

Moving trihedral

- Tangent and normal vectors: κ is called the **curvature**

$$\hat{\mathbf{e}}_T(s) \equiv \frac{d\mathbf{r}(s)}{ds} \quad \Rightarrow \quad \mathbf{v}(s) = v(s) \hat{\mathbf{e}}_T(s)$$

$$\hat{\mathbf{e}}_N(s) \equiv \frac{1}{\kappa(s)} \frac{d\hat{\mathbf{e}}_T(s)}{ds}$$

- Binormal vector: τ is called the **torsion**

$$\hat{\mathbf{e}}_B(s) \equiv \hat{\mathbf{e}}_T(s) \times \hat{\mathbf{e}}_N(s), \quad \frac{d\hat{\mathbf{e}}_B(s)}{ds} \equiv -\tau(s) \hat{\mathbf{e}}_N(s)$$

EXERCISE 1.2: Show that the acceleration of a particle moving along a trajectory $\mathbf{r}(t)$ is give by

$$\mathbf{a}(t) = \frac{dv(t)}{dt} \hat{\mathbf{e}}_T + \frac{v^2(t)}{\rho} \hat{\mathbf{e}}_N,$$

where $\rho \equiv 1/\kappa$ is its radius of curvature.

$$\mathbf{v}(t) = \frac{d\mathbf{r}(t)}{dt} = \frac{ds(t)}{dt} \frac{d\mathbf{r}(s)}{ds} = v(t) \hat{\mathbf{e}}_T \quad \blacksquare$$

$$\begin{aligned} \mathbf{a}(t) &= \frac{d\mathbf{v}(t)}{dt} = \frac{dv(t)}{dt} \hat{\mathbf{e}}_T + v(t) \frac{d\hat{\mathbf{e}}_T}{dt} \\ &= \frac{dv(t)}{dt} \hat{\mathbf{e}}_T + v(t) \frac{ds(t)}{dt} \frac{d\hat{\mathbf{e}}_T}{ds} \\ &= \frac{dv(t)}{dt} \hat{\mathbf{e}}_T + v^2(t) \kappa \hat{\mathbf{e}}_N \\ &= \frac{dv(t)}{dt} \hat{\mathbf{e}}_T + \frac{v^2(t)}{\rho} \hat{\mathbf{e}}_N \quad \blacksquare \end{aligned}$$

Example: Circular helix

- Position vector: a , b and ω are constants

$$\mathbf{r}(t) = a \cos \omega t \hat{\mathbf{e}}_x + a \sin \omega t \hat{\mathbf{e}}_y + b\omega t \hat{\mathbf{e}}_z$$

- Curvature and torsion: circular helix is the unique curve with non-zero constant curvature and torsion

$$\kappa(t) = \frac{a}{a^2 + b^2}, \quad \tau(t) = \frac{b}{a^2 + b^2}$$

EXERCISE 1.3: Find the tangent, normal and binormal vectors, as well as, curvature and torsion for the circular helix.

$$\mathbf{r}(t) = a \cos \omega t \hat{\mathbf{e}}_x + a \sin \omega t \hat{\mathbf{e}}_y + b \omega t \hat{\mathbf{e}}_z$$

$$\dot{\mathbf{r}}(t) = -a\omega \sin \omega t \hat{\mathbf{e}}_x + a\omega \cos \omega t \hat{\mathbf{e}}_y + b\omega \hat{\mathbf{e}}_z$$

$$s(t) = \int_0^t |\dot{\mathbf{r}}(t)| \, dt = \omega \sqrt{a^2 + b^2} t \quad \Rightarrow \quad \frac{ds(t)}{dt} = \omega \sqrt{a^2 + b^2}$$

$$\hat{\mathbf{e}}_T(t) = \frac{d\mathbf{r}(s)}{ds} = \frac{\frac{d\mathbf{r}(t)}{dt}}{\frac{ds(t)}{dt}} = \frac{\dot{\mathbf{r}}(t)}{\dot{s}(t)} = \frac{1}{\sqrt{a^2 + b^2}} (-a \sin \omega t \hat{\mathbf{e}}_x + a \cos \omega t \hat{\mathbf{e}}_y + b \hat{\mathbf{e}}_z) \quad \blacksquare$$

$$\hat{\mathbf{e}}_T(t) = \frac{1}{\sqrt{a^2 + b^2}} (-a \sin \omega t \hat{\mathbf{e}}_x + a \cos \omega t \hat{\mathbf{e}}_y + b \hat{\mathbf{e}}_z)$$

