PC3261: Classical Mechanics II

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Semester II, 2024/25

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Notes

Lecture 3: Linear Momentum

Linear momentum of two-particle system

ullet Forces are *assumed* to obey principle of superposition of forces: ${f f}_{12}$ is the force acting on m_1 due to m_2

$$\begin{cases} \mathbf{F}_1(t) = \mathbf{F}_1^{\mathsf{ext}}(t) + \mathbf{f}_{12}(t) \\ \mathbf{F}_2(t) = \mathbf{F}_2^{\mathsf{ext}}(t) + \mathbf{f}_{21}(t) \end{cases}$$

 Total linear momentum of the system: forces between particles are assumed to obey Newton's third law of motion

$$\mathbf{P}(t) \equiv \mathbf{p}_1(t) + \mathbf{p}_2(t) \quad \Rightarrow \quad \dot{\mathbf{P}}(t) = \mathbf{F}_1^{\mathsf{ext}}(t) + \mathbf{F}_2^{\mathsf{ext}}(t)$$

 Newton's second law: the time rate of change of total linear momentum of the two-particle system equals to the total external force acting upon it

$$\dot{\mathbf{P}}(t) = \mathbf{F}^{\mathsf{ext}}(t), \qquad \mathbf{F}^{\mathsf{ext}}(t) \equiv \mathbf{F}_1^{\mathsf{ext}}(t) + \mathbf{F}_2^{\mathsf{ext}}(t)$$

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Linear momentum of multi-particle system

ullet Total force acting on the lpha-particle: ${f f}_{lphaeta}$ is the force acting on m_lpha due to m_eta

$$\mathbf{F}_{lpha}(t) = \mathbf{F}^{\mathsf{ext}}_{lpha}(t) + \sum_{eta=1, eta
eq lpha}^{N} \mathbf{f}_{lphaeta}(t), \qquad lpha=1, 2, 3, \cdots, N$$

• Total linear momentum of multi-particle system:

$$\mathbf{P}(t) \equiv \sum_{lpha=1}^{N} \mathbf{p}_{lpha}(t)$$

• Newton's second law: the time rate of change of total linear momentum of multi-particle system equals to the total *external* force acting upon it

$$\dot{\mathbf{P}}(t) = \mathbf{F}^{\mathrm{ext}}(t)\,, \qquad \mathbf{F}^{\mathrm{ext}}(t) \equiv \sum_{lpha=1}^{N} \mathbf{F}_{lpha}^{\mathrm{ext}}(t)$$

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Impuse-Momentum theorem

• Impulse-Momentum theorem (integral form of the Newton's second law): change of total linear momentum equals to the time integral of the total external force

$$\mathbf{F}^{\mathsf{ext}}(t) = \frac{\mathrm{d}\mathbf{P}(t)}{\mathrm{d}t} \qquad o \qquad \int_{t_1}^{t_2} \mathbf{F}^{\mathsf{ext}}(t) \, \mathrm{d}t = \mathbf{P}(t_2) - \mathbf{P}(t_1)$$

• Conservation law of linear momentum: if the total external force on a multiparticle system is zero, then the total linear momentum of the multi-particle is a constant

$$\dot{\mathbf{P}}(t) = \mathbf{F}^{\mathrm{ext}}(t) = \mathbf{0} \quad \Rightarrow \quad \mathbf{P}(t) = \mathrm{constant} \quad \Rightarrow \quad \mathbf{P}(t_1) = \mathbf{P}(t_2) \quad \forall \quad t_1, t_2$$

• The validity of the conservation law of linear momentum depends crucially on the *experimental* basis of the Newton's third law!

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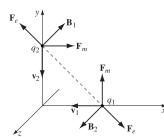
A violation of Newton's third law???

