# Announcement



From: Engineering Undergraduate Society

To celebrate the end of midterms, we are putting on one more event called "PP semi" which is essentially a semi formal dance for first year engineering students and their plus ones. The event will take place on the 27th of March from 7:30-10:00pm.

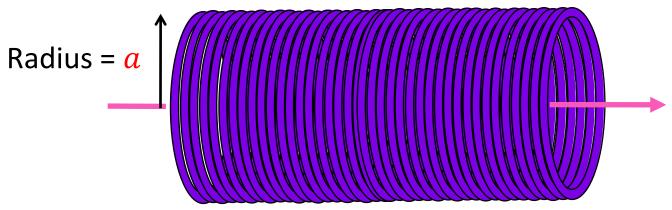
Lecture 31.

Real-life solenoids.

Faraday's law & Lenz's law: making sense

## Ideal solenoid

# Last Time:

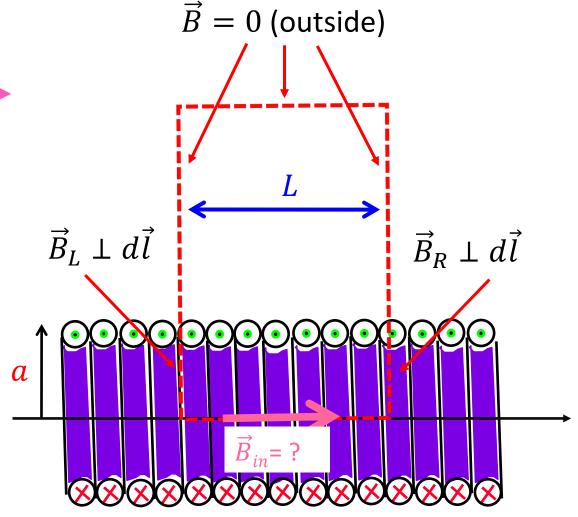


• Ampere's law:  $\oint \vec{B} \cdot d\vec{l} = \mu_0 I_{\rm encl}$ 

$$\oint \vec{B} \cdot d\vec{l} = B_{in}L$$

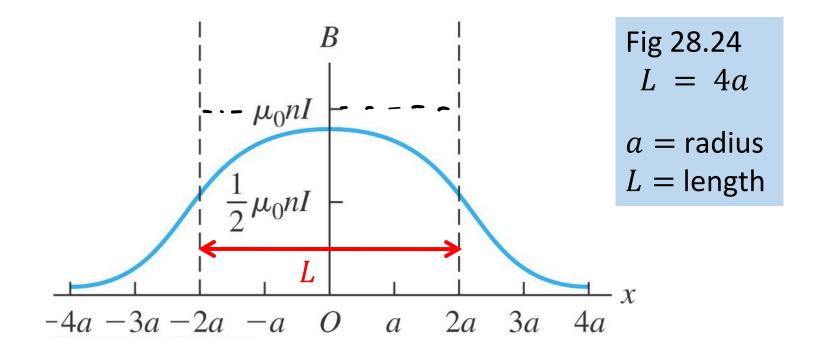
$$I_{\text{encl}} = InL$$

 $B_{in} = \mu_0 In$ 



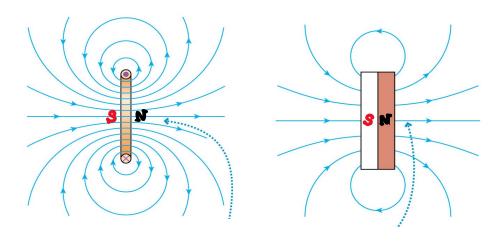
#### Finite Length Solenoid

Unfortunately, Engineers have to deal with finite length solenoids!



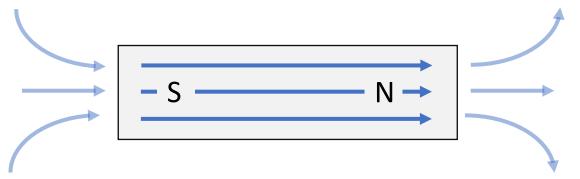
#### Fringe effects

To better understand fringe effects, let's look at the B-field of one ring

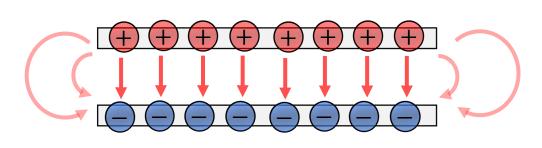


- If we look at the field of a current-carrying ring, we see that it is equivalent to the field of a tiny magnet!
  - You can even define North and South poles for such a loop...

