#### PC3261: Classical Mechanics II

#### Kenneth HONG Chong Ming

Office: S16-07-06 Email: phyhcmk@nus.edu.sg

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### **Lecture 3: Linear Momentum**

## Linear momentum of two-particle system

ullet Forces are *assumed* to obey principle of superposition of forces:  ${f f}_{12}$  is the force acting on  $m_1$  due to  $m_2$ 

$$\left\{ \begin{array}{l} \mathbf{F}_1(t) = \mathbf{F}_1^{\mathsf{ext}}(t) + \mathbf{f}_{12}(t) \\ \\ \mathbf{F}_2(t) = \mathbf{F}_2^{\mathsf{ext}}(t) + \mathbf{f}_{21}(t) \end{array} \right.$$

 Total linear momentum of the system: forces between particles are assumed to obey Newton's third law of motion

$$\mathbf{P}(t) \equiv \mathbf{p}_1(t) + \mathbf{p}_2(t) \quad \Rightarrow \quad \dot{\mathbf{P}}(t) = \mathbf{F}_1^{\mathsf{ext}}(t) + \mathbf{F}_2^{\mathsf{ext}}(t)$$

• Newton's second law: the time rate of change of total linear momentum of the two-particle system equals to the total *external* force acting upon it

$$\dot{\mathbf{P}}(t) = \mathbf{F}^{\rm ext}(t)\,, \qquad \mathbf{F}^{\rm ext}(t) \equiv \mathbf{F}_1^{\rm ext}(t) + \mathbf{F}_2^{\rm ext}(t)$$

### Linear momentum of multi-particle system

• Total force acting on the lpha-particle:  $\mathbf{f}_{lphaeta}$  is the force acting on  $m_{lpha}$  due to  $m_{eta}$ 

$$\mathbf{F}_{\alpha}(t) = \mathbf{F}_{\alpha}^{\mathsf{ext}}(t) + \sum_{\beta=1, \beta \neq \alpha}^{N} \mathbf{f}_{\alpha\beta}(t), \qquad \alpha = 1, 2, 3, \dots, N$$

• Total linear momentum of multi-particle system:

$$\mathbf{P}(t) \equiv \sum_{\alpha=1}^{N} \mathbf{p}_{\alpha}(t)$$

 Newton's second law: the time rate of change of total linear momentum of multi-particle system equals to the total external force acting upon it

$$\dot{\mathbf{P}}(t) = \mathbf{F}^{\mathrm{ext}}(t) \,, \qquad \mathbf{F}^{\mathrm{ext}}(t) \equiv \sum_{\alpha=1}^{N} \mathbf{F}_{\alpha}^{\mathrm{ext}}(t)$$

# Impuse-Momentum theorem

 Impulse-Momentum theorem (integral form of the Newton's second law): change of total linear momentum equals to the time integral of the total external force

$$\mathbf{F}^{\mathsf{ext}}(t) = \frac{\mathrm{d}\mathbf{P}(t)}{\mathrm{d}t} \longrightarrow \int_{t_1}^{t_2} \mathbf{F}^{\mathsf{ext}}(t) \, \mathrm{d}t = \mathbf{P}(t_2) - \mathbf{P}(t_1)$$

• Conservation law of linear momentum: if the total external force on a multiparticle system is zero, then the total linear momentum of the multi-particle is a constant

$$\dot{\mathbf{P}}(t) = \mathbf{F}^{\mathrm{ext}}(t) = \mathbf{0} \quad \Rightarrow \quad \mathbf{P}(t) = \mathrm{constant} \quad \Rightarrow \quad \mathbf{P}(t_1) = \mathbf{P}(t_2) \quad \forall \quad t_1, t_2$$

• The validity of the conservation law of linear momentum depends crucially on the *experimental* basis of the Newton's third law!

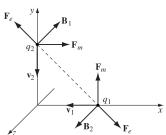
### A violation of Newton's third law???

