_1 I + \psi_2 J + \psi_3 K)$  through the condition  $\psi_0^2 = |V|^2 = 4(\psi_1^2 + \psi_2^2 + \psi_3^2)$ . Normalizing  $\psi_1^2 + \psi_2^2 + \psi_3^2 = \frac{1}{4}$  ensures  $\psi_0 = 1$ , reproducing SM masses such as the Higgs boson at 125 GeV. The SU(4) matrices project onto the SM gauge groups—SU(3) for the strong force, SU(2) for the weak force, and U(1) for electromagnetism—fully quantized within a QFT framework. Gravity emerges as a collective effect of SM fields through the stress-energy tensor  $T_{\mu\nu} = \text{Re}(Q^\dagger Q)$ , which is quantized as a massless spin-2 field  $h_{\mu\nu}$ , identified as the graviton. An operator  $q$ , also a 4x4 real matrix, acts on  $Q$  via left multiplication ( $Qq$ ), generating quaternion-valued state vectors, with the vector components of  $q^2$  serving as pilot waves to guide fermion dynamics in a de Broglie-Bohm-like manner.

This paper formalizes the theory across multiple dimensions: Section 2 defines the mathematical formalism, detailing the quaternion structure, SU(4) projections, and the operator  $q$ 's role in state transformations and data storage. Section 3 explores the Higgs and fermion fields, elucidating mass generation via symmetry breaking and fermion quantization. Section 4 presents the quantization process yielding gravitons, unifying SM interactions with gravity. Section 5 addresses renormalization to ensure finite observables, tackling the challenges of quantizing gravity. Section 6 examines cosmological evolution and structure formation, predicting a universe without dark matter particles. Section 7 investigates the potential for superluminal wave packet propagation, constrained by relativistic causality. Finally, Sections 8 and 9 provide intuitive actions and experimental requirements, respectively, framing the theory as a testable unification of quantum mechanics and gravity within 4D spacetime.

## 2 Theory Definition and Formalism

### 2.1 Quaternion Wave Function

The foundational entity of this theory is the quaternion wave function  $\psi$ , defined as:

$$\psi = \psi_0 + \psi_1 I + \psi_2 J + \psi_3 K, \quad (1)$$

where  $\psi_0, \psi_1, \psi_2, \psi_3 \in \mathbb{R}$  are real-valued coefficients representing the scalar and vector components, respectively, and  $I, J, K$  are quaternion basis elements satisfying the algebraic relations:

$$I^2 = J^2 = K^2 = IJK = -1, \quad IJ = K, \quad JI = -K, \quad JK = I, \quad KJ = -I, \quad KI = J, \quad IK = -J. \quad (2)$$

These relations define the non-commutative quaternion algebra, which underpins the theory's unification of matter and spacetime in four dimensions. In matrix representation,  $\psi$  is expressed as a 4x4 real matrix acting on 4D vectors:

$$\psi = \begin{pmatrix} \psi_0 & -\psi_1 & -\psi_2 & -\psi_3 \\ \psi_1 & \psi_0 & -\psi_3 & \psi_2 \\ \psi_2 & \psi_3 & \psi_0 & -\psi_1 \\ \psi_3 & -\psi_2 & \psi_1 & \psi_0 \end{pmatrix}. \quad (3)$$

This matrix form arises from the quaternion basis mapped to 4x4 real matrices: the scalar 1 corresponds to the identity matrix, while  $I, J, K$  are represented by specific antisymmetric matrices ensuring the algebraic properties hold. The components  $\psi_j(x)$

( $j = 0, 1, 2, 3$ ) are spacetime-dependent fields, with  $\psi_j = R_j(x)e^{iS_j(x)/\hbar}$ , where  $R_j(x) \in \mathbb{R}_+$  is the amplitude and  $S_j(x) \in \mathbb{R}$  is the phase. This allows  $\psi$  to encode a relativistic four-momentum:

$$p^\mu = (\psi_0, p_1, p_2, p_3), \quad (4)$$

where  $p_j = \nabla S_j$  represents the spatial momentum components derived from the phase gradients, and  $\psi_0$  relates to the energy or scalar potential. Thus,  $\psi$  serves as the measurable state, unifying particle and field properties in a single quaternion structure, setting the stage for the theory's broader formalism.

## 2.2 Vector Part and Condition

The vector part of the quaternion wave function  $\psi$  is defined as the fermion component:

$$V = 2(\psi_1 I + \psi_2 J + \psi_3 K), \quad (5)$$

where  $V$  encapsulates the spatial degrees of freedom associated with fermion fields, distinct from the scalar component  $\psi_0$ . The factor of 2 is introduced to align the magnitude of  $V$  with physical scales, as will be evident in the normalization condition. The squared magnitude of  $V$  is computed using the quaternion conjugate  $V^\dagger = 2(\psi_1 I^\dagger + \psi_2 J^\dagger + \psi_3 K^\dagger)$ , noting that  $I^\dagger = -I$ ,  $J^\dagger = -J$ , and  $K^\dagger = -K$  in the matrix representation, yielding:

$$|V|^2 = V^\dagger V = 4(\psi_1^2 + \psi_2^2 + \psi_3^2). \quad (6)$$

This magnitude reflects the total contribution of the vector components  $\psi_1, \psi_2, \psi_3$ , squared and scaled by the factor of 4 due to the coefficient in  $V$ . A fundamental condition ties the scalar  $\psi_0$  to the vector part:

$$\psi_0^2 = |V|^2 = 4(\psi_1^2 + \psi_2^2 + \psi_3^2), \quad (7)$$

implying:

$$\psi_0 = 2\sqrt{\psi_1^2 + \psi_2^2 + \psi_3^2}. \quad (8)$$

This dynamic coupling ensures that the scalar component  $\psi_0$  is not an independent parameter but is determined by the magnitude of the fermion vector  $V$ , establishing a direct relationship between the scalar field (related to energy or mass scales) and the vectorial fermion content. To align with Standard Model (SM) phenomenology, particularly the Higgs vacuum expectation value (VEV) of 246 GeV, a normalization is imposed:

$$\psi_1^2 + \psi_2^2 + \psi_3^2 = \frac{1}{4}, \quad (9)$$

so:

$$|V|^2 = 4 \times \frac{1}{4} = 1, \quad \psi_0 = 2\sqrt{\frac{1}{4}} = 1. \quad (10)$$

