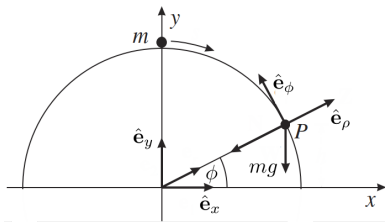


# Example: Particle on a hemisphere

- A particle of mass  $m$  is located at the “North pole” of a smooth hemisphere of radius  $R$  fixed on the ground. The particle slides down the hemisphere after a small kick.

- Particle is constrained to move on the hemisphere before breaking off:

$$\rho(t) = R \quad \Rightarrow \quad \begin{cases} \dot{\rho}(t) = 0 \\ \ddot{\rho}(t) = 0 \end{cases}$$



**EXERCISE 2.3:** Find the angle and the speed at which the particle breaks off from the hemisphere.

$$\mathbf{F}(t) = [N(t) - mg \sin \phi(t)] \hat{\mathbf{e}}_\rho(t) - mg \cos \phi(t) \hat{\mathbf{e}}_\phi(t)$$

$$\begin{cases} N(t) - mg \sin \phi(t) = -mR \dot{\phi}^2(t) \\ -mg \cos \phi(t) = mR \ddot{\phi}(t) \end{cases}$$

$$-mg \sin \phi(t_0) = -mR \dot{\phi}^2(t_0) \quad \Rightarrow \quad \dot{\phi}^2(t_0) = \frac{g}{R} \sin \phi(t_0) \quad \blacksquare$$

$$\ddot{\phi}(t) = -\frac{g}{R} \cos \phi(t) \quad \Rightarrow \quad \frac{d\dot{\phi}(\phi)}{d\phi} \frac{d\phi(t)}{dt} = -\frac{g}{R} \cos \phi(t)$$

$$\Rightarrow \int_{\dot{\phi}'=0}^{\dot{\phi}(t_0)} \dot{\phi}' d\dot{\phi}' = -\frac{g}{R} \int_{\phi'=\pi/2}^{\phi(t_0)} \cos \phi' d\phi'$$

$$\Rightarrow \frac{1}{2} \dot{\phi}^2(t_0) = -\frac{g}{R} [\sin \phi(t_0) - 1] \quad \blacksquare$$

$$\begin{cases} \dot{\phi}^2(t_0) = \frac{g}{R} \sin \phi(t_0) \\ \frac{1}{2} \dot{\phi}^2(t_0) = -\frac{g}{R} [\sin \phi(t_0) - 1] \end{cases}$$

$$\Rightarrow \begin{cases} \sin \phi(t_0) = \frac{2}{3} \\ \dot{\phi}^2(t_0) = \frac{2g}{3R} \end{cases} \quad \blacksquare$$

$$\phi(t_0) = \sin^{-1} \frac{2}{3} \approx 42^\circ \quad \blacksquare$$

$$\mathbf{v}(t_0) = \dot{\rho}(t_0) \hat{\mathbf{e}}_\rho + \rho(t_0) \dot{\phi}(t_0) \hat{\mathbf{e}}_\phi = -\sqrt{\frac{2Rg}{3}} \hat{\mathbf{e}}_\phi \quad \Rightarrow \quad v(t_0) = \sqrt{\frac{2Rg}{3}} \quad \blacksquare$$

# Projectile with resistance

- Linear resistance:  $\mathbf{F} = -mk\mathbf{v}$ ,  $k \geq 0$

- Equation of motion:

$$\frac{d^2\mathbf{r}(t)}{dt^2} = -g\hat{\mathbf{e}}_z - k\mathbf{v}(t)$$

- Initial conditions:

$$\mathbf{r}(0) = (x_0, y_0, z_0), \quad \mathbf{v}(0) = (0, v_0 \cos \theta_0, v_0 \sin \theta_0)$$

