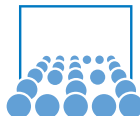


CFD Lab

The Lattice-Boltzmann Method

Nikola Tchipev, Friedrich Menhorn

19.05.2017



Outline

Intro

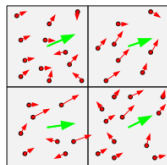
Molecular Dynamics

From MD to LBM

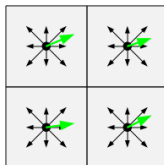
Lattice-Boltzmann Method

LBM - a different story

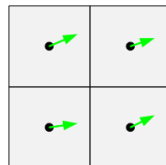
- Macroscale:
 - Finite Difference Methods (FDM)
- Mesoscale:
 - Lattice-Boltzmann Method (LBM)
- Microscale:
 - Molecular Dynamics (MD)



MD
microscale



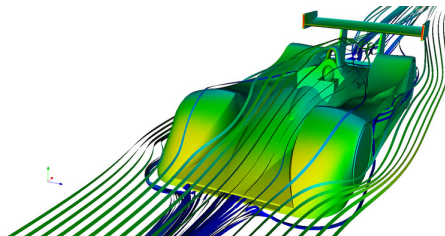
LBM
mesoscale



FDM
macroscale

LBM

- fluid solver, but we don't solve NSE
- based on statistical mechanics
- new (and still evolving) method
- easy to program
- already a factor in the automotive industry



Some assumptions

What assumptions did we have in Worksheet 1?

Some assumptions

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- incompressible, isothermal, Newtonian, ...

Some assumptions

What assumptions did we have in Worksheet 1?

- incompressible, isothermal, Newtonian, ...
- continuum assumption

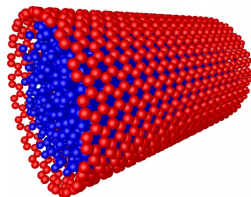
Some assumptions

What assumptions did we have in Worksheet 1?

- incompressible, isothermal, Newtonian, ...
- continuum assumption

Can we solve *any* flow problem with NSE?

Flow in a carbon nanotube?



The continuum assumption

Fluids in reality

composed of atoms and molecules, empty space in between

Fluids under the continuum assumption

composed of **continuous matter, filling the entire space**

When is the continuum assumption valid?

The continuum assumption

Continuum assumption is valid for

$$Kn \ll 1. \quad (1)$$

Kn : Knudsen number

$$Kn = \frac{\lambda}{L_c}, \quad (2)$$

L_c : characteristic length

λ : mean free path

- air at STP: $\lambda \approx \mathcal{O}(nm)$

Thought experiment

Small particle in fluid at rest

- $\Rightarrow \vec{u} = 0$ identically
- L_c - diameter of particle.
- as L_c decreases, Kn increases
- as Kn approaches 1, the particle begins to feel collisions with individual molecules
- Brownian motion kicks in!

But NSE (or the Stokes Eq.) predict no motion of the particle!

Outline

Intro

Molecular Dynamics

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Molecular Dynamics

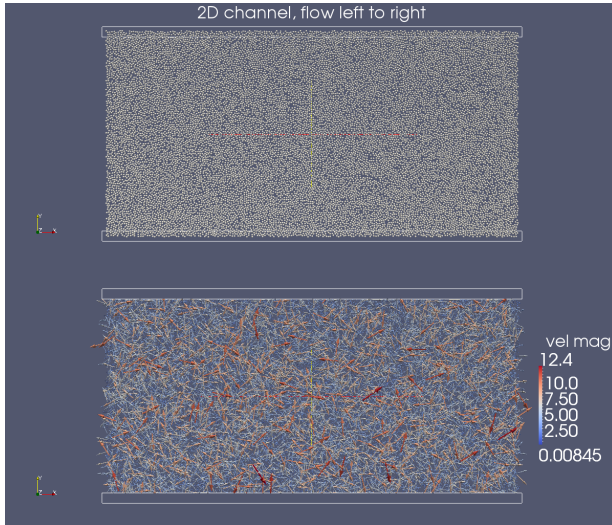
Nano-, Micro- things:

- Nano-, Microfluidics

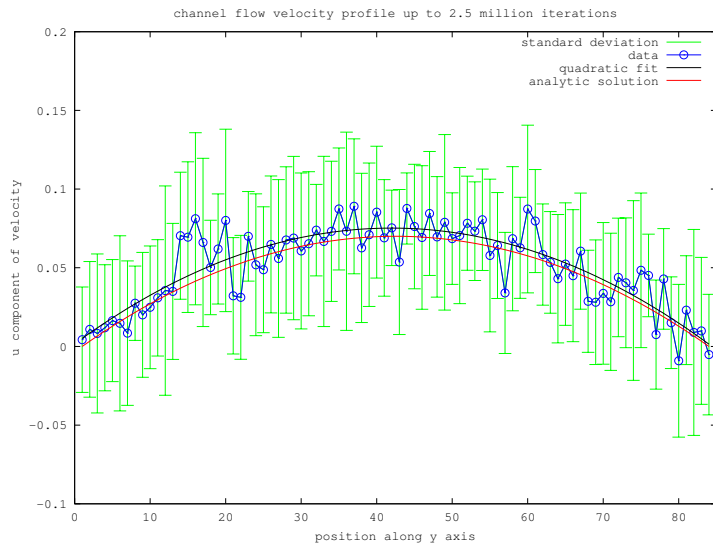
Applications

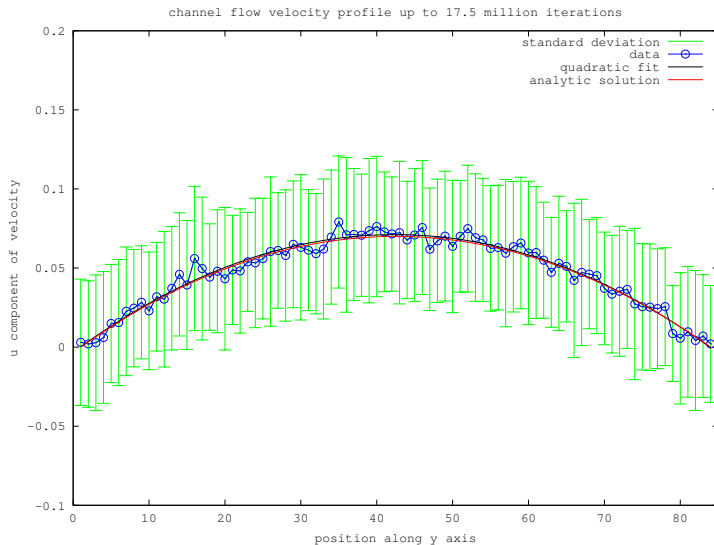
- nanotubes, -pores, -filters...
- Lab-on-a-chip

(Videos - flow in a channel)



mean speed: 0.17, max speed: 12.4





Can we solve any flow problem with MD?

theoretically yes. Practically:

- Largest MD simulation: 4×10^{12} particles.
- 1 milliliter of water: 3×10^{22} particles..
- timesteps in MD: $\mathcal{O}(10^{-15}\text{s})$...

very compute intensive

Statistical noise

Need to sample:

- in space
- in time
- ...

Outline

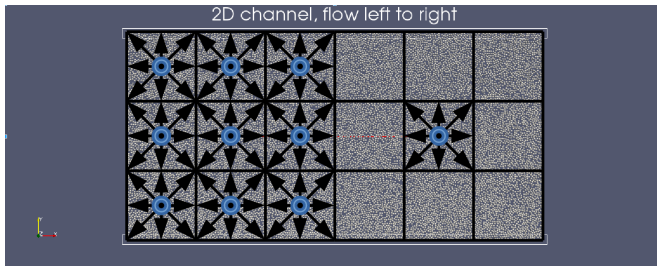
Intro

Molecular Dynamics

From MD to LBM

Lattice-Boltzmann Method

From MD to LBM



1. introduce cells
2. what is the probability that between two timesteps
 - a particle travels (**streams**) from cell i to cell j ?
 - particles **collide**?

