

Clebsch-Gordan Coefficients

Addition of angular momenta

$$\mathbf{J}_1 \quad \mathbf{J}_2$$

$$|\alpha_1 j_1 m_1\rangle = \text{basis for } J_1^2 \text{ and } J_{1z}$$

$$|\alpha_2 j_2 m_2\rangle = \text{basis for } J_2^2 \text{ and } J_{2z}$$

The base vectors

$$|\alpha j_1 j_2 m_1 m_2\rangle \equiv |\alpha_1 j_1 m_1\rangle |\alpha_2 j_2 m_2\rangle$$

α, j_1, j_2 fixed,

m_1, m_2 vary

$$-j_1 \leq m_1 \leq j_1$$

$$-j_2 \leq m_2 \leq j_2$$

span the subspace $\xi(\alpha, j_1, j_2)$.

$$\tilde{J}^2 = (\tilde{J}_1 + \tilde{J}_2)^2 \quad \text{and } J_z \text{ act on } \xi(\alpha, j_1, j_2)$$

Since \tilde{J}_1^2 and \tilde{J}_2^2 commute with \tilde{J}^2 and J_z , can also use the base

$$\begin{aligned} &|\alpha \ j_1 \ j_2 \ j \ m\rangle, \quad \alpha, j_1, j_2 \text{ fixed} \\ &\quad j, m \text{ vary} \\ &\quad |j_1 - j_2| \leq j \leq (j_1 + j_2) \\ &\quad -j \leq m \leq j \end{aligned}$$

to generate the same subspace $\xi(\alpha, j_1, j_2)$.

The two bases are related:

$$|\alpha \ j_1 \ j_2 \ j \ m\rangle = \sum_{m_2=-j_2}^{j_2} \sum_{m_1=-j_1}^{j_1} |\alpha \ j_1 \ j_2 \ m_1 \ m_2\rangle \langle j_1 \ j_2 \ m_1 \ m_2 | j \ m\rangle$$

$$|\alpha \ j_1 \ j_2 \ m_1 \ m_2\rangle = \sum_{m=-j}^j \sum_{j=|j_1-j_1|}^{(j_1+j_2)} |\alpha \ j_1 \ j_2 \ j \ m\rangle \langle j \ m | j_1 \ j_2 \ m_1 \ m_2\rangle$$

$$\langle j_1 \ j_2 \ m_1 \ m_2 | j \ m\rangle = \langle j \ m | j_1 \ j_2 \ m_1 \ m_2\rangle^* \quad \equiv \quad \text{Clebsch-Gordan coefficients}$$

Meaning of C.G. coeffs

- (i) relating two basis vectors (just like Fourier transform)
- (ii) $\langle j_1 j_2 m_1 m_2 | j m \rangle =$ probability amplitude of finding the state $| j_1 j_2 m_1 m_2 \rangle$ when the system is in state $| j m \rangle$

Properties of C.G. coeffs

- (1) Selection rule:

$$\langle \alpha j_1 j_2 m_1 m_2 | j m \rangle = 0 \text{ unless}$$

$$m_1 + m_2 = m \text{ and } |j_1 - j_2| \leq j \leq (j_1 + j_2)$$

- (2) Phase convention: require

$$\langle j_1 j_2 j_1 m_2 | j j \rangle \text{ real and } \geq 0$$

$$m_2 = j - j_1$$

$$j = |j_1 - j_2|, |j_1 - j_2| + 1, \dots, (j_1 + j_2)$$

Note: When $m_1 = j_1$ and $m = j$, it does not necessarily imply $m_2 = j_2$ since $j \neq (j_1 + j_2)$ in general

(3) Reality: All C.G. coeffs can be obtained from

$$\langle j_1 j_2 j_1 m_2 | j j \rangle$$

\therefore all C.G. coeffs are real

(4) Orthogonality

$$\sum_{m_1 m_2} \langle j_1 j_2 m_1 m_2 | j m \rangle \langle j_1 j_2 m_1 m_2 | j' m' \rangle = \delta_{jj'} \delta_{mm'}$$

$$\sum_{j m} \langle j_1 j_2 m_1 m_2 | j m \rangle \langle j_1 j_2 m_1' m_2' | j m \rangle = \delta_{m_1 m_1'} \delta_{m_2 m_2'}$$

Wigner-Eckart theorem

In a standard representation $\{J^2, J_z\}$ whose basis vectors are denoted by $|\tau j m\rangle$,

The matrix element $\langle \tau j m | T_g^{(k)} | \tau' j' m' \rangle$

of the q^{th} standard component of a given k^{th} order irreducible tensor operator, $T^{(k)}$, is equal to the product of the Clebsch-Gordan coefficient.

$$\langle j' k m' q | j m \rangle$$

by a quantity independent of m, m' and q ($q = -k, -k+1, \dots, +k$)

$$\langle \tau j m | T_q^{(k)} | \tau' j' m' \rangle = \frac{1}{\sqrt{2j+1}} \langle \tau j || T^{(k)} || \tau' j' \rangle \langle j' k m' q | j m \rangle$$

$\langle \tau j || T^{(k)} || \tau' j' \rangle$ = reduced matrix element

$\langle j' k m' q | j m \rangle$ = Clebsch-Gordan coefficient

$\neq 0$ only if $m = m' + q$ and $|j - j'| \leq k \leq j + j'$

For a scalar operator S $\langle \tau j m | S | \tau' j' m' \rangle = \delta_{jj'} \delta_{mm'} S_{\tau\tau'}^{(j)}$