

**Tutorial 10** (Block encoding techniques<sup>1</sup>)

Qubitization consists of two fundamental subroutines: block encoding and signal processing. In this tutorial, we discuss strategies for block encoding, for which there are various problem-specific methods available.

The goal is to realize a linear map  $\mathcal{H} \in \mathbb{C}^{N \times N}$  with spectral norm  $\|\mathcal{H}\| \leq 1$  by a quantum circuit, with  $N = 2^n$  for  $n$  qubits. Block encoding refers to embedding  $\mathcal{H}$  into a larger unitary matrix  $U$ , which is necessary since  $\mathcal{H}$  is not unitary in general. Specifically, we use additional ancillary qubits, such that  $U$  maps  $\mathbb{C}^K \otimes \mathbb{C}^N \rightarrow \mathbb{C}^K \otimes \mathbb{C}^N$ , with  $K = 2^k$  for  $k$  ancillary qubits. The actual encoding is then realized using a special ancillary state  $|G\rangle \in \mathbb{C}^K$ , such that

$$\mathcal{H} = (\langle G| \otimes I_N) U (|G\rangle \otimes I_N),$$

where  $I_N$  refers to the  $N \times N$  identity matrix. For simplicity, we take  $|G\rangle = |0, \dots, 0\rangle$  in the following. The matrix representation of  $U$  is then of the form

$$U = \begin{pmatrix} \mathcal{H} & * \\ * & * \end{pmatrix},$$

where the stars refer to unspecified blocks.

- (a) Find a block encoding for  $\mathcal{H}$  Hermitian and with spectral norm  $\|\mathcal{H}\| \leq 1$  using only a single ancillary qubit.

The construction for  $U$  so far is general, but still leaves open the question of how to express  $U$  in terms of elementary gates. As an example for a more explicit circuit construction, we investigate the case that  $\mathcal{H}$  is a linear combination of unitaries (LCU) in the following. Such a decomposition always exists (but might contain many terms): for example, one can use the set of Pauli strings as basis for the LCU.

- (b) Consider the following Hamiltonian written as a sum of Pauli strings:

$$\mathcal{H} = \frac{1}{4}(XX + YZ + YY + ZX).$$

Construct a block encoding for  $\mathcal{H}$  using elementary quantum circuit gates.

- (c) Generalize the approach in (b), assuming that there exists a LCU decomposition of the Hamiltonian of the form

$$\mathcal{H} = \sum_{i=0}^{m-1} \alpha_i V_i. \quad (1)$$

Here each  $V_i$  is a unitary matrix, and without loss of generality, we can assume that the coefficients  $\alpha_i$  are positive numbers, since phase factors can be absorbed into the unitaries. How does the number of ancillary qubits scale?

**Exercise 10.1** (Numerical simulation of block encoding)

In this exercise, we programmatically realize the block encoding techniques discussed in the tutorial.

- (a) Write a Python function which constructs

$$U = \begin{pmatrix} \mathcal{H} & \sqrt{I - \mathcal{H}^2} \\ \sqrt{I - \mathcal{H}^2} & -\mathcal{H} \end{pmatrix} \quad (2)$$

for a given Hermitian matrix  $\mathcal{H}$ . To test your implementation, first sample a random  $3 \times 3$  complex Hermitian matrix  $\mathcal{H}$ , and rescale  $\mathcal{H}$  such that all eigenvalues are in the interval  $[-1, 1]$ . Use your code to compute the corresponding  $U$ , and verify numerically that  $U$  is unitary and contains  $\mathcal{H}$  in its upper left block.

Hint: The representation  $U = Z \otimes \mathcal{H} + X \otimes \sqrt{I - \mathcal{H}^2}$  and `scipy.linalg.sqrtm` for the matrix square root might be helpful.

- (b) Inserting the spectral decomposition  $\mathcal{H} = \sum_j \lambda_j |\psi_j\rangle \langle \psi_j|$  (with eigenvalues  $\lambda_j \in [-1, 1]$ ) into Eq. (2) leads to

$$U = \sum_j \begin{pmatrix} \lambda_j & \sqrt{1 - \lambda_j^2} \\ \sqrt{1 - \lambda_j^2} & -\lambda_j \end{pmatrix} \otimes |\psi_j\rangle \langle \psi_j| = \sum_j R(\lambda_j) \otimes |\psi_j\rangle \langle \psi_j|, \quad R(a) = \begin{pmatrix} a & \sqrt{1 - a^2} \\ \sqrt{1 - a^2} & -a \end{pmatrix}.$$