$$\frac{d\hat{\mathbf{e}}_T(t)}{dt} = \frac{a\omega}{\sqrt{a^2 + b^2}} (-\cos \omega t \hat{\mathbf{e}}_x - \sin \omega t \hat{\mathbf{e}}_y)$$

$$\frac{d\mathbf{e}_T(t)}{ds} = \frac{\frac{d\mathbf{e}_T(t)}{dt}}{\frac{ds(t)}{dt}} = \frac{a}{a^2 + b^2} (-\cos \omega t \hat{\mathbf{e}}_x - \sin \omega t \hat{\mathbf{e}}_y) \quad \Rightarrow \quad \left| \frac{d\hat{\mathbf{e}}_T(t)}{ds} \right| = \frac{a}{a^2 + b^2}$$

$$\hat{\mathbf{e}}_N(t) = \frac{1}{\kappa(t)} \frac{d\hat{\mathbf{e}}_T(t)}{ds} \quad \Rightarrow \quad \kappa(t) = \left| \frac{d\hat{\mathbf{e}}_T(t)}{ds} \right| = \frac{a}{a^2 + b^2} \quad \blacksquare$$

$$\hat{\mathbf{e}}_N(t) = \frac{1}{\kappa(t)} \frac{d\hat{\mathbf{e}}_T(t)}{ds} = -\cos \omega t \hat{\mathbf{e}}_x - \sin \omega t \hat{\mathbf{e}}_y \quad \blacksquare$$

$$\hat{\mathbf{e}}_T(t) = \frac{1}{\sqrt{a^2 + b^2}} (-a \sin \omega t \hat{\mathbf{e}}_x + a \cos \omega t \hat{\mathbf{e}}_y + b \hat{\mathbf{e}}_z), \quad \hat{\mathbf{e}}_N(t) = -\cos \omega t \hat{\mathbf{e}}_x - \sin \omega t \hat{\mathbf{e}}_y$$

$$\hat{\mathbf{e}}_B(t) = \hat{\mathbf{e}}_T(t) \times \hat{\mathbf{e}}_N(t) = \frac{1}{\sqrt{a^2 + b^2}} (b \sin \omega t \hat{\mathbf{e}}_x - b \cos \omega t \hat{\mathbf{e}}_y + a \hat{\mathbf{e}}_z) \quad \blacksquare$$

$$\frac{d\hat{\mathbf{e}}_B(t)}{dt} = \frac{b\omega}{\sqrt{a^2 + b^2}} (\cos \omega t \hat{\mathbf{e}}_x + \sin \omega t \hat{\mathbf{e}}_y)$$

$$\frac{d\hat{\mathbf{e}}_B(t)}{ds} = \frac{\frac{d\hat{\mathbf{e}}_B(t)}{dt}}{\frac{ds(t)}{dt}} = \frac{b}{a^2 + b^2} (\cos \omega t \hat{\mathbf{e}}_x + \sin \omega t \hat{\mathbf{e}}_y)$$

$$\frac{d\hat{\mathbf{e}}_B(t)}{ds} = -\tau(t) \hat{\mathbf{e}}_N(t) \quad \Rightarrow \quad \tau(t) = -\hat{\mathbf{e}}_N(t) \cdot \frac{d\hat{\mathbf{e}}_B(t)}{ds} = \frac{b}{a^2 + b^2} \quad \blacksquare$$

$$\hat{\mathbf{e}}_N(t) = -\cos \omega t \hat{\mathbf{e}}_x - \sin \omega t \hat{\mathbf{e}}_y, \quad \hat{\mathbf{e}}_B(t) = \frac{1}{\sqrt{a^2 + b^2}} (b \sin \omega t \hat{\mathbf{e}}_x - b \cos \omega t \hat{\mathbf{e}}_y + a \hat{\mathbf{e}}_z)$$

$$\frac{d\hat{\mathbf{e}}_N(t)}{dt} = \omega (\sin \omega t \hat{\mathbf{e}}_x - \cos \omega t \hat{\mathbf{e}}_y)$$

$$\frac{d\hat{\mathbf{e}}_N(t)}{ds} = \frac{\frac{d\hat{\mathbf{e}}_N(t)}{dt}}{\frac{ds(t)}{dt}} = \frac{1}{\sqrt{a^2 + b^2}} (\sin \omega t \hat{\mathbf{e}}_x - \cos \omega t \hat{\mathbf{e}}_y)$$

