

2024. 3. 19

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(D)

Transition prob. per unit time

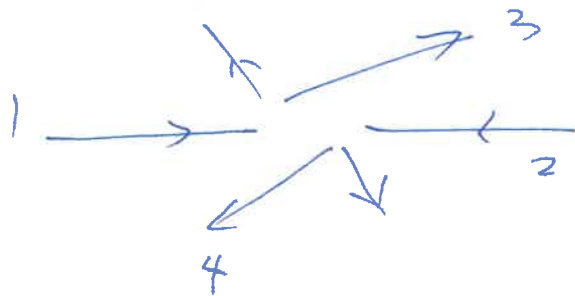
$$= \frac{2\pi}{\hbar} \cdot |M|^2 \cdot \text{phase space factor}$$

Decay rate Γ $1 \rightarrow 2 + 3 + \dots + N$

Differential decay rate

$$d\Gamma = \frac{S}{2k \cdot m_1} |M|^2 (2\pi)^4 \delta^{(4)}(P_1 - P_2 - \dots - P_N)$$

$$\frac{N}{\prod_{j=2}^N} \frac{d^4 P_j}{(2\pi)^4} (2\pi) \delta(P_j^2 - m_j^2 c^2) \Theta(P_j^0)$$

Differential cross section $1 + 2 \rightarrow 3 + 4 + \dots + N$ 

$$d\sigma = \frac{S \hbar^2}{4 \cdot \sqrt{(P_1 \cdot P_2)^2 - (m_1 m_2 c^2)^2}} |M|^2 (2\pi)^4 \delta^{(4)}(P_1 + P_2 - P_3 - \dots - P_N)$$

$$\frac{N}{\prod_{j=3}^N} \frac{d^4 P_j}{(2\pi)^4} \Theta(P_j^0) (2\pi) \delta(P_j^2 - m_j^2 c^2) \quad \Theta(P_j^0) = \text{step function}$$

Chapter 6 Griffiths

Decays and cross section (scattering)

Experimentally, the spectroscopic investigation (spectral lines) provides information about bound states of the particles (e.g. hydrogen atom H as bound state of e^- and p). Another approach is scattering of the particles, observe decay of these particles.

scattering can reveal the nature of particle interaction
→ Formulate decay process and scattering mathematically.

A typical decay process: $1 \rightarrow 2 + 3 + 4 \dots + N$

Define decay rate = probability a particle decay
per unit time = Γ

The probability a particle will decay in time $\delta t = \Gamma \delta t$

If there are N particles at time t , then the number of
particles will decay in time $\delta t = N \cdot \Gamma \cdot \delta t$

$$\rightarrow \delta N = -N \Gamma \delta t$$

$$\frac{dN}{dt} = -\Gamma N$$

Solving

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$$\rightarrow N = N_0 e^{-\Gamma t}$$

N_0 = # of particles at time $t=0$.

Mean life time of a particle = τ

$$= \frac{\text{Sum of the lifetimes of all the decayed particles}}{\text{Sum of all the decayed particles}}$$

