Lecture 29.

B-field due to a wire (short segment and long wire). "Two-step logic" in magnetic problems.

Ampere's law.

Announcement

Midterm 2 – next Monday

What is on MT2?

- All electricity (Coulomb law and on)
- Cutoff: Up to Lorentz force (Week 9)

Resources?

- Lectures, textbook, HW, tutorial problems...
- Posted practice exams

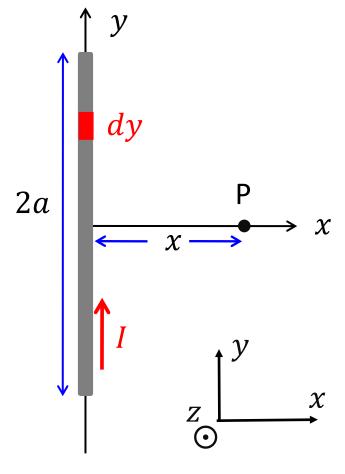
Help sessions?

- Monday 17:30-18:30 (Henn 318)
- Wednesday 17:00 18:00 (Henn 318)
- Friday 17:00 18:00 (Henn 318) & on Zoom
- Saturday 14:00 16:00 (Hebb 112)

Magnetic field of a short straight wire

Last Time

Q: Use the Biot-Savart Law to compute magnetic field B created by a short current segment of length 2a with current I at the symmetry axis of the wire. Your answer should be a vector.



Exercise: Before you do the math, think about how to solve the problem and write a few sentences outlining your strategy.

• First consider a small wire segment dy => find the field $d\vec{B}$ produced by it at P => considering symmetry, integrate its components to get the resultant field at P

You might need:

$$\int \frac{x \, dy}{(x^2 + y^2)^{3/2}} = \frac{y}{x \sqrt{x^2 + y^2}}$$

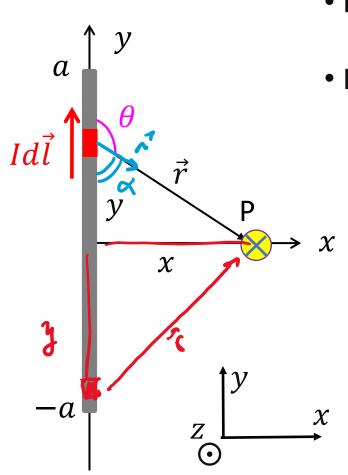
Biot-Savart law:

$$d\vec{B} = \frac{\mu_0}{4\pi} \frac{I \ d\vec{l} \times \hat{r}}{r^2}$$

Magnetic field of a short straight wire $\frac{1}{dB} = \frac{1}{\sqrt{L}} \frac{1}{\sqrt{L}}$

$$d\vec{B} = \frac{r_0}{\sqrt{\pi}} \frac{1}{r^2} \sin \theta = \sin(\pi - \theta) = x/r$$

Q: Use the Biot-Savart Law to compute magnetic field B created by a short current segment of length 2a with current I at the symmetry axis of the wire. Your answer should be a vector.



• Direction?
$$d\vec{l} \times \hat{r} \Rightarrow \text{along } (-\vec{k}) = \text{into the page}$$

• Magnitude?
$$\left| I \ d\vec{l} \times \hat{r} \right| = I \ dl \ |\hat{r}| \sin \theta = I \ dy \sin \theta$$

$$\vec{B}_{\text{at P}} = \int d\vec{B} = \int \frac{\mu_0}{4\pi} \frac{I \, d\vec{l} \times \hat{r}}{r^2} = (-\hat{k}) \frac{\mu_0 I}{4\pi} \int \frac{\sin \theta \, dy}{r^2}$$

