

MESA

Modules for Experiments in Stellar Astrophysics

- [MESA home](#)
- [code capabilities](#)
- [prereqs & installation](#)
- [getting started](#)
- [using pgstar](#)
- [using MESA output](#)
- [beyond inlists \(extending MESA\)](#)
- [troubleshooting](#)
- [FAQ](#)
- [star_job defaults](#)
- [controls defaults](#)
- [pgstar defaults](#)
- [binary_controls defaults](#)
- [news archive](#)
- [documentation archive](#)

Latest News

- 21 Mar 2018
 - » [Release 10398](#)
- 28 Jan 2018
 - » [New MESA SDK Version](#)
- 13 Jan 2018
 - » [Summer School 2018](#)
- 23 Oct 2017
 - » [Release 10108](#)
- 23 Oct 2017
 - » [Instrument Paper 4](#)
- 21 Sep 2017
 - » [New MESA SDK Version](#)
- 12 Sep 2017
 - » [Release 10000](#)
- 01 Sep 2017
 - » [Docker, Python and Web](#)
- 14 Aug 2017
 - » [New mesa-users list](#)
- 02 Aug 2017
 - » [New MESA SDK Version](#)



This web documentation corresponds to the most recent MESA release (r10398). Documentation for past versions can be [found here](#).

If you find errors or formatting issues, please email Josiah Schwab using [this link](#).

This page documents the MESA options that are part of the controls namelist. It is autogenerated from the file \$MESA_DIR/star/defaults/controls.defaults.

Boxes like

```
option_name = 'default'
```

show the default value of each option. To override the default values, add an entry to the controls namelist in your inlist.

Contents

- [1 specifications for starting model](#)
- [2 controls for output](#)
- [3 definition of core overshooting boundary for output](#)
- [4 definition of core boundaries](#)
- [5 when to stop](#)
- [6 mixing parameters](#)
 - [6.1 overshooting](#)
 - [6.2 Predictive mixing](#)
 - [6.3 Self-Driving Overshoot & Semiconvection \(S-DOS\) mixing](#)
 - [6.4 New overshooting](#)
 - [6.5 turbulence](#)
 - [6.6 mixing misc](#)
 - [6.7 glitches](#)
- [7 rotation controls](#)
- [8 atmosphere boundary conditions](#)
- [9 mass gain or loss](#)
 - [9.1 implicit wind computation](#)
 - [9.2 remove H wind](#)
 - [9.3 controls for adjust mass](#)
- [10 composition controls](#)
- [11 mesh adjustment](#)
- [12 nuclear reaction controls](#)
- [13 element diffusion](#)
 - [13.1 parameters for ionization solver](#)
- [14 eos controls](#)
- [15 opacity controls](#)
 - [15.1 OP mono opacities](#)
- [16 asteroseismology controls](#)
 - [16.1 Brunt controls](#)
- [17 structure equations](#)
 - [17.1 eps_grav](#)
- [18 solver controls](#)
- [19 split mixing](#)
- [20 timestep controls](#)
 - [20.1 limits based on relative changes in variables L, P, Rho, T, R, eps_nuc](#)
 - [20.2 limits based on integrated power at each point for each category of nuclear reaction](#)
 - [20.3 limits based on total integrated power at surface for all nuclear reactions](#)
 - [20.4 limits based on changes in mass of the star](#)
 - [20.5 limits based on changes in log total angular momentum](#)
- [21 debugging controls](#)

- [22 miscellaneous controls](#)
 - [22.1 mixing diffusion coeffs](#)

specifications for starting model ¶

NOTE: if you are loading a saved model, then the following initial values are NOT USED to modify the model. in particular, you cannot use these to change Y or Z of an existing model. if you want to do that, see `star_job.defaults` controls such as `change_Y`. however, these are reported in output as the initial values for the star.

initial_mass ¶

initial mass in Msun units. can be any value you'd like when you are creating a pre-main sequence model.

not used when loading a saved model. however is reported in output as the initial mass of the star.

if you are loading a ZAMS model and the requested mass is in the range of prebuilt models, the code will interpolate in mass using the closest prebuilt models. if the requested mass is beyond the range of the prebuilt models, the code will load the closest one and then call "relax mass" to create a model to match the request. the prebuilt range is 0.08 Msun to 100 Msun, so the `relax_mass` method is only used for extreme cases. there are enough prebuilt models that the interpolation in mass seems to work fine for many applications.

```
initial_mass = 1
```

initial_z ¶

initial metallicity for create pre-ms and create initial model `initial_z` can be any value from 0 to 0.04

not used when loading a saved model. however is reported in output as the initial Z of the star.

however, if you are loading a zams model, then `initial_z` must match one of the prebuilt values. look in the 'data/star_data/zams_models' directory to see what prebuilt zams Z's are available. at time of writing, only 0.02 was included in the standard version of star.

```
initial_z = 0.02d0
```

initial_y ¶

initial helium mass fraction for create pre-ms and create initial (`< 0` means use default which is $0.24 + 2 * \text{initial_z}$)

not used when loading a saved model or a zams model. however is reported in output as the initial Y of the star.

NOTE: this is only used for create pre-main-sequence model and create initial model, and not when loading a zams model.

```
initial_y = -1
```

controls for output ¶

terminal_interval ¶

write info to terminal when $\text{mod}(\text{model_number}, \text{terminal_interval}) = 0$. note: this replaces the obsolete control `terminal_cnt`.

```
terminal_interval = 1
```

write_header_frequency ¶

output the log header info to the terminal when $\text{mod}(\text{model_number}, \text{write_header_frequency} * \text{terminal_interval}) = 0$.

```
write_header_frequency = 10
```

extra_terminal_output_file ¶

if not empty, output terminal info to this file in addition to terminal. this does not capture all of the terminal output – just the common items. it is intended for use in situations where you cannot directly see the terminal output such as when running on a cluster. if you want to be able to monitor the progress for such cases, you can set `extra_terminal_output_file = 'log'` and then do `tail -f log` to view the terminal output as it is recorded in the file.

```
extra_terminal_output_file = ''
```

terminal_show_age_in_years ¶

if false, then show in seconds

```
terminal_show_age_in_years = .true.
```

terminal_show_Age_in_days ¶

```
terminal_show_Age_in_days = .false.
```

num_trace_history_values ¶

any valid name for a history data column, such as `surf_v_rot` for example if you have rapid rotation at the surface, you might want to try something like this:

```
num_trace_history_values = 7
trace_history_value_name(1) = 'surf_v_rot'
trace_history_value_name(2) = 'surf_omega_div_omega_crit'
```

```
trace_history_value_name(3) = 'log_rotational_mdot_boost'  
trace_history_value_name(4) = 'log_total_angular_momentum'  
trace_history_value_name(5) = 'center n14'  
trace_history_value_name(6) = 'surface n14'  
trace_history_value_name(7) = 'average n14'
```

value must be less than or equal to 10

```
num_trace_history_values = 0
```

trace_history_value_name(:) ¶

write values to terminal

```
trace_history_value_name(:) = ''
```

photo_directory ¶

directory for binary snapshots used in restarts

```
photo_directory = 'photos'
```

photo_interval ¶

save a photo file for possible restarting when $\text{mod}(\text{model_number}, \text{photo_interval}) = 0$. note: this replaces the obsolete control photostep.

```
photo_interval = 50
```

photo_digits ¶

use this many digits from the end of the model_number for the photo name

```
photo_digits = 3
```

log_directory ¶

for data files about the run

```
log_directory = 'LOGS'
```

do_history_file ¶

history file is created if this is true

```
do_history_file = .true.
```

history_interval ¶

append an entry to the history.data file when $\text{mod}(\text{model_number}, \text{history_interval}) = 0$.

```
history_interval = 5
```

star_history_name ¶

name of history file

```
star_history_name = 'history.data'
```

star_history_header_name ¶

If not empty, then put star history header info in star_history_name file. In this case the history file has only data, making it easier to use with some plotting packages.

```
star_history_header_name = ''
```

star_history_dbl_format ¶

format for writing reals to star_history_name file

```
star_history_dbl_format = '(1pes40.16e3, 1x)'
```

star_history_int_format ¶

format for writing integer to star_history_name file

```
star_history_int_format = '(i40, 1x)'
```

star_history_txt_format ¶

format for writing characters to star_history_name file

```
star_history_txt_format = '(a40, 1x)'
```

write_profiles_flag ¶

profiles are written only if this is true

```
write_profiles_flag = .true.
```

profile_interval ¶

save a model profile info when $\text{mod}(\text{model_number}, \text{profile_interval}) = 0$.

```
profile_interval = 50
```

priority_profile_interval ¶

give saved profile a higher priority for retention when $\text{mod}(\text{model_number}, \text{priority_profile_interval}) = 0$.

```
priority_profile_interval = 1000
```

profiles_index_name ¶

name of the profile index file

```
profiles_index_name = 'profiles.index'
```

profiles_data_prefix ¶

prefix of the profile data

```
profile_data_prefix = 'profile'
```

profiles_data_suffix ¶

suffix of the profile data

```
profile_data_suffix = '.data'
```

profile_data_header_suffix ¶

If not empty, then put profile data header info here. In this case the profile data file has only data, making it easier to use with some plotting packages.

```
profile_data_header_suffix = ''
```

profile_dbl_format ¶

format for writing reals to profile file

```
profile_dbl_format = '(1pes40.16e3, 1x)'
```

profile_int_format ¶

format for writing integers to profile file

```
profile_int_format = '(i40, 1x)'
```

profile_txt_format ¶

format for writing characters to profile file

```
profile_txt_format = '(a40, 1x)'
```

max_num_profile_zones ¶

if `nz > this`, then only write a subsample of the zones. only used if `> 1`

```
max_num_profile_zones = -1
```

max_num_profile_models ¶

Maximum number of saved profiles. If there's no limit on the number of profiles saved, you can fill up your disk – I've done it. So it's a good idea to set this limit to a reasonable number such as 20 or 30. Once that many have been saved during a run, old ones will be discarded to make room for new ones. Profiles that were saved for key events are given priority and aren't removed as long as there is a lower priority profile that can be discarded instead. Less than zero means no limit.

```
max_num_profile_models = 100
```

profile_model ¶

save profile when `model_number` equals this

```
profile_model = -1
```

write_model_with_profile ¶

if this is true, models are written at same time as profiles

```
write_model_with_profile = .false.
```

model_data_prefix ¶

prefix of the model data files

```
model_data_prefix = 'profile'
```

model_data_suffix ¶

suffix of the model data files

```
model_data_suffix = '.mod'
```

write_controls_info_with_profile ¶

if this is true, the values of the options in the controls inlist are written at same time as profiles

```
write_controls_info_with_profile = .false.
```

controls_data_prefix ¶

prefix of the control data files

```
controls_data_prefix = 'controls'
```

controls_data_suffix ¶

suffix of the control data files

```
controls_data_suffix = '.data'
```

mixing_D_limit_for_log ¶

if max D_{mix} in mixing region is less than this, don't include the region in the log doesn't apply to thermohaline or semiconvective regions

```
mixing_D_limit_for_log = 1d4
```

eta_RTI_limit_for_log ¶

if eta_{RTI} is less than this, treat it as zero for history log

```
eta_RTI_limit_for_log = 1d4
```

mass_loc_for_extra_log_info ¶

log contains info about this mass location in the model negative value means “don’t bother”

```
mass_loc_for_extra_log_info = -1
```

write_pulse_data_with_profile ¶

if true, write pulse info file when writing profile

```
write_pulse_data_with_profile = .false.
```

pulse_data_format ¶

pulsation code format, e.g., ‘FGONG’, ‘OSC’, ‘GYRE’

```
pulse_data_format = 'FGONG'
```

add_atmosphere_to_pulse_data ¶

if true, write atmosphere to pulse files

```
add_atmosphere_to_pulse_data = .false.
```

add_center_point_to_pulse_data ¶

if true, add point for r=0 to pulse files

```
add_center_point_to_pulse_data = .true.
```

keep_surface_point_for_pulse_data ¶

if true, add k=1 cell to pulse files

```
keep_surface_point_for_pulse_data = .false.
```

add_double_points_to_pulse_data ¶

add double points at discontinuities

```
add_double_points_to_pulse_data = .false.
```

interpolate_rho_for_pulse_data ¶

If true, then get `rho_face` by interpolating `rho` at cell center. If false, then calculate `rho_face` by $dm/(4*\pi*r^2*dr)$.

```
interpolate_rho_for_pulse_data = .true.
```

threshold_grad_mu_for_double_point

threshold in `grad_mu = dln(mu)/dln(P)` for a double point to be written

```
threshold_grad_mu_for_double_point = 10d0
```

max_number_of_double_points

maximum number of double points to be written (0 = no limit); when this limit is set, double points are chosen in order of decreasing `|grad_mu|`

```
max_number_of_double_points = 0
```

format_for_FGONG_data

This is the ‘wide’ FGONG format, as agreed on at the 5th Aarhus RGB workshop (University of Birmingham, UK, October 2015)

```
format_for_FGONG_data = '(1P,5(X,E26.18E3))'
```

format_for_OSC_data

[FGONG Format Documentation] (http://www.astro.up.pt/corot/ntools/docs/CoRoT_ESTA_Files.pdf)

```
format_for_OSC_data = '(1P5E19.12,x)'
```

write_pulsation_plot_data

if true and saving pulsation info, also write out text file in column format for plotting

```
write_pulsation_plot_data = .false.
```

max_num_gyre_points

limit gyre output files to at most this number of points only used when > 1

```
max_num_gyre_points = -1
```

fgong_zero_A_inside_r

when writing FGONG, if $r < \text{this}$ and cell has mixing of some kind, force $A = 0$ R_{sun} units

```
fgong_zero_A_inside_r = 0d0
```

trace_mass_location

location for trace_mass_radius, trace_mass_logT, etc. (M_{sun} units)

```
trace_mass_location = 0
```

min_tau_for_max_abs_v_location

controls choice of location in model for max_abs_v history info. can use this to exclude locations too close to surface. ignore if ≤ 0

```
min_tau_for_max_abs_v_location = 0
```

min_q_for_inner_mach1_location

controls choice of location in model for innermost mach 1 history info. can use this to exclude locations too close to center.

```
min_q_for_inner_mach1_location = 0
```

max_q_for_outer_mach1_location

controls choice of location in model for outermost mach 1 history info. can use this to exclude locations too close to surface.

```
max_q_for_outer_mach1_location = 1
```

burn_min1

used for reporting where burning zone occur, for example in the pgstar TRho profiles. see star/public/star_data.inc for details. must be $< \text{burn_min2}$. In ergs/g/sec.

```
burn_min1 = 50
```

burn_min2

used for reporting where burning zone occur, for example in the pgstar TRho profiles. see star/public/star_data.inc for details. In ergs/g/sec.

```
burn_min2 = 1000
```

max_conv_vel_div_csound_maxq ¶

only consider from center out to this location

```
max_conv_vel_div_csound_maxq = 1
```

width_for_limit_conv_vel ¶

look this number of cells on either side of boundary to see if any boundary k in that range has $s\% \text{csound}(k) < s\% v(k) \leq s\% \text{csound}(k-1)$ i.e. transition from subsonic to supersonic as go inward if find any such transition then don't allow increase in convection velocity. this implies no change from radiative to convective. the purpose of this is to prevent convective energy transport from moving energy from behind a shock to in front of the shock.

```
width_for_limit_conv_vel = 3
```

max_q_for_limit_conv_vel ¶

for $q(k) \leq$ this, don't allow conv_vel to grow

```
max_q_for_limit_conv_vel = -1
```

max_r_in_cm_for_limit_conv_vel ¶

for $r(k) \leq$ this, don't allow conv_vel to grow

```
max_r_in_cm_for_limit_conv_vel = -1
```

max_mass_in_gm_for_limit_conv_vel ¶

for $m(k) \leq$ this, don't allow conv_vel to grow

```
max_mass_in_gm_for_limit_conv_vel = -1
```

center_avg_value_dq ¶

reported center values are averages over this fraction of star mass

```
center_avg_value_dq = 1d-8
```

surface_avg_abundance_dq ¶

reported surface abundances are averages over this fraction of star mass

```
surface_avg_abundance_dq = 1d-8
```

mass_depth_for_L_surf ¶

only if use_flux_limiting_with_dPrad_dm_form

```
mass_depth_for_L_surf = 0d0
```

conv_core_gap_dq_limit ¶

skip non-convective gaps of less than this limit when reporting convective core size

```
conv_core_gap_dq_limit = 0d0
```

definition of core overshooting boundary for output ¶

alpha_bdy_core_overshooting ¶

```
bdy = core_overshoot_r0 + core_overshoot_Hp* &
      (alpha_bdy_core_overshooting*core_overshoot_f - core_overshoot_f0)
```

```
alpha_bdy_core_overshooting = 5
```

definition of core boundaries ¶

he_core_boundary_h1_fraction ¶

If ≥ 0 , boundary is outermost location where h1 mass fraction is \leq this value, and he4 mass fraction \geq min_boundary_fraction (see below). If < 0 , boundary is outermost location where he4 is the most abundant species.

```
he_core_boundary_h1_fraction = 0.01d0
```

c_core_boundary_he4_fraction ¶

If ≥ 0 , boundary is outermost location where he4 mass fraction is \leq this value, and c12 mass fraction \geq min_boundary_fraction (see below). If < 0 , boundary is outermost location where c12 is the most

abundant species.

```
c_core_boundary_he4_fraction = 0.01d0
```

o_core_boundary_c12_fraction ¶

If ≥ 0 , boundary is outermost location where c12 mass fraction is \leq this value, and o16 mass fraction \geq `min_boundary_fraction` (see below). If < 0 , boundary is outermost location where o16 is the most abundant species.

```
o_core_boundary_c12_fraction = 0.01d0
```

si_core_boundary_o16_fraction ¶

If ≥ 0 , boundary is outermost location where o16 mass fraction is \leq this value, and si28 mass fraction \geq `min_boundary_fraction` (see below). If < 0 , boundary is outermost location where si28 is the most abundant species.

```
si_core_boundary_o16_fraction = 0.01d0
```

fe_core_boundary_si28_fraction ¶

For this case, “iron” includes any species with $A > 46$. If ≥ 0 , boundary is outermost location where si28 mass fraction is \leq this value, and “iron” mass fraction \geq `min_boundary_fraction` (see below). If < 0 , boundary is outermost location where “iron” is the most abundant species.

```
fe_core_boundary_si28_fraction = 0.01d0
```

neutron_rich_core_boundary_Ye_max ¶

Boundary is outermost location where Ye is \leq this value.

```
neutron_rich_core_boundary_Ye_max = 0.48d0
```

min_boundary_fraction ¶

Value for deciding boundary regions.

```
min_boundary_fraction = 0.1d0
```

when to stop ¶

max_model_number ¶

The code will stop when it reaches this model number. Negative means no maximum.

```
max_model_number = -1
```

when_to_stop_rtol ¶

Relative error criteria when hitting stop target time. The system will automatically redo with a smaller timestep to hit a stopping target. It calculates the following “error” term and retries if it is > 1 .

```
error = abs(value - target_value)/ &  
        (when_to_stop_atol + when_to_stop_rtol*max(abs(value), abs(target_value)))
```

```
when_to_stop_rtol = 1d99
```

when_to_stop_atol ¶

Absolute error criteria when hitting stop target time. The system will automatically redo with a smaller timestep to hit a stopping target. It calculates the following “error” term and retries if it is > 1 .

```
error = abs(value - target_value)/ &  
        (when_to_stop_atol + when_to_stop_rtol*max(abs(value), abs(target_value)))
```

```
when_to_stop_atol = 1d99
```

max_age ¶

Stop when the age of the star exceeds this value (in years). only applies when > 0 .

```
max_age = 1d36
```

max_age_in_seconds ¶

Stop when the age of the star exceeds this value (in seconds). only applies when > 0 .

```
max_age_in_seconds = -1
```

num_adjusted_dt_steps_before_max_age ¶

This adjusts `max_years_for_timestep` so that hit `max_age` exactly, without needing possibly large change in timestep at end of run. only used if > 0

number of time steps to adjust to prior to hitting max age only used if > 0

```
num_adjusted_dt_steps_before_max_age = 0
```

dt_years_for_steps_before_max_age ¶

timestep in years

```
dt_years_for_steps_before_max_age = 1d6
```

reduction_factor_for_max_timestep ¶

per time step reduction limited to this

```
reduction_factor_for_max_timestep = 0.98d0
```

gamma_center_limit ¶

gamma is the plasma interaction parameter. Stop when the center value of gamma exceeds this limit.

```
gamma_center_limit = 1d99
```

eta_center_limit ¶

eta is the electron chemical potential in units of $k*T$. Stop when the center value of eta exceeds this limit.

```
eta_center_limit = 1d99
```

log_center_density_limit ¶

Stop when log10 of the center density exceeds this limit.

```
log_center_density_limit = 11
```

log_center_density_lower_limit ¶

Stop when log10 of the center density is below this limit.

```
log_center_density_lower_limit = -1d99
```

log_center_temp_limit ¶

Stop when log10 of the center temperature exceeds this limit.

```
log_center_temp_limit = 11
```

log_center_temp_lower_limit ¶

Stop when log10 of the center temperature is below this limit.

```
log_center_temp_lower_limit = -1d99
```

surface_accel_div_grav_limit = -1 ¶

This is used when do not have a velocity variable. The acceleration ratio is $\text{abs}(\text{accel})/\text{grav}$ at surface, where accel is $(\text{rdot} - \text{rdot_old})/\text{dt}$ and grav is $G \cdot m / r^2$. Stop if the ratio becomes larger than this limit. Ignored if ≤ 0 .

```
surface_accel_div_grav_limit = -1
```

log_max_temp_upper_limit ¶

stop when log10 of the maximum temperature rises above this limit.

```
log_max_temp_upper_limit = 99
```

log_max_temp_lower_limit ¶

stop when log10 of the maximum temperature drops below this limit.

```
log_max_temp_lower_limit = -99
```

center_entropy_limit ¶

stop when the center entropy exceeds this limit. in kerg per baryon

```
center_entropy_limit = 1d99
```

center_entropy_lower_limit ¶

stop when the center entropy is below this limit. in kerg per baryon

```
center_entropy_lower_limit = -1d99
```

max_entropy_limit ¶

stop when the max entropy exceeds this limit. in kerg per baryon

```
max_entropy_limit = 1d99
```

max_entropy_lower_limit ¶

stop when the max entropy is below this limit. in kerg per baryon

```
max_entropy_lower_limit = -1d99
```

xa_central_lower_limit_species ¶

xa_central_lower_limit ¶

Lower limits on central mass fractions. Stop when central abundance drops below this limit. Can have up to `num_xa_central_limits` of these (see `star_def.inc` for value).

`xa_central_lower_limit_species` contains an isotope name as defined in `chem_def.f`.

`xa_central_lower_limit` contains the lower limit value.

```
xa_central_lower_limit_species(1) = ''
xa_central_lower_limit(1) = 0
```

xa_central_upper_limit_species ¶

xa_central_upper_limit ¶

Upper limits on central mass fractions. Stop when central abundance rises above this limit. Can have up to `num_xa_central_limits` of these (see `star_def.inc` for value). E.g., to stop when center c12 abundance reaches 0.5, set

```
xa_central_upper_limit_species(1) = 'c12'
xa_central_upper_limit(1) = 0.5
```

```
xa_central_upper_limit_species(1) = ''
xa_central_upper_limit(1) = 0
```

xa_surface_lower_limit_species ¶

xa_surface_lower_limit ¶

Lower limits on surface mass fractions. Stop when surface abundance drops below this limit. Can have up to `num_xa_surface_limits` of these (see `star_def` for value)

`xa_surface_lower_limit_species` contains an isotope name as defined in `chem_def.f`

`xa_surface_lower_limit` contains the lower limit value

```
xa_surface_lower_limit_species(1) = ''
xa_surface_lower_limit(1) = 0
```

xa_surface_upper_limit_species

xa_surface_upper_limit

upper limits on surface mass fractions stop when surface abundance rises above this limit can have up to `num_xa_surface_limits` of these (see `star_def` for value) e.g., to stop when surface c12 abundance reaches 0.5, set

```
xa_surface_upper_limit_species(1) = 'c12'
xa_surface_upper_limit(1) = 0.5
```

```
xa_surface_upper_limit_species(1) = ''
xa_surface_upper_limit(1) = 0
```

xa_average_lower_limit_species

xa_average_lower_limit

lower limits on average mass fractions stop when average abundance drops below this limit can have up to `num_xa_average_limits` of these (see `star_def` for value)

```
xa_average_lower_limit_species(1) = ''
xa_average_lower_limit(1) = 0
```

xa_average_upper_limit_species

xa_average_upper_limit

upper limits on average mass fractions stop when average abundance rises above this limit can have up to `num_xa_average_limits` of these (see `star_def` for value)

```
xa_average_upper_limit_species(1) = ''
xa_average_upper_limit(1) = 0
```

HB_limit

For detecting horizontal branch. Only applies when center abundance by mass of h1 is $< 1d-4$. Stop when the center abundance by mass of he4 drops below this limit.

```
HB_limit = 0
```

stop_at_TP

If true, stop at next AGB thermal pulse. This is defined as having a convective zone with helium burning when central helium is depleted and $\text{he_core_mass} - \text{c_core_mass} \leq \text{TP_he_shell_max}$.

```
stop_at_TP = .false.
```

TP_he_shell_max ¶

Stop when thermal pulse helium shell mass reaches this value, in Msun units

```
TP_he_shell_max = 0.2d0
```

star_mass_min_limit ¶

Stop when star mass in Msun units is < this. <= 0 means no limit.

```
star_mass_min_limit = 0
```

star_mass_max_limit ¶

Stop when star mass in Msun units is > this. <= 0 means no limit.

```
star_mass_max_limit = 0
```

star_H_mass_min_limit ¶

Stop when star hydrogen mass in Msun units is < this. <= 0 means no limit.

```
star_H_mass_min_limit = 0
```

star_H_mass_max_limit ¶

Stop when star hydrogen mass in Msun units is > this. <= 0 means no limit.

```
star_H_mass_max_limit = 0
```

star_He_mass_min_limit ¶

Stop when star he3+he4 mass in Msun units is < this. <= 0 means no limit.

```
star_He_mass_min_limit = 0
```

star_He_mass_max_limit = 0 ¶

Stop when star he3+he4 mass in Msun units is > this. <= 0 means no limit.

```
star_He_mass_max_limit = 0
```

star_C_mass_min_limit ¶

Stop when star c12 mass in Msun units is < this. <= 0 means no limit.

```
star_C_mass_min_limit = 0
```

star_C_mass_max_limit = 0 ¶

Stop when star c12 mass in Msun units is > this. <= 0 means no limit.

```
star_C_mass_max_limit = 0
```

envelope_mass_limit ¶

```
envelope_mass = star_mass - he_core_mass
```

Stop when envelope_mass drops below this limit, in Msun units.

```
envelope_mass_limit = 0
```

envelope_fraction_left_limit ¶

```
envelope_fraction_left = (star_mass - he_core_mass)/(initial_mass - he
```

```
envelope_fraction_left_limit = 0
```

xmstar_min_limit ¶

! xmstar = mstar - M_center stop when xmstar in grams is < this. <= 0 means no limit.

```
xmstar_min_limit = 0
```

xmstar_max_limit ¶

xmstar = mstar - M_center stop when xmstar in grams is > this. <= 0 means no limit.

```
xmstar_max_limit = 0
```

he_core_mass_limit ¶

stop when helium core reaches this mass, in Msun units

```
he_core_mass_limit = 1d99
```

c_core_mass_limit ¶

stop when carbon core reaches this mass, in Msun units

```
c_core_mass_limit = 1d99
```

o_core_mass_limit ¶

stop when oxygen core reaches this mass, in Msun units

```
o_core_mass_limit = 1d99
```

si_core_mass_limit ¶

stop when silicon core reaches this mass, in Msun units

```
si_core_mass_limit = 1d99
```

fe_core_mass_limit ¶

stop when iron core reaches this mass, in Msun units

```
fe_core_mass_limit = 1d99
```

neutron_rich_core_mass_limit ¶

stop when neutron rich core reaches this mass, in Msun units

```
neutron_rich_core_mass_limit = 1d99
```

he_layer_mass_lower_limit ¶

he layer mass is defined as $\text{he_core_mass} - \text{c_core_mass}$ stop when $\text{c_core_mass} > 0$ and he layer mass $<$ this limit (Msun units).

```
he_layer_mass_lower_limit = 0
```

abs_diff_lg_LH_lg_Ls_limit ¶

stop when $\text{abs}(\text{lg_LH} - \text{lg_Ls}) \leq \text{abs_diff_LH_Lsurf_limit}$ can be useful for deciding when pre-main sequence star has reached ZAMS set to negative value to disable

```
abs_diff_lg_LH_lg_Ls_limit = -1
```

Teff_upper_limit ¶

stop when Teff is greater than this limit.

```
Teff_upper_limit = 1d99
```

Teff_lower_limit ¶

stop when Teff is less than this limit.

```
Teff_lower_limit = -1d99
```

photosphere_r_upper_limit ¶

stop when photosphere_r is greater than this limit, in R_{sun} units

```
photosphere_r_upper_limit = 1d99
```

photosphere_r_lower_limit ¶

stop when photosphere_r is less than this limit, in R_{sun} units

```
photosphere_r_lower_limit = -1d99
```

photosphere_m_upper_limit ¶

stop when photosphere_m is greater than this limit, in M_{sun} units

```
photosphere_m_upper_limit = 1d99
```

photosphere_m_lower_limit ¶

stop when photosphere_m is less than this limit, in M_{sun} units


```
photosphere_m_lower_limit = -1d99
```

photosphere_m_sub_M_center_limit ¶

stop when photosphere_m is less than this limit above M_center, in Msun units

```
photosphere_m_sub_M_center_limit = -1d99
```

log_Teff_upper_limit ¶

stop when log10 of Teff is greater than this limit.

```
log_Teff_upper_limit = 1d99
```

log_Teff_lower_limit ¶

stop when log10 of Teff is less than this limit.

```
log_Teff_lower_limit = -1d99
```

log_Tsurf_upper_limit ¶

stop when log10 of T in outermost cell is greater than this limit.

```
log_Tsurf_upper_limit = 1d99
```

log_Tsurf_lower_limit ¶

stop when log10 of T in outermost cell is less than this limit.

```
log_Tsurf_lower_limit = -1d99
```

log_L_upper_limit ¶

stop when log10(total luminosity in Lsun units) is greater than this limit. in order to skip pre-ms, this limit only applies when $L_{\text{nuc}} > 0.01 * L$

```
log_L_upper_limit = 1d99
```

log_L_lower_limit ¶

stop when log10(total luminosity in Lsun units) is less than this limit.

```
log_L_lower_limit = -1d99
```

log_g_upper_limit ¶

stop when $\log_{10}(\text{gravity at surface})$ is greater than this limit.

```
log_g_upper_limit = 1d99
```

log_g_lower_limit ¶

stop when $\log_{10}(\text{gravity at surface})$ is less than this limit.

```
log_g_lower_limit = -1d99
```

log_Psurf_upper_limit ¶

stop when \log_{10} of surface pressure is greater than this limit.

```
log_Psurf_upper_limit = 1d99
```

log_Psurf_lower_limit ¶

stop when \log_{10} of surface pressure is less than this limit.

```
log_Psurf_lower_limit = -1d99
```

log_Dsurf_upper_limit ¶

stop when \log_{10} of surface density is greater than this limit.

```
log_Dsurf_upper_limit = 1d99
```

log_Dsurf_lower_limit ¶

stop when \log_{10} of surface density is less than this limit.

```
log_Dsurf_lower_limit = -1d99
```

power_nuc_burn_upper_limit ¶

stop when total power from all nuclear reactions (in L_{sun} units) is $>$ this.

```
power_nuc_burn_upper_limit = 1d99
```

power_h_burn_upper_limit ¶

stop when total power from hydrogen-consuming reactions (in Lsun units) is > this.

```
power_h_burn_upper_limit = 1d99
```

power_he_burn_upper_limit ¶

stop when total power from reactions burning helium (in Lsun units) is > this.

```
power_he_burn_upper_limit = 1d99
```

power_c_burn_upper_limit ¶

stop when total power from reactions burning carbon (in Lsun units) is > this

```
power_c_burn_upper_limit = 1d99
```

power_nuc_burn_lower_limit ¶

stop when total power from all nuclear reactions (in Lsun units) is < this.

```
power_nuc_burn_lower_limit = -1d99
```

power_h_burn_lower_limit ¶

stop when total power from hydrogen consuming reactions (in Lsun units) is < this.

```
power_h_burn_lower_limit = -1d99
```

power_he_burn_lower_limit ¶

stop when total power from reactions burning helium (in Lsun units) is < this.

```
power_he_burn_lower_limit = -1d99
```

power_c_burn_lower_limit ¶

stop when total power from reactions burning carbon (in Lsun units) is < this.

```
power_c_burn_lower_limit = -1d99
```

max_number_backups ¶

Stop if the number of backups exceeds this value. Ignore if < 0 .

```
max_number_backups = -1
```

max_number_retries ¶

Stop if the number of retries exceeds this value. Ignore if < 0 .

```
max_number_retries = -1
```

max_backups_in_a_row ¶

if do more than this many without a successful step, then terminate the run.

```
max_backups_in_a_row = 15
```

relax_max_number_backups ¶

Stop if the number of backups during a “relax” evolution exceeds this value. ignore if < 0

```
relax_max_number_backups = 100
```

relax_max_number_retries ¶

Stop if the number of retries during a “relax” evolution exceeds this value. ignore if < 0

```
relax_max_number_retries = 300
```

min_timestep_limit ¶

stop if need timestep smaller than this limit, in seconds

```
min_timestep_limit = 1d-6
```

logQ_limit ¶

$\log Q = \log \rho - 2 * \log T + 12$. stop if $\log Q$ at any zone is larger than this limit. 5 is a reasonable limit for the current mesa/eos.

```
logQ_limit = 5d0
```

logQ_min_limit ¶

$\log Q = \log \rho - 2 \log T + 12$. stop if $\log Q$ at any zone is smaller than this limit. -10 is a reasonable limit for the current mesa/eos.

```
logQ_min_limit = -10d0
```

center_Ye_lower_limit ¶

stop if center_ye drops below this limit

```
center_Ye_lower_limit = -1
```

center_R_lower_limit ¶

stop if R_center drops below this limit (in cm)

```
center_R_lower_limit = -1
```

fe_core_infall_limit ¶

stop if max infall velocity at any location interior to fe_core_mass, in cm/s

```
fe_core_infall_limit = 1d8
```

non_fe_core_infall_limit ¶

stop if max infall velocity at any location interior to he_core_mass. in cm/s

```
non_fe_core_infall_limit = 1d99
```

v_div_csound_max_limit ¶

stop if any $v/c_{\text{sound}} >$ this limit

```
v_div_csound_max_limit = 1d99
```

v_div_csound_surf_limit ¶

stop if $v_{\text{surf}}/c_{\text{sound_surf}} >$ this limit

```
v_div_csound_surf_limit = 1d99
```

v_surf_div_v_kh_upper_limit ¶

stop if $\text{abs}(v_{\text{surf}}/v_{\text{kh}}) > \text{this limit}$, where $v_{\text{kh}} = \text{photosphere_r}/\text{kh_timescale}$

```
v_surf_div_v_kh_upper_limit = 1d99
```

v_surf_div_v_kh_lower_limit ¶

stop if $\text{abs}(v_{\text{surf}}/v_{\text{kh}}) < \text{this limit}$, where $v_{\text{kh}} = \text{photosphere_r}/\text{kh_timescale}$

```
v_surf_div_v_kh_lower_limit = -1d99
```

v_surf_div_v_esc_limit ¶

stop if $v_{\text{surf}}/v_{\text{esc}} > \text{this limit}$

```
v_surf_div_v_esc_limit = 1d99
```

v_surf_kms_limit ¶

stop if v_{surf} in km/s $> \text{this limit}$

```
v_surf_kms_limit = 1d99
```

Lnuc_div_L_zams_limit ¶

defines “near zams” – note: must also set `stop_near_zams`

```
Lnuc_div_L_zams_limit = 0.9d0
```

stop_near_zams ¶

if true, stop if $\text{Lnuc}/L > \text{Lnuc_div_L_zams_limit}$

```
stop_near_zams = .false.
```