$$\hat{\mathbf{e}}_N(s) \cdot \hat{\mathbf{e}}_B(s) = 0 \quad \Rightarrow \quad \hat{\mathbf{e}}_N(s) \cdot \frac{d\hat{\mathbf{e}}_B(s)}{ds} + \frac{d\hat{\mathbf{e}}_N(s)}{ds} \cdot \hat{\mathbf{e}}_B(s) = 0$$

$$\Rightarrow \quad -\tau(s) \hat{\mathbf{e}}_N(s) \cdot \hat{\mathbf{e}}_N(s) + \frac{d\hat{\mathbf{e}}_N(s)}{ds} \cdot \hat{\mathbf{e}}_B(s) = 0 \quad \Rightarrow \quad \tau(s) = \hat{\mathbf{e}}_B(s) \cdot \frac{d\hat{\mathbf{e}}_N(s)}{ds}$$

$$\tau(t) = \hat{\mathbf{e}}_B(t) \cdot \frac{d\hat{\mathbf{e}}_N(t)}{ds} = \frac{b}{a^2 + b^2} \quad \blacksquare$$

2D polar coordinate system

- Polar coordinates: $(u_1, u_2) = (\rho, \phi)$

ρ : distance from the origin, $0 \leq \rho < \infty$

ϕ : azimuthal angle from $+x$ -axis, $0 \leq \phi < 2\pi$

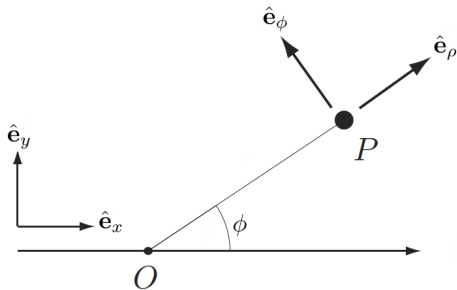
- Coordinate transformation between polar and Cartesian coordinates:

$$\begin{cases} x = \rho \cos \phi \\ y = \rho \sin \phi \end{cases} \Leftrightarrow \begin{cases} \rho = \sqrt{x^2 + y^2} \\ \phi = \tan^{-1} \left(\frac{y}{x} \right) \end{cases}$$

- Unit basis vectors $(\hat{e}_\rho, \hat{e}_\phi)$ are *not* constant!

EXERCISE 1.4: Establish the relationship between unit basis vectors $(\hat{e}_\rho, \hat{e}_\phi)$ of the polar coordinate system and the unit basis vectors (\hat{e}_x, \hat{e}_y) of the Cartesian coordinate system.

$$\begin{cases} \hat{\mathbf{e}}_\rho = \cos \phi \hat{\mathbf{e}}_x + \sin \phi \hat{\mathbf{e}}_y \\ \hat{\mathbf{e}}_\phi = -\sin \phi \hat{\mathbf{e}}_x + \cos \phi \hat{\mathbf{e}}_y \end{cases} \quad \blacksquare$$



$$\begin{aligned}
\begin{pmatrix} \hat{\mathbf{e}}_\rho \\ \hat{\mathbf{e}}_\phi \end{pmatrix} &= \begin{pmatrix} \cos \phi & \sin \phi \\ -\sin \phi & \cos \phi \end{pmatrix} \begin{pmatrix} \hat{\mathbf{e}}_x \\ \hat{\mathbf{e}}_y \end{pmatrix} \\
\Rightarrow \begin{pmatrix} \hat{\mathbf{e}}_x \\ \hat{\mathbf{e}}_y \end{pmatrix} &= \begin{pmatrix} \cos \phi & \sin \phi \\ -\sin \phi & \cos \phi \end{pmatrix}^{-1} \begin{pmatrix} \hat{\mathbf{e}}_\rho \\ \hat{\mathbf{e}}_\phi \end{pmatrix} \\
&= \begin{pmatrix} \cos \phi & -\sin \phi \\ \sin \phi & \cos \phi \end{pmatrix} \begin{pmatrix} \hat{\mathbf{e}}_\rho \\ \hat{\mathbf{e}}_\phi \end{pmatrix} \\
\Rightarrow \begin{cases} \hat{\mathbf{e}}_x = \cos \phi \hat{\mathbf{e}}_\rho - \sin \phi \hat{\mathbf{e}}_\phi \\ \hat{\mathbf{e}}_y = \sin \phi \hat{\mathbf{e}}_\rho + \cos \phi \hat{\mathbf{e}}_\phi \end{cases} \quad \blacksquare
\end{aligned}$$