This normalization sets  $\psi_0 = 1$ , which, when multiplied by the SM Higgs VEV scale  $v_0 = 246$  GeV, yields the physical VEV  $v' = \psi_0 v_0 = 246$  GeV, ensuring consistency with observed particle masses (e.g.,  $m_W \approx 80.4$  GeV). The condition  $\psi_0^2 = |V|^2$  thus serves as a cornerstone of the theory, linking the scalar and vector components in a unified quaternion framework, with  $\psi_0$  dynamically adjusting to spacetime variations in  $V$ .

## 2.3 Total State $Q$

While the quaternion wave function  $\psi$  represents the measurable state, encompassing the scalar  $\psi_0$  and fermion vector  $V$ , the total quantum state of the theory is encapsulated by  $Q$ , a quaternion-like field extending over 4D spacetime. It is defined as:

$$Q = (M_1, M_2) = M_1 + M_2 k, \quad (11)$$

where  $M_1, M_2 \in \text{SU}(4)$  are 4x4 complex unitary matrices with determinant 1, and  $k$  is a Cayley-Dickson extension satisfying  $k^2 = -1$ . Classically,  $Q(x)$  is a quaternion-valued function:

$$Q(x) = Q_0(x) + Q_1(x)I + Q_2(x)J + Q_3(x)K, \quad (12)$$

where  $Q_0(x), Q_1(x), Q_2(x), Q_3(x) \in \mathbb{R}$  are real scalar functions, represented as a 4x4 real matrix acting on 4D vectors:

$$Q(x) = \begin{pmatrix} Q_0(x) & -Q_1(x) & -Q_2(x) & -Q_3(x) \\ Q_1(x) & Q_0(x) & -Q_3(x) & Q_2(x) \\ Q_2(x) & Q_3(x) & Q_0(x) & -Q_1(x) \\ Q_3(x) & -Q_2(x) & Q_1(x) & Q_0(x) \end{pmatrix}. \quad (13)$$

In the quantum field theory (QFT) context,  $Q(x)$  becomes a 4x4 matrix-valued operator acting on a Hilbert space, unifying SM fields (via  $\text{SU}(4)$  projections) and gravitational contributions. Canonical quantization replaces the classical coefficients with operator-valued fields:

$$Q(x) = \int \frac{d^3 k}{(2\pi)^3 2\omega_k} \sum_{j=0}^3 \left[ a_j(k) Q_j(k) e^{-ik \cdot x} + a_j^\dagger(k) Q_j^\dagger(k) e^{ik \cdot x} \right], \quad (14)$$

where: -  $\omega_k = \sqrt{c^2 |\mathbf{k}|^2 + m^2}$  is the dispersion relation, with  $m$  a mass scale tied to the Higgs VEV ( $v' = \psi_0 v_0$ ,  $v_0 = 246 \text{ GeV}$  when  $\psi_0 = 1$ ), -  $a_j(k)$  and  $a_j^\dagger(k)$  are annihilation and creation operators for the  $j$ -th quaternion component, satisfying:

$$[a_j(k), a_{j'}^\dagger(k')] = (2\pi)^3 2\omega_k \delta_{jj'} \delta^3(\mathbf{k} - \mathbf{k}'), \quad [a_j(k), a_{j'}(k')] = 0,$$

-  $Q_j(k)$  are 4x4 basis matrices (e.g.,  $Q_0 = I_{4 \times 4}$ ,  $Q_1 = -I$  on off-diagonal terms, etc.), adjusted for momentum  $k$ .

The operator  $Q(x)$  retains its 4x4 structure, but each element becomes a field operator, enabling creation and annihilation of states (e.g., gluons, W bosons, fermions, gravitons). Classically,  $Q_j(x) \in \mathbb{R}$  ensures a deterministic configuration, while quantization introduces fluctuations, with the classical limit recovered via expectation values  $\langle Q(x) \rangle$ .  $Q$  thus represents the full quantum state space spanned by  $M_1$  and  $M_2$ , dynamically evolving to encode all physical degrees of freedom in the TOE.

## 2.4 $\text{SU}(4)$ Projections

The unification of Standard Model (SM) forces—strong ( $\text{SU}(3)$ ), weak ( $\text{SU}(2)$ ), and electromagnetic ( $\text{U}(1)$ )—is achieved through projections from the  $\text{SU}(4)$  gauge group inherent in the total quantum state  $Q = (M_1, M_2)$ , where  $M_1$  and  $M_2$  are 4x4 complex  $\text{SU}(4)$

matrices with determinant 1. The  $SU(4)$  group, with 15 independent generators corresponding to its Lie algebra  $\mathfrak{su}(4)$ , is represented by Hermitian, traceless 4x4 matrices  $T^a$  ( $a = 1, \dots, 15$ ) satisfying:

$$[T^a, T^b] = if^{abc}T^c, \quad (15)$$

where  $f^{abc}$  are structure constants, and the adjoint representation has dimension  $15 = 4^2 - 1$ . These generators span the SM gauge groups as subgroups, which are projected from  $M_1$  and  $M_2$  to recover the SM's gauge fields.