- Equation of motion in Cartesian coordinates:

$$\frac{d^2x(t)}{dt^2} = -kv_x(t), \quad \frac{d^2y(t)}{dt^2} = -kv_y(t), \quad \frac{d^2z(t)}{dt^2} = -g - kv_z(t)$$

# Projectile with resistance: $x$ -direction

$$\frac{d^2x(t)}{dt^2} = -kv_x(t), \quad x(0) = x_0, \quad v_x(0) = 0$$

- Solving for  $v_x(t)$ :

$$\frac{dv_x(t)}{dt} = -kv_x(t) \quad \Rightarrow \quad v_x(t) = 0$$

- Solving for  $x(t)$ :

$$v_x(t) = 0 \quad \Rightarrow \quad \frac{dx(t)}{dt} = 0 \quad \Rightarrow \quad x(t) = x_0$$

- Motion along the  $x$ -direction is essentially stationary

# Projectile with resistance: $y$ -direction

$$\frac{d^2 y(t)}{dt^2} = -k v_y(t), \quad y(0) = y_0, \quad v_y(0) = v_0 \cos \theta_0$$

- Solving:

$$v_y(t) = v_0 \cos \theta_0 e^{-kt}, \quad y(t) = y_0 + \frac{v_0 \cos \theta_0}{k} (1 - e^{-kt})$$

- Zero-friction limit:  $k \rightarrow 0$

$$v_y(t) \rightarrow v_0 \cos \theta_0, \quad y(t) \rightarrow y_0 + v_0 (\cos \theta_0) t$$

**EXERCISE 2.4:** Obtain short-time and long-time behaviours for  $v_y(t)$  and  $y(t)$ .

$$\frac{d^2y(t)}{dt^2} = -kv_y(t), \quad y(0) = y_0, \quad v_y(0) = v_0 \cos \theta_0$$

$$\frac{dv_y(t)}{dt} = -kv_y(t) \Rightarrow \int_{v'_y=v_0 \cos \theta_0}^{v_y} \frac{dv'_y}{v'_y} = -k \int_{t'=0}^t dt$$

$$\Rightarrow v_y(t) = v_0 \cos \theta_0 e^{-kt} \quad \blacksquare$$

$$\frac{dy(t)}{dt} = v_0 \cos \theta_0 e^{-kt} \Rightarrow \int_{y'=y_0}^y dy' = \int_{t'=0}^t v_0 \cos \theta_0 e^{-kt'} dt'$$

$$\Rightarrow y(t) = y_0 + \frac{v_0 \cos \theta_0}{k} (1 - e^{-kt}) \quad \blacksquare$$

$$v_y(t) = v_0 \cos \theta_0 e^{-kt}, \quad y(t) = y_0 + \frac{v_0 \cos \theta_0}{k} (1 - e^{-kt})$$

$$t \ll \frac{1}{k} \quad \rightarrow \quad \begin{cases} v_y(t) \rightarrow v_0 \cos \theta_0 (1 - kt) \\ y(t) \rightarrow y_0 + v_0 (\cos \theta_0) t - \frac{1}{2} k v_0 (\cos \theta_0) t^2 \end{cases} \quad \blacksquare$$

$$t \gg \frac{1}{k} \quad \rightarrow \quad \begin{cases} v_y(t) \rightarrow 0 \\ y(t) \rightarrow y_0 + \frac{v_0 \cos \theta_0}{k} \end{cases} \quad \blacksquare$$

$$\frac{dv_y(t)}{dt} = -k v_y(t) \quad \Rightarrow \quad \frac{dv_y(y)}{dy} \frac{dy(t)}{dt} = -k v_y(t)$$

$$\Rightarrow \int_{v'_y = v_0 \cos \theta_0}^{v_y} dv'_y = - \int_{y' = y_0}^y k dy' \quad \Rightarrow \quad v_y(y) = v_0 \cos \theta_0 - k (y - y_0) \quad \blacksquare$$



