

# Homework 1 solution

1.  $\vec{r} = r \hat{r}$  in spherical coordinates

$\vec{r} = x \hat{x} + y \hat{y} + z \hat{z}$  in Cartesian coordinates

Can switch freely to calculate the expressions.

$$(1) \quad \nabla \times \left( \frac{\vec{r}}{r} \right) = \nabla \times (1 \cdot \hat{r})$$

↓ check the formula for  $\nabla \times V$  in spherical coordinates

$$= \frac{1}{r \sin \theta} \left[ \frac{\partial (0)}{\partial \theta} - \frac{\partial (0)}{\partial \phi} \right] \hat{r} + \frac{1}{r} \left[ \frac{1}{\sin \theta} \frac{\partial (1)}{\partial \phi} - \frac{\partial (0)}{\partial r} \right] \hat{\theta} + \frac{1}{r} \left[ \frac{\partial (0)}{\partial r} - \frac{\partial (1)}{\partial \theta} \right] \hat{\phi}$$

$$= 0$$

$$(2) \quad \nabla \cdot \left( \frac{\vec{r}}{r} \right) = \nabla \cdot (1 \cdot \hat{r})$$

↓ check formula for  $\nabla \cdot V$  in spherical coordinates

$$= \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \cdot 1) + \frac{1}{r \sin \theta} \frac{\partial (0)}{\partial \theta} + \frac{1}{r \sin \theta} \frac{\partial (0)}{\partial \phi} = \frac{2}{r}$$

$$(3) \quad \nabla (\vec{a} \cdot \vec{r})$$

↓ product rule

$$= \underbrace{\vec{a} \times (\nabla \times \vec{r})}_{(1)} + \underbrace{\vec{r} \times (\nabla \times \vec{a})}_{(2)} + \underbrace{(\vec{a} \cdot \nabla) \vec{r}}_{(3)} + \underbrace{(\vec{r} \cdot \nabla) \vec{a}}_{(4)}$$

(2) & (4) both zero, because  $\vec{a}$  is a constant vector

For (1),  $\nabla \times \vec{r} = 0$  (verify with expression of  $\nabla \times$  in spherical coordinates)

$$= (\vec{a} \cdot \nabla) \vec{r}$$

$$= \left( a_x \frac{\partial}{\partial x} + a_y \frac{\partial}{\partial y} + a_z \frac{\partial}{\partial z} \right) (x \hat{x} + y \hat{y} + z \hat{z})$$

(1)

$$= a_x \hat{x} + a_y \hat{y} + a_z \hat{z} = \vec{a}$$

Another way to show  $(\vec{a} \cdot \nabla) \vec{r} = \vec{a}$  is to use spherical coordinates

$$(\vec{a} \cdot \nabla) \vec{r}$$

$$= \left( a_r \frac{\partial}{\partial r} + \frac{a_\theta}{r} \frac{\partial}{\partial \theta} + \frac{a_\phi}{r \sin \theta} \frac{\partial}{\partial \phi} \right) (r \hat{r})$$

$$= \underbrace{\hat{r} a_r \frac{\partial r}{\partial r}}_{a_r \hat{r}} + \underbrace{\frac{a_\theta}{r} \cdot r \frac{\partial \hat{r}}{\partial \theta}}_{a_\theta \hat{\theta}} + \underbrace{\frac{a_\phi}{r \sin \theta} \cdot r \frac{\partial \hat{r}}{\partial \phi}}_{a_\phi \hat{\phi}}$$

$$= a_r \hat{r} + a_\theta \hat{\theta} + a_\phi \hat{\phi} = \vec{a}$$

$$(4) \nabla \cdot [(\vec{a} \cdot \vec{r}) \vec{b}]$$

↓ product rule, treat  $(\vec{a} \cdot \vec{r})$  as scalar field  $f$

$$= (\vec{a} \cdot \vec{r}) (\nabla \cdot \vec{b}) + \vec{b} \cdot [\nabla (\vec{a} \cdot \vec{r})]$$

$$\downarrow \nabla \cdot \vec{b} = 0, \text{ \& } \nabla (\vec{a} \cdot \vec{r}) = \vec{a}$$