$$= \frac{1}{\Gamma} \quad (\text{as shown below})$$

Suppose at time t , we have N particles, and

at time $t + \delta t$, δN particles decay away,

that means life time of all δN particles

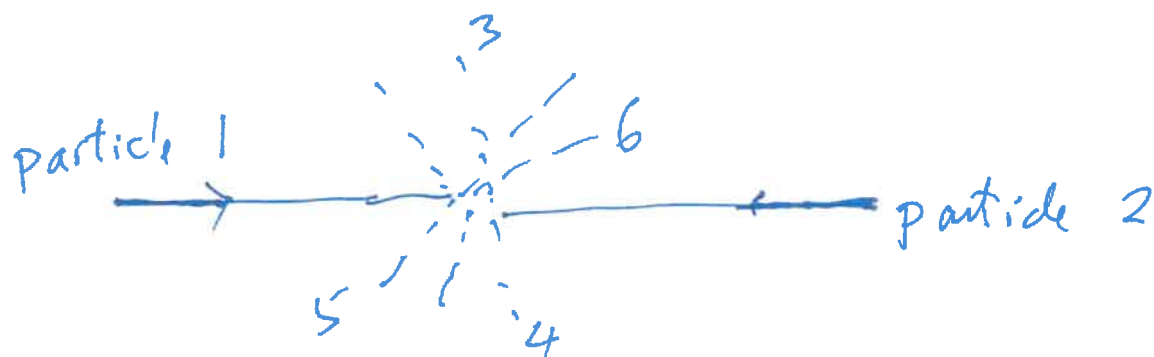
$$= t \cdot \delta N \quad (\text{each of the } \delta N \text{ particles has a lifetime } t)$$

$$\tau = \frac{\int_{N_0}^0 t dN}{\int_{N_0}^0 dN} = \frac{1}{\Gamma} \quad (H W)$$

$$\left(\text{Using } N = N_0 e^{-\Gamma t}, \quad H W \right)$$

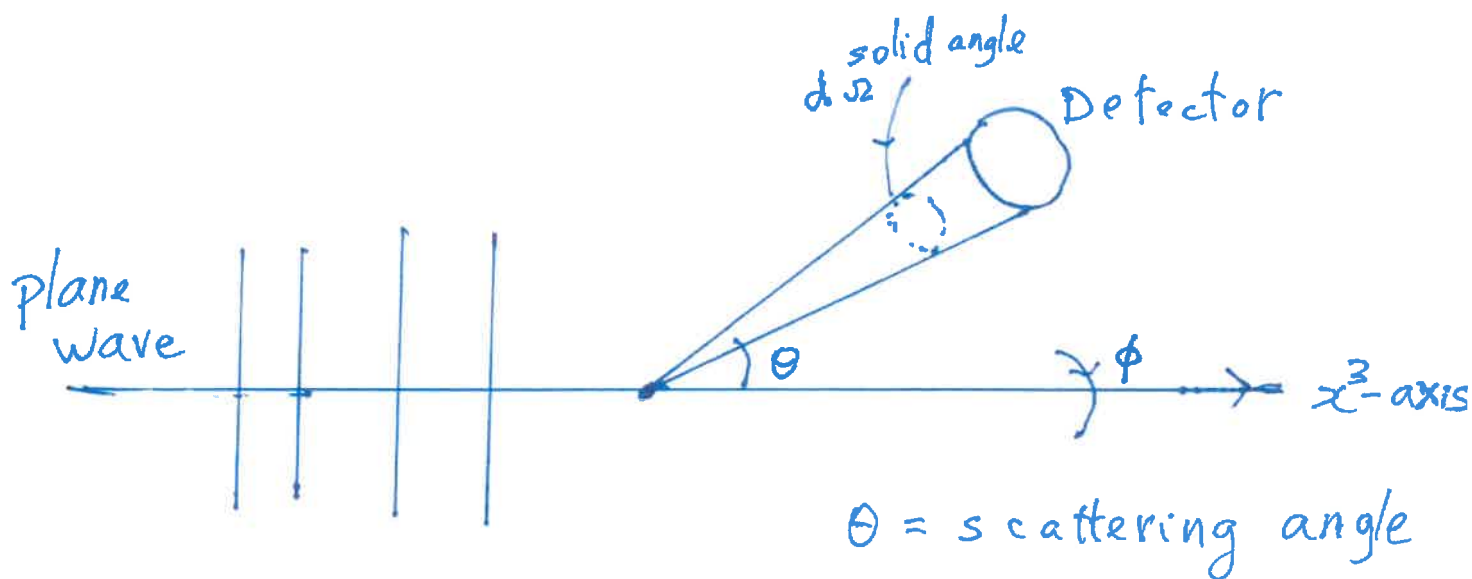
Scattering

$$1 + 2 \rightarrow 3 + 4 + 5 + \dots \quad (3)$$



lab frame

incident particle \rightarrow \times target



scattering amplitude

$$\begin{aligned} \text{Differential cross section} &= \frac{\text{scattered flux}}{\text{Incident flux}} \\ &= \frac{d\sigma}{d\Omega} \end{aligned}$$