# Projectile with resistance: $z$ -direction

$$\frac{d^2 z(t)}{dt^2} = -g - k v_z(t), \quad z(0) = z_0, \quad v_z(0) = v_0 \sin \theta_0$$

- Solving:

$$v_z(t) = \left( v_0 \sin \theta_0 + \frac{g}{k} \right) e^{-kt} - \frac{g}{k}, \quad z(t) = z_0 + \frac{1}{k} \left( v_0 \sin \theta_0 + \frac{g}{k} \right) (1 - e^{-kt}) - \frac{gt}{k}$$

- Short-time behaviour:

$$v_z(t) \rightarrow v_0 \sin \theta_0 - (g + k v_0 \sin \theta_0) t, \quad z(t) \rightarrow z_0 + v_0 (\sin \theta_0) t - \frac{1}{2} (g + k v_0 \sin \theta_0) t^2$$

- Long-time behaviour:

$$v_z(t) \rightarrow -\frac{g}{k}, \quad z(t) \rightarrow z_0 + \frac{1}{k} \left( v_0 \sin \theta_0 + \frac{g}{k} \right) - \frac{gt}{k}$$

$$\frac{d^2 z(t)}{dt^2} = -g - kv_z(t), \quad z(0) = z_0, \quad v_z(0) = v_0 \sin \theta_0$$

$$\frac{dv_z(t)}{dt} = -g - kv_z(t) \Rightarrow \int_{v'_z=v_0 \sin \theta_0}^{v_z} \frac{dv'_z}{g + kv'_z} = - \int_{t'=0}^t dt'$$

$$\Rightarrow \frac{1}{k} \ln \frac{g + kv_z(t)}{g + kv_0 \sin \theta_0} = -t \Rightarrow v_z(t) = \left( v_0 \sin \theta_0 + \frac{g}{k} \right) e^{-kt} - \frac{g}{k} \quad \blacksquare$$

$$\frac{dz(t)}{dt} = \left( v_0 \sin \theta_0 + \frac{g}{k} \right) e^{-kt} - \frac{g}{k}$$

$$\Rightarrow \int_{z'=z_0}^z dz' = \int_{t'=0}^t \left[ \left( v_0 \sin \theta_0 + \frac{g}{k} \right) e^{-kt'} - \frac{g}{k} \right] dt'$$

$$\Rightarrow z(t) = z_0 + \frac{1}{k} \left( v_0 \sin \theta_0 + \frac{g}{k} \right) (1 - e^{-kt}) - \frac{gt}{k} \quad \blacksquare$$

# Projectile with resistance: horizontal range

- Time of the flight:  $z_0 = 0$

$$z(T) = 0 \quad \Rightarrow \quad (kv_0 \sin \theta_0 + g) (1 - e^{-kT}) - kgT = 0$$

- Dimensionless resistance parameter:

$$\epsilon \equiv \frac{kv_0}{g} \quad \Rightarrow \quad (\epsilon \sin \theta_0 + 1) (1 - e^{-kT}) - kT = 0$$

- Perturbation calculation for *weak* friction:  $\epsilon \ll 1$

$$T = \frac{2v_0 \sin \theta_0}{g} [1 + c_1 \epsilon + c_2 \epsilon^2 + \mathcal{O}(\epsilon^3)]$$

- Values for  $c_1$  and  $c_2$  are to be determined

# Projectile with resistance: horizontal range – cont'd

- Substitution, series expansion and solving:

$$T = \frac{2v_0 \sin \theta_0}{g} \left[ 1 - \frac{1}{3} \epsilon \sin \theta_0 + \frac{2}{9} \epsilon^2 \sin^2 \theta_0 + \mathcal{O}(\epsilon^3) \right]$$

- Horizontal range:  $y_0 = 0$

$$R \equiv y(T) = \frac{v_0 \cos \theta_0}{k} (1 - e^{-kT})$$

- Substitutions and series expansion:

$$R = \frac{2v_0^2 \sin \theta_0 \cos \theta_0}{g} \left[ 1 - \frac{4}{3} \epsilon \sin \theta_0 + \frac{14}{9} \epsilon^2 \sin^2 \theta_0 + \mathcal{O}(\epsilon^3) \right]$$