Discretize

- particle position
- particle velocity
- time

From MD to LBM

Lattice Gas Cellular Automata

- historically - Lattice Gas Cellular Automata as intermediate step
- theoretically - from Boltzmann Equation (Thermodynamics)

Outline

Intro

Molecular Dynamics

From MD to LBM

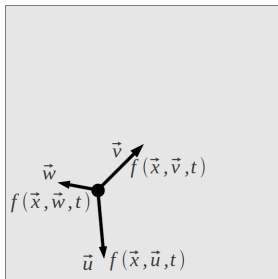
Lattice-Boltzmann Method

Replace boolean n_i with real f_i .

Algorithm

```
while ( $t \neq t_{end}$ ) {  
  1. collide - handle  $f_i$ 's at the same site  
  2. stream - travel the respective edge  
  3.  $t = t + \Delta t$   
}
```

The f_i



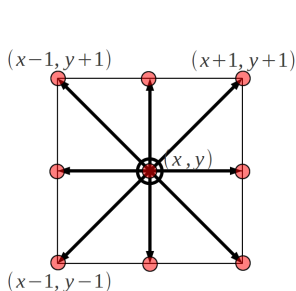
$f(\vec{x}, \vec{v}, t)$: probability density function for finding particles with velocity \vec{v} at (\vec{x}, t)

$$f \in \mathbb{R}, f \in (0, 1)$$

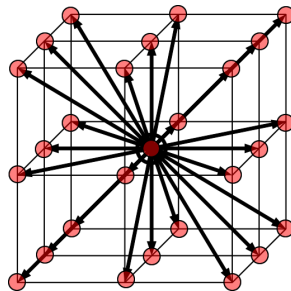
The LBM lattices

DnQm notation:

- n: number of dimensions
- m: number of directions



D2Q9



D3Q27

other possibilities: D2Q5, D2Q7, D3Q15, D3Q19

Mesoscopic to macroscopic quantities

Mesoscopic \rightarrow macroscopic

- (where do we need this conversion?)

Mesoscopic \leftarrow macroscopic

- (where do we need this conversion?)

Mesoscopic \rightarrow macroscopic

Given $\{f_i\}$, compute $\{\rho, u, p\}$:

- density:

$$\rho(\vec{x}, t) = \sum_{i=0}^{Q-1} f_i \approx 1$$

- momentum:

$$\vec{u}(\vec{x}, t) \rho(\vec{x}, t) = \sum_{i=0}^{Q-1} f_i \cdot \vec{c}_i$$

- pressure:

$$p = \rho \cdot c_s^2$$

c_i : velocity associated to f_i

(e.g. for D2Q9, $c_i = (\alpha_i, \beta_i)$, $\alpha_i, \beta_i \in \{-1, 0, 1\}$)

$c_s = \frac{1}{\sqrt{3}}$: speed of sound

local operations

Mesoscopic \leftarrow macroscopic

Given $\{\rho, u\}$, compute $\{f_i\}$:

The equilibrium distribution function $f_i^{eq}(\rho, u)$:
a specific mapping from $\{\rho, u\}$ to $\{f_i\}$:

$$f_i^{eq} = w_i \rho \left(1 + \frac{\vec{c}_i \cdot \vec{u}}{c_s^2} + \frac{1}{2c_s^4} (\vec{c}_i \cdot \vec{u})^2 - \frac{1}{2c_s^2} \vec{u} \cdot \vec{u} \right),$$

w_i - weights, depending on the chosen lattice.

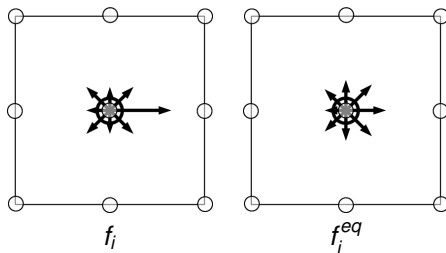
E.g. D2Q9:

$$w_i = \begin{cases} \frac{4}{9} & \text{if } \|c_i\| = 0 \\ \frac{1}{9} & \text{if } \|c_i\| = 1 \\ \frac{1}{36} & \text{if } \|c_i\| = \sqrt{2} \end{cases}$$