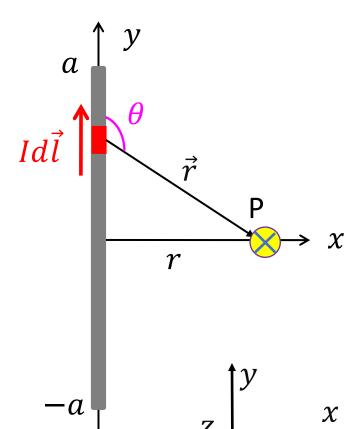
$$= (-\hat{k}) \frac{\mu_0 I}{4\pi} \int \frac{x \, dy}{r^3} = (-\hat{k}) \frac{\mu_0 I \, x}{4\pi} \int_{-a}^{a} \frac{dy}{(x^2 + y^2)^{3/2}}$$

$$\vec{B}_{\text{at P}} = -\frac{\mu_0 I}{4\pi} \frac{2a}{x\sqrt{x^2 + a^2}} \,\hat{k}$$

- y = integration variable
- x = parameter

Magnetic field of an infinitely long straight wire

Q: Use the Biot-Savart Law to compute magnetic field B created by an infinitely long currentcarrying wire with current I at a perpendicular distance r from the wire.



• Finite wire:
$$\vec{B}_{\text{at P}} = -\frac{\mu_0 I}{4\pi} \frac{2a}{r\sqrt{\sqrt{2} + a^2}} \, \hat{k}$$

• Take limit: $a\gg r$ (neglect r^2 in comparison with a^2)

$$\frac{\mu_0 I}{4\pi} \frac{2a}{r\sqrt{r^2 + a^2}} \to \frac{\mu_0 I}{4\pi} \frac{2a}{r\sqrt{a^2}} = \to \frac{\mu_0 I}{2\pi r}$$

$$\vec{B}_{\text{at P}} = -\frac{\mu_0 I}{2\pi r} \; \hat{k}$$

We will see very shortly that there is a much easier way to compute the B-field for a long straight wire – Ampere's Law

A wire carries current *I* into the junction. What is the magnitude of the B-field at point P? (NOTE: both wire segments are very, very long)



A.
$$B = \frac{2\mu_0 R}{\pi x}$$

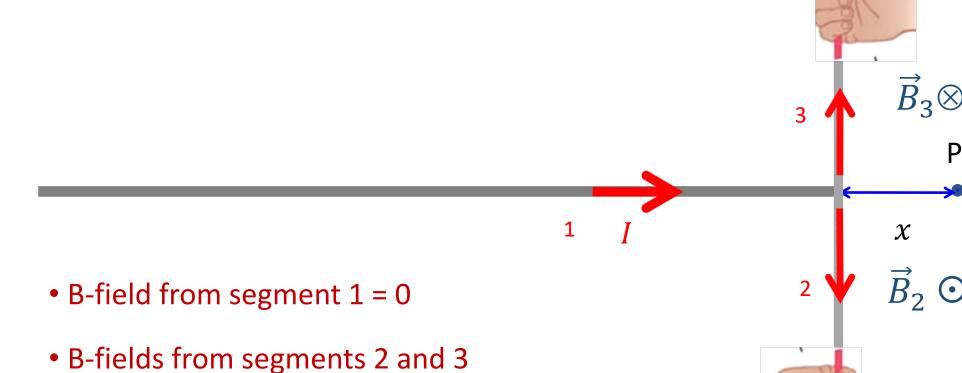
$$B. B = \frac{\mu_0 I}{\pi x}$$

C.
$$B = \frac{\mu_0 I}{2\pi x}$$

D.
$$B = \frac{\mu_0 I}{4\pi x}$$

E.
$$B = 0$$

A wire carries current *I* into the junction. What is the magnitude of the B-field at point P? (NOTE: both wire segments are very, very long)



A.
$$B = \frac{2\mu_0 R}{\pi x}$$

$$B. B = \frac{\mu_0 I}{\pi x}$$

C.
$$B = \frac{\mu_0 I}{2\pi x}$$

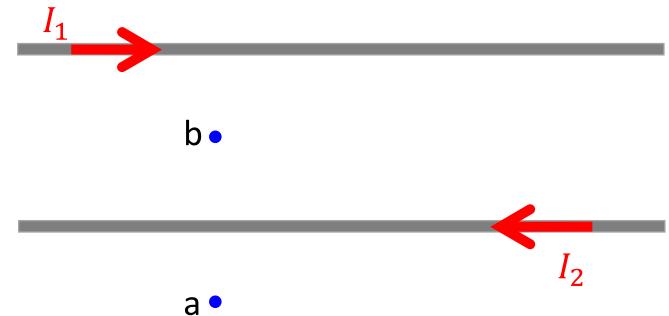
D.
$$B = \frac{\mu_0 I}{4\pi x}$$

$$(E.)B = 0$$

• Half-wire B-field: coming soon!