Total cross section σ

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Incident flux = $\frac{\text{number}}{\text{per unit area per unit time}}$ of particles incident
 $= I_0$ (probability current density \vec{j})

Scattered flux = $\frac{\text{number}}{\text{per unit solid angle per unit time}}$ of particles scattered
 into the solid angle direction (θ, ϕ) per solid angle per unit time [the detector is at the angular position (θ, ϕ)]
 $= I(\theta, \phi)$

Differential cross section
 $\frac{d\sigma}{d\Omega} = \frac{I(\theta, \phi)}{I_0}$

Total cross section $\sigma = \int_{\text{all angles}} \frac{d\sigma}{d\Omega} d\Omega$

To day discuss how to compute Γ and $\frac{d\sigma}{d\Omega}$
 $\Gamma = \text{decay rate}$

To do that we need a formula from quantum mechanics, the Fermi golden rule. We quote:

→ Transition $\frac{\text{probability}}{\text{per unit time}} = \frac{2\pi}{\hbar} |M|^2 \cdot \text{phase space factor}$

Compute $|M|^2$ from dynamics, $M = \text{scattering amplitude}$
 phase space factor from kinematics.

M = scattering amplitude (matrix element)

can be obtained by solving equation of motion
or using Feynman diagrams with Feynman rules.

Phase space factor denotes the states available for
the finally produced particles to occupy

The larger the phase space factor, the more likely
the process will be.

Using this transition probability formula, one
can derive, the differential decay rate $d\Gamma$

For decay of a single particle $\underline{p}_1 = (m, c, 0)$

$$d\Gamma = \frac{S}{2\hbar m_1} |M|^2 \quad (\text{stationary})$$

$$(2\pi)^4 \delta^{(4)}(\underline{p}_1 - \underline{p}_2 - \underline{p}_3 \dots \underline{p}_N) = \text{product of 4 Dirac delta functions for overall conservation of 4-momentum}$$



$$\prod_{j=2}^N \left(\frac{d^4 \underline{p}_j}{(2\pi)^4} \theta(p_j^0) (2\pi) \delta(\underline{p}_j^2 - m_j^2 c^2) \right)$$

mass shell
condition
for each particle

S = statistical factor

$= \frac{1}{j!}$ if there are j identical particles
produced

Decay of a single particle (at rest, $\underline{P}_1 = (m, c, 0)$)

$$d\Gamma = \frac{S}{2\hbar m_1} |\mathcal{M}|^2 \cdot (2\pi)^4 \delta^{(4)}(\underline{P}_1 - \underline{P}_2 - \underline{P}_3 - \dots - \underline{P}_n) \cdot \prod_{j=2}^n \frac{1}{j!} \frac{(d^4 \underline{P}_j)}{(2\pi)^4} \theta(p_j^0) (2\pi) \delta(\underline{P}_j^2 - m_j^2 c^2)$$



S = statistical factor



$= \frac{1}{j!}$ if there are j identical particles produced

e.g. if there are $3 \pi^0$, $4 \pi^-$, $5 \pi^+$ in the final produced particles, then $S = \frac{1}{3! 4! 5!}$

$\theta(x)$ = step function, $\theta(x) = \begin{cases} 1 & \text{if } x > 0 \\ 0 & \text{if } x < 0 \end{cases}$