Lnuc_div_L_upper_limit ¶

stop when Lnuc/L is greater than this limit.

```
Lnuc_div_L_upper_limit = 1d99
```

Lnuc_div_L_lower_limit ¶

stop when Lnuc/L is less than this limit.

```
lnuc_div_L_lower_limit = -1d99
```

Pgas_div_P_limit ¶

criteria for stopping on Pgas/P

```
Pgas_div_P_limit = 0
```

Pgas_div_P_limit_max_q ¶

stop if Pgas/P < this limit at any location with $q \leq \text{Pgas_div_P_limit_max_q}$ values near unity skip the outer envelope

```
Pgas_div_P_limit_max_q = 0.95d0
```

peak_burn_vconv_div_cs_limit ¶

limits ratio of convection velocity to sound speed at location of peak `eps_nuc`

```
peak_burn_vconv_div_cs_limit = 1d99
```

omega_div_omega_crit_limit ¶

stop if $\omega/\omega_{\text{crit}}$ is > this anywhere in star ignore if < 0

```
omega_div_omega_crit_limit = -1
```

delta_nu_lower_limit ¶

stop when asteroseismology $\delta\nu$ in micro Hz is < this. ≤ 0 means no limit.

```
delta_nu_lower_limit = 0
```

delta_nu_upper_limit ¶

stop when asteroseismology $\delta\nu$ in micro Hz is > this. ≤ 0 means no limit.

```
delta_nu_upper_limit = 0
```

shock_mass_upper_limit ¶

stop when shock_mass is > this. ≤ 0 means no limit.

```
shock_mass_upper_limit = -1
```

mach1_mass_upper_limit

stop when outer location of mach 1 is > this. <= 0 means no limit.

```
mach1_mass_upper_limit = -1
```

delta_Pg_lower_limit

stop when delta_Pg in micro Hz is < this. <= 0 means no limit.

```
delta_Pg_lower_limit = 0
```

delta_Pg_upper_limit

stop when delta_Pg in micro Hz is > this. <= 0 means no limit.

```
delta_Pg_upper_limit = 0
```

stop_when_reach_this_cumulative_extra_heating

(ignore if <= 0)

```
stop_when_reach_this_cumulative_extra_heating = 0d0
```

mixing parameters

mixing_length_alpha

The mixing length is this parameter times a local pressure scale height. To increase R vs. L, decrease mixing_length_alpha.

```
mixing_length_alpha = 2
```

remove_small_D_limit

If MLT diffusion coeff D (cm²/sec) is less than this limit, then set D to zero and change the point to mixing_type == no_mixing.

```
remove_small_D_limit = 1d-6
```

use_Ledoux_criterion

a location in the model is Schwarzschild stable when $\text{gradr} < \text{grada}$ it is Ledoux stable when $\text{gradr} < \text{gradL}$, where $\text{gradL} = \text{grada} + \text{composition_gradient}$ note that these are the same when $\text{composition_gradient} = 0$ so you can force the use of the Schwarzschild criterion by passing 0 for the `composition_gradient` argument to the `mlt` routine. that's what happens if you set the control `"use_Ledoux_criterion"` to false.

overshooting and rotational mixing are dealt with separately and are added after the MLT classifications are made.

```
use_Ledoux_criterion = .false.
```

num_cells_for_smooth_gradL_composition_term ¶

Number of cells on either side to use in weighted smoothing of `gradL_composition_term`. `gradL_composition_term` is set to the "raw" unsmoothed `brunt_B` and then optionally smoothed according `num_cells_for_smooth_gradL_composition_term`. In cases where the Ledoux criterion is used to evaluate the boundary for burning convective cores, you may need to set `num_cells_for_smooth_gradL_composition_term = 0` to avoid smoothing the stabilizing composition jump into the convection zone and unphysically causing it to shrink. See section 3.2 in Moore, K., & Garaud, P. 2016, APJ, 817, 54

```
num_cells_for_smooth_gradL_composition_term = 3
```

threshold_for_smooth_gradL_composition_term ¶

Threshold for weighted smoothing of `gradL_composition_term`. Only apply smoothing (controlled by `num_cells_for_smooth_gradL_composition_term`) for contiguous regions where $|\text{gradL}|$ exceeds this threshold. Might be useful for preventing narrow composition jumps from being excessively broadened by smoothing

```
threshold_for_smooth_gradL_composition_term = 0
```

alpha_semiconvection ¶

Determines efficiency of semiconvective mixing. Semiconvection only applies if `use_Ledoux_criterion` is true.

```
alpha_semiconvection = 0
```

semiconvection_upper_limit_center_h1 ¶

Turn off semiconvection when `center_h1` > this limit. This let's you delay semiconvection until helium burning. E.g., you can do overshooting for core hydrogen burning, then switch to semiconvection after core h is gone.

```
semiconvection_upper_limit_center_h1 = 1d99
```

semiconvection_option ¶

- 'Langer_85 mixing; gradT = gradr' : uses Langer scheme for mixing but sets gradT = gradr
- 'Langer_85' : this calculates special gradT as well as doing mixing.

```
semiconvection_option = 'Langer_85 mixing; gradT = gradr'
```

thermohaline_coeff ¶

Determines efficiency of thermohaline mixing. was previously named thermo_haline_coeff. thermohaline mixing only applies if use_Ledoux_criterion is true.

```
thermohaline_coeff = 0
```

thermohaline_option ¶

determines which method to use for calculating thermohaline diffusion coef:

- 'Kippenhahn' : use method of Kippenhahn, R., Ruschenplatt, G., & Thomas, H.-C. 1980, A&A, 91, 175.
- 'Traxler_Garaud_Stellmach_11' : use method of Traxler, Garaud, & Stellmach, ApJ Letters, 728:L29 (2011).
- 'Brown_Garaud_Stellmach_13' : use method of Brown, Garaud, & Stellmach, (2013). Recommends thermohaline_coeff = 1, but it can nevertheless be changed.

```
thermohaline_option = 'Kippenhahn'
```

alt_scale_height_flag ¶

If false, then stick to the usual definition – $P/(g \cdot \rho)$. If true, use min of the usual and sound speed * hydro time scale, $\sqrt{P/G}/\rho$.

```
alt_scale_height_flag = .true.
```

mlt_use_rotation_correction ¶

When doing rotation, multiply grad_rad by ft_rot/ft_rot if this flag is true.

```
mlt_use_rotation_correction = .true.
```

MLT_option ¶

Options are:

- 'none' : just give radiative values with no mixing.
- 'Cox' : MLT as developed in Cox & Giuli 1968, Chapter 14.
- 'ML1' : Bohm-Vitense 1958

- ‘ML2’ : Bohm and Cassinelli 1971
- ‘Mihalas’ : Mihalas 1978, Kurucz 1979
- ‘Henyey’ : Henyey, Vardya, and Bodenheimer 1965

‘Cox’ option assumes optically thick material. The other options are various ways of extending to include optically thin material.

```
MLT_option = 'Cox'
```

Henyey_MLT_y_param ¶

Henyey_MLT_nu_param ¶

Values of the $f_1..f_4$ coefficients are taken from Table 1 of Ludwig et al. 1999, A&A, 346, 111 with the following exception: their value of f_3 for Henyey convection is $f_4/8$ when it should be $8*f_4$, i.e., $f_3=32*\pi^{**2/3}$ and $f_4=4*\pi^{**2/3}$. f_3 and f_4 are related to the henyey y parameter, so for the ‘Henyey’ case they are set based on the value of Henyey_y_param.

```
Heney_MLT_y_param = 0.33333333d0
Heney_MLT_nu_param = 8
```

make_gradr_sticky_in_newton_iters ¶

if true, then location that becomes radiative during newton iterations, stays radiative for rest of the newton iterations. to avoid flip-flopping between radiative and convective.

```
make_gradr_sticky_in_newton_iters = .false.
```

no_MLT_below_shock ¶

if true, then no MLT below an outward going shock (just radiative).

```
no_MLT_below_shock = .false.
```

no_MLT_below_T_max ¶

if true, then no MLT below location of max T (just radiative).

```
no_MLT_below_T_max = .false.
```

T_mix_limit ¶

If there is any convection in surface zones with $T < T_mix_limit$, then extend the innermost such convective region outward all the way to the surface. For example,

- $T_mix_limit \leq 0$ means omit this operation.
- $T_mix_limit = 1d5$ will effectively make the star convective down to the He++ region.

units in Kelvin

```
T_mix_limit = 0
```

conv_dP_term_factor ¶

Set to 0 to turn off effect of pressure from convective turbulence. The convective turbulence factor is based on Cox&Giuli (14.69) Multiplier for conv_dP_term P is increased by factor $(1 + \text{conv_dP_term})$ by inclusion of convective turbulence.

```
conv_dP_term_factor = 0
```

mlt_gradT_fraction ¶

let $f := \text{mlt_gradT_fraction}$ if $f \geq 0$ and $f \leq 1$, then gradT from mlt is replaced by $f * \text{grada_at_face}(k) + (1-f) * \text{gradr}(k)$ see also the vector control `adjust_mlt_gradT_fraction` for fine grain control

```
mlt_gradT_fraction = -1
```

okay_to_reduce_gradT_excess ¶

`gradT_excess` = `gradT_sub_grada` = superadiabaticity.

Inefficient convection => large gradT excess and steep T gradient to enhance radiative transport. Reduce gradT excess by making gradT closer to adiabatic gradient. If true, code is allowed to adjust gradT to boost efficiency of energy transport See `gradT_excess_f1`, `gradT_excess_f2`, and `gradT_excess_age_fraction` below.

```
okay_to_reduce_gradT_excess = .false.
```

gradT_excess_f1 ¶

gradT_excess_f2 ¶

These are for calculation of efficiency boosted gradT.

```
gradT_excess_f1 = 1d-4
gradT_excess_f2 = 1d-3
```

gradT_excess_age_fraction ¶

These are for calculation of efficiency boosted gradT. Fraction of old to mix with new to get next.

```
gradT_excess_age_fraction = 0.9d0
```

gradT_excess_max_change ¶

These are for calculation of efficiency boosted gradT. Maximum change allowed in one timestep for gradT_excess_alpha. Ignored if negative.

```
gradT_excess_max_change = -1d0
```

gradT_excess_lambda1 ¶**gradT_excess_beta1 ¶**

In some situations you might want to force $\alpha = 1$. You can do that by setting gradT_excess_lambda1 < 0. The following are for the normal calculation of gradT_excess_alpha

```
gradT_excess_lambda1 = 1.0d0
gradT_excess_beta1 = 0.35d0
```

gradT_excess_lambda2 ¶**gradT_excess_beta2 ¶**

The following are for the normal calculation of gradT_excess_alpha.

```
gradT_excess_lambda2 = 0.5d0
gradT_excess_beta2 = 0.25d0
```

gradT_excess_dlambda ¶**gradT_excess_dbeta ¶**

The following are for the normal calculation of gradT_excess_alpha.

```
gradT_excess_dlambda = 0.1d0
gradT_excess_dbeta = 0.1d0
```

gradT_excess_max_center_h1 ¶

No boost if center H1 > this limit.

```
gradT_excess_max_center_h1 = 1d0
```

gradT_excess_min_center_he4 ¶

No boost if center He4 < this limit.

```
gradT_excess_min_center_he4 = 0d0
```

gradT_excess_max_logT ¶

No local boost if local logT > this limit.

```
gradT_excess_max_logT = 8
```

gradT_excess_min_log_tau_full_on ¶

gradT_excess_max_log_tau_full_off ¶

No local boost if local log_tau < gradT_excess_max_log_tau_full_off. Reduced local boost if local log_tau < gradT_excess_min_log_tau_full_on.

```
gradT_excess_min_log_tau_full_on = -99
gradT_excess_max_log_tau_full_off = -99
```

smooth_gradT ¶

use_grada_for_smooth_gradT ¶

gradT_smooth_low ¶

gradT_smooth_mid ¶

gradT_smooth_high ¶

gradT_smooth_factor ¶

EXPERIMENTAL: soften gradT at the boundaries of convective zones to help convergence

```
smooth_gradT = .false.
use_grada_for_smooth_gradT = .false.
gradT_smooth_low = -0.005d0
gradT_smooth_mid = 0d0
gradT_smooth_high = 0.01d0
gradT_smooth_factor = 1.0d0
```

max_logT_for_mlt ¶

No mlt at cell if local logT > this limit.

```
max_logT_for_mlt = 99
```

overshooting ¶

Overshooting depends on the classification of the convective zone and can be different at the top and the bottom of the zone.

min_overshoot_q

Overshooting is only allowed at locations with mass $m \geq \text{min_overshoot_q} * m_{\text{star}}$. E.g., if $\text{min_overshoot_q} = 0.1$, then only the outer 90% by mass can have overshooting. This provides a simple way of suppressing bogus center overshooting in which a small convective region at the core can produce excessively large overshooting because of a large pressure scale height at the center.

```
min_overshoot_q = 1d-3
```

D_mix_ov_limit

Overshooting shuts off when the exponential decay has dropped the diffusion coefficient to this level.

```
D_mix_ov_limit = 1d2
```

max_brunt_B_for_overshoot

Terminate overshoot region when encounter stabilizing composition gradient where (unsmoothed) brunt_B is greater than this limit. (≤ 0 means ignore this limit) note: both brunt_B and $\text{gradL_composition_term}$ come from $\text{unsmoothed_brunt_B}$ and differ only in optional smoothing. (see $\text{num_cells_for_smooth_brunt_B}$ and $\text{num_cells_for_smooth_gradL_composition_term}$).

```
max_brunt_B_for_overshoot = 0
```

Parameters for exponential diffusive overshoot are described in the paper by Falk Herwig, “The evolution of AGB stars with convective overshoot”, A&A, 360, 952-968 (2000).

NOTE: In addition to giving these ‘f’ parameters non-zero values, you should also check the settings for $\text{mass_for_overshoot_full_on}$ and $\text{mass_for_overshoot_full_off}$.

The switch from convective mixing to overshooting happens at a distance $f_0 * H_p$ into the convection zone from the estimated location where $\text{grad_ad} == \text{grad_rad}$, where H_p is the pressure scale height at that location. A value ≤ 0 for f_0 is a mistake – you are required to set f_0 as well as f . take a look at the following from an email concerning this:

Overshooting works by taking the diffusion mixing coefficient at the edge of the convection zone and extending it beyond the zone. But – and here’s the issue – at the exact edge of the zone the mixing coefficient goes to 0. So we don’t want that. Instead we want the value of the mixing coeff NEAR the edge, but not AT the edge. The “ f_0 ” parameter determines the exact meaning of “near” for this. It tells the code how far back into the zone to go in terms of scale height. The overshooting actually begins at the location determined by f_0 back into the convection zone rather than at the edge where the diffusion coeff is ill-defined. So, for example, if you want overshooting of 0.2 scale heights beyond the normal edge, you might want to back up 0.05 scale heights to get the diffusion coeff from near the edge and then go out by 0.25 scale heights from there to reach 0.2 H_p beyond the old boundary. In the inlist this would mean setting the “ f_0 ” to 0.05 and the “ f ” to 0.25.

There is no default value for f_0 ; if you set $f > 0$ then you must get $f_0 > 0$ as well.

overshoot_alpha ¶

The value of H_p for overshooting is limited to the radial thickness of the convection zone divided by `overshoot_alpha`. only used when > 0 . if ≤ 0 , then use `mixing_length_alpha` instead.

```
overshoot_alpha = -1
```

limit_overshoot_Hp_using_size_of_convection_zone ¶

if false, allow large distance of overshoot for small convective zones.

```
limit_overshoot_Hp_using_size_of_convection_zone = .true.
```

overshoot_f_above_nonburn_core ¶

overshoot_f0_above_nonburn_core ¶

overshoot_f_above_nonburn_shell ¶

overshoot_f0_above_nonburn_shell ¶

overshoot_f_below_nonburn_shell ¶

overshoot_f0_below_nonburn_shell ¶

For nonburning regions.

```
overshoot_f_above_nonburn_core = 0
overshoot_f0_above_nonburn_core = -1
overshoot_f_above_nonburn_shell = 0
overshoot_f0_above_nonburn_shell = -1
overshoot_f_below_nonburn_shell = 0
overshoot_f0_below_nonburn_shell = -1
```

overshoot_f_above_burn_h_core ¶

overshoot_f0_above_burn_h_core ¶

overshoot_f_above_burn_h_shell ¶

overshoot_f0_above_burn_h_shell ¶

overshoot_f_below_burn_h_shell ¶

overshoot_f0_below_burn_h_shell ¶

For hydrogen burning regions.


```
overshoot_f_above_burn_h_core = 0
overshoot_f0_above_burn_h_core = -1
overshoot_f_above_burn_h_shell = 0
overshoot_f0_above_burn_h_shell = -1
overshoot_f_below_burn_h_shell = 0
overshoot_f0_below_burn_h_shell = -1
```

overshoot_f_above_burn_he_core ¶

overshoot_f0_above_burn_he_core ¶

overshoot_f_above_burn_he_shell ¶

overshoot_f0_above_burn_he_shell ¶

overshoot_f_below_burn_he_shell ¶

overshoot_f0_below_burn_he_shell ¶

For helium burning regions.

```
overshoot_f_above_burn_he_core = 0
overshoot_f0_above_burn_he_core = -1
overshoot_f_above_burn_he_shell = 0
overshoot_f0_above_burn_he_shell = -1
overshoot_f_below_burn_he_shell = 0
overshoot_f0_below_burn_he_shell = -1
```

overshoot_f_above_burn_z_core ¶

overshoot_f0_above_burn_z_core ¶

overshoot_f_above_burn_z_shell ¶

overshoot_f0_above_burn_z_shell ¶

overshoot_f_below_burn_z_shell ¶

overshoot_f0_below_burn_z_shell ¶

For metals burning regions.

```
overshoot_f_above_burn_z_core = 0
overshoot_f0_above_burn_z_core = -1
overshoot_f_above_burn_z_shell = 0
overshoot_f0_above_burn_z_shell = -1
overshoot_f_below_burn_z_shell = 0
overshoot_f0_below_burn_z_shell = -1
```

overshoot_below_noburn_shell_factor ¶

Multiply `overshoot_f_below_nonburn_shell` by this factor only during dredge up phase of AGB thermal pulse.

```
overshoot_below_noburn_shell_factor = 1
```

Optional step function for overshooting. This can be used simultaneously with exponential overshooting. When using step overshoot, you must set `overshoot_f0` as well as `f`.

A convective region is considered a shell if it doesn't reach the center.

```
step_overshoot_f_above_nonburn_core = 0
step_overshoot_f_above_nonburn_shell = 0
step_overshoot_f_below_nonburn_shell = 0
```

```
step_overshoot_f_above_burn_h_core = 0
step_overshoot_f_above_burn_h_shell = 0
step_overshoot_f_below_burn_h_shell = 0
```

```
step_overshoot_f_above_burn_he_core = 0
step_overshoot_f_above_burn_he_shell = 0
step_overshoot_f_below_burn_he_shell = 0
```

```
step_overshoot_f_above_burn_z_core = 0
step_overshoot_f_above_burn_z_shell = 0
step_overshoot_f_below_burn_z_shell = 0
```

step_overshoot_D ¶

step_overshoot_D0_coeff ¶

As above, $f_0 \cdot H_p$ determines r_0 where switch from convection to overshooting. Overshooting extends a distance $\text{step_f} \cdot H_{p0}$ from r_0 with constant diffusion coeff $D = \text{step_D} + \text{step_D0_coeff} \cdot D_0$ where D_0 = diffusion coefficient D at point r_0 .

```
step_overshoot_D = 0
step_overshoot_D0_coeff = 1
```

mass_for_overshoot_full_on ¶

You can specify a range of star masses over which overshooting above H burning zones is gradually enabled. Do specified overshooting above H burning zone if `star_mass` \geq this (M_{sun}).

```
mass_for_overshoot_full_on = 0
```

mass_for_overshoot_full_off ¶

You can specify a range of star masses over which overshooting above H burning zones is gradually enabled. No overshooting above H burning zone if `star_mass` <= this (Msun).

```
mass_for_overshoot_full_off = 0
```

DUP_varcontrol_factor ¶

```
DUP_varcontrol_factor = 1d0
```

max_DUP_counter ¶

For deciding when to terminate use of `overshoot_below_noburn_shell_factor`.

```
max_DUP_counter = 200
```

ovr_below_burn_he_shell_factor ¶

Multiply `overshoot_f_below_burn_he_shell` by this factor after the first AGB thermal pulse.

```
ovr_below_burn_he_shell_factor = 1
```

overshoot_D2 ¶

overshoot_f2 ¶

optional 2nd scale length for exponential overshooting

$f_0 \cdot H_p$ determines location r_0 where we switch from convection to overshooting. Let D_0 = diffusion coefficient D at point r_0 . Let H_{p0} = the scale height at r_0 . In the standard version of exponential overshooting, there is a single length scale = $f \cdot H_{p0}$ and at a distance dr from r_0 , $D(dr) = D_0 \cdot \exp(-2 \cdot dr / (f \cdot H_{p0}))$.

In the extended version there is a second length scale = $f_2 \cdot H_{p0}$. The second length scale takes effect for distances $dr > dr_2$ where dr_2 is defined by $D_2 = D_0 \cdot \exp(-2 \cdot dr_2 / (f \cdot H_{p0}))$.

- for $dr \leq dr_2$, $D(dr) = D_0 \cdot \exp(-2 \cdot dr / (f \cdot H_{p0}))$
- for $dr > dr_2$, $D(dr) = D_2 \cdot \exp(-2 \cdot (dr - dr_2) / (f_2 \cdot H_{p0})) = D_0 \cdot \exp(-2 \cdot dr_2 / (f \cdot H_{p0})) \cdot \exp(-2 \cdot (dr - dr_2) / (f_2 \cdot H_{p0})) = D_0 \cdot \exp(-2 \cdot (dr_2 / (f \cdot H_{p0}) + (dr - dr_2) / (f_2 \cdot H_{p0})))$

```
overshoot_D2_above_nonburn = -1d0
overshoot_D2_below_nonburn = -1d0
```

```
overshoot_D2_above_burn_h = -1d0
overshoot_D2_below_burn_h = -1d0
```

```
overshoot_D2_above_burn_he = -1d0
overshoot_D2_below_burn_he = -1d0
```

```
overshoot_D2_above_burn_z = -1d0
overshoot_D2_below_burn_z = -1d0
```

```
overshoot_f2_above_nonburn = 1d0
overshoot_f2_below_nonburn = 1d0
```

```
overshoot_f2_above_burn_h = 1d0
overshoot_f2_below_burn_h = 1d0
```

```
overshoot_f2_above_burn_he = 1d0
overshoot_f2_below_burn_he = 1d0
```

```
overshoot_f2_above_burn_z = 1d0
overshoot_f2_below_burn_z = 1d0
```

RGB_to_AGB_cbm_switch ¶

If center hydrogen abundance is < 0.01 and center helium abundance by mass is less than `RGB_to_AGB_cbm_switch`, then system will include `overshoot_D2` and `overshoot_f2` parameters to describe convective boundary mixing during the AGB phase if set to positive values. See Battino et al. 2016, APJ, 827:30

“Application of a theory and simulation-based convective boundary mixing model for AGB star evolution and nucleosynthesis” `RGB_to_AGB_cbm_switch = 1d-4` is used in the paper above.

```
RGB_to_AGB_cbm_switch = -1
```

Predictive mixing ¶

Predictive mixing is an approach for expanding convective boundaries until `gradr = grada` on the convective side of the boundary (as required by the criterion that the convective velocity and luminosity vanish at the boundary). It is discussed in detail in Paxton et al. 2018, ApJ, in press: “Modules for Experiments in Stellar Astrophysics (MESA): Convective boundaries, element diffusion, and massive star explosions”

Predictive mixing is controlled by specifying a set of parameters, which combines matching criteria (determining which boundaries to apply the predictive mixing to) together with values (determining how the predictive mixing should operate at those boundaries). Up to `NUM_PREDICTIVE_PARAM_SETS` of these parameter sets can be defined (see `star_def.inc` for value).

predictive_mix ¶

Set to `.true.` to enable this set of parameters

```
predictive_mix(1) = .false.
```

predictive_zone_type ¶

Matching criterion for the type of the convection zone. Possible values are `burn_H` (hydrogen burning), `burn_He` (helium burning), `burn_Z` (metal burning), `nonburn` (no burning) or `any` (which matches any type of zone).

```
predictive_zone_type(1) = ''
```

predictive_zone_loc ¶

Matching criterion for the location of the convection zone. Possible values are `core` (the core convection zone), `shell` (a convective shell), `surf` (the surface convection zone) or `any` (which matches any location).

```
predictive_zone_loc(1) = ''
```

predictive_bdy_loc ¶

Matching criterion for the location of the convective boundary. Possible values are `top` (the top of the convection zone), `bottom` (the bottom of the convection zone) or `any` (which matches any location).

```
predictive_bdy_loc(1) = ''
```

predictive_bdy_q_min ¶

Matching criterion for the minimum fractional mass coordinate of the convective boundary

```
predictive_bdy_q_min(1) = 0d0
```

predictive_bdy_q_max ¶

Matching criterion for the maximum fractional mass coordinate of the convective boundary

```
predictive_bdy_q_max(1) = 1d0
```

predictive_superad_thresh ¶

Threshold for minimum superadiabaticity in the predictive mixing scheme; boundary expansion stops when `gradr/grada-1` drops below this threshold. Default value is usually good for main-sequence evolution; for

core He-burning, set to 0.005, 0.01 or larger to prevent splitting of the core convection zone and/or core breathing pulses.

```
predictive_superad_thresh(1) = 0d0
```

predictive_avoid_reversal ¶

Species to monitor for reversals in abundance evolution. If this is set to the name of a species, then the predictive mixing scheme will try to avoid causing reversals in the abundance of that species (e.g., changing the abundance evolution from decreasing to increasing). Set to 'he4' during core He-burning to prevent splitting of the core convection zone and/or core breathing pulses.

```
predictive_avoid_reversal(1) = ''
```

predictive_limit_ingestion ¶

predictive_ingestion_factor ¶

Limit the rate of ingestion of a species, following the prescription given in equation (2) of Constantino, Campbell & Lattanzio (2017, MNRAS, 472, 4900). The control `predictive_limit_ingestion` specifies which species to limit, and the control `predictive_ingestion_factor` gives the multiplying factor. Setting this factor to 5/12 is the same as choosing $\alpha_i = 1$ in their equation (2).

```
predictive_limit_ingestion(1) = ''
predictive_ingestion_factor(1) = 0d0
```

Self-Driving Overshoot & Semiconvection (S-DOS) mixing ¶

```
do_sdos_mix = .false.
sdos_avoid_reversal = .false.
```

New overshooting ¶

```
overshoot_new = .false.
overshoot_f = 0d0
overshoot_f0 = 0d0
overshoot_f2 = 0d0
overshoot_D0 = 0d0
overshoot_D2 = 0d0
overshoot_Delta0 = 1d0
overshoot_c1 = 100d0
overshoot_c2 = 90d0
overshoot_mass_full_on = 0d0
overshoot_mass_full_off = 0d0
overshoot_scheme = ''
overshoot_zone_type = ''
overshoot_zone_loc = ''
overshoot_bdy_loc = ''
```

```
overshoot_D_min = 1d2
overshoot_brunt_B_max = 0d0
```

turbulence ¶

RTI_max_time_full_off ¶

RTI_min_time_full_on ¶

```
RTI_max_time_full_off = 0d0
RTI_min_time_full_on = 0d0
```

RTI_smooth_mass ¶

RTI_smooth_iterations ¶

RTI_smooth_fraction ¶

smoothing for dPdr_dRhodr_info done at start of step

```
RTI_smooth_mass = 0d0
RTI_smooth_iterations = 0
RTI_smooth_fraction = 1d0
```

alpha_RTI_diffusion_factor ¶

dudt_RTI_diffusion_factor ¶

dedt_RTI_diffusion_factor ¶

dlnddt_RTI_diffusion_factor ¶

composition_RTI_diffusion_factor ¶

max_M_RTI_factors_full_on ¶

min_M_RTI_factors_full_off ¶

```
alpha_RTI_diffusion_factor = 1d0
dudt_RTI_diffusion_factor = 1d0
dedt_RTI_diffusion_factor = 1d0
dlnddt_RTI_diffusion_factor = 1d0
composition_RTI_diffusion_factor = 1d0
max_M_RTI_factors_full_on = 1d99
min_M_RTI_factors_full_off = 1d99
```

alpha_RTI_src_max_q ¶

alpha_RTI_src_min_q ¶

option to set alpha_RTI source term to zero when cell q out of bounds. to turn off RTI near surface or center

```
alpha_RTI_src_max_q = 1d0
alpha_RTI_src_min_q = 0d0
```

alpha_RTI_src_min_v_div_cs ¶

option to set alpha_RTI source term to zero when $v/cs <$ this min. e.g. to filter out false sources ahead of shock

```
alpha_RTI_src_min_v_div_cs = 1d0
```

radiation_turbulence_coeff ¶

To counter depletion of h and metals in outer envelope of stars with $M > 1.4 M_{\text{sun}}$. Morel, P., and Thevenin, F., Atomic diffusion in stellar models of type earlier than G., A&A, 390:611-620 (2002)

```
D = radiation_turbulence_coeff * 4*crad*T^4/(15*clight*opacity*rho^2)
```

1 is reasonable value for this coefficient

```
radiation_turbulence_coeff = 0
```

turbulent_diffusion_D0 ¶

Turbulent diffusion below outer convection zone. Similar effect to overshooting. Proffitt, C.R., and Michaud, G., GRAVITATIONAL SETTLING IN SOLAR MODELS, ApJ, 380:238-290, 1991. e.g., $8000 \text{ cm}^2 \text{ s}^{-1}$

```
turbulent_diffusion_D0 = 0
```

turbulent_diffusion_rho_max ¶

Only have turbulent diffusion if $\rho <$ this. Diffusion coef $D = D0 * (\rho/\rho_{\text{base_cz}})^{-3}$, but only if $\rho_{\text{base_cz}} \leq \text{turbulent_diffusion_rho_max}$.

```
turbulent_diffusion_rho_max = 1d99
```

turbulent_diffusion_Dmin ¶

Only set turbulent diffusion if it gives $D \geq$ this.


```
turbulent_diffusion_Dmin = 1d1
```

mixing misc ¶

such as smoothing and editing of diffusion coefficients

mix_factor ¶

Mixing coefficients are multiplied by this factor. The `mix_factor` is applied in subroutine `get_convection_sigmas` in `star/private/mix_info.f90` – the lagrangian diffusion coefficient $\sigma(k)$ at cell boundary k is set to $\text{mix_factor} * D * (4 * \pi * r(k)^2 * \rho_{\text{face}}(k))^2$. Note that the value of D is not changed – it is just used as a term in calculating σ .

```
mix_factor = 1
```

min_dt_for_increases_in_convection_velocity ¶

convective velocities are not increased if $dt < \text{this value}$ (in seconds)

```
min_dt_for_increases_in_convection_velocity = -1d0
```

max_conv_vel_div_csound ¶

convective velocities are limited to local sound speed times this factor

```
max_conv_vel_div_csound = 1d99
```

max_v_div_cs_for_convection ¶

disable convection for locations with $\text{abs}(v)/c_s > \text{this limit}$

```
max_v_div_cs_for_convection = 1d99
```

max_abs_du_div_cs_for_convection ¶

main purpose is to force radiative in shock face

```
max_abs_du_div_cs_for_convection = 0.03
```

min_T_for_acceleration_limited_conv_velocity ¶

Acceleration limiting based on Wood 1974 and Arnett 1969. Wood, P.R., ApJ, 190:609-630, 1974. (Appendix V, eqns 1-3) Arnett, W.D., 1969, Ap. and Space Sci, 5, 180.

```
min_T_for_acceleration_limited_conv_velocity = 99e9
```

max_T_for_acceleration_limited_conv_velocity

```
max_T_for_acceleration_limited_conv_velocity = 99e9
```

mlt_accel_g_theta

use this (if > 0) for limiting the acceleration of convection velocities.

```
mlt_accel_g_theta = -1
```

prune_bad_cz_min_Hp_height

Lower limit on radial extent of cz (≤ 0 to disable). Remove tiny convection zones unless have strong nuclear burning i.e., remove if $\text{size} < \text{prune_bad_cz_min_Hp_height}$ and $\text{max_log_eps} < \text{prune_bad_cz_min_log_eps_nuc}$.

```
prune_bad_cz_min_Hp_height = 0
```

prune_bad_cz_min_log_eps_nuc

Lower limit on max log eps nuc in cz. In units of average pressure scale height at top and bottom of region. This allows emergence of very small cz at site of the core flash, for example.

```
prune_bad_cz_min_log_eps_nuc = -99
```

redo_conv_for_dr_lt_mixing_length

Check for small convection zones with total height less than mixing length and redo with reduced $\text{mixing_length_alpha}$ to make $\text{mixing_length} \leq \text{dr}$.

```
redo_conv_for_dr_lt_mixing_length = .false.
```

limit_mixing_length_by_dist_to_bdy

reduce local value of mixing length alpha if necessary in order to make $\text{mixing_length} \leq \text{distance to convective boundary}$ times this value only applies when value is > 0 setting this value = 1 implements the restriction that near a convective boundary, the mixing length doesn't exceed the distance to the boundary. Peter Eggleton, "Composition Changes during Stellar Evolution", MNRAS 156, 361-376, 1972.

