PC3261: Classical Mechanics II

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Lecture 6: d'Alembert's Principle

Constraints

- Positions of all particles of a mechanical system define the **configuration** of the system at any instant of time
- **Constraints** are limitations of a *kinematic* nature, such as *possible* positions or velocities of the particles, imposed on the mechanical system
- Given that $(\xi_1, \xi_2, \cdots, \xi_n)$ are *arbitrary* coordinates employed to specify the configuration of a mechanical system, a constraint is said to be **holonomic** if it can be expressed by a functional relation among the coordinates alone with a possible explicit time dependence

$$f(\xi_1, \xi_2, \cdots, \xi_n, t) = 0$$

 \bullet Example: a system of two particles connected by a rigid rod of length ℓ moving in the space

$$|\mathbf{r}_2 - \mathbf{r}_1|^2 - \ell^2 = 0 \quad \Rightarrow \quad (x_2 - x_1)^2 + (y_2 - y_1)^2 + (z_2 - z_1)^2 - \ell^2 = 0$$

Velocity-dependent constraints

• Velocity-dependent constraints are given by functional relations among the coordinates and velocities which leading to a set of differential equations

$$g(\xi_1,\cdots,\xi_n,\dot{\xi}_1,\cdots,\dot{\xi}_n,t)=0$$

- Velocity-dependent constraints restrict the possible displacements of the mechanical system but do not impose any limits on the possible configurations
- If the velocity-dependent constraints can be reduced to the functional relations among coordinates alone with possible explicit time dependence by integration (solving differential equation), then it is said to be holonomic
- Differential equations from the velocity-dependent constraints are generally *non-integrable* which prevents them being holonomic

Example: Cylinder rolls on a straight track

- Denote the position of the center of masss of the cylinder as x and the angle of rotation about the center of mass as ϕ
- ullet Condition for roll without slipping: R is the radius of the cylinder

$$\dot{x}(t) = R \,\dot{\phi}(t)$$

• Velocity-dependent constraint is integrable leading to a functional relation among coordinates alone

$$x(t) = R \phi(t) \implies x(t) - R \phi(t) = 0$$

Example: Upright disc rolls on a horizontal plane

- Denote (x,y) as the position of the disc's center of mass on the plane, θ be the angle of the plane of the disc made with the x-axis and ϕ be the angle of rotation of the disc about its symmetry axis
- ullet Condition for roll without slipping: R is the radius of the disc

$$\left\{ \begin{array}{l} \dot{x}(t) = R\,\dot{\phi}(t)\cos\theta(t) \\ \dot{y}(t) = R\,\dot{\phi}(t)\sin\theta(t) \end{array} \right.$$

- Velocity-dependent constraint is non-integrable (Frobenius integrability condition); the constraints represent restrictions on the velocities alone
- ullet The coordinates, $x,\ y,\ heta,\ \phi$, employed to specify the instantaneous configuration of the disc are functionally independent

Virtual displacements

- Virtual displacements are infinitesimal displacements that change a possible configuration to another possible configuration compatible with the constraints at a fixed instant of time
- ullet Virtual displacements of a system of N particles:

$$f(\mathbf{r}_1(t), \cdots, \mathbf{r}_N(t), t) = 0 \quad \Rightarrow \quad f(\mathbf{r}_1(t) + \delta \mathbf{r}_1, \cdots, \mathbf{r}_N(t) + \delta \mathbf{r}_N, t) = 0$$

 \bullet Virtual displacements $\delta {\bf r}_i$ coincide with the real displacements ${\rm d} {\bf r}_i$ if all constraints are time independent

EXERCISE 6.1: A particle is confined to a moving surface where the equation of the surface is given by $f(\mathbf{r},t)=0$. Show that the virtual displacement is tangent to the surface at the same time.

Virtual work

 Virtual work by the force: no integration is required as the the virtual displacement is infinitesimal

$$\delta W \equiv \mathbf{F} \cdot \delta \mathbf{r}$$

- The total virtual work of the constraint forces is zero in most physically interesting cases
- ullet Example: a system of two particles connected by a rigid rod of length ℓ moving in the space; ${f f}_1$ and ${f f}_2$ are constrained forces on the particles along the line connecting the particles

$$\delta W = \mathbf{f}_1 \cdot \delta \mathbf{r}_1 + \mathbf{f}_2 \cdot \delta \mathbf{r}_2 = 0$$

EXERCISE 6.2: Show that the total virtual work by the constrained forces on the two particles connected by a rigid rod moving in the space is zero.

Principle of virtual work

• Newton's second law: $\mathbf{F}_{\alpha}^{(\mathsf{A})}(t)$ and $\mathbf{F}_{\alpha}^{(\mathsf{C})}(t)$ are applied and constraint forces

$$\mathbf{F}_{\alpha}(t) = m_{\alpha}\ddot{\mathbf{r}}_{\alpha}(t), \qquad \mathbf{F}_{\alpha}(t) = \mathbf{F}_{\alpha}^{(\mathsf{A})}(t) + \mathbf{F}_{\alpha}^{(\mathsf{C})}(t)$$

