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HOW TO TRAIN YOUR ELECTRON

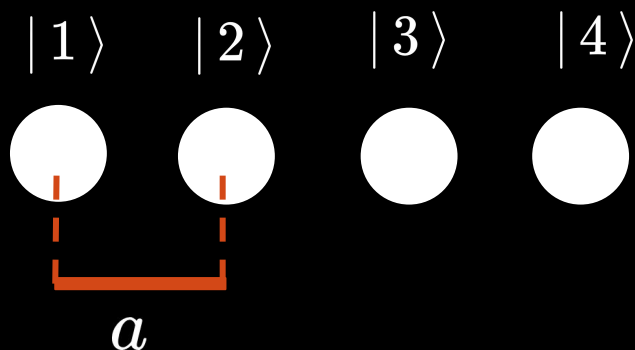
Anderson localization on the lattice & random graphs

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The Tight Binding Model

- **Goal:** understand how an electron behaves in a material
- For simplicity, think of an electron on an atomic ring:
 - The n^{th} atom has a single orbital, labelled $|n\rangle$
 - We represent the wavefunction in the site basis
 - The electron can hop between nearby sites



The Tight Binding Hamiltonian

- The tight binding hamiltonian:

$$\langle n|H|m\rangle = H_{n,m} = \underbrace{\epsilon_0\delta_{n,m}}_{\text{binding}} - \underbrace{t(\delta_{n+1,m} + \delta_{n-1,m})}_{\text{hopping}}$$

- where ϵ_0 refers to the energy at a site and t is the hopping parameter
- Hamiltonian \sim Time Evolution
- This is typically diffusive!!

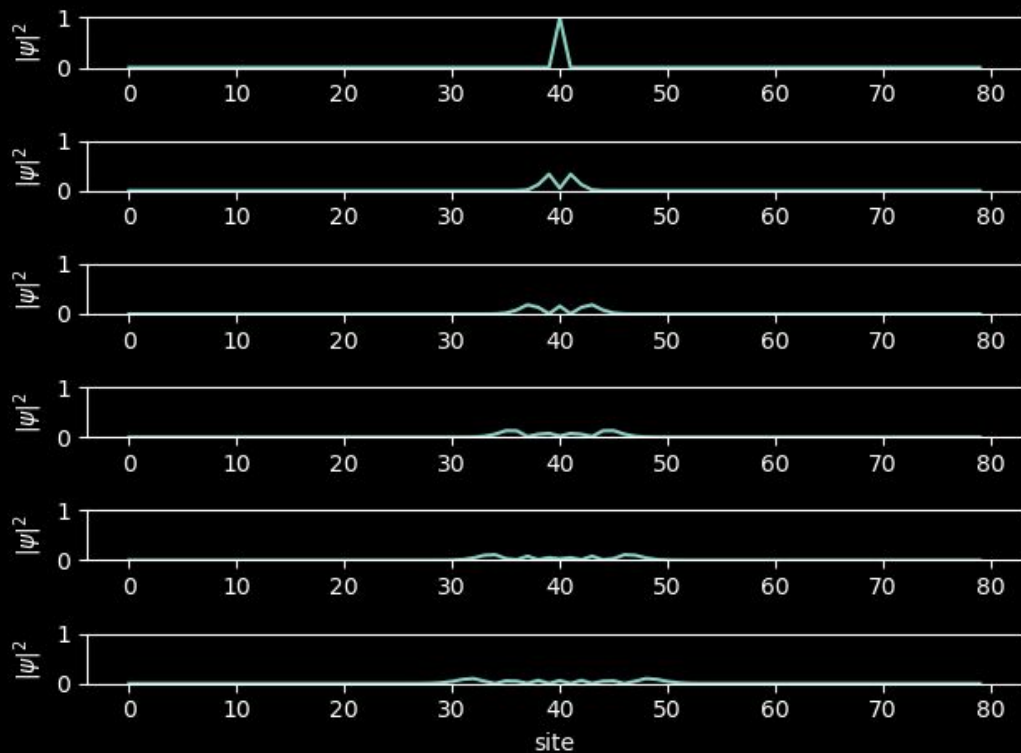
Numerical Approach

- Given a hamiltonian, we can time evolve the wave function:

