

PC3261: Classical Mechanics II

Kenneth HONG Chong Ming

Office: S16-07-06

Email: phyhcmk@nus.edu.sg

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Department of Physics
Faculty of Science

Lecture 3: Linear Momentum

Linear momentum of two-particle system

- Forces are *assumed* to obey principle of superposition of forces: \mathbf{f}_{12} is the force acting on m_1 due to m_2

$$\begin{cases} \mathbf{F}_1(t) = \mathbf{F}_1^{\text{ext}}(t) + \mathbf{f}_{12}(t) \\ \mathbf{F}_2(t) = \mathbf{F}_2^{\text{ext}}(t) + \mathbf{f}_{21}(t) \end{cases}$$

- Total linear momentum of the system: forces between particles are assumed to obey Newton's third law of motion

$$\mathbf{P}(t) \equiv \mathbf{p}_1(t) + \mathbf{p}_2(t) \quad \Rightarrow \quad \dot{\mathbf{P}}(t) = \mathbf{F}_1^{\text{ext}}(t) + \mathbf{F}_2^{\text{ext}}(t)$$

- Newton's second law: the time rate of change of total linear momentum of the two-particle system equals to the total *external* force acting upon it

$$\dot{\mathbf{P}}(t) = \mathbf{F}^{\text{ext}}(t), \quad \mathbf{F}^{\text{ext}}(t) \equiv \mathbf{F}_1^{\text{ext}}(t) + \mathbf{F}_2^{\text{ext}}(t)$$

Linear momentum of multi-particle system

- Total force acting on the α -particle: $\mathbf{f}_{\alpha\beta}$ is the force acting on m_α due to m_β

$$\mathbf{F}_\alpha(t) = \mathbf{F}_\alpha^{\text{ext}}(t) + \sum_{\beta=1, \beta \neq \alpha}^N \mathbf{f}_{\alpha\beta}(t), \quad \alpha = 1, 2, 3, \dots, N$$

- Total linear momentum of multi-particle system:

$$\mathbf{P}(t) \equiv \sum_{\alpha=1}^N \mathbf{p}_\alpha(t)$$

- Newton's second law: the time rate of change of total linear momentum of multi-particle system equals to the total *external* force acting upon it

$$\dot{\mathbf{P}}(t) = \mathbf{F}^{\text{ext}}(t), \quad \mathbf{F}^{\text{ext}}(t) \equiv \sum_{\alpha=1}^N \mathbf{F}_\alpha^{\text{ext}}(t)$$

$$\mathbf{F}_\alpha(t) = \mathbf{F}_\alpha^{\text{ext}}(t) + \sum_{\beta=1, \beta \neq \alpha}^N \mathbf{f}_{\alpha\beta}(t) \quad \Rightarrow \quad \dot{\mathbf{p}}_\alpha(t) = \mathbf{F}_\alpha^{\text{ext}}(t) + \sum_{\beta=1, \beta \neq \alpha}^N \mathbf{f}_{\alpha\beta}(t)$$

$$\mathbf{P}(t) \equiv \sum_{\alpha=1}^N \mathbf{p}_\alpha(t) \quad \Rightarrow \quad \dot{\mathbf{P}}(t) \equiv \sum_{\alpha=1}^N \dot{\mathbf{p}}_\alpha(t)$$

$$\begin{aligned} \dot{\mathbf{P}}(t) &= \sum_{\alpha=1}^N \mathbf{F}_\alpha^{\text{ext}}(t) + \sum_{\alpha=1}^N \sum_{\beta=1, \beta \neq \alpha}^N \mathbf{f}_{\alpha\beta}(t) \\ &= \sum_{\alpha=1}^N \mathbf{F}_\alpha^{\text{ext}}(t) + \sum_{\alpha=1}^N \sum_{\beta > \alpha}^N [\mathbf{f}_{\alpha\beta}(t) + \mathbf{f}_{\beta\alpha}(t)] \\ &= \sum_{\alpha=1}^N \mathbf{F}_\alpha^{\text{ext}}(t) \quad \blacksquare \end{aligned}$$