<sup>1</sup>G. H. Low and I. L. Chuang: *Hamiltonian simulation by qubitization*. Quantum 3, 163 (2019), and L. Lin: *Lecture notes on quantum algorithms for scientific computation*. <https://math.berkeley.edu/~linlin/qasc/>

# hw10

July 1, 2023

```
[ ]: # Ex 10.1
# a)
import numpy as np
import scipy

X = np.array([[0, 1], [1, 0]])
Y = np.array([[0, -1j], [1j, 0]])
Z = np.array([[1, 0], [0, -1]])

def U(H):
    A = np.outer(Z, H)
    Id = np.eye((int)(np.sqrt(H.size)))
    B = scipy.linalg.sqrtm(Id - np.dot(H, H))
    C = np.outer(X, B)
    return A + C

H = np.array([[1, 0], [0, 1]])

phi = np.array([[1, 0, 0], [0, 1, 0], [0, 0, 1]])
# print(phi[0].T)
eigen_states = []
for i in phi:
    eigen_states.append(np.outer(i.T, i))

R = np.array([1, 0.5, 0.1])
H = R[0]*eigen_states[0] + R[1] *eigen_states[1] + R[2]*eigen_states[2]
print(H)
print(U(H))
```

```
[[1.  0.  0. ]
 [0.  0.5 0. ]
 [0.  0.  0.1]]
[[ 1.          0.          0.          0.5          0.
   0.          0.          0.1          ]
 [ 0.          0.          0.          0.          0.8660254  0.
   0.          0.          0.99498744]]
```

```
[ 0.      0.      0.      0.      0.8660254  0.
  0.      0.      0.99498744]
[-1.      0.      0.      0.      -0.5      0.
  0.      0.     -0.1      ]]
```

```
[ ]: # Ex 10.1
# b)
phi = np.array([[1, 0, 0],[0, 1, 0], [0, 0, 1]])
# print(phi[0].T)
eigen_states = []
for i in phi:
    eigen_states.append(np.outer(i.T, i))

H_second = np.zeros((4,9))
for i in range(3):
    H_second += np.outer(U[R[i]], eigen_states[i])

print(H_second)
```

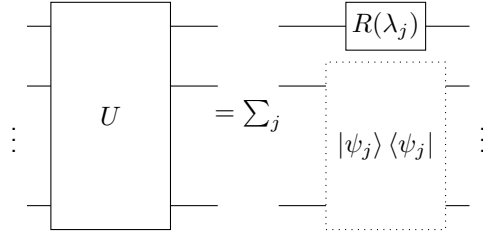
```
[ [ 1.      0.      0.      0.      0.5      0.
   0.      0.      0.1      ]
  [ 0.      0.      0.      0.      0.8660254  0.
   0.      0.      0.99498744]
  [ 0.      0.      0.      0.      0.8660254  0.
   0.      0.      0.99498744]
 [-1.      0.      0.      0.      -0.5      0.
   0.      0.     -0.1      ]]
```

```
[ ]: print(U(H) == H_second)
```

```
[ [ True True True True True True True True True]
  [ True True True True True True True True True]
  [ True True True True True True True True True]
  [ True True True True True True True True True]]
```