- $\bullet$  Two point charges,  $q_1$  and  $q_2,$  are moving at uniform velocities  $\mathbf{v}_1$  and  $\mathbf{v}_2$  respectively
- Electric fields and forces:

$$\begin{cases} \mathbf{E}_{1}(\mathbf{r}_{2}) = \frac{q_{1}}{4\pi\epsilon_{0}} \frac{\mathbf{r}_{2} - \mathbf{r}_{1}}{|\mathbf{r}_{2} - \mathbf{r}_{1}|^{3}} \\ \mathbf{E}_{2}(\mathbf{r}_{1}) = \frac{q_{2}}{4\pi\epsilon_{0}} \frac{\mathbf{r}_{1} - \mathbf{r}_{2}}{|\mathbf{r}_{1} - \mathbf{r}_{2}|^{3}} \end{cases} \Rightarrow \begin{cases} \mathbf{F}_{e,12} = q_{1}\mathbf{E}_{2}(\mathbf{r}_{1}) \\ \mathbf{F}_{e,21} = q_{2}\mathbf{E}_{1}(\mathbf{r}_{2}) \end{cases}$$

• Electric forces obey Newton's third law

$$\mathbf{F}_{e,12} = -\mathbf{F}_{e,21}$$



### A violation of Newton's third law??? - cont'd

• Magnetic fields and forces:

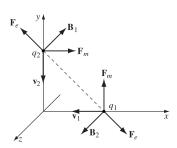
$$\begin{cases} \mathbf{B}_{1}(\mathbf{r}_{2}) = \frac{\mu_{0}q_{1}}{4\pi} \frac{\mathbf{v}_{1} \times (\mathbf{r}_{2} - \mathbf{r}_{1})}{\left|\mathbf{r}_{2} - \mathbf{r}_{1}\right|^{3}} \\ \mathbf{B}_{2}(\mathbf{r}_{1}) = \frac{\mu_{0}q_{2}}{4\pi} \frac{\mathbf{v}_{2} \times (\mathbf{r}_{1} - \mathbf{r}_{2})}{\left|\mathbf{r}_{1} - \mathbf{r}_{2}\right|^{3}} \end{cases} \Rightarrow \begin{cases} \mathbf{F}_{m,12} = q_{1}\mathbf{v}_{1} \times \mathbf{B}_{2}(\mathbf{r}_{1}) \\ \mathbf{F}_{m,21} = q_{2}\mathbf{v}_{2} \times \mathbf{B}_{1}(\mathbf{r}_{2}) \end{cases}$$

 Magnetic forces do not obey Newton's third law!

$$\mathbf{F}_{m,12} \neq -\mathbf{F}_{m,21}$$

• Electromagnetic linear momentum density: fields also possesses linear momentum!

$$\mathbf{g}_{\mathsf{EM}}(\mathbf{r}) = \epsilon_0 \, \mathbf{E}(\mathbf{r}) \times \mathbf{B}(\mathbf{r})$$



# System with variable mass

• Newton's second law with variable mass:

$$\mathbf{F}^{\mathsf{ext}}(t) = \frac{\mathrm{d}}{\mathrm{d}t} \left[ m(t) \, \mathbf{v}(t) \right] \quad \xrightarrow{???} \quad \mathbf{F}^{\mathsf{ext}}(t) = \frac{\mathrm{d}m(t)}{\mathrm{d}t} \, \mathbf{v}(t) + m \, \frac{\mathrm{d}\mathbf{v}(t)}{\mathrm{d}t}$$

• Galilean velocity transformation:  $\mathbf{v}'(t) = \mathbf{v}(t) - \mathbf{V}$ 

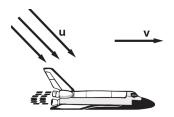
$$\frac{\mathrm{d}}{\mathrm{d}t} \left[ m(t)\mathbf{v}'(t) \right] = \mathbf{F}^{\mathsf{ext}}(t) \quad \Leftrightarrow \quad \frac{\mathrm{d}}{\mathrm{d}t} \left[ m(t)\mathbf{v}(t) \right] = \mathbf{F}^{\mathsf{ext}}(t)$$

 $\bullet$  There is no fundamental difficulty in handling any system with variable mass provided the same set of particles is included throughout the time interval  $t_1$  to  $t_2$ 

$$\int_{t_1}^{t_2} \mathbf{F}^{\mathsf{ext}}(t) \, \mathrm{d}t = \mathbf{P}(t_2) - \mathbf{P}(t_1)$$

### **Example: Spacecraft and dust particles**

• A spacecraft with mass M moves through space with constant velocity  ${\bf v}$ . The spacecraft encounters a stream of dust particles that embed themselves in the hull at rate dm/dt. The dust has velocity  ${\bf u}$  just before it hits.