Impulse-Momentum theorem

- **Impulse-Momentum theorem** (integral form of the Newton's second law): change of total linear momentum equals to the time integral of the total external force

$$\mathbf{F}^{\text{ext}}(t) = \frac{d\mathbf{P}(t)}{dt} \quad \rightarrow \quad \int_{t_1}^{t_2} \mathbf{F}^{\text{ext}}(t) dt = \mathbf{P}(t_2) - \mathbf{P}(t_1)$$

- **Conservation law of linear momentum:** if the total external force on a multi-particle system is zero, then the total linear momentum of the multi-particle is a constant

$$\dot{\mathbf{P}}(t) = \mathbf{F}^{\text{ext}}(t) = \mathbf{0} \quad \Rightarrow \quad \mathbf{P}(t) = \text{constant} \quad \Rightarrow \quad \mathbf{P}(t_1) = \mathbf{P}(t_2) \quad \forall \quad t_1, t_2$$

- The validity of the conservation law of linear momentum depends crucially on the *experimental* basis of the Newton's third law!

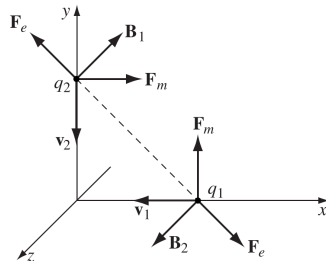
A violation of Newton's third law???

- Two point charges, q_1 and q_2 , are moving at uniform velocities \mathbf{v}_1 and \mathbf{v}_2 respectively
- Electric fields and forces:

$$\left\{ \begin{array}{l} \mathbf{E}_1(\mathbf{r}_2) = \frac{q_1}{4\pi\epsilon_0} \frac{\mathbf{r}_2 - \mathbf{r}_1}{|\mathbf{r}_2 - \mathbf{r}_1|^3} \\ \mathbf{E}_2(\mathbf{r}_1) = \frac{q_2}{4\pi\epsilon_0} \frac{\mathbf{r}_1 - \mathbf{r}_2}{|\mathbf{r}_1 - \mathbf{r}_2|^3} \end{array} \right. \Rightarrow \left\{ \begin{array}{l} \mathbf{F}_{e,12} = q_1 \mathbf{E}_2(\mathbf{r}_1) \\ \mathbf{F}_{e,21} = q_2 \mathbf{E}_1(\mathbf{r}_2) \end{array} \right.$$

- Electric forces obey Newton's third law

$$\mathbf{F}_{e,12} = -\mathbf{F}_{e,21}$$



A violation of Newton's third law??? – cont'd

- Magnetic fields and forces:

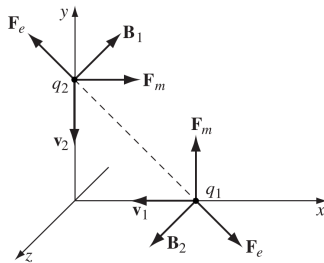
$$\begin{cases} \mathbf{B}_1(\mathbf{r}_2) = \frac{\mu_0 q_1}{4\pi} \frac{\mathbf{v}_1 \times (\mathbf{r}_2 - \mathbf{r}_1)}{|\mathbf{r}_2 - \mathbf{r}_1|^3} \\ \mathbf{B}_2(\mathbf{r}_1) = \frac{\mu_0 q_2}{4\pi} \frac{\mathbf{v}_2 \times (\mathbf{r}_1 - \mathbf{r}_2)}{|\mathbf{r}_1 - \mathbf{r}_2|^3} \end{cases} \Rightarrow \begin{cases} \mathbf{F}_{m,12} = q_1 \mathbf{v}_1 \times \mathbf{B}_2(\mathbf{r}_1) \\ \mathbf{F}_{m,21} = q_2 \mathbf{v}_2 \times \mathbf{B}_1(\mathbf{r}_2) \end{cases}$$

- Magnetic forces do not obey Newton's third law!

$$\mathbf{F}_{m,12} \neq -\mathbf{F}_{m,21}$$

- Electromagnetic linear momentum density: fields also possess linear momentum!

$$\mathbf{g}_{\text{EM}}(\mathbf{r}) = \epsilon_0 \mathbf{E}(\mathbf{r}) \times \mathbf{B}(\mathbf{r})$$