• Static equilibrium: $\mathbf{F}_{\alpha}(t) = \mathbf{0}$

$$\sum_{\alpha} \mathbf{F}_{\alpha}(t) \cdot \delta \mathbf{r}_{\alpha} = 0 \quad \Rightarrow \quad \sum_{\alpha} \mathbf{F}_{\alpha}^{(\mathsf{A})}(t) \cdot \delta \mathbf{r}_{\alpha} + \sum_{\alpha} \mathbf{F}_{\alpha}^{(\mathsf{C})}(t) \cdot \delta \mathbf{r}_{\alpha} = 0$$

• **Principle of virtual work**: allowing one to express the equilibrium conditions for a constrained system in terms of the applied forces alone

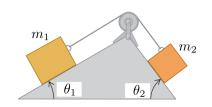
$$\sum_{\alpha} \mathbf{F}_{\alpha}^{(\mathsf{A})}(t) \cdot \delta \mathbf{r}_{\alpha} = 0$$

Example: Two masses on double inclined plane

• Two masses m_1 and m_2 are located each on a smooth double inclined plane with angles θ_1 and θ_2 respectively. The masses are connected by a massless and inextensible string running over a massless and frictionless pulley

Holonomic constraint:

$$\begin{cases} f_1(\mathbf{r}_1, \mathbf{r}_2) = x_1 + x_2 - \ell = 0 \\ f_2(\mathbf{r}_1, \mathbf{r}_2) = y_1 = 0 \\ f_3(\mathbf{r}_1, \mathbf{r}_2) = y_2 = 0 \end{cases}$$



Applied forces:

$$\mathbf{F}_{1}^{(\mathsf{A})}(t) = m_{1}g\sin\theta_{1}\,\hat{\mathbf{e}}_{x_{1}} - m_{1}g\cos\theta_{1}\,\hat{\mathbf{e}}_{y_{1}}\,,\,\,\mathbf{F}_{2}^{(\mathsf{A})}(t) = m_{2}g\sin\theta_{2}\,\hat{\mathbf{e}}_{x_{2}} - m_{2}g\cos\theta_{2}\,\hat{\mathbf{e}}_{y_{2}}$$

EXERCISE 6.3: Establish the condition for equilibrium from the principle of virtual work.

d'Alembert's principle

Newton's second law:

$$\mathbf{F}_{\alpha}(t) = m_{\alpha}\ddot{\mathbf{r}}_{\alpha}(t) \quad \Rightarrow \quad \mathbf{F}_{\alpha}(t) - m_{\alpha}\ddot{\mathbf{r}}_{\alpha}(t) = 0$$

• d'Alembert's principle: an extension of the principle of virtual work to mechanical systems in motion

$$\sum_{\alpha} \left[\mathbf{F}_{\alpha}^{(\mathsf{A})}(t) - m_{\alpha} \ddot{\mathbf{r}}_{\alpha}(t) \right] \cdot \delta \mathbf{r}_{\alpha} = 0$$

- d'Alembert's principle is a substantial leap forward with respect to the Newtonian approach as it excludes any reference to the constraint forces
- It is important to take note that the virtual displacements $\delta {\bf r}_{\alpha}$ are *not* independent as they have to be in harmony with the constraints

Example: Atwood's machine

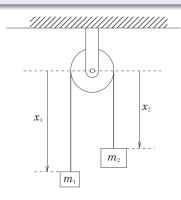
ullet Two masses m_1 and m_2 are suspended by an inextensible string which passes over a massless and frictionless pulley

Holonomic constraint:

$$f(\mathbf{r}_1, \mathbf{r}_2) = x_1 + x_2 - \ell = 0$$

• Applied forces:

$$\mathbf{F}_1^{(\mathsf{A})}(t) = m_1 g \,\hat{\mathbf{e}}_1 \,, \qquad \mathbf{F}_2^{(\mathsf{A})}(t) = m_2 g \,\hat{\mathbf{e}}_2$$

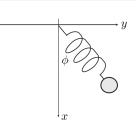


EXERCISE 6.4: Use d'Alembert's principle to find the accelerations of the masses $\ddot{x}_1(t)$ and $\ddot{x}_2(t)$.

Example: Pendulum with spring

- \bullet A point particle of mass m attached to a massless spring of original length ℓ_0 and spring constant k rotates about a frictionless pivot in a plane
- Applied forces:

$$\left\{ \begin{array}{l} \mathbf{F}_{\mathrm{gravity}}(t) = mg\cos\phi(t)\,\hat{\mathbf{e}}_{\rho} - mg\sin\phi(t)\,\hat{\mathbf{e}}_{\phi} \\ \\ \mathbf{F}_{\mathrm{spring}}(t) = -k\left[\rho(t) - \ell_{0}\right]\,\hat{\mathbf{e}}_{\rho} \end{array} \right.$$



EXERCISE 6.5: Use d'Alembert's principle to obtain equations of motion for $\rho(t)$ and $\phi(t)$.

Example: Spherical pendulum

 \bullet A particle of mass m is suspended by a massless wire of length r(t) to move on the surface of the spherical of radius r(t)

$$r(t) = a + b\cos\omega t$$
, $a > b > 0$

Holonomic constraint:

$$f(\mathbf{r},t) = r(t) - a - b\cos\omega t = 0$$

Applied force:

$$\mathbf{F}_{\text{gravity}}(t) = -mg\cos\theta(t)\,\hat{\mathbf{e}}_r + mg\sin\theta(t)\,\hat{\mathbf{e}}_\theta$$

EXERCISE 6.6: Use d'Alembert's principle to obtain equations of motion for $\theta(t)$ and $\phi(t)$.