**EXERCISE 2.5:** Complete the perturbation calculations to obtain the expression for  $R$  up to  $\epsilon^2$ .

$$(\epsilon \sin \theta_0 + 1) (1 - e^{-kT}) - kT = 0, \quad \epsilon \equiv \frac{kv_0}{g}, \quad T = \frac{2v_0 \sin \theta_0}{g} (1 + c_1\epsilon + c_2\epsilon^2)$$

$$(\epsilon \sin \theta_0 + 1) (1 - e^{-kT}) - kT = 0$$

$$\Rightarrow (\epsilon \sin \theta_0 + 1) (1 - e^{-\epsilon g T / v_0}) - \epsilon \frac{gT}{v_0} = 0$$

$$\Rightarrow (\epsilon \sin \theta_0 + 1) \left\{ 1 - \exp \left[ -\frac{\epsilon g}{v_0} \frac{2v_0 \sin \theta_0}{g} (1 + c_1\epsilon + c_2\epsilon^2) \right] \right\} - \frac{\epsilon g}{v_0} \frac{2v_0 \sin \theta_0}{g} (1 + c_1\epsilon + c_2\epsilon^2) = 0$$

$$\Rightarrow (\epsilon \sin \theta_0 + 1) \left\{ 1 - \exp \left[ -2\epsilon \sin \theta_0 (1 + c_1\epsilon + c_2\epsilon^2) \right] \right\} - 2\epsilon \sin \theta_0 (1 + c_1\epsilon + c_2\epsilon^2) = 0$$

$$f(\epsilon) \equiv (\epsilon \sin \theta_0 + 1) \left\{ 1 - \exp \left[ -2\epsilon \sin \theta_0 (1 + c_1 \epsilon + c_2 \epsilon^2) \right] \right\} - 2\epsilon \sin \theta_0 (1 + c_1 \epsilon + c_2 \epsilon^2)$$

$$\Rightarrow \begin{cases} f^{(0)}(0) = 0 \\ f^{(1)}(0) = 0 \\ f^{(2)}(0) = 0 \\ f^{(3)}(0) = -4 \sin^2 \theta_0 (3c_1 + \sin \theta_0) \\ f^{(4)}(0) = 16 \sin^2 \theta_0 (\sin^2 \theta_0 - 3c_1^2 - 3c_2) \end{cases}$$

$$\Rightarrow f(\epsilon) = \frac{1}{3!} (-4 \sin^2 \theta_0) (3c_1 + \sin \theta_0) \epsilon^3 + \frac{1}{4!} (16 \sin^2 \theta_0) (\sin^2 \theta_0 - 3c_1^2 - 3c_2) \epsilon^4 + \mathcal{O}(\epsilon^5)$$

$$f(\epsilon) = 0 \quad \Rightarrow \quad \begin{cases} 3c_1 + \sin \theta_0 = 0 \\ \sin^2 \theta_0 - 3c_1^2 - 3c_2 = 0 \end{cases} \quad \Rightarrow \quad \begin{cases} c_1 = -\frac{1}{3} \sin \theta_0 \\ c_2 = \frac{2}{9} \sin^2 \theta_0 \end{cases} \quad \blacksquare$$

$$R = \frac{v_0 \cos \theta_0}{k} (1 - e^{-kT}) , \quad \epsilon \equiv \frac{kv_0}{g} , \quad T = \frac{2v_0 \sin \theta_0}{g} (1 + c_1\epsilon + c_2\epsilon^2)$$