 Putting many loops together (solenoid / magnet):

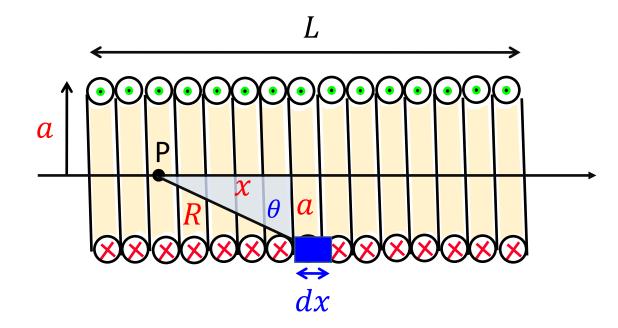


Doesn't it remind a capacitor?



• Store magnetic / electric energy, uniform field...

# Finite Length Solenoid



• Same way as with a finite wire: introduce  $\theta$ 

• Recall  $B_{\chi}$  on the axis of a circular loop:

$$B_{x}(x) = \frac{\mu_0}{2} \frac{I a^2}{(x^2 + a^2)^{3/2}}$$

• Consider magnetic field created at the observation point P by a tiny segment of solenoid (thickness dx):

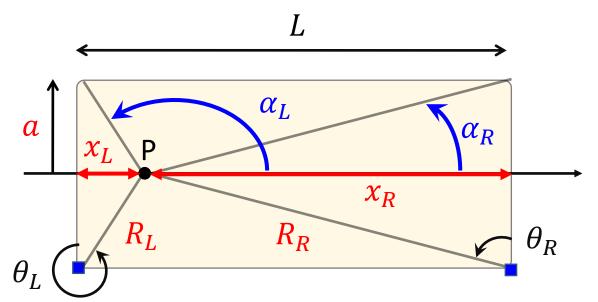
$$dB_{x}(x) = \frac{\mu_0}{2} \frac{Ndx}{L} \frac{Ia^2}{(x^2 + a^2)^{3/2}}$$

$$x = a \tan \theta \implies dx = \frac{a \ d\theta}{\cos^2 \theta}$$
, and  $R \cos(\theta) = a \implies dB_x = \frac{\mu_0 N}{2 L} I \cos(\theta) \ d\theta$  it!

$$B_{\text{real solenoid}}(P) = \int dB_x = \frac{\mu_0}{2} \frac{N}{L} I \left\{ \sin(\theta) = \pm \frac{x}{\sqrt{x^2 + a^2}} \right\} \Big|_{x_L}^{x_R}$$

## Finite Length Solenoid & Trigonometry

$$B_{\text{real solenoid}}(@P) = \frac{\mu_0}{2} \frac{N}{L} I(\sin \theta_R - \sin \theta_L)$$



• Let's brush it up:

$$\sin \theta_L = -\sin(2\pi - \theta_L) = -\frac{x_L}{R_L} = \cos \alpha_R$$

$$\sin \theta_R = \frac{x_R}{R_R} = \cos \alpha_R$$

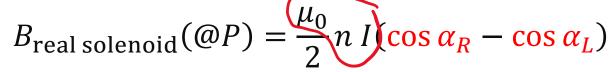
• We get:

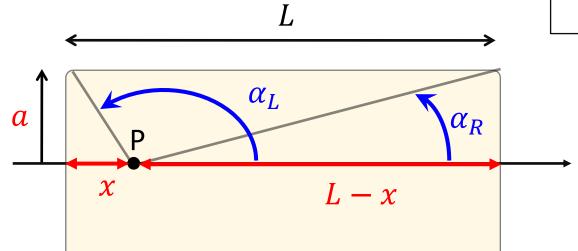
$$B_{\text{real solenoid}}(@P) = \frac{\mu_0}{2} n I(\cos \alpha_R - \cos \alpha_L)$$

- n = N/L is the density of coils
- $\alpha_L$  and  $\alpha_R$  are the angles at which you see the edges of the solenoid from point P.
- For a long solenoid:  $\alpha_R \to 0$  &  $\alpha_L \to \pi$   $\Rightarrow$

 $B_{\text{ideal solenoid}} = \mu_0 \ n \ I = \text{const}$ 

## Finite Length Solenoid & Trigonometry





 $B_{\text{ideal solenoid}} = \mu_0 \, n \, I = \text{const}$ 

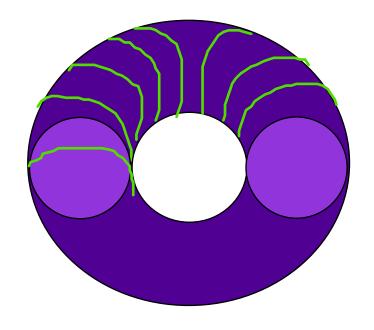
$$B_{\text{real solenoid}}(x) = B_{\text{ideal}} \times \frac{1}{2} \left( \frac{L - x}{\sqrt{a^2 + (L - x)^2}} + \frac{x}{\sqrt{a^2 + x^2}} \right)$$

