

# PC3261: Classical Mechanics II

Kenneth HONG Chong Ming

Office: S16-07-06

Email: [phyhcmk@nus.edu.sg](mailto:phyhcmk@nus.edu.sg)

Semester II, 2024/25

Latest update: January 21, 2025 10:48am



Department of Physics  
Faculty of Science

## Lecture 2: Newton's Laws of Motion

# Newton's first law and inertia

- **Newton's first law:** a particle remains at rest or in uniform motion unless acted upon a force
- **Inertia** is the *resistance* of any particle to any change in its velocity and the quantitative measure of inertia is **mass**
- A mathematical description of the motion of a particle requires the choice of a **frame of reference** – a set of coordinates in space that can be used to specify the position, velocity and acceleration of the particle at any given instant of time
- A frame of reference at which Newton's first law is valid is called an **inertial frame of reference**

# Newton's second law

- **Linear momentum** of a particle is defined as the product of its mass and velocity

$$\mathbf{p}(t) \equiv m\mathbf{v}(t)$$

- **Newton's second law:** a particle acted upon a force moves in such a manner that the time rate of change of linear momentum equals the force

$$\mathbf{F}(t) = \frac{d\mathbf{p}(t)}{dt}$$

- Both Newton's first and second laws remain exactly true in special relativity with a *suitably* redefinition of linear momentum

# Newton's third law

- **Newton's third law:** if two particles exert forces on each other, these forces are equal in magnitude and opposite in direction
- **Central forces** are the forces acting along the line connecting two particles
- Velocity-dependent forces are non-central and Newton's third law *may* not apply
- Newton's third law is not valid in special relativity as the concept of absolute time is abandoned

# Galilean relativity

- Two inertial frames,  $\mathcal{O}$  and  $\mathcal{O}'$ , are oriented such that their spatial coordinate axes are parallel, their spatial origins are coincided when  $t = t' = 0$  and  $\mathcal{O}'$  moves at *uniform velocity*  $\mathbf{V}$  with respect to  $\mathcal{O}$

- Galilean boost:**

$$\begin{cases} t' = t \\ \mathbf{r}'(t) = \mathbf{r}(t) - \mathbf{V}t \end{cases}$$

- Galilean velocity transformation:

$$\mathbf{v}'(t) = \mathbf{v}(t) - \mathbf{V}$$

- Newton's laws are **Galilean invariance**

$$\begin{cases} \mathbf{r}(t) = x(t) \hat{\mathbf{e}}_x + y(t) \hat{\mathbf{e}}_y + z(t) \hat{\mathbf{e}}_z \\ \mathbf{r}'(t') = x'(t') \hat{\mathbf{e}}_{x'} + y'(t') \hat{\mathbf{e}}_{y'} + z'(t') \hat{\mathbf{e}}_{z'} \end{cases}, \quad \begin{cases} \hat{\mathbf{e}}_x = \hat{\mathbf{e}}_{x'} \\ \hat{\mathbf{e}}_y = \hat{\mathbf{e}}_{y'} \\ \hat{\mathbf{e}}_z = \hat{\mathbf{e}}_{z'} \end{cases}$$

$$t' = t \quad \Rightarrow \quad \mathbf{r}'(t') = \mathbf{r}'(t) = x'(t) \hat{\mathbf{e}}_x + y'(t) \hat{\mathbf{e}}_y + z'(t) \hat{\mathbf{e}}_z$$

$$\mathbf{v}'(t') \equiv \frac{d\mathbf{r}'(t')}{dt'} = \frac{d\mathbf{r}'(t)}{dt} = \frac{dx'(t)}{dt} \hat{\mathbf{e}}_x + \frac{dy'(t)}{dt} \hat{\mathbf{e}}_y + \frac{dz'(t)}{dt} \hat{\mathbf{e}}_z \equiv \mathbf{v}'(t)$$

$$\mathbf{r}'(t) = \mathbf{r}(t) - \mathbf{V}t \quad \Rightarrow \quad \frac{d\mathbf{r}'(t)}{dt} = \frac{d\mathbf{r}(t)}{dt} - \frac{d}{dt}(\mathbf{V}t) \quad \Rightarrow \quad \mathbf{v}'(t) = \mathbf{v}(t) - \mathbf{V} \quad \blacksquare$$

$$\Rightarrow \quad \begin{cases} \mathbf{v}(t) \equiv \frac{d\mathbf{r}(t)}{dt} = \frac{dx(t)}{dt} \hat{\mathbf{e}}_x + \frac{dy(t)}{dt} \hat{\mathbf{e}}_y + \frac{dz(t)}{dt} \hat{\mathbf{e}}_z \\ \mathbf{v}'(t) \equiv \frac{d\mathbf{r}'(t)}{dt} = \frac{dx'(t)}{dt} \hat{\mathbf{e}}_{x'} + \frac{dy'(t)}{dt} \hat{\mathbf{e}}_{y'} + \frac{dz'(t)}{dt} \hat{\mathbf{e}}_{z'} \end{cases} \quad \blacksquare$$