$$\begin{aligned} R &= \frac{v_0 \cos \theta_0}{k} (1 - e^{-kT}) = \frac{v_0^2 \cos \theta_0}{\epsilon g} (1 - e^{-\epsilon g T / v_0}) \\ &= \frac{v_0^2 \cos \theta_0}{\epsilon g} \left\{ 1 - \exp \left[ -\frac{\epsilon g}{v_0} \frac{2v_0 \sin \theta_0}{g} (1 + c_1\epsilon + c_2\epsilon^2) \right] \right\} \\ &= \frac{v_0^2 \cos \theta_0}{\epsilon g} \left\{ 1 - \exp \left[ -2\epsilon \sin \theta_0 (1 + c_1\epsilon + c_2\epsilon^2) \right] \right\} \quad \blacksquare \end{aligned}$$

$$g(\epsilon) \equiv \frac{v_0^2 \cos \theta_0}{\epsilon g} \left\{ 1 - \exp \left[ -2\epsilon \sin \theta_0 (1 + c_1 \epsilon + c_2 \epsilon^2) \right] \right\}$$

$$\Rightarrow \begin{cases} g^{(0)}(0) = \frac{2v_0^2 \sin \theta_0 \cos \theta_0}{g} \\ g^{(1)}(0) = \frac{2v_0^2 \sin \theta_0 \cos \theta_0}{g} (c_1 - \sin \theta_0) \\ g^{(2)}(0) = \frac{4v_0^2 \sin \theta_0 \cos \theta_0}{3g} (2 + 3c_2 - 6c_1 \sin \theta_0 - 2 \cos^2 \theta_0) \end{cases}$$

$$\begin{cases} c_1 = -\frac{1}{3} \sin \theta_0 \\ c_2 = \frac{2}{9} \sin^2 \theta_0 \end{cases} \Rightarrow \begin{cases} g^{(0)}(0) = \frac{2v_0^2 \sin \theta_0 \cos \theta_0}{g} \\ g^{(1)}(0) = -\frac{8v_0^2 \sin^2 \theta_0 \cos \theta_0}{3g} \\ g^{(2)}(0) = \frac{56v_0^2 \sin^3 \theta_0 \cos \theta_0}{9g} \end{cases} \quad \blacksquare$$

$$\Rightarrow R = \frac{2v_0^2 \sin \theta_0 \cos \theta_0}{g} \left[ 1 - \frac{4}{3} \epsilon \sin \theta_0 + \frac{14}{9} \epsilon^2 \sin^2 \theta_0 + \mathcal{O}(\epsilon^3) \right] \quad \blacksquare$$



# Linear homogeneous ODEs

- $n$ -order homogeneous equation with constant coefficients:  $a_n \neq 0$

$$a_n y^{(n)}(x) + a_{n-1} y^{(n-1)}(x) + \cdots + a_1 y^{(1)}(x) + a_0 y^{(0)}(x) = 0$$

- Characteristics equation:  $n$ -degree polynomial of  $\lambda$

$$y(x) = e^{\lambda x} \quad \Rightarrow \quad a_n \lambda^n + a_{n-1} \lambda^{n-1} + \cdots + a_1 \lambda + a_0 = 0$$

- Characteristic roots give linearly independent solutions:
  - $\lambda$  is a real root with no degeneracy:  $e^{\lambda x}$  is the solution
  - $\lambda$  is a real root with doubly degeneracy:  $e^{\lambda x}$  and  $x e^{\lambda x}$  are solutions
  - $\lambda = \alpha \pm i\beta$  are complex root with no degeneracy:  $e^{\alpha x} \sin \beta x$  and  $e^{\alpha x} \cos \beta x$  are solutions
  - $\lambda = \alpha \pm i\beta$  are complex root with doubly degeneracy:  $e^{\alpha x} \sin \beta x$ ,  $x e^{\alpha x} \sin \beta x$ ,  $e^{\alpha x} \cos \beta x$  and  $x e^{\alpha x} \cos \beta x$  are solutions