System with variable mass

- Newton's second law with variable mass:

$$\mathbf{F}^{\text{ext}}(t) = \frac{d}{dt} [m(t) \mathbf{v}(t)] \quad \xrightarrow{???} \quad \mathbf{F}^{\text{ext}}(t) = \frac{dm(t)}{dt} \mathbf{v}(t) + m \frac{d\mathbf{v}(t)}{dt}$$

- Galilean velocity transformation: $\mathbf{v}'(t) = \mathbf{v}(t) - \mathbf{V}$

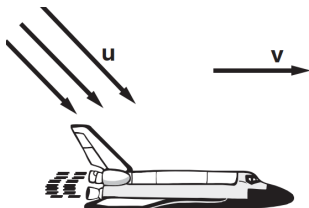
$$\frac{d}{dt} [m(t) \mathbf{v}'(t)] = \mathbf{F}^{\text{ext}}(t) \quad \nleftrightarrow \quad \frac{d}{dt} [m(t) \mathbf{v}(t)] = \mathbf{F}^{\text{ext}}(t)$$

- There is *no* fundamental difficulty in handling any system with variable mass provided the same set of particles is included *throughout* the time interval t_1 to t_2

$$\int_{t_1}^{t_2} \mathbf{F}^{\text{ext}}(t) dt = \mathbf{P}(t_2) - \mathbf{P}(t_1)$$

Example: Spacecraft and dust particles

- A spacecraft with mass M moves through space with constant velocity \mathbf{v} . The spacecraft encounters a stream of dust particles that embed themselves in the hull at rate dm/dt . The dust has velocity \mathbf{u} just before it hits.



EXERCISE 3.1: Find the external force necessary to keep the spacecraft moving uniformly.

$$\begin{cases} \mathbf{P}(t) = M(t) \mathbf{v} + (\Delta m) \mathbf{u} \\ \mathbf{P}(t + \Delta t) = M(t) \mathbf{v} + \Delta m \mathbf{v} \end{cases}$$

$$\Rightarrow \Delta \mathbf{P} = \mathbf{P}(t + \Delta t) - \mathbf{P}(t) = (\mathbf{v} - \mathbf{u}) \Delta m$$

$$\mathbf{F}(t) = \frac{d\mathbf{P}(t)}{dt} = \lim_{\Delta t \rightarrow 0} \frac{\Delta \mathbf{P}}{\Delta t} = (\mathbf{v} - \mathbf{u}) \frac{dm(t)}{dt} \quad \blacksquare$$

Newton's second law with variable mass

- A system with mass $m(t)$ moves at velocity $\mathbf{v}(t)$. Particles are added to the system at a rate $dm(t)/dt$. These particles have velocity $\mathbf{u}(t)$ just before entering the system.
- Newton's second law:

$$\mathbf{F}^{\text{ext}}(t) = \frac{d}{dt} [m(t) \mathbf{v}(t)] - \frac{dm(t)}{dt} \mathbf{u}(t)$$
$$\Rightarrow m(t) \dot{\mathbf{v}}(t) = \mathbf{F}^{\text{ext}}(t) + \dot{m}(t) [\mathbf{u}(t) - \mathbf{v}(t)]$$

- Galilean invariance is preserved:

$$\mathbf{F}^{\text{ext}}(t) = \frac{d}{dt} [m(t) \mathbf{v}(t)] - \dot{m}(t) \mathbf{u}(t) \quad \leftrightarrow \quad \mathbf{F}^{\text{ext}}(t) = \frac{d}{dt} [m(t) \mathbf{v}'(t)] - \dot{m}(t) \mathbf{u}'(t)$$

$$m(t)\dot{\mathbf{v}}(t) = \mathbf{F}(t) + \dot{m}(t) [\mathbf{u}(t) - \mathbf{v}(t)]$$

$$\begin{cases} \mathbf{P}(t) = m(t)\mathbf{v}(t) + \Delta m \mathbf{u}(t) \\ \mathbf{P}(t + \Delta t) = [m(t) + \Delta m] [\mathbf{v}(t) + \Delta \mathbf{v}] \end{cases}$$