- ullet Two point charges, q_1 and q_2 , are moving at uniform velocities ${f v}_1$ and ${f v}_2$ respectively
- Electric fields and forces:

$$\begin{cases} \mathbf{E}_{1}(\mathbf{r}_{2}) = \frac{q_{1}}{4\pi\epsilon_{0}} \frac{\mathbf{r}_{2} - \mathbf{r}_{1}}{|\mathbf{r}_{2} - \mathbf{r}_{1}|^{3}} \\ \mathbf{E}_{2}(\mathbf{r}_{1}) = \frac{q_{2}}{4\pi\epsilon_{0}} \frac{\mathbf{r}_{1} - \mathbf{r}_{2}}{|\mathbf{r}_{1} - \mathbf{r}_{2}|^{3}} \end{cases} \Rightarrow \begin{cases} \mathbf{F}_{e,12} = q_{1}\mathbf{E}_{2}(\mathbf{r}_{1}) \\ \mathbf{F}_{e,21} = q_{2}\mathbf{E}_{1}(\mathbf{r}_{2}) \end{cases}$$

• Electric forces obey Newton's third law

$$\mathbf{F}_{e,12} = -\mathbf{F}_{e,21}$$



Notes

A violation of Newton's third law??? - cont'd

• Magnetic fields and forces:

$$\begin{cases}
\mathbf{B}_{1}(\mathbf{r}_{2}) = \frac{\mu_{0}q_{1}}{4\pi} \frac{\mathbf{v}_{1} \times (\mathbf{r}_{2} - \mathbf{r}_{1})}{|\mathbf{r}_{2} - \mathbf{r}_{1}|^{3}} \\
\mathbf{B}_{2}(\mathbf{r}_{1}) = \frac{\mu_{0}q_{2}}{4\pi} \frac{\mathbf{v}_{2} \times (\mathbf{r}_{1} - \mathbf{r}_{2})}{|\mathbf{r}_{1} - \mathbf{r}_{2}|^{3}}
\end{cases} \Rightarrow \begin{cases}
\mathbf{F}_{m,12} = q_{1}\mathbf{v}_{1} \times \mathbf{B}_{2}(\mathbf{r}_{1}) \\
\mathbf{F}_{m,21} = q_{2}\mathbf{v}_{2} \times \mathbf{B}_{1}(\mathbf{r}_{2})
\end{cases}$$

 Magnetic forces do not obey Newton's third law!

$$\mathbf{F}_{m,12} \neq -\mathbf{F}_{m,21}$$

• Electromagnetic linear momentum density: fields also possess linear momentum!

$$\mathbf{g}_{\mathsf{EM}}(\mathbf{r}) = \epsilon_0 \, \mathbf{E}(\mathbf{r}) \times \mathbf{B}(\mathbf{r})$$

 \mathbf{F}_{m} \mathbf{F}_{m} \mathbf{F}_{m} \mathbf{F}_{m} \mathbf{F}_{m} \mathbf{F}_{m} \mathbf{F}_{m} \mathbf{F}_{m} \mathbf{F}_{m}

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System with variable mass

Newton's second law with variable mass:

$$\mathbf{F}^{\mathsf{ext}}(t) = \frac{\mathrm{d}}{\mathrm{d}t} \left[m(t) \, \mathbf{v}(t) \right] \quad \xrightarrow{???} \quad \mathbf{F}^{\mathsf{ext}}(t) = \frac{\mathrm{d}m(t)}{\mathrm{d}t} \, \mathbf{v}(t) + m \, \frac{\mathrm{d}\mathbf{v}(t)}{\mathrm{d}t}$$

ullet Galilean velocity transformation: $\mathbf{v}'(t) = \mathbf{v}(t) - \mathbf{V}$

$$\frac{\mathrm{d}}{\mathrm{d}t} \left[m(t) \mathbf{v}'(t) \right] = \mathbf{F}^{\mathsf{ext}}(t) \quad \not\leftrightarrow \quad \frac{\mathrm{d}}{\mathrm{d}t} \left[m(t) \mathbf{v}(t) \right] = \mathbf{F}^{\mathsf{ext}}(t)$$

 \bullet There is no fundamental difficulty in handling any system with variable mass provided the same set of particles is included throughout the time interval t_1 to t_2

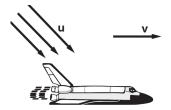
$$\int_{t_1}^{t_2} \mathbf{F}^{\mathsf{ext}}(t) \, \mathrm{d}t = \mathbf{P}(t_2) - \mathbf{P}(t_1)$$

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Example: Spacecraft and dust particles

• A spacecraft with mass M moves through space with constant velocity ${\bf v}$. The spacecraft encounters a stream of dust particles that embed themselves in the hull at rate dm/dt. The dust has velocity ${\bf u}$ just before it hits.