WARNING: I've seen problems with 25M before He core burn when using this. bp.

```
limit_mixing_length_by_dist_to_bdy = 0
```

conv_bdy_mix_softening_f0

conv_bdy_mix_softening_f

conv_bdy_mix_softening_min_D_mix

These controls cause the convective mixing coefficient to drop off “softly” near the boundary of the convective zone – i.e, they prevent situations where the mixing coefficient drops from 10^{10} or more to zero in a distance covered by only one or two cells as can happen at jumps in composition. Such a sharp edge is no problem when it is not adjacent to a convective region. But when it shows up at a convective boundary, it is problematic. It may not be physical, and it is certainly bad news numerically. This has been discussed as early as 1970’s – see for example, Peter Eggleton, “Composition Changes during Stellar Evolution”, MNRAS 156, 361-376, 1972.

The implementation of softening at convective boundaries is like overshooting but the distances are typically smaller by an order of magnitude or more. Also, the softening is primarily inside the convective region rather than penetrating strongly into the area beyond. This is done by backing up from the boundary into the convection region to start the softening of the mixing coefficient so that most of the effect takes place before reaching the exterior of the region. The softening extends a short way into the exterior with a decreasing value of mixing. For example, it might start a distance of $0.003 * H_p$ into the convective region (H_p = pressure scale height at the boundary), and then project a mixing coefficient outward from there decreasing exponentially with distance scale of $0.001 * H_p$. For those numbers, there are 3 e-foldings before reaching the boundary, so most of the drop has happened inside the convective region. The strength of the mixing then continues to drop exponentially until it reaches some given limit.

The effect of this will be to soften the jump in abundances immediately adjacent to the convective region. That will in turn soften the jump in opacity and the corresponding jump in `grad_rad` so that there will not be a large jump from a convective point with `grad_rad >> grad_ad` to a neighboring non-convective point with `grad_rad << grad_ad`.

```
conv_bdy_mix_softening_f0 = 0
conv_bdy_mix_softening_f = 0
conv_bdy_mix_softening_min_D_mix = 0
```

smooth_convective_bdy

This is an option to smooth composition gradients in newly non-convective regions trailing behind a retreating convection zone. This effectively erases (most) of the stair-casing that happens without it. But you should be aware that the smoothing process does not conserve species mass – e.g., if have retreating He burning core below H shell, then the smoothing will convert some H into He in the newly non-convective region (this can be hand waved away as modeling partial burning of those regions during the substep period before the convection had retreated past the location).

set this true to have the stair-casing removed at the price of some changes in abundances.

```
smooth_convective_bdy = .true.
```

max_dR_div_Hp_for_smooth

Don’t smooth across a newly nonconvective region larger than this limit where `dR` is radial thickness of region and `Hp` is min pressure scale height in region.

```
max_dR_div_Hp_for_smooth = 10
```

max_delta_limit_for_smooth ¶

Don't smooth across a newly nonconvective region where any mass fraction changes by more than this limit.

```
max_delta_limit_for_smooth = 0.1d0
```

remove_mixing_glitches ¶

If true, then okay to remove gaps and singletons.

```
remove_mixing_glitches = .true.
```

glitches ¶

The following controls are for different kinds of “glitches” that can be removed.

okay_to_remove_mixing_singleton ¶

If true, remove singetons.

```
okay_to_remove_mixing_singleton = .true.
```

clip_D_limit ¶

Zero mixing diffusion coeffs that are smaller than this.

```
clip_D_limit = 0
```

min_convective_gap ¶

Close gap between convective regions if smaller than this (< 0 means skip this). Gap measured radially in units of pressure scale height.

```
min_convective_gap = -1
```

min_thermohaline_gap ¶

Close gap between thermohaline mixing regions if smaller than this (< 0 means skip this). Gap measured radially in units of pressure scale height.

```
min_thermohaline_gap = -1
```

min_thermohaline_dropout ¶

max_dropout_gradL_sub_grada ¶

If find radiative region embedded in thermohaline, and $\max(\text{gradL} - \text{grada})$ in region is everywhere $< \text{max_dropout_gradL_sub_grada}$ and region height is $< \text{min_thermohaline_dropout}$ then convert the region to thermohaline. $\text{min_thermohaline_dropout} \leq 0$ disables.

```
min_thermohaline_dropout = -1
max_dropout_gradL_sub_grada = 1d-3
```

min_semiconvection_gap ¶

Close gap between semiconvective mixing regions if smaller than this (< 0 means skip this). Gap measured radially in units of pressure scale height.

```
min_semiconvection_gap = -1
```

remove_embedded_semiconvection ¶

If have a semiconvection region bounded on each side by convection, convert it to be convective too.

```
remove_embedded_semiconvection = .false.
```

set_min_D_mix ¶

mass_lower_limit_for_min_D_mix ¶

mass_upper_limit_for_min_D_mix ¶

min_D_mix ¶

D_{mix} will be at least this large if set_min_D_mix is true. doesn't apply for $\text{mass} < \text{lower limit}$ or $\text{mass} > \text{upper limit}$.

```
set_min_D_mix = .false.
mass_lower_limit_for_min_D_mix = 0d0
mass_upper_limit_for_min_D_mix = 1d99
min_D_mix = 1d3
```

set_min_D_mix_in_H_He ¶

min_D_mix_in_H_He ¶

D_{mix} will be at least this large in regions where max mass fractions of H and He add to more that 0.5 if $\text{set_min_D_mix_in_H_He}$ is true.

```
set_min_D_mix_in_H_He = .false.
min_D_mix_in_H_He = 1d3
```

set_min_D_mix_below_Tmax ¶

min_D_mix_below_Tmax ¶

D_mix will be at least this large for cells below location of max temperature if set_min_D_mix_below_Tmax is true.

```
set_min_D_mix_below_Tmax = .false.
min_D_mix_below_Tmax = 1d3
```

min_center_Ye_for_min_D_mix ¶

min_D_mix is only used when center_ye >= this i.e., when center_ye drops below this, min_D_mix = 0.

```
min_center_Ye_for_min_D_mix = 0.47d0
```

smooth_outer_xa_big ¶

smooth_outer_xa_small ¶

Soften composition jumps in outer layers. If smooth_outer_xa_big and smooth_outer_xa_small are bigger than 0, then starting from the outermost grid point, homogeneously mix a region of size smooth_outer_xa_small (in solar masses), and proceed inwards, linearly reducing the size of the homogeneously mixed region in such a way that it becomes zero. After going smooth_outer_xa_big solar masses in. In this way, the outer smooth_outer_xa_big solar masses are “cleaned” of composition jumps.

```
smooth_outer_xa_big = -1d0
smooth_outer_xa_small = -1d0
```

rotation controls ¶

In the following “am” stands for “angular momentum”.

the mesa implementation of rotation closely follows these papers:

- Heger, Langer, & Woosley, ApJ, 528, 368. 2000
- Heger, Woosley, & Spruit, ApJ, 626, 350. 2005
- D_DSI = dynamical shear instability
- D_SH = Solberg-Hoiland
- D_SSI = secular shear instability
- D_ES = Eddington-Sweet circulation

- D_GSF = Goldreich-Schubert-Fricke
- D_ST = Spruit-Tayler dynamo

skip_rotation_in_convection_zones ¶

if true, then set rotational diffusion coefficients to 0 in convective regions. This applies both for material mixing and diffusion of angular momentum.

```
skip_rotation_in_convection_zones = .false.
```

am_D_mix_factor ¶

Rotation and mixing of material. D_mix = diffusion coefficient for mixing of material. It is sum of non-rotational and rotational components. The rotational part is multiplied by this factor.

```
D_mix = D_mix_non_rotation + f*am_D_mix_factor*(
    D_DSI_factor * D_DSI +
    D_SH_factor * D_SH +
    D_SSI_factor * D_SSI +
    D_ES_factor * D_ES +
    D_GSF_factor * D_GSF +
    D_ST_factor * D_ST)

f = 1  when logT <= D_mix_rotation_max_logT_full_on = full_on
    = 0  when logT >= D_mix_rotation_max_logT_full_on = full_off
    = (log(T)-full_on)/(full_off-full_on) else
```

note that for regions with brunt $N^2 < 0$, we set Richardson number to 1 which is $> Ri_{critical}$ and therefore turns off DSI and SSI

according to Heger et al 2000 : 1/30d0 by default : 0

```
am_D_mix_factor = 0
```

am_nu_factor ¶

am_nu_non_rotation_factor ¶

diffusion of angular momentum

am_nu = diffusion coefficient for angular momentum

```
am_nu_non_rot = am_nu_factor*am_nu_non_rotation_factor*D_mix_non_rotat
am_nu_rot = am_nu_factor*(
    am_nu_visc_factor * D_visc +
    am_nu_DSI_factor * D_DSI +
    am_nu_SH_factor * D_SH +
    am_nu_SSI_factor * D_SSI +
    am_nu_ES_factor * D_ES +
```

```
    am_nu_GSF_factor * D_GSF +  
    am_nu_ST_factor * nu_ST)  
am_nu = am_nu_non_rot + am_nu_rot
```

Note that for regions with $\text{brunt } N^2 < 0$, we set Richardson number to 1 which is $> \text{Ri_critical}$ and therefore turns off DSI and SSI.

see also `star_job` controls for `am_nu_rot_flag`

```
am_nu_factor = 1  
am_nu_non_rotation_factor = 1
```

am_nu_DSI_factor ¶

< 0 means use `D_DSI_factor`

```
am_nu_DSI_factor = -1
```

am_nu_SSI_factor ¶

< 0 means use `D_SSI_factor`

```
am_nu_SSI_factor = -1
```

am_nu_SH_factor ¶

< 0 means use `D_SH_factor`

```
am_nu_SH_factor = -1
```

am_nu_ES_factor ¶

< 0 means use `D_ES_factor`

```
am_nu_ES_factor = -1
```

am_nu_GSF_factor ¶

< 0 means use `D_GSF_factor`

```
am_nu_GSF_factor = -1
```

am_nu_ST_factor ¶

< 0 means use D_ST_factor

```
am_nu_ST_factor = -1
```

am_nu_visc_factor ¶

< 0 means use D_visc_factor. By default = 1 to mix angular momentum.

```
am_nu_visc_factor = 1
```

am_nu_omega_rot_factor ¶

am_nu_omega_non_rot_factor ¶

```
dj/dt = d/dm((4 pi r^2 rho)^2*(am_nu_omega*i_rot*domega/dm + am_nu_j*d
am_nu_omega = am_nu_omega_non_rot_factor*am_nu_non_rot + am_nu_omega_r
```

```
am_nu_omega_rot_factor = 1
am_nu_omega_non_rot_factor = 1
```

am_nu_j_rot_factor ¶

am_nu_j_non_rot_factor ¶

```
dj/dt = d/dm((4 pi r^2 rho)^2*(am_nu_omega*i_rot*domega/dm + am_nu_j*d
am_nu_j = am_nu_j_non_rot_factor*am_nu_non_rot + am_nu_j_rot_factor*am
```

```
am_nu_j_rot_factor = 0
am_nu_j_non_rot_factor = 0
```

set_uniform_am_nu_non_rot ¶

uniform_am_nu_non_rot ¶

You can specify a uniform value for am_nu_non_rot by setting this flag true. A large uniform am_nu will produce a uniform omega.

```
set_uniform_am_nu_non_rot = .false.
uniform_am_nu_non_rot = 1d20
```

set_min_am_nu_non_rot ¶**min_am_nu_non_rot ¶**

You can also specify a minimum am_nu_non_rot. am_nu will be at least this large.

```
set_min_am_nu_non_rot = .false.
min_am_nu_non_rot = 1d8
```

min_center_Ye_for_min_am_nu_non_rot ¶

min_am_nu_non_rot is only used when center Ye >= this.

```
min_center_Ye_for_min_am_nu_non_rot = 0.47d0
```

Each rotationally induced diffusion coefficient has a factor that lets you control it. Value of 1 gives normal strength; value of 0 turns it off.

Note that for regions with brunt $N^2 < 0$, we set Richardson number to 1, which is $> Ri_{critical}$ and therefore turns off DSI and SSI.

```
D_DSI_factor = 0
D_SH_factor = 0
D_SSI_factor = 0
D_ES_factor = 0
D_GSF_factor = 0
D_ST_factor = 0
```

D_visc_factor ¶

Kinematic shear viscosity. Should be = 0 because viscosity doesn't mix chemical elements.

```
D_visc_factor = 0
```

am_gradmu_factor ¶

Sensitivity to composition gradients. In calculation of rotational induced mixing, grad_mu is multiplied by am_gradmu_factor. Value from from Heger et al 2000.

```
am_gradmu_factor = 0.05d0
```

Spatial smoothing is used in calculations of diffusion coefficients. These control the smoothing window widths (number of cells on each side).

```
smooth_D_DSI = 0
smooth_D_SH = 0
```

```
smooth_D_SSI = 0
smooth_D_ES = 0
smooth_D_GSF = 0
smooth_D_ST = 0
smooth_nu_ST = 0
```

time smoothing. Set to 0 to turn off time smoothing.

```
angsm_t_D_DSI = 0.0d0
angsm_t_D_SH = 0.0d0
angsm_t_D_SSI = 0.0d0
angsm_t_D_ES = 0.0d0
angsm_t_D_GSF = 0.0d0
angsm_t_D_ST = 0.2d0
angsm_t_nu_ST = 0.2d0
angsm_l = 1d-3
```

am_time_average ¶

If true, then $D = (D_{\text{new}} + D_{\text{old}})/2$, where D_{old} is D from previous step and D_{new} is D as calculated for current as if no time smoothing.

```
am_time_average = .false.
```

simple_i_rot_flag ¶

If true, $i_{\text{rot}} = (2/3) * r^2$. If false, use slightly more complex expression that takes into account finite shell thickness. In practice, there doesn't seem to be a significant difference.

```
simple_i_rot_flag = .true.
```

fitted_i_rot_flag ¶

EXPERIMENTAL: If true and `simple_i_rot_flag = .false.`, use fit to i_{rot} that depends on rotation rate.

```
fitted_i_rot_flag = .false.
```

do_adjust_J_lost ¶

adjust_J_fraction ¶

adjust angular momentum With `do_adjust_J_lost = .false.`, the angular momentum removed via winds from the star corresponds to that contained in the removed layers. However, since j_{rot} can increase steeply in the very outer layers, very small steps are required to obtain a convergent solution. To avoid this, the `do_adjust_J_lost` option adjusts the angular momentum content of layers below those removed, such that

```
actual_J_lost = &
  adjust_J_fraction*mass_lost*s% j_rot_avg_surf + &
  (1d0 - adjust_J_fraction)*s% angular_momentum_removed
```

where `s% angular_momentum_removed` is the angular momentum contained in the removed layers of the star in that step. Note that `s% angular_momentum_removed` is set to `actual_J_lost` after this.

The region from which angular momentum is removed is chosen such that at its bottom `q<min_q_for_adjust_J_lost`, it contains at least `min_J_div_delta_J` times the angular momentum that needs to be accounted for, and it is at an optical depth below `min_tau_for_adjust_J_lost`. Angular momentum in these regions is adjusted in such a way that no artificial shear is produced at the inner boundary.

This can also be used to model mass loss mechanisms that remove more angular momentum than `mass_lost*s% j_rot_avg_surf`, for instance magnetic braking or wind mass loss. In that case, you can use the `use_other_j_for_adjust_J_lost` option to specify a specific angular momentum of removed material different from `j_rot_avg_surf`

```
do_adjust_J_lost = .true.
adjust_J_fraction = 1d0
min_q_for_adjust_J_lost = 0.99d0
min_J_div_delta_J = 3d0
min_tau_for_adjust_J_lost = 300d0
```

premix_omega ¶

if `premix_omega` is true, then do 1/2 of the transport of angular momentum before updating the structure and 1/2 after. otherwise, do all of the transport after updating the structure. RECOMMENDED to turn it on when modelling an accreting star or when using `do_adjust_J_lost`.

```
premix_omega = .true.
```

recalc_mixing_info_each_substep ¶

if `recalc_mixing_info_each_substep` is true, then recalculate the omega mixing coefficients after each substep of the solve omega mix process.

```
recalc_mixing_info_each_substep = .false.
```

use_fitted_fpft ¶

Use analytical fits to the rotational corrections `fp` and `ft`, computed using the Roche potential for a single particle.

```
use_fitted_fpft = .false.
```

w_div_wcrit_min_for_fpft ¶

When `use_fitted_fpft = .true.`, limit `fp` and `ft` to their values at this `w_div_wcrit`

```
w_div_wcrit_min_for_fpft = 0.9999
```

FP_min ¶

FT_min ¶

Lower limits for rotational distortion corrections factors `FP` and `FT`. Used for the calculation when `use_fitted_fpft = .false.`, otherwise the limits are set using `w_div_wcrit_min_for_fpft`

```
FP_min = 0.75d0  
FT_min = 0.95d0
```

FP_error_limit ¶

If calculate an `fp < this`, treat it as an error. Used for the calculation when `use_fitted_fpft = .false.`

```
FP_error_limit = 0d0
```

FT_error_limit ¶

If calculate an `ft < this`, treat it as an error. Used for the calculation when `use_fitted_fpft = .false.`

```
FT_error_limit = 0d0
```

D_mix_rotation_max_logT_full_on ¶

Use rotational components of `D_mix` for locations where `logT <= this`. For numerical stability, turn off rotational part of `D_mix` at very high `T`.

```
D_mix_rotation_max_logT_full_on = 9.4d0
```

D_mix_rotation_min_logT_full_off ¶

Drop rotational components of `D_mix` for locations where `logT >= this`. For numerical stability, turn off rotational part of `D_mix` at very high `T`.

```
D_mix_rotation_min_logT_full_off = 9.5d0
```

D_omega_max_replacement_fraction ¶

D_omega_growth_rate ¶

D_omega_mixing_rate ¶

D_omega_mixing_across_convection_boundary (previously called D_omega_mixing_in_convection_regions) ¶

```
D_omega_max_replacement_fraction = 0.5d0
D_omega_growth_rate = 1d0
D_omega_mixing_rate = 1d0
D_omega_mixing_across_convection_boundary = .false.
max_q_for_D_omega_zero_in_convection_region = 0.8d0
```

nu_omega_max_replacement_fraction ¶

nu_omega_growth_rate ¶

nu_omega_mixing_rate ¶

nu_omega_mixing_across_convection_boundary ¶

```
nu_omega_max_replacement_fraction = 0.5d0
nu_omega_growth_rate = 1d0
nu_omega_mixing_rate = 1d0
nu_omega_mixing_across_convection_boundary = .false.
max_q_for_nu_omega_zero_in_convection_region = 0.8d0
```

atmosphere boundary conditions ¶

which_atm_option ¶

- 'simple_photosphere' : don't integrate, just estimate for $\tau=2/3$
- 'Eddington_grey' : Eddington T-tau integration
- 'Krishna_Swamy' : Krishna Swamy T-tau integration
- 'solar_Hopf_grey' : another T(τ), this one tuned to solar data.
- 'tau_100_tables' : use model atmosphere tables for Pgas and T at $\tau=100$; solar Z only.
- 'tau_10_tables' : use model atmosphere tables for Pgas and T at $\tau=10$; solar Z only.
- 'tau_1_tables' : use model atmosphere tables for Pgas and T at $\tau=1$; solar Z only.
- 'tau_1m1_tables' : use model atmosphere tables for Pgas and T at $\tau=1e-1$; solar Z only.
- 'photosphere_tables' : use model atmosphere tables for photosphere; range of Z's.
- 'grey_and_kap' : iterate simple grey to find consistent P, T, and kap at surface
- 'grey_irradiated' : based on Guillot, T, and Havel, M., A&A 527, A20 (2011).
- 'Paczynski_grey' : create an atmosphere for given base conditions. inspired by B. Paczynski, 1969, Acta Astr., vol. 19, 1. takes into account dilution when $\tau < 2/3$, and calls mlt to get gradT allowing for convection in atmosphere.
- 'WD_tau_25_tables' : hydrogen atmosphere tables for cool white dwarfs giving Pgas and T at $\log_{10}(\tau) = 1.4$ ($\tau = 25.11886$) Teff goes from 40,000 K down to 2,000K with step of 100 K Log10(g) goes from 9.5 down to 5.5 with step of 0.1. R.D. Rohrmann, L.G. Althaus, and S.O. Kepler, Lyman α wing absorption in cool white dwarf stars, Mon. Not. R. Astron. Soc. 411, 781–791 (2011)
- 'fixed_Teff' : set Tsurf from Eddington T-tau relation for current surface tau and Teff = atm_fixed_Teff. set Psurf = Radiation_Pressure(Tsurf)

- 'fixed_Tsurf' : get value of Tsurf from control parameter atm_fixed_Tsurf. set Teff from Eddington T-tau relation for given Tsurf and tau=2/3 set Psurf = Radiation_Pressure(Tsurf)
- 'fixed_Psurf' : get value of Psurf from control parameter atm_fixed_Psurf. set Tsurf from L and R using $L = 4 \cdot \pi \cdot R^2 \cdot \text{boltz_sigma} \cdot T^4$. set Teff using Eddington T-tau relation for tau=2/3 and T=Tsurf.
- 'fixed_Psurf_and_Tsurf' : get value of Psurf from control parameter atm_fixed_Psurf. get value of Tsurf from control parameter atm_fixed_Tsurf.

```
which_atm_option = 'simple_photosphere'
```

which_atm_off_table_option ¶

If have selected an atm table as your option, fallback to using this if the args are off the table. 'simple_photosphere' or 'grey_and_kap'.

```
which_atm_off_table_option = 'simple_photosphere'
```

atm_fixed_Teff ¶

Set this when using atm_option = 'fixed_Teff'

```
atm_fixed_Teff = 0
```

atm_fixed_Tsurf ¶

Set this when using atm_option = 'fixed_Tsurf'

```
atm_fixed_Tsurf = 0
```

atm_fixed_Psurf ¶

Set this when using which_atm_option = 'fixed_Psurf'

```
atm_fixed_Psurf = -1
```

atm_switch_to_grey_as_backup ¶

If you select a table option, but the args are out of the range of the tables, then this flag determines whether you get an error or the code automatically switches to option = atm_simple_photosphere as a backup.

```
atm_switch_to_grey_as_backup = .true.
```

Pextra_factor ¶

Parameter for extra pressure in surface boundary conditions. Pressure at optical depth τ is calculated as $P = \tau * g / \kappa * (1 + P_{\text{extra}})$. P_{extra} takes into account nonzero radiation pressure at $\tau=0$. The equation for P_{extra} includes $P_{\text{extra_factor}}$

```
Pextra = Pextra_factor*(kap/tau)*(L/M)/(6d0*pi*cflight*cgrav)
```

For certain situations such as super eddington L , you may need to increase P_{extra} to help convergence. e.g. try $P_{\text{extra_factor}} = 2$. $P_{\text{extra_factor}} < 0$ means use (incorrect) old form $1.6d-4 * \kappa * (L/L_{\text{sun}})/(M/M_{\text{sun}})$.

```
Pextra_factor = 1
```

atm_grey_and_kap_atol ¶

atm_grey_and_kap_rtol ¶

Relative and absolute tolerance parameters for the `grey_and_kap` option. Iterates on κ until $\text{err} = |\Delta \kappa| / (\text{atol} + \text{rtol} * \kappa) < 1$.

```
atm_grey_and_kap_atol = 1d-7
atm_grey_and_kap_rtol = 1d-7
```

atm_grey_and_kap_max_tries ¶

trace_atm_grey_and_kap ¶

Limit on iterations and trace.

```
atm_grey_and_kap_max_tries = 50
trace_atm_grey_and_kap = .false.
```

`atm_grey_irradiated_atol` and `atm_grey_irradiated_rtol`. Parameters for the `grey_irradiated` option. Absolute and relative error tolerances.

```
atm_grey_irradiated_atol = 1d-4
atm_grey_irradiated_rtol = 1d-4
```

atm_grey_irradiated_T_eq ¶

Equilibrium temperature based on irradiation.

```
irrad_flux = Lstar/(4*pi*orbit**2)
```


- Area of planet in plane perpendicular to `irrad_flux` = $\pi \cdot R_{\text{planet}}^2$.
- Stellar luminosity received by planet = `irrad_flux`*area.
- This luminosity determines `T_eq`: $T_{\text{eq}}^4 = \text{irrad_flux} / (4 \cdot \sigma)$.

```
atm_grey_irradiated_T_eq = 1000
```

atm_grey_irradiated_kap_v

atm_grey_irradiated_simple_kap_th

Opacity for irradiation. The $T(\tau)$ relation for this option depends on the ratio `kap_v/kap_th` where `kap_v` is the planet atmosphere opacity for stellar irradiation, and `kap_th` is the thermal opacity for internally produced radiation. You can either specify the ratio of `kap_v/kap_th`, or you can specify `kap_v` and have the code calc `kap_th` to get the ratio.

```
atm_grey_irradiated_kap_v = 4d-3
atm_grey_irradiated_simple_kap_th = .false.
```

atm_grey_irr_kap_v_div_kap_th

If `atm_grey_irradiated_simple_kap_th` is true, then just set `kap_th` = `kap_v/kap_v_div_kap_th`. Only used if > 0.

```
atm_grey_irr_kap_v_div_kap_th = 0
```

atm_grey_irradiated_P_surf

Surface pressure ; set to 1 bar in cgs units.

```
atm_grey_irradiated_P_surf = 1d6
```

atm_grey_irradiated_max_tries

Limit on iterations.

```
atm_grey_irradiated_max_tries = 50
```

trace_atm_grey_irradiated

Trace the grey atmosphere.

```
trace_atm_grey_irradiated = .false.
```

atm_int_errtol

dump_int_atm_info_model_number ¶

Parameters for integrate T(τ) and write atm structure at model number to terminal.

```
atm_int_errtol = 1d-7
dump_int_atm_info_model_number = -1111
```

Parameters for Paczynski_grey. create_atm_max_step_size in units of log10_tau.

```
trace_atm_Paczynski_grey = .false.
Paczynski_atm_R_surf_errtol = 3d-4
create_atm_max_step_size = 0.1d0
```

surface_extra_Pgas ¶

Extra gas pressure at surface. Added to surface pressure from atm. In ergs/cm³.

```
surface_extra_Pgas = 0d0
```

use_atm_PT_at_center_of_surface_cell ¶

The surface boundary conditions for pressure and temperature, compare the model values at the center of the surface cell to values derived from the P and T returned by the atm module. If this flag is true, then the atm values are directly used. If false, then the values from the atm are treated as being for the outer boundary of the surface cell, and those values are used to estimate corresponding values for the cell center for comparison to the model values at the cell center.

Most cases will have this flag false. An example that sets this flag true is a case in which are using a special boundary condition (BC) routine to force a certain entropy for the surface cell. In that situation, it is better to have the special BC directly return P and T for the center of cell 1 to produce the desired entropy.

```
use_atm_PT_at_center_of_surface_cell = .false.
```

use_compression_outer_BC ¶

gradient of compression vanishes at surface

```
see Grott, Chernigovski, Glatzel, 2005.
d_dm(d_dm(r^2*v)) = 0 at surface
by continuity, this is d_dm(d_dt(1/rho)) = 0 at surface
finite volume form is
(1/rho(1) - 1/rho_start(1)) = (1/rho(2) - 1/rho_start(2))
this BC determines the density for surface cell.
```

```
use_compression_outer_BC = .false.
```

use_momentum_outer_BC ¶

use `P_surf` from atm to set pressure gradient at surface in momentum equation calculate `v(1)` based on pressure difference `P_surf - P(1)`

```
use_momentum_outer_BC = .false.
```

use_zero_Pgas_outer_BC ¶

use `Psurf = Radiation_Pressure(T_start(1))`

```
use_zero_Pgas_outer_BC = .false.
```

use_zero_dLdm_outer_BC ¶

use `L(1) = L(2)` for T outer BC

```
use_zero_dLdm_outer_BC = .false.
```

use_T_Paczynski_outer_BC ¶

`T_surf^4` is set to $L / (8 * \pi * \text{boltz_sig} * R^2)$

```
use_T_Paczynski_outer_BC = .false.
```

use_T_black_body_outer_BC ¶

`T_surf` is set to `Tsurf_factor * T_black_body(L_surf, R_surf)`

```
use_T_black_body_outer_BC = .false.
```

use_fixed_vsurf_outer_BC ¶**fixed_vsurf ¶**

`v` at outer boundary of model is set to be `fixed_vsurf`

```
use_fixed_vsurf_outer_BC = .false.  
fixed_vsurf = 0
```

use_fixed_L_for_BB_outer_BC ¶**fixed_L_for_BB_outer_BC ¶**

for use_T_black_body_outer_BC and use_T_Paczynski_outer_BC

```
use_fixed_L_for_BB_outer_BC = .false.
fixed_L_for_BB_outer_BC = 0
```

tau_for_L_BB ¶

determines location where get L for use with BB outer BCs use the L at outermost location where tau >= tau_for_L_BB

```
tau_for_L_BB = -1
```

Tsurf_factor ¶

used when use_momentum_outer_BC T_surf is set to Tsurf_factor*T_black_body(L_surf,R_surf)

```
Tsurf_factor = 1
```

irradiation_flux ¶

column_depth_for_irradiation ¶

```
irradiation_flux = 0
column_depth_for_irradiation = -1
```

mass gain or loss ¶

mass_change ¶

Rate of accretion (Msun/year). Negative for mass loss. This only applies when the wind_scheme = ''.

```
mass_change = 0d0
```

Enhanced mass loss due to rotation as in Heger, Langer, and Woosley, 2000, ApJ, 528:368-396.

```
Mdot = Mdot_no_rotation/(1 - Osurf/Osurf_crit)^mdot_omega_power
```

where

```
Osurf = angular velocity at surface
Osurf_crit^2 = (1 - Gamma_edd)*G*M/R^3
```

```
Gamma_edd = kappa*L/(4 pi c G M), Eddington factor
```

Typical value for `mdot_omega_power` = 0.43.

mdot_omega_power ¶

Set to 0 to disable this feature.

```
mdot_omega_power = 0.43d0
```

max_rotational_mdot_boost ¶

This limits the rotational boost.

```
max_rotational_mdot_boost = 1d4
```

max_mdot_jump_for_rotation ¶

Don't increase prev mdot by more than this. NOTE: use `vcrit_max_years_for_timestep` with this.

```
max_mdot_jump_for_rotation = 2
```

lim_trace_rotational_mdot_boost ¶

Output to terminal if boost > this.

```
lim_trace_rotational_mdot_boost = 1d99
```

rotational_mdot_boost_fac ¶

Increase mdot.

```
rotational_mdot_boost_fac = 1d5
```

rotational_mdot_kh_fac ¶

Kelvin-helmholtz boost.

```
rotational_mdot_kh_fac = 0.3d0
```

surf_avg_tau_min ¶

Use mass avg starting from this optical depth.

```
surf_avg_tau_min = 1
```

surf_avg_tau ¶

Use mass avg down to this optical depth.

```
surf_avg_tau = 100
```

hot_wind_scheme ¶

hot_wind_Wolf-Rayet_scheme ¶

cool_wind_RGB_scheme ¶

cool_wind_AGB_scheme ¶

This section replaces the old “RGB_wind_scheme” and “AGB_wind_scheme” with temperature-dependent hot_wind and cool_wind. You can still use the RGB and AGB wind scheme as before, the functionality remains.

