

PC3130

Quantum Mechanics II

Time-Independent Perturbation Theory  
Application – Fine Structure of the Hydrogen Atom

# Dirac Hamiltonian: Existence of Spin

The Pauli Hamiltonian contains the interaction of spin with an external magnetic field. For a particle in an external potential, or interacting with other particles, a spin-dependence can also be present via the exchange term. But the spin up and spin down Hamiltonians can be separated and solved independently.

$$H = H_{Pauli}$$

$$- \frac{p^4}{8m^3c^2} \quad \text{mass - velocity}$$

$$+ \frac{\hbar^2 q}{8m^2c^2} \nabla \cdot \nabla \phi(\mathbf{r}) \quad \text{Darwin}$$

$$- \frac{\hbar q}{4m^2c^2} \boldsymbol{\sigma} \cdot [\boldsymbol{\pi} \times \nabla \phi(\mathbf{r})] \quad \text{spin - orbit}$$

these terms are  
independent of spin

Only the spin-orbit term couples  
the spin up and spin down  
Hamiltonians, and also relates the  
spin to the lattice degrees of  
freedom. Without the spin-orbit  
term, we do not need two-  
component spinors.

Solutions are two-component spinors

# Fine Structure of Hydrogen

Kinetic energy term

potential energy term

$$\hat{H}^0 = \frac{\hat{p}^2}{2m_e} - \frac{e^2}{4\pi\epsilon_0} \frac{1}{r}$$

reduced mass of the electron-proton pair,  
to be more precise

Exact wavefunction:

$$\langle r, \theta, \phi | \tilde{\psi}_{n,l,m}^0 \rangle = R_{nl}(r) Y_l^m(\theta, \phi)$$

Exact eigenvalues:

$$E_{n,l,m}^0 = - \left[ \frac{m_e}{2\hbar^2} \left( \frac{e^2}{4\pi\epsilon_0} \right)^2 \right] \frac{1}{n^2}$$

highly degenerate because it is independent of  $(l, m)$  !!!

# Fine Structure of Hydrogen Atom

in terms of the **fine structure constant**  $\alpha \equiv \frac{e^2}{4\pi\epsilon_0\hbar c} \approx \frac{1}{137.036}$   $E_{n,l,m}^0 = - \left[ \frac{m_e}{2\hbar^2} \left( \frac{e^2}{4\pi\epsilon_0} \right)^2 \right] \frac{1}{n^2}$

Bohr formula

Eigenvalues from the Bohr formula

$$E_n = \alpha^2 (m_e c^2)$$

From Dirac's equation  
(without external fields)

Fine structure

$$E_n \left( 1 + \frac{\alpha^2}{n^2} \left( \frac{n}{j + 1/2} - \frac{3}{4} \right) \right)$$

depends on  $j$  only,  
i.e. for the same  $j$ , all  
 $m_j$  states have the  
same energy.

- mass-velocity term

$$E_n \frac{\alpha^2}{n^2} \left( \frac{n}{\ell + 1/2} - \frac{3}{4} \right)$$

- spin-orbit coupling term

$$E_n \frac{\alpha^2}{n^2} \frac{n \left[ \frac{3}{4} + \ell(\ell + 1) - j(j + 1) \right]}{2\ell(\ell + \frac{1}{2})(\ell + 1)} \quad \ell \neq 0$$