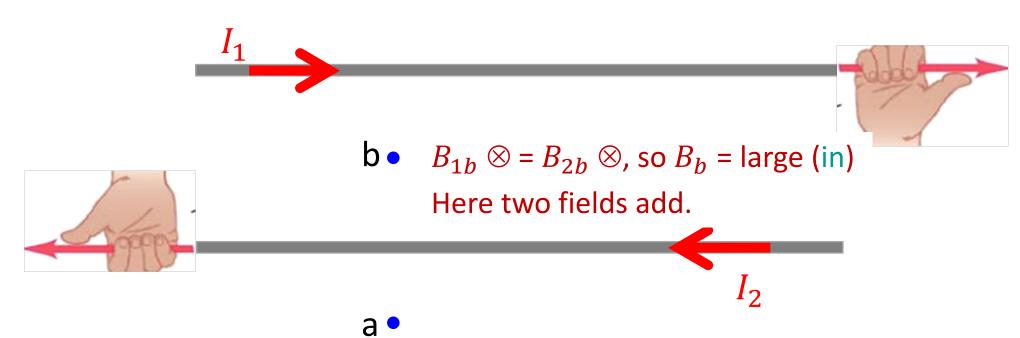
cancel one another by symmetry

Q: Compare the magnitude and direction of the B-fields at points a and b. ("out"=out of the page, "in"=into the page). NOTE: both wire segments are infinitely long and $I_1 = I_2$.



- A. B_a (in) $< B_b$ (in)
- B. $B_a(out) < B_b(out)$
- C. B_a (out) > B_b (in)
- D. B_a (out) $< B_b$ (in)
- E. B_a (in) > B_b (out)

Q: Compare the magnitude and direction of the B-fields at points a and b. ("out"=out of the page, "in"=into the page). NOTE: both wire segments are infinitely long and $I_1 = I_2$.



- A. B_a (in) $< B_b$ (in)
- B. $B_a(out) < B_b(out)$
- C. B_a (out) > B_b (in)
- D. B_a (out) $< B_b$ (in)
- E. B_a (in) > B_b (out)

$$B_{1a} \otimes \langle B_{2a} \odot \rangle$$
, so $B_a = \text{small (out)}$

Here two fields partially subtract.

$$B_{\text{wire}}(r) = \frac{\mu_0 I}{2\pi r}$$

Q: A long straight wire carries current I out of the page.

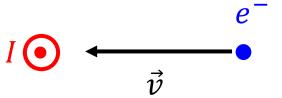
An electron travels towards the wire from the right. What is the direction of the <u>force</u> on the electron?











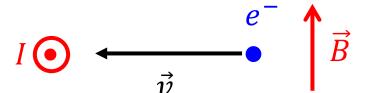
Q: A long straight wire carries current I out of the page.

An electron travels towards the wire from the right. What is the direction of the <u>force</u> on the electron?

 $B = \frac{\mu_0 I}{2\pi r}$

Force on the electron:

$$\vec{F}_m = q_+ \, \vec{v} \times \vec{B}$$





C. ←

• Why there is a force on the electron?? – Because it is a charged particle moving in a magnetic field.

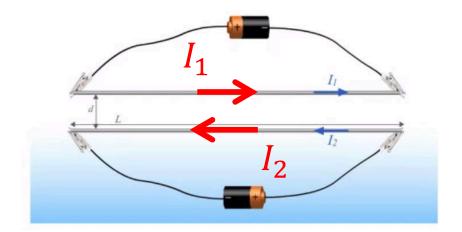
 $D. \longrightarrow$

Who creates this magnetic field?? – The wire!