- **SU(3) Projection (Strong Force):** The  $SU(3)$  subgroup, governing the strong interaction, is embedded in the upper 3x3 block of  $M_1$ :

$$M_1 = \begin{pmatrix} M_{SU(3)} & * \\ * & * \end{pmatrix},$$

where  $M_{SU(3)}$  is a 3x3  $SU(3)$  matrix with 8 generators, the Gell-Mann matrices  $\lambda^a$  ( $a = 1, \dots, 8$ ):

$$\lambda^1 = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad \lambda^2 = \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad \dots, \quad \lambda^8 = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -2 \end{pmatrix}.$$

This yields 8 gluon fields  $G_\mu^a$  ( $a = 1, \dots, 8$ ):

$$G_{\mu\nu}^a = \partial_\mu G_\nu^a - \partial_\nu G_\mu^a + g_s f^{abc} G_\mu^b G_\nu^c,$$

with field strength Lagrangian  $\mathcal{L}_{SU(3)} = -\frac{1}{4}G_{\mu\nu}^a G^{a,\mu\nu}$ , and coupling  $g_s \approx 1$ . Quantization follows:

$$G_\mu^a(x) = \int \frac{d^3k}{(2\pi)^3 2\omega_k} \sum_\lambda \left[ c_{a,\lambda}(k) \epsilon_\mu^\lambda e^{-ik \cdot x} + c_{a,\lambda}^\dagger(k) \epsilon_\mu^{\lambda*} e^{ik \cdot x} \right],$$

where  $\omega_k = c|\mathbf{k}|$ , reflecting massless gluons.

- **SU(2) Projection (Weak Force):** The  $SU(2)$  subgroup, governing the weak interaction, is embedded within  $M_1$ , typically in a 2x2 block (e.g., rows/columns 1-2):

$$M_1 \supset \begin{pmatrix} M_{SU(2)} & 0 \\ 0 & I_{2 \times 2} \end{pmatrix},$$

with 3 generators, the Pauli matrices  $\tau^a$  ( $a = 1, 2, 3$ ):

$$\tau^1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \tau^2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \tau^3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

This produces 3 W boson fields  $W_\mu^a$ :

$$W_{\mu\nu}^a = \partial_\mu W_\nu^a - \partial_\nu W_\mu^a + g\epsilon^{abc} W_\mu^b W_\nu^c,$$

$\mathcal{L}_{SU(2)} = -\frac{1}{4}W_{\mu\nu}^a W^{a,\mu\nu}$ ,  $g \approx 0.65$ , quantized similarly to gluons pre-symmetry breaking.

- **U(1) Projection (Electromagnetic Force):** The U(1) subgroup, tied to hypercharge pre-breaking (and electromagnetism post-breaking), is a diagonal generator in SU(4):

$$M_1 \supset e^{i\theta Y}, \quad Y = \text{diag}(y_1, y_2, y_3, y_4),$$

traceless within SU(4), yielding the photon field  $B_\mu$ :

$$B_{\mu\nu} = \partial_\mu B_\nu - \partial_\nu B_\mu,$$

$\mathcal{L}_{\text{U}(1)} = -\frac{1}{4}B_{\mu\nu}B^{\mu\nu}$ ,  $g' \approx 0.36$ . Post-Higgs breaking,  $B_\mu$  mixes with  $W_\mu^3$  to form the photon  $A_\mu$ .

These projections from  $M_1$  (and similarly  $M_2$ ) recover the SM gauge fields, with  $M_1, M_2$  spanning dual Hilbert spaces  $(\mathcal{H}_1, \mathcal{H}_2)$  within  $Q$ , dynamically evolving to unify SM interactions and gravity via  $T_{\mu\nu}$ .

## 2.5 Operator $q$ and $q^2$

The operator  $q$  is a pivotal entity in the theory, acting on the total quantum state  $Q$  to generate transformed states and guide fermion dynamics via pilot waves. It is defined as a quaternion-valued operator:

$$q = q_0 + q_1 I + q_2 J + q_3 K, \quad (16)$$

where  $q_0, q_1, q_2, q_3 \in \mathbb{R}$  are spacetime-dependent coefficients, represented as a 4x4 real matrix:

$$q = \begin{pmatrix} q_0 & -q_1 & -q_2 & -q_3 \\ q_1 & q_0 & -q_3 & q_2 \\ q_2 & q_3 & q_0 & -q_1 \\ q_3 & -q_2 & q_1 & q_0 \end{pmatrix}. \quad (17)$$

In the quantum context,  $q(x)$  is promoted to a field operator:

$$q(x) = \int \frac{d^3 k}{(2\pi)^3 2\omega_k} \sum_{j=0}^3 \left[ b_j(k) q_j(k) e^{-ik \cdot x} + b_j^\dagger(k) q_j^\dagger(k) e^{ik \cdot x} \right], \quad (18)$$

where  $\omega_k = \sqrt{c^2 |\mathbf{k}|^2 + m^2}$ ,  $b_j(k)$  and  $b_j^\dagger(k)$  are annihilation and creation operators satisfying:

$$[b_j(k), b_{j'}^\dagger(k')] = (2\pi)^3 2\omega_k \delta_{jj'} \delta^3(\mathbf{k} - \mathbf{k}'), \quad [b_j(k), b_{j'}(k')] = 0, \quad (19)$$

and  $q_j(k)$  are basis matrices akin to those of  $Q$ . The operator  $q$  acts on  $Q$  via left multiplication:

$$\tilde{q} = Qq, \quad (20)$$

transforming a 4D vector  $q = (q_0, q_1, q_2, q_3)^T$  into  $\tilde{q} = (\tilde{q}_0, \tilde{q}_1, \tilde{q}_2, \tilde{q}_3)^T$ :

$$\tilde{q}_0 = Q_0 q_0 - Q_1 q_1 - Q_2 q_2 - Q_3 q_3, \quad (21)$$

$$\tilde{q}_1 = Q_1 q_0 + Q_0 q_1 - Q_3 q_2 + Q_2 q_3, \quad (22)$$

$$\tilde{q}_2 = Q_2 q_0 + Q_3 q_1 + Q_0 q_2 - Q_1 q_3, \quad (23)$$

$$\tilde{q}_3 = Q_3 q_0 - Q_2 q_1 + Q_1 q_2 + Q_0 q_3, \quad (24)$$

forming a new quaternion-valued vector  $\tilde{q} = \tilde{q}_0 + \tilde{q}_1 I + \tilde{q}_2 J + \tilde{q}_3 K$ , representing transformed SM field states (e.g., fermions or gauge bosons).

The square of the operator,  $q^2 = qq$ , is computed via matrix multiplication:

$$q^2 = \begin{pmatrix} q_0^2 - q_1^2 - q_2^2 - q_3^2 & -2q_0q_1 & -2q_0q_2 & -2q_0q_3 \\ 2q_0q_1 & q_0^2 - q_1^2 - q_2^2 - q_3^2 & -2q_0q_3 & 2q_0q_2 \\ 2q_0q_2 & 2q_0q_3 & q_0^2 - q_1^2 - q_2^2 - q_3^2 & -2q_0q_1 \\ 2q_0q_3 & -2q_0q_2 & 2q_0q_1 & q_0^2 - q_1^2 - q_2^2 - q_3^2 \end{pmatrix}, \quad (25)$$

or in quaternion form:

$$q^2 = (q_0^2 - q_1^2 - q_2^2 - q_3^2) + 2q_0q_1I + 2q_0q_2J + 2q_0q_3K, \quad (26)$$

verified by quaternion algebra: scalar part  $q_0^2 + q_1^2(-1) + q_2^2(-1) + q_3^2(-1) = q_0^2 - q_1^2 - q_2^2 - q_3^2$ , vector part  $2q_0(q_1I + q_2J + q_3K)$ , with cross terms (e.g.,  $q_1q_2K - q_2q_1K$ ) canceling. The vector component:

$$V_{q^2} = 2q_0(q_1I + q_2J + q_3K), \quad (27)$$

serves as pilot wave components, analogous to  $V$ , guiding fermion dynamics in a de Broglie-Bohm-like manner via phase gradients  $\nabla S_j$ . The scalar part influences energy scales, tied to  $\psi_0$ , modulating masses in  $Q$ 's dispersion relation.

## 2.6 Time Parameterization and Group Velocity

In this theory, the vector part  $V = 2(\psi_1I + \psi_2J + \psi_3K)$  parameterizes time through the phase evolution of its components, influencing relativistic mass and group velocity of particles via its connection to the Higgs mechanism. The quaternion components  $\psi_j$  ( $j = 1, 2, 3$ ) are expressed as:

$$\psi_j(x) = R_j(x)e^{iS_j(x)/\hbar}, \quad (28)$$

where  $R_j(x) \in \mathbb{R}_+$  is the amplitude,  $S_j(x) \in \mathbb{R}$  is the phase, and  $\hbar$  is the reduced Planck constant. The phase  $S_j$  evolves with proper time  $\tau$ , defining temporal dynamics relativistically. Momentum is derived from spatial gradients:

$$p_j = \nabla S_j, \quad j = 1, 2, 3, \quad (29)$$

forming the spatial components of the four-momentum  $p^\mu = (\psi_0, p_1, p_2, p_3)$ , with total momentum magnitude:

$$|\mathbf{p}| = \sqrt{p_1^2 + p_2^2 + p_3^2}. \quad (30)$$

Relativistic mass is  $m_{\text{rel}} = \gamma m$ , where  $\gamma = (1 - v^2/c^2)^{-1/2}$ , and  $m$  is the rest mass from the Higgs mechanism,  $m = y_f v' / \sqrt{2}$ , with  $v' = \psi_0 v_0$ ,  $v_0 = 246 \text{ GeV}$ . The dispersion relation for fields in  $Q$ , such as fermions, is:

$$\omega = \sqrt{c^2 |\mathbf{k}|^2 + m^2}, \quad (31)$$

yielding group velocity:

$$v_g = \frac{d\omega}{d|\mathbf{k}|} = \frac{c^2|\mathbf{k}|}{\sqrt{c^2|\mathbf{k}|^2 + m^2}}, \quad (32)$$

where  $|\mathbf{k}|$  is the wave number magnitude, tied to momentum  $|\mathbf{p}| = |\mathbf{k}|$  ( $\hbar = 1$ ). Along the direction of travel, assuming  $|\mathbf{p}| = mv$  classically:

$$v_g = \frac{c^2|\mathbf{p}|}{\sqrt{c^2|\mathbf{p}|^2 + m^2c^4}} = \frac{c^2(mv)}{\sqrt{c^2(mv)^2 + m^2c^4}} = \frac{vc}{\sqrt{v^2 + c^2}}. \quad (33)$$

For an electron ( $y_e \approx 2.1 \times 10^{-6}$ ): -  $\psi_0 = 1$ :  $v' = 246$  GeV,  $m_e = 0.511$  MeV,  $v_g = \frac{|\mathbf{k}|}{\sqrt{|\mathbf{k}|^2 + (0.511)^2}}$ , e.g., at  $|\mathbf{k}| = 1$  MeV,  $v_g \approx 0.859c$ , -  $\psi_0 = 2$ :  $v' = 492$  GeV,  $m_e = 1.022$  MeV,  $v_g = \frac{|\mathbf{k}|}{\sqrt{|\mathbf{k}|^2 + (1.022)^2}}$ ,  $v_g \approx 0.700c$ .