$\underline{P}_j^2 = m_j^2 c^2$ means particle j is in its mass shell. $(p_j^{02} - \underline{p}_j^2 = m_j^2 c^2)$

$$\int d p_j^0 \theta(p_j^0) \delta(\underline{P}_j^2 - m_j^2 c^2)$$

$$= \frac{1}{2 p_j^0}$$

using $\delta(x^2 - a^2)$

$$= \frac{1}{2|a|} (\delta(x-a) + \delta(x+a))$$

(shown later)

(6a)

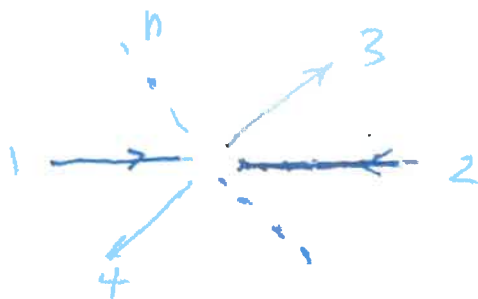
After integrating away $\int dP_j^0$, the differential decay rate is

$$\Gamma = \frac{S}{2\pi m_1} \int |M|^2 (2\pi)^4 \delta^{(4)}(P_1 - P_2 - P_3 \cdots P_n) \cdot \frac{n}{\pi} \left(\frac{1}{2 P_j^0} \frac{d^3 P_j}{(2\pi)^3} \right),$$

$$P_j^0 = \sqrt{\vec{P}_j^2 + m_j^2 c^2}$$

scattering

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$$1 + 2 \rightarrow 3 + 4 + \dots + n$$

The formula is

$$d\sigma = \frac{s \hbar^2}{4 \cdot \sqrt{(\underline{P}_1 \cdot \underline{P}_2)^2 - (m_1 m_2 c^2)^2}} \cdot |\mathcal{M}|^2.$$

$$(2\pi)^4 \delta^{(4)}(\underline{P}_1 + \underline{P}_2 - \underline{P}_3 - \underline{P}_4 - \dots - \underline{P}_n).$$

$$\prod_{j=3}^n \frac{d^4 P_j}{(2\pi)^4} \theta(P_j^0) \cdot 2\pi \delta(\underline{P}_j^2 - m_j^2 c^2)$$

Integrating
 $\rightarrow \int dP_j^0$

$$d\sigma = \frac{s \hbar^2}{4 \sqrt{(\underline{P}_1 \cdot \underline{P}_2)^2 - (m_1 m_2 c^2)^2}} \cdot |\mathcal{M}|^2.$$

$$(2\pi)^4 \delta^{(4)}(\underline{P}_1 + \underline{P}_2 - \underline{P}_3 - \underline{P}_4 - \dots - \underline{P}_n).$$

$$\prod_{j=3}^n \frac{d^3 P_j}{(2\pi)^3 2P_j^0}$$

Basically, learn how to reduce 4-dimensional integral to 3-dimensional, then to 1-dimensional integral

To show

$$\delta(x^2 - a^2) = \frac{1}{2|a|} (\delta(x-a) + \delta(x+a))$$

Proof

By definition, for a smooth function $f(x)$,

$$\int_{-b}^b f(x) \delta(x-a) dx = \begin{cases} f(a) & \text{if } a \in [-b, b] \\ 0 & \text{if } a \notin [-b, b] \end{cases}$$

LHS

$$= \int_{-\infty}^{\infty} f(x) \delta(x^2 - a^2) dx$$

$$= \int_{-\infty}^0 f(x) \delta(x^2 - a^2) dx + \int_0^{\infty} f(x) \delta(x^2 - a^2) dx$$

$$= \int_0^{\infty} f(-x) \delta(x^2 - a^2) dx + \int_0^{\infty} f(x) \delta(x^2 - a^2) dx$$

$$= \int_0^{\infty} f(-\sqrt{y}) \delta(y - a^2) \frac{dy}{2\sqrt{y}} + \int_0^{\infty} f(\sqrt{y}) \delta(y - a^2) \frac{dy}{2\sqrt{y}}$$

$$= f(-a) \frac{1}{2a} + f(a) \frac{1}{2a} \quad \text{assuming } a \geq 0$$

$$= \frac{1}{2|a|} (f(-a) + f(a))$$

$$\text{RHS} = \int_{-\infty}^{\infty} f(x) \cdot \frac{1}{2|a|} (\delta(x-a) + \delta(x+a)) dx$$