$$\begin{cases} \mathbf{r}(t) = x(t) \hat{\mathbf{e}}_x + y(t) \hat{\mathbf{e}}_y + z(t) \hat{\mathbf{e}}_z \\ \mathbf{r}'(t') = x'(t') \hat{\mathbf{e}}_{x'} + y'(t') \hat{\mathbf{e}}_{y'} + z'(t') \hat{\mathbf{e}}_{z'} \end{cases}, \quad \begin{cases} \hat{\mathbf{e}}_x = \hat{\mathbf{e}}_{x'} \\ \hat{\mathbf{e}}_y = \hat{\mathbf{e}}_{y'} \\ \hat{\mathbf{e}}_z = \hat{\mathbf{e}}_{z'} \end{cases}$$

$$t' = t \quad \Rightarrow \quad \mathbf{r}'(t') = \mathbf{r}'(t) = x'(t) \hat{\mathbf{e}}_x + y'(t) \hat{\mathbf{e}}_y + z'(t) \hat{\mathbf{e}}_z$$

$$\mathbf{v}'(t) = \mathbf{v}(t) - \mathbf{V}$$

$$\Rightarrow \quad \frac{d\mathbf{v}'(t)}{dt} = \frac{d\mathbf{v}(t)}{dt} - \frac{d\mathbf{V}}{dt} \quad \Rightarrow \quad \mathbf{a}'(t) = \mathbf{a}(t) \quad \blacksquare$$

$$\Rightarrow \quad \begin{cases} \mathbf{a}(t) \equiv \frac{d\mathbf{v}(t)}{dt} = \frac{d^2x(t)}{dt^2} \hat{\mathbf{e}}_x + \frac{d^2y(t)}{dt^2} \hat{\mathbf{e}}_y + \frac{d^2z(t)}{dt^2} \hat{\mathbf{e}}_z \\ \mathbf{a}'(t) \equiv \frac{d\mathbf{v}'(t)}{dt} = \frac{d^2x'(t)}{dt^2} \hat{\mathbf{e}}_{x'} + \frac{d^2y'(t)}{dt^2} \hat{\mathbf{e}}_{y'} + \frac{d^2z'(t)}{dt^2} \hat{\mathbf{e}}_{z'} \end{cases} \quad \blacksquare$$



# Equation of motion

- Second order ordinary differential equation:  $\mathbf{r}(0) = \mathbf{r}_0, \dot{\mathbf{r}}(0) = \mathbf{v}_0$

$$m\ddot{\mathbf{r}}(t) = \mathbf{F}(\mathbf{r}(t), \dot{\mathbf{r}}(t), t) \quad \rightarrow \quad \begin{cases} \mathbf{r}(t) = ? \\ \dot{\mathbf{r}}(t) = ?? \end{cases}$$

- Cartesian coordinates:

$$m\ddot{\mathbf{r}}(t) = \mathbf{F}(\mathbf{r}(t), \dot{\mathbf{r}}(t), t) \quad \Rightarrow \quad \begin{cases} m\ddot{x}(t) = F_x(\mathbf{r}(t), \dot{\mathbf{r}}(t), t) \\ m\ddot{y}(t) = F_y(\mathbf{r}(t), \dot{\mathbf{r}}(t), t) \\ m\ddot{z}(t) = F_z(\mathbf{r}(t), \dot{\mathbf{r}}(t), t) \end{cases}$$

- Polar coordinates:

$$m\ddot{\mathbf{r}}(t) = \mathbf{F}(\mathbf{r}(t), \dot{\mathbf{r}}(t), t) \quad \Rightarrow \quad \begin{cases} m [\ddot{\rho}(t) - \rho(t) \dot{\phi}^2(t)] = F_\rho(\mathbf{r}(t), \dot{\mathbf{r}}(t), t) \\ m [\rho(t) \ddot{\phi}(t) + 2\dot{\rho}(t) \dot{\phi}(t)] = F_\phi(\mathbf{r}(t), \dot{\mathbf{r}}(t), t) \end{cases}$$

# First order separable ordinary differential equation

- General form:

$$\frac{dy(x)}{dx} = f(x) g(y)$$

- Implicit **general solution**: existence of an *arbitrary* constant in the solution

$$\int \frac{1}{g(y)} dy = \int f(x) dx$$

# First order linear ordinary differential equation

- Standard form:  $a_1(x) \neq 0$

$$a_1(x) \frac{dy(x)}{dx} + a_0(x) y(x) = f(x)$$

- **Integrating factor**  $\mu(x)$ : integration constant is irrelevant

$$\mu(x) a_1(x) \frac{dy(x)}{dx} + \mu(x) a_0(x) y(x) \equiv \frac{d}{dx} [\mu(x) a_1(x) y(x)]$$

$$\Rightarrow \mu(x) = \frac{1}{a_1(x)} \exp \left[ \int^x \frac{a_0(\xi)}{a_1(\xi)} d\xi \right]$$