$$= \vec{b} \cdot \vec{a} = \vec{a} \cdot \vec{b}$$

$$(5) \nabla \times [(\vec{a} \cdot \vec{r}) \vec{r}]$$

$$= (\vec{a} \cdot \vec{r}) (\nabla \times \vec{r}) - \vec{r} \times [\nabla (\vec{a} \cdot \vec{r})]$$

$$\downarrow \nabla \times \vec{r} = 0, \quad \nabla (\vec{a} \cdot \vec{r}) = \vec{a}$$

$$= -\vec{r} \times \vec{a} = \vec{a} \times \vec{r}$$

2. The divergence theorem reads as

$$\int_V (\nabla \cdot \vec{v}) d\tau = \oint_S \vec{v} \cdot d\vec{a}$$

$$\text{For } \vec{v} = x^2 \hat{x} + y^2 \hat{y} + z \hat{z}$$

$$\nabla \cdot \vec{v} = \frac{\partial}{\partial x} (x^2) + \frac{\partial}{\partial y} (y^2) + \frac{\partial}{\partial z} (z) = 2x + 2y + 1$$

In the defined volume

$$\int_V (\nabla \cdot \vec{v}) d\tau = \int_0^a \int_0^a \int_0^a (2x + 2y + 1) dx dy dz$$

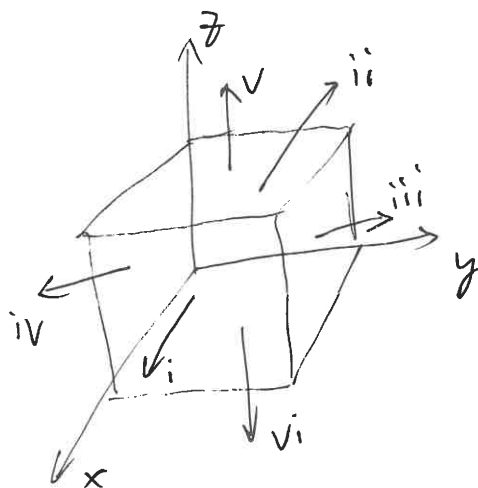
$$= \iiint 2x dx dy dz + \iiint 2y dx dy dz + \iiint 1 dx dy dz$$

$$= a \cdot a \cdot x^2 \Big|_0^a + a \cdot a \cdot y^2 \Big|_0^a + a \cdot a \cdot a$$

$$= 2a^4 + a^3$$

For the surface integral

$$\oint_S \vec{v} \cdot d\vec{a} = \int_i \vec{v} \cdot d\vec{a} + \int_{ii} \vec{v} \cdot d\vec{a} + \dots + \int_{vi} \vec{v} \cdot d\vec{a}$$



6 surfaces intotal defined in the graph

$$\int_i \vec{v} \cdot d\vec{a} = \int_0^a \int_0^a (x^2 \hat{x} + y^2 \hat{y} + z \hat{z}) \cdot (\hat{x} dy dz) = \int_0^a \int_0^a x^2 dy dz$$

$$x=a \text{ on surface } i \Rightarrow \int_i \vec{v} \cdot d\vec{a} = a^4$$

$$\int_{ii} \vec{v} \cdot d\vec{a} = \int_0^a \int_0^a (x^2 \hat{x} + y^2 \hat{y} + z \hat{z}) \cdot (-\hat{x} dy dz) = - \int_0^a \int_0^a x^2 dy dz$$

$$x=0 \text{ on surface ii} \Rightarrow \int_{ii} \vec{v} \cdot d\vec{a} = 0$$

$$\int_{iii} \vec{v} \cdot d\vec{a} = \int_0^a \int_0^a (x^2 \hat{x} + y^2 \hat{y} + z \hat{z}) \cdot (\hat{y} dx dz) = \int_0^a \int_0^a y^2 dx dz$$

$$y=a \text{ on surface iii} \Rightarrow \int_{iii} \vec{v} \cdot d\vec{a} = a^4$$

$$\int_{iv} \vec{v} \cdot d\vec{a} = \int_0^a \int_0^a -y^2 dx dz \text{ while } y=0 \Rightarrow \int_{iv} \vec{v} \cdot d\vec{a} = 0$$

$$\int_v \vec{v} \cdot d\vec{a} = \int_0^a \int_0^a (x^2 \hat{x} + y^2 \hat{y} + z \hat{z}) \cdot (\hat{z} dx dy) = \int_0^a \int_0^a z dx dy$$