Variations in  $|V|^2$  modulate  $\psi_0$ , thus  $m$ , affecting  $v_g$  and connecting  $V$ 's time parameterization to relativistic dynamics and gravity via  $T_{\mu\nu}$ .

## 2.7 Orthogonality and Data Storage in $Q$

The total quantum state  $Q$  not only unifies Standard Model (SM) fields and gravity but also serves as a structured data repository characterized by perfect orthogonality, enabled by its quaternion basis and interaction with the operator  $q$ . The data space within  $Q$  is divisible and mutually acted upon by  $q$ , forming a system where orthogonality is preserved across transformations. We conceptualize  $Q$  as a 4x4x4 tensor—a 64-element cube—where each 4x4 slice corresponds to one of the quaternion basis elements  $(1, I, J, K)$ :

$$Q = \sum_{j=0}^3 Q_j \otimes e_j, \quad (34)$$

where  $Q_j$  are 4x4 matrices (e.g., from Eq. 15 in Section 2.3), and  $e_j = (1, I, J, K)$  are basis vectors forming a quaternion set. This tensorial representation extends  $Q$ 's 16-element 4x4 matrix into a richer structure, with each slice encoding independent physical data (e.g., SM field states or gravitational contributions).

The operator  $q$  interacts with  $Q$  bidirectionally through left and right multiplication:

$$Q' = qQ, \quad Q'' = Qq, \quad (35)$$

where  $Q' = \sum_{j=0}^3 (qQ_j) \otimes e_j$  and  $Q'' = \sum_{j=0}^3 (Q_jq) \otimes e_j$ . These operations rotate the 4x4x4 cube along any 3D direction defined by  $q = q_0 + q_1I + q_2J + q_3K$ , preserving the orthogonality of the basis elements, as  $e_i \cdot e_j = \delta_{ij}$  (where the dot product is the quaternion inner product, zero for  $i \neq j$  due to  $I^2 = J^2 = K^2 = -1$ ). This mutual action reflects a perfect symmetry:  $Qq$  transforms state vectors (as in Section 2.5), while  $qQ$  adjusts  $Q$ 's internal configuration, maintaining its unitary  $SU(4)$  properties.

This orthogonality enables  $Q$  to store data—quantum states, field configurations, or gravitational degrees of freedom—as a perfectly orthogonal system. Each 4x4 slice, tied to a basis element, can represent distinct physical entities (e.g., gluon states in one slice, fermion states in another), with  $q$ 's rotations ensuring independence across transformations. The 64 elements (4 basis elements  $\times$  16 matrix entries) provide a compact, high-dimensional data structure within 4D spacetime, enhancing  $Q$ 's role as a unified field



while aligning with the theory's SU(4)-based SM projections and gravitational sourcing via  $T_{\mu\nu}$ .

### 3 Higgs and Fermion Fields

#### 3.1 Higgs Field

The Higgs field  $\phi$  plays a critical role in this theory, mediating electroweak symmetry breaking and endowing particles with mass through its vacuum expectation value (VEV), integrated within the quaternion framework via  $\psi_0$ . It is an SU(2) doublet embedded in  $Q$ :

$$\phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}, \quad (36)$$

where  $\phi^+ = (\phi_1 + i\phi_2)/\sqrt{2}$  and  $\phi^0 = (\phi_3 + i\phi_4)/\sqrt{2}$  are complex charged and neutral components, transforming under SU(2)<sub>L</sub> (weak isospin) and U(1)<sub>Y</sub> (hypercharge  $Y = 1/2$ ). Unlike the SM's standalone field,  $\phi$  is dynamically linked to  $\psi$ :

$$\mathcal{L}_\phi = (D_\mu \phi)^\dagger (D^\mu \phi) - V(\phi), \quad (37)$$

with covariant derivative:

$$D_\mu = \partial_\mu - igW_\mu^a \tau^a - ig' \frac{Y}{2} B_\mu, \quad (38)$$

where  $W_\mu^a$  ( $a = 1, 2, 3$ ) are SU(2) gauge fields,  $B_\mu$  is the U(1) field,  $g \approx 0.65$ ,  $g' \approx 0.36$ , and  $\tau^a = \sigma^a/2$  (Pauli matrices). The potential is:

$$V(\phi) = -\mu^2 \phi^\dagger \phi + \lambda (\phi^\dagger \phi)^2, \quad \mu^2 > 0, \quad \lambda \approx 0.129, \quad (39)$$

minimized at:

$$\phi^\dagger \phi = \frac{\mu^2}{2\lambda}, \quad v' = \sqrt{\frac{\mu^2}{\lambda}}, \quad (40)$$

where  $v' = \psi_0 v_0$ , and  $v_0 = 246 \text{ GeV}$  when  $\psi_0 = 1$ . The VEV breaks SU(2)<sub>L</sub>  $\times$  U(1)<sub>Y</sub> to U(1)<sub>EM</sub>:

$$\langle \phi \rangle = \begin{pmatrix} 0 \\ \frac{v'}{\sqrt{2}} \end{pmatrix}, \quad (41)$$

mixing  $W_\mu^3$  and  $B_\mu$ :

$$W_\mu^3 = \cos \theta_W Z_\mu + \sin \theta_W A_\mu, \quad B_\mu = -\sin \theta_W Z_\mu + \cos \theta_W A_\mu, \quad (42)$$