Kinematics in 2D polar coordinates

- Position vector:

$$\mathbf{r}(t) = \rho(t) \hat{\mathbf{e}}_\rho$$

- Velocity:

$$\mathbf{v}(t) = \dot{\rho}(t) \hat{\mathbf{e}}_\rho + \rho(t) \dot{\phi}(t) \hat{\mathbf{e}}_\phi$$

- Acceleration:

$$\mathbf{a}(t) = [\ddot{\rho}(t) - \rho(t) \dot{\phi}^2(t)] \hat{\mathbf{e}}_\rho + [\rho(t) \ddot{\phi}(t) + 2\dot{\rho}(t) \dot{\phi}(t)] \hat{\mathbf{e}}_\phi$$

EXERCISE 1.5: Express the velocity and acceleration vectors in 2D polar coordinates.

$$\left\{ \begin{array}{l} x = \rho \cos \phi \\ y = \rho \sin \phi \end{array} \right., \quad \left\{ \begin{array}{l} \hat{\mathbf{e}}_\rho = \cos \phi(t) \hat{\mathbf{e}}_x + \sin \phi(t) \hat{\mathbf{e}}_y \\ \hat{\mathbf{e}}_\phi = -\sin \phi(t) \hat{\mathbf{e}}_x + \cos \phi(t) \hat{\mathbf{e}}_y \end{array} \right.$$

$$\mathbf{r}(t) = x(t) \hat{\mathbf{e}}_x + y(t) \hat{\mathbf{e}}_y = r_\rho \hat{\mathbf{e}}_\rho + r_\phi \hat{\mathbf{e}}_\phi$$

$$\left\{ \begin{array}{l} r_\rho = \hat{\mathbf{e}}_\rho \cdot \mathbf{r}(t) = x(t) \cos \phi(t) + y(t) \sin \phi(t) = \rho(t) \\ r_\phi = \hat{\mathbf{e}}_\phi \cdot \mathbf{r}(t) = -x(t) \sin \phi(t) + y(t) \cos \phi(t) = 0 \end{array} \right.$$

$$\Rightarrow \quad \mathbf{r}(t) = \rho(t) \hat{\mathbf{e}}_\rho \quad \blacksquare$$

$$\begin{cases} \hat{\mathbf{e}}_\rho = \cos \phi(t) \hat{\mathbf{e}}_x + \sin \phi(t) \hat{\mathbf{e}}_y \\ \hat{\mathbf{e}}_\phi = -\sin \phi(t) \hat{\mathbf{e}}_x + \cos \phi(t) \hat{\mathbf{e}}_y \end{cases}$$

$$\Rightarrow \begin{cases} \frac{d\hat{\mathbf{e}}_\rho}{dt} = -\dot{\phi}(t) \sin \phi(t) \hat{\mathbf{e}}_x + \dot{\phi}(t) \cos \phi(t) \hat{\mathbf{e}}_y = \dot{\phi}(t) \hat{\mathbf{e}}_\phi \\ \frac{d\hat{\mathbf{e}}_\phi}{dt} = -\dot{\phi}(t) \cos \phi(t) \hat{\mathbf{e}}_x - \dot{\phi}(t) \sin \phi(t) \hat{\mathbf{e}}_y = -\dot{\phi}(t) \hat{\mathbf{e}}_\rho \end{cases}$$

$$\begin{aligned} \mathbf{v}(t) &= \frac{d\mathbf{r}(t)}{dt} = \frac{d}{dt} [\rho(t) \hat{\mathbf{e}}_\rho] \\ &= \dot{\rho}(t) \hat{\mathbf{e}}_\rho + \rho(t) \dot{\phi}(t) \hat{\mathbf{e}}_\phi \quad \blacksquare \end{aligned}$$

$$\mathbf{v}(t) = \dot{\rho}(t) \hat{\mathbf{e}}_\rho + \rho(t) \dot{\phi}(t) \hat{\mathbf{e}}_\phi$$

$$\begin{cases} \frac{d\hat{\mathbf{e}}_\rho}{dt} = \dot{\phi}(t) \hat{\mathbf{e}}_\phi \\ \frac{d\hat{\mathbf{e}}_\phi}{dt} = -\dot{\phi}(t) \hat{\mathbf{e}}_\rho \end{cases}$$

$$\begin{aligned} \mathbf{a}(t) &= \frac{d\mathbf{v}(t)}{dt} = \frac{d}{dt} [\dot{\rho}(t) \hat{\mathbf{e}}_\rho + \rho(t) \dot{\phi}(t) \hat{\mathbf{e}}_\phi] \\ &= [\ddot{\rho}(t) - \rho(t) \dot{\phi}^2(t)] \hat{\mathbf{e}}_\rho + [\rho(t) \ddot{\phi}(t) + 2\dot{\rho}(t) \dot{\phi}(t)] \hat{\mathbf{e}}_\phi \quad \blacksquare \end{aligned}$$