E. ↓

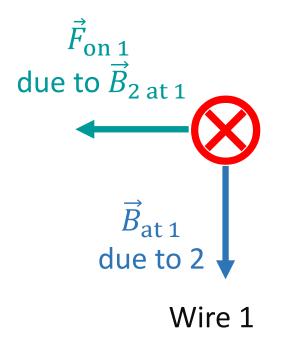
 "Two-step logic": object #1 creates magnetic field at the location of object #2 => object #2 experiences magnetic force! Q: Two wires carry current in opposite directions. What the wires will do?

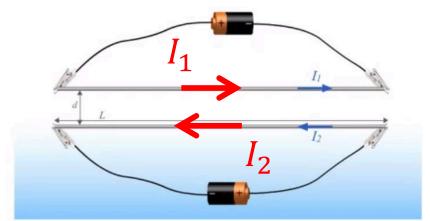
- A. Attract
- B. Repel
- C. No effect on each other

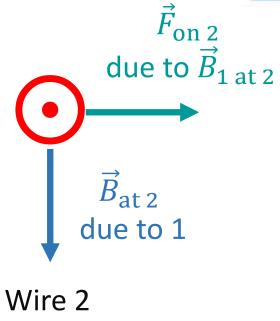


Q: Two wires carry current in opposite directions. What the wires will do?

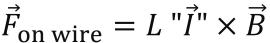
- A. Attract
- (B.) Repel
- C. No effect on each other





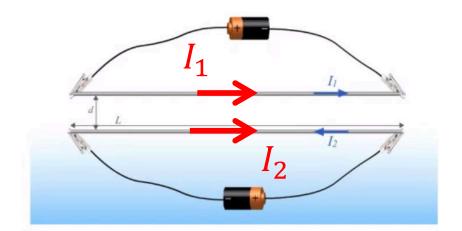






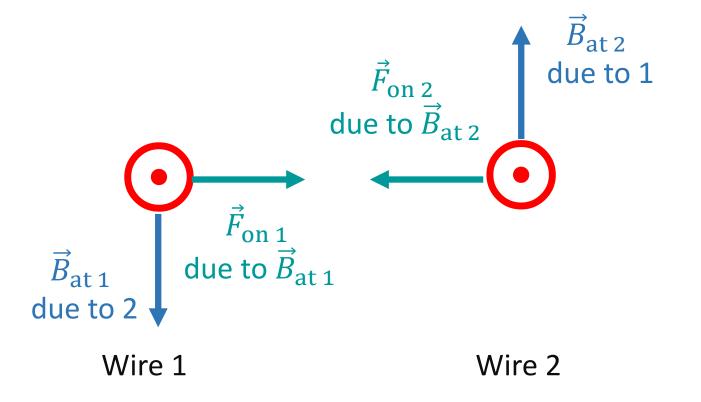
Q: Two wires carry current in the same directions. What the wires will do?

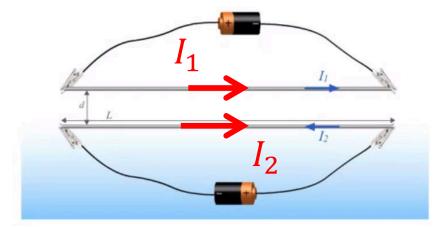
- A. Attract
- B. Repel
- C. No effect on each other



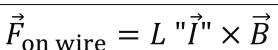
Q: Two wires carry current in opposite directions. What the wires will do?

- (A.) Attract
 - B. Repel
 - C. No effect on each other



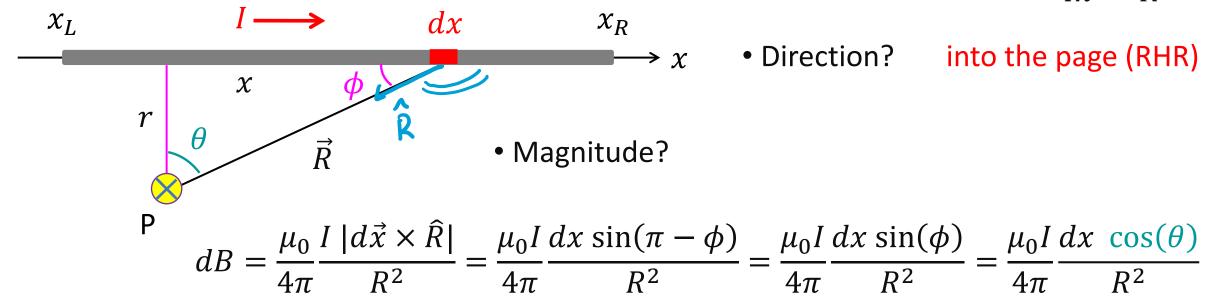






B-Field of a short straight wire (off-center)