• Show that the approximation of ideal solenoid works well when  $L\gg a$ 

L~ 4a-fringe effects

#### Magnetic Applications – Toroid Solenoid

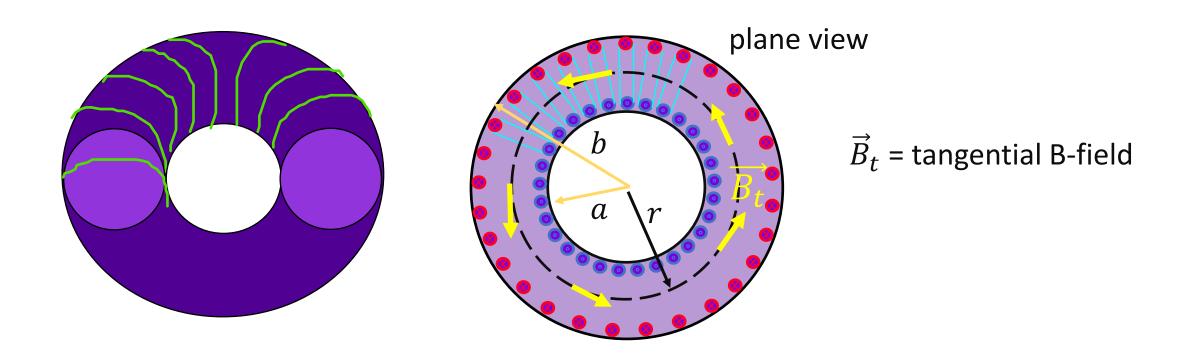
• Let's get rid of fringes at the ends – by getting rid of the ends themselves!



- Here we show a <u>toroid</u> with a circular cross-section, i.e. a "doughnut" wrapped with N "turns" of wire.
- The wire carries current I
- This is a solenoid twisted into a circle.

• What the magnetic field in this structure is?

## Magnetic Applications – Toroid Solenoid



$$\oint \vec{B} \cdot d\vec{l} = 2\pi r \, B_t = \mu_0 I_{encl} = \mu_0 IN$$

Note -- 
$$B \cong 0$$
 for  $r < a$  or  $r > b$ 

$$B_t = \frac{\mu_0 IN}{2\pi r}$$

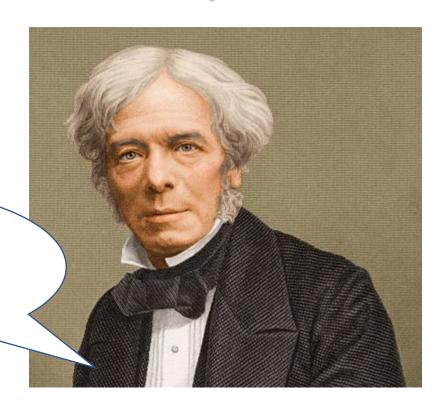
# Lenz's law



$$\varepsilon = -\frac{d\Phi_B}{dt}$$

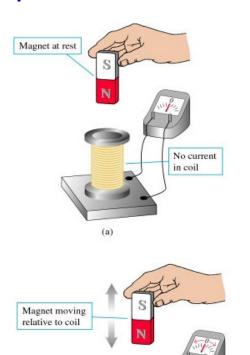
# Faraday's law

$$\varepsilon = \left| \frac{d\Phi_B}{dt} \right|$$

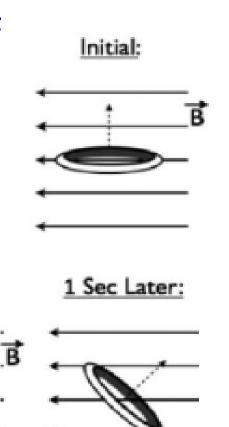


# **Experiments on induced EMF**

Current induced

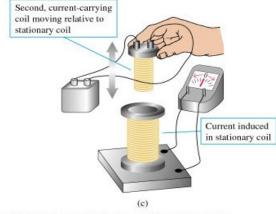






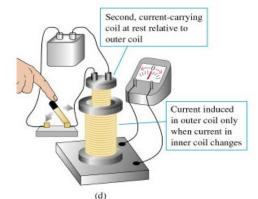
 Rotating a loop in an external magnetic field creates current in the loop

 Moving a current-carrying solenoid around a coil creates current in the solenoid



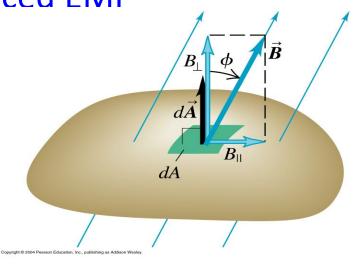
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 Switching a current on in the inner coil momentarily creates current in the outer coil



• This experiment works for any closed (so that it can support electric current) loop

#### Induced EMF



$$\Phi_B = \oint \vec{B} \cdot d\vec{A}$$

- Common denominator:
  - ➤ In these experiments, there in no battery in the loop. The current is excited literally by changing the magnetic flux through the loop!