EXERCISE 3.1: Find the external force necessary to keep the spacecraft moving uniformly.

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Newton's second law with variable mass

- A system with mass m(t) moves at velovity $\mathbf{v}(t)$. Particles are added to the system at a rate $\mathrm{d}m(t)/\mathrm{d}t$. These particles have velocity $\mathbf{u}(t)$ just before entering the system.
- Newton's second law:

$$\begin{split} \mathbf{F}^{\mathsf{ext}}(t) &= \frac{\mathrm{d}}{\mathrm{d}t} \left[m(t) \, \mathbf{v}(t) \right] - \frac{\mathrm{d}m(t)}{\mathrm{d}t} \, \mathbf{u}(t) \\ \Rightarrow & m(t) \dot{\mathbf{v}}(t) = \mathbf{F}^{\mathsf{ext}}(t) + \dot{m}(t) \left[\mathbf{u}(t) - \mathbf{v}(t) \right] \end{split}$$

Galilean invariance is preserved:

$$\mathbf{F}^{\mathsf{ext}}(t) = \frac{\mathrm{d}}{\mathrm{d}t} \left[m(t)\mathbf{v}(t) \right] - \dot{m}(t)\mathbf{u}(t) \quad \leftrightarrow \quad \mathbf{F}^{\mathsf{ext}}(t) = \frac{\mathrm{d}}{\mathrm{d}t} \left[m(t)\mathbf{v}'(t) \right] - \dot{m}(t)\mathbf{u}'(t)$$

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Example: Rocket in a constant gravitational field

- A rocket is taking off from rest in a uniform gravitation field $\mathbf{g}=-g\,\hat{\mathbf{e}}_z$. The fuel is ejected at a constant rate $\dot{m}(t)=-k$ at a constant exhaust speed u relative to the rocket.
- Linear momentum of the system:

$$\begin{cases} \mathbf{P}(t) = m(t)\mathbf{v}(t) \\ \mathbf{P}(t + \Delta t) = \left[m(t) + \Delta m\right]\left[\mathbf{v}(t) + \Delta \mathbf{v}\right] + \left(-\Delta m\right)\left[\mathbf{v}(t) + \Delta \mathbf{v} + \mathbf{u}(t + \Delta t)\right] \end{cases}$$

Newton's second law:

$$m(t) \dot{\mathbf{v}}(t) - \mathbf{u}(t) \dot{m}(t) = \mathbf{F}^{\mathsf{ext}}(t)$$

EXERCISE 3.2: Find $\mathbf{v}(t)$ and $\mathbf{r}(t)$ for the rocket in its subsequent motion given that the initial mass of the rocket is m_0 .

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Center of mass

• Position vector of the center of mass of a multi-particle system:

$${f R}_{\sf CM}(t) \equiv rac{1}{M} \sum_{lpha=1}^N m_lpha {f r}_lpha(t) \,, \qquad M \equiv \sum_{lpha=1}^N m_lpha$$

• Velocity of the center of mass: total linear momentum of the system is equal to the linear momentum of the center of mass $as\ if$ it were a particle of mass M with velocity $\mathbf{V}_{\mathsf{CM}}(t)$

$$\mathbf{V}_{\mathsf{CM}}(t) \equiv \dot{\mathbf{R}}_{\mathsf{CM}}(t) = \frac{1}{M} \sum_{\alpha=1}^{N} m_{\alpha} \dot{\mathbf{r}}_{\alpha}(t) \quad \Rightarrow \quad \mathbf{P}(t) = M \mathbf{V}_{\mathsf{CM}}(t)$$