Cylindrical coordinate system

- Cylindrical coordinates: $(u_1, u_2, u_3) = (\rho, \phi, z)$

ρ : polar distance from the z axis, $0 \leq \rho < \infty$

ϕ : azimuthal angle from the x axis on the xy -plane, $0 \leq \phi < 2\pi$

z : coordinate along the z axis, $-\infty < z < \infty$

- Coordinate transformation between cylindrical and Cartesian coordinates:

$$\begin{cases} x = \rho \cos \phi \\ y = \rho \sin \phi \\ z = z \end{cases} \Leftrightarrow \begin{cases} \rho = \sqrt{x^2 + y^2} \\ \phi = \tan^{-1}(y/x) \\ z = z \end{cases}$$

- Velocity and acceleration:

$$\begin{cases} \mathbf{v}(t) = \dot{\rho}(t) \hat{\mathbf{e}}_\rho + \rho(t) \dot{\phi}(t) \hat{\mathbf{e}}_\phi + \dot{z}(t) \hat{\mathbf{e}}_z \\ \mathbf{a}(t) = [\ddot{\rho}(t) - \rho(t) \dot{\phi}^2(t)] \hat{\mathbf{e}}_\rho + [\rho(t) \ddot{\phi}(t) + 2\dot{\rho}(t) \dot{\phi}(t)] \hat{\mathbf{e}}_\phi + \ddot{z}(t) \hat{\mathbf{e}}_z \end{cases}$$

Spherical coordinate system

- Spherical coordinates: $(u_1, u_2, u_3) = (r, \theta, \phi)$

r : radial distance from the origin, $0 \leq r < \infty$

θ : polar angle from the z axis, $0 \leq \theta \leq \pi$

ϕ : azimuthal angle from the x axis on the xy -plane, $0 \leq \phi < 2\pi$

- Coordinate transformation between spherical and Cartesian coordinates:

$$\begin{cases} x = r \sin \theta \cos \phi \\ y = r \sin \theta \sin \phi \\ z = r \cos \theta \end{cases} \Leftrightarrow \begin{cases} r = \sqrt{x^2 + y^2 + z^2} \\ \theta = \tan^{-1} \left(\sqrt{x^2 + y^2} / z \right) \\ \phi = \tan^{-1} (y/x) \end{cases}$$

EXERCISE 1.6: Express the spherical unit basis vectors $(\hat{\mathbf{e}}_r, \hat{\mathbf{e}}_\theta, \hat{\mathbf{e}}_\phi)$ in terms of Cartesian unit basis vectors $(\hat{\mathbf{e}}_x, \hat{\mathbf{e}}_y, \hat{\mathbf{e}}_z)$.

$$\mathbf{r} = x \hat{\mathbf{e}}_x + y \hat{\mathbf{e}}_y + z \hat{\mathbf{e}}_z = r \sin \theta \cos \phi \hat{\mathbf{e}}_x + r \sin \theta \sin \phi \hat{\mathbf{e}}_y + r \cos \theta \hat{\mathbf{e}}_z$$

$$\Rightarrow \left\{ \begin{array}{l} \frac{\partial \mathbf{r}}{\partial r} = \sin \theta \cos \phi \hat{\mathbf{e}}_x + \sin \theta \sin \phi \hat{\mathbf{e}}_y + \cos \theta \hat{\mathbf{e}}_z \\ \frac{\partial \mathbf{r}}{\partial \theta} = r \cos \theta \cos \phi \hat{\mathbf{e}}_x + r \cos \theta \sin \phi \hat{\mathbf{e}}_y - r \sin \theta \hat{\mathbf{e}}_z \\ \frac{\partial \mathbf{r}}{\partial \phi} = -r \sin \theta \sin \phi \hat{\mathbf{e}}_x + r \sin \theta \cos \phi \hat{\mathbf{e}}_y \end{array} \right.$$

