

# Quantum Tunneling through a Single Potential Barrier: Equivalence of the Transfer Matrix and Analytical Solutions

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11 November 2025

## Problem Setup

Suppose the potential  $V(x)$  is given by:

$$V(x) = \begin{cases} V_0, & 0 \leq x \leq L \\ 0, & \text{otherwise} \end{cases}$$

We analyze the case for a particle with energy  $E > 0$  incident from the left. This paper focuses on the quantum tunneling, so we assume the particle's energy is less than the barrier height:  $0 < E < V_0$ .

The Time-independent Schrödinger equation is:

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + V\psi = E\psi$$

## Wave Function Solutions in Each Region

We define  $k_1^2 = \frac{2mE}{\hbar^2}$  and  $k_2^2 = \frac{2m(V_0-E)}{\hbar^2}$ .

- **Region I** ( $x < 0, V = 0$ ):

$$\psi_I(x) = Ae^{ik_1x} + Be^{-ik_1x}$$

- **Region II** ( $0 \leq x \leq L, V = V_0$ ):

$$\psi_{II}(x) = Ce^{k_2x} + De^{-k_2x}$$

- **Region III** ( $x > L, V = 0$ ):

$$\psi_{III}(x) = Fe^{ik_1x}$$

## Boundary Conditions

We apply continuity for  $\psi(x)$  and  $\frac{d\psi}{dx}$  at both  $x = 0$  and  $x = L$ . This gives four conditions:

$$\psi_I(0) = \psi_{II}(0) \implies A + B = C + D \quad (1)$$

$$\psi_{II}(L) = \psi_{III}(L) \implies Ce^{k_2 L} + De^{-k_2 L} = Fe^{ik_1 L} \quad (2)$$

$$\left. \frac{d\psi_I}{dx} \right|_{x=0} = \left. \frac{d\psi_{II}}{dx} \right|_{x=0} \implies ik_1(A - B) = k_2(C - D) \quad (3)$$

$$\left. \frac{d\psi_{II}}{dx} \right|_{x=L} = \left. \frac{d\psi_{III}}{dx} \right|_{x=L} \implies k_2(Ce^{k_2 L} - De^{-k_2 L}) = ik_1 F e^{ik_1 L} \quad (4)$$

## Method 1: Analytical Approach (Solving the System)

Solving the system of four equations leads to expressions for the coefficients. By solving (2) and (4) for  $C$  and  $D$ :

$$C = \left(1 + \frac{ik_1}{k_2}\right) \frac{F}{2} e^{ik_1 L - k_2 L} \quad (5)$$

$$D = \left(1 - \frac{ik_1}{k_2}\right) \frac{F}{2} e^{ik_1 L + k_2 L} \quad (6)$$

And by solving (1) and (3) for  $A$  and  $B$ :

$$A = \left(1 - \frac{ik_2}{k_1}\right) \frac{C}{2} + \left(1 + \frac{ik_2}{k_1}\right) \frac{D}{2} \quad (7)$$

$$B = \left(1 + \frac{ik_2}{k_1}\right) \frac{C}{2} + \left(1 - \frac{ik_2}{k_1}\right) \frac{D}{2} \quad (8)$$

Substituting (5) and (6) into (7) and solving for the transmission amplitude  $F/A$ , we find the transmission probability  $T = \left|\frac{F}{A}\right|^2$ . This yields the well-known analytical result:

$$T = \frac{1}{\cosh^2(k_2 L) + \frac{\gamma^2}{4} \sinh^2(k_2 L)}$$

where  $\gamma = \frac{k_1^2 - k_2^2}{k_1 k_2}$ .

## Method 2: Transfer Matrix Approach

We can rewrite the equations from the boundary conditions in matrix form. The transfer matrix operator  $\hat{O}$  relates the coefficients in Region I to the final coefficients in Region III

$$\begin{pmatrix} A \\ B \end{pmatrix} = \hat{O} \begin{pmatrix} F \\ 0 \end{pmatrix}$$

The operator  $\hat{O}$  is a product of three operators:  $\hat{O} = \hat{M}_\alpha \cdot \hat{M}_\beta \cdot \hat{M}_\eta$ .

- **Operator  $\hat{M}_\alpha$  (Interface at  $x = 0$ ):** From equations (7) and (8), we relate (A, B) to (C, D):

$$\begin{pmatrix} A \\ B \end{pmatrix} = \hat{M}_\alpha \begin{pmatrix} C \\ D \end{pmatrix} \quad \text{where} \quad \hat{M}_\alpha = \frac{1}{2} \begin{pmatrix} 1 - \frac{ik_2}{k_1} & 1 + \frac{ik_2}{k_1} \\ 1 + \frac{ik_2}{k_1} & 1 - \frac{ik_2}{k_1} \end{pmatrix}$$