**EXERCISE 3.1:** Find the external force necessary to keep the spacecraft moving uniformly.

### Newton's second law with variable mass

ullet A system with mass m(t) moves at velovity  ${f v}(t)$ . Particles are added to the system at a rate  ${
m d} m(t)/{
m d} t$ . These particles have velocity  ${f u}(t)$  just before entering the system.

Newton's second law:

$$\begin{split} \mathbf{F}^{\mathsf{ext}}(t) &= \frac{\mathrm{d}}{\mathrm{d}t} \left[ m(t) \, \mathbf{v}(t) \right] - \frac{\mathrm{d}m(t)}{\mathrm{d}t} \, \mathbf{u}(t) \\ \Rightarrow & m(t) \dot{\mathbf{v}}(t) = \mathbf{F}^{\mathsf{ext}}(t) + \dot{m}(t) \left[ \mathbf{u}(t) - \mathbf{v}(t) \right] \end{split}$$

• Galilean invariance is preserved:

$$\mathbf{F}^{\mathsf{ext}}(t) = \frac{\mathrm{d}}{\mathrm{d}t} \left[ m(t)\mathbf{v}(t) \right] - \dot{m}(t)\mathbf{u}(t) \quad \leftrightarrow \quad \mathbf{F}^{\mathsf{ext}}(t) = \frac{\mathrm{d}}{\mathrm{d}t} \left[ m(t)\mathbf{v}'(t) \right] - \dot{m}(t)\mathbf{u}'(t)$$

# Example: Rocket in a constant gravitational field

- A rocket is taking off from rest in a uniform gravitation field  $\mathbf{g}=-g\,\hat{\mathbf{e}}_z$ . The fuel is ejected at a constant rate  $\dot{m}(t)=-k$  at a constant exhaust speed u relative to the rocket.
- Linear momentum of the system:

$$\begin{cases} \mathbf{P}(t) = m(t)\mathbf{v}(t) \\ \mathbf{P}(t+dt) = [m(t) + dm] [\mathbf{v}(t) + d\mathbf{v}] + (-dm) [\mathbf{v}(t) + d\mathbf{v} + \mathbf{u}(t+dt)] \end{cases}$$

Newton's second law:

$$m(t)\dot{\mathbf{v}}(t) - \mathbf{u}(t)\dot{m}(t) = \mathbf{F}^{\mathsf{ext}}(t)$$

**EXERCISE 3.2:** Find  $\mathbf{v}(t)$  and  $\mathbf{r}(t)$  for the rocket in its subsequent motion given that the initial mass of the rocket is  $m_0$ .

### Center of mass

• Position vector of the **center of mass** of a multi-particle system:

$$\mathbf{R}_{\mathsf{CM}}(t) \equiv \frac{1}{M} \sum_{\alpha=1}^{N} m_{\alpha} \mathbf{r}_{\alpha}(t) \,, \qquad M \equiv \sum_{\alpha=1}^{N} m_{\alpha}$$

• Velocity of the center of mass: total linear momentum of the system is equal to the linear momentum of the center of mass as if it were a particle of mass M with velocity  $\mathbf{V}_{\mathsf{CM}}(t)$ 

$$\mathbf{V}_{\mathsf{CM}}(t) \equiv \dot{\mathbf{R}}_{\mathsf{CM}}(t) = \frac{1}{M} \sum_{\alpha=1}^{N} m_{\alpha} \dot{\mathbf{r}}_{\alpha}(t) \quad \Rightarrow \quad \mathbf{P}(t) = M \mathbf{V}_{\mathsf{CM}}(t)$$

