PHYS 170

Week 7: Kinematics I

Section 201 (Mon Wed Fri 12:00 – 13:00)

Intro remarks:

- Kinematics: Considers how (but not why) objects move
 - ...in contrast to Kinetics (Weeks 9&10), which analyses forces causing accelerated motion & motion per se
- We will need calculus!
- ...and vector analysis, and algebra, and trigonometry.









Rectilinear (1D) motion



Text: 12.1-12.3

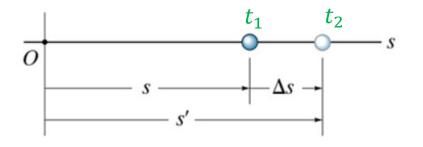
Content:

- Position, velocity, acceleration
- Velocity vs average velocity vs speed
- Connections $x(t) \Leftrightarrow v(t) \Leftrightarrow a(t)$
- Initial conditions: when and why do we need them?
- Special case: Motion with constant acceleration

Average velocity, Velocity, Speed & Average acceleration, Acceleration

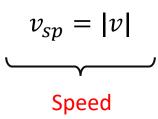
• Position of the object can change as the time goes:

- Velocity: shows how fast the coordinate changes
- Units: m/s, ft/s, km/h...



$$v_{av} = \frac{\Delta s}{\Delta t}$$
 $v = \lim_{\Delta t \to 0} \frac{\Delta s}{\Delta t} = \frac{ds}{dt}$

Velocity (can be positive or negative)



- Velocity of the object can change as the time goes:
 - $\begin{array}{c|cccc}
 & t_1 & t_2 \\
 \hline
 & v_1 & v_2
 \end{array}$

 Acceleration: shows how fast the velocity changes • Units: m/s², ft/s²...

$$a_{av} = \frac{\Delta v}{\Delta t}$$
 $a = \lim_{\Delta t \to 0} \frac{\Delta v}{\Delta t} = \frac{dv}{dt}$

SIGNS OF VELOCITY AND ACCELERATION

• Acceleration:

- The object is speeding up if $\vec{a} \uparrow \uparrow \vec{v}$
- The object is slowing down if $\vec{a} \uparrow \downarrow \vec{v}$

$$a_{av} = \frac{\Delta v}{\Delta t} = \frac{v_2 - v_1}{t_2 - t_1}$$

Assume $\Delta t = t_2 - t_1 = 1 \, s$, and that the object is slowing down from 3 m/s at t_1 to 1 m/s at t_2 .

Moves in:	v_1	v_2	а
$\rightarrow +s$	3 m/s	1 m/s	-2 m/s ²
<i>-s</i> ←	-3 m/s	-1 m/s	2 m/s ²

Indeed, a and v always have opposite signs.

When is the object slowing down?

When its speed decreases.

$$(\text{speed}) = \begin{cases} v, & \text{if } v > 0 \\ -v, & \text{if } v < 0 \end{cases}$$

$$\frac{d}{dt}(\text{speed}) = \begin{cases} \frac{d}{dt}(v) = a, & \text{if } v > 0\\ \frac{d}{dt}(-v) = -a, & \text{if } v < 0 \end{cases}$$

$$d(\text{speed})/dt < 0$$
 if: $\begin{cases} a < 0 \text{ and } v > 0 \\ a > 0 \text{ and } v < 0 \end{cases}$

... and a similar analysis for the case when an object is speeding up.

$$s(t) \Rightarrow v(t) \Rightarrow a(t)$$

•
$$s(t) \Rightarrow v(t)$$
: $v = \frac{ds}{dt} \equiv \dot{s}$

$$a ds = \frac{d\sigma}{dt} \cdot ds = d\sigma \left(\frac{ds}{dt}\right) = d\sigma \cdot \sigma$$

•
$$v(t) \Rightarrow a(t)$$
: $a = \frac{dv}{dt} = \dot{v}$

•
$$s(t) \Rightarrow a(t)$$
: $a = \frac{dv}{dt} = \frac{d}{dt} \left(\frac{ds}{dt}\right) = \frac{d^2s}{dt^2} \equiv \ddot{s}$

W7-1. Prove that for any 1D motion (with arbitrary acceleration, not necessarily constant!)

$$a ds = v dv$$

This equation is convenient when you want to determine the velocity of a particle, and the acceleration is given as a function of position, a = a(s).

$$a(t) \Rightarrow v(t) \Rightarrow s(t)$$

•
$$v(t) \Rightarrow s(t)$$
: $v = \frac{ds}{dt}$ $\Rightarrow s(t) = \int v \, dt + C_s$

•
$$a(t) \Rightarrow v(t)$$
: $a = \frac{dv}{dt}$ $\Rightarrow v(t) = \int a \, dt + C_v$