Now you can also select a hot wind scheme that takes effect *above* some temperature, set by hot_wind_full_on_T. Similarly, the cool wind scheme has temperature controls that set the temperature *below* which they are relevant (cool_wind_full_on_T).

As before, an empty string ‘’ means no wind.

The wind “eta” values, which are constant scaling factors, have all renamed *_wind_eta -> *_scaling_factor.

Here is an example of how to translate an existing inlist from the old style to the new:

Before	After
RGB_wind_scheme = 'Reimers'	cool_wind_RGB_scheme = 'Reimers'
Reimers_wind_eta = 0.1	Reimers_scaling_factor = 0.1
AGB_wind_scheme = 'Blocker'	cool_wind_AGB_scheme = 'Blocker'
Blocker_wind_eta = 0.5	Blocker_scaling_factor = 0.5
RGB_to_AGB_wind_switch = 1d-4	RGB_to_AGB_wind_switch = 1d-4
	! only use the cool_wind_scheme
	cool_wind_full_on_T = 1d10 !K
	hot_wind_full_on_T = 1.1d10 !K
	hot_wind_scheme = ''

suggested hot and cool wind schemes follow but any valid wind option will work for either hot or cool.

Empty string means no wind

Suggested hot wind options:

- ‘Kudritzki’

- ‘Vink’

Suggested cool wind options:

- ‘Reimers’
- ‘Blöcker’
- ‘de Jager’
- ‘van Loon’
- ‘Nieuwenhuijzen’

For now the ‘Dutch’ scheme can be used in either capacity.

```
hot_wind_scheme = ''
cool_wind_RGB_scheme = ''
cool_wind_AGB_scheme = ''
```

cool_wind_full_on_T ¶

hot_wind_full_on_T ¶

use only cool wind schemes for $T_{\text{phot}} < \text{cool_wind_full_on_T}$ use only hot wind schemes for $T_{\text{phot}} > \text{hot_wind_full_on_T}$ if $\text{cool_wind_full_on_T} \neq \text{hot_wind_full_on_T}$ then ramp between these limits requires $\text{hot_wind_full_on_T} > \text{cool_wind_full_on_T}$

```
cool_wind_full_on_T = 0.8d4
hot_wind_full_on_T = 1.2d4
```

RGB_to_AGB_wind_switch ¶

If center hydrogen abundance is < 0.01 and center helium abundance by mass is less than `RGB_to_AGB_wind_switch`, then system will use `AGB_wind_scheme` rather than `RGB_wind_scheme`.

```
RGB_to_AGB_wind_switch = 1d-4
```

The code will automatically choose between an RGB wind and an AGB wind. The following names for the different schemes are recognized:

- ‘Reimers’
- ‘Blocker’
- ‘de Jager’
- ‘van Loon’
- ‘Nieuwenhuijzen’
- ‘Kudritzki’
- ‘Vink’
- ‘Dutch’
- ‘Stern51’
- ‘Grafener’
- ‘other’ — experimental

Reimers_scaling_factor ¶

Reimers mass loss for red giants.

D. Reimers “Problems in Stellar Atmospheres and Envelopes” Baschek, Kegel, Traving (eds), Springer, Berlin, 1975, p. 229.

Parameter for mass loss by Reimers wind prescription. Reimers mdot is $\eta \cdot 4 \cdot 10^{-13} \cdot L \cdot R / M$ (Msun/year), with L, R, and M in solar units. Typical value is 0.5.

```
Reimers_scaling_factor = 0
```

Blocker_scaling_factor = 0 ¶

Blocker’s mass loss for AGB stars.

T. Blocker “Stellar evolution of low and intermediate-mass stars” A&A 297, 727-738 (1995).

Parameter for mass loss by Blocker’s wind prescription. Blocker mdot is $\eta \cdot 4 \cdot 10^{-9} \cdot M^{-2.1} \cdot L^{2.7} \cdot 4 \cdot 10^{-13} \cdot L \cdot R / M$ (Msun/year), with L, R, and M in solar units. Typical value is 0.1d0.

```
Blocker_scaling_factor = 0
```

de_Jager_scaling_factor ¶

de Jager mass loss for various applications. de Jager, C., Nieuwenhuijzen, H., & van der Hucht, K. A. 1988, A&AS, 72, 259. Parameter for mass loss by de Jager wind prescription.

```
de_Jager_scaling_factor = 0d0
```

van_Loon_scaling_factor ¶

see van Loon et al. 2005, A&A, 438, 273 “An empirical formula for the mass-loss rates of dust-enshrouded red supergiants and oxygen-rich Asymptotic Giant Branch stars”

```
van_Loon_scaling_factor = 0d0
```

Kudritzki_scaling_factor ¶

Radiation driven winds of hot stars. See Kudritzki et al, Astron. Astrophys. 219, 205-218 (1989).

```
Kudritzki_scaling_factor = 0d0
```

Nieuwenhuijzen_scaling_factor ¶

See Nieuwenhuijzen, H.; de Jager, C. 1990, A&A, 231, 134.

```
Nieuwenhuijzen_scaling_factor = 0d0
```


Vink_scaling_factor ¶

Vink, J.S., de Koter, A., & Lamers, H.J.G.L.M., 2001, A&A, 369, 574. “Mass-loss predictions for O and B stars as a function of metallicity”

```
Vink_scaling_factor = 0d0
```

Grafener_scaling_factor ¶

Grafener, G. & Hamann, W.-R. 2008, A&A 482, 945 contributed to mesa by Nilou Afsari

```
Grafener_scaling_factor = 0d0
```

Dutch_scaling_factor ¶

The “Dutch” wind scheme for massive stars combines results from several papers, all with authors mostly from the Netherlands.

The particular combination we use is based on Glebbeek, E., et al, A&A 497, 255-264 (2009) [more Dutch authors!]

For $T_{\text{eff}} > 1e4$ and surface $H > 0.4$ by mass, use Vink et al 2001 Vink, J.S., de Koter, A., & Lamers, H.J.G.L.M., 2001, A&A, 369, 574.

For $T_{\text{eff}} > 1e4$ and surface $H < 0.4$ by mass, use Nugis & Lamers 2000 Nugis, T., & Lamers, H.J.G.L.M., 2000, A&A, 360, 227 Some folks use 0.8 for non-rotating mdoels (Maeder & Meynet, 2001).

```
Dutch_scaling_factor = 0d0
```

Dutch_wind_lowT_scheme ¶

For $T_{\text{eff}} < 1e4$

Use de Jager if Dutch_wind_lowT_scheme = 'de Jager' de Jager, C., Nieuwenhuijzen, H., & van der Hucht, K. A. 1988, A&AS, 72, 259.

Use van Loon if Dutch_wind_lowT_scheme = 'van Loon' van Loon et al. 2005, A&A, 438, 273.

Use Nieuwenhuijzen if Dutch_wind_lowT_scheme = 'Nieuwenhuijzen' Nieuwenhuijzen, H.; de Jager, C. 1990, A&A, 231, 134

```
Dutch_wind_lowT_scheme = 'de Jager'
```

Stern51_scaling_factor ¶

wind scheme from Stern

```
Stern51_scaling_factor = 0d0
```

use_accreted_material_j ¶

Angular momentum of accreted material.

```
use_accreted_material_j = .false.
```

If false, then accreted material is given j so that it is rotating at the same angular velocity as the surface. If true, then accreted material is given j = accreted_material_j.

```
accreted_material_j = 0
```

no_wind_if_no_rotation ¶

Use this to delay start of wind until after have started rotation.

```
no_wind_if_no_rotation = .false.
```

min_wind ¶

Min wind in Msun/year > 0; ignore this limit if it is <= 0. e.g., might have low level wind even when normal scheme doesn't call for any.

```
min_wind = 0d0
```

max_wind ¶

Max wind in Msun/year > 0; ignore this limit if it is <= 0.

```
max_wind = 0d0
```

For critical rotation mass loss Redo step as needed to find mdot that brings model to just below critical. if max_mdot_redo_cnt > 0, and surf_w_div_w_crit > surf_w_div_w_crit_limit, then recompute the step while increasing mdot, until surf_w_div_w_crit < surf_w_div_w_crit_limit. Once an upper limit for mdot is found, the solution for mdot is further refined by bisection until it is computed to a tolerance of surf_w_div_w_crit_tol. During iterations, mdot is adjusted alternately by multiplication by mdot_revise_factor, and by adjusting it by implicit_mdot_boost*mdot_initial, where mdot_initial is the value of mdot at the first iteration. This is done to deal with mass accreting stars, where mdot might need to change sign for the star to remain below critical.

```
max_mdot_redo_cnt = 0
min_years_dt_for_redo_mdot = 0
surf_w_div_w_crit_limit = 0.99d0
```

```
surf_w_div_w_crit_tol = 0.05d0
mdot_revise_factor = 1.1d0
implicit_mdot_boost = 0.1d0
```

implicit wind computation. ¶

max_tries_for_implicit_wind ¶

The implicit method will modify the mass transfer rate and redo the step until it either finds a solution, or the number of tries hits `max_tries_for_implicit_wind`. If `max_tries_for_implicit_wind = 0`, the wind computation is explicit, meaning that the value of `mdot` is set using values at the start of the step. This only applies when `mdot < 0`.

```
max_tries_for_implicit_wind = 0
```

iwind_tolerance ¶

Tolerance for which a solution is considered valid. A solution is valid if

```
abs(explicit_mdot - implicit_mdot) <
    abs(implicit_mdot)*iwind_tolerance
```

where

```
explicit_mdot = mstar_dot at start of step
implicit_mdot = mstar_dot at end of step
```

```
iwind_tolerance = 1d-3
```

iwind_lambda ¶

If do not satisfy tolerance, redo with a different `mdot` as follows:

```
mstar_dot = explicit_mdot + &
    iwind_lambda*(implicit_mdot - explicit_mdot)
```

```
iwind_lambda = 1d0
```

remove H wind ¶

remove_H_wind_mdot ¶

This wind removes surface material until reaching a target total H mass for the star. Max rate of removal in Msun/year ; only applies if this is > 0.

```
remove_H_wind_mdot = 0d0
```

remove_H_wind_H_mass_limit ¶

This wind removes surface material until reaching a target total H mass for the star. Turn off this wind when total H mass < this limit (Msun units).

```
remove_H_wind_H_mass_limit = 0
```

super_eddington_scaling_factor ¶

For super eddington wind we use L_{edd} averaged by mass to optical depth $\tau = \text{surf_avg_tau}$.

```
super_eddington_scaling_factor = 0
```

super_eddington_wind_Ledd_factor ¶

Parameter for mass loss driven by super Eddington luminosity. Divide L by this factor when computing super Eddington wind, e.g., if this is 2, then only get wind when $L/2 > L_{\text{edd}}$.

```
super_eddington_wind_Ledd_factor = 1
```

wind_boost_full_off_L_div_Ledd ¶

Boost off for $L/L_{\text{edd}} \leq$ this (set large to disable this). This alternative form is used when `super_eddington_scaling_factor == 0`.

```
wind_boost_full_off_L_div_Ledd = 1.5d0
```

wind_boost_full_on_L_div_Ledd ¶

Do max boost for $L/L_{\text{edd}} \geq$ this. This alternative form is used when `super_eddington_scaling_factor == 0`.

```
wind_boost_full_on_L_div_Ledd = 5
```

super_eddington_wind_max_boost ¶

Multiply wind \dot{m} by up to this amount. This alternative form is used when `super_eddington_scaling_factor == 0`.

```
super_eddington_wind_max_boost = 1
```

trace_super_eddington_wind_boost ¶

Send super eddington wind information to terminal.

```
trace_super_eddington_wind_boost = .false.
```

mass_change_full_on_dt ¶

mass_change_full_off_dt ¶

These params provide the option to turn off mass change when have very small timesteps. Between mass_change_full_on_dt and mass_change_full_off_dt mass change is gradually reduced. Units in seconds.

```
mass_change_full_on_dt = 1d-99  
mass_change_full_off_dt = 1d-99
```

trace_dt_control_mass_change ¶

```
trace_dt_control_mass_change = .false.
```

min_abs_mdot_for_change_limits ¶

Only apply limits if $\text{abs}(\text{prev mdot}) > \text{this limit}$. These limit the change in mdot from one step to the next.

```
min_abs_mdot_for_change_limits = 1d-14
```

max_abs_mdot_factor ¶

Only allow $\text{abs}(\text{mdot})$ to increase by this factor per timestep.

```
max_abs_mdot_factor = 2
```

min_abs_mdot_factor ¶

Only allow $\text{abs}(\text{mdot})$ to decrease by this factor per timestep.

```
min_abs_mdot_factor = 0.5d0
```

max_star_mass_for_gain ¶

Automatic stops for mass loss/gain in Msun units (negative means ignore this parameter). Turn off mass gain when star mass reaches this limit.

```
max_star_mass_for_gain = -1
```

min_star_mass_for_loss ¶

Automatic stops for mass loss/gain in Msun units (negative means ignore this parameter). Turn off mass loss when star mass reaches this limit.

```
min_star_mass_for_loss = -1
```

max_T_center_for_any_mass_loss ¶

No mass loss for T center > this.

```
max_T_center_for_any_mass_loss = 2d9
```

max_T_center_for_full_mass_loss ¶

No reduction in mass loss for T center <= this. This must be <= max_T_center_for_full_mass_loss. Reduce mass loss rate to 0 as T center climbs from max_for_full to max_for_any. The idea behind this is that during final stages of burning, there is so little time left in the life of the star, that any mass loss to winds will be negligible, but the inclusion of that insignificant mass loss can actually make convergence more difficult, so you are better off without it.

```
max_T_center_for_full_mass_loss = 1d9
```

wind_envelope_limit ¶

Winds automatically shut off when the hydrogen rich envelope mass is less than this limit. The value of h1_boundary_limit defines what is considered to be hydrogen poor. Mass in Msun units.

```
wind_envelope_limit = -1
```

rlo_scaling_factor ¶

Amplitude of mass loss. “rlo” wind scheme provides a simple radius-determined-wind with exponential increase.

```
rlo_scaling_factor = 0
```

rlo_wind_min_L ¶

Only on when L > this limit. (Lsun)

```
rlo_wind_min_L = 1d-6
```

rlo_wind_max_Teff ¶

Only on when Teff < this limit.

```
rlo_wind_max_Teff = 1d99
```

rlo_wind_roche_lobe_radius ¶

Only on when R > this (Rsun).

```
rlo_wind_roche_lobe_radius = 0.40d0
```

rlo_wind_base_mdot ¶

Base rate of mass loss when R = roche lobe radius (Msun/year).

```
rlo_wind_base_mdot = 1d-3
```

rlo_wind_scale_height ¶

Determines exponential growth rate of mass loss (Rsun).

```
rlo_wind_scale_height = 1d-1
```

roche_lobe_xfer_full_on ¶

Full accretion when R/RL <= this. Limit accretion when Roche lobe is nearing full (only with rlo_scaling_factor > 0).

```
roche_lobe_xfer_full_on = 0.5d0
```

roche_lobe_xfer_full_off ¶

No accretion when R/RL >= this.

```
roche_lobe_xfer_full_off = 1.0d0
```

nova_scaling_factor ¶

Amplitude of wind. “nova” wind is scheme used in Kato and Hachisu, ApJ 437:802-826, 1994. (eqn 23). This only applies when nova_scaling_factor > 0.

```
nova_scaling_factor = 0
```

nova_wind_b

Wind parameter

```
nova_wind_b = 0
```

nova_wind_max_Teff

Only on when Teff < this limit.

```
nova_wind_max_Teff = 0
```

nova_wind_min_L

only on when L > this limit. (Lsun)

```
nova_wind_min_L = 0
```

nova_min_Teff_for_accretion

When nova_scaling_factor $\neq 0$ and Teff < this and L > nova_wind_min_L, no accretion.

```
nova_min_Teff_for_accretion = 0
```

nova_roche_lobe_radius

units in R_{sun}

```
nova_roche_lobe_radius = 0
```

nova_RLO_mdot

roche lobe overflow mdot, Msun/year

```
nova_RLO_mdot = 0
```

flash_wind_mdot

Rate of mass ejection in Msun/year. “flash” wind is scheme used in Kato, Saio, and Hachisu, ApJ 340:509-517, 1989. This only applies when flash_wind_mdot > 0.


```
flash_wind_mdot = -1
```

flash_wind_starts ¶

Wind starts when $R \geq$ this limit (Rsun units).

```
flash_wind_starts = -1
```

flash_wind_declines ¶

Wind starts to decline when $R \leq$ this limit (Rsun units).

```
flash_wind_declines = -1
```

flash_wind_full_off ¶

Wind full off when $R \leq$ this limit (Rsun units).

```
flash_wind_full_off = -1
```

controls for adjust_mass ¶

max_logT_for_k_below_const_q ¶

max_q_for_k_below_const_q ¶

min_q_for_k_below_const_q ¶

Move `k_below_const_q` inward from surface until $q(k) \leq \text{max_q}$. Then continue moving inward until reach $\log T(k) \geq \text{max_logT}$ or $q(k) \leq \text{min_q}$.

```
max_logT_for_k_below_const_q = 5
max_q_for_k_below_const_q = 1.0d0
min_q_for_k_below_const_q = 0.999d0
```

max_logT_for_k_const_mass ¶

max_q_for_k_const_mass ¶

min_q_for_k_const_mass ¶

Move `k_below_const_q` inward from `k_below_const_q+1` until $q(k) \leq \text{max_q}$. Then continue moving inward until reach $\log T(k) \geq \text{max_logT}$ or $q(k) \leq \text{min_q}$.

```
max_logT_for_k_const_mass = 6
max_q_for_k_const_mass = 1.0d0
min_q_for_k_const_mass = 0.995d0
```

composition controls ¶

accrete_same_as_surface ¶

If true, composition of accreted material is identical to the current surface composition.

```
accrete_same_as_surface = .true.
```

accrete_given_mass_fractions ¶

If true, use the following mass fractions – they must add to 1.0.

```
accrete_given_mass_fractions = .false.
```

num_accretion_species ¶

Up to max_num_accretion_species.

```
num_accretion_species = 0
```

accretion_species_id ¶

Isotope name as defined in chem_def.

```
accretion_species_id(1) = ''
```

accretion_species_xa ¶

mass fraction

```
accretion_species_xa(1) = 0
```

otherwise, use the following composition

accretion_h1 ¶

Hydrogen mass fraction.

```
accretion_h1 = 0
```

accretion_h2 ¶

If no h2 in current net, then this is automatically added to h1.

```
accretion_h2 = 0
```

accretion_he3 ¶

he3 mass fraction

```
accretion_he3 = 0
```

accretion_he4 ¶

he4 mass fraction

```
accretion_he4 = 0
```

accretion_zfracs = ¶

One of the following identifiers for different Z fractions from chem_def.

- AG89_zfracs = 1, Anders & Grevesse 1989
- GN93_zfracs = 2, Grevesse & Noels 1993
- GS98_zfracs = 3, Grevesse & Sauval 1998
- L03_zfracs = 4, Lodders 2003
- AGS05_zfracs = 5, Asplund, Grevesse & Sauval 2005

or set `accretion_zfracs = 0` to use the following list of z fractions

```
accretion_zfracs = -1
```

accretion_dump_missing_metals_into_heaviest ¶

this controls the treatment metals that are not included in the current net. if this flag is true, then the mass fractions of missing metals are added to the mass fraction of the most massive metal included in the net. if this flag is false, then the mass fractions of the metals in the net are renormalized to make up for the total mass fraction of missing metals.

```
accretion_dump_missing_metals_into_heaviest = .true.
```

Special list of z fractions. If you use these, they must add to 1.0.

```
z_fraction_li = 0  
z_fraction_be = 0
```

```

z_fraction_b = 0
z_fraction_c = 0
z_fraction_n = 0
z_fraction_o = 0
z_fraction_f = 0
z_fraction_ne = 0
z_fraction_na = 0
z_fraction_mg = 0
z_fraction_al = 0
z_fraction_si = 0
z_fraction_p = 0
z_fraction_s = 0
z_fraction_cl = 0
z_fraction_ar = 0
z_fraction_k = 0
z_fraction_ca = 0
z_fraction_sc = 0
z_fraction_ti = 0
z_fraction_v = 0
z_fraction_cr = 0
z_fraction_mn = 0
z_fraction_fe = 0
z_fraction_co = 0
z_fraction_ni = 0
z_fraction_cu = 0
z_fraction_zn = 0

```

lgT_lo_for_set_new_abundances ¶

lgT_hi_for_set_new_abundances ¶

Composition controls for set_new_abundances.

```

lgT_lo_for_set_new_abundances = 5.2d0
lgT_hi_for_set_new_abundances = 5.5d0

```

pure_fe56_limit ¶

Pure fe56 for base of ns envelope. If mass fraction of fe56 > this, convert cell to pure fe56.

```

pure_fe56_limit = 0.999999d0

```

mesh adjustment ¶

max_allowed_nz ¶

Maximum number of grid points allowed.

```

max_allowed_nz = 8000

```

remesh_max_allowed_logT ¶

Turn off remesh if any cell has logT > this.

```
remesh_max_allowed_logT = 1d99
```

mesh_max_allowed_ratio ¶

Must be >= 2.5. Max ratio for mass of adjacent cells. If have ratio exceeding this, split the larger cell.

```
mesh_max_allowed_ratio = 2.5d0
```

max_delta_x_for_merge ¶

Don't merge neighboring cells if any abundance differs by more than this.

```
max_delta_x_for_merge = 0.1d0
```

mesh_delta_coeff ¶

A larger value increases the max allowed deltas and decreases the number of grid points. and a smaller does the opposite. E.g., you'll roughly double the number of grid points if you cut mesh_delta_coeff in half. Don't expect it to exactly double the number however since other parameters in addition to gradients also influence the details of the grid spacing.

```
mesh_delta_coeff = 1.0d0
```

mesh_delta_coeff_for_highT ¶

Use different mesh_delta_coeff at higher temperatures.

```
mesh_delta_coeff_for_highT = 3.0d0
```

logT_max_for_standard_mesh_delta_coeff ¶

Use mesh_delta_coeff for center logT <= this. This value should be less than logT_min_for_highT_mesh_delta_coeff.

```
logT_max_for_standard_mesh_delta_coeff = 9.0d0
```

logT_min_for_highT_mesh_delta_coeff ¶

Use mesh_delta_coeff_for_highT for center logT >= this. Linearly interpolate in logT for intermediate center temperatures.

```
logT_min_for_highT_mesh_delta_coeff = 9.5d0
```

mesh_Pgas_div_P_exponent ¶

Multiply mesh_delta_coeff by (Pgas/Ptotal) to this power.

```
mesh_Pgas_div_P_exponent = 0
```

mesh_delta_coeff_pre_ms ¶

Multiply mesh_delta_coeff by this when center $X_H > 0.5$ and $\lg_{LH} < \lg_L - 1$.

```
mesh_delta_coeff_pre_ms = 1
```

max_dq ¶

Max size for cell as fraction of total mass.

```
max_dq = 1d-2
```

min_dq ¶

Min size for cell as fraction of total mass.

```
min_dq = 1d-14
```

min_dq_for_xa ¶

Min size for splitting because of composition gradient.

```
min_dq_for_xa = 1d-14
```

mesh_min_dlnR ¶

Limit on difference in $\ln R$ across cell for mesh refinement. Do not make this smaller than about 1d-14 or will fail with numerical problems.

```
mesh_min_dlnR = 1d-9
```

merge_if_dlnR_too_small ¶

If true, mesh adjustment will force merge if difference in $\ln R$ across cell is too small.

```
merge_if_dlnR_too_small = .false.
```

mesh_min_dr_div_dRstar ¶

Limit on relative radial extent for mesh refinement. $dRstar = s\% r(1) - s\% R_center$ Don't split if $dr/dRstar$ would drop below this limit.

```
mesh_min_dr_div_dRstar = -1
```

merge_if_dr_div_dRstar_too_small ¶

If true, mesh adjustment will force merge if dr_div_dRstar too small.

```
merge_if_dr_div_dRstar_too_small = .true.
```

mesh_min_dr_div_cs ¶

Limit (in seconds) on sound crossing time for mesh refinement. Don't split if sound crossing time would drop below this limit.

```
mesh_min_dr_div_cs = -1
```

merge_if_dr_div_cs_too_small ¶

If true, mesh adjustment will force merge if dr_div_cs too small.

```
merge_if_dr_div_cs_too_small = .true.
```

max_center_cell_dq ¶

Largest allowed dq at center.

```
max_center_cell_dq = 1d-7
```

max_surface_cell_dq ¶

Largest allowed dq at surface.

```
max_surface_cell_dq = 1d-12
```

max_num_subcells ¶

Limits number of new cells from 1 old one.

```
max_num_subcells = 2
```

max_num_merge_cells ¶

Limits number of old cells to merge into 1 new one.

```
max_num_merge_cells = 2
```

mesh_adjust_use_quadratic ¶

Linear or quadratic reconstruction polynomials for mesh adjustments.

```
mesh_adjust_use_quadratic = .true.
```

mesh_adjust_get_T_from_E ¶

If true, then use internal energy conservation to set new temperature. If false, just use average temperature based on reconstruction polynomials.

```
mesh_adjust_get_T_from_E = .true.
```

P_function_weight ¶

Pressure gradient, $P_function = P_function_weight * \log_{10}(P)$.

```
P_function_weight = 40
```

T_function1_weight ¶

Temperature gradient, $T_function1 = T_function1_weight * \log_{10}(T)$. NOTE: The T gradient mesh controls below seems to be necessary to allow burning that starts off center to be able to reach the center. You can see this in the `pre_zahb test_suite` case if you try running it without the T function. The center temperature will fail to rise.

```
T_function1_weight = 110
```

T_function2_weight ¶

T_function2_param ¶

```
T_function2 = T_function2_weight * log10(T / (T + T_function2_param))
```


Largest change in `T_function2` happens around $T = T_function2_param$. Default value puts this in the envelope ionization region.

```
T_function2_weight = 0
T_function2_param = 2d4
```

R_function_weight ¶

R_function_param ¶

log radius gradient

```
R_function = R_function_weight*log10(1 + (r/Rsun)/R_function_param)
```

```
R_function_weight = 0
R_function_param = 1d-4
```

R_function2_weight ¶

R_function2_param1 ¶

R_function2_param2 ¶

```
R_function2 = R_function2_weight*min(R_function2_param1,max(R_function
```

where R_{star} = radius of outer edge of model.

```
R_function2_weight = 0
R_function2_param1 = 0.4d0
R_function2_param2 = 0
```

R_function3_weight ¶

radius gradient

```
R_function3 = R_function3_weight*(r/Rstar)
```

```
R_function3_weight = 0
```

M_function_weight ¶

M_function_param ¶

log mass gradient

```
M_function = M_function_weight*log10(1 + (m/Msun)/M_function_param)
```

```
M_function_weight = 0  
M_function_param = 1d-6
```

gradT_function_weight ¶

gradT gradient, gradT_function = gradT_function_weight*gradT

```
gradT_function_weight = 0
```

log_tau_function_weight ¶

log_tau gradient (optical depth)

```
log_tau_function = log_tau_function_weight*log10(tau)
```

```
log_tau_function_weight = 0
```

log_kap_function_weight ¶

log_kap gradient (optical depth)

```
log_kap_function = log_kap_function_weight*log10(kap)
```

```
log_kap_function_weight = 0
```

omega_function_weight ¶

omega gradient (rotation omega in rad/sec)

```
omega_function = omega_function_weight*log10(omega)
```

```
omega_function_weight = 0
```

gam_function_weight**gam_function_param1****gam_function_param2**

For extra resolution around liquid/solid transition.

```
gam = plasma interaction parameter
gam_function = gam_function_weight*tanh((gam - gam_function_param1)/ga
```

```
gam_function_weight = 0
gam_function_param1 = 170
gam_function_param2 = 20
```

xa_function_species**xa_function_weight**

Mass fraction gradients.

```
xa_function = xa_function_weight*log10(xa + xa_function_param),
```

Up to num_xa_function of these - see star_def for value of num_xa_function. 0 length string means skip, otherwise name of nuclide as defined in chem_def. weight <= 0 means skip.

```
xa_function_species(:) = ''
xa_function_weight(:) = 0
```

```
xa_function_species(1) = 'he4'
xa_function_weight(1) = 30
xa_function_param(1) = 1d-2
```

xa_mesh_delta_coeff

Useful if you want to increase mesh_delta_coeff during advanced burning. If

xa_function_species(j) has the largest atomic number in current set of species, then multiply mesh_delta_coeff by xa_mesh_delta_coeff(j).

```
xa_mesh_delta_coeff(:) = 1
```

“Indirect” mesh controls work by increasing sensitivity in selected regions. They work in the same way as `mesh_delta_coeff` – values less than 1.0 mean smaller allowed jumps in mesh functions and hence smaller grid points and higher resolution. But whereas `mesh_delta_coeff` applies uniformly to all cells, the “extra” coefficients can vary in value from one cell to the next.

`xtra_coef_above_xtrans`

`xtra_coef_below_xtrans`

Multiply `mesh_delta_coeff` near any change in most abundant species by this factor. Value < 1 gives increased resolution.

```
xtra_coef_above_xtrans = 1
xtra_coef_below_xtrans = 1
```

`xtra_dist_above_xtrans`

`xtra_dist_below_xtrans`

Increase resolution up to this distance away from the abundance transition with distance measured in units of the pressure scale height at the boundary.

```
xtra_dist_above_xtrans = 0.2d0
xtra_dist_below_xtrans = 0.2d0
```

`mesh_logX_species`

`mesh_logX_min_for_extra`

Increase resolution at points with large $\text{abs}(\text{dlogX}/\text{dlogP})$; $\text{logX} = \text{log}_{10}(\text{X mass fraction})$.

```
mesh_logX_species(1) = ''
mesh_logX_min_for_extra(1) = -6
```

`mesh_dlogX_dlogP_extra(1)`

`mesh_dlogX_dlogP_full_on(1)`

`mesh_dlogX_dlogP_full_off(1)`

Only increase resolution if $\text{logX} \geq \text{mesh_logX_min_for_extra}$. Make $\text{mesh_dlogX_dlogP_extra} < 1$ for smaller allowed change in logP and hence higher resolution. Full effect if $\text{abs}(\text{dlogX}/\text{dlogP}) \geq \text{mesh_dlogX_dlogP_full_on}$. No effect if $\text{abs}(\text{dlogX}/\text{dlogP}) \leq \text{mesh_dlogX_dlogP_full_off}$. Up to `num_mesh_logX` of these (see `star_def` for value of `num_mesh_logX`).

```
mesh_dlogX_dlogP_extra(1) = 1
mesh_dlogX_dlogP_full_on(1) = 2
mesh_dlogX_dlogP_full_off(1) = 1
```

Multiply `mesh_delta_coeff` near convection zone boundary (czb) by the following factors. Value < 1 gives increased resolution.

`xtra_coef_czb_full_on`

`xtra_coef_czb_full_off`

The center mass fraction of he4 is used to control this extra coefficient. The default settings limit the application to after center he4 is depleted.

- if center he4 < `xtra_coef_czb_full_on`, then use xtra coef's
- if center he4 > `xtra_coef_czb_full_off`, then don't use xtra coef's

a more verbose form of these names might be the following:

```
xtra_coef_czb_full_on_if_center_he4_below_this
xtra_coef_czb_full_off_if_center_he4_above_this
```

```
xtra_coef_czb_full_on = 1d-4
xtra_coef_czb_full_off = 0.1d0
```

`xtra_coef_{above | below}_{lower | upper}_{nonburn | hburn | heburn | zburn}_czb`

Make these < 1 to increase resolution.