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[ ]:
```

A circuit diagram representation of this equation is given by



Verify that  $R(\lambda_j) \otimes |\psi_j\rangle\langle\psi_j|$  agrees with  $U$  from (a) (for the same random Hamiltonian matrix  $\mathcal{H}$ ).

- (c) Write a Python/NumPy program to simulate the block encoding method in tutorial 10 (c) for a linear combination of unitaries (LCU) in Eq. (1) with  $m = 4$  and  $\|\alpha\|_1 = 1$ . You can realize each controlled- $V_i$  gate in the SELECT operation via

$$\text{controlled-}V_i = |i\rangle\langle i| \otimes V_i + (I_m - |i\rangle\langle i|) \otimes I_N,$$

where  $|i\rangle$  is the  $i$ -th unit vector in  $\mathbb{R}^m$ . The second summand corresponds to inactive control. To construct the PREPARE unitary matrix, you can start from a  $m \times 1$  matrix containing  $\sqrt{\alpha}$  as its single column, and then use a “complete” QR-decomposition (`np.linalg.qr`) to extend it to a  $m \times m$  matrix. The resulting  $Q$  matrix could contain  $-\sqrt{\alpha}$  (instead of  $\sqrt{\alpha}$ ) in its first column; in this case, use  $-Q$  (instead of  $Q$ ) as PREPARE gate.

To test your implementation, first draw four random unitary  $3 \times 3$  matrices  $\{V_i\}$  (Haar random distribution) and corresponding random non-negative coefficients  $\alpha_i$ , normalized such that  $\|\alpha\|_1 = 1$ . Then compute the matrix representation of the circuit from tutorial 10 (c), and verify that it is unitary and contains  $\mathcal{H}$  defined in Eq. (1) in its upper left  $3 \times 3$  block.

Hint: You can use `scipy.stats.unitary_group` to sample a Haar random unitary matrix.

### Exercise 10.2 (General signal-processing operation)

The quantum eigenvalue transformation combines the block encoding of a Hamiltonian  $\mathcal{H}$  with quantum signal processing. In case of a single auxiliary qubit for block encoding as in tutorial 10 (a), the corresponding quantum circuit is specified by (with  $\vec{\varphi} \in \mathbb{R}^{d+1}$ )

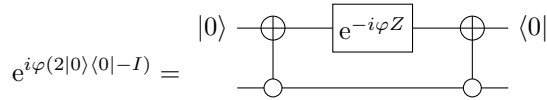
$$U_{\vec{\varphi}} = e^{i\varphi_0 Z} \dots U^\dagger e^{i\varphi_{d-1} Z} U e^{i\varphi_d Z},$$

where  $U$  is the block encoding gate, the alternation between  $U$  and  $U^\dagger$  is analogous to the amplitude amplification example, and the signal-processing gates  $e^{i\varphi_k Z}$  act on the auxiliary qubit. Note the representation  $e^{i\varphi Z} = e^{i\varphi(2|0\rangle\langle 0| - I)}$ . For a general block encoding gate, the signal-processing phase operation generalizes to (for  $\varphi \in \mathbb{R}$ )

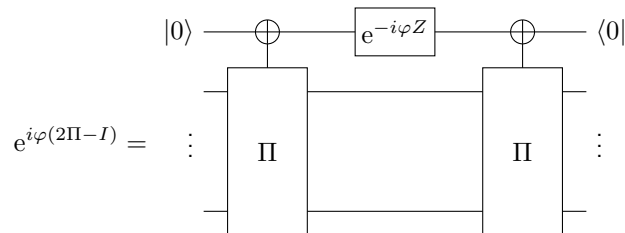
$$\Pi_\varphi = e^{i\varphi(2\Pi - I)},$$

where  $\Pi$  is the projector onto the block in  $U$  containing  $\mathcal{H}$ . In the context of tutorial 10,  $\Pi = |G\rangle\langle G| \otimes I_N$ . In this exercise, our goal is to find quantum circuits realizing such signal-processing operations.

- (a) Prove that the following circuit (using an additional auxiliary qubit) implements the projector-controlled phase operation for  $\Pi = |0\rangle\langle 0|$ :



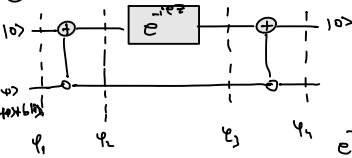
- (b) A circuit construction for a general projector  $\Pi$  is shown below, where the control is activated by states from the subspace which  $\Pi$  projects to. Argue why this indeed implements  $\Pi_\varphi$ :



- (c) Explicitly specify the circuit from (b) for  $\Pi = |G\rangle\langle G|$  with  $|G\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$  using only elementary gates. Hint: Think about a circuit that transforms  $|G\rangle \mapsto |00\rangle$ . You are allowed to use a multi-controlled NOT gate as elementary gate.

Ex 10.2.

6.



$$e^{-i\theta Z} = \cos(\theta)I - i\sin(\theta)Z$$

$$= \begin{pmatrix} e^{-i\theta} & 0 \\ 0 & e^{i\theta} \end{pmatrix}$$

$$\psi_1 = |0\rangle \otimes (a|0\rangle + b|1\rangle)$$

$$\psi_2 = a|00\rangle + b|11\rangle$$

$$\psi_3 = a \cdot e^{-i\theta} |00\rangle + b \cdot e^{i\theta} |11\rangle$$

$$\psi_4 = a \cdot e^{-i\theta} |00\rangle + b \cdot e^{i\theta} |11\rangle$$

$$= |0\rangle \otimes (a e^{-i\theta} |0\rangle + b e^{i\theta} |1\rangle)$$

$$= |0\rangle \otimes \left( \begin{pmatrix} e^{-i\theta} & 0 \\ 0 & e^{i\theta} \end{pmatrix} \cdot (a|0\rangle + b|1\rangle) \right)$$

$$= |0\rangle \otimes (e^{i\theta Z} (a|0\rangle + b|1\rangle))$$

$$= |0\rangle \otimes e^{i\theta (2|0\rangle\langle 0| - I)} \cdot (a|0\rangle + b|1\rangle)$$

7.  $\pi_{\psi'}$