 \bullet Acceleration of the center of mass: center of mass moves exactly as if it were a single particle of mass M subjected to the total external force on the system

$$\mathbf{A}_{\mathsf{CM}}(t) \equiv \ddot{\mathbf{R}}_{\mathsf{CM}}(t) = \frac{1}{M} \sum_{\alpha=1}^{N} m_{\alpha} \ddot{\mathbf{r}}_{\alpha}(t) \quad \Rightarrow \quad \dot{\mathbf{P}}(t) = M \mathbf{A}_{\mathsf{CM}}(t)$$

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Example: Projectile motion

- A rigid object consists of two masses m_1 and m_2 separated by a light rod of length L. It is thrown into the air.
- Center of mass:

$$\mathbf{R}_{\mathsf{CM}}(t) = \frac{m_1 \mathbf{r}_1(t) + m_2 \mathbf{r}_2(t)}{m_1 + m_2}$$

• Equation of motion of the center of mass:

$$\mathbf{F}^{\mathsf{ext}}(t) = (m_1 + m_2) \, \ddot{\mathbf{R}}_{\mathsf{CM}}(t) \quad \Rightarrow \quad \ddot{\mathbf{R}}_{\mathsf{CM}}(t) = \mathbf{g}$$

• The center of mass follows the parabolic trajectory of a single mass, m_1+m_2 , in a uniform gravitational field (motions of m_1 and m_2 about the center of mass are to be analyzed separately)

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Center of mass of extended body

• Visualize mass element ${\rm d}m$ of volume ${\rm d}V$ located at position ${\bf r}(t)$ with mass density $\rho({\bf r})$:

$$\mathbf{R}_{\mathsf{CM}}(t) = \frac{1}{M} \iiint_{V} \mathbf{r}(t) \, \rho(\mathbf{r}) \, \mathrm{d}V$$

 $\, \bullet \,$ Center of mass of a uniform solid (upper) hemisphere: mass M and radius R

$$\mathbf{R}_{\mathsf{CM}} = \frac{3}{8} R \,\hat{\mathbf{e}}_z$$

EXERCISE 3.3: A thin non-uniform plates lies on the xy-plane with corners (0,0), (a,0), (0,b) and (a,b). Its surface mass density is $\sigma(x,y) = \sigma_0 xy/ab$ where σ_0 is a constant. Find its center of mass.

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Center-of-mass frame

• Center-of-mass frame is a reference frame at which the center of mass remains at the origin:

$$\mathbf{r}'_{\alpha}(t) \equiv \mathbf{r}_{\alpha}(t) - \mathbf{R}_{\mathsf{CM}}(t) \quad \Rightarrow \quad \mathbf{R}'_{\mathsf{CM}}(t) = \mathbf{0}$$

• Velocity of the center of mass in the center-of-mass frame: center of mass is stationary in the center-of-mass frame

$$\mathbf{V}_{\mathsf{CM}}'(t) = \frac{1}{M} \sum_{\alpha=1}^{N} m_{\alpha} \dot{\mathbf{r}}_{\alpha}'(t) = \mathbf{0}$$

• Acceleration of the center of mass in the center-of-mass frame:

$$\mathbf{A}'_{\mathsf{CM}}(t) = \frac{1}{M} \sum_{\alpha=1}^{N} m_{\alpha} \ddot{\mathbf{r}}'_{\alpha}(t) = \mathbf{0}$$

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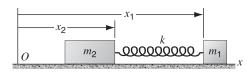
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Example: Two-body oscillations

- ullet Two idential blocks 1 and 2 each of mass m slide without friction on a straight track. They are connected by a massless spring with unstretched length L_0 and spring constant k. Initially, the system is at rest. At t=0, block 1 is hit sharply giving it an instanteneous velocity v_0 to the right.
- Equations of motion in the center-of-mass frame:

$$\begin{cases} m\ddot{x}_1'(t) = -k\left[x_1'(t) - x_2'(t) - L_0\right] \\ m\ddot{x}_2'(t) = +k\left[x_1'(t) - x_2'(t) - L_0\right] \end{cases} \xrightarrow{x_2} \xrightarrow{m_2}$$



EXERCISE 3.4: Find the velocities of each block at later times with respect to the track.

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