$$\Delta \mathbf{P} \equiv \mathbf{P}(t + \Delta t) - \mathbf{P}(t) = m(t)\Delta \mathbf{v} + \Delta m \mathbf{v}(t) + \Delta m \Delta \mathbf{v} - \Delta m \mathbf{u}(t)$$

$$\begin{aligned} \frac{d\mathbf{P}(t)}{dt} &\equiv \lim_{\Delta t \rightarrow 0} \frac{\Delta \mathbf{P}}{\Delta t} \\ &= m(t) \lim_{\Delta t \rightarrow 0} \frac{\Delta \mathbf{v}}{\Delta t} + \left(\lim_{\Delta t \rightarrow 0} \frac{\Delta m}{\Delta t} \right) \mathbf{v}(t) + \left(\lim_{\Delta t \rightarrow 0} \Delta m \right) \left(\lim_{\Delta t \rightarrow 0} \frac{\Delta \mathbf{v}}{\Delta t} \right) - \left(\lim_{\Delta t \rightarrow 0} \frac{\Delta m}{\Delta t} \right) \mathbf{u}(t) \\ &= m(t)\dot{\mathbf{v}}(t) + \dot{m}(t)\mathbf{v}(t) - \dot{m}(t)\mathbf{u}(t) \quad \blacksquare \end{aligned}$$

$$\mathbf{F}^{\text{ext}}(t) = \frac{d\mathbf{P}(t)}{dt} \quad \Rightarrow \quad m(t)\dot{\mathbf{v}}(t) = \mathbf{F}(t) + \dot{m}(t) [\mathbf{u}(t) - \mathbf{v}(t)] \quad \blacksquare$$

$$\mathbf{v}'(t) = \mathbf{v}(t) - \mathbf{V} \quad \Rightarrow \quad \dot{\mathbf{v}}'(t) = \dot{\mathbf{v}}(t)$$

$$m(t)\dot{\mathbf{v}}'(t) = \mathbf{F}(t) + \dot{m}(t) [\mathbf{u}'(t) - \mathbf{v}'(t)]$$

$$\Rightarrow m(t)\dot{\mathbf{v}}(t) = \mathbf{F}(t) + \dot{m}(t) \{[\mathbf{u}(t) - \mathbf{V}] - [\mathbf{v}(t) - \mathbf{V}]\}$$

$$\Rightarrow m(t)\dot{\mathbf{v}}(t) = \mathbf{F}(t) + \dot{m}(t) [\mathbf{u}(t) - \mathbf{v}(t)] \quad \blacksquare$$

Example: Rocket in a constant gravitational field

- A rocket is taking off from rest in a uniform gravitation field $\mathbf{g} = -g \hat{\mathbf{e}}_z$. The fuel is ejected at a constant rate $\dot{m}(t) = -k$ at a constant exhaust speed u relative to the rocket.
- Linear momentum of the system:

$$\begin{cases} \mathbf{P}(t) = m(t)\mathbf{v}(t) \\ \mathbf{P}(t + \Delta t) = [m(t) + \Delta m] [\mathbf{v}(t) + \Delta \mathbf{v}] + (-\Delta m) [\mathbf{v}(t) + \Delta \mathbf{v} + \mathbf{u}(t + \Delta t)] \end{cases}$$

- Newton's second law:

$$m(t) \dot{\mathbf{v}}(t) - \mathbf{u}(t) \dot{m}(t) = \mathbf{F}^{\text{ext}}(t)$$

EXERCISE 3.2: Find $\mathbf{v}(t)$ and $\mathbf{r}(t)$ for the rocket in its subsequent motion given that the initial mass of the rocket is m_0 .

$$\mathbf{g} = -g \hat{\mathbf{e}}_z, \quad \mathbf{u} = -u \hat{\mathbf{e}}_z, \quad \dot{m}(t) = -k \quad \Rightarrow \quad m(t) = m_0 - kt$$

$$m(t) \dot{\mathbf{v}}(t) - \mathbf{u}(t) \dot{m}(t) = \mathbf{F}^{\text{ext}}(t)$$