- needed for IC, BC and collisions

The equilibrium distribution

$f^{eq} \approx$ Normal distribution



The equilibrium distribution

f^{eq} - Maxwell Boltzmann equilibrium distribution

$$f^{eq}(\vec{x}, \vec{c}_i, t) = \left(\frac{m}{2\pi k_B T} \right)^{D/2} \cdot \exp \left(- \frac{m(\vec{c}_i - \vec{u}(\vec{x}, t))^2}{2k_B T} \right),$$

it **assumes** an ideal gas:

$$pV = Nk_b T$$

$$p = \frac{N}{V} \times const = \rho \times const$$

Moreover, the form

$$f_i^{eq} = w_i \rho \left(1 + \frac{\vec{c}_i \cdot \vec{u}}{c_s^2} + \frac{1}{2c_s^4} (\vec{c}_i \cdot \vec{u})^2 - \frac{1}{2c_s^2} \vec{u} \cdot \vec{u} \right),$$

assumes $u \ll c_i$.

So, we assume a **low Mach number**:

$$Ma = \frac{u}{c_s} \ll 1$$

c_s - speed of sound, information transfer

$$c_s = \text{const} \times \sqrt{k_b T}$$

weakly compressible flow

Collision

BGK approximation

Bhatnagar-Gross-Krook approximation

$$\Omega_i = -\frac{1}{\tau}(f_i - f_i^{eq}(\rho, u))$$

$\tau \in (0.5, 2.0)$ - relaxation “time”

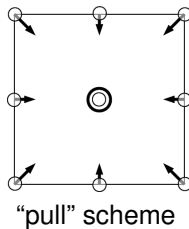
$$\tau = \frac{\nu}{c_s^2} + 0.5$$

ν - kinematic viscosity

we model “weak departure from equilibrium state of an ideal gas”

Collision II

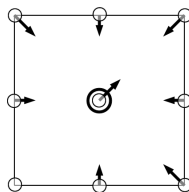
Why collision is local



Distributions "pass through" each other.

Collision II

Why collision is local

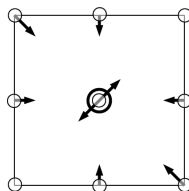


“pull” scheme

Distributions “pass through” each other.

Collision II

Why collision is local

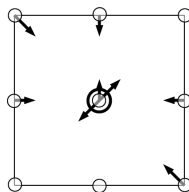


“pull” scheme

Distributions “pass through” each other.

Collision II

Why collision is local

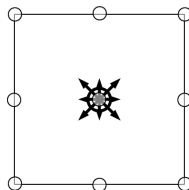


“pull” scheme

Distributions “pass through” each other.

Collision II

Why collision is local



“pull” scheme

Distributions “pass through” each other.

Collision III

Collision with intuition

$$\begin{aligned}\text{collision:} \quad f_i &:= f_i + \Omega_i \\ f_i &:= f_i - \frac{1}{\tau} (f_i - f_i^{eq})\end{aligned}$$

Let $\omega = \frac{1}{\tau} \in (0.5, 2.0)$:

$$f_i := (1 - \omega)f_i + \omega f_i^{eq},$$

Collision III

Collision with intuition

$$\begin{aligned}\text{collision:} \quad f_i &:= f_i + \Omega_i \\ f_i &:= f_i - \frac{1}{\tau} (f_i - f_i^{eq})\end{aligned}$$

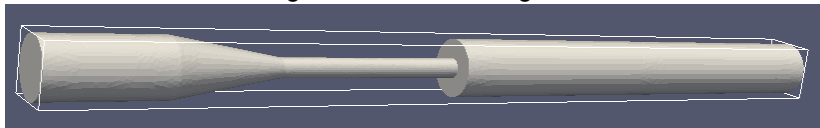
Let $\omega = \frac{1}{\tau} \in (0.5, 2.0)$:

$$f_i := (1 - \omega)f_i + \omega f_i^{eq},$$

SOR update rule?