$$z=a \text{ on surface v} \Rightarrow \int_v \vec{v} \cdot d\vec{a} = a^3$$

$$\int_{vi} \vec{v} \cdot d\vec{a} = \int_0^a \int_0^a (-z) dx dy \text{ while } z=0 \Rightarrow \int_{vi} \vec{v} \cdot d\vec{a} = 0$$

$$\text{In total, } \oint_S \vec{v} \cdot d\vec{a} = a^4 + 0 + a^4 + 0 + a^3 + 0 = 2a^4 + a^3$$

$$\text{Noticing } \oint_V (\nabla \cdot \vec{v}) d\tau = 2a^4 + a^3, \text{ proved } \checkmark$$

3. The Stoke's theorem reads as

$$\int_S (\nabla \times \vec{v}) \cdot d\vec{a} = \oint_P \vec{v} \cdot d\vec{l}$$

$$\text{For } \vec{v} = s^2 \hat{\phi} + \phi \hat{z}$$

$$\nabla \times \vec{v} = \left[ \frac{1}{s} \frac{\partial(\phi)}{\partial \phi} - \frac{\partial(s^2)}{\partial z} \right] \hat{s} + \left( -\frac{\partial \phi}{\partial s} \right) \hat{\phi} + \frac{1}{s} \left[ \frac{\partial(s \cdot s^2)}{\partial s} \right] \hat{z}$$

$$= \frac{1}{s} \hat{s} + 3s \hat{z}$$

In the defined surface  $d\vec{a} = \hat{z} s \cdot ds d\phi$

$$\Rightarrow \int_S (\nabla \times \vec{v}) \cdot d\vec{a} = \int_0^{2\pi} \int_0^a \left( \frac{1}{s} \hat{s} + 3s \hat{z} \right) \cdot \hat{z} s ds d\phi$$

$$= \int_0^{2\pi} \int_0^a 3s^2 ds d\phi = 2\pi s^3 \Big|_0^a = 2\pi a^3$$

Next we examine  $\oint_P \vec{v} \cdot d\vec{u}$  along the boundary of the surface

$d\vec{u} = ds \hat{s} + s d\phi \hat{\phi} + dz \hat{z}$  in cylindrical coordinates

On the circular boundary line,  $s=a$ ,  $z=0$ ,  $\Rightarrow ds=dz=0$

$$\Rightarrow d\vec{u} = s d\phi \hat{\phi}$$

$$\oint_P \vec{v} \cdot d\vec{u} = \int_0^{2\pi} (s^2 \hat{\phi} + \phi \hat{z}) \cdot \hat{\phi} s d\phi$$

$$= \int_0^{2\pi} s^3 d\phi = a^3 \int_0^{2\pi} d\phi = 2\pi a^3$$

Noticing  $\int_S (\nabla \times \vec{v}) \cdot d\vec{a} = 2\pi a^3$  above, proved  $\checkmark$ .



$$4. \quad \vec{E} = \frac{1}{4\pi\epsilon_0} \int \frac{\lambda(\vec{r}')}{r^2} \hat{r} \, dl'$$

In cylindrical coordinates

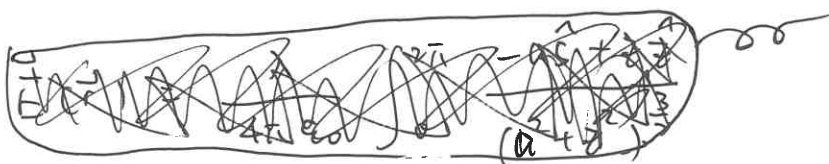
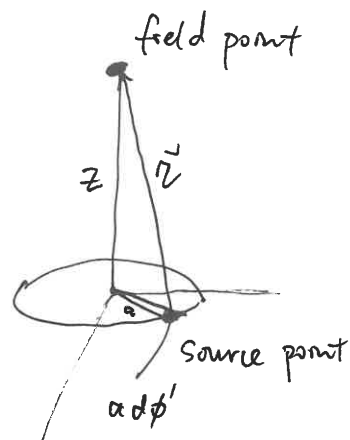
Source point  $\vec{r}' = a\hat{s}$

field point  $\vec{r} = z\hat{z}$

separation  $\vec{r} = \vec{r} - \vec{r}' = -a\hat{s} + z\hat{z}$

$$r = \sqrt{a^2 + z^2}$$

On the other hand  $dl' = a \, d\phi'$ ,  $\lambda(\vec{r}') = \lambda$ , constant



