

Quantum mechanics

Particle in the box

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Problem Setup

Suppose the potential $V(x)$ is given by:

$$V(x) = \begin{cases} 0, & \text{if } 0 \leq x \leq L \\ \infty, & \text{otherwise} \end{cases}$$

Time-independent Schrödinger equation

The equation for the spatial part $\psi(x)$ is:

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + V\psi = E\psi$$

Inside the box, $V = 0$ (for $0 \leq x \leq L$):

$$\begin{aligned} -\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} &= E\psi \\ \frac{d^2\psi}{dx^2} &= -\frac{2mE}{\hbar^2}\psi \\ \frac{d^2\psi}{dx^2} + \frac{2mE}{\hbar^2}\psi &= 0 \end{aligned}$$

To simplify, we define a constant k^2 :

$$k^2 = \frac{2mE}{\hbar^2}$$

This transforms the equation into the second-order ODE:

$$\frac{d^2\psi}{dx^2} + k^2\psi = 0$$

The general solution to this equation is:

$$\psi(x) = A \sin(kx) + B \cos(kx)$$

Applying Boundary Conditions

The wave function $\psi(x)$ must be zero at the boundaries (where $V = \infty$).

At $x = 0$:

$$\begin{aligned} \psi(0) &= 0 \\ A \sin(0) + B \cos(0) &= 0 \\ A(0) + B(1) &= 0 \\ \therefore B &= 0 \end{aligned}$$

At $x = L$:

$$\begin{aligned}\psi(L) &= 0 \\ A \sin(kL) &= 0\end{aligned}$$

Since $A = 0$ would mean that there would be no particle, we must have:

$$kL = n\pi, \quad \text{where } n = 1, 2, 3, \dots$$

This quantizes the possible values of k :

$$k_n = \frac{n\pi}{L}$$

Find A (Normalization)

We normalize the wave function $\psi(x)$ so that the total probability of finding the particle in the box is 1.

$$\begin{aligned}\int_0^L |\psi(x)|^2 dx &= \int_0^L (A \sin(kx))^2 dx = 1 \\ A^2 \int_0^L \sin^2(kx) dx &= 1\end{aligned}$$

To solve the integral, use the identity:

$$\cos(2\theta) = 1 - 2 \sin^2(\theta) \implies \sin^2(\theta) = \frac{1 - \cos(2\theta)}{2}$$

So, $\sin^2(kx) = \frac{1 - \cos(2kx)}{2}$.

$$\begin{aligned}A^2 \int_0^L \left(\frac{1 - \cos(2kx)}{2} \right) dx &= 1 \\ \frac{A^2}{2} \left[\int_0^L 1 dx - \int_0^L \cos(2kx) dx \right] &= 1 \\ \frac{A^2}{2} \left[[x]_0^L - \left[\frac{1}{2k} \sin(2kx) \right]_0^L \right] &= 1 \\ \frac{A^2}{2} \left[(L - 0) - \left(\frac{1}{2k} \sin(2kL) - \frac{1}{2k} \sin(0) \right) \right] &= 1 \\ \frac{A^2}{2} \left[L - \frac{1}{2k} \sin(2kL) \right] &= 1\end{aligned}$$

Substitute our condition $kL = n\pi$:

$$\begin{aligned}\frac{A^2}{2} \left[L - \frac{1}{2k} \sin(2n\pi) \right] &= 1 \\ \frac{A^2}{2} [L - 0] &= 1 \\ \frac{A^2 L}{2} &= 1 \\ A^2 &= \frac{2}{L} \\ \therefore A &= \sqrt{\frac{2}{L}}\end{aligned}$$

Energy eigenstates

From our definitions $k^2 = \frac{2mE}{\hbar^2}$ and $k_n = \frac{n\pi}{L}$:

$$\left(\frac{n\pi}{L} \right)^2 = \frac{2mE_n}{\hbar^2}$$

Solving for energy E_n gives the quantized energy levels:

$$E_n = \frac{n^2 \pi^2 \hbar^2}{2mL^2}$$

Stationary State Wave function

These are the normalized energy eigenfunctions (or stationary states):

$$\psi_n(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right)$$

General Wave function (Linear Combination)

The most general solution for the *initial state* $\Psi(x, t = 0)$ is a superposition:

$$\Psi(x, 0) = \sum_{n=1}^{\infty} C_n \psi_n(x) = \sum_{n=1}^{\infty} C_n \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right)$$

Time-Dependent Schrödinger Equation

We use the separation of variables: $\Psi(x, t) = \psi(x)\varphi(t)$. The time-dependent part of the Schrödinger equation is the following.

$$\begin{aligned} i\hbar \frac{d\varphi}{dt} &= E\varphi \\ \frac{d\varphi}{\varphi} &= -\frac{iE}{\hbar} dt \\ \int \frac{d\varphi}{\varphi} &= \int -\frac{iE}{\hbar} dt \\ \ln \varphi &= -\frac{iEt}{\hbar} \\ \varphi(t) &= e^{-\frac{iEt}{\hbar}} \end{aligned}$$

We combine the parts $\Psi(x, t) = \sum C_n \psi_n(x)\varphi_n(t)$, using the specific energy E_n for each state ψ_n :

$$\Psi(x, t) = \sum_{n=1}^{\infty} C_n \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right) e^{-\frac{iE_n t}{\hbar}}$$

Substituting the formula for E_n :

$$\Psi(x, t) = \sum_{n=1}^{\infty} C_n \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right) e^{-\frac{i}{\hbar} \left(\frac{n^2 \pi^2 \hbar^2}{2mL^2}\right) t}$$