

NATIONAL UNIVERSITY OF SINGAPORE

PC4245 Particle Physics
(Semester II: AY 2021-22)

Time Allowed: 2 Hours

INSTRUCTIONS TO STUDENTS

1. Write your Matric Number on the front cover page of each answer book.
2. This examination paper contains **4** questions and comprises **6** printed pages. Answer any **3** questions.
3. All questions carry equal marks.
4. Students should write the answers for each question on a new page.
5. This is a **CLOSED BOOK** examination.

1. (a) The kaon mesons are produced via strong interaction but decay by weak interaction. Describe briefly the Cronin-Fitch experiment on K^0 decay that evidenced CP violation.

Find the ratio of K_S (K short) and K_L (K long) in a beam of 10 GeV neutral kaons at a distance of 20 meters from where the beam is produced.

Note that the probability of getting $|K_L\rangle$ at time t is $P_L(t) = \frac{1}{2}e^{-\Gamma_L t/\hbar}$ and that of getting $|K_S\rangle$ at time t is $P_S(t) = \frac{1}{2}e^{-\Gamma_S t/\hbar}$, where $\Gamma_L = \frac{\hbar}{\tau_L}$, $\Gamma_S = \frac{\hbar}{\tau_S}$ and $\tau_L = 5.2 \times 10^{-8}$ sec, $\tau_S = 0.86 \times 10^{-10}$ sec.

(b) Consider the elastic scattering ($m_3 = m_1, m_4 = m_2$) in the lab frame (particle 2 at rest), $1+2 \rightarrow 3+4$.

Derive an expression of the differential cross section

$$\frac{d\sigma}{d\Omega} = \left(\frac{\hbar}{8\pi}\right)^2 \frac{p_3^2 S |M|^2}{m_2 |\tilde{p}_1| [E_1 + m_2 c^2] |\tilde{p}_3| - |\tilde{p}_1| E_3 \cos \theta}$$

in the usual notations...

Note: The following formula can be used

$$d\sigma = S |\mathcal{M}|^2 \frac{\hbar^2}{4} \left[(\underline{p}_1 \cdot \underline{p}_2)^2 - (m_1 m_2 c^2)^2 \right]^{-1/2} \frac{d^3 \tilde{p}_3}{(2\pi)^3 2 p_3^0} \times \frac{d^3 \tilde{p}_4}{(2\pi)^3 2 p_4^0} (2\pi)^4 \delta^{(4)}(\underline{p}_1 + \underline{p}_2 - \underline{p}_3 - \underline{p}_4)$$

and for the lab frame (particle 2 at rest)

$$\sqrt{(\underline{p}_1 \cdot \underline{p}_2)^2 - (m_1 m_2 c^2)^2} = m_2 |\tilde{p}_1| c$$

2. (a) The spin/flavor wave function for a baryon Λ with spin up is given by

$$\left| \Lambda; \frac{1}{2} \frac{1}{2} \right\rangle = \frac{1}{2\sqrt{3}} [u(\uparrow)d(\downarrow)s(\uparrow) - u(\downarrow)d(\uparrow)s(\uparrow) + 6 \text{ permutations}]$$

Calculate the magnetic dipole moment of Λ by considering only the first term of the wave function, namely,

$$\frac{1}{2\sqrt{3}} (u(\uparrow)d(\downarrow)s(\uparrow)).$$

Your answer should be expressed in terms of the quark magnetic dipole moments. In the usual notations, these are

$$\mu_u = \frac{2}{3} \frac{e\hbar}{2m_u c}, \quad \mu_d = -\frac{1}{3} \frac{e\hbar}{2m_d c}, \quad \mu_s = -\frac{1}{3} \frac{e\hbar}{2m_s c}.$$

Note: The magnetic dipole moment is related to spin by the relation

$$\mu_{\sim} = \frac{q}{mc} s_{\sim}.$$

The magnetic dipole moment of a baryon in the usual notations can be written as $\mu_B = \langle B \uparrow | (\mu_1 + \mu_2 + \mu_3) | B \uparrow \rangle$

(b) The mass of a baryon in terms of the masses of its constituent

quarks and the their spins, in the usual notations, is given by

$$M(\text{baryon}) = m_1 + m_2 + m_3 + A' \left[\frac{\vec{s}_1 \cdot \vec{s}_2}{m_1 m_2} + \frac{\vec{s}_1 \cdot \vec{s}_3}{m_1 m_3} + \frac{\vec{s}_2 \cdot \vec{s}_3}{m_2 m_3} \right]$$

where $m_u = m_d = 363 \text{ MeV}/c^2$, $m_s = 538 \text{ MeV}/c^2$, and

$$A' = \left(\frac{2m_u}{\hbar} \right)^2 50 \text{ MeV}/c^2.$$

✓(i) Prove that

$$\vec{s}_1 \cdot \vec{s}_2 + \vec{s}_1 \cdot \vec{s}_3 + \vec{s}_2 \cdot \vec{s}_3 = \begin{cases} -\frac{3}{4} \hbar^2, & \text{for } j = 1/2 \text{ (octet)} \\ \frac{3}{4} \hbar^2, & \text{for } j = 3/2 \text{ (decuplet)} \end{cases}$$