`xtra_dist_{above | below}_{lower | upper}_{nonburn | hburn | heburn | zburn}_czb`

Increase resolution up to this distance away from the convective zone boundary, with distance measured in units of the pressure scale height at the boundary.

```
xtra_coef_a_l_nb_czb = 1
xtra_dist_a_l_nb_czb = 0.2d0
xtra_coef_b_l_nb_czb = 1
xtra_dist_b_l_nb_czb = 0.2d0
```

```
xtra_coef_a_u_nb_czb = 1
xtra_dist_a_u_nb_czb = 0.2d0
xtra_coef_b_u_nb_czb = 1
xtra_dist_b_u_nb_czb = 0.2d0
```

```
xtra_coef_a_l_hb_czb = 1
xtra_dist_a_l_hb_czb = 0.2d0
xtra_coef_b_l_hb_czb = 1
xtra_dist_b_l_hb_czb = 0.2d0
```

```
xtra_coef_a_u_hb_czb = 1
xtra_dist_a_u_hb_czb = 0.2d0
xtra_coef_b_u_hb_czb = 1
xtra_dist_b_u_hb_czb = 0.2d0
```

```
xtra_coef_a_l_heb_czb = 1
xtra_dist_a_l_heb_czb = 0.2d0
xtra_coef_b_l_heb_czb = 1
xtra_dist_b_l_heb_czb = 0.2d0
```

```
xtra_coef_a_u_heb_czb = 1
xtra_dist_a_u_heb_czb = 0.2d0
xtra_coef_b_u_heb_czb = 1
xtra_dist_b_u_heb_czb = 0.2d0
```

```
xtra_coef_a_l_zb_czb = 1
xtra_dist_a_l_zb_czb = 0.2d0
xtra_coef_b_l_zb_czb = 1
xtra_dist_b_l_zb_czb = 0.2d0
```

```
xtra_coef_a_u_zb_czb = 1
xtra_dist_a_u_zb_czb = 0.2d0
xtra_coef_b_u_zb_czb = 1
xtra_dist_b_u_zb_czb = 0.2d0
```

xtra_coef_scz_above_{nonburn | hburn | heburn | zburn}_cz ¶

Make these < 1 to increase resolution in semiconvective region adjacent to convective region. e.g., xtra_coef_scz_above_nb_cz is extra coef for semiconvectize zone above a non-burn convective zone.

```
xtra_coef_scz_above_nb_cz = 1
xtra_coef_scz_above_hb_cz = 1
xtra_coef_scz_above_heb_cz = 1
xtra_coef_scz_above_zb_cz = 1
```

Multiply mesh_delta_coeff in overshooting regions by the following factors. Value < 1 gives increased resolution.

xtra_coef_os_full_on ¶

xtra_coef_os_full_off ¶

The center mass fraction of he4 is used to control this extra coefficient. The default settings limit the application to after center he4 is depleted.

- if center he4 < xtra_coef_os_full_on, then use xtra_coef_coef's

- if center he4 > xtra_coef_os_full_off, then don't use xtra_coef_coef's

```
xtra_coef_os_full_on = 1d-4
xtra_coef_os_full_off = 0.1d0
```

xtra_coef_os_{above | below}_{nonburn | hburn | heburn | zburn} ¶

Make these < 1 to increase resolution.

```
xtra_coef_os_above_nonburn = 1
xtra_coef_os_below_nonburn = 1
xtra_coef_os_above_burn_h = 1
xtra_coef_os_below_burn_h = 1
xtra_coef_os_above_burn_he = 1
xtra_coef_os_below_burn_he = 1
xtra_coef_os_above_burn_z = 1
xtra_coef_os_below_burn_z = 1
```

xtra_dist_os_{above | below}_{nonburn | hburn | heburn | zburn} ¶

Continue to increase resolution for this distance beyond the edge of the overshooting region, with distance measured in units of the pressure scale height at the edge of the overshooting region. This applies to both edges of the overshooting region.

```
xtra_dist_os_above_nonburn = 0.2d0
xtra_dist_os_below_nonburn = 0.2d0
xtra_dist_os_above_burn_h = 0.2d0
xtra_dist_os_below_burn_h = 0.2d0
xtra_dist_os_above_burn_he = 0.2d0
xtra_dist_os_below_burn_he = 0.2d0
xtra_dist_os_above_burn_z = 0.2d0
xtra_dist_os_below_burn_z = 0.2d0
```

Increase resolution at points with large $\text{abs}(\text{dlog_eps}/\text{dlogP})$ for nuclear power eps (ergs/g/sec). At any particular location, only use eps nuc category with max local value e.g., only use mesh_dlog_pp_dlogP_extra at points where pp is the max burn source.

mesh_dlog_eps_min_for_extra ¶

Only increase resolution if $\text{log_eps} \geq \text{mesh_dlog_eps_min_for_extra}$.

```
mesh_dlog_eps_min_for_extra = -2
```

mesh_dlog_eps_dlogP_full_on ¶

Full effect if $\text{abs}(\text{dlog_eps}/\text{dlogP}) \geq \text{mesh_dlog_eps_dlogP_full_on}$.

```
mesh_dlog_eps_dlogP_full_on = 4
```

mesh_dlog_eps_dlogP_full_off ¶

No effect if $\text{abs}(\text{dlog_eps}/\text{dlogP}) \leq \text{mesh_dlog_eps_dlogP_full_off}$.

```
mesh_dlog_eps_dlogP_full_off = 1
```

Multiply the allowed change between adjacent cells by the following factors; (small factor => smaller allowed change => more cells).

pp and cno burning

```
mesh_dlog_pp_dlogP_extra = 0.25d0
mesh_dlog_cno_dlogP_extra = 0.25d0
```

triple alpha, c, n, and o burning

```
mesh_dlog_3alf_dlogP_extra = 0.25d0
mesh_dlog_burn_c_dlogP_extra = 0.25d0
mesh_dlog_burn_n_dlogP_extra = 0.25d0
mesh_dlog_burn_o_dlogP_extra = 0.25d0
```

ne, na, and mg burning

```
mesh_dlog_burn_ne_dlogP_extra = 0.25d0
mesh_dlog_burn_na_dlogP_extra = 0.25d0
mesh_dlog_burn_mg_dlogP_extra = 0.25d0
```

c12+c12, c12+o16, and o16+o16 burning

```
mesh_dlog_cc_dlogP_extra = 0.25d0
mesh_dlog_co_dlogP_extra = 0.25d0
mesh_dlog_oo_dlogP_extra = 0.25d0
```

si to iron alog alpha chain burning

```
mesh_dlog_burn_si_dlogP_extra = 0.25d0
mesh_dlog_burn_s_dlogP_extra = 0.25d0
mesh_dlog_burn_ar_dlogP_extra = 0.25d0
mesh_dlog_burn_ca_dlogP_extra = 0.25d0
mesh_dlog_burn_ti_dlogP_extra = 0.25d0
mesh_dlog_burn_cr_dlogP_extra = 0.25d0
mesh_dlog_burn_fe_dlogP_extra = 0.25d0
```

photodisintegration burning


```
mesh_dlog_pnhe4_dlogP_extra = 0.25d0
mesh_dlog_other_dlogP_extra = 0.25d0
mesh_dlog_photo_dlogP_extra = 1
```

convective_bdy_weight ¶

convective_bdy_dq_limit ¶

convective_bdy_min_dt_yrs ¶

Mesh function to enhance resolution near convective boundaries including regions that are newly nonconvective because of moving boundary. EXPERIMENTAL

```
convective_bdy_weight = 0
convective_bdy_dq_limit = 1d-4
convective_bdy_min_dt_yrs = 1d-3
```

trace_mesh_adjust_error_in_conservation ¶

If true, report relative errors for total PE, KE, and IE. (potential, kinetic, internal).

```
trace_mesh_adjust_error_in_conservation = .false.
```

okay_to_remesh ¶

If false, then no remeshing.

```
okay_to_remesh = .true.
```

remesh_dt_limit ¶

No remesh if $dt < \text{remesh_dt_limit}$, in seconds.

```
remesh_dt_limit = -1
```

remesh_log_L_nuc_burn_min ¶

No mesh adjustments when $\log_{10}(\text{L_nuc_burn_total})$ is less than this. By default, this turns off mesh changes during the early pre-MS.

```
remesh_log_L_nuc_burn_min = -50
```

use_split_merge_amr ¶

```
use_split_merge_amr = .false.
```

use_split_merge_amr_log_zoning ¶

if true, target is even grid spacing in logr if false, target is even grid spacing in r

```
split_merge_amr_log_zoning = .true.
```

split_merge_amr_nz_baseline ¶

```
split_merge_amr_nz_baseline = 1000
```

split_merge_amr_MaxLong ¶

split cell if ratio of actual/desired dr is > this; ignore if <= 0

```
split_merge_amr_MaxLong = 1.5d0
```

split_merge_amr_MaxShort ¶

merge cell if ratio of desired/actual dr is > this; ignore if <= 0

```
split_merge_amr_MaxShort = 4d0
```

merge_amr_max_abs_du_div_cs ¶

```
merge_amr_max_abs_du_div_cs = 0.1d0
```

merge_amr_inhibit_at_jumps ¶

```
merge_amr_inhibit_at_jumps = .false.
```

split_merge_amr_dq_min ¶

do not split if dq for cell is < this

```
split_merge_amr_dq_min = 1d-14
```

split_merge_amr_max_iters ¶

```
split_merge_amr_max_iters = 100
```

split_merge_amr_okay_to_split_1 ¶

split_merge_amr_okay_to_split_nz ¶

```
split_merge_amr_okay_to_split_1 = .true.  
split_merge_amr_okay_to_split_nz = .true.
```

equal_split_density_amr ¶

```
equal_split_density_amr = .false.
```

trace_split_merge_amr ¶

```
trace_split_merge_amr = .false.
```

nuclear reaction controls ¶

default_net_name ¶

Name of base reaction network. Each net corresponds to a file in \$MESA_DIR/data/net_data/nets. Look in that directory to see your network options, or learn how to create your own net.

```
default_net_name = 'basic.net'
```

screening_mode ¶

- empty string means no screening
- 'classic' : DeWitt, Graboske, Cooper, “Screening Factors for Nuclear Reactions. I. General Theory”, ApJ, 181:439-456, 1973. Graboske, DeWitt, Grossman, Cooper, “Screening Factors for Nuclear Reactions. II. Intermediate Screening and Astrophysical Applications”, ApJ, 181:457-474, 1973.
- 'extended' : extends the Graboske method using results from Alastuey and Jancovici (1978), along with plasma parameters from Itoh et al (1979) for strong screening.
- 'salpeter' : weak screening only. following Salpeter (1954), with equations (4-215) and (4-221) of Clayton (1968).

```
screening_mode = 'extended'
```

net_logTcut_lo ¶

strong rates are zero $\log T < \log T_{\text{cut_lo}}$ use default from net if this is ≤ 0

```
net_logTcut_lo = -1
```

net_logTcut_lim

strong rates cutoff smoothly for $\log T < \log T_{\text{cut_lim}}$ use default from net if this is ≤ 0

```
net_logTcut_lim = -1
```

max_abar_for_burning

if $\bar{a} > \text{this}$, suppress all burning e.g., if want an “inert” core heavy elements, set this to 55 or, if want to turn off the net, set this to -1

```
max_abar_for_burning = 199
```

dxdt_nuc_factor

Control for abundance changes by burning. Changes dxdt_nuc (rate of change of abundances) without changing the rates or eps_nuc (rate of energy generation).

```
dxdt_nuc_factor = 1
```

weak_rate_factor

all weak rates are multiplied by this factor

```
weak_rate_factor = 1
```

reaction_neuQs_factor

all neutrino Q factors are multiplied by this factor

```
reaction_neuQs_factor = 1
```

nonlocal_NiCo_kap_gamma

```
nonlocal_NiCo_kap_gamma = 0
```

nonlocal_NiCo_decay_heat

if true, do non-local deposition of gamma-ray energy from Ni56 and Co56 decays. only for approx nets including co56. intended for use with stripped envelope supernovae.

```
nonlocal_NiCo_decay_heat = .false.
```

dtau_gamma_NiCo_decay_heat

```
dtau_gamma_NiCo_decay_heat = 1d0
```

max_logT_for_net

```
max_logT_for_net = 10.2d0
```

element diffusion

gravitational settling and chemical diffusion.

show_diffusion_info

terminal output for diffusion

```
show_diffusion_info = .false.
```

show_diffusion_substep_info

terminal output for diffusion

```
show_diffusion_substep_info = .false.
```

show_diffusion_timing

show time for each call on diffusion

```
show_diffusion_timing = .false.
```

do_element_diffusion

determines whether or not we do element diffusion

```
do_element_diffusion = .false.
```

diffusion_dt_limit

no element diffusion if dt < this limit (in seconds)

```
diffusion_dt_limit = 3.15d7
```

diffusion_use_paquette

if true, use atomic diffusion coefficients according to Paquette et al. (1986). if false, use Stanton & Murillo (2016) for diffusion coefficients. (Paquette coefficients still used for electron-ion because Stanton & Murillo did not do calculations for attractive potentials.)

```
diffusion_use_paquette = .false.
```

diffusion_use_iben_macdonald

if true, use diffusion coefficients similar to Iben & MacDonald (1985). if false, use Stanton & Murillo (2016) for diffusion coefficients. this was previously called `diffusion_use_pure_coulomb`.

```
diffusion_use_iben_macdonald = .false.
```

diffusion_use_cgs_solver

if false, solve the system of equations descibed by Thoul et al. (1994) if true, solve the unmodified Burgers equations in cgs units

```
diffusion_use_cgs_solver = .true.
```

cgs_thermal_diffusion_eta_full_on

cgs_thermal_diffusion_eta_full_off

When `diffusion_use_cgs_solver = .true.` for `eta < cgs_thermal_diffusion_eta_full_on`, includes the heat flow vector terms in the Burgers equations. Then smoothly turns off use of these terms so that they are not included for `eta > cgs_thermal_diffusion_eta_full_off`, since these terms are problematic when distribution function become non-Maxwellian.

```
cgs_thermal_diffusion_eta_full_on = 0d0  
cgs_thermal_diffusion_eta_full_off = 2d0
```

do_Ne22_sedimentation_heating

if true, include heating from sedimentation of Ne22 when element diffusion is on. Your net must include Ne22 for this to work. For best results, Ne22 should be treated as its own diffusion class. This will affect white dwarf cooling times. See also `eps_Ne22_sedimentation_factor`

```
do_Ne22_sedimentation_heating = .false.
```

diffusion_min_dq_at_surface

treat at least this much at surface as a single cell for purposes of diffusion

```
diffusion_min_dq_at_surface = 1d-9
```

diffusion_min_T_at_surface ¶

treat cells at surface with $T < \text{this}$ as a single cell for purposes of diffusion default should be large enough to ensure hydrogen ionization

```
diffusion_min_T_at_surface = 1d4
```

diffusion_min_dq_ratio_at_surface ¶

combine cells at surface until have total mass \geq this factor times the next cell below them this helps with surface boundary condition for diffusion by putting large cell at surface

```
diffusion_min_dq_ratio_at_surface = 10
```

diffusion_dt_div_timescale ¶

dt is at most this fraction of timescale. Each stellar evolution step can be divided into many substeps for diffusion. The substep timescale is set by rates of flow in and out for each species in each cell. The substep size, dt, is initially set to $\text{timescale} * \text{diffusion_dt_div_timescale}$.

```
diffusion_dt_div_timescale = 1
```

diffusion_min_num_substeps ¶

Max substep dt is total time divided by this.

```
diffusion_min_num_substeps = 1
```

diffusion_max_iters_per_substep ¶

If the substep requires too many iterations, the substep time is decreased for a retry.

```
diffusion_max_iters_per_substep = 10
```

diffusion_max_retries_per_substep ¶

If the substep requires too many retries, diffusion fails and forces a retry for the star.

```
diffusion_max_retries_per_substep = 10
```

diffusion_tol_correction_max ¶

diffusion_tol_correction_norm ¶

Tolerances for newton iterations. Corrections smaller will be treated as converged. Corrections larger will cause another newton iteration.

```
diffusion_tol_correction_max = 1d-1  
diffusion_tol_correction_norm = 1d-3
```

diffusion_min_X_hard_limit ¶

tolerance for negative mass fraction errors errors larger will cause retry; errors smaller will be corrected.

```
diffusion_min_X_hard_limit = -1d-3
```

diffusion_X_total_atol ¶

diffusion_X_total_rtol ¶

tolerances for errors in total species conservation errors larger will cause retry; errors smaller will be corrected.

```
diffusion_X_total_atol = 1d-9  
diffusion_X_total_rtol = 1d-6
```

diffusion_upwind_abs_v_limit ¶

switch to upwind for i at face k if $\text{abs}(v(i,k)) > \text{this limit}$ mainly for use with radiative levitation where get very much higher velocities

```
diffusion_upwind_abs_v_limit = 1d99
```

diffusion_v_max ¶

Max velocity (cm/sec). We can get extremely large velocities in the extreme outer envelope that cause problems numerically without really effecting the results, so we allow a max for the velocities that should help the numerics without changing the results. Note: change `diffusion_v_max` to at least 1d-2 when using radiative levitation.

```
diffusion_v_max = 1d-3
```

D_mix_ignore_diffusion ¶

Diffusion is turned off in core and surface convection zones, since it is overwhelmed by other mixing there. `D_mix_ignore_diffusion` roughly defines the mixing coefficient below which diffusion is included again. The code finds the location where `D_mix` falls to this value, backs up some, and turns on diffusion from there onward.


```
D_mix_ignore_diffusion = 1d5
```

diffusion_gamma_full_off ¶

diffusion_gamma_full_on ¶

gamma_full_on <= gamma_full_off Shut off diffusion for large gamma (i.e. for gamma >= gamma_full_off). Gradually decrease diffusion as gamma increases from full_on to full_off. Allow normal diffusion for gamma <= gamma_full_on. Default is diffusion off when get well into liquid regime.

```
diffusion_gamma_full_off = 175  
diffusion_gamma_full_on = 150
```

diffusion_T_full_on ¶

diffusion_T_full_off ¶

T_full_on >= T_full_off Shut off diffusion for small T (i.e., for T <= T_full_off) Gradually decrease diffusion as T decreases from T_full_on to T_full_off. Allow normal diffusion for T >= T_full_on.

```
diffusion_T_full_on = 1d3  
diffusion_T_full_off = 1d3
```

diffusion_calculates_ionization ¶

If diffusion_calculates_ionization is false, MESA uses typical charges for a set of representative species as defined in diffusion_class_typical_charge and diffusion_class_representative for all points rather than calculating the ionization from the local conditions.

```
diffusion_calculates_ionization = .true.
```

diffusion_nsmooth_typical_charge ¶

smoothing over charge

```
diffusion_nsmooth_typical_charge = 10
```

diffusion_SIG_factor ¶

diffusion_GT_factor ¶

factors for playing with SIG and GT terms for concentration diffusion and advection

```
diffusion_SIG_factor = 1d0
diffusion_GT_factor = 1d0
```

diffusion_AD_dm_full_on ¶

diffusion_AD_dm_full_off ¶

diffusion_AD_boost_factor ¶

artificial concentration diffusion near surface (mainly for radiative levitation) Msun units for `full_on` and `full_off` boost only used if > 0

```
diffusion_AD_dm_full_on = -1
diffusion_AD_dm_full_off = -1
diffusion_AD_boost_factor = 0
```

diffusion_Vlimit_dm_full_on ¶

diffusion_Vlimit_dm_full_off ¶

in Msun units artificial velocity limitation near surface (mainly for radiative levitation)

```
diffusion_Vlimit_dm_full_on = -1
diffusion_Vlimit_dm_full_off = -1
```

diffusion_Vlimit ¶

In units of local cell crossing velocity (only used if > 0). When full on, limit $\text{abs}(v) \leq V_{\text{limit}} \cdot dr/dt$, cell size dr , substep time dt .

```
diffusion_Vlimit = 0
```

diffusion_min_T_for_radaccel ¶

diffusion_max_T_for_radaccel ¶

If T between these limits, then include radiative levitation at that location. Calculation of radiative levitation is costly, so only use it where necessary. Note: change `diffusion_v_max` to at least $1d-2$ when using radiative levitation.

Note that radiative levitation requires OP calculations of g_{rad} for each class, and only 17 elements are supported (H, He, C, N, O, Ne, Na, Mg, Al, Si, S, Ar, Ca, Cr, Mn, Fe, Ni). If you want to include radiative levitation, your options are:

- Define diffusion classes such that all class representatives are among the 17 elements listed above.
- Use a net with only elements from the 17 above, and set `diffusion_use_full_net = .true.`

```
diffusion_min_T_for_radaccel = 0
diffusion_max_T_for_radaccel = 0
```

diffusion_min_Z_for_radaccel ¶

diffusion_max_Z_for_radaccel ¶

If Z between these limits, then include radiative levitation for that element. Calculation of radiative levitation is costly, so only use it where necessary. e.g., limit to Fe and Ni by `min_Z = 26` and `max_Z = 28`

```
diffusion_min_Z_for_radaccel = 0
diffusion_max_Z_for_radaccel = 1000
```

diffusion_screening_for_radaccel ¶

Include screening for radiative levitation.

```
diffusion_screening_for_radaccel = .true.
```

diffusion_use_full_net ¶

If true, don't lump elements into classes for diffusion. Instead, each isotope in the network is treated as its own separate class. This can cause significant slowdowns for large nets, so it is off by default. This works for nets with up to 100 isotopes; larger nets require lumping into classes.

```
diffusion_use_full_net = .false.
```

diffusion_num_classes ¶

Number of representative classes of species for diffusion calculations. (maximum of 100)

```
diffusion_num_classes = 5
```

diffusion_class_representative(:) ¶

isotope names for diffusion representatives

```
diffusion_class_representative(1) = 'h1'
diffusion_class_representative(2) = 'he3'
diffusion_class_representative(3) = 'he4'
diffusion_class_representative(4) = 'o16'
diffusion_class_representative(5) = 'fe56'
```

diffusion_class_A_max(:) ¶

atomic number A. in ascending order. species goes into 1st class with $A_{\text{max}} \geq$ species A

```
diffusion_class_A_max(1) = 2
diffusion_class_A_max(2) = 3
diffusion_class_A_max(3) = 4
diffusion_class_A_max(4) = 16
diffusion_class_A_max(5) = 10000
```

diffusion_class_typical_charge(:) ¶

Typical charges for use if `diffusion_calculates_ionization` is false Use charge 21 for Fe in the sun, from Thoul, Bahcall, and Loeb (1994), ApJ, 421, 828.

```
diffusion_class_typical_charge(1) = 1
diffusion_class_typical_charge(2) = 2
diffusion_class_typical_charge(3) = 2
diffusion_class_typical_charge(4) = 8
diffusion_class_typical_charge(5) = 21
```

diffusion_class_factor(:) ¶

Arbitrarily enhance or inhibit diffusion effects by class.

```
diffusion_class_factor(:) = 1d0
```

parameters for ionization solver ¶

diffusion_use_isolve ¶

Activate iterative solver.

```
diffusion_use_isolve = .false.
```

diffusion_rtol_for_isolve ¶

diffusion_atol_for_isolve ¶

Relative and absolute error parameters for iterative solver.

```
diffusion_rtol_for_isolve = 1d-4
diffusion_atol_for_isolve = 1d-5
```

diffusion_maxsteps_for_isolve ¶

Maximum number of steps to take in iterative solver.

```
diffusion_maxsteps_for_iseolve = 1000
```

diffusion_iseolve_solver ¶

Which ode solver to use for iterative.

Options include:

- 'ros2_solver'
- 'rose2_solver'
- 'ros3p_solver'
- 'ros3pl_solver'
- 'rodas3_solver'
- 'rodas4_solver'
- 'rodasp_solver'

```
diffusion_iseolve_solver = 'ros2_solver'
```

diffusion_dump_call_number ¶

debugging info of diffusion at call number

```
diffusion_dump_call_number = -1
```

eos controls ¶

more eos controls can be found in `star_job.defaults`

use_eosDT_ideal_gas ¶

if true, then eos is ideal gas eos as implemented by HELMEOS

```
use_eosDT_ideal_gas = .false.
```

use_eosDT_HELMEOS ¶

if true, then eos is as implemented by HELMEOS alone; no blending.

```
use_eosDT_HELMEOS = .false.
```

eosDT_HELMEOS_include_radiation ¶

if `use_eosDT_HELMEOS`, then this flag is passed as arg to control whether include radiation.

```
eosDT_HELMEOS_include_radiation = .true.
```

eosDT_HELMEOS_always_skip_elec_pos

if use_eosDT_HELMEOS, then this flag is passed as arg to control whether skip electrons and positrons. if true, then always skip them else skip only for low T, low density situations

```
eosDT_HELMEOS_always_skip_elec_pos = .false.
```

eosDT_HELMEOS_always_include_elec_pos

if use_eosDT_HELMEOS, then this flag is passed as arg to control whether skip electrons and positrons. if true, then always include them else include only for low T, low density situations

```
eosDT_HELMEOS_always_include_elec_pos = .false.
```

use_eosPTEH_for_low_density

use_eosPTEH_for_high_Z

Z_for_all_PTEH

Z_for_any_PTEH

overwrites the control of same name in the `Eos_General_Info` structure if use_eosPTEH_for_low_density is true, use PTEH for $\log Rho < -10$ or so. i.e., regions not covered by OPAL/SCVH. if use_eosPTEH_for_high_Z is true, use PTEH in place of HELM for Z above OPAL/SCVH range $Z \geq Z_{\text{for_all_PTEH}}$, then use pure PTEH with no blend $Z < Z_{\text{for_any_PTEH}}$, then do not use PTEH blend with PTEH for intermediate Z. if use_eosDT2, blend with DT2. else blend with OPAL/SCVH data.

```
use_eosPTEH_for_low_density = .true.
use_eosPTEH_for_high_Z = .true.
Z_for_all_PTEH = 0.040d0
Z_for_any_PTEH = 0.039d0
```

use_eosDT2

overwrites the control of same name in the `Eos_General_Info` structure if use_eosPTEH_for_high_Z and use_eosDT2, use DT2 in place of OPAL/SCVH for cases with $Z < Z_{\text{for_all_PTEH}}$

```
use_eosDT2 = .false.
```

use_eosELM

logT_max_for_ELM

overwrites the control of same name in the `Eos_General_Info` structure if `use_eosELM` is true, use ELM in high Rho and T regions where otherwise use HELM. limit use of ELM to $\log T \leq \log T_{\text{max_for_ELM}}$

```
use_eosELM = .false.  
logT_max_for_ELM = 8.8d0
```

use_fixed_XZ_for_eos ¶

for debugging

```
use_fixed_XZ_for_eos = .false.
```

fixed_X_for_eos ¶

if `use_fixed_XZ_for_eos`, then pass this value to eos instead of actual X.

```
fixed_X_for_eos = -1
```

fixed_Z_for_eos ¶

if `use_fixed_XZ_for_eos`, then pass this value to eos instead of actual Z.

```
fixed_Z_for_eos = -1
```

opacity controls ¶

more opacity controls can be found in `star_job.defaults`

cubic_interpolation_in_X ¶

type of interpolation in X

```
cubic_interpolation_in_X = .false.
```

cubic_interpolation_in_Z ¶

type of interpolation in Z

```
cubic_interpolation_in_Z = .false.
```

include_electron_conduction ¶

add conduction opacities to radiative opacities

```
include_electron_conduction = .true.
```

use_simple_es_for_kap ¶

for experiments with simple electron scattering if true, $\text{opacity} = 0.2 * (1 + x)$

```
use_simple_es_for_kap = .false.
```

use_Type2_opacities ¶

Type2 opacities for extra C/O during and after He burning. To use Type2 opacities one needs to specify a base metallicity, Zbase, which gives the metal abundances previous to any CO enhancement. In regions where central hydrogen is above a given threshold, or the metallicity is not significantly higher than Zbase, Type1 tables are used instead, with blending regions to smoothly transition from one to the other. Why not just use Type2 all the time? Is it a performance reason to still support Type1s?

No – the Type1 tables cover a wider range of X and have a higher resolution in Z for each X.

The Type1 tables are for (X,Z) pairs from the following sets:

The 10 Type1 X's are 0.0, 0.1, 0.2, 0.35, 0.5, 0.7, 0.8, 0.9, 0.95, 1-Z The 13 Type1 Z's are 0.0, 1e-4, 3e-4, 1e-3, 2e-3, 4e-3, 1e-2, 2e-2, 3e-2, 4e-2, 6e-2, 8e-1, 1e-1

The Type2 tables are for (X,Z) pairs from the following more limits sets: The 5 Type2 X's are 0.0, 0.03, 0.10, 0.35, 0.70 The 8 Type2 Z's are 0.00, 0.001, 0.004, 0.01, 0.02, 0.03, 0.05, 0.1

There are 130 (X,Z) combinations for Type1 and only 40 for Type2. So Type2 gives you C/O enhancement at a cost of lower resolution in (X,Z). The Type1 tables also cover the full possible range of X from 0.0 to 1-Z, whereas the Type2 tables stop at a max X of 0.70. Extrapolating the Type2 tables to higher X is not reliable, so we switch over to using Type1 data instead. In this case, Type1 opacities are computed using Zbase instead of the actual metallicity.

```
use_Type2_opacities = .false.
```

Zbase ¶

the base metallicity for the Type2 kap evaluations.

```
Zbase = -1
```

use_Zbase_for_Type1_blend ¶

If true, then if use_Type2_opacities = .true. Type1 opacities will be computed using Zbase instead of Z. Ignored if use_Type2_opacities = .false.

```
use_Zbase_for_Type1_blend = .true.
```

kap_Type2_full_off_X ¶

kap_Type2_full_on_X ¶

switch to Type1 if X too large Type2 is full off for $X \geq \text{kap_Type2_full_off_X}$ Type2 can be full on for $X \leq \text{kap_Type2_full_on_X}$

```
kap_Type2_full_off_X = 0.71d0
kap_Type2_full_on_X = 0.70d0
```

kap_Type2_full_off_dZ ¶

kap_Type2_full_on_dZ ¶

switch to Type1 if dZ too small ($dZ = Z - Z_{\text{base}}$) Type2 is full off for $dZ \leq \text{kap_Type2_full_off_dZ}$ Type2 can be full on for $dZ \geq \text{kap_Type2_full_on_dZ}$.

```
kap_Type2_full_off_dZ = 0.001d0
kap_Type2_full_on_dZ = 0.01d0
```

X and dZ terms are multiplied to get actual fraction of Type2. The fraction of Type2 is calculated for each cell depending on the X and dZ for that cell. So you can be using Type1 in cells where X is large or dZ is small, while at the same time you can be using Type2 where X is small and dZ is large. When `frac_Type2` is > 0 and < 1 , then both Type1 and Type2 are evaluated and combined linearly as $(1 - \text{frac_Type2}) * \text{kap_type1} + \text{frac_Type2} * \text{kap_type2}$. Add `kap_frac_Type2` to your profile columns list to see `frac_Type2` for each cell.

opacity_max ¶

limit opacities to this value (ignore this is value is < 0)

```
opacity_max = -1
```

opacity_factor ¶

opacities are multiplied by this value

```
opacity_factor = 1
```

min_logT_for_opacity_factor_off ¶

min_logT_for_opacity_factor_on and ¶

max_logT_for_opacity_factor_on ¶

max_logT_for_opacity_factor_off ¶

temperature controls for where the `opacity_factor` is applied if, for example, you only want the opacity factor to apply in the iron bump region you can give a logT range such as

```
min_logT_for_opacity_factor_off = 5.2
min_logT_for_opacity_factor_on = 5.3
max_logT_for_opacity_factor_on = 5.7
max_logT_for_opacity_factor_off = 5.8
```

ignore these if < 0 .

```
min_logT_for_opacity_factor_off = -1
min_logT_for_opacity_factor_on = -1
max_logT_for_opacity_factor_on = -1
max_logT_for_opacity_factor_off = -1
```

if you need cell-by-cell control of opacity factor, set the vector “extra_opacity_factor” using the routine “other_opacity_factor”

OP mono opacities ¶

The OP_mono opacities use data and code from the OP website as modified by Haili Hu. Since the tar.gz file is large (656 MB), it is not included in the standard mesa download.