$z$ -component of  $\vec{E}$

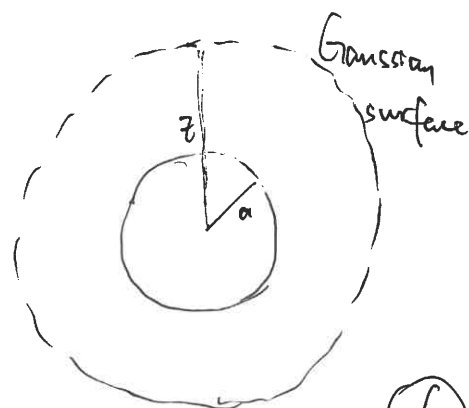
$$\vec{E}(\vec{r}) \cdot \hat{z} = \frac{a\lambda}{4\pi\epsilon_0} \int_0^{2\pi} \frac{z}{(a^2 + z^2)^{3/2}} \, d\phi'$$

$$= \frac{a\lambda}{4\pi\epsilon_0} \frac{z}{(a^2 + z^2)^{3/2}} \cdot 2\pi = \frac{a\lambda}{2\epsilon_0} \frac{z}{(a^2 + z^2)^{3/2}}$$

5. For use of the Gauss's law

Construct a spherical shell with radius  $z$

as the Gaussian surface for integration



(6)

$$\oint \vec{E} \cdot d\vec{a} = \frac{1}{\epsilon_0} Q_{enc}$$

Spherical symmetry  $\Rightarrow \vec{E} = E \hat{r}$

$$\oint \vec{E} \cdot d\vec{a} = E \oint da = E \cdot 4\pi r^2$$

On the other hand,  $Q_{enc} = \sigma \cdot \underset{\substack{\uparrow \\ \text{Surface area} \\ \text{of charge-carrying ball}}}{A} = \sigma \cdot 4\pi a^2$

$$\Rightarrow E \cdot 4\pi r^2 = \frac{\sigma}{\epsilon_0} 4\pi a^2 \Rightarrow E = \frac{\sigma a^2}{\epsilon_0 r^2}$$

Now we use Coulomb's law  $\vec{E} = \frac{1}{4\pi\epsilon_0} \int \frac{\sigma(r')}{r^2} \hat{r} da'$

Using the law of cosines:

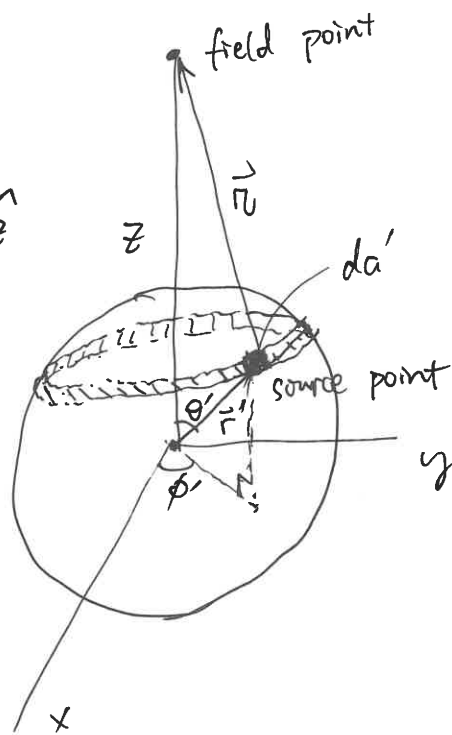
$$r^2 = z^2 + a^2 - 2za \cos \theta'$$

$$\hat{r} = a \sin \theta' \cos \phi' \hat{x} + a \sin \theta' \sin \phi' \hat{y} + (z - a \cos \theta') \hat{z}$$