✓(ii) If the three quarks have the same mass m , show that for an octet baryon,

$$M(\text{baryon}) = 3m - \frac{3}{4} \frac{\hbar^2}{m^2} A'$$

and for a decuplet baryon

$$M(\text{baryon}) = 3m + \frac{3}{4} \frac{\hbar^2}{m^2} A'$$

✓(ii) If two of the three quarks have the same mass, say $m_1 = m_2$, show that for an octet baryon,

$$M(\text{baryon}) = 2m_1 + m_3 + A' \left[\frac{\tilde{s}_1 \cdot \tilde{s}_2}{m_1^2} - \frac{\tilde{s}_1 \cdot \tilde{s}_2 + \frac{3}{4} \hbar^2}{m_1 m_3} \right]$$

and for a decuplet baryon,

$$M(\text{baryon}) = 2m_1 + m_3 + A' \left[\frac{\tilde{s}_1 \cdot \tilde{s}_2}{m_1^2} - \frac{\tilde{s}_1 \cdot \tilde{s}_2 - \frac{3}{4} \hbar^2}{m_1 m_3} \right] .$$

Hence or otherwise, obtain the mass of the octet baryon Λ

$$M(\Lambda) = 2m_u + m_s - A' \frac{3}{4} \frac{\hbar^2}{m_u^2} .$$

Note: The baryon Λ is an isosinglet,

$$\Lambda = \frac{1}{\sqrt{12}} [2(USD - DSU) + UDS - SDU - DUS + SUD]$$

3. ✓(a) Show that the helicity operator $s(p) \equiv \tilde{\Sigma} \cdot \tilde{p}/|\tilde{p}|$ commutes with the Hamiltonian of a Dirac particle, $H = c\tilde{\alpha} \cdot \tilde{p} + \beta mc^2$. Here $\tilde{\Sigma}$ is the spin operator of the Dirac particle.

Explain qualitatively why the helicity of a particle is, in general, not an invariant.


$$\text{Note: } \tilde{\Sigma} = \begin{pmatrix} \tilde{\sigma} & 0 \\ 0 & \tilde{\sigma} \end{pmatrix}, \quad \tilde{\alpha} = \begin{pmatrix} 0 & \tilde{\sigma} \\ \tilde{\sigma} & 0 \end{pmatrix}, \quad \beta = \begin{pmatrix} I & 0 \\ 0 & -I \end{pmatrix}, \quad I = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

and $\tilde{\sigma}$ are the Pauli matrices $\sigma^3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$, $\sigma^1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$, $\sigma^2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$


(b) For a massive fermion, show that handedness is not a good quantum number. That is, show that γ^5 does not commute with H .

Note: $\gamma^5 = i\gamma^0\gamma^1\gamma^2\gamma^3 = \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix}$, $I = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$.


(c) Describe briefly two experimental evidences that suggest each quark flavor must come in three color varieties.

4.  Draw the lowest-order Feynman diagrams for the electron-electron scattering

$$e^- + e^- \rightarrow e^- + e^-.$$

 Using the Feynman rules for quantum electrodynamics, obtain the scattering amplitude \mathcal{M} for the above process.

Note: For vertex, $ig\gamma^\mu$; for propagators, $\frac{-ig^{\mu\nu}}{q^2}$, $\frac{i}{q_\mu\gamma^\mu - mc}$

 Consider the electron-electron scattering at very high energy so that the mass of the electron can be ignored (i.e., set $m = 0$).

Define the spin-averaged quantity $\langle |\mathcal{M}|^2 \rangle$.

The scattering amplitude \mathcal{M} can be written as $\mathcal{M} = \mathcal{M}_1 + \mathcal{M}_2$.

Using the Casimir trick, show that

$$(i) \quad \langle |\mathcal{M}_1|^2 \rangle = \frac{g^4}{4(p_1 - p_3)^4} \text{Tr}(\gamma^\mu \not{p}_1 \gamma^\nu \not{p}_3) \cdot \text{Tr}(\gamma_\mu \not{p}_2 \gamma_\nu \not{p}_4),$$

$$(ii) \quad \langle |\mathcal{M}_1 \mathcal{M}_2^*| \rangle = \frac{-g^4}{16(p_1 - p_3)(p_1 - p_4)} \text{Tr}(\gamma^\mu \not{p}_1 \gamma^\nu \not{p}_4 \gamma_\mu \not{p}_2 \gamma_\nu \not{p}_3).$$

Hence or otherwise obtain an expression of $\langle |\mathcal{M}|^2 \rangle$.

Note: (I) for massless particles the conservation of momentum

($\underline{p}_1 + \underline{p}_2 = \underline{p}_3 + \underline{p}_4$) implies that

$$\underline{p}_1 \cdot \underline{p}_2 = \underline{p}_3 \cdot \underline{p}_4, \quad \underline{p}_1 \cdot \underline{p}_3 = \underline{p}_2 \cdot \underline{p}_4, \quad \underline{p}_1 \cdot \underline{p}_4 = \underline{p}_2 \cdot \underline{p}_3.$$

$$(II) \sum_s \bar{u}^{(s)}(\underline{p}) u^{(s)}(\underline{p}) = \not{p} + mc. \quad \not{p} \equiv \gamma^\mu p_\mu$$

(OCH)

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