You can get OP4STARS_1.3.tar.gz [here](#)

Put it any place you want on your disk.

```
gunzip OP4STARS_1.3.tar.gz
tar -xvf OP4STARS_1.3.tar
```

Set the inlist controls for the “mono” directory with the data files. For example, in my case it looks like the following, but you can put the directory anywhere you like – it doesn’t need to be in the mesa/data directory. And the cache file doesn’t need to be in the mono directory.

```
op_mono_data_path = '/Users/bpaxton/OP4STARS_1.3/mono'
op_mono_data_cache_filename = '/Users/bpaxton/OP4STARS_1.3/mono/op_mon
```

op_mono_data_path ¶

if this path is set to the empty string, “”, then it defaults to the environment variable $\$(MESA_OP_MONO_DATA_PATH)$

```
op_mono_data_path = ''
```

op_mono_data_cache_filename ¶

if this is set to the empty string, '', then it defaults to the environment variable \$(MESA_OP_MONO_DATA_CACHE_FILENAME)

```
op_mono_data_cache_filename = ''
```

high_logT_op_mono_full_off

high_logT_op_mono_full_on

low_logT_op_mono_full_off

low_logT_op_mono_full_on

you can select a range of log10T for using op_mono opacities outside that range, the code will use standard opacity tables. for example, you might only use high T limits so that op_mono is only used in the envelope, or you might set both low and high T limits so that op_mono is used around the Fe peak logT but not for other locations in the star.

```
high_logT_op_mono_full_off >= high_logT_op_mono_full_on
high_logT_op_mono_full_on >= low_logT_op_mono_full_on
low_logT_op_mono_full_on >= low_logT_op_mono_full_off
```

```
op_mono opacities full on if
log10T <= high_logT_op_mono_full_on
and
log10T >= low_logT_op_mono_full_on
```

```
op_mono opacities full off if
log10T >= high_logT_op_mono_full_off
or
log10T <= low_logT_op_mono_full_off
```

partially on for other cases

```
high_logT_op_mono_full_off = -1d99
high_logT_op_mono_full_on = -1d99
```

```
low_logT_op_mono_full_off = -1d99
low_logT_op_mono_full_on = -1d99
```

op_mono_min_X_to_include

skip iso if mass fraction < this

```
op_mono_min_X_to_include = 1d-20
```

use_op_mono_alt_get_kap

if true, call the op_mono_alt_get_kap routine instead of op_mono_get_kap. see mesa/kap/public/kap_lib.f for details about these routines.

```
use_op_mono_alt_get_kap = .false.
```

kap_phot_factor_for_kap_floor

kap_phot_step_factor

min_kap_floor_step_factor

min_for_kap_floor

tau_use_kap_floor

```
kap_phot_factor_for_kap_floor = 1d0
kap_phot_step_factor = 0.05d0
min_kap_floor_step_factor = 0.333d0
min_for_kap_floor = 1d-4
tau_use_kap_floor = .false.
```

kap_min_Z_0pt02

kap_min_Z_1pt0

```
kap_min_Z_0pt02 = 1d-20
kap_min_Z_1pt0 = 1d-20
```

kap_fac_X_lo

kap_fac_X_lo_fac

kap_fac_X_hi

kap_fac_X_hi_fac

```
kap_fac_X_lo = -1
kap_fac_X_lo_fac = -1
kap_fac_X_hi = -1
kap_fac_X_hi_fac = -1
```

asteroseismology controls

get_delta_nu_from_scaled_solar

use scaled solar values

```
get_delta_nu_from_scaled_solar = .false.
```

nu_max_sun ¶

solar value of nu_max

```
nu_max_sun = 3100d0
```

delta_nu_sun ¶

solar value of delta_nu

```
delta_nu_sun = 135d0
```

Teff_sun ¶

solar value of Teff

```
Teff_sun = 5777d0
```

delta_Pg_mode_freq ¶

uHz. if <=0, use nu_max from scaled solar value

```
delta_Pg_mode_freq = 0d0
```

Brunt controls ¶

calculate_Brunt_N2 ¶

Only calculate Brunt_N2 if this is true.

```
calculate_Brunt_N2 = .true.
```

brunt_N2_coefficient ¶

Standard N2 is multiplied by this value.

```
brunt_N2_coefficient = 1
```

num_cells_for_smooth_brunt_B ¶

Number of cells on either side to use in weighted smoothing of `brunt_B`.

```
num_cells_for_smooth_brunt_B = 2
```

threshold_for_smooth_brunt_B ¶

Threshold for weighted smoothing of `brunt_B`. Only apply smoothing (controlled by `num_cells_for_smooth_brunt_B`) for contiguous regions where `|brunt_B|` exceeds this threshold. Might be useful for preventing narrow peaks from being excessively broadened by smoothing

```
threshold_for_smooth_brunt_B = 0d0
```

use_brunt_gradmuX_form ¶

For comparison to older codes. Assumes ideal gas plus radiation for `brunt_B`. Uses hydrogen mass fraction to estimate $d\ln\mu = dX/(X + 0.6)$.

```
use_brunt_gradmuX_form = .false.
```

interpolate_rho_for_pulsation_info ¶

If true, then get `rho_face` by interpolating `rho` at cell center. If false, then calculate `rho_face` by $dm/(4*\pi*r^2*dr)$.

```
interpolate_rho_for_pulsation_info = .true.
```

min_magnitude_brunt_B ¶

If set `brunt_B` to 0 if absolute value is $<$ this.

```
min_magnitude_brunt_B = -1d99
```

structure equations ¶

velocity_q_upper_bound ¶

Local override for global `v_flag`. If local `q >` this bound, local `v_flag` is set false, else local `v_flag` is set to global `v_flag`. this lets you force `v = 0` in outer envelope.

```
velocity_q_upper_bound = 1d99
```

velocity_logT_lower_bound ¶

Local override for global `v_flag`. If local `logT` < this bound, local `v_flag` is set false, else local `v_flag` is set to global `v_flag`. this lets you force `v = 0` in outer envelope.

```
velocity_logT_lower_bound = -1d99
```

sponge_max_q_full_on ¶

sponge_min_q_full_off ¶

“sponge” soaks up velocities; full on mean velocities forced to zero. sponge full on for `q >= sponge_max_q_full_on` (== velocities full off) sponge full off for `q <= sponge_min_q_full_off` `sponge_min_q_full_off < sponge_max_q_full_on`

```
sponge_max_q_full_on = -1d99  
sponge_min_q_full_off = -1d99
```

max_dt_yrs_for_velocity_logT_lower_bound ¶

Only apply `velocity_logT_lower_bound` when timestep < this limit.

```
max_dt_yrs_for_velocity_logT_lower_bound = 1d99
```

use_dP_dm_rotation_correction ¶

With rotation, multiply `dP/dm` by `fp_rot` if this flag is true.

```
use_dP_dm_rotation_correction = .true.
```

use_mass_corrections ¶

Gravitational vs baryonic mass corrections. If false, then no distinction between gravitational and baryonic mass. If true, then gravitational mass is calculated using mass corrections. Note: may need to wait for pre-ms model to converged before turning this on.

```
use_mass_corrections = .false.
```

use_sr_sound_speed ¶

SR correction for sound speed.

```
use_sr_sound_speed = .false.
```

use_gr_factors ¶

GR corrections. Currently just for pressure equation.

```
use_gr_factors = .false.
```

use_ODE_var_eqn_pairing ¶

changes the pairing of equations and variables helps with numerical issues in hydro matrix solves

```
use_ODE_var_eqn_pairing = .false.
```

use_dvdt_form_of_momentum_eqn ¶

if true, use $dv/dt = \dots$ form of momentum equation. this replaces the default pressure gradient form. only when `v_flag` is true.

```
use_dvdt_form_of_momentum_eqn = .false.
```

use_Paczynski_term_in_dvdt_eqn ¶

if true, then the $dv/dt = \dots$ form of momentum equation. includes Paczynski correction term in optically thin regions. B. Paczynski, 1969, Acta Astr., vol. 19

```
use_Paczynski_term_in_dvdt_eqn = .false.
```

use_dedt_form_of_energy_eqn ¶

if true, use $de/dt = \dots$ form of energy equation. this replaces the default dL/dm and `eps_grav` form.

```
use_dedt_form_of_energy_eqn = .false.
```

use_ODE_form_of_density_eqn ¶

if true, use ODE $d\ln\rho/dt = \dots$ if false, use algebraic relation $\rho = \text{cell_mass}/\text{cell_vol}$

```
use_ODE_form_of_density_eqn = .false.
```

non_nuc_neu_factor ¶

Multiplies power from non-nuclear reaction neutrinos. i.e., thermal neutrinos such as computed by mesa/neu.

```
non_nuc_neu_factor = 1
```

eps_nuc_factor ¶

Multiplies `eps_nuc` without changing rates or `dxd_t_nuc`. Thus controls energy production without modifying the amount of change in abundances.

```
eps_nuc_factor = 1
```

eps_Ne22_sedimentation_factor ¶

This controls energy production from element diffusion sedimentation of Ne22.

```
eps_Ne22_sedimentation_factor = 1
```

max_abs_eps_nuc ¶

Limit magnitude of `eps_nuc` to this.

```
max_abs_eps_nuc = 1d99
```

fe56ec_fake_factor ¶

min_T_for_fe56ec_fake_factor ¶

Multiplier on ni56 electron capture rate to take isotopes in hardwired networks to more neutron rich isotopes.

```
fe56ec_fake_factor = 1d-7
min_T_for_fe56ec_fake_factor = 3d9
```

eps_grav ¶

In mesa, “`eps_grav`” means $-T \cdot dS/dt$ which is equivalent to $-(dE/dt + P \cdot dV/dt)$ where S is specific entropy, E is specific internal energy, and $V = 1/\rho$. There are several options for how `eps_grav` is calculated. These alternatives are equivalent from the “ideal” physics viewpoint, but they can be very different numerically depending on the situation.

The standard default forms are used if you don’t set any of the following flags. The default when using `ln d` instead of `ln Pgas` as primary variable is

```
eps_grav = -T*cp*((1-grada)*chiT*dlnT_dt - grada*chiRho*dln d_dt)
```

When using `ln Pgas` instead of `ln d`, the default is

```
eps_grav = -T*cp((1-grada*4*Prad/P)*dlnT_dt - grada*Pgas/P*dln Pgas_dt)
```

use_dEdRho_form_for_eps_grav ¶

If true, use $\text{eps_grav} = -(\text{cv} \cdot T \cdot d\ln T/dt + (\rho \cdot dE_{\text{dRho}} - P/\rho) \cdot d\ln \rho/dt)$

```
use_dEdRho_form_for_eps_grav = .false.
```

use_dln_d_t_form_for_eps_grav ¶

If true, use $\text{eps_grav} = -d\epsilon/dt + P/\rho \cdot d\ln \rho/dt$.

```
use_dln_d_t_form_for_eps_grav = .false.
```

use_PdVdt_form_for_eps_grav ¶

If true, use $\text{eps_grav} = -(d\epsilon/dt + P \cdot d(1/\rho)/dt)$. [1/rho = V, specific vol.] With time centering for 1/rho, $P \cdot d(1/\rho)/dt$ becomes $P \cdot (1/\rho - 1/\rho_{\text{start}})/dt$

```
use_PdVdt_form_for_eps_grav = .false.
```

use_lnS_for_eps_grav ¶

If true, use $\text{eps_grav} = -T \cdot DS/Dt$. Note: while this seems like the obvious way to go, it has problems numerically. $\ln S$ is not a basic variable in the way that $\ln T$ and $\ln \rho$ (or $\ln P_{\text{gas}}$) are. I.e., we get $\ln T$ and $\ln \rho$ from the newton solver directly, but we get $\ln S$ by calling the eos using $\ln T$ and $\ln \rho$ as args. Also, in many cases, the cell mass coordinates don't change for the step, making it possible to use the solver value for the increment in the $\ln T$ or $\ln \rho$ directly in estimating the Lagrangian time derivative (e.g., $D\ln T/dt = \Delta \ln T / \Delta t$ instead of $= (\text{new } \ln T - \text{old } \ln T)/dt$). By avoiding the roundoff error in this way, we also get a boost numerically. With $\ln S$ we cannot do that -- we're stuck with $D\ln S/Dt = (\text{new } \ln S - \text{old } \ln S)/dt$. In a perfect world with infinite precision, none of this would matter. But in practice, it turns out to make a significant difference, so unless there are strong reasons otherwise, you should use one of the other schemes for getting ``eps_grav``. One situation where you need to use the $\ln S$ version is deep in a white dwarf where you can have a phase transition that the $\ln S$ form will take care of but the others won't. That particular case is handled using ``Gamma_lnS_eps_grav_full_off`/on`.

```
use_lnS_for_eps_grav = .false.
```

include_dmu_dt_in_eps_grav ¶

The above do not include the contribution from composition changes. In most cases, that is okay (at least it is a common practice!), but for high T, high density situations, you may want to full the full form. to do that, set `include_dmu_dt_in_eps_grav` to true.

This only is relevant when you are not using the $\ln S$ form of `eps_grav`. Since when using $\text{eps_grav} = -T \cdot dS/dt$, the composition effects are already included. otherwise, we calculate the composition term in `eps_grav` as $-dE_{\text{dmu}} \cdot d\mu/dt$ with μ approximated by $\bar{a} / (1 + \bar{z})$ corresponding to complete ionization and dE_{dmu} approximated by $-3/2 \cdot c_{\text{gas}} \cdot T / \mu^2$ (c_{gas} = ideal gas constant; erg/K/mole) where $d\mu/dt$ is 1st order approximation Lagrangian time derivative, $(\mu - \text{prev_mu})/dt$, prev_mu interpolated at same mass coordinate in start-of-step model.

```
include_dmu_dt_in_eps_grav = .false.
```

Gamma_lnS_eps_grav_full_off ¶

Gamma_lnS_eps_grav_full_on ¶

Automatic switch to lnS form for regions with high Gamma (plasma interaction parameter). Set `use_lnS_for_eps_grav` false to use these controls. These are ignored when `use_lnS_for_eps_grav` is true.

```
Gamma_lnS_eps_grav_full_on = 150d0
Gamma_lnS_eps_grav_full_off = 120d0
```

eps_grav_factor ¶

multiply `eps_grav` by this factor

```
eps_grav_factor = 1
```

eps_grav_dt_use_start_values ¶

set true if must use values for `lnT`, `lnd`, or `lnP` from start of step in `d_dt`. e.g., if have made significant changes in abundance profiles in diffusion.

```
eps_grav_dt_use_start_values = .false.
```

eps_grav_time_deriv_separation ¶

Separation (in grid cells) over which `eps_grav` can be time-differenced when `Mstar` changes. The mesh has two major regions - an interior region where the cells are Lagrangian and an outer region where they are homologous (constant $dq = dm/M$). There is also a small transition region between these two. In the Lagrangian region Ds/dt is evaluated with a Lagrangian finite difference in time, while in the homologous region Ds/dt is evaluated with a finite difference in time at constant q plus an advection-like term accounting for the movement of the q boundaries in mass. In the transition region, these two derivatives are combined. This means that at the edges of the transition region, finite differences may cross cell boundaries. This control determines how the mesh for the end of the current timestep is placed to ensure that finite differences cross no more than this many cell boundaries.

```
eps_grav_time_deriv_separation = 1.5d0
```

zero_eps_grav_in_just_added_material ¶

If true, set `eps_grav(k) = 0` for `k < k_below_just_added`. NOTE: this does not mean that Ds/Dt is forced to be 0 in cell k . Instead it simply means we ignore Ds/Dt in the cell's energy calculation.

```
zero_eps_grav_in_just_added_material = .false.
```

min_dxm_Eulerian_div_dxm_removed ¶

Controls for Eulerian or Lagrangian forms of `eps_grav`. Only for mass loss. Specifies a minimum value for the ratio of the mass layer at the surface using Eulerian `eps_grav` (`dxm_Eulerian`) divided by the mass removed in the current step.

```
min_dxm_Eulerian_div_dxm_removed = 2
```

min_dxm_Eulerian_div_dxm_added ¶

Controls for Eulerian or Lagrangian forms of `eps_grav`. Only for mass gain. Specifies a minimum value for the ratio of the mass layer at the surface using Eulerian `eps_grav` (`dxm_Eulerian`) divided by the mass added in the current step.

```
min_dxm_Eulerian_div_dxm_added = 5
```

min_cells_for_Eulerian_to_Lagrangian_transition ¶

Width of eulerian to lagrangian transition region.

```
min_cells_for_Eulerian_to_Lagrangian_transition = 10
```

fix_eps_grav_transition_to_grid ¶

If true, fix the transition region for the computation of `eps_grav` to the transition from Lagrangian to constant in `q` of the grid.

```
fix_eps_grav_transition_to_grid = .false.
```

min_del_T_div_dt ¶

Controls for Lagrangian time derivatives in newly added material only applies to cells with `k < k_below_just_added`. If `del_t_for_just_added(k)/dt < this limit`, then set `del_t_for_just_added(k) = dt*this limit`.

```
min_del_T_div_dt = 1d-10
```

max_num_surf_revisions ¶

Max number of forced reconverges for changes in `surf_lnS`.

```
max_num_surf_revisions = 1
```

max_abs_rel_change_surf_lnS ¶

Force newton reconverge if surf_lnS changed more than this.

```
max_abs_rel_change_surf_lnS = 5d-4
```

trace_force_another_iteration ¶

If true, report when force another iter.

```
trace_force_another_iteration = .false.
```

accel_factor ¶

coefficient for acceleration term in the momentum equation

```
accel_factor = 1
```

extra_power_source ¶

erg/g/sec applied uniformly throughout the model This can be used to push a pre-ms model up the track to lower center temperatures. Can be used simultaneously with inject_extra_ergs_sec and inject_uniform_extra_heat

```
extra_power_source = 0
```

inject_uniform_extra_heat ¶

extra heat in erg g⁻¹ s⁻¹ Added to cells in range min_q_for_uniform_extra_heat to max. Can be used simultaneously with inject_extra_ergs_sec and extra_power_source.

```
inject_uniform_extra_heat = 0
```

min_q_for_uniform_extra_heat ¶

sets bottom of region for inject_uniform_extra_heat

```
min_q_for_uniform_extra_heat = 0
```

max_q_for_uniform_extra_heat ¶

sets top of region for inject_uniform_extra_heat

```
max_q_for_uniform_extra_heat = 1
```

inject_extra_ergs_sec ¶

added to mass equal to grams_for_inject_extra_core_ergs_sec can be used simultaneously with extra_power_source and inject_uniform_extra_heat

```
inject_extra_ergs_sec = 0
```

base_of_inject_extra_ergs_sec ¶

(units: Msun) sets bottom of region for inject_extra_ergs_sec note: actual base is at max of this and the center of the model

```
base_of_inject_extra_ergs_sec = 0
```

total_mass_for_inject_extra_ergs_sec ¶

(units: Msun) sets size of region for inject_extra_ergs_sec

```
total_mass_for_inject_extra_ergs_sec = 0
```

start_time_for_inject_extra_ergs_sec ¶

(units: sec) start time for injecting extra ergs/s

```
start_time_for_inject_extra_ergs_sec = -1d99
```

duration_for_inject_extra_ergs_sec ¶

(units: sec) length of time for injecting extra ergs/s set to negative value to keep injecting indefinitely or until reach target

```
duration_for_inject_extra_ergs_sec = -1
```

inject_until_reach_model_with_total_energy ¶

(units: ergs) target for model total energy usually want to set

duration_for_inject_extra_ergs_sec = -1 for this option. see also:

inject_until_reach_delta_total_energy continue injecting until total energy of model reaches min of inject_until_reach_model_with_total_energy, and

inject_until_reach_delta_total_energy + initial total energy

```
inject_until_reach_model_with_total_energy = 1d99
```

inject_until_reach_delta_total_energy ¶

(units: ergs) target for change in total energy stop injecting when `total_energy - total_energy_initial > this`. usually want to set `duration_for_inject_extra_ergs_sec = -1` for this option. see also: `inject_until_reach_model_with_total_energy` continue injecting until total energy of model reaches min of `inject_until_reach_model_with_total_energy`, and `inject_until_reach_delta_total_energy + initial total energy`

```
inject_until_reach_delta_total_energy = 1d99
```

max_inject_velocity_km_per_sec ¶

```
max_inject_velocity_km_per_sec = 0
```

theta_P ¶

for time weighting P in energy and momentum equations

```
<P> = theta_P*P + (1 - theta_P)*P_start
```

```
theta_P = 1d0
```

use_porosity_with_dPrad_dm_form ¶

```
use_porosity_with_dPrad_dm_form = .false.
```

qmax_zero_non_radiative_luminosity ¶**qmin_freeze_non_radiative_luminosity ¶****use_dPrad_dm_form_of_T_gradient_eqn ¶****use_flux_limiting_with_dPrad_dm_form ¶**

These are for alternatives ways to determine the T gradient. The standard form of the equation is

```
dT/dm = dP/dm * T/P * grad_T, grad_T = dlnT/dlnP from MLT.
```

use hydrostatic value for `dP/dm` in this. this is because of limitations of MLT for calculating `grad_T`. (MLT assumes hydrostatic equilibrium) see comment in K&W chpt 9.1.

The alternatives forms are for dynamic situations where the use of hydrostatic dP/dm is inappropriate. In order of priority,

```
if q(k) > qmin_freeze_non_radiative_luminosity then
  use L_conv from start of step to get L_rad = L - L_conv_start
else if q(k) <= qmax_zero_non_radiative_luminosity then
  simply use L_rad = L
else if (use_dPrad_dm_form_of_T_gradient_eqn)
  if (gradT < gradr) then
    use L_rad = L*gradT/gradr (see, e.g., Cox&Giuli 14.109)
  else
    use L_rad = L
```

With the resulting L_{rad} , determine the expected dT/dm by

```
d_Prada/dm = -kap*L_rad/(clight*area^2) -- see, e.g., K&W (5.12)
```

```
qmax_zero_non_radiative_luminosity = 0d0
qmin_freeze_non_radiative_luminosity = 1d0
use_dPrad_dm_form_of_T_gradient_eqn = .false.
use_flux_limiting_with_dPrad_dm_form = .false.
```

center_energy_pulses ¶

pulse_step_interval ¶

pulse_max_step ¶

pulse_ergs_even_steps ¶

pulse_ergs_odd_steps ¶

```
center_energy_pulses = .false.
pulse_step_interval = 11
pulse_max_step = 1000
pulse_ergs_even_steps = 1d20
pulse_ergs_odd_steps = -1d20
```

for hydro comparison tests (e.g., Sedov)

gamma_low_hydro ¶

off as long as value is ≤ 0

```
gamma_low_hydro = 0d0
```


zero_gravity

if true, then set G to zero

```
zero_gravity = .false.
```

disable_riemann_reconstruction

if true, then set just use bounding cell average P and u at face

```
disable_riemann_reconstruction = .false.
```

constant_L

if true, then $L(k) = L_{\text{center}}$ for all k. disable $dl_n T_{dm}$ equation.

```
constant_L = .false.
```

drag_per_step**drag_per_second**

```
drag_per_step = -1d0  
drag_per_second = -1d0
```

Rayleigh-Taylor Instability**RTI_A****RTI_B****RTI_C****RTI_D****RTI_C_X_factor****RTI_C_X0****RTI_max_alpha****RTI_min_dm_behind_shock_for_full_on****RTI_dm_for_center_alpha_nondecreasing****RTI_energy_floor**

RTI_D_mix_floor

RTI_min_m_for_D_mix_floor

RTI_log_max_boost

RTI_m_full_boost

RTI_m_no_boost

Note that these parameters are not exactly the same

as used by Paul Duffell.

His calibrated D is 2, where mesa uses D = 3 (see mesaIV paper).

Users should try both values since the choice is not clear cut.

```
RTI_A = 1d-3
RTI_B = 2.5d0
RTI_C = 0.2d0
RTI_D = 3d0
```

```
RTI_C_X0_frac = 0.9d0
RTI_C_X_factor = 0d0
```

```
RTI_max_alpha = 0.5d0
RTI_min_dm_behind_shock_for_full_on = 0d0
RTI_dm_for_center_eta_nondecreasing = 0.02d0
RTI_energy_floor = 0d0
RTI_D_mix_floor = 0d0
RTI_min_m_for_D_mix_floor = 0d0
```

```
RTI_log_max_boost = 3d0
RTI_m_full_boost = 4d0
RTI_m_no_boost = 5d0
```

solver controls

the following is from a response on mesa-users to a question about controls for solver tolerances:

The “residual” is the left over difference between the left and right hand sides of the equation we are trying to solve. We do iterations to reduce that, but we are limited by the non-linearity of the problem and the quality of the estimates for the derivatives.

The “correction” is the change in the primary variable that is calculated using good-old Newton’s rule in multiple dimensions — so Jacobian and residuals give a correction that would make the next residual vanish

if the problem were linear and the Jacobian was exact, neither of which are true. So the best we can hope for is that the corrections will get smaller next time.

The “norm” is the average; the “max” is the max. Sometimes you mainly care about the norm and will accept a few outliers. But sometimes you don’t want any really bad outliers, so you want to set a low limit for the max residual or correction as well as the norm.

You might want to try for several iterations with strict tolerances, and then relax them if things are still not converged. For example, you might be willing to live with the larger tolerances, but you’d like to give it a good try at the smaller ones before switching. Also, you might be willing to settle for any-old residual if the corrections have become small enough. You can do that too by relaxing the residual tolerances after a few iterations.

Hope that at least helps with the nomenclature.

I agree with Frank that you should consider the effects of smaller timesteps and more grid points as your main technique — tightening up the tolerances for the solver won’t help if you are taking timesteps that are too large or if you have inadequate grid resolution.

tol_correction_norm ¶

tol_max_correction ¶

“Correction” for variable $x(i,k)$ is scaled change, $dx(i,k)/xscale(i,k)$. these tolerances are for the magnitude of the scaled corrections.

```
tol_correction_norm = 3d-5
tol_max_correction = 3d-3
```

tol_correction_high_T_limit ¶

For very late stages of massive star evolution, need to relax tolerances. If $\max T \geq$ this limit, switch scaling factors.

```
tol_correction_high_T_limit = 1d9
```

tol_correction_norm_high_T ¶

tol_max_correction_high_T ¶

Above `tol_correction_high_T_limit` use these scaling factors.

```
tol_correction_norm_high_T = 3d-3
tol_max_correction_high_T = 3d-1
```

tol_correction_extreme_T_limit ¶

For very late stages of massive star evolution, need to relax tolerances. If $\text{center } T \geq$ this limit, switch scaling factors.

```
tol_correction_extreme_T_limit = 6d9
```

tol_correction_norm_extreme_T ¶

tol_max_correction_extreme_T ¶

For very late stages of massive star evolution, need to relax tolerances. If center T >= this limit, switch scaling factors.

```
tol_correction_norm_extreme_T = 8d-3  
tol_max_correction_extreme_T = 8d-1
```

tol_bad_max_correction ¶

if max_correction > tol_max_correction and no more iterations allowed, then still accept the solution if max_correction <= tol_bad_max_correction. but if max_correction > tol_bad_max_correction, then reject the solution.

```
tol_bad_max_correction = 0d0
```

bad_max_correction_series_limit ¶

If have this many steps in a row with max_correction > tol_max_correction, then do a retry with a smaller timestep.

```
bad_max_correction_series_limit = 2
```

relax_use_gold_tolerances ¶

```
relax_use_gold_tolerances = .false.
```

relax_newton_iterations_limit ¶

relax_newton_iterations_hard_limit ¶

relax_tol_correction_norm ¶

relax_tol_max_correction ¶

relax_tol_residual_norm1 ¶

relax_tol_max_residual1 ¶

relax_iter_for_resid_tol2 ¶

relax_tol_residual_norm2**relax_tol_max_residual2****relax_iter_for_resid_tol3****relax_tol_residual_norm3****relax_tol_max_residual3****relax_maxT_for_gold_tolerances****relax_max_eosPC_frac_for_gold_tolerances**

For use during relax operations. Only used if /= 0.

```
relax_newton_iterations_limit = 0
relax_newton_iterations_hard_limit = 0
```

```
relax_tol_correction_norm = 0d0
relax_tol_max_correction = 0d0
```

```
relax_tol_residual_norm1 = 0d0
relax_tol_max_residual1 = 0d0
relax_iter_for_resid_tol2 = 3
```

```
relax_tol_residual_norm2 = 0d0
relax_tol_max_residual2 = 0d0
relax_iter_for_resid_tol3 = 0
```

```
relax_tol_residual_norm3 = 0d0
relax_tol_max_residual3 = 0d0
relax_maxT_for_gold_tolerances = -1d0
relax_max_eosPC_frac_for_gold_tolerances = 0d0
```

include_L_in_error_est**include_v_in_error_est****include_u_in_error_est**

Some variables can be excluded from calculation of correction norm and max.

```
include_L_in_error_est = .false.
include_v_in_error_est = .false.
include_u_in_error_est = .false.
```

tol_correction_norm_alt

tol_max_correction_alt

If you have several backups in a row, your run is having a near death experience. So as a last hope, try relaxing the correction tolerances. It might help. The code will use these tolerances after 3 or more backups in a row. Once there is a step without a backup, it goes back to the normal tolerances.

```
tol_correction_norm_alt = 1d-3  
tol_max_correction_alt = 1d-2
```

correction_xa_limit

Ignore correction to abundance when calculating correction norm and max if current mass fraction is less than this limit.

```
correction_xa_limit = 5d-3
```

xa_scale

Scaling for abundance variables is $\max(\text{xa_scale}, \text{current mass fraction})$.

```
xa_scale = 1d-5
```

tol_residual_norm1

tol_max_residual1

iter_for_resid_tol2

“residual” for equation is the difference between left and right sides use `tol_residual_norm1` & `tol_max_residual1` at iteration number `iter_for_resid_tol2`, switch to next tolerances.

```
tol_residual_norm1 = 1d-10  
tol_max_residual1 = 1d-9  
iter_for_resid_tol2 = 6
```

tol_residual_norm2

tol_max_residual2

iter_for_resid_tol3

Use `tol_residual_norm2` & `tol_max_residual2` these apply starting at iteration number `iter_for_resid_tol2`. at iteration number `iter_for_resid_tol3`, switch to next tolerances.

```
tol_residual_norm2 = 1d99  
tol_max_residual2 = 1d99  
iter_for_resid_tol3 = 15
```

tol_residual_norm3 ¶

tol_max_residual3 ¶

Use `tol_residual_norm3` & `tol_max_residual3` these apply starting at iteration number `iter_for_resid_tol3`.

```
tol_residual_norm3 = 1d99  
tol_max_residual3 = 1d99
```

If things get worse from one iteration to next, give up. The following are the limits that define “getting worse enough to stop”.

corr_norm_jump_limit ¶

If correction norm increases by this factor or more, quit.

```
corr_norm_jump_limit = 1d99
```

max_corr_jump_limit ¶

If correction max increases by this factor or more, quit.

```
max_corr_jump_limit = 1d6
```

resid_norm_jump_limit ¶

If residual norm increases by this factor or more, quit.

```
resid_norm_jump_limit = 1d99
```

max_resid_jump_limit ¶

If residual max increases by this factor or more, quit.

```
max_resid_jump_limit = 1d6
```

max_iterations_for_jacobian ¶

EXPERIEMENTAL: not working at present. leave at 1. Jacobian is always created fresh for 1st iteration. If this param > 1, then will try to reuse jacobian. After use jacobian this many times, remake it. E.g., if = 2, then

will make a new jacobian for every other iteration. This is automatically = 1 immediately following a backup.

```
max_iterations_for_jacobian = 1
```

convergence_ignore_equL_residuals ¶

```
convergence_ignore_equL_residuals = .false.
```

trace_newton_damping ¶

Send newton damping data to screen.

```
trace_newton_damping = .false.
```

hydro_decsol_switch ¶

small_mtx_decsol ¶

large_mtx_decsol ¶

If current `nvar` \leq `hydro_decsol_switch`, (recall `nvar` = `nvar_hydro` + `species`) then use `small_mtx_decsol` for current step, else use `large_mtx_decsol`.

Options for `small_mtx_decsol` are 'block_thomas_dble' or 'bcyclic_dble'.