$\sigma(r') = \sigma$ , uniform throughout the shell

$$da' = a^2 \sin \theta' d\theta' d\phi'$$

Therefore



$$\vec{E} \cdot \hat{z} = \frac{\sigma}{4\pi\epsilon_0} \int_0^{2\pi} \int_0^\pi \frac{a^2 (z - a \cos \theta') \sin \theta'}{(z^2 + a^2 - 2za \cos \theta')^{\frac{3}{2}}} d\theta' d\phi'$$

$$= \frac{\sigma \cdot 2\pi}{4\pi\epsilon_0} \int_0^\pi \frac{a^2 (z - a \cos \theta') \sin \theta'}{(z^2 + a^2 - 2za \cos \theta')^{\frac{3}{2}}} d\theta'$$

(7)

Define  $u \equiv z^2 + a^2 - 2za \cos \theta'$

$$\Rightarrow du = -2za \sin \theta' d\theta'$$

$z - a \cos \theta' = \frac{1}{2z} (u + z^2 - a^2)$  appears in the integral

$$\vec{E} \cdot \hat{z} = \frac{\sigma \cdot 2\pi}{4\pi\epsilon_0} \int_{(z-a)^2}^{(z+a)^2} a^2 \cdot \frac{1}{2z} (u + z^2 - a^2) \cdot \frac{1}{2za} \cdot u^{-\frac{3}{2}} du$$

$$= \frac{\sigma}{2\epsilon_0} \cdot \frac{a}{4z^2} \int_{(z-a)^2}^{(z+a)^2} (u + z^2 - a^2) u^{-\frac{3}{2}} du$$

$$= \frac{\sigma a}{8\epsilon_0 z^2} \cdot \left[ \int_{(z-a)^2}^{(z+a)^2} u^{-\frac{1}{2}} du + (z^2 - a^2) \int_{(z-a)^2}^{(z+a)^2} u^{-\frac{3}{2}} du \right]$$

$$= \frac{\sigma a}{8\epsilon_0 z^2} \cdot \left[ 2u^{\frac{1}{2}} \Big|_{(z-a)^2}^{(z+a)^2} + (z^2 - a^2)(-2) u^{-\frac{1}{2}} \Big|_{(z-a)^2}^{(z+a)^2} \right]$$

$$= \frac{\sigma a}{8\epsilon_0 z^2} \cdot [4a + 4a] = \frac{\sigma a^2}{\epsilon_0 z^2}$$

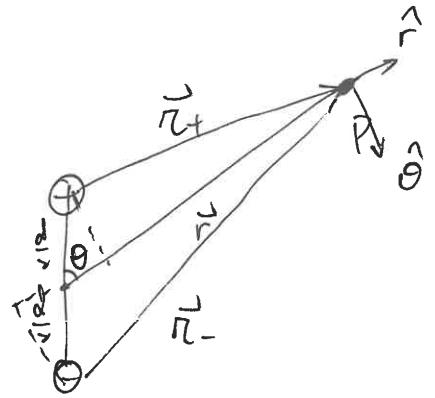
This result agrees with that using the Gauss's law.



6. In spherical coordinates

$$\vec{r}_+ = (r - \frac{d}{2} \cos \theta) \hat{r} + \frac{d}{2} \sin \theta \hat{\theta}$$

$$\vec{r}_- = (r + \frac{d}{2} \cos \theta) \hat{r} - \frac{d}{2} \sin \theta \hat{\theta}$$



$$r_+^2 = r^2 + (\frac{d}{2})^2 - dr \cos \theta \approx r^2 - dr \cos \theta$$

$$r_-^2 = r^2 + (\frac{d}{2})^2 + dr \cos \theta \approx r^2 + dr \cos \theta$$

$$\vec{E} = \frac{q}{4\pi\epsilon_0} \left( \frac{\hat{r}_+}{r_+^2} - \frac{\hat{r}_-}{r_-^2} \right)$$

$$= \frac{q}{4\pi\epsilon_0} \left( \frac{\vec{r}_+}{r_+^3} - \frac{\vec{r}_-}{r_-^3} \right)$$

$$\vec{E} \cdot \hat{r} = \frac{q}{4\pi\epsilon_0} \left[ \frac{r - \frac{d}{2} \cos \theta}{(r^2 - dr \cos \theta)^{\frac{3}{2}}} - \frac{r + \frac{d}{2} \cos \theta}{(r^2 + dr \cos \theta)^{\frac{3}{2}}} \right]$$

$$\approx \frac{q}{4\pi\epsilon_0} \left[ \frac{r - \frac{d}{2} \cos \theta}{(r - \frac{d}{2} \cos \theta)^3} - \frac{r + \frac{d}{2} \cos \theta}{(r + \frac{d}{2} \cos \theta)^3} \right]$$

$$\approx \frac{q}{4\pi\epsilon_0} \left[ \frac{1}{(r - \frac{d}{2} \cos \theta)^2} - \frac{1}{(r + \frac{d}{2} \cos \theta)^2} \right]$$