- General solution:  $c$  is an arbitrary integration constant

$$\frac{d}{dx} [\mu(x) a_1(x) y(x)] = \mu(x) f(x) \quad \Rightarrow \quad y(x) = \frac{1}{\mu(x) a_1(x)} \left[ \int^x \mu(\xi) f(\xi) d\xi + c \right]$$

$$a_1(x) \frac{dy(x)}{dx} + a_0(x) y(x) = f(x)$$

$$\mu(x) a_1(x) \frac{dy(x)}{dx} + \mu(x) a_0(x) y(x) \equiv \frac{d}{dx} [\mu(x) a_1(x) y(x)]$$

$$\Rightarrow \mu(x) a_0(x) = \mu(x) \frac{da_1(x)}{dx} + \frac{d\mu(x)}{dx} a_1(x)$$

$$\Rightarrow \frac{d\mu}{\mu(x)} = \frac{a_0(x)}{a_1(x)} dx - \frac{da_1}{a_1(x)}$$

$$\Rightarrow \ln \mu(x) = \int^x \frac{a_0(\xi)}{a_1(\xi)} d\xi - \ln a_1(x)$$

$$\Rightarrow \mu(x) = \frac{1}{a_1(x)} \exp \left[ \int^x \frac{a_0(\xi)}{a_1(\xi)} d\xi \right] \quad \blacksquare$$

## Special case: $F_x = F_x(t)$

- Solving for  $v_x(t)$ :  $v_x(0) = v_{x0}$

$$\begin{aligned} m\ddot{x}(t) = F_x(t) &\Rightarrow m \frac{dv_x(t)}{dt} = F_x(t) \Rightarrow m \int_{v'_x=v_{x0}}^{v_x} dv'_x = \int_{t'=0}^t F_x(t') dt' \\ &\Rightarrow v_x(t) = v_{x0} + \frac{1}{m} \int_{t'=0}^t F_x(t') dt' \end{aligned}$$

- Solving for  $x(t)$ :  $x(0) = x_0$

$$\begin{aligned} \frac{dx(t)}{dt} = v_x(t) &\Rightarrow \int_{x'=x_0}^x dx' = \int_{t'=0}^t v_x(t') dt' \\ &\Rightarrow x(t) = x_0 + v_{x0}t + \frac{1}{m} \int_{t'=0}^t \left[ \int_{t''=0}^{t'} F_x(t'') dt'' \right] dt' \end{aligned}$$

## Special case: $F_x = F_x(x)$

- Solving for  $v_x(x)$ :  $x = x(t) \leftrightarrow t = t(x)$

$$\begin{aligned} m\ddot{x}(t) = F_x(x) &\Rightarrow m \frac{dv_x(t)}{dt} = F_x(x) \Rightarrow m \frac{dv_x(x)}{dx} \frac{dx(t)}{dt} = F_x(x) \\ \Rightarrow mv_x(x) \frac{dv_x(x)}{dx} = F_x(x) &\Rightarrow m \int_{v'_x=v_{x0}}^{v_x} v'_x dv'_x = \int_{x'=x_0}^x F_x(x') dx' \\ \Rightarrow v_x^2(x) = v_{x0}^2 + \frac{2}{m} \int_{x'=x_0}^x F_x(x') dx' \end{aligned}$$

- Solving for  $x(t)$ :  $x = x(t) \leftrightarrow t = t(x)$

$$\begin{aligned} \frac{dx(t)}{dt} = v_x(x) &\Rightarrow \int_{x'=x_0}^x \frac{dx'}{v_x(x')} = \int_{t'=0}^t dt' \\ \Rightarrow t = \int_{x'=x_0}^x \frac{dx'}{v_x(x')} &\Rightarrow x(t) \end{aligned}$$

## Special case: $F_x = F_x(v_x)$

- Solving for  $v_x(t)$ :

$$\begin{aligned} m\ddot{x}(t) = F_x(v_x) &\Rightarrow m \frac{dv_x(t)}{dt} = F_x(v_x) \\ \Rightarrow m \int_{v'_x=v_{x0}}^{v_x} \frac{dv'_x}{F_x(v'_x)} = \int_{t'=0}^t dt' &\Rightarrow v_x(t) \Rightarrow x(t) \end{aligned}$$

- Solving for  $v_x(x)$ :

$$\begin{aligned} m\ddot{x}(t) = F_x(v_x) &\Rightarrow m \frac{dv_x(t)}{dt} = F_x(v_x) \Rightarrow m \frac{dv_x(x)}{dx} \frac{dx(t)}{dt} = F_x(v_x) \\ \Rightarrow mv_x(x) \frac{dv_x(x)}{dx} = F_x(v_x) &\Rightarrow m \int_{v'_x=v_{x0}}^{v_x} \frac{v'_x}{F_x(v'_x)} dv'_x = \int_{x'=x_0}^x dx' \\ &\Rightarrow v_x(x) \Rightarrow x(t) \end{aligned}$$