Options for `large_mtx_decsol` are 'bcyclic_klu'.

```
hydro_decsol_switch = 99999999  
small_mtx_decsol = 'bcyclic_dble'  
large_mtx_decsol = 'bcyclic_klu'
```

star_bcyclic_do_pivot ¶

Controls whether or not do pivoting in matrix solves in star bcyclic.

```
star_bcyclic_do_pivot = .true.
```

max_tries ¶

Max number newton iterations before give up.

```
max_tries = 25
```

max_tries1 ¶

Max tries on 1st model.

```
max_tries1 = 250
```

max_tries_for_retry ¶

Normal number of retries.

```
max_tries_for_retry = 25
```

max_tries_after_5_retries ¶

Increase number of tries after 5 failed ones.

```
max_tries_after_5_retries = 35
```

max_tries_after_10_retries ¶

Increase number of tries after 10 failed ones.

```
max_tries_after_10_retries = 50
```

max_tries_after_20_retries ¶

Increase number of tries after 20 failed ones.

```
max_tries_after_20_retries = 75
```

max_tries_after_backup ¶

Max tries after first backup.

```
max_tries_after_backup = 25
```

max_tries_after_backup2 ¶

Max tries after second backup.

```
max_tries_after_backup2 = 25
```

retry_limit ¶

Only use if > 0 . In case the solver fails for some reason, it will retry with a smaller timestep. It does up to this many retries for the current step before doing a backup to the previous step.

```
retry_limit = 2
```

redo_limit ¶

Only use if > 0 . Do up to this many redo's for the current step before doing a backup to the previous step.

```
redo_limit = 100
```

newton_itermin ¶

Use at least this many iterations in newton for hydro solve.

```
newton_itermin = 2
```

newton_itermin_until_reduce_min_corr_coeff ¶

Use at least this many iterations in newton before try using small `min_corr_coeff`

```
newton_itermin_until_reduce_min_corr_coeff = 8
```

newton_reduced_min_corr_coeff ¶

For use with `newton_itermin_for_reduce_min_corr_coeff`.

```
newton_reduced_min_corr_coeff = 0.1d0
```

tiny_corr_coeff_limit ¶

scale_correction_norm ¶

corr_param_factor ¶

scale_max_correction ¶

corr_coeff_limit ¶

tiny_corr_factor ¶

see `star/private/star_newton` for info about these

```
tiny_corr_coeff_limit = 5  
scale_correction_norm = 0.1
```

```
corr_param_factor = 10
scale_max_correction = 1d99
corr_coeff_limit = 1d-2
tiny_corr_factor = 2
```

min_xa_hard_limit ¶

min_xa_hard_limit_for_highT ¶

If solver produces mass fraction < this limit, then reject the trial solution. Can optionally relax this limit at high T.

```
min_xa_hard_limit = -1d-5
min_xa_hard_limit_for_highT = -3d-5
```

logT_max_for_xa_hard_limit ¶

Use min_xa_hard_limit for center logT <= this.

```
logT_max_for_min_xa_hard_limit = 9.49d0
```

logT_min_for_xa_hard_limit_for_highT ¶

Use min_xa_hard_limit_for_highT for center logT >= this. Linear interpolate in logT for intermediate center temperatures.

```
logT_min_for_min_xa_hard_limit_for_highT = 9.51d0
```

sum_xa_hard_limit ¶

sum_xa_hard_limit_for_highT ¶

If solver produces any cell with $\text{abs}(\text{sum}(\text{xa})-1) >$ this limit, then reject the trial solution. Can optionally relax this limit at high T.

```
sum_xa_hard_limit = 5d-4
sum_xa_hard_limit_for_highT = 1d-3
```

logT_max_for_sum_xa_hard_limit ¶

Use sum_xa_hard_limit for center logT <= this.

```
logT_max_for_sum_xa_hard_limit = 9.49d0
```

logT_min_for_sum_xa_hard_limit_for_highT ¶

Use `sum_xa_hard_limit_for_highT` for center `logT` \geq this. Linear interpolate in `logT` for intermediate center temperatures.

```
logT_min_for_sum_xa_hard_limit_for_highT = 9.51d0
```

do_newton_damping_for_neg_xa ¶

If true, uniformly reduce newton corrections if necessary to avoid neg abundances.

```
do_newton_damping_for_neg_xa = .true.
```

min_logT_for_quad ¶

```
min_logT_for_quad = 99
```

min_chem_eqn_scale ¶

```
min_chem_eqn_scale = 1d0
```

hydro_mtx_max_allowed_{abs}{dlogT | dlogRho | dlogPgas | logT | logRho | logPgas} ¶

Force retry with smaller timestep if hydro solves change T, Rho, or Pgas by too much or make them too large.

```
hydro_mtx_max_allowed_abs_dlogT = 99d0
hydro_mtx_max_allowed_abs_dlogE = 99d0
hydro_mtx_max_allowed_abs_dlogRho = 99d0
hydro_mtx_max_allowed_abs_dlogPgas = 99d0
min_logT_for_hydro_mtx_max_allowed = -1d99
hydro_mtx_max_allowed_logT = 99d0
hydro_mtx_max_allowed_logE = 99d0
hydro_mtx_max_allowed_logRho = 99d0
hydro_mtx_max_allowed_logPgas = 99d0
```

solver_clip_dlogT ¶

solver_clip_dlogRho ¶

solver_clip_dlogPgas ¶

solver_clip_dlogR ¶

Limit magnitude of relative changes per iteration in solver. Ignore if limit is ≤ 0 .

```
solver_clip_dlogT = -1d0
solver_clip_dlogRho = -1d0
```

```
solver_clip_dlogPgas = -1d0  
solver_clip_dlogR = -1d0
```

use_time_centering ¶

if true, then use time centered velocity and time weighted area.

```
use_time_centering = .false.
```

gold tolerances for newton solver

use_gold_tolerances ¶

gold_newton_iterations_limit ¶

gold_newton_iterations_hard_limit ¶

maxT_for_gold_tolerances ¶

max_eosPC_frac_for_gold_tolerances ¶

gold_tol_residual_norm1 ¶

gold_iter_for_resid_tol2 ¶

gold_tol_residual_norm2 ¶

gold_tol_max_residual2 ¶

gold_iter_for_resid_tol3 ¶

gold_tol_residual_norm3 ¶

gold_tol_max_residual3 ¶

```
use_gold_tolerances = .false.
```

```
maxT_for_gold_tolerances = 7e8  
max_eosPC_frac_for_gold_tolerances = 1d-9  
gold_tol_residual_norm1 = 1d-11  
gold_tol_max_residual1 = 1d-9  
gold_iter_for_resid_tol2 = 9  
gold_tol_residual_norm2 = 1d-8  
gold_tol_max_residual2 = 1d-6  
gold_iter_for_resid_tol3 = 14  
gold_tol_residual_norm3 = 1d-6  
gold_tol_max_residual3 = 1d-4
```

```
gold_newton_iterations_limit = 14
gold_newton_iterations_hard_limit = -1
```

artificial viscosity

eps_visc_factor ¶

multiply the eps_visc term in energy equation by this factor.

```
eps_visc_factor = 1d0
```

dvdt_visc_factor ¶

multiply the dvdt_visc term in momentum equation by this factor.

```
dvdt_visc_factor = 1d0
```

use_artificial_viscosity ¶

artificial viscosity – only applies when using velocity variables

```
use_artificial_viscosity = .false.
```

artificial_viscosity_Q_shift ¶

```
Qvisc = min(0d0, Qvisc + artificial_viscosity_Q_shift)
```

This serves to filter out use of artificial viscosity in low compression regions. e.g., try
`artificial_viscosity_Q_shift = 1d34`.

```
artificial_viscosity_Q_shift = -1
```

post_shock_viscosity_decay_factor ¶

Exponential decrease in artificial viscosity inward from Mach 1 location. The value of Qvisc is multiplied by $\exp(-\text{dist_to_Mach1}/Hq)$ where Hq = this factor times the local radius and `dist_to_Mach1` is the distance at the start of the current step outward to the nearest Mach 1 location if using both pre and post shock decay, use the closer Mach1

```
post_shock_viscosity_decay_factor = -1
```

pre_shock_viscosity_decay_factor ¶

Exponential decrease in artificial viscosity outward from Mach 1 location. The value of `Qvisc` is multiplied by $\exp(-\text{dist_to_Mach1}/Hq)$ where `Hq` = this factor times the local radius and `dist_to_Mach1` is the distance at the start of the current step inward to the nearest Mach 1 location if using both pre and post shock decay, use the closer Mach1

```
pre_shock_viscosity_decay_factor = -1
```

shock_spread_quadratic ¶

the artificial viscosity coefficient includes a quadratic term that is proportional to $(\text{shock_spread_quadratic} * r)^2$ where `r` is the local radius.

```
shock_spread_quadratic = 1d-3
```

shock_spread_linear ¶

the artificial viscosity coefficient includes a linear term that is proportional to $(\text{shock_spread_linear} * r * cs)$ where `cs` is the local sound speed and `r` is the local radius.

```
shock_spread_linear = 0
```

art_visc_full_on_logRho_ge_this ¶

art_visc_full_off_logRho_le_this ¶

```
art_visc_full_on_logRho_ge_this = -99
art_visc_full_off_logRho_le_this = -99
```

split mixing ¶

split_mixing_choice ¶

- 0 = no split = mixing coupled to burn and structure.
- -1 = mix for full dt before burn+struct
- -2 = mix for full dt after burn+struct
- -3 = mix for dt/2 before and dt/2 after

```
split_mixing_choice = 0
```

reset_mixing_info_before_final_mix ¶

Relevant for `split_mixing_choice` options -2 and -3.

```
reset_mixing_info_before_final_mix = .true.
```

op_split_mix_atol ¶

op_split_mix_rtol ¶

Absolute and relative error tolerances.

```
op_split_mix_atol = 1d-5
op_split_mix_rtol = 1d-6
```

controls related to split mixing for timesteps

max_fixup_for_mix_limit ¶

If split_mix_fixup is bigger than this, reduce the next timestep.

```
max_fixup_for_mix_limit = 1d-3
```

max_fixup_for_mix_hard_limit ¶

If split_mix_fixup is bigger than this, retry.

```
max_fixup_for_mix_hard_limit = 1d99
```

op_split_mix_trace ¶

```
op_split_mix_trace = .false.
```

op_split_burn ¶

```
op_split_burn = .false.
```

timestep controls ¶

The terminal output during evolution includes a short string for the `dt_limit`. This is to give you some indication of what is limiting the time steps. Here's a dictionary mapping those terminal strings to the corresponding control parameters. (There is a similar table in `mesa/binary/defaults/binary_controls.defaults`.)

terminal output	related parameter
'avg lgE resid'	limit_for_avg_lgE_residual
'CpT_absMdot_div_L'	CpT_absMdot_div_L_limit
'Lnuc'	delta_lgL_nuc_limit
'Lnuc_cat'	delta_lgL_nuc_cat_limit
'Lnuc_H'	delta_lgL_H_limit

'Lnuc_He'	delta_lgL_He_limit
'Lnuc_photo'	delta_lgL_photo_limit
'Lnuc_z'	delta_lgL_z_limit
'bad_X_sum'	(solver found bad mass sum)
'dH'	dH_limit
'dH/H'	dH_div_H_limit
'dHe'	dHe_limit
'dHe/He'	dHe_div_He_limit
'dHe3'	dHe3_limit
'dHe3/He3'	dHe3_div_He3_limit
'dL/L'	dL_div_L_limit
'dX'	dX_limit
'dX/X'	dX_div_X_limit
'dX_nuc_drop'	dX_nuc_drop_limit
'd_delR_grow'	d_deltaR_grow_limit
'd_delR_shrink'	d_deltaR_shrink_limit
'delta Ye'	delta_Ye_limit
'delta mdot'	delta_mdot_limit
'delta total J'	delta_lg_total_J_limit
'delta_HR'	delta_HR_limit
'delta_mstar'	delta_lg_star_mass_limit
'diff iters'	diffusion_iters_limit
'diff steps'	diffusion_steps_limit
'min_dr_div_cs'	dt_div_min_dr_div_cs_limit
'dt_acoustic'	dt_div_dt_acoustic_limit
'dt_collapse'	dt_div_dt_cell_collapse_limit
'dt_dynamic'	dt_div_dt_dynamic_limit
'dt_mass_loss'	dt_div_dt_mass_loss_limit
'dt_thermal'	dt_div_dt_thermal_limit
'eps_nuc_cntr'	delta_log_eps_nuc_cntr_limit
'error enrg'	limit_for_rel_error_in_energy_conservation
'error rate'	limit_for_log_rel_rate_in_energy_conservation
'highT del Ye'	delta_Ye_highT_limit
'hold'	(recent backup, so no increase in dt)
'lgL'	delta_lgL_limit
'lgL_phot'	delta_lgL_phot_limit
'lgP'	delta_lgP_limit
'lgP_cntr'	delta_lgP_cntr_limit
'lgR'	delta_lgR_limit
'lgRho'	delta_lgRho_limit
'lgRho_cntr'	delta_lgRho_cntr_limit
'lgRho_max'	delta_lgRho_max_limit
'lgT'	delta_lgT_limit
'lgT_cntr'	delta_lgT_cntr_limit
'lgT_max'	delta_lgT_max_limit
'lgTeff'	delta_lgTeff_limit
'lg_XC_cntr'	delta_lg_XC_cntr_limit
'lg_XH_cntr'	delta_lg_XH_cntr_limit
'lg_XHe_cntr'	delta_lg_XHe_cntr_limit
'lg_XNe_cntr'	delta_lg_XNe_cntr_limit
'lg_XO_cntr'	delta_lg_XO_cntr_limit
'lg_XSi_cntr'	delta_lg_XSi_cntr_limit
'log_eps_nuc'	delta_log_eps_nuc_limit
'max E resid'	limit_for_max_E_residual
'max lgE resid'	limit_for_max_abs_lgE_residual
'max_dt'	max_years_for_timestep
'max dt change'	max_timestep_factor
'min dt change'	min_timestep_factor

```
'neg_mass_frac'      (solver found neg mass frac)
'newton_iters'        newton_iterations_limit
'rotation_steps'      rotation_steps_limit
'v/v_crit'            v_div_v_crit_limit
'varcontrol'          varcontrol_target
'b_****'              see binary/defaults/binary_controls.default
```

max_timestep ¶

In seconds. max_timestep <= 0 means no upper limit.

```
max_timestep = 0
```

max_years_for_timestep ¶

max_years_for_timestep <= 0 means no upper limit. Note: max_timestep is the control that is used by most of the code. max_years_for_timestep is just provided as a convenience. At the start of each step, the evolve routine checks to see if max_years_for_timestep > 0, and if so, it sets max_timestep = max_years_for_timestep*secyer.

```
max_years_for_timestep = 0
```

max_timestep_hi_T_limit ¶

If max T >= this, then switch to hi_T_max_years_for_timestep. Ignore if <= 0.

```
max_timestep_hi_T_limit = -1
```

hi_T_max_years_for_timestep ¶

Max years for timestep if max_timestep_hi_T_limit is active.

```
hi_T_max_years_for_timestep = 0
```

min_timestep_factor ¶

Lower limit for ratio of new timestep to previous timestep. i.e., allow dt to get smaller by no more than this factor – 0 means no limit.

```
min_timestep_factor = 0.8d0
```

force_timestep_min ¶

In seconds. force_timestep_min <= 0 means no forced lower limit.

```
force_timestep_min = 0
```

force_timestep_min_years

`force_timestep_min_years <= 0` means no forced lower limit. Note: `force_timestep_min` is the control that is used by most of the code. `force_timestep_min_years` is just provided as a convenience. At the start of each step, the `evolve` routine checks if `force_timestep_min_years > 0`, and if so, it sets `force_timestep_min = force_timestep_min_years*secyer`.

```
force_timestep_min_years = 0
```

force_timestep_min_factor

If `dt` is `< force_timestep_min`, then replace `dt` by `min(dt*force_timestep_min_factor, force_timestep_min)`

```
force_timestep_min_factor = 2d0
```

max_timestep_factor

Upper limit for ratio of new timestep to previous timestep. i.e., allow `dt` to get larger by no more than this factor – 0 means no limit.

```
max_timestep_factor = 1.2d0
```

timestep_factor_for_retries

Before retry, decrease `dt` by this.

```
timestep_factor_for_retries = 0.5d0
```

timestep_factor_for_backups

Before backup, decrease `dt` by this (or more if multiple backups in a row).

```
timestep_factor_for_backups = 0.5d0
```

backup_hold

No increases in timestep for `backup_hold` steps after a backup.

```
backup_hold = 2
```

retry_hold ¶

No increases in timestep for `retry_hold` steps after a retry.

```
retry_hold = 1
```

neg_mass_fraction_hold ¶

No increases in timestep for `neg_mass_fraction_hold` steps after a retry or backup caused by a negative mass fraction.

```
neg_mass_fraction_hold = 2
```

timestep_dt_factor = 0.9 ¶

dt reduction factor exceed timestep limits.

```
timestep_dt_factor = 0.9d0
```

dt_limit_ratio_target ¶

Aim for this ratio on dt limited timesteps.

```
dt_limit_ratio_target = 1d0
```

use_dt_low_pass_controller ¶

Enable low pass filter for smoother timestep variations.

```
use_dt_low_pass_controller = .true.
```

varcontrol_target ¶

This is the target value for relative variation in the structure from one model to the next. The default timestep adjustment is to increase or reduce the timestep depending on whether the actual variation was smaller or greater than this value.

```
varcontrol_target = 1d-4
```

varcontrol_dt_limit_ratio_hard_max ¶

`varcontrol_dt_limit_ratio` is the actual `varcontrol` value divided by the target. if that ratio exceeds this limit, then retry with a smaller timestep. this let's you prevent large changes from happening in a single step.

```
varcontrol_dt_limit_ratio_hard_max = 1d99
```

relax_hard_limits_after_backup ¶

If true, then don't enforce hard limits immediately after a backup.

```
relax_hard_limits_after_backup = .true.
```

relax_hard_limits_after_retry ¶

If true, then don't enforce hard limits immediately after a retry.

```
relax_hard_limits_after_retry = .true.
```

limits based on iterations required by various solvers

newton_iterations_limit ¶

If newton solve uses more `newton_iterations` than this, reduce the next timestep.

```
newton_iterations_limit = 7
```

newton_iterations_hard_limit ¶

If uses more iterations than this, retry.

```
newton_iterations_hard_limit = -1
```

rotation_steps_limit ¶

If rotation solver uses more steps than this, reduce the next timestep.

```
rotation_steps_limit = 500
```

rotation_steps_hard_limit ¶

If rotation solver uses more steps than this, retry.

```
rotation_steps_hard_limit = 700
```

diffusion_steps_limit ¶

If diffusion solver uses more steps than this, reduce the next timestep.

```
diffusion_steps_limit = 500
```

diffusion_steps_hard_limit ¶

If diffusion solver uses more steps than this, retry.

```
diffusion_steps_hard_limit = 700
```

diffusion_iters_limit ¶

If use a total number of iters > this, reduce the next timestep.

```
diffusion_iters_limit = 600
```

diffusion_iters_hard_limit ¶

If use a total number of iters > this, retry.

```
diffusion_iters_hard_limit = 800
```

limits based on max decrease in mass fraction at any location in star

dX_mix_dist_limit ¶

Option to ignore decreases in abundance in non-mixed cells near mixing boundaries. Ignore abundance changes if nearest mixing boundary is closer than this in Msun units. This applies to dH, dH_div_H, dHe, dHe_d_He, dX, and dX_div_X limits.

```
dX_mix_dist_limit = 1d-4
```

Limit on magnitude of decrease in any cell hydrogen abundance during a single timestep. dH here is $\text{abs}(x_a(h1, k) - x_{a_old}(h1, k))$ for any cell k. Considers all cells except where have convective mixing.

dH_limit_min_H ¶

dH limits only apply where $x_a(h1, k) \geq$ this limit.

```
dH_limit_min_H = 1d99
```

dH_limit ¶

If max dH is greater than this, reduce the next timestep by $\text{dH_limit}/\text{max_dH}$.

```
dH_limit = 1d99
```

dH_hard_limit

If max dH is greater than this, retry with smaller timestep.

```
dH_hard_limit = 1d99
```

dH_decreases_only

If true, then only consider decreases in abundance.

```
dH_decreases_only = .true.
```

Limit on magnitude of relative decrease in any cell hydrogen abundance. dH_div_H here is $\text{abs}(x_a(h1, k) - x_{a_old}(h1, k)) / x_a(h1, k)$ considers all cells except where have convective mixing. dH_decreases_only applies to dH_div_H also.

dH_div_H_limit_min_H

dH_div_H limits only apply where $x_a(h1, k) \geq$ this limit.

```
dH_div_H_limit_min_H = 1d-3
```

dH_div_H_limit

If max dH_div_H is greater than this, reduce the next timestep by dH_limit/max_dH.

```
dH_div_H_limit = 0.5d0
```

dH_div_H_hard_limit

If max dH_div_H is greater than this, retry with smaller timestep.

```
dH_div_H_hard_limit = 1d99
```

Limit on magnitude of decrease in any cell helium abundance during a single timestep. dHe here is $\text{abs}(x_a(\text{he4}, k) - x_{a_old}(\text{he4}, k))$ for any cell k. Considers all cells except where have convective mixing.

dHe_limit_min_He

dHe limits only apply where $x_a(\text{he4}, k) \geq$ this limit.

```
dHe_limit_min_He = 1d99
```

dHe_limit ¶

If max dHe is greater than this, reduce the next timestep by $\text{dHe_limit}/\text{max_dHe}$.

```
dHe_limit = 1d99
```

dHe_hard_limit ¶

If max dHe is greater than this, retry with smaller timestep.

```
dHe_hard_limit = 1d99
```

dHe_decreases_only ¶

If true, then only consider decreases in abundance. `dHe_decreases_only` applies to `dHe_div_He` also.

```
dHe_decreases_only = .true.
```

Limit on magnitude of relative decrease in any cell helium abundance. `dHe_div_He` here is $\text{abs}(\text{xa}(\text{he4}, k) - \text{xa_old}(\text{he4}, k)) / \text{xa}(\text{he4}, k)$. Considers all cells except where have convective mixing.

dHe_div_He_limit_min_He ¶

`dHe_div_He` limits only apply where $\text{xa}(\text{he4}, k) \geq$ this limit.

```
dHe_div_He_limit_min_He = 1d-3
```

dHe_div_He_limit ¶

If max `dHe_div_He` is greater than this, reduce the next timestep by $\text{dHe_limit}/\text{max_dHe}$.

```
dHe_div_He_limit = 0.5d0
```

dHe_div_He_hard_limit ¶

If max `dHe_div_He` is greater than this, retry with smaller timestep.

```
dHe_div_He_hard_limit = 1d99
```

Limit on magnitude of decrease in any cell helium abundance during a single timestep. `dHe3` here is $\text{abs}(\text{xa}(\text{he4}, k) - \text{xa_old}(\text{he3}, k))$ for any cell k . Considers all cells except where have convective mixing.

dHe3_limit_min_He3 ¶

dHe3 limits only apply where $xa(\text{he3}, k) \geq$ this limit.

```
dHe3_limit_min_He3 = 1d99
```

dHe3_limit ¶

If max dHe3 is greater than this, reduce the next timestep by $dHe3_limit / \max_dHe3$.

```
dHe3_limit = 1d99
```

dHe3_hard_limit ¶

If max dHe3 is greater than this, retry with smaller timestep.

```
dHe3_hard_limit = 1d99
```

dHe3_decreases_only ¶

If true, then only consider decreases in abundance. dHe3_decreases_only applies to dHe3_div_He3 also.

```
dHe3_decreases_only = .true.
```

Limit on magnitude of relative decrease in any cell helium abundance. dHe3_div_He3 here is $\text{abs}(xa(\text{he3}, k) - xa_old(\text{he3}, k)) / xa(\text{he3}, k)$. Considers all cells except where have convective mixing.

dHe3_div_He3_limit_min_He3 ¶

dHe3_div_He3 limits only apply where $xa(\text{he3}, k) \geq$ this limit.

```
dHe3_div_He3_limit_min_He3 = 1d99
```

dHe3_div_He3_limit ¶

if max dHe3_div_He3 is greater than this, reduce the next timestep by $dHe3_limit / \max_dHe3$.

```
dHe3_div_He3_limit = 1d99
```

dHe3_div_He3_hard_limit ¶

If max dHe3_div_He3 is greater than this, retry with smaller timestep.

```
dHe3_div_He3_hard_limit = 1d99
```

Limit on magnitude of decrease in any cell nonH, nonHe abundance. dX here is $\text{abs}(x_a(j, k) - x_{a_old}(j, k))$ for any cell k and any species j other except hydrogen or helium. Considers all cells except where have convective mixing.

`dX_limit_min_X`

dX limits only apply where $x_a(j, k) \geq$ this limit.

```
dX_limit_min_X = 1d99
```

`dX_limit`

If max dX is greater than this, reduce the next timestep by $dX_limit/\text{max_dX}$.

```
dX_limit = 1d99
```

`dX_hard_limit`

If max dX is greater than this, retry with smaller timestep.

```
dX_hard_limit = 1d99
```

`dX_decreases_only`

If true, then only consider decreases in abundance. `dX_decreases_only` applies to `dX_div_X` also.

```
dX_decreases_only = .true.
```

Limit on magnitude of relative decrease in any cell nonH, nonHe abundance. dX_div_X here is $\text{abs}(x_a(j, k) - x_{a_old}(j, k))/x_a(j, k)$ for any cell k and any species j other except hydrogen or helium. Considers all cells except where have convective mixing.

`dX_div_X_limit_min_X`

dX_div_X limits only apply where $x_a(j, k) \geq$ this limit.

```
dX_div_X_limit_min_X = 1d99
```

`dX_div_X_limit`

If max dX_div_X is greater than this, reduce the next timestep by $dX_limit/\text{max_dX}$.

```
dX_div_X_limit = 1d99
```

dX_div_X_hard_limit ¶

If max dX_div_X is greater than this, retry with smaller timestep.

```
dX_div_X_hard_limit = 1d99
```

Limits on max drop in abundance mass fraction from burning with possible mixing inflow. This considers both nuclear reactions and offsetting effect of mixing inflow.

dX_nuc_drop_min_X_limit ¶

dX_nuc_drop_limit only for X > dX_nuc_drop_min_X_limit.

```
dX_nuc_drop_min_X_limit = 1d-4
```

dX_nuc_drop_max_A_limit ¶

dX_nuc_drop_limit only for species with A <= dX_nuc_drop_max_A_limit.

```
dX_nuc_drop_max_A_limit = 52
```

dX_nuc_drop_limit_at_high_T ¶

Negative means use value for dX_nuc_drop_limit, else use this limit when center logT > 9.45.

```
dX_nuc_drop_limit_at_high_T = -1
```

dX_nuc_drop_limit ¶

If max dX_nuc_drop is greater than dX_nuc_drop_limit, reduce the next timestep by dX_nuc_drop_limit/max_dX_nuc_drop.

```
dX_nuc_drop_limit = 5d-2
```

dX_nuc_drop_hard_limit ¶

If max dX_nuc_drop is greater than dX_nuc_drop_hard_limit, retry with smaller timestep.

```
dX_nuc_drop_hard_limit = 1d99
```

dX_nuc_drop_min_yrs_for_dt ¶

Don't let dX_nuc_drop change dt to smaller than this.

```
dX_nuc_drop_min_yrs_for_dt = 1d-9
```

limits based on relative changes in variables L, P, Rho, T, R, eps_nuc ¶

limit on magnitude of relative change in L at any grid point

```
dL_div_L = abs(L(k) - L_old(k))/L(k)
```

dL_div_L_limit ¶

If max abs dL_div_L is greater than this, reduce the next timestep.

```
dL_div_L_limit = -1
```

dL_div_L_hard_limit ¶

If max abs dL_div_L is greater than this, retry with smaller timestep.

```
dL_div_L_hard_limit = -1
```

dL_div_L_limit_min_L ¶

In Lsun units. dL_div_L limits only apply where $L(k) \geq L_{\text{sun}} \cdot \text{dL_limit_min_L}$

```
dL_div_L_limit_min_L = 1d99
```

delta_lgP_limit ¶

Limit for magnitude of max change in log10 total pressure in any cell.

```
delta_lgP_limit = 1
```

delta_lgP_hard_limit ¶

If max delta_lgP is greater than delta_lgP_hard_limit, retry with smaller timestep.

```
delta_lgP_hard_limit = -1
```

delta_lgP_limit_min_lgP ¶

`delta_lgP_limit` limits only apply where $\log_{10}P(k) \geq \text{delta_lgP_limit_min_lgP}$

```
delta_lgP_limit_min_lgP = 1d99
```

delta_lgRho_limit ¶

Limit for magnitude of max change in \log_{10} density in any cell.

```
delta_lgRho_limit = 1
```

delta_lgRho_hard_limit = -1 ¶

If max `delta_lgRho` is greater than `delta_lgRho_hard_limit`, retry with smaller timestep.

```
delta_lgRho_hard_limit = -1
```

delta_lgRho_limit_min_lgRho ¶

`delta_lgRho_limit` limits only apply where $\log_{10}Rho(k) \geq \text{delta_lgRho_limit_min_lgRho}$.

```
delta_lgRho_limit_min_lgRho = 1d99
```

delta_lgT_limit ¶

Limit for magnitude of max change in \log_{10} temperature in any cell.

```
delta_lgT_limit = 0.5d0
```

delta_lgT_hard_limit ¶

If max `delta_lgT` is greater than `delta_lgT_hard_limit`, retry with smaller timestep.

```
delta_lgT_hard_limit = -1
```

delta_lgT_limit_min_lgT ¶

`delta_lgT_limit` limits only apply where $\log_{10}T(k) \geq \text{delta_lgT_limit_min_lgT}$.

```
delta_lgT_limit_min_lgT = 1d99
```

delta_lgE_limit ¶

Limit for magnitude of max change in log10 internal energy in any cell.

```
delta_lgE_limit = 0.1d0
```

delta_lgE_hard_limit ¶

If max delta_lgE is greater than delta_lgE_hard_limit, retry with smaller timestep.

```
delta_lgE_hard_limit = -1
```

delta_lgE_limit_min_lgE ¶

delta_lgE_limit limits only apply where $\log_{10}(E(k)) \geq \text{delta_lgE_limit_min_lgE}$.

```
delta_lgE_limit_min_lgE = 1d99
```

delta_lgR_limit ¶

Limit for magnitude of max change in log10 radius at any cell boundary.

```
delta_lgR_limit = 0.5d0
```

delta_lgR_hard_limit ¶

If max delta_lgR is greater than delta_lgR_hard_limit, retry with smaller timestep.

```
delta_lgR_hard_limit = -1
```

delta_lgR_limit_min_lgR ¶

delta_lgR_limit limits only apply where $\log_{10}R(k) \geq \text{delta_lgR_limit_min_lgR}$.

```
delta_lgR_limit_min_lgR = 1d99
```

delta_Ye_limit ¶

Limit for magnitude of max change in Ye in any cell.

```
delta_Ye_limit = 1
```

delta_Ye_hard_limit ¶

If max delta_Ye is greater than delta_Ye_hard_limit, retry with smaller timestep.

```
delta_Ye_hard_limit = -1
```

delta_Ye_highT_limit ¶

Limit for magnitude of max change in Ye in high T cells.

```
delta_Ye_highT_limit = 99
```

Limit testing for max delta_ye to cells with $T \geq \text{minT_for_highT_Ye_limit}$. If this high T max delta_Ye is greater than delta_Ye_highT_limit, reduce the next timestep by delta_Ye_highT_limit/max_delta_Ye.

```
delta_Ye_highT_hard_limit = -1
```

minT_for_highT_Ye_limit ¶

Limit testing for max delta_ye to cells with $T \geq \text{minT_for_highT_Ye_limit}$. If this high T max delta_Ye is greater than delta_Ye_highT_limit, retry with smaller timestep.

```
minT_for_highT_Ye_limit = 7d9
```

delta_log_eps_nuc_limit ¶

Limit for magnitude of max change in log10 eps_nuc in any cell. Only applies to increases in non-convective zones.

```
delta_log_eps_nuc_limit = -1
```

delta_log_eps_nuc_hard_limit ¶

If max delta_log_eps_nuc is greater than delta_log_eps_nuc_hard_limit, retry with smaller timestep.

```
delta_log_eps_nuc_hard_limit = -1
```

d_deltaR_shrink_limit ¶

Limit for relative decrease in radial thickness of any zone.

```
d_deltaR_shrink_limit = -1
```

d_deltaR_shrink_hard_limit ¶

If max `d_deltaR_shrink` is greater than `d_deltaR_shrink_hard_limit`, retry with smaller timestep.

```
d_deltaR_shrink_hard_limit = -1
```

d_deltaR_grow_limit ¶

Limit for relative increase in radial thickness of any zone.

```
d_deltaR_grow_limit = -1
```

d_deltaR_grow_hard_limit ¶

If max `d_deltaR_grow` is greater than `d_deltaR_grow_hard_limit`, retry with smaller timestep.

```
d_deltaR_grow_hard_limit = -1
```

limits based on integrated power at each point for each category of nuclear reaction ¶

`lgL_nuc_cat` = nuclear reaction energy release for a particular category of reaction (Lsun units). Energy release here excludes neutrinos.

delta_lgL_nuc_cat_limit ¶

Limit for magnitude of change in `lgL_nuc` for category.

```
delta_lgL_nuc_cat_limit = -1
```

delta_lgL_nuc_cat_hard_limit ¶

If max `delta` is greater than `delta_lgL_nuc_cat_hard_limit`, retry with smaller timestep.

```
delta_lgL_nuc_cat_hard_limit = -1
```

lgL_nuc_cat_burn_min ¶

Ignore changes in `lgL_nuc` for category if value is less than this.

```
lgL_nuc_cat_burn_min = -1
```

lgL_nuc_mix_dist_limit ¶

Ignore if nearest boundary is closer than this. Ignore changes in `lgL` in cells near mixing boundaries.


```
lgL_nuc_mix_dist_limit = 1d-6
```

check_delta_lgL_{burning_category} ¶

Flags determining which reaction categories are considered.

```
check_delta_lgL_pp = .true.
check_delta_lgL_cno = .true.
check_delta_lgL_3alf = .true.
check_delta_lgL_burn_c = .true.
check_delta_lgL_burn_n = .true.
check_delta_lgL_burn_o = .true.
check_delta_lgL_burn_ne = .true.
check_delta_lgL_burn_na = .true.
check_delta_lgL_burn_mg = .true.
check_delta_lgL_burn_si = .true.
check_delta_lgL_burn_s = .true.
check_delta_lgL_burn_ar = .true.
check_delta_lgL_burn_ca = .true.
check_delta_lgL_burn_ti = .true.
check_delta_lgL_burn_cr = .true.
check_delta_lgL_burn_fe = .true.
```

c12 + c12, c12 + o16, and o16 + o16

```
check_delta_lgL_cc = .true.
check_delta_lgL_co = .true.
check_delta_lgL_oo = .true.
```