$\sin^2 \theta_W \approx 0.231$ , yielding masses:

$$m_W = \frac{gv'}{2} \approx 80.4 \text{ GeV}, \quad m_Z = \frac{v' \sqrt{g^2 + g'^2}}{2} \approx 91.2 \text{ GeV}, \quad m_A = 0, \quad (43)$$

and Higgs boson mass:

$$m_H = \sqrt{2\lambda} v' \approx 125 \text{ GeV}, \quad \text{when } \psi_0 = 1. \quad (44)$$

Quantization gives:

$$\phi(x) = \begin{pmatrix} 0 \\ \frac{v' + h(x)}{\sqrt{2}} \end{pmatrix}, \quad h(x) = \int \frac{d^3k}{(2\pi)^3 2\omega_k} \left[ a_h(k) e^{-ik \cdot x} + a_h^\dagger(k) e^{ik \cdot x} \right], \quad (45)$$

$\omega_k = \sqrt{|\mathbf{k}|^2 + m_H^2}$ . The scalar  $\psi_0 = 2\sqrt{\psi_1^2 + \psi_2^2 + \psi_3^2}$  sources  $v'$ , normalizing to  $\psi_0 = 1$  for SM consistency, with variations (e.g.,  $\psi_0 = 2$ ) scaling masses (e.g.,  $m_H = 250 \text{ GeV}$ ), impacting  $T_{\mu\nu}$  and gravitational effects.

### 3.2 Fermion Fields

Fermions in this theory emerge from the vector part  $V = 2(\psi_1 I + \psi_2 J + \psi_3 K)$  of the quaternion wave function  $\psi$ , representing the matter content of the SM within  $Q$ . They are organized into representations consistent with  $\text{SU}(3)_C$ ,  $\text{SU}(2)_L$ , and  $\text{U}(1)_Y$ :

- **Quarks:** Left-handed doublets  $q_L = \begin{pmatrix} u_L \\ d_L \end{pmatrix}$ , right-handed singlets  $u_R, d_R$ , transforming as  $\text{SU}(3)$  triplets ( $q = (q_r, q_g, q_b)$ ), hypercharge  $Y = 1/6$ ,
- **Leptons:** Left-handed doublets  $l_L = \begin{pmatrix} \nu_{eL} \\ e_L \end{pmatrix}$ , right-handed singlets  $e_R$ ,  $Y = -1/2$  for doublets,  $Y = -1$  for  $e_R$ .

Quantization follows Dirac field formalism within  $Q$ :

$$\psi_f(x) = \int \frac{d^3k}{(2\pi)^3 2\omega_k} \sum_s \left[ b_{k,s} u_{k,s} e^{-ik \cdot x} + d_{k,s}^\dagger v_{k,s} e^{ik \cdot x} \right], \quad (46)$$

where  $\omega_k = \sqrt{|\mathbf{k}|^2 + m_f^2}$ ,  $b_{k,s}, d_{k,s}^\dagger$  are annihilation and creation operators for particles and antiparticles, satisfying:

$$\{b_{k,s}, b_{k',s'}^\dagger\} = (2\pi)^3 2\omega_k \delta_{ss'} \delta^3(\mathbf{k} - \mathbf{k}'), \quad \{d_{k,s}, d_{k',s'}^\dagger\} = (2\pi)^3 2\omega_k \delta_{ss'} \delta^3(\mathbf{k} - \mathbf{k}'), \quad (47)$$

$u_{k,s}, v_{k,s}$  are Dirac spinors, and  $m_f = y_f v' / \sqrt{2}$  from Yukawa coupling to  $\phi$ . The Lagrangian is:

$$\mathcal{L}_f = \bar{\psi}_f i D_\mu \gamma^\mu \psi_f - y_f \bar{\psi}_f \phi \psi_f, \quad (48)$$

with covariant derivative:

$$D_\mu = \partial_\mu - ig_s G_\mu^a T^a - ig W_\mu^a \tau^a - ig' Y B_\mu, \quad (49)$$

coupling fermions to  $\text{SU}(3)_C$  gluons ( $T^a = \lambda^a/2$ ),  $\text{SU}(2)_L$  W bosons, and  $\text{U}(1)_Y$  photons, e.g.,  $m_e = 0.511 \text{ MeV}$ ,  $m_t = 173 \text{ GeV}$ .

The vector  $V$  sources fermion states, with  $\psi_1, \psi_2, \psi_3$  providing three orthogonal directions, naturally yielding three generations (e.g.,  $e, \mu, \tau$ ) via wave interference, as  $V_{q^2} = 2q_0(q_1 I + q_2 J + q_3 K)$  guides their dynamics (Section 2.5). Contributions to  $Q$  include  $Q_V \sim V$ , with  $T_{\mu\nu} \propto |V|^2 = 1$  when  $\psi_0 = 1$ , dominating matter density in cosmic evolution.

## 4 Quantization Leading to Gravitons

### 4.1 Classical Gravity

In the classical regime, gravity arises from the stress-energy tensor  $T_{\mu\nu} = \text{Re}(Q^\dagger Q)$ , sourced by the total quantum state  $Q$ . The Einstein field equations govern this interaction:

$$G_{\mu\nu} = 8\pi G T_{\mu\nu}, \quad (50)$$

where  $G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu}$  is the Einstein tensor,  $R_{\mu\nu}$  the Ricci curvature tensor,  $R$  the scalar curvature, and  $g_{\mu\nu}$  the metric tensor. Here,  $Q(x) = Q_0(x) + Q_1(x)I + Q_2(x)J + Q_3(x)K$  is initially treated classically, with components  $Q_j(x) \in \mathbb{R}$ , encoding SM fields and their gravitational influence. This formulation aligns with general relativity, where  $T_{\mu\nu}$  dictates spacetime curvature.

### 4.2 Quantization

To quantize gravity, we promote  $Q(x)$  to a field operator:

$$Q(x) = \int \frac{d^3k}{(2\pi)^3 2\omega_k} \sum_{j=0}^3 \left[ a_j(k) Q_j(k) e^{-ik \cdot x} + a_j^\dagger(k) Q_j^\dagger(k) e^{ik \cdot x} \right], \quad (51)$$

rendering  $T_{\mu\nu} = \text{Re}(Q^\dagger Q)$  a Hermitian operator. We expand the metric in the weak-field limit:

$$g_{\mu\nu} = \eta_{\mu\nu} + \kappa h_{\mu\nu}, \quad (52)$$

where  $\eta_{\mu\nu} = \text{diag}(-1, 1, 1, 1)$ ,  $\kappa = \sqrt{8\pi G}$ , and  $h_{\mu\nu}$  is the graviton field. In harmonic gauge ( $\partial^\mu \bar{h}_{\mu\nu} = 0$ , where  $\bar{h}_{\mu\nu} = h_{\mu\nu} - \frac{1}{2}\eta_{\mu\nu} h^\lambda_\lambda$ ), the linearized Einstein equations become:

$$\square \bar{h}_{\mu\nu} = -2\kappa T_{\mu\nu}, \quad (53)$$

with  $\square = \partial^\mu \partial_\mu$ . The graviton field is quantized as:

$$h_{\mu\nu}(x) = \int \frac{d^3k}{(2\pi)^3 2\omega_k} \sum_{\lambda=\pm 2} \left[ a_\lambda(k) \epsilon_{\mu\nu}^\lambda(k) e^{-ik \cdot x} + a_\lambda^\dagger(k) \epsilon_{\mu\nu}^{\lambda*}(k) e^{ik \cdot x} \right], \quad (54)$$

where  $\omega_k = c|\mathbf{k}|$  (massless),  $a_\lambda(k)$  and  $a_\lambda^\dagger(k)$  satisfy:

$$[a_\lambda(k), a_{\lambda'}^\dagger(k')] = (2\pi)^3 2\omega_k \delta_{\lambda\lambda'} \delta^3(\mathbf{k} - \mathbf{k}'), \quad (55)$$

and  $\epsilon_{\mu\nu}^\lambda$  are symmetric, traceless polarization tensors for the two graviton helicity states.

### 4.3 Properties of Polarization Tensors

The polarization tensors  $\epsilon_{\mu\nu}^\lambda(k)$  characterize the two helicity states ( $\lambda = \pm 2$ ) of the massless spin-2 graviton field  $h_{\mu\nu}$ , defining its tensorial structure in the quantized theory. These tensors are essential to ensure that  $h_{\mu\nu}$  represents the physical degrees of freedom of gravitational waves, consistent with general relativity and quantum field theory. Their

properties are derived from the requirements of a massless, spin-2 particle propagating at the speed of light  $c$  and are outlined as follows:

1. **\*\*Symmetry\*\***: The tensors are symmetric in their spacetime indices:

$$\epsilon_{\mu\nu}^\lambda = \epsilon_{\nu\mu}^\lambda.$$

This property reflects the symmetry of the graviton field  $h_{\mu\nu}$ , which corresponds to perturbations of the metric tensor  $g_{\mu\nu}$ , inherently symmetric in general relativity.

2. **\*\*Tracelessness\*\***: The contraction with the Minkowski metric  $\eta_{\mu\nu} = \text{diag}(-1, 1, 1, 1)$  vanishes:

$$\eta^{\mu\nu} \epsilon_{\mu\nu}^\lambda = 0.$$

This ensures that  $h_{\mu\nu}$  has no scalar component (e.g., no breathing mode), restricting it to pure tensor modes, a hallmark of a massless spin-2 field with only two physical degrees of freedom.

3. **\*\*Transversality\*\***: The tensors are orthogonal to the wave vector  $k^\mu = (\omega_k/c, \mathbf{k})$ , satisfying the on-shell condition  $k^\mu k_\mu = 0$  (where  $\omega_k = c|\mathbf{k}|$ ):

$$k^\mu \epsilon_{\mu\nu}^\lambda = 0.$$

This condition guarantees that the graviton propagates transverse to its momentum direction, eliminating longitudinal modes and aligning with the massless nature of gravitational waves.

4. **\*\*Two Helicity States\*\***: The index  $\lambda = \pm 2$  labels the two independent polarization states, often referred to as the "plus" (+2) and "cross" (-2) modes in gravitational wave astronomy. These correspond to the quadrupolar deformations of spacetime, distinguishing the graviton from spin-1 particles (e.g., photons with  $\lambda = \pm 1$ ).

For a graviton propagating along the  $z$ -axis ( $k^\mu = (\omega/c, 0, 0, k_z)$ ), explicit forms in the transverse-traceless (TT) gauge are: -  $\lambda = +2$  (plus polarization):

$$\epsilon_{\mu\nu}^{+2} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix},$$

-  $\lambda = -2$  (cross polarization):

$$\epsilon_{\mu\nu}^{-2} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}.$$

These satisfy all required properties: symmetry ( $\epsilon_{ij} = \epsilon_{ji}$ ), tracelessness ( $\eta^{ij} \epsilon_{ij}^{+2} = 1 - 1 = 0$ ), and transversality ( $k^z \epsilon_{z\mu} = 0$ , as  $\epsilon_{z\mu} = 0$ ). In the quantization of  $h_{\mu\nu}$  (Eq. 47), these tensors ensure that only the two physical helicity states contribute, aligning with gravitational wave observations (e.g., LIGO detections) and the theory's prediction of gravitons as emergent from  $Q$ .

## 4.4 Implications

The quanta of  $h_{\mu\nu}$  are gravitons, massless spin-2 particles mediating gravity, emerging naturally from  $Q$ 's quantization. This unifies SM fields with gravity without extra dimensions, with  $T_{\mu\nu}$  coupling gravitons to matter. The absence of a bare cosmological constant suggests late-time acceleration arises from  $Q$ 's dynamics, testable via gravitational wave observations (e.g., LIGO).

## 5 Renormalization

### 5.1 Bare Lagrangian and Divergences

The bare Lagrangian is:

$$\mathcal{L}_0 = \mathcal{L}_{\text{SM}} + \mathcal{L}_Q + \mathcal{L}_h + \mathcal{L}_{\text{int}}, \quad (56)$$

with terms as defined previously. Perturbative loop diagrams yield UV divergences: logarithmic in SM and  $Q$  sectors (e.g., gauge self-energies), and power-law in gravity due to  $\kappa$ 's dimension  $(-1)$ .

### 5.2 Counterterms

The counterterm Lagrangian is:

$$\mathcal{L}_{\text{ct}} = \delta_Z \text{Tr}(D_\mu Q^\dagger D^\mu Q) + \delta_m V(Q) + \delta_{Z_h} \mathcal{L}_h + \delta_\kappa h_{\mu\nu} T^{\mu\nu} + \text{SM terms}, \quad (57)$$

absorbing divergences into redefined parameters.

### 5.3 Procedure

1. Compute loop amplitudes. 2. Use dimensional regularization ( $d = 4 - \epsilon$ ). 3. Identify poles (e.g.,  $1/\epsilon$ ). 4. Tune counterterms to cancel divergences. 5. Set physical parameters at scale  $\mu$  (e.g.,  $m_H = 125 \text{ GeV}$ ).

### 5.4 Gravity Challenges

Gravity's non-renormalizability arises from  $\kappa$ 's dimension, but  $\psi_0^2 = |V|^2$  may constrain divergences, suggesting an energy-dependent  $\kappa$  or finite counterterms, requiring loop calculations to confirm.

## 6 Cosmology and Structure Formation

### 6.1 Dynamics

The Friedmann equations:

$$\left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho, \quad \frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3p), \quad (58)$$

are sourced by  $T_{\mu\nu}$ , with early inflation from  $\psi_0 > 1$ , transitioning to matter domination ( $\rho \propto a^{-3}$ ) as  $\psi_0 = 1$ .

## 6.2 Structures

Perturbations in  $Q$  via  $V_{q^2}$  amplify clustering:

$$\ddot{\delta} + 2H\dot{\delta} - 4\pi G\rho\delta = S_q, \quad (59)$$

replacing dark matter with non-local effects.

## 6.3 Predictions

- CMB: Enhanced low- $\ell$  power. - Matter spectrum: Steeper small-scale slope. - Growth rate: Increased, testable via lensing.

# 7 Superluminal Propagation

## 7.1 Wave Packet

Wave packets in  $Q$ :

$$v_g = \frac{c^2|\mathbf{k}|}{\sqrt{c^2|\mathbf{k}|^2 + m^2}}, \quad (60)$$

potentially modified by  $q^2$ .

## 7.2 Hypothesis

$V_{q^2}$  pilot waves may yield:

$$\omega_k = \sqrt{c^2|\mathbf{k}|^2 + m^2} + f(V_{q^2}), \quad (61)$$

allowing  $v_g > c$  if  $f > 0$ .

## 7.3 Feasibility

Superluminal effects, constrained by causality, are testable via LHC or cosmological signals, likely non-physical.

# 8 Actions for Intuition

To conceptualize the theory intuitively, envision  $Q$  as a 4D rotating cube, each face encoding SM fields or gravity. The operator  $q$  rotates this cube, transforming states like a multidimensional prism, with  $q^2$ 's pilot waves steering particle trajectories. This analogy captures the unification of forces and spacetime dynamics, where  $\psi_0$  adjusts the cube's "size" (mass scale), and  $V$  directs its "motion" (fermion paths). Visualizing  $T_{\mu\nu}$  as the cube's "shadow" on spacetime illustrates gravity's emergence, offering an accessible bridge to the theory's mathematical formalism.

## 9 Interpretation Requirements

Validating the theory requires experimental tests leveraging current and future facilities:

- **Higgs Properties**: Precision measurements at the Large Hadron Collider (LHC) of  $m_H$  and couplings to probe  $\psi_0$  variations (e.g.,  $m_H = 250 \text{ GeV}$  if  $\psi_0 = 2$ ), testing deviations from SM predictions. - **Gravitational Waves**: LIGO/Virgo detection of early universe tensor modes from  $h_{\mu\nu}$ , verifying quantization and  $T_{\mu\nu}$  sourcing, with signatures differing from  $\Lambda\text{CDM}$ . - **Cosmological Probes**: Planck, Euclid, and LSST data to confirm no dark matter particles, assessing CMB power spectrum enhancements, matter clustering, and growth rates aligned with  $Q$  perturbations.

These experiments will constrain  $\psi_0$ ,  $q^2$ , and gravitational dynamics, potentially confirming the TOE's unification or necessitating refinements.