$$|\psi(t)\rangle = \exp(-iHt/\hbar)|\psi(0)\rangle$$

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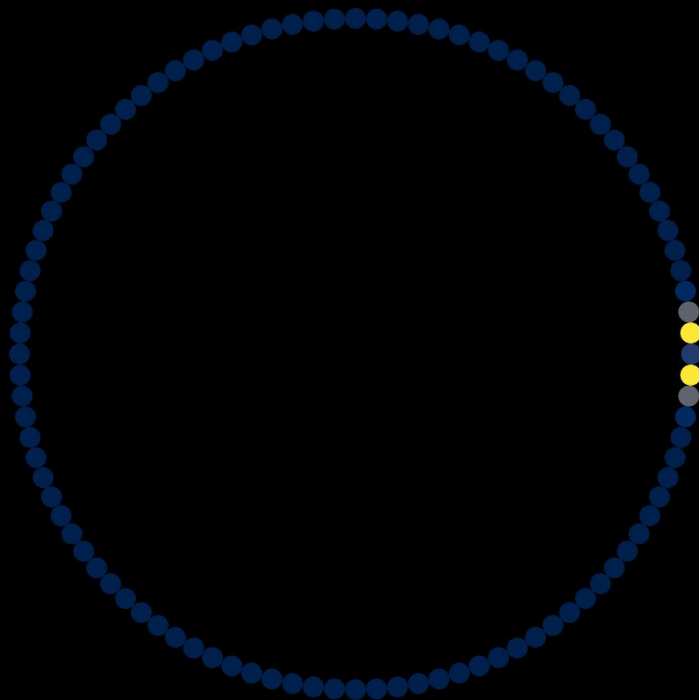
Anderson Localization

- Counterintuitively, we can get localization by modifying tight binding.

$$\langle n|H|m\rangle = H_{n,m} = \underbrace{\epsilon_0\delta_{n,m}}_{\text{binding}} - \underbrace{t(\delta_{n+1,m} + \delta_{n-1,m})}_{\text{hopping}}$$

- Key idea: randomly sample values of ϵ from a uniform distribution from $[-W, W]$
 - Randomizing the diagonal of H
 - Represents a disordered material
 - Larger W = more localized

Anderson Localization on the Ring



$W = 0$



*One ring to rule them all,
one ring to find them,
One ring to bring them all
and in the darkness bind them.*

$W = 10$

Anderson Localization on a Graph

- We can identify the hopping term with the adjacency matrix of a graph

$$\langle n|H|m\rangle = H_{n,m} = \underbrace{\epsilon_0\delta_{n,m}}_{\text{binding}} - \underbrace{t(\delta_{n+1,m} + \delta_{n-1,m})}_{\text{hopping}}$$

- This allows us to extend to more interesting configurations*

*we assume periodic boundary conditions

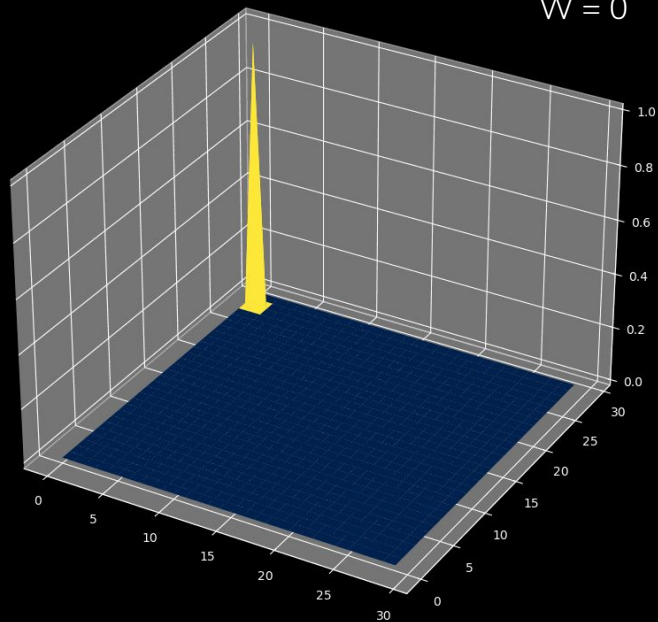
Anderson Localization on a Periodic Lattice

- Consider a 2D lattice with periodic boundary conditions.
- For low W , the wave function diffuses. For high W , we see localization