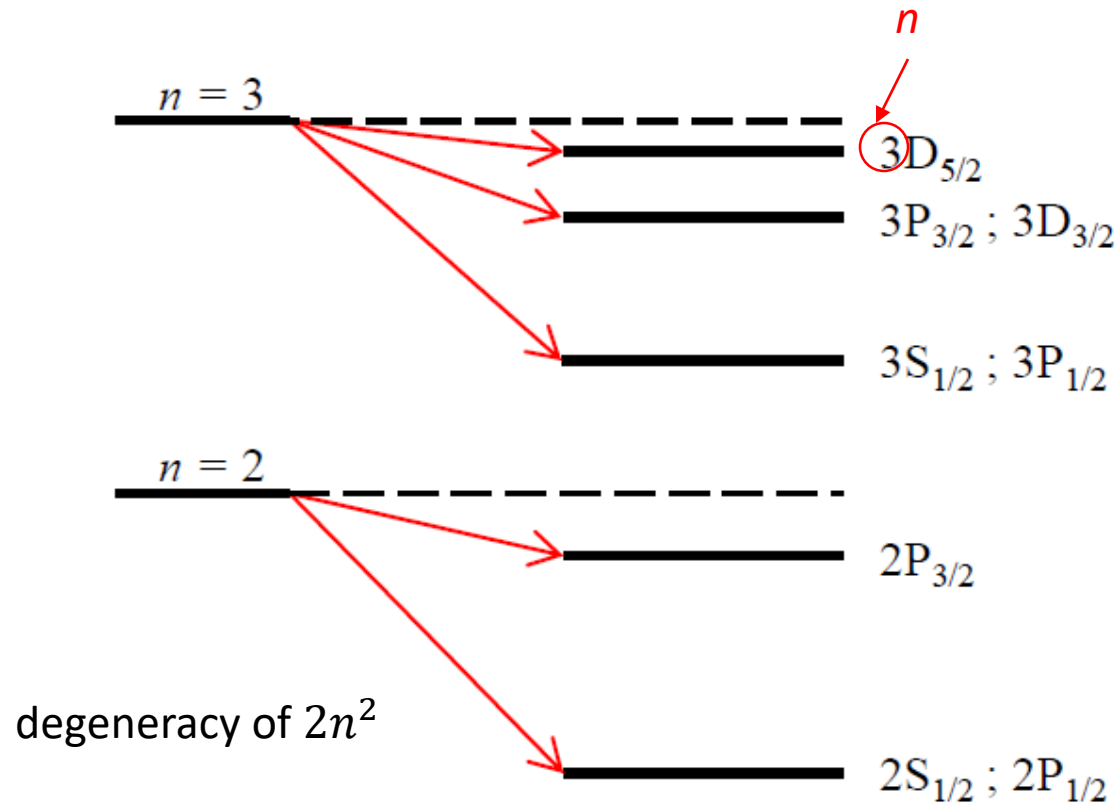
- Darwin term

$$-E_n \frac{\alpha^2}{n} \quad \text{Only for } \ell = 0$$

# Term symbols and Hydrogen fine structure

$nL_j$

For the single electron state,  $s = \frac{1}{2}$   
Omitted in the notation below.



degeneracy is partially lifted

states with the same  $j$  but different  $m_j$  or  $l$  have the same energy.

# Besides hydrogen atoms...

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these terms are  
independent of spin

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Only the spin-orbit term couples the spin up and spin down Hamiltonians, and also relates the spin to the lattice degrees of freedom. Without the spin-orbit term, we do not need two-component spinors.

Solutions are two-component spinors

# Spin-orbit term as a perturbation

Given an external potential and an approximation to the Hamiltonian, e.g. Hartree-Fock, how do we include the relativistic corrections?

$$H = H_{\text{Pauli}} - \frac{p^4}{8m^3c^2} \quad \text{mass - velocity} + \frac{\hbar^2 q}{8m^2c^2} \nabla \cdot \nabla \phi(\mathbf{r}) \quad \text{Darwin} - \frac{\hbar q}{4m^2c^2} \boldsymbol{\sigma} \cdot [\boldsymbol{\pi} \times \nabla \phi(\mathbf{r})] \quad \text{spin - orbit}$$

**scalar relativistic approximation**  
includes these terms only. No spinors are needed.

**The spin-orbit term couples the spin up and spin down Hamiltonians. It can be included as a perturbation to the scalar relativistic Hamiltonian.**

**Solutions are two-component spinors**