 $\bullet$  Acceleration of the center of mass: center of mass moves exactly as if it were a single particle of mass M subjected to the total external force on the system

$$\mathbf{A}_{\mathsf{CM}}(t) \equiv \ddot{\mathbf{R}}_{\mathsf{CM}}(t) = \frac{1}{M} \sum_{\alpha=1}^{N} m_{\alpha} \ddot{\mathbf{r}}_{\alpha}(t) \quad \Rightarrow \quad \dot{\mathbf{P}}(t) = M \mathbf{A}_{\mathsf{CM}}(t)$$

## **Example: Projectile motion**

ullet A rigid object consists of two masses  $m_1$  and  $m_2$  separated by a light rod of length L. It is thrown into the air.

Center of mass:

$$\mathbf{R}_{CM}(t) = \frac{m_1 \mathbf{r}_1(t) + m_2 \mathbf{r}_2(t)}{m_1 + m_2}$$

• Equation of motion of the center of mass:

$$\mathbf{F}^{\text{ext}}(t) = (m_1 + m_2) \, \ddot{\mathbf{R}}_{\text{CM}}(t) \quad \Rightarrow \quad \ddot{\mathbf{R}}_{\text{CM}}(t) = \mathbf{g}$$

• The center of mass follows the parabolic trajectory of a single mass,  $m_1+m_2$ , in a uniform gravitational field (motions of  $m_1$  and  $m_2$  about the center of mass are to be analyzed separately)

# Center of mass of extended body

• Visualize mass element  ${\rm d} m$  of volume  ${\rm d} V$  located at position  ${\bf r}$  with mass density  $\rho({\bf r})$ :

$$\mathbf{R}_{\mathsf{CM}} = \frac{1}{M} \iiint_{V} \mathbf{r} \, \rho(\mathbf{r}) \, \mathrm{d}V$$

 $\bullet$  Center of mass of a uniform solid (upper) hemisphere: mass M and radius R

$$\mathbf{R}_{\mathsf{CM}} = \frac{3}{8} R \,\hat{\mathbf{e}}_z$$

**EXERCISE 3.3:** A thin non-uniform plates lies on the xy-plane with corners (0,0), (a,0), (0,b) and (a,b). Its surface mass density is  $\sigma(x,y) = \sigma_0 xy/ab$  where  $\sigma_0$  is a constant. Find its center of mass.

#### Center-of-mass frame

• Center-of-mass frame is a reference frame at which the center of mass remains at the origin:

$$\mathbf{r}'_{\alpha}(t) \equiv \mathbf{r}_{\alpha}(t) - \mathbf{R}_{\mathsf{CM}}(t) \quad \Rightarrow \quad \mathbf{R}'_{\mathsf{CM}}(t) = \mathbf{0}$$

• Velocity of the center of mass in the center-of-mass frame: center of mass is stationary in the center-of-mass frame

$$\mathbf{V}'_{\mathsf{CM}}(t) = \frac{1}{M} \sum_{\alpha=1}^{N} m_{\alpha} \dot{\mathbf{r}}'_{\alpha}(t) = \mathbf{0}$$

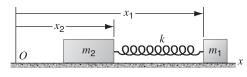
• Acceleration of the center of mass in the center-of-mass frame:

$$\mathbf{A}'_{\mathsf{CM}}(t) = \frac{1}{M} \sum_{\alpha=1}^{N} m_{\alpha} \ddot{\mathbf{r}}'_{\alpha}(t) = \mathbf{0}$$

# **Example: Two-body oscillations**

- ullet Two idential blocks 1 and 2 each of mass m slide without friction on a straight track. They are connected by a massless spring with unstretched length  $L_0$  and spring constant k. Initially, the system is at rest. At t=0, block 1 is hit sharply giving it an instanteneous velocity  $v_0$  to the right.
- Equations of motion in the center-of-mass frame:

$$\begin{cases} m\ddot{x}_1'(t) = -k \left[ x_1'(t) - x_2'(t) - L_0 \right] \\ m\ddot{x}_2'(t) = +k \left[ x_1'(t) - x_2'(t) - L_0 \right] \end{cases}$$



**EXERCISE 3.4:** Find the velocities of each block at later times with respect to the track.