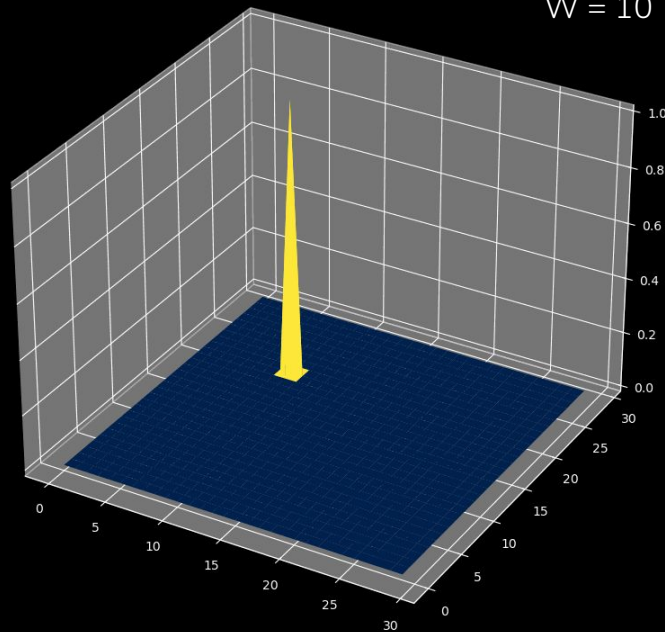
Wave function probability density at time 0

Wave function probability density at time 0

$W = 0$



$W = 10$

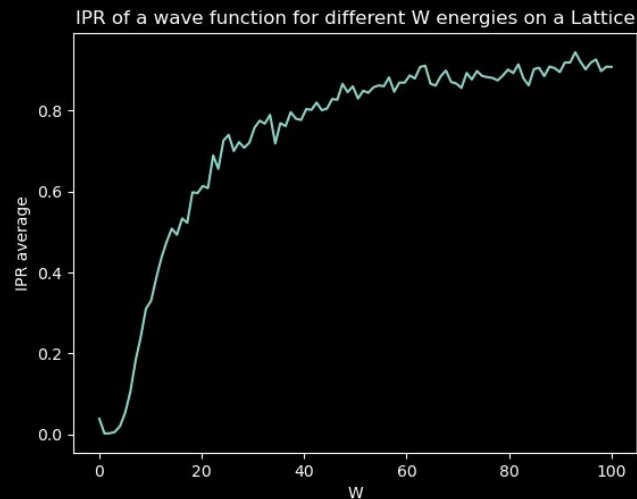
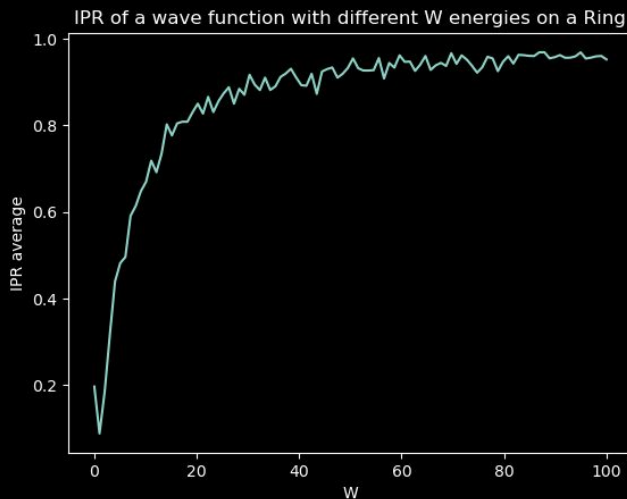


Inverse Participation Ratio (IPR)

- IPR is a rough measure of localization. IPR of 1 means fully localized.

$$IPR = \sum_n |\psi(n)|^4$$

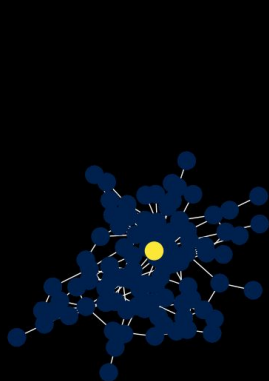
- As W increases, the localization also tends to increase.



Random Graphs

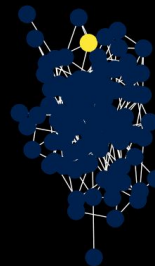
- We generalize even more by simulating the behavior of a wave function on Erdős–Rényi graphs (where you're given a node set and each possible edge is generated with probability p)