- **Propagation Operator  $\hat{M}_\beta$ :** This operator propagates the wave function through Region II, from  $x = 0$  to  $x = L$ . Let  $C' = Ce^{k_2 L}$  and  $D' = De^{-k_2 L}$  be the amplitudes at  $x = L$ . We relate  $\begin{pmatrix} C \\ D \end{pmatrix}$  to  $\begin{pmatrix} C' \\ D' \end{pmatrix}$ :

$$\begin{pmatrix} C \\ D \end{pmatrix} = \hat{M}_\beta \begin{pmatrix} C' \\ D' \end{pmatrix} \quad \text{where} \quad \hat{M}_\beta = \begin{pmatrix} e^{-k_2 L} & 0 \\ 0 & e^{k_2 L} \end{pmatrix}$$

- **Operator  $\hat{M}_\eta$  (Interface at  $x = L$ ):** This operator relates the amplitudes in Region II at  $x = L$  ( $C', D'$ ) to the amplitudes in Region III at  $x = L$ .  $\psi_{III}(x) = Fe^{ik_1 x}$ . At  $x = L$ , the amplitudes are  $Fe^{ik_1 L}$ .

$$\begin{pmatrix} C' \\ D' \end{pmatrix} = \hat{M}_\eta \begin{pmatrix} Fe^{ik_1 L} \\ 0 \end{pmatrix}$$

From solving equations (2) and (4), this boundary operator is:

$$\hat{M}_\eta = \frac{1}{2} \begin{pmatrix} 1 + \frac{ik_1}{k_2} & 1 - \frac{ik_1}{k_2} \\ 1 - \frac{ik_1}{k_2} & 1 + \frac{ik_1}{k_2} \end{pmatrix}$$

The total relationship is:

$$\begin{pmatrix} A \\ B \end{pmatrix} = (\hat{M}_\alpha \hat{M}_\beta \hat{M}_\eta) \begin{pmatrix} Fe^{ik_1 L} \\ 0 \end{pmatrix}$$

Let  $\hat{O} = \hat{M}_\alpha \hat{M}_\beta \hat{M}_\eta$ . We only need the  $O_{11}$  element:

$$A = O_{11}(Fe^{ik_1 L})$$

## Operator $\hat{O}$ Elements

$$O_{11} = (M_{\alpha,11})(M_{\beta,11})(M_{\eta,11}) + (M_{\alpha,12})(M_{\beta,22})(M_{\eta,21})$$

$$\begin{aligned} O_{11} &= \left[ \frac{1}{2} \left( 1 - \frac{ik_2}{k_1} \right) \right] [e^{-k_2 L}] \left[ \frac{1}{2} \left( 1 + \frac{ik_1}{k_2} \right) \right] \\ &\quad + \left[ \frac{1}{2} \left( 1 + \frac{ik_2}{k_1} \right) \right] [e^{k_2 L}] \left[ \frac{1}{2} \left( 1 - \frac{ik_1}{k_2} \right) \right] \end{aligned}$$

$$O_{11} = \frac{1}{4} \left[ \left( 1 - \frac{ik_2}{k_1} \right) \left( 1 + \frac{ik_1}{k_2} \right) e^{-k_2 L} + \left( 1 + \frac{ik_2}{k_1} \right) \left( 1 - \frac{ik_1}{k_2} \right) e^{k_2 L} \right]$$

where  $\gamma = \frac{k_1^2 - k_2^2}{k_1 k_2}$ :

- $\left( 1 - \frac{ik_2}{k_1} \right) \left( 1 + \frac{ik_1}{k_2} \right) = 2 + i\gamma$
- $\left( 1 + \frac{ik_2}{k_1} \right) \left( 1 - \frac{ik_1}{k_2} \right) = 2 - i\gamma$

We get:

$$O_{11} = \frac{1}{4} [(2 + i\gamma)e^{-k_2 L} + (2 - i\gamma)e^{k_2 L}]$$

$$\begin{aligned} O_{11} &= \frac{1}{4} [2(e^{k_2 L} + e^{-k_2 L}) - i\gamma(e^{k_2 L} - e^{-k_2 L})] \\ &= \frac{1}{4} [4 \cosh(k_2 L) - 2i\gamma \sinh(k_2 L)] \\ &= \cosh(k_2 L) - \frac{i\gamma}{2} \sinh(k_2 L) \end{aligned}$$

## Conclusion: Equivalence

From  $A = O_{11}(F e^{ik_1 L})$ , the transmission amplitude is

$$\frac{F}{A} = \frac{1}{O_{11} e^{ik_1 L}} = \frac{e^{-ik_1 L}}{O_{11}}$$

$$\frac{F}{A} = \frac{e^{-ik_1 L}}{\cosh(k_2 L) - \frac{i\gamma}{2} \sinh(k_2 L)}$$

This exactly matches the analytical result derived earlier. Calculating the transmission probability  $T = |\frac{F}{A}|^2$ :

$$\begin{aligned} T &= \left| \frac{e^{-ik_1 L}}{\cosh(k_2 L) - \frac{i\gamma}{2} \sinh(k_2 L)} \right|^2 = \frac{|e^{-ik_1 L}|^2}{|\cosh(k_2 L) - \frac{i\gamma}{2} \sinh(k_2 L)|^2} \\ T &= \frac{1}{\cosh^2(k_2 L) + \frac{\gamma^2}{4} \sinh^2(k_2 L)} \end{aligned}$$

This demonstrates that the standard Transfer Matrix method and the standard Analytical Solution are equivalent, yielding the identical result.