# Linear homogeneous ODEs – cont'd

- **Wronskian** of a set of  $n$  functions  $\{f_1(x), \dots, f_n(x)\}$ :

$$W[f_1, f_2, \dots, f_n](x) \equiv \begin{vmatrix} f_1(x) & f_2(x) & \cdots & f_n(x) \\ f_1'(x) & f_2'(x) & \cdots & f_n'(x) \\ \vdots & \vdots & \ddots & \vdots \\ f_1^{(n-1)}(x) & f_2^{(n-1)}(x) & \cdots & f_n^{(n-1)}(x) \end{vmatrix}$$

- General solution:  $\{y_n(x)\}$  is a set of linearly independent solutions

$$W[y_1, y_2, \dots, y_n](x) \neq 0$$

$$y(x) = C_1 y_1(x) + C_2 y_2(x) + \cdots + C_{n-1} y_{n-1}(x) + C_n y_n(x)$$

- Constants  $C_i$  are to be determined from initial/boundary conditions

# Charge in magnetic field

- A point charge of mass  $m$  and charge  $q$  is moving in a region of uniform magnetic field  $\mathbf{B} = B_0 \hat{\mathbf{e}}_y$
- Equation of motion in Cartesian coordinates:  $\omega \equiv qB_0/m$

$$m \frac{d^2 \mathbf{r}(t)}{dt^2} = q \mathbf{v}(t) \times \mathbf{B} \quad \Rightarrow \quad \begin{cases} \frac{d^2 x(t)}{dt^2} = -\omega \frac{dz(t)}{dt} \\ \frac{d^2 y(t)}{dt^2} = 0 \\ \frac{d^2 z(t)}{dt^2} = \omega \frac{dx(t)}{dt} \end{cases}$$

- Initial conditions:

$$\mathbf{r}(0) = (x_0, y_0, z_0) , \quad \mathbf{v}(0) = (0, v_{y0}, v_{z0})$$

# Charge in magnetic field: $y$ -direction

$$\frac{d^2y(t)}{dt^2} = 0, \quad y(0) = y_0, \quad v_y(0) = v_{y0}$$

- Solving for  $v_y(t)$ :

$$\frac{dv_y(t)}{dt} = 0 \quad \Rightarrow \quad v_y(t) = v_{y0}$$

- Solving for  $y(t)$ :

$$v_y(t) = v_{y0} \quad \Rightarrow \quad \frac{dy(t)}{dt} = v_{y0} \quad \Rightarrow \quad v_y(t) = v_{y0}t + y_0$$

- Motion along the  $y$ -direction is essentially uniform

# Charge in magnetic field: $x$ and $z$ -directions

- Coupled differential equations:

$$\left\{ \begin{array}{l} \frac{d^2 x(t)}{dt^2} = -\omega \frac{dz(t)}{dt} \\ \frac{d^2 z(t)}{dt^2} = \omega \frac{dx(t)}{dt} \end{array} \right., \quad \left\{ \begin{array}{l} x(0) = x_0, \quad v_x(0) = 0 \\ z(0) = z_0, \quad v_z(0) = v_{z0} \end{array} \right.$$

- Decoupling and solving:

$$\left\{ \begin{array}{l} x(t) = C_1 \cos \omega t + C_2 \sin \omega t + C_0 \\ z(t) = D_1 \cos \omega t + D_2 \sin \omega t + D_0 \end{array} \right.$$

- Question: Are  $C_1$ ,  $C_2$ ,  $C_0$ ,  $D_1$ ,  $D_2$  and  $D_0$  all independent from each other?