- Okay, we have current => there must be EMF
- We can define the so-called "motional emf": (emf = electromotive force)

$$\varepsilon = \left| \frac{d\Phi_B}{dt} \right|$$

an emf induced in a conducting loop if the magnetic flux  $\Phi_B$  through it changes



"emf responsible for that current"

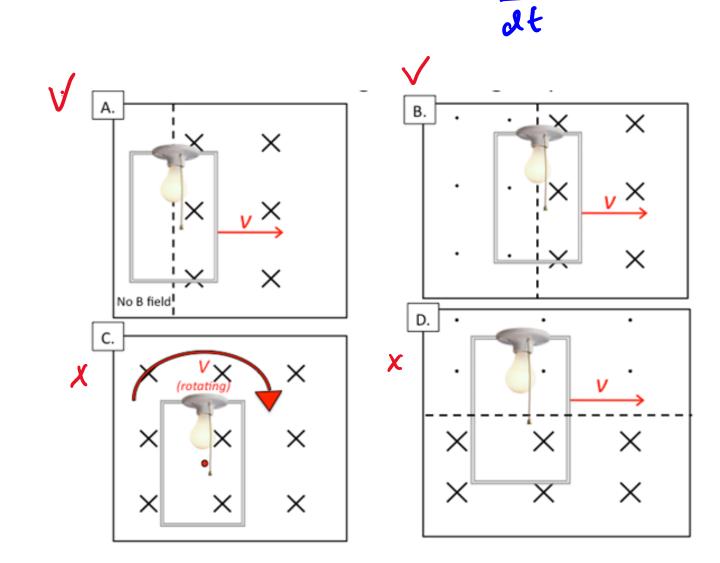
Q: Which of the following could light up the bulb?

g = mar

- A. A only
- (B) A and B
- C. A, B and C
- D. A and C
- E. All of them

Q: Which of these will light the bulb most strongly?

- A. A
- B B
  - C. Some



#### **Direction of Induced EMF**

• So  $\varepsilon = \left| \frac{d\Phi_B}{dt} \right|$  is the magnitude of induced EMF. We need to know its direction, CW or CCW

- Induced current is a current, and hence it creates (induces) magnetic field!
- Their directions are related by the RHR.

• The direction of the induced magnetic field (and, therefore, the direction of the induced current!) is determined by the Lenz law:



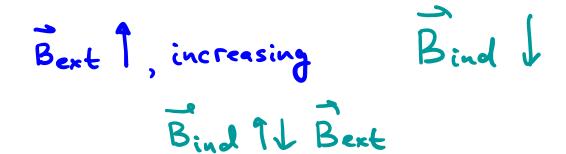
The direction of the induced current is such that the induced magnetic field opposes the change in the flux.

$$\varepsilon = -\frac{d\Phi_m}{dt}$$

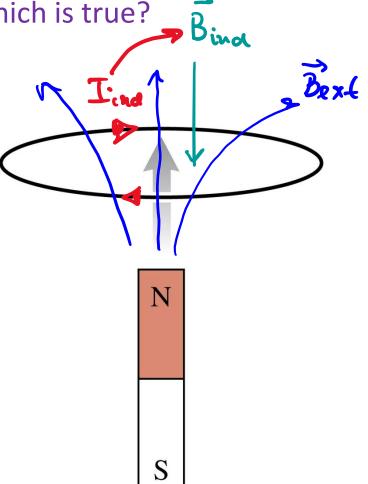
Q: The bar magnet is pushed toward the center of a wire loop. Which is true?

A. There is a clockwise induced current in the loop.

- B. There is a counterclockwise induced current in the loop.
- C. There is no induced current in the loop.
- D. Not sure



(assume we look at the loop from above)



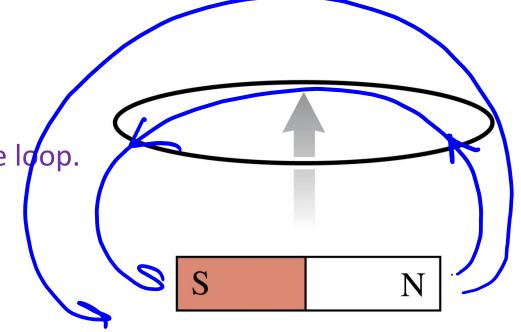
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