$$\Rightarrow \quad m(t) \frac{dv_z(t)}{dt} + \dot{m}(t) u = -m(t)g$$

$$\Rightarrow \quad \frac{dv_z(t)}{dt} = -g - u \frac{\dot{m}(t)}{m(t)}$$

$$\Rightarrow \quad v_z(t) = -gt - u \ln \frac{m(t)}{m_0}$$

$$\Rightarrow \quad v_z(t) = -gt - u \ln \left(\frac{m_0 - kt}{m_0} \right) \quad \blacksquare$$

$$\frac{dz(t)}{dt} = v_z(t) \quad \Rightarrow \quad z(t) = -\frac{1}{2}gt^2 - u \int_0^t \ln \left(\frac{m_0 - kt'}{m_0} \right) dt'$$

$$\Rightarrow \quad z(t) = -\frac{1}{2}gt^2 - ut + \frac{u}{k} (m_0 - kt) \ln \left(\frac{m_0 - kt}{m_0} \right) \quad \blacksquare$$

Center of mass

- Position vector of the **center of mass** of a multi-particle system:

$$\mathbf{R}_{\text{CM}}(t) \equiv \frac{1}{M} \sum_{\alpha=1}^N m_{\alpha} \mathbf{r}_{\alpha}(t), \quad M \equiv \sum_{\alpha=1}^N m_{\alpha}$$

- Velocity of the center of mass: total linear momentum of the system is equal to the linear momentum of the center of mass *as if* it were a particle of mass M with velocity $\mathbf{V}_{\text{CM}}(t)$

$$\mathbf{V}_{\text{CM}}(t) \equiv \dot{\mathbf{R}}_{\text{CM}}(t) = \frac{1}{M} \sum_{\alpha=1}^N m_{\alpha} \dot{\mathbf{r}}_{\alpha}(t) \quad \Rightarrow \quad \mathbf{P}(t) = M \mathbf{V}_{\text{CM}}(t)$$

- Acceleration of the center of mass: center of mass moves exactly *as if* it were a single particle of mass M subjected to the total external force on the system

$$\mathbf{A}_{\text{CM}}(t) \equiv \ddot{\mathbf{R}}_{\text{CM}}(t) = \frac{1}{M} \sum_{\alpha=1}^N m_{\alpha} \ddot{\mathbf{r}}_{\alpha}(t) \quad \Rightarrow \quad \dot{\mathbf{P}}(t) = M \mathbf{A}_{\text{CM}}(t)$$

$$\mathbf{R}_{\text{CM}}(t) = \frac{1}{M} \sum_{\alpha=1}^N m_{\alpha} \mathbf{r}_{\alpha}(t), \quad \mathbf{V}_{\text{CM}}(t) = \frac{1}{M} \sum_{\alpha=1}^N m_{\alpha} \dot{\mathbf{r}}_{\alpha}(t)$$

$$\mathbf{P}(t) = \sum_{\alpha=1}^N m_{\alpha} \dot{\mathbf{r}}_{\alpha}(t) = M \mathbf{V}_{\text{CM}}(t) \quad \blacksquare$$

Example: Projectile motion

- A rigid object consists of two masses m_1 and m_2 separated by a light rod of length L . It is thrown into the air.

- Center of mass:

$$\mathbf{R}_{\text{CM}}(t) = \frac{m_1 \mathbf{r}_1(t) + m_2 \mathbf{r}_2(t)}{m_1 + m_2}$$

- Equation of motion of the center of mass:

$$\mathbf{F}^{\text{ext}}(t) = (m_1 + m_2) \ddot{\mathbf{R}}_{\text{CM}}(t) \quad \Rightarrow \quad \ddot{\mathbf{R}}_{\text{CM}}(t) = \mathbf{g}$$

- The center of mass follows the parabolic trajectory of a single mass, $m_1 + m_2$, in a uniform gravitational field (motions of m_1 and m_2 about the center of mass are to be analyzed separately)

Center of mass of extended body

- Visualize mass element dm of volume dV located at position $\mathbf{r}(t)$ with mass density $\rho(\mathbf{r})$:

$$\mathbf{R}_{\text{CM}}(t) = \frac{1}{M} \iiint_V \mathbf{r}(t) \rho(\mathbf{r}) dV$$