$$\Rightarrow \left\{ \begin{array}{l} \hat{\mathbf{e}}_r \equiv \frac{\frac{\partial \mathbf{r}}{\partial r}}{\left| \frac{\partial \mathbf{r}}{\partial r} \right|} = \sin \theta \cos \phi \hat{\mathbf{e}}_x + \sin \theta \sin \phi \hat{\mathbf{e}}_y + \cos \theta \hat{\mathbf{e}}_z \\ \hat{\mathbf{e}}_\theta \equiv \frac{\frac{\partial \mathbf{r}}{\partial \theta}}{\left| \frac{\partial \mathbf{r}}{\partial \theta} \right|} = \cos \theta \cos \phi \hat{\mathbf{e}}_x + \cos \theta \sin \phi \hat{\mathbf{e}}_y - \sin \theta \hat{\mathbf{e}}_z \\ \hat{\mathbf{e}}_\phi \equiv \frac{\frac{\partial \mathbf{r}}{\partial \phi}}{\left| \frac{\partial \mathbf{r}}{\partial \phi} \right|} = -\sin \phi \hat{\mathbf{e}}_x + \cos \phi \hat{\mathbf{e}}_y \end{array} \right. \quad \blacksquare$$

$$\begin{cases} \hat{\mathbf{e}}_r = \sin \theta \cos \phi \hat{\mathbf{e}}_x + \sin \theta \sin \phi \hat{\mathbf{e}}_y + \cos \theta \hat{\mathbf{e}}_z \\ \hat{\mathbf{e}}_\theta = \cos \theta \cos \phi \hat{\mathbf{e}}_x + \cos \theta \sin \phi \hat{\mathbf{e}}_y - \sin \theta \hat{\mathbf{e}}_z \\ \hat{\mathbf{e}}_\phi = -\sin \phi \hat{\mathbf{e}}_x + \cos \phi \hat{\mathbf{e}}_z \end{cases}$$

$$\begin{aligned} \hat{\mathbf{e}}_r \cdot (\hat{\mathbf{e}}_\theta \times \hat{\mathbf{e}}_\phi) &= \begin{vmatrix} \sin \theta \cos \phi & \sin \theta \sin \phi & \cos \theta \\ \cos \theta \cos \phi & \cos \theta \sin \phi & -\sin \theta \\ -\sin \phi & \cos \phi & 0 \end{vmatrix} \\ &= -\sin \phi \begin{vmatrix} \sin \theta \sin \phi & \cos \theta \\ \cos \theta \sin \phi & -\sin \theta \end{vmatrix} - \cos \phi \begin{vmatrix} \sin \theta \cos \phi & \cos \theta \\ \cos \theta \cos \phi & -\sin \theta \end{vmatrix} \\ &= -\sin \phi (-\sin^2 \theta \sin \phi - \cos^2 \theta \sin \phi) - \cos \phi (-\sin^2 \theta \cos \phi - \cos^2 \theta \cos \phi) \\ &= 1 \quad \blacksquare \end{aligned}$$

Kinematics in spherical coordinates

- Position vector:

$$\mathbf{r}(t) = r(t) \hat{\mathbf{e}}_r$$

- Velocity vector:

$$\mathbf{v}(t) = \dot{r}(t) \hat{\mathbf{e}}_r + r(t) \dot{\theta}(t) \hat{\mathbf{e}}_\theta + r(t) \dot{\phi}(t) \sin \theta(t) \hat{\mathbf{e}}_\phi$$

- Acceleration vector:

$$\begin{aligned} \mathbf{a}(t) = & \left[\ddot{r}(t) - r(t) \dot{\phi}^2(t) \sin^2 \theta(t) - r(t) \dot{\theta}^2(t) \right] \hat{\mathbf{e}}_r \\ & + \left[r(t) \ddot{\theta}(t) + 2\dot{r}(t) \dot{\theta}(t) - r(t) \dot{\phi}^2(t) \sin \theta(t) \cos \theta(t) \right] \hat{\mathbf{e}}_\theta \\ & + \left[r(t) \ddot{\phi}(t) \sin \theta(t) + 2\dot{r}(t) \dot{\phi}(t) \sin \theta(t) + 2r(t) \dot{\theta}(t) \dot{\phi}(t) \cos \theta(t) \right] \hat{\mathbf{e}}_\phi \end{aligned}$$