$$d\vec{B} = \frac{\mu_0}{4\pi} \frac{I \ d\vec{x} \times \hat{R}}{R^2}$$

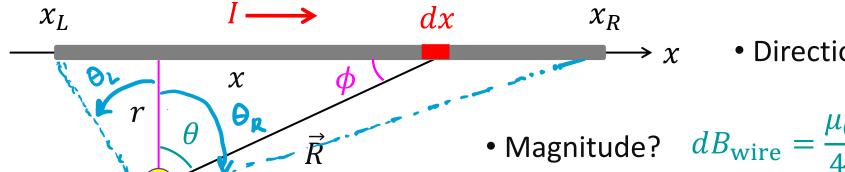


• Note: integration is much easier using θ :

$$x = r \tan \theta \implies \frac{dx}{d\theta} = \frac{r}{\cos^2 \theta} \implies dx = \frac{r d\theta}{\cos^2 \theta} \implies dB = \frac{\mu_0 I \sqrt{d\theta \cos(\theta)}}{4\pi} \frac{1}{\cos^2(\theta) R^2}$$
... and $R \cos(\theta) = r \implies dB_{\text{wire}} = \frac{\mu_0 I}{4\pi} \frac{d\theta \cos(\theta)}{r}$

B-Field of a short straight wire (off-center)

$$d\vec{B} = \frac{\mu_0}{4\pi} \frac{I \ d\vec{x} \times \vec{R}}{R^2}$$



• Direction? into the page (RHR)

• Magnitude?
$$dB_{\rm wire} = \frac{\mu_0 I}{4\pi} \frac{d\theta \cos(\theta)}{r}$$

$$B_{\text{wire}} = \frac{\mu_0}{4\pi} \frac{I}{r} \int_{\theta_L}^{\theta_R} d\theta \cos(\theta) \qquad B_{\text{wire}} = \frac{\mu_0}{4\pi} \frac{I}{r} (\sin(\theta_R) - \sin(\theta_L)) \qquad \text{etc} \qquad \sin(\theta_R) = \frac{x_R}{\sqrt{x_R^2 + r^2}},$$

$$B_{\text{wire}} = \frac{\mu_0}{4\pi} \frac{I}{r} (\sin(\theta_R) - \sin(\theta_L))$$

$$\sin(\theta_R) = \frac{x_R}{\sqrt{x_R^2 + r^2}}$$
 etc.

• Infinite wire: $\sin(\theta_R) = +1$, $\sin(\theta_L) = -1$, and we get:

$$B_{\text{LOOOONG wire}} = \frac{\mu_0}{2\pi} \frac{I}{r}$$

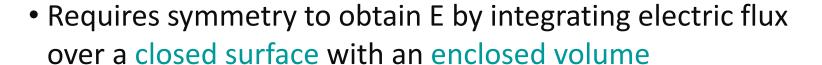
• Semi-Infinite wire with $x_L = 0$: $\sin(\theta_R) = +1$, $\sin(\theta_L) = 0$:

$$B_{\text{half-wire}} = \frac{\mu_0}{4\pi} \frac{I}{r}$$

Did you enjoy this derivation?

• Let's recall: we could derive the electric field of a very long line of charge by integrating Coulomb law – but there was an alternative:

$$\Phi_e = \oint \vec{E} \cdot d\vec{A} = \frac{Q_{in}}{\varepsilon_0}$$



• Solving for B_p for a very long current-carrying wire (i.e. $a \gg x$) using Biot-Savart's Law was also hard, but there is a MUCH easier alternative!

$$\Phi_m = \oint \vec{B} \cdot d\vec{l} = \mu_0 I_{in}$$

• Requires symmetry to obtain B by integrating magnetic flux over a closed path with an enclosed area