- Center of mass of a uniform solid (upper) hemisphere: mass M and radius R

$$\mathbf{R}_{\text{CM}} = \frac{3}{8}R \hat{\mathbf{e}}_z$$

EXERCISE 3.3: A thin non-uniform plates lies on the xy -plane with corners $(0, 0)$, $(a, 0)$, $(0, b)$ and (a, b) . Its surface mass density is $\sigma(x, y) = \sigma_0 xy/ab$ where σ_0 is a constant. Find its center of mass.

$$\mathbf{R}_{\text{CM}}(t) = \frac{1}{M} \iiint_V \mathbf{r}(t) \rho(\mathbf{r}) \, dV, \quad \mathbf{r} = x \hat{\mathbf{e}}_x + y \hat{\mathbf{e}}_y + z \hat{\mathbf{e}}_z$$

$$\begin{aligned} \mathbf{R}_{\text{CM}} &= \frac{1}{M} \iiint_V \mathbf{r} \rho(\mathbf{r}) \, dV \\ &= \frac{1}{M} \iiint_V (x \hat{\mathbf{e}}_x + y \hat{\mathbf{e}}_y + z \hat{\mathbf{e}}_z) \rho(\mathbf{r}) \, dV \\ &= \hat{\mathbf{e}}_x \frac{1}{M} \iiint_V x \rho(\mathbf{r}) \, dV + \hat{\mathbf{e}}_y \frac{1}{M} \iiint_V y \rho(\mathbf{r}) \, dV + \hat{\mathbf{e}}_z \frac{1}{M} \iiint_V z \rho(\mathbf{r}) \, dV \\ &\equiv X_{\text{CM}} \hat{\mathbf{e}}_x + Y_{\text{CM}} \hat{\mathbf{e}}_y + Z_{\text{CM}} \hat{\mathbf{e}}_z \quad \blacksquare \end{aligned}$$

$$\begin{aligned}
 Z_{\text{CM}} &= \frac{1}{M} \int z \, dm = \frac{1}{M} \int \frac{M}{V} z \, dV = \frac{1}{V} \int z \pi r^2 \, dz \\
 &= \frac{1}{V} \int_0^R \pi z (R^2 - z^2) \, dz = \frac{3}{8} R \quad \blacksquare
 \end{aligned}$$

$$Z_{\text{CM}} = \frac{1}{M} \iiint_V z \rho \, dV = \frac{1}{V} \int_{r=0}^R \int_{\theta=0}^{\pi/2} \int_{\phi=0}^{2\pi} (r \cos \theta) r^2 \sin \theta \, dr \, d\theta \, d\phi = \frac{3}{8} R \quad \blacksquare$$

$$Y_{\text{CM}} = \frac{1}{M} \iiint_V y \rho \, dV = \frac{1}{V} \int_{r=0}^R \int_{\theta=0}^{\pi/2} \int_{\phi=0}^{2\pi} (r \sin \theta \sin \phi) r^2 \sin \theta \, dr \, d\theta \, d\phi = 0 \quad \blacksquare$$

$$\mathbf{r} = x \hat{\mathbf{e}}_x + y \hat{\mathbf{e}}_y, \quad \sigma(\mathbf{r}) = \sigma(x, y) = \sigma_0 \frac{xy}{ab}$$

$$M = \iint \sigma(\mathbf{r}) \, dA = \int_{x=0}^a \int_{y=0}^b \frac{\sigma_0}{ab} xy \, dx \, dy = \frac{1}{4} \sigma_0 ab \quad \blacksquare$$

$$X_{\text{CM}} = \frac{1}{M} \iint x \sigma(\mathbf{r}) \, dA = \frac{1}{M} \int_{x=0}^a \int_{y=0}^b x \left(\frac{\sigma_0}{ab} xy \right) \, dx \, dy = \frac{2}{3} a \quad \blacksquare$$

$$Y_{\text{CM}} = \frac{1}{M} \iint y \sigma(\mathbf{r}) \, dA = \frac{1}{M} \int_{x=0}^a \int_{y=0}^b y \left(\frac{\sigma_0}{ab} xy \right) \, dx \, dy = \frac{2}{3} b \quad \blacksquare$$

$$\mathbf{R}_{\text{CM}} = \frac{2}{3} a \hat{\mathbf{e}}_x + \frac{2}{3} b \hat{\mathbf{e}}_y \quad \blacksquare$$