$$\hat{\mathbf{e}}_r = \sin \theta(t) \cos \phi(t) \hat{\mathbf{e}}_x + \sin \theta(t) \sin \phi(t) \hat{\mathbf{e}}_y + \cos \theta(t) \hat{\mathbf{e}}_z$$

$$\begin{aligned} \frac{d\hat{\mathbf{e}}_r}{dt} &= \frac{\partial \hat{\mathbf{e}}_r}{\partial \theta} \dot{\theta} + \frac{\partial \hat{\mathbf{e}}_r}{\partial \phi} \dot{\phi} \\ &= (\cos \theta \cos \phi \hat{\mathbf{e}}_x + \cos \theta \sin \phi \hat{\mathbf{e}}_y - \sin \theta \hat{\mathbf{e}}_z) \dot{\theta} + (-\sin \theta \sin \phi \hat{\mathbf{e}}_x + \sin \theta \cos \phi \hat{\mathbf{e}}_y) \dot{\phi} \\ &= \dot{\theta} \hat{\mathbf{e}}_\theta + \sin \theta \dot{\phi} \hat{\mathbf{e}}_\phi \quad \blacksquare \end{aligned}$$

$$\begin{aligned} \mathbf{v}(t) &\equiv \frac{d\mathbf{r}(t)}{dt} = \frac{d}{dt} [r(t) \hat{\mathbf{e}}_r] \\ &= \dot{r}(t) \hat{\mathbf{e}}_r + r(t) \frac{d\hat{\mathbf{e}}_r}{dt} \\ &= \dot{r}(t) \hat{\mathbf{e}}_r + r(t) \dot{\theta}(t) \hat{\mathbf{e}}_\theta + r(t) \dot{\phi}(t) \sin \theta(t) \hat{\mathbf{e}}_\phi \quad \blacksquare \end{aligned}$$

PC3261: Classical Mechanics II

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Semester I, 2023/24

Latest update: August 18, 2023 2:11pm



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Lecture 2: Newton's Laws of Motion

Newton's first law and inertia

- **Newton's first law:** a particle remains at rest or in uniform motion unless acted upon a force
- **Inertia** is the *resistance* of any particle to any change in its velocity and the quantitative measure of inertia is **mass**
- A mathematical description of the motion of a particle requires the choice of a **frame of reference** – a set of coordinates in space that can be used to specify the position, velocity and acceleration of the particle at any given instant of time
- A frame of reference at which Newton's first law is valid is called an **inertial frame of reference**

Newton's second law

- **Linear momentum** of a particle is defined as the product of its mass and velocity

$$\mathbf{p}(t) \equiv m\mathbf{v}(t)$$

- **Newton's second law:** a particle acted upon a force moves in such a manner that the time rate of change of linear momentum equals the force

$$\mathbf{F}(t) = \frac{d\mathbf{p}(t)}{dt}$$

- Both Newton's first and second laws remain exactly true in special relativity with a *suitably* redefinition of linear momentum

Newton's third law

- **Newton's third law:** if two particles exert forces on each other, these forces are equal in magnitude and opposite in direction
- **Central forces** are the forces acting along the line connecting two particles
- Velocity-dependent forces are non-central and Newton's third law *may* not apply
- Newton's third law is not valid in special relativity as the concept of absolute time is abandoned

Galilean relativity

- Two inertial frames, \mathcal{O} and \mathcal{O}' , are oriented such that their spatial coordinate axes are parallel, their spatial origins are coincided when $t = t' = 0$ and \mathcal{O}' moves at *uniform velocity* \mathbf{V} with respect to \mathcal{O}

- Galilean boost:**

$$\begin{cases} t' = t \\ \mathbf{r}'(t) = \mathbf{r}(t) - \mathbf{V}t \end{cases}$$

- Galilean velocity transformation:

$$\mathbf{v}'(t) = \mathbf{v}(t) - \mathbf{V}$$

- Newton's laws are **Galilean invariance**

$$\begin{cases} \mathbf{r}(t) = x(t) \hat{\mathbf{e}}_x + y(t) \hat{\mathbf{e}}_y + z(t) \hat{\mathbf{e}}_z \\ \mathbf{r}'(t') = x'(t') \hat{\mathbf{e}}_{x'} + y'(t') \hat{\mathbf{e}}_{y'} + z'(t') \hat{\mathbf{e}}_{z'} \end{cases}, \quad \begin{cases} \hat{\mathbf{e}}_x = \hat{\mathbf{e}}_{x'} \\ \hat{\mathbf{e}}_y = \hat{\mathbf{e}}_{y'} \\ \hat{\mathbf{e}}_z = \hat{\mathbf{e}}_{z'} \end{cases}$$