Wave function probability density at time 0



$W = 10, p = 0.03$

Wave function probability density at time 0

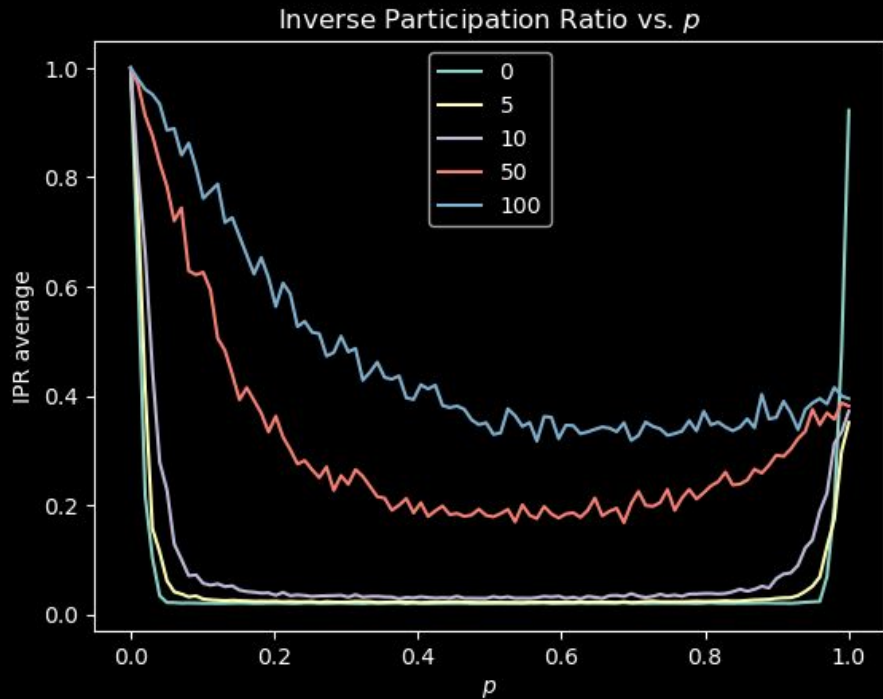
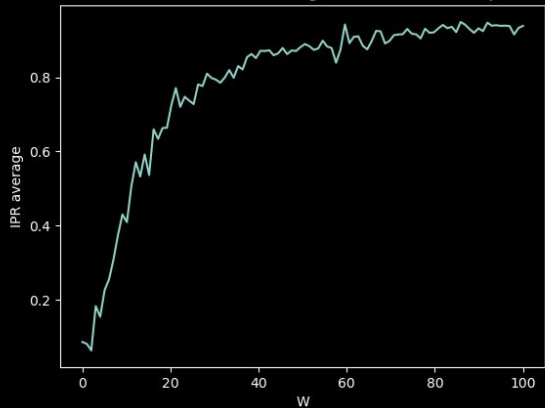


$W = 10, p = 0.05$

IPR vs. p for different W values

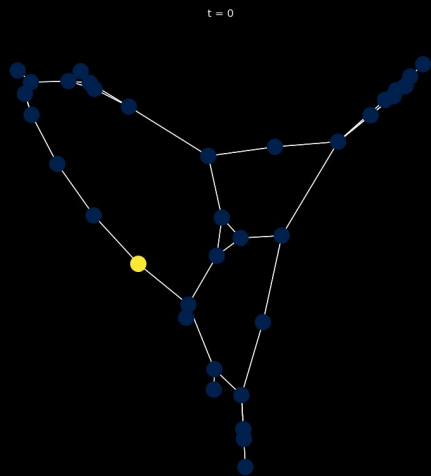
- It's harder to see localization for random graphs, so we lean on IPR more heavily.
- We're interested in how localization changes as the graph becomes more connected (i.e. p increases).

IPR of a wave function with different W energies on a Random Graph with Probability p

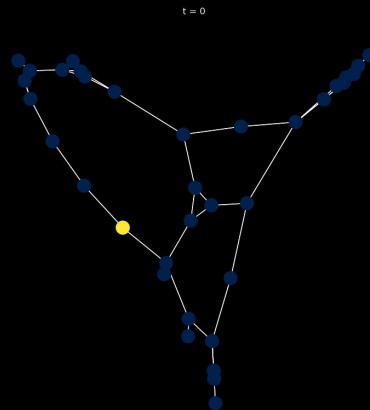


Sparse Graphs

- In a restricted domain, do structural (e.g., symmetry) properties of the graph influence localization parameters?



$$W = 0$$



$$W = 18.7$$

Future Directions

- Correlating & extending analytical predictions made for regular random graphs to specific properties of adjacency matrices.
- Why does the IPR vs. W graph dip in the beginning instead of just increasing?
- Why does the IPR vs. p graph increase at the end?
- Investigate time-dependence of IPR, e.g. test whether noise in IPR vs. W graph disappears over IPR averaged over multiple time points
- Investigate different boundary conditions and connections between nodes

Questions?

- Thank you Dr. Gilpin, Anthony, and Edoardo for the fun semester!