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$$RHS = \int_{-\infty}^0 f(x) \frac{1}{2|a|} (\delta(x-a) + \delta(x+a)) dx$$

$$+ \int_0^{\infty} f(x) \frac{1}{2|a|} (\delta(x-a) + \delta(x+a)) dx$$

$$= \frac{1}{2|a|} \int_{-\infty}^0 f(x) \delta(x+a) dx + \frac{1}{2|a|} \int_0^{\infty} f(x) \delta(x-a) dx$$

assume $a \geq 0$

$$= \frac{1}{2a} f(-a) + \frac{1}{2a} f(a)$$

$$= \frac{1}{2|a|} (f(-a) + f(a))$$

q.e.d.

$$\int_{-\infty}^{\infty} dp^0 \theta(p^0) \delta(\underline{p}^2 - m^2 c^2) f(p^0)$$

$$\underline{p}^2 = p^0{}^2 - \underline{p}^2$$

$$= \int_{-\infty}^{\infty} dp^0 \theta(p^0) \frac{1}{2|a|} [\delta(p^0 - a) + \delta(p^0 + a)] f(p^0)$$

$$p^0 = a = (\underline{p}^2 + m^2 c^2)^{\frac{1}{2}}$$

$$= \frac{1}{2|a|} \int_{-\infty}^{\infty} dp^0 \delta(p^0 - a) f(p^0)$$

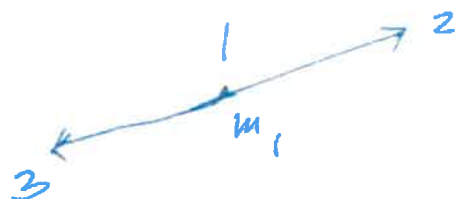
$$= \frac{1}{2|a|} f(a), \quad p^0 = a = \sqrt{\underline{p}^2 + m^2 c^2}$$

$$\therefore \int_{-\infty}^{\infty} dp^0 \theta(p^0) \delta(\underline{p}^2 - m^2 c^2) = \frac{1}{2p^0}, \quad \begin{matrix} a \geq 0 \\ p^0 = a \end{matrix}$$

Consider 2-particle decays

$$1 \rightarrow 2 + 3$$

Assume particle 1 at rest and decays



The decay rate is given by (page 6a)

$$\Gamma = \frac{S}{2\hbar m_1} \int |\mathcal{M}|^2 (2\pi)^4 \delta^{(4)}(\underline{P}_1 - \underline{P}_2 - \underline{P}_3) \cdot$$

$$\cdot \prod_{j=2}^3 \left(\frac{1}{2P_j^0} \frac{d^3 \underline{P}_j}{(2\pi)^3} \right)$$

scattering amplitude
 \mathcal{M}

$$= \mathcal{M}(\underline{P}_1, \underline{P}_2, \underline{P}_3 \dots)$$

$$= \frac{S}{8\pi^2 \hbar m_1} \int |\mathcal{M}|^2 \delta^{(4)}(\underline{P}_1 - \underline{P}_2 - \underline{P}_3) \frac{d^3 \underline{P}_2}{2P_2^0} \cdot \frac{d^3 \underline{P}_3}{2P_3^0}$$

$$= \frac{S}{8\pi^2 \hbar m_1} \int |\mathcal{M}|^2 \delta(P_1^0 - P_2^0 - P_3^0) \delta^{(3)}(\underline{P}_1 - \underline{P}_2 - \underline{P}_3) \cdot \frac{d^3 \underline{P}_2}{2P_2^0} \frac{d^3 \underline{P}_3}{2P_3^0}$$