$$t' = t \quad \Rightarrow \quad \mathbf{r}'(t') = \mathbf{r}'(t) = x'(t) \hat{\mathbf{e}}_x + y'(t) \hat{\mathbf{e}}_y + z'(t) \hat{\mathbf{e}}_z$$

$$\mathbf{v}'(t') \equiv \frac{d\mathbf{r}'(t')}{dt'} = \frac{d\mathbf{r}'(t)}{dt} = \frac{dx'(t)}{dt} \hat{\mathbf{e}}_x + \frac{dy'(t)}{dt} \hat{\mathbf{e}}_y + \frac{dz'(t)}{dt} \hat{\mathbf{e}}_z \equiv \mathbf{v}'(t)$$

$$\mathbf{r}'(t) = \mathbf{r}(t) - \mathbf{V}t \quad \Rightarrow \quad \frac{d\mathbf{r}'(t)}{dt} = \frac{d\mathbf{r}(t)}{dt} - \frac{d}{dt}(\mathbf{V}t) \quad \Rightarrow \quad \mathbf{v}'(t) = \mathbf{v}(t) - \mathbf{V} \quad \blacksquare$$

$$\Rightarrow \quad \begin{cases} \mathbf{v}(t) \equiv \frac{d\mathbf{r}(t)}{dt} = \frac{dx(t)}{dt} \hat{\mathbf{e}}_x + \frac{dy(t)}{dt} \hat{\mathbf{e}}_y + \frac{dz(t)}{dt} \hat{\mathbf{e}}_z \\ \mathbf{v}'(t) \equiv \frac{d\mathbf{r}'(t)}{dt} = \frac{dx'(t)}{dt} \hat{\mathbf{e}}_{x'} + \frac{dy'(t)}{dt} \hat{\mathbf{e}}_{y'} + \frac{dz'(t)}{dt} \hat{\mathbf{e}}_{z'} \end{cases} \quad \blacksquare$$

$$\begin{cases} \mathbf{r}(t) = x(t) \hat{\mathbf{e}}_x + y(t) \hat{\mathbf{e}}_y + z(t) \hat{\mathbf{e}}_z \\ \mathbf{r}'(t') = x'(t') \hat{\mathbf{e}}_{x'} + y'(t') \hat{\mathbf{e}}_{y'} + z'(t') \hat{\mathbf{e}}_{z'} \end{cases}, \quad \begin{cases} \hat{\mathbf{e}}_x = \hat{\mathbf{e}}_{x'} \\ \hat{\mathbf{e}}_y = \hat{\mathbf{e}}_{y'} \\ \hat{\mathbf{e}}_z = \hat{\mathbf{e}}_{z'} \end{cases}$$

$$t' = t \quad \Rightarrow \quad \mathbf{r}'(t') = \mathbf{r}'(t) = x'(t) \hat{\mathbf{e}}_x + y'(t) \hat{\mathbf{e}}_y + z'(t) \hat{\mathbf{e}}_z$$

$$\mathbf{v}'(t) = \mathbf{v}(t) - \mathbf{V}$$

$$\Rightarrow \quad \frac{d\mathbf{v}'(t)}{dt} = \frac{d\mathbf{v}(t)}{dt} - \frac{d\mathbf{V}}{dt} \quad \Rightarrow \quad \mathbf{a}'(t) = \mathbf{a}(t) \quad \blacksquare$$

$$\Rightarrow \quad \begin{cases} \mathbf{a}(t) \equiv \frac{d\mathbf{v}(t)}{dt} = \frac{dx^2(t)}{dt^2} \hat{\mathbf{e}}_x + \frac{d^2y(t)}{dt^2} \hat{\mathbf{e}}_y + \frac{d^2z(t)}{dt^2} \hat{\mathbf{e}}_z \\ \mathbf{a}'(t) \equiv \frac{d\mathbf{v}'(t)}{dt} = \frac{d^2x'(t)}{dt^2} \hat{\mathbf{e}}_{x'} + \frac{d^2y'(t)}{dt^2} \hat{\mathbf{e}}_{y'} + \frac{d^2z'(t)}{dt^2} \hat{\mathbf{e}}_{z'} \end{cases} \quad \blacksquare$$