- This shows that when we are computing s(t) using known v(t), or computing v(t) using known a(t), we need to specify an <u>initial condition</u>: for example, if we know that $s(t = t_0) = s_0$, we will be able to find C_s . Likewise, if we know that $v(t = t_0) = v_0$, we will be able to find C_v .
- Actually, it is more convenient to rewrite these equations using definite integrals, as it is done on the next slide

Example

$$v(t) \Rightarrow s(t): \qquad v = \frac{ds}{dt} \qquad \Rightarrow \qquad ds = v dt$$

$$ds = \int v dt \qquad \Rightarrow \qquad S = \int v(t) dt$$

$$S(t) = \int v(t) dt$$

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$$S(t) = \int v(t) dt$$

$$\begin{cases} ds = v \, dt = v(t) dt, \\ s(t_0) = s_0 \end{cases}$$

Then:
$$\int_{s_0}^{s} ds = \int_{t_0}^{t} v(t) dt$$

$$s(t) - s_0 = \int_{t_0}^t v(t) dt$$

$$s(t) = s_0 + \int_{t_0}^t v(t) dt$$

$$\begin{cases} dv = a \, dt = a(t) dt, \\ v(t_0) = v_0 \end{cases}$$

Then:
$$\int_{v_0}^{v} dv = \int_{t_0}^{t} a(t) dt$$

$$v(t) - v_0 = \int_{t_0}^t a(t) dt$$

$$v(t) = v_0 + \int_{t_0}^t a(t) dt$$

$$\begin{cases} ds = v \, dt = v(t) \, dt \\ s(t_0) = s_0 \end{cases}$$

Then:
$$\int ds = \int v(t) dt$$

$$s(t) = \int v(t) dt + C$$

and we will choose C in such a way, that at $t=t_0$ we would get: $s(t_0)=s_0$.

Let's summarize:

• What we will do in this chapter: $s(t) \Leftrightarrow v(t) \Leftrightarrow a(t)$

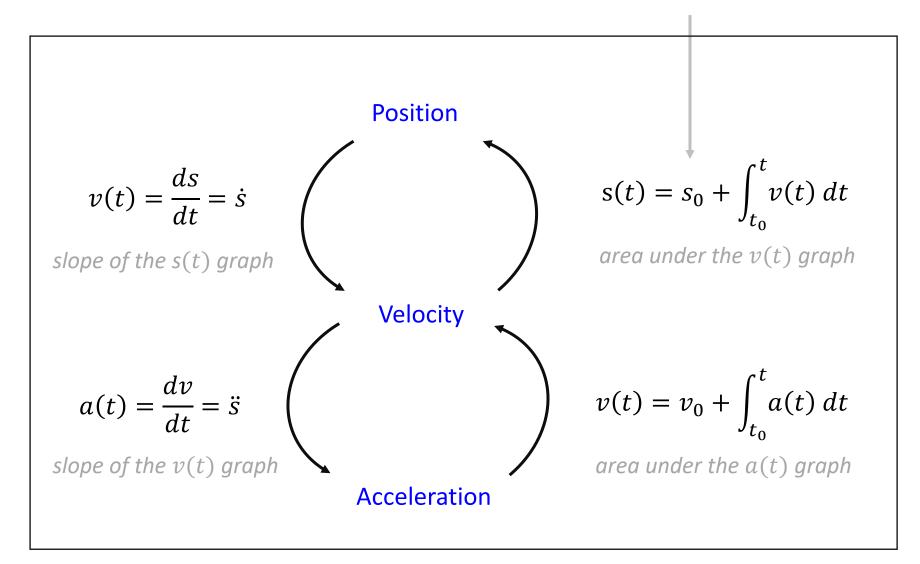
• When integrating, we will need initial conditions (s_0, v_0)

- Toolkit:
 - Derivatives
 - > Integration
- Notations:

$$\frac{ds}{dt} \equiv \dot{s}$$

$$\frac{d^2s}{dt^2} \equiv \ddot{s}$$

$$\frac{d^3s}{dt^3} \equiv \ddot{s}$$



. . .

What to expect:

- We will learn how, knowing one of the functions s(t), v(t), a(t), to derive all other functions.
- We will work in 1D, 2D and 3D, and in a variety of coordinate systems.