**EXERCISE 2.6:** Obtain the general solutions for the coupled differential equations for  $x(t)$  and  $z(t)$ .

$$\begin{cases} \frac{d^2 x(t)}{dt^2} = -\omega \frac{dz(t)}{dt} \\ \frac{d^2 z(t)}{dt^2} = \omega \frac{dx(t)}{dt} \end{cases} \Rightarrow \begin{cases} \frac{d^3 x(t)}{dt^3} = -\omega^2 \frac{dx(t)}{dt} \\ \frac{d^3 z(t)}{dt^3} = -\omega^2 \frac{dz(t)}{dt} \end{cases}$$

$$x(t) = e^{\lambda t} \Rightarrow \frac{d^3 x(t)}{dt^3} = -\omega^2 \frac{dx(t)}{dt} \Rightarrow \lambda (\lambda^2 + \omega^2) = 0 \Rightarrow \lambda = 0, \pm i\omega \quad \blacksquare$$

$$x(t) = Ae^{+i\omega t} + B + Ce^{-i\omega t} = (A + C) \cos \omega t + i(A - C) \sin \omega t + B \quad \blacksquare$$

# Charge in magnetic field: $x$ and $z$ -directions – cont'd

- Eliminating dependencies:

$$\begin{cases} x(t) = C_1 \cos \omega t + C_2 \sin \omega t + C_0 \\ z(t) = -C_2 \cos \omega t + C_1 \sin \omega t + D_0 \end{cases}$$

- Imposing initial conditions for  $x(t)$  and  $z(t)$ :

$$\begin{cases} x(0) = x_0 & \Rightarrow & C_0 = x_0 - C_1 \\ z(0) = z_0 & \Rightarrow & D_0 = z_0 \end{cases}$$

- Imposing initial conditions for  $\dot{x}(t)$  and  $\dot{z}(t)$ :

$$\begin{cases} \dot{x}(0) = 0 & \Rightarrow & C_2 = 0 \\ \dot{z}(0) = v_{z0} & \Rightarrow & C_1 = \frac{v_{z0}}{\omega} \end{cases}$$

$$\begin{cases} \frac{d^2x(t)}{dt^2} = -\omega \frac{dz(t)}{dt} \\ \frac{d^2z(t)}{dt^2} = \omega \frac{dx(t)}{dt} \end{cases}, \quad \begin{cases} x(t) = C_1 \cos \omega t + C_2 \sin \omega t + C_0 \\ z(t) = D_1 \cos \omega t + D_2 \sin \omega t + D_0 \end{cases}$$

$$\frac{d^2x(t)}{dt^2} = -\omega \frac{dz(t)}{dt}$$

$$\Rightarrow -\omega^2 C_1 \cos \omega t - \omega^2 C_2 \sin \omega t = -\omega (-\omega D_1 \sin \omega t + \omega D_2 \cos \omega t)$$

$$\Rightarrow \begin{cases} -\omega^2 C_1 = -\omega^2 D_2 \\ -\omega^2 C_2 = \omega^2 D_1 \end{cases} \Rightarrow \begin{cases} D_2 = C_1 \\ D_1 = -C_2 \end{cases} \quad \blacksquare$$

$$\begin{vmatrix} \sin \omega t & \cos \omega t \\ \omega \cos \omega t & -\omega \sin \omega t \end{vmatrix} = -\omega \neq 0 \quad \blacksquare$$



# Charge in magnetic field: trajectory

- Position and velocity:

$$\left\{ \begin{array}{l} x(t) = \frac{v_{z0}}{\omega} \cos \omega t + x_0 - \frac{v_{z0}}{\omega} \\ y(t) = v_{y0}t + y_0 \\ z(t) = \frac{v_{z0}}{\omega} \sin \omega t + z_0 \end{array} \right. , \quad \left\{ \begin{array}{l} v_x(t) = -v_{z0} \sin \omega t \\ v_y(t) = v_{y0} \\ v_z(t) = v_{z0} \cos \omega t \end{array} \right.$$

- Trajectory of the point charge is a circular helix of radius  $mv_{z0}/qB_0$  centered at  $(x, z) = (x_0 - mv_{z0}/qB_0, z_0)$

$$\left[ x(t) - \left( x_0 - \frac{mv_{z0}}{qB_0} \right) \right]^2 + [z(t) - z_0]^2 = \left( \frac{mv_{z0}}{qB_0} \right)^2$$