Note that  $(r - \frac{d}{2} \cos \theta)^2 = r^2 (1 - \frac{d}{2r} \cos \theta)^2 \approx r^2 (1 - \frac{d}{r} \cos \theta)$

$$\vec{E} \cdot \hat{r} \approx \frac{q}{4\pi\epsilon_0} \frac{1}{r^2} \left[ \frac{1}{1 - \frac{d}{r} \cos \theta} - \frac{1}{1 + \frac{d}{r} \cos \theta} \right]$$

$$\approx \frac{q}{4\pi\epsilon_0} \frac{1}{r^2} \frac{2d}{r} \cos \theta = \frac{2qd}{4\pi\epsilon_0 r^3} \cos \theta$$

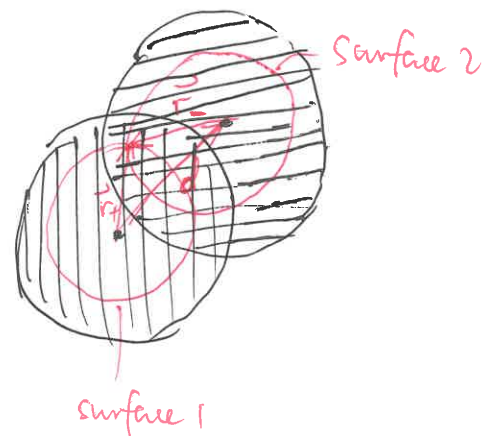
$$\vec{E} \cdot \hat{\theta} = \frac{q}{4\pi\epsilon_0} \left[ \frac{\frac{d}{2} \sin \theta}{(r^2 - dr \cos \theta)^{\frac{3}{2}}} - \frac{-\frac{d}{2} \sin \theta}{(r^2 + dr \cos \theta)^{\frac{3}{2}}} \right]$$

$$\approx \frac{q}{4\pi\epsilon_0} \frac{\frac{d}{2} \sin \theta - (-\frac{d}{2} \sin \theta)}{r^3}$$

$$= \frac{qd \sin \theta}{4\pi\epsilon_0 r^3}$$

Summarizing above,  $\vec{E} = \frac{2qd}{4\pi\epsilon_0 r^3} \cos \theta \hat{r} + \frac{qd \sin \theta}{4\pi\epsilon_0 r^3} \hat{\theta}$

7. First apply Gauss's law to "+" charge cloud on surface 1 (which intersects that point in the overlapping region).



$\vec{r}_{\pm}$  denotes vector going from the "+" charge center to the point of interest

Gauss's law for surface 1 on "+" charge

$$\oint \vec{E} \cdot d\vec{a} = \frac{Q_{enc}}{\epsilon_0}$$

$$\oint \vec{E} \cdot d\vec{a} = E \cdot 4\pi r_+^2, \quad \frac{Q_{enc}}{\epsilon_0} = \frac{\frac{4}{3}\pi r_+^3 \cdot \rho}{\epsilon_0}$$

$$\Rightarrow E_+ = \frac{1}{4\pi r_+^2} \frac{\frac{4}{3}\pi r_+^3 \cdot \rho}{\epsilon_0} = \frac{\rho r_+}{3\epsilon_0}$$

$$\vec{E}_+ = \frac{\rho r_+}{3\epsilon_0} \hat{r}_+ = \frac{\rho \vec{r}_+}{3\epsilon_0}$$

Gauss's law for surface 2 on "-" charge

$$\vec{E}_- = \frac{-\rho \vec{r}_-}{3\epsilon_0}$$

According to superposition principle,

$$\vec{E}_{tot} = \vec{E}_+ + \vec{E}_- = \frac{\rho}{3\epsilon_0} (\vec{r}_+ - \vec{r}_-) = \frac{\rho}{3\epsilon_0} \vec{d}$$

Which is a constant