L_H_burn = integrated power at surface from PP and CNO (in Lsun units)

values for lgL_H are $\log_{10}(\max(1, L_H_burn))$

delta_lgL_H_limit ¶

limit for magnitude of change in lgL_H

```
delta_lgL_H_limit = -1
```

delta_lgL_H_hard_limit ¶

if max delta is greater than delta_lgL_H_hard_limit, retry with smaller timestep

```
delta_lgL_H_hard_limit = -1
```

lgL_H_burn_min ¶

ignore changes in lgL_H if value is less than this

```
lgL_H_burn_min = 1.5d0
```

lgL_H_drop_factor ¶

when L_H is dropping, multiply limits by this factor

```
lgL_H_drop_factor = 1
```

lgL_H_burn_relative_limit ¶

ignore changes in lgL_H if $\max(\text{lgL_He}, \text{lgL_z}) - \text{lgL_H} > \text{this}$

```
lgL_H_burn_relative_limit = 3
```

L_He_burn = integrated power at surface from triple alpha (in Lsun units)

values for lgL_He are $\log_{10}(\max(1, L_{\text{He_burn}}))$

delta_lgL_He_limit ¶

Limit for magnitude of change in lgL_He.

```
delta_lgL_He_limit = 0.025d0
```

delta_lgL_He_hard_limit ¶

If max delta is greater than delta_lgL_He_hard_limit, retry with smaller timestep.

```
delta_lgL_He_hard_limit = -1
```

lgL_He_burn_min ¶

Ignore changes in lgL_He if value is less than this.

```
lgL_He_burn_min = 2.5d0
```

lgL_He_drop_factor ¶

When L_He is dropping, multiply limits by this factor.

```
lgL_He_drop_factor = 1
```

lgL_He_burn_relative_limit ¶

Ignore changes in `lgL_He` if $\max(\text{lgL}_H, \text{lgL}_Z) - \text{lgL}_{He} > \text{this}$.

```
lgL_He_burn_relative_limit = 3
```

`L_z_burn` = integrated power at surface from nuclear burning other than H, He, or C (in `Lsun` units) excluding photodistintegrations

values for `lgL_z` are $\log_{10}(\max(1, L_z_{\text{burn}}))$

delta_lgL_z_limit ¶

Limit for magnitude of change in `lgL_z`.

```
delta_lgL_z_limit = -1
```

delta_lgL_z_hard_limit ¶

If max delta is greater than `delta_lgL_z_hard_limit`, retry with smaller timestep.

```
delta_lgL_z_hard_limit = -1
```

lgL_z_burn_min ¶

Ignore changes in `lgL_z` if value is less than this.

```
lgL_z_burn_min = 2.5d0
```

lgL_z_drop_factor ¶

When `L_z` is dropping, multiply limits by this factor.

```
lgL_z_drop_factor = 1
```

lgL_z_burn_relative_limit ¶

Ignore changes in `lgL_z` if $\max(\text{lgL}_H, \text{lgL}_{He}) - \text{lgL}_Z > \text{this}$.

```
lgL_z_burn_relative_limit = 3
```

`L_photo_burn` = magnitude of integrated power at surface from photodistintegrations

values for `lgL_photo` are based on `L_by_category(iphoto)`

delta_lgL_photo_limit ¶

Limit for magnitude of change in `lgL_photo`.

```
delta_lgL_photo_limit = -1
```

delta_lgL_photo_hard_limit ¶

If max delta is greater than `delta_lgL_photo_hard_limit`, retry with smaller timestep.

```
delta_lgL_photo_hard_limit = -1
```

lgL_photo_burn_min ¶

Ignore changes in `lgL_photo` if value is less than this.

```
lgL_photo_burn_min = 2.5d0
```

lgL_photo_drop_factor ¶

When `L_photo` is dropping, multiply limits by this factor.

```
lgL_photo_drop_factor = 1
```

limits based on total integrated power at surface for all nuclear reactions ¶

excluding photodistintegrations

`L_nuc` = nuclear reaction total energy release for all nuclear reactions (Lsun units)

delta_lgL_nuc_limit ¶

limit for magnitude of change in `lgL_nuc`

```
delta_lgL_nuc_limit = -1
```

delta_lgL_nuc_hard_limit ¶

if max delta is greater than `delta_lgL_nuc_hard_limit`, retry with smaller timestep

```
delta_lgL_nuc_hard_limit = -1
```

lgL_nuc_burn_min ¶

ignore changes in `lgL_nuc` if value is less than this

```
lgL_nuc_burn_min = 0.5d0
```

lgL_nuc_drop_factor ¶

When L_nuc is dropping, multiply limits by this factor.

```
lgL_nuc_drop_factor = 1
```

limits based on changes at photosphere

delta_lgTeff_limit ¶

delta_lgTeff_hard_limit ¶

Limit for magnitude of max change in log10 temperature at photosphere.

```
delta_lgTeff_limit = 0.01d0  
delta_lgTeff_hard_limit = -1
```

delta_lgL_limit_L_min ¶

delta_lgL_limit ¶

delta_lgL_hard_limit ¶

Limit for magnitude of change in log10(L/Lsun). Only apply this limit when L >= delta_lgL_limit_L_min (in Lsun units).

```
delta_lgL_limit_L_min = -100  
delta_lgL_limit = 0.1d0  
delta_lgL_hard_limit = -1
```

delta_lgL_phot_limit_L_min ¶

delta_lgL_phot_limit ¶

delta_lgL_phot_hard_limit ¶

Limit for magnitude of change in log10(L_phot/Lsun). Only apply this limit when L_phot >= delta_lgL_phot_limit_L_min (in Lsun units).

```
delta_lgL_phot_limit_L_min = -100  
delta_lgL_phot_limit = 0.1d0  
delta_lgL_phot_hard_limit = -1
```

v_div_v_crit_limit ¶

v_div_v_crit_hard_limit ¶

Limit surface rotational velocity div critical velocity (v_div_v_crit_avg_surf).

```
v_div_v_crit_limit = -1
v_div_v_crit_hard_limit = -1
```

dt_div_dt_thermal_limit ¶**dt_div_dt_thermal_hard_limit ¶**

limit for dt compared to thermal timescale (negative means no limit)

```
dt_thermal = (3/4)*G*M^2/(R*L); Kelvin-Helmholtz time
```

```
dt_div_dt_thermal_limit = -1
dt_div_dt_thermal_hard_limit = -1
```

dt_div_dt_dynamic_limit ¶**dt_div_dt_dynamic_hard_limit ¶**

limit for dt compared to dynamic timescale (negative means no limit)

```
dt_dynamic = 2*Pi*sqrt(R^3/(G*M))
```

```
dt_div_dt_dynamic_limit = -1
dt_div_dt_dynamic_hard_limit = -1
```

dt_div_dt_acoustic_limit ¶**dt_div_dt_acoustic_hard_limit ¶**

limit for dt compared to dt_acoustic (negative means no limit)

dt_acoustic = time for sound from center to photosphere = sum over shells of local sound crossing time dr/csound.

```
dt_div_dt_acoustic_limit = -1
dt_div_dt_acoustic_hard_limit = -1
```

dt_div_dt_mass_loss_limit ¶

dt_div_dt_mass_loss_hard_limit ¶

limit for dt compared to mass loss timescale (negative means no limit)

```
dt_mass_loss = -M/Mdot; only applies when Mdot < 0
```

```
dt_div_dt_mass_loss_limit = -1
dt_div_dt_mass_loss_hard_limit = -1
```

dt_div_min_dr_div_cs_limit ¶**dt_div_min_dr_div_cs_hard_limit ¶**

limit for dt compared to explicit solver timescale (negative means no limit)

```
min_dr_div_cs = min over all cells of dr/csound (seconds)
```

```
dt_div_min_dr_div_cs_limit = -1
dt_div_min_dr_div_cs_hard_limit = -1
```

min_k_for_dt_div_min_dr_div_cs_limit ¶**min_q_for_dt_div_min_dr_div_cs_limit ¶****max_q_for_dt_div_min_dr_div_cs_limit ¶**

```
min_k_for_dt_div_min_dr_div_cs_limit = 20
min_q_for_dt_div_min_dr_div_cs_limit = 0.005d0
max_q_for_dt_div_min_dr_div_cs_limit = 0.995d0
```

min_abs_du_div_cs_for_dt_div_min_dr_div_cs_limit ¶

only use dt_div_min_dr_div_cs_limit at cells where abs_du_div_cs > this limit. allow focus on regions near shock face.

```
min_abs_du_div_cs_for_dt_div_min_dr_div_cs_limit = 0.01d0
```

dt_div_dt_cell_collapse_limit ¶**dt_div_dt_cell_collapse_hard_limit ¶**

limit for dt compared to cell_collapse timescale (negative means no limit)

```
dt_cell_collapse = min over shells k that have  $v(k+1) > v(k)$  of  $(r(k)-r(k+1))/(v(k+1)-v(k))$ , the time for the cell to collapse to zero thickness at current velocities.
```

```
dt_div_dt_cell_collapse_limit = -1
dt_div_dt_cell_collapse_hard_limit = -1
```

limits based on changes in location on HR diagram

delta_HR_ds_L ¶

delta_HR_ds_Teff ¶

```
dlgL = log10(L/L_prev)
dlgTeff = log10(Teff/Teff_prev)
```

```
delta_HR_ds_L = 1
delta_HR_ds_Teff = 1
```

delta_HR_limit ¶

delta_HR_hard_limit ¶

limit for dHR (negative means no limit)

```
dHR = sqrt((delta_HR_ds_L*dlgL)**2 + (delta_HR_ds_Teff*dlgTeff)**2)
```

```
delta_HR_limit = -1
delta_HR_hard_limit = -1
```

limits based on change in max temperature or density

delta_lgT_max_limit ¶

delta_lgT_max_hard_limit ¶

limit for magnitude of change in log10 max temperature

```
delta_lgT_max_limit = -1
delta_lgT_max_hard_limit = -1
```

delta_lgRho_max_limit ¶

delta_lgRho_max_hard_limit ¶

limit for magnitude of change in log10 max density

```
delta_lgRho_max_limit = -1  
delta_lgRho_max_hard_limit = -1
```

limits based on changes at center

delta_lgT_cntr_limit ¶

delta_lgT_cntr_hard_limit ¶

limit for magnitude of change in log10 temperature at center

```
delta_lgT_cntr_limit = 0.01d0  
delta_lgT_cntr_hard_limit = -1
```

delta_lgP_cntr_limit ¶

delta_lgP_cntr_hard_limit ¶

limit for magnitude of change in log10 pressure at center

```
delta_lgP_cntr_limit = -1  
delta_lgP_cntr_hard_limit = -1
```

delta_lgRho_cntr_limit ¶

delta_lgRho_cntr_hard_limit ¶

limit for magnitude of change in log10 density at center

```
delta_lgRho_cntr_limit = 0.05d0  
delta_lgRho_cntr_hard_limit = -1
```

delta_log_eps_nuc_cntr_limit ¶

delta_log_eps_nuc_cntr_hard_limit ¶

Limit for magnitude of change in log10 eps_nuc at center. Only applies to increase in eps_nuc in non-convective core.. This can help to catch the start of core convection..

```
delta_log_eps_nuc_cntr_limit = 1  
delta_log_eps_nuc_cntr_hard_limit = -1
```

lg_XH_cntr is log10(h1 mass fraction at center). Small timesteps as the center hydrogen is exhausted.

delta_lg_XH_cntr_min

Ignore changes in lg_XH_cntr if value is less than this.

```
delta_lg_XH_cntr_min = -6
```

delta_lg_XH_cntr_max

Ignore changes in lg_XH_cntr if value is more than this.

```
delta_lg_XH_cntr_max = -3
```

delta_lg_XH_cntr_limit

If max delta is greater than this, reduce the next timestep by delta_lg_XH_cntr_limit/max_delta.

```
delta_lg_XH_cntr_limit = 0.05d0
```

delta_lg_XH_cntr_hard_limit

If max delta is greater than delta_lg_XH_cntr_hard_limit, retry with smaller timestep.

```
delta_lg_XH_cntr_hard_limit = -1
```

lg_XHe_cntr is log10(he4 mass fraction at center) small timesteps as the center helium is exhausted.

delta_lg_XHe_cntr_min

Ignore changes in lg_XHe_cntr if value is less than this.

```
delta_lg_XHe_cntr_min = -6
```

delta_lg_XHe_cntr_max

Ignore changes in lg_XHe_cntr if value is more than this.

```
delta_lg_XHe_cntr_max = -3
```

delta_lg_XHe_cntr_limit

If max delta is greater than delta_lg_XHe_cntr_limit, reduce the next timestep by delta_lg_XHe_cntr_limit/max_delta.

```
delta_lg_XHe_cntr_limit = 0.1d0
```

delta_lg_XHe_cntr_hard_limit ¶

If max delta is greater than `delta_lg_XHe_cntr_hard_limit`, retry with smaller timestep.

```
delta_lg_XHe_cntr_hard_limit = -1
```

`lg_XC_cntr` is $\log_{10}(\text{c12 mass fraction at center})$. Small timesteps as the center carbon is exhausted.

delta_lg_XC_cntr_min ¶

Ignore changes in `lg_XC_cntr` if value is less than this.

```
delta_lg_XC_cntr_min = -5
```

delta_lg_XC_cntr_max ¶

Ignore changes in `lg_XC_cntr` if value is more than this.

```
delta_lg_XC_cntr_max = -3
```

delta_lg_XC_cntr_limit ¶

If max delta is greater than `delta_lg_XC_cntr_limit`, reduce the next timestep by `delta_lg_XC_cntr_limit/max_delta`.

```
delta_lg_XC_cntr_limit = 0.1d0
```

delta_lg_XC_cntr_hard_limit ¶

If max delta is greater than `delta_lg_XC_cntr_hard_limit`, retry with smaller timestep.

```
delta_lg_XC_cntr_hard_limit = -1
```

`lg_XNe_cntr` is $\log_{10}(\text{ne20 mass fraction at center})$ Small timesteps as the center neon is exhausted.

delta_lg_XNe_cntr_min ¶

Ignore changes in `lg_XNe_cntr` if value is less than this.

```
delta_lg_XNe_cntr_min = -5
```

delta_lg_XNe_cntr_max

Ignore changes in lg_XNe_cntr if value is more than this.

```
delta_lg_XNe_cntr_max = 0
```

delta_lg_XNe_cntr_limit

If max delta is greater than delta_lg_XNe_cntr_limit, reduce the next timestep by delta_lg_XNe_cntr_limit/max_delta.

```
delta_lg_XNe_cntr_limit = 1d99
```

delta_lg_XNe_cntr_hard_limit

If max delta is greater than delta_lg_XNe_cntr_hard_limit, retry with smaller timestep.

```
delta_lg_XNe_cntr_hard_limit = -1
```

lg_XO_cntr is log10(o16 mass fraction at center) Small timesteps as the center oxygen is exhausted.

delta_lg_XO_cntr_min

Ignore changes in lg_XO_cntr if value is less than this.

```
delta_lg_XO_cntr_min = -5
```

delta_lg_XO_cntr_max

Ignore changes in lg_XO_cntr if value is more than this.

```
delta_lg_XO_cntr_max = 0
```

delta_lg_XO_cntr_limit

If max delta is greater than delta_lg_XO_cntr_limit, reduce the next timestep by delta_lg_XO_cntr_limit/max_delta.

```
delta_lg_XO_cntr_limit = 1d99
```

delta_lg_XO_cntr_hard_limit

If max delta is greater than delta_lg_XO_cntr_hard_limit, retry with smaller timestep.

```
delta_lg_XO_cntr_hard_limit = -1
```

lg_XSi_cntr is log10(si28 mass fraction at center) Small timesteps as the center silicon is exhausted.

delta_lg_XSi_cntr_min ¶

Ignore changes in lg_XSi_cntr if value is less than this.

```
delta_lg_XSi_cntr_min = -5
```

delta_lg_XSi_cntr_max ¶

Ignore changes in lg_XSi_cntr if value is more than this.

```
delta_lg_XSi_cntr_max = 0
```

delta_lg_XSi_cntr_limit ¶

If max delta is greater than delta_lg_XSi_cntr_limit, reduce the next timestep by delta_lg_XSi_cntr_limit/max_delta.

```
delta_lg_XSi_cntr_limit = 1d99
```

delta_lg_XSi_cntr_hard_limit ¶

If max delta is greater than delta_lg_XSi_cntr_hard_limit, retry with smaller timestep.

```
delta_lg_XSi_cntr_hard_limit = -1
```

limits based on changes in mass of the star ¶

delta_lg_star_mass_limit ¶

delta_lg_star_mass_hard_limit ¶

Limit for magnitude of change in log10(M/Msun).

```
delta_lg_star_mass_limit = 5d-3
delta_lg_star_mass_hard_limit = -1
```

limit for change in mdot in Msun/yr

- delta_mdot_atol tolerance for absolute changes
- delta_mdot_rtol tolerance for relative changes

```
delta_mdot_atol = 1d-3
delta_mdot_rtol = 0.5d0
```

delta_mdot_limit ¶

delta_mdot_hard_limit ¶

```
delta_mdot = abs(mdot - mdot_old)/ (delta_mdot_atol*Msun/secyer + &
    delta_mdot_rtol*max(abs(mdot),abs(mdot_old)))
```

ignore if < 0

```
delta_mdot_limit = -1
delta_mdot_hard_limit = -1
```

factor_for_test_CpT_absMdot_div_L ¶

Limit on ratio $C_p(k) * T(k) * \text{abs}(\text{mstar_dot}) / L(k)$ at $k = k_for_CpT_absMdot_div_L$. Cell index $k_for_CpT_absMdot_div_L$ is set by the `adjust_mass` routine as follows: Let delta_m be $\text{mdot} * dt$, the change in mass for this step. Let $\text{delta_m_for_limit} = \text{abs}(\text{delta_m}) * \text{factor_for_test_CpT_absMdot_div_L}$. Then $k_for_CpT_absMdot_div_L$ is the outermost cell boundary k , where the mass exterior to k is $\geq \text{delta_m_for_limit}$.

```
factor_for_test_CpT_absMdot_div_L = 1
```

CpT_absMdot_div_L_limit ¶

Only use if > 0. Reduce next timestep if ratio is greater than this limit.

```
CpT_absMdot_div_L_limit = -1
```

CpT_absMdot_div_L_hard_limit ¶

Only use if > 0. Retry if ratio exceeds this limit.

```
CpT_absMdot_div_L_hard_limit = -1
```

limits based on changes in log total angular momentum ¶

delta_lg_total_J_limit ¶

If max delta is greater than `delta_lg_total_J_limit`, reduce the next timestep by `delta_lg_total_J_limit/max_delta`.

```
delta_lg_total_J_limit = 0.1d0
```

delta_lg_total_J_hard_limit ¶

If max delta is greater than delta_lg_total_J_hard_limit, retry with smaller timestep.

```
delta_lg_total_J_hard_limit = 0.5d0
```

limit_for_rel_error_in_energy_conservation ¶

hard_limit_for_rel_error_in_energy_conservation ¶

```
rel_error_in_energy_conservation = abs(error_in_energy_conservation/to
```

```
limit_for_rel_error_in_energy_conservation = 1d99
hard_limit_for_rel_error_in_energy_conservation = 1d99
```

limit_for_rel_rate_in_energy_conservation ¶

hard_limit_for_rel_rate_in_energy_conservation ¶

```
rel_rate_in_energy_conservation = abs(error_in_energy_conservation/tot
```

```
limit_for_rel_rate_in_energy_conservation = 1d99
hard_limit_for_rel_rate_in_energy_conservation = 1d99
```

limit_for_avg_lgE_residual ¶

hard_limit_for_avg_lgE_residual ¶

```
avg_lgE_residual = abs(dot_product(s% dq(1:s% nz),s% lnE_residual(1:s%
```

```
limit_for_avg_lgE_residual = 1d99
hard_limit_for_avg_lgE_residual = 1d99
```

limit_for_max_abs_lgE_residual ¶

hard_limit_for_max_abs_lgE_residual ¶

```
max_abs_lgE_residual = maxval(abs(s% lnE_residual(1:s% nz)))/ln10
```

```
limit_for_max_abs_lgE_residual = 1d99
hard_limit_for_max_abs_lgE_residual = 1d99
```

limit_for_avg_v_residual ¶**hard_limit_for_avg_v_residual ¶**

```
avg_v_residual = abs(dot_product(s% dq(1:s% nz),s% v_residual(1:s% nz))
```

```
limit_for_avg_v_residual = 1d99
hard_limit_for_avg_v_residual = 1d99
```

limit_for_max_abs_v_residual ¶**hard_limit_for_max_abs_v_residual ¶**

```
max_abs_v_residual = maxval(abs(s% v_residual(1:s% nz)))
```

```
limit_for_max_abs_v_residual = 1d99
hard_limit_for_max_abs_v_residual = 1d99
```

limit_for_abs_rel_E_err ¶**hard_limit_for_abs_rel_E_err ¶**

```
abs_rel_E_err = maxval(abs(s% rel_E_err(1:s% nz)))
```

```
limit_for_abs_rel_E_err = 1d99
hard_limit_for_abs_rel_E_err = 1d99
```

limit_for_max_E_residual ¶**hard_limit_for_max_E_residual ¶**


```
max_E_residual = maxval(abs(s% E_residual(1:s% nz)))
```

```
limit_for_max_E_residual = 1d99  
hard_limit_for_max_E_residual = 1d99
```

report_why_dt_limits ¶

If true, produce terminal output about choice of timestep.

```
report_why_dt_limits = .false.
```

report_all_dt_limits ¶

If true, produce terminal output about all influences for choice of timestep.

```
report_all_dt_limits = .false.
```

report_hydro_dt_info ¶

If true, produce terminal output about choice of timestep based on varcontrol_target.

```
report_hydro_dt_info = .false.
```

report_dX_nuc_drop_dt_limits ¶

If true, report timestep limits from drop in abundance from nuclear reactions.

```
report_dX_nuc_drop_dt_limits = .false.
```

debugging controls ¶

report_hydro_solver_progress ¶

Set true to see info about newton iterations.

```
report_hydro_solver_progress = .false.
```

report_ierr ¶

If true, produce terminal output when have some internal error.

```
report_ierr = .false.
```

stop_for_NaN

If true and report_ierr is also true, then stop for NaNs.

```
stop_for_NaN = .false.
```

trace_newton_bcyclic_solve_input

Input is “B” j k iter B(j,k).

```
trace_newton_bcyclic_solve_input = .false.
```

trace_newton_bcyclic_solve_output

Output is “X” j k iter X(j,k).

```
trace_newton_bcyclic_solve_output = .false.
```

trace_newton_bcyclic_matrix_input

Matrix before factor.

```
trace_newton_bcyclic_matrix_input = .false.
```

trace_newton_bcyclic_matrix_output

Matrix after factor.

```
trace_newton_bcyclic_matrix_output = .false.
```

trace_newton_bcyclic_steplo

1st model number to trace.

```
trace_newton_bcyclic_steplo = 1
```

trace_newton_bcyclic_stephi

Last model number to trace.

```
trace_newton_bcyclic_stephi = -1
```

trace_newton_bcyclic_iterlo ¶

1st newton iter to trace.

```
trace_newton_bcyclic_iterlo = 1
```

trace_newton_bcyclic_iterhi ¶

Last newton iter to trace.

```
trace_newton_bcyclic_iterhi = -1
```

trace_newton_bcyclic_nzlo ¶

1st cell to trace.

```
trace_newton_bcyclic_nzlo = 1
```

trace_newton_bcyclic_nzhi ¶

Last cell to trace; if < 0 , then use nz as nzhi.

```
trace_newton_bcyclic_nzhi = -1
```

trace_newton_bcyclic_jlo ¶

1st var to trace.

```
trace_newton_bcyclic_jlo = 1
```

trace_newton_bcyclic_jhi ¶

Last var to trace; if < 0 , then use nvar as jhi.

```
trace_newton_bcyclic_jhi = -1
```

To get info about the mesh set `show_mesh_changes = .true..` Restart and get the `mesh_call_number` from terminal output. Set `mesh_dump_call_number = mesh_call_number`. Restart and it will write data files to `mesh_plot_data`. view with `test/mesh.rb` and `test/mesh_plan.rb`.

show_mesh_changes ¶

When `show_mesh_changes` is true, the terminal output includes the `mesh_call_number`.

```
show_mesh_changes = .false.
```

mesh_dump_call_number ¶

When `mesh_call_number == mesh_dump_call_number`, various plotting information is written..

```
mesh_dump_call_number = -1
```

trace_evolve ¶

Send evolve output to screen.

```
trace_evolve = .false.
```

variety of output from the hydro solver

hydro solver

```
hydro_numerical_jacobian = .false.  
hydro_jacobian_nzlo = 1  
hydro_jacobian_nzhi = -1  
hydro_check_everything = .false.  
hydro_inspectB_flag = .false.  
hydro_sizequ_flag = .false.  
hydro_get_a_numerical_partial = -1d0  
hydro_testpartials_k = -1  
hydro_show_correction_info = .false.  
hydro_save_numjac_plot_data = .false.  
hydro_dump_call_number = -1  
hydro_dump_iter_number = -1  
hydro_epsder_struct = 1d-5  
hydro_epsder_chem = 1d-5
```

hydro_save_photo ¶

Saves a photo when `hydro_call_number = hydro_dump_call_number - 1`

```
hydro_save_photo = .false.
```

xa_clip_limit ¶

Abundances smaller than this limit are set to 0.

```
xa_clip_limit = 1d-99
```

trace_k ¶

Print out trace information about cell with number = trace_k.

```
trace_k = -1
```

fill_arrays_with_NaNs ¶

initialize arrays with NaNs to trap reads of uninitialized entries.

```
fill_arrays_with_NaNs = .true.
```

zero_when_allocate ¶

initialize arrays with zeros.

```
zero_when_allocate = .false.
```

miscellaneous controls ¶

relax_dY ¶

Change Y by this amount per step when relaxing Y.

```
relax_dY = 0.005d0
```

relax_dlnZ ¶

Change lnZ by this amount per step when relaxing Z. Default is ln10/10.

```
relax_dlnZ = 2.3025850929940459d-1
```

zams_filename ¶

Default is for Z=0.02, Y=0.28.

```
zams_filename = 'zams_z2m2_y28.data'
```

```
set_rho_to_dm_div_dV = .true.
```

use_other_{hook} ¶

Logicals to deploy the use_other routines.

```
use_other_eos = .false.  
use_other_kap = .false.  
use_other_atm = .false.  
use_other_diffusion = .false.  
use_other_mlt = .false.
```

```
use_other_adjust_net = .false.  
use_other_adjust_mdot = .false.  
use_other_j_for_adjust_J_lost = .false.  
use_other_alpha_mlt = .false.  
use_other_am_mixing = .false.  
use_other_brunt = .false.  
use_other_brunt_smoothing = .false.  
use_other_build_initial_model = .false.  
use_other_cgrav = .false.
```

```
use_other_energy_implicit = .false.  
use_other_energy = .false.  
use_other_momentum = .false.  
use_other_eps_grav = .false.  
use_other_mesh_functions = .false.  
use_other_D_mix = .false.
```

```
use_other_neu = .false.  
use_other_net_get = .false.  
use_other_newton_monitor = .false.  
use_other_opacity_factor = .false.  
use_other_paquette_coefficients = .false.  
use_other_pgstar_plots = .false.  
use_other_porosity_factor = .false.  
use_other_eval_fp_ft = .false.  
use_other_torque = .false.
```

```
use_other_torque_implicit = .false.  
use_other_eta_visc = .false.  
use_other_wind = .false.  
use_other_split_mix = .false.  
use_other_after_struct_burn_mix = .false.  
use_other_before_struct_burn_mix = .false.  
use_other_surface_PT = .false.
```

```
use_other_export_pulse_data = .false.  
use_other_get_pulse_data = .false.  
use_other_edit_pulse_data = .false.
```

```
use_other_astero_freq_corr = .false.
```

mixing diffusion coeffs ¶

sig_term_limit ¶

Limit on coefficients in convective mixing equations. Consider a diffusion eqn of form:

$$x(k) - x_0(k) = c_1 * (x(k-1) - x(k)) - c_2 * (x(k) - x(k+1))$$

Simplify for $c_1=c_2=c$, $x(k-1)=x(k+1)=x_0(k)=x_0$, $x(k)=x_0+dx$ Then eqn becomes

$$(1+2*c)*(x_0+dx) - 2*c*x_0 = x_0$$

If $2*c \gg 1$, then eqn becomes ill-conditioned, so we enforce $c \leq \text{sig_term_limit}$ In physical terms c is $dt * \text{sig} / dm$, where $\text{sig} = (4 \pi r^2 \rho)^2 * D$ and D = diffusion coeff (cm^2/s), so c can get large when dt/dm is large.

```
sig_term_limit = 1d13
```

am_sig_term_limit ¶

Limit on coefficients in angular momentum transport equations. Necessary for numerical stability. Plays same role as `sig_term_limit` for material mixing.

```
am_sig_term_limit = 1d13
```

sig_min_factor_for_high_Tcenter ¶

High center T limit to avoid negative mass fractions. If $T_{\text{center}} \geq T_{\text{center_min_for_sig_min_factor_full_on}}$, then okay to reduce sig by as much as this factor as needed to prevent causing negative abundances. Inactive when $\geq 1d0$.

```
sig_min_factor_for_high_Tcenter = 0.01d0
```

Tcenter_min_for_sig_min_factor_full_on ¶

If $T_{\text{center}} \geq$ this, $\text{factor} = \text{sig_min_factor_for_neg_abundances}$, this should be $> T_{\text{center_max_for_sig_min_factor_full_off}}$.

```
Tcenter_min_for_sig_min_factor_full_on = 3.2d9
```

Tcenter_max_for_sig_min_factor_full_off ¶

If Tcenter <= this, factor = 1, so has no effect this should be < Tcenter_min_for_sig_min_factor_full_on. For T > full_off and < full_on, factor changes linearly with Tcenter.

```
Tcenter_max_for_sig_min_factor_full_off = 2.8d9
```

max_delta_m_to_bdy_for_sig_min_factor ¶

sig_min factor goes to 1 as distance (in Msun units) from boundary of mixing region reaches this value

```
max_delta_m_to_bdy_for_sig_min_factor = 0.5d0
```

delta_m_upper_for_sig_min_factor ¶

okay to change sig min factor to 1 for mix region larger than this

```
delta_m_upper_for_sig_min_factor = 0.3d0
```

delta_m_lower_for_sig_min_factor ¶

don't change sig min factor for mix region smaller than this

```
delta_m_lower_for_sig_min_factor = 0.1d0
```

Tcenter_max_for_dble_bccyclic ¶

if Tcenter <= this, use dble precision version of bccyclic. if Tcenter > this, use quad precision.

```
Tcenter_max_for_dble_bccyclic = 1d99
```

extra params as a convenience for developing new features note: the parameter num_x_ctrls is defined in star_def.inc

```
x_ctrl(1:num_x_ctrls) = 0d0
x_integer_ctrl(1:num_x_ctrls) = 0
x_logical_ctrl(1:num_x_ctrls) = .false.
```

One can split controls inlist into pieces using the following parameters. BTW: it works recursively, so the extras can read extras too.

read_extra_controls_inlist1 ¶**extra_controls_inlist1_name ¶**

If `read_extra_controls_inlist1` is true, then read &controls from this namelist file.

```
read_extra_controls_inlist1 = .false.  
extra_controls_inlist1_name = 'undefined'
```

If you try one of the following prebuilt extras, you must also set `read_extra_star_job_inlist1` true and change the `extra_star_job_inlist1_name` to match `extra_controls_inlist1_name`.

evolve 1 Msun from pre-ms to white dwarf

```
read_extra_controls_inlist1 = .true.  
extra_controls_inlist1_name = 'inlist_extras_1M_lifecycle'
```

for debugging

```
extra_controls_inlist1_name = 'inlist_debug'
```

read_extra_controls_inlist2 ¶

extra_controls_inlist2_name ¶

If `read_extra_controls_inlist2` is true, then read &controls from this namelist file.

```
read_extra_controls_inlist2 = .false.  
extra_controls_inlist2_name = 'undefined'
```

read_extra_controls_inlist3 ¶

extra_controls_inlist3_name ¶

If `read_extra_controls_inlist3` is true, then read &controls from this namelist file.

```
read_extra_controls_inlist3 = .false.  
extra_controls_inlist3_name = 'undefined'
```

read_extra_controls_inlist4 ¶

extra_controls_inlist4_name ¶

If `read_extra_controls_inlist4` is true, then read &controls from this namelist file.

```
read_extra_controls_inlist4 = .false.  
extra_controls_inlist4_name = 'undefined'
```

read_extra_controls_inlist5 ¶

extra_controls_inlist5_name ¶

If read_extra_controls_inlist5 is true, then read &controls from this namelist file.

```
read_extra_controls_inlist5 = .false.  
extra_controls_inlist5_name = 'undefined'
```