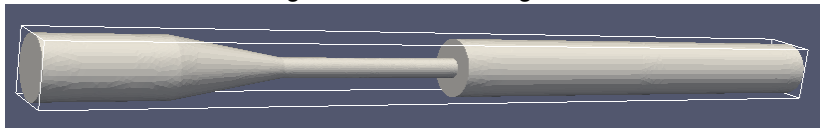
Collision III

Channel with a narrowing, flow from left to right:

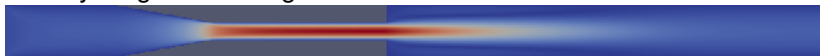


Collision III

Channel with a narrowing, flow from left to right:

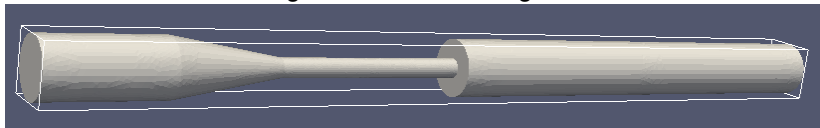


Velocity magnitude at high Re :

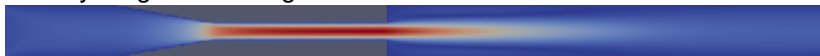


Collision III

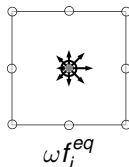
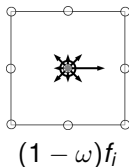
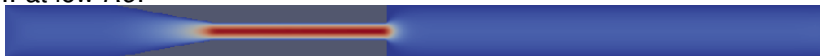
Channel with a narrowing, flow from left to right:



Velocity magnitude at high Re :



.. at low Re :



$$\omega = \omega(\nu) = \omega(Re)$$

Collision and streaming

Combined update rule

$$f_i(\vec{x} + c_i \Delta t, t + \Delta t) = f_i(\vec{x}, t) - \frac{1}{\tau} (f_i(\vec{x}, t) - f_i^{eq}(\rho(\vec{x}, t), u(\vec{x}, t)))$$

Implementation

collide:

$$f_i^*(\vec{x}, t) := f_i(\vec{x}, t) - \frac{1}{\tau} (f_i(\vec{x}, t) - f_i^{eq}(\rho(\vec{x}, t), u(\vec{x}, t)))$$

stream:

$$f_i(\vec{x} + c_i \Delta t, t + \Delta t) = f_i^*(\vec{x}, t)$$

NSE can be recovered!

multiscale Chapman Enskog analysis

The continuous case

We are solving the continuous Boltzmann equation

$$\frac{\partial f}{\partial t} + \vec{v} \cdot \nabla = \Omega$$

with our particular Lattice space- and velocity discretization.

- density:

$$\rho(\vec{x}, t) = \int_{\mathbb{R}^D} f(\vec{v}) d\vec{v}$$

- momentum:

$$\vec{u}(\vec{x}, t) \rho(\vec{x}, t) = \int_{\mathbb{R}^D} f(\vec{v}) \cdot \vec{v} d\vec{v}$$

Stability

We made many assumptions ...

- $\rho \approx 1$
- $u \ll 1$
- ν can't get arbitrarily small
 \Rightarrow how to control Re ?

Pros

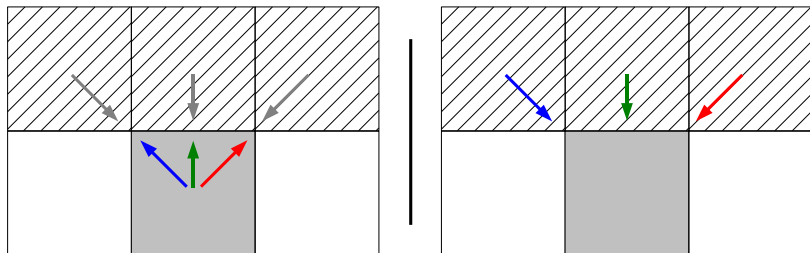
- can handle continuum problems (e.g. cars), but also higher Kn numbers
- coupling to MD
- closer to true physical description than NSE:
turbulence, diffusion, multi-component flows

Cons

- Homework: compare with NSE

Boundary Conditions

No-slip "bounce back"



Moving Wall

how to impose a specific velocity?

Questions?