- One specific kinematic case known to most of you from high school is motion with a constant acceleration, such as free fall. We will start with it.
- However, our accent will be on a general situation, when a = a(t). It is only natural that objects are not bound to move keeping their acceleration constant; think of a car which first speeds up $(\vec{a} \uparrow \uparrow \vec{v})$, then goes with a constant speed $(\vec{a} = 0)$, then slows down $(\vec{a} \uparrow \downarrow \vec{v})$. Here $\vec{a} \neq const$.
 - $> j(t) = \frac{da}{dt} = \frac{d^3s}{dt^3}$: "jerk" (important for roller coasters)
 - $\sigma(t) = \frac{dj}{dt} = \frac{d^4s}{dt^4}$: "snap" (and next ones are "crackle" and "pop"... seriously)
- Important to remember when integrating and differentiating:
 - > Differentiation is about 1 point (local in time domain)
 - > Integration requires knowing initial conditions (result depends on pre-history)
- Review calculus !!!

Q: An object travels with velocity $v = 3t^2$ m/s, where t is measured in seconds. How far does it travel between t = 1 s and t = 2 s?

- A. 12 m
- B. 11 m
- C. 9 m
- D. 7 m
- E. None of the above

$$S(t=2s) \qquad t=2s$$

$$\int ds = \int \sigma(t) dt$$

$$S(t=1s) \qquad t=1s$$

IMPORTANT SPECIAL CASE: Motion with constant acceleration

W7-2. Sometimes acceleration does not change. Important **example** is a free fall; near the surface of the Earth its magnitude is g = 9.81 m/s² and it acts downwards.

Prove that in this case (a is the acceleration, s_0 and v_0 are the object's position and velocity at t=0):

$$s(t) = s_0 + v_0 t + \frac{a t^2}{2}$$

$$s(t) = s_0 + v_0 t + \frac{a t^2}{2}$$
 (1)

$$v(t) = v_0 + a t$$

(2)
$$S_0$$
 S_0 S_0

$$2a(s-s_0)=v^2-s_0^2$$

$$v^2(t) = v_0^2 + 2a(s - s_0)$$

$$v^{2}(t) = v_{0}^{2} + 2a(s - s_{0})$$
(2) $a = \frac{d\sigma}{dt}$

$$\sigma = \int_{\sigma} d\sigma = \int_{\sigma} d\sigma = \int_{\sigma} d\tau =$$

(1)
$$V = \frac{dS}{dt} \rightarrow \int dS = \int v \, dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt + \int at \, dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int v_0 dt \rightarrow S - So = \int \left[v_0 + at \right] dt = \int \left[$$

Curvilinear (2D, 3D) motion: Cartesian components

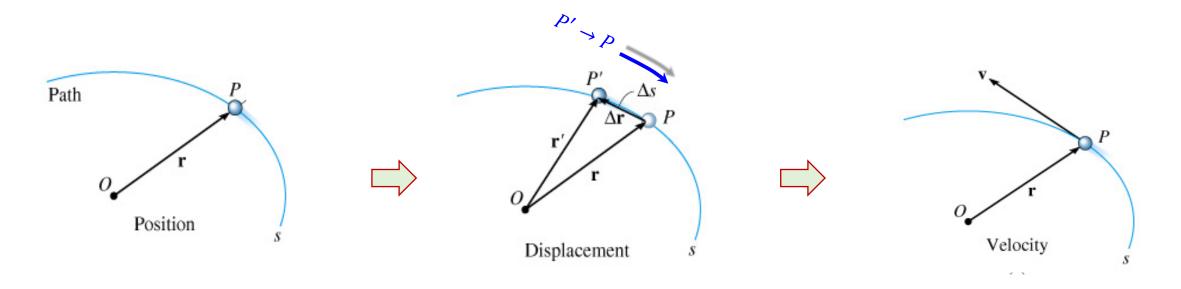


Text: 12.4-12.5

Content:

- Beyond 1D: Definitions of s(t), v(t), a(t)
- Graphical interpretation
- Equations of motion in Cartesian components

Velocity & Acceleration in 2D: Graphical Approach



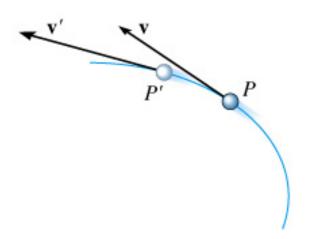
• Trajectory: P = r(t)

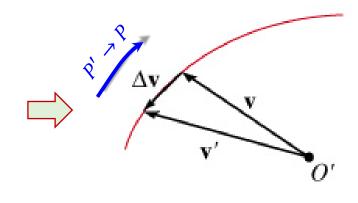
- Draw a small shift along the trajectory, $\Delta \vec{r}$
- $\Delta \vec{r}$ = displacement vector, Δs = its magnitude
- Let $P' \to P$: $\Delta \vec{r} \to d\vec{r}$, $\Delta s \to ds$

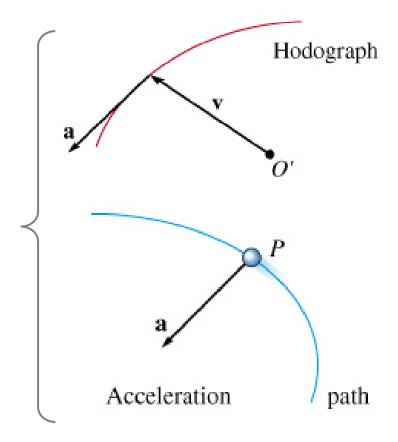
- Velocity:
 - > Tangent to the trajectory
 - ightharpoonup Magnitude: v = ds/dt
 - > We can also say that:

$$\vec{v} = \frac{d\vec{r}}{dt}$$

Velocity & Acceleration in 2D: Graphical Approach







- Draw trajectory, P = r(t), and add the velocity vectors at each point
- Draw the hodograph (the curve which the velocity arrowheads touch)
- $\Delta \vec{v} =$ difference between two close velocities (at moments t' and t)
- Let $t' \to t$: $\Delta \vec{v} \to d\vec{v}$
- Then we will get:

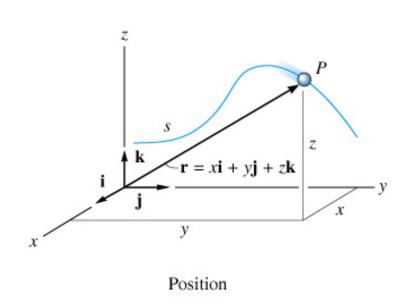
$$\vec{a} = \frac{d\vec{v}}{dt}$$

Acceleration:

- > Tangent to the hodograph
- Points inwards the trajectory
- Must account for the change in the direction of v, and also for change in its magnitude!

Velocity & Acceleration in Rectangular Components

• These pictures are nice, but it is difficult to work with them. Let us come up with something else.



• If
$$\vec{r} = x\vec{i} + y\vec{j} + z\vec{k}$$
, what is $\frac{d\vec{r}}{dt} = \vec{v}$?

• Product rule:
$$\frac{d(ab)}{dt} = a\frac{db}{dt} + b\frac{da}{dt}$$

• Product rule:
$$\frac{d(ab)}{dt} = a\frac{db}{dt} + b\frac{da}{dt}$$
• Then:
$$\frac{d(x\vec{i})}{dt} = x\frac{d\vec{i}}{dt} + \vec{i}\frac{dx}{dt}$$

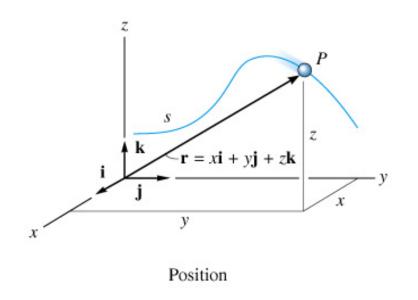
• Note that: $\frac{d\vec{i}}{dt} = 0$ (\vec{i} does not change with t)

$$\vec{v} = v_x \vec{i} + v_y \vec{j} + v_z \vec{k}$$

where

$$v_x = \frac{dx}{dt} = \dot{x}, \qquad v_y = \frac{dy}{dt} = \dot{y}, \qquad v_z = \frac{dz}{dt} = \dot{z}$$

Velocity & Acceleration in Rectangular Components



• If
$$\vec{r} = x\vec{i} + y\vec{j} + z\vec{k}$$
:

$$\vec{v} = v_x \vec{i} + v_y \vec{j} + v_z \vec{k}$$

$$\vec{a} = a_x \vec{i} + a_y \vec{j} + a_z \vec{k}$$

$$v_{x} = \frac{dx}{dt} = \dot{x},$$

$$v_x = \frac{dx}{dt} = \dot{x}, \qquad v_y = \frac{dy}{dt} = \dot{y}, \qquad v_z = \frac{dz}{dt} = \dot{z}$$

$$v_z = \frac{dz}{dt} = \dot{z}$$

$$a_{x} = \frac{dv_{x}}{dt} = \ddot{x}$$

$$a_{y} = \frac{av_{y}}{dt} = \ddot{y},$$

$$a_x = \frac{dv_x}{dt} = \ddot{x}, \qquad a_y = \frac{dv_y}{dt} = \ddot{y}, \qquad a_z = \frac{dv_z}{dt} = \ddot{z}$$

- Note: we now have three one-dimensional problems (which we already know how to work with!)
- We can use these algebraic equations to find the components of $\vec{r}(t)$, $\vec{v}(t)$ and $\vec{a}(t)$
 - Q: These quasi